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# Forearc high uplift by deep crustal flow during growth of the Anatolian margin

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## Abstract

We present a model for the dynamic formation of the forearc high of southern Anatolia where sedimentation in the forearc basin leads to thermally-activated deformation in the lower crust. Our thermo-mechanical models demonstrate that forearc sedimentation increases the temperature of the underlying crust by “blanketing” the heat flux and increasing Moho depth. Deformation switches from frictional to viscous with a higher strain rate led by increased temperature. Viscous deformation changes large-wavelength subsidence into coeval, short-wavelength uplift and subsidence. Models show that forearc highs are intrinsic to accretionary wedges and can grow dynamically and non-linearly at rates dependent on sediment accretion, sedimentation and temperature. The mechanism explains Neogene first-order upper-plate strain and vertical motions in the Anatolian margin along Central Cyprus, and in the orogenic plateau margin in South Turkey. This system is analogous to forearc highs in other mature accretionary margins, like Cascadia, Lesser Antilles or Nankai.

**keywords:** forearc high; plateau margin; subduction zone; subduction wedge; uplift; Central Anatolian Plateau

# 1 **1 Introduction**

2 Climatic and geodynamic processes are the first-order drivers of topography in orogenic  
3 plateaus and plateau margins. However, mechanisms for detailed patterns of uplift in orogenic  
4 plateau systems, such as Himalaya-Tibet and Puna-Altiplano (e.g., Allmendinger et al., 1997;  
5 Molnar, 1984) remain diverse and difficult to generalize. This is also true for the history of  
6 topography growth of the orogenic plateau of Central Anatolia and its margins. While  
7 continental delamination (Bartol and Govers, 2014) or lithospheric drip (Göğüş et al., 2017)  
8 have been suggested to sustain Central Anatolia low relief at ~1 km, its plateau margins are  
9 geodynamically different; transpressional orogenic uplift may have formed the northern margin  
10 (Yildirim et al., 2011) whereas the southern margin is strongly influenced by the Cyprus  
11 subduction zone to the south. The formation mechanism of the latter is of particular interest, for  
12 the uplift in the southern margin has a limited N-S extent and occurred in the absence of  
13 regional accommodating faults (Fernández-Blanco et al., 2019).

14 The southern margin of the Central Anatolian Plateau (SCAP) (Fig. 1) is suggested to form by  
15 shallow break-off of the Cyprus slab (e.g., Schildgen et al., 2014). This geodynamic scenario is  
16 compatible with stratigraphic and paleontological records (Cosentino et al., 2012) and the  
17 thinned Central Anatolian lithosphere inferred from seismic tomography (Mutlu and Karabulut,  
18 2011). However, the presence of the African slab (e.g., Abgarmi et al., 2017) and the thick  
19 Anatolian mantle lithosphere (Delph et al., 2017) under the modern Central Taurides suggest  
20 that this mechanism is not likely to be active over the entire E-W extent of the Taurides. In  
21 addition, the large wavelength motions expected at the surface during slab break-off (e.g., Göğüş  
22 and Pysklywec, 2008) are at odds with the sharp bound and short wavelength between margin  
23 uplift and seaward subsidence in South Turkey (Fernández-Blanco et al., 2019; Walsh-Kennedy  
24 et al., 2014). Therefore, we ought to look for other causal mechanisms, such as crustal  
25 thickening, to explain the elevation of the plateau margin (Fernández-Blanco, 2014; Fernández-  
26 Blanco et al., 2019; Meijers et al., 2018).

27 Forearc deformation of an orogenic wedge (Platt, 1986; Willett et al. 1993), where deep-  
28 seated crustal flow results in forearc high uplift has been proposed as a general mechanism for  
29 forearc uplift with examples from the Lesser Antilles, Cascadia and other accretionary margins  
30 (e.g., Pavlis and Bruhn, 1983; Williams et al., 1994). The development of a preceding forearc  
31 basin that is later fragmented by the uplift of the high (McNeill et al., 2000) has been suggested  
32 in more detailed mechanical models. If the forearc basin forms as a negative-alpha basin,  
33 stabilizing the underlying wedge (Willett & Schlunegger, 2010), basin growth may promote  
34 ductile deformation in the lower crust, and a later stage, uplift of the forearc high (Fuller et al.,  
35 2006). We demonstrate that the mechanism of deep crustal thickening by thermally-activated  
36 viscous flow, as proposed by Fuller et al. (2006) and expanded here, reproduces the first-order  
37 spatiotemporal pattern of deformation and vertical motion (uplift and subsidence) across the  
38 SCAP during its growth as the forearc high of the Central Cyprus forearc.

39 Here, we explore the role of wedge-top sedimentation on forearc dynamics for conditions  
40 applicable to the Cyprus-Anatolia margin. To do this, we use transects of the Central Cyprus  
41 subduction margin and its forearc, derived by integration of geophysical and geological data, to  
42 constraint coupled thermo-mechanical, visco-plastic numerical models. An important transition  
43 emerges in these models as accretionary growth and sediment deposition produce a “thermal  
44 blanketing” effect restricting heat flux from underneath the forearc basin and leading to thermal  
45 weakening of the upper plate. In this context, a forearc high grows dynamically and non-linearly  
46 as an integral part of an accreting wedge, often seaward of any continental backstop. This  
47 thermo-mechanical interplay can provide an important uplift mechanism applicable to the  
48 Anatolian margin that can be generalised to similar accretionary margins.

## 49 **2 Background**

50 The Africa/Arabia-Eurasian plate convergence and subduction dynamics that dominate the  
51 Mediterranean controls the evolution of the Anatolian margin (Wortel and Spakman, 2000),  
52 where the SCAP occupies the forearc high of the Central Cyprus subduction at present (Fig. 1).

53 Early to Late Miocene subsidence broadened a wide forearc basin that spanned from South  
54 Turkey to Central Cyprus, and led to protracted growth of a laterally continuous carbonate  
55 platform (e.g., Bassant et al., 2005; Karabıyıköğlü et al., 2000). Late Miocene regional vertical  
56 motions of short wavelength and opposite sense led to coeval South Turkey uplift and seaward  
57 subsidence (Walsh-Kennedy et al., 2014). These short-wavelength motions of opposite sense  
58 fragmented the antecedent forearc basin while forming the SCAP as a flexural monocline absent  
59 of regional surface-reaching faults (Fernández-Blanco et al., 2019). The vertical motions led to  
60 truncation and erosion of Late Miocene rocks in the uplifting sectors of the SCAP while  
61 subsiding sectors of the monocline sustained deposition, leading to maximum sediment  
62 thicknesses near the Turkish coast (Aksu et al., 2005; Walsh-Kennedy et al., 2014). At present,  
63 the monocline is delineated by Late Miocene shallow marine rocks laying at ~2 km elevation in  
64 the modern Central Taurides (Cosentino et al., 2012) and at ~-2 km depth in the Cilicia Basin  
65 (Aksu et al., 2005).

66 Mutually contradictory evidence along the plateau margin keeps open the debate on when  
67 and how it was formed. Biostratigraphy and polarity chrons report an age of 8.35–8.108 Ma for  
68 marine rocks at the top of the sequence in the present-day Central Taurides hinterland  
69 (Cosentino et al., 2012). The lateral continuity of these marine rocks across the plateau margin  
70 and their contact relationships with the basement, together with the absence of Messinian  
71 deposits, suggest km-scale topography by >5 Ma (Fernández-Blanco et al., 2019). This is  
72 consistent with stable isotope paleoaltimetry estimates pointing at ~2 km of relief at ~5 Ma  
73 (Meijers et al., 2018), and the protracted sedimentary deposition in thick, stacked delta lobes in  
74 the Göksu Delta during latest Messinian - Recent seen in seismic reflection lines (Aksu et al.,  
75 2014). Cosmogenic datation of Göksu River terraces and uplifted marine fossil assemblages  
76 suggest surface uplift onset between 8 and 5.45 Ma with average uplift rates of 0.25 to 0.37  
77 mm/yr, and uplift rates of 0.72 to 0.74 mm/yr leading to 1.2 km of surface uplift since 1.66 to  
78 1.62 Ma (Schildgen et al., 2012). Paleontological evidence in younger marine rocks at the  
79 margins of the Göksu River or closer to the coast points to uplift rates as fast as 3.21-3.42

80 mm/yr and suggests 1200-1500 m of topographic growth in the last ~450 ka (Öğretmen et al.,  
81 2018).

82 Studies using the ages of onshore uplifted features as markers of different phases of uplift  
83 suggest the SCAP formed by delamination and/or slab break-off events that vary in time and/or  
84 depth (e.g., Cosentino et al., 2012; Öğretmen et al., 2018; Schildgen et al., 2014, 2012). However,  
85 these studies constrain depositional age or subaerial exposure of rocks at specific sites, and the  
86 proposed geodynamic scenarios of multiple uplift phases overlook the presence of the African  
87 slab below the modern Central Taurides and the relevant thickness of the Anatolian crust and  
88 lithosphere there (e.g., Abgarmi et al., 2017; Bakırcı et al., 2012; Biryol et al., 2011; Delph et al.,  
89 2017). Studies accounting for aforementioned geological constraints and geophysical  
90 observables suggest crustal thickening as an alternative mechanism of plateau margin  
91 formation (Fernández-Blanco, 2014; Fernández-Blanco et al., 2019; Meijers et al., 2018). Yet,  
92 these latter studies seem to be at odds with the youngest uplifted marine rocks and provide no  
93 detail on the mechanism leading to crustal thickening.

### 94 **3 The Anatolian margin**

#### 95 **3.1 South Anatolian margin transect: Lithospheric and crustal structure**

96 We reconstruct a plate-scale transect spanning from the East Mediterranean to the Central  
97 Anatolian Plateau interior along 33°30'E longitude (Fig. 2). To portray the lithospheric  
98 structure, we integrate constraints from Biryol et al. (2011) and Bakırcı et al. (2012) into the  
99 TransMED transect VII (Stephenson et al., 2004). To derive the crustal structure and constrain  
100 the thicknesses of the African and Anatolian crust as well as the dip of the Cyprus slab, we  
101 collectively interpret data from 10+ geophysical studies along the section (Fig. 2B,-C, see  
102 caption).

103 Along the Central Cyprus subduction zone, the African lithosphere under-thrusts  
104 northwards below the Anatolian plate (Fig. 2D). The overriding Anatolian lithosphere has

105 maximum thicknesses of ~110 km at the contact with the Cyprus slab below the modern  
106 Central Taurides, and thins northwards down to ~85 km in Central Anatolia (Fig. 2D).

107 In the southern sectors of the transect, crustal thicknesses are well detected by the  
108 gravimetric signal of Ergün et al. (2005) and the Moho models of Koulakov and Sobolev (2006).  
109 In the African plate, crustal thicknesses of ~28 km are observed at the site of the Eratosthenes  
110 Seamount, south of Cyprus, where the African lithosphere is ~40 km thicker than northwards  
111 (Fig. 2D). The oceanic crust is the thinnest (~25 km) below the trench area. Northward and  
112 between the subducting and overriding plates, thickening occurs in relation to the Troodos  
113 Ophiolite, possibly as a result of thrust doubling due to its emplacement. The locked  
114 underthrust of the Eratosthenes Seamount is underneath this location, and the detachment  
115 depth of the Troodos Ophiolite is uncertain. Similarly, the extent of the continental crust  
116 underneath the Troodos Ophiolite and the position of its transition to oceanic crust farther to  
117 the north remains enigmatic. The Anatolian plate has maximum crustal thicknesses of ~45 km  
118 (Luccio and Pasyanos, 2007) below the Central Taurides that decrease gently towards the plate  
119 interior to ~35 km (Fig. 2D). For this interpretation, we used Pn tomography from Mutlu and  
120 Karabulut (2011) instead of gravity data (Özeren and Holt, 2010), which points to crustal  
121 thickness values up to 10 km thicker (Fig. 2C).

122 All geophysical models concur on a northwards increase in Moho depth from ~28 km to >40  
123 km between 34°30'N and 37°N (Fig. 2C) that we correlate with the steepening of the subducting  
124 slab (up to 40°) (Fig. 2D). An overall subduction angle of 45° is observed until ~60 km depth at  
125 36°30', where angles of ~60° are reached as the slab deepens. Northward prolongation at  
126 similar dips fit appearances of the slab at ~300 km in the interior of Central Anatolia (A-A  
127 section of Biryol et al., 2011).

### 128 **3.2 Central Cyprus forearc transect: Structural and stratigraphic relationships**

129 We reproduce uppermost crustal structures and the geometry of Miocene and younger rocks  
130 (Fig. 3) integrating own findings (Fernández-Blanco, 2014; Fernández-Blanco et al., 2019) with

131 published data in regional studies (e.g., Calon et al., 2005; Robertson, 1998; Stephenson et al.,  
132 2004). We assembled the interpretations of each area as shown in their original sources and the  
133 reader is referred there for details.

134 Compressional, regional-scale structures along the Cyprus forearc become older  
135 northwards (Fig. 3A). South-verging thrusts rooted in the subduction megathrust are presently  
136 active in the trench and pass northwards into thrust culminations covered by Quaternary and  
137 Pleistocene rocks in North Cyprus (Fig. 3A-b). The north-verging thrust in Central Cilicia Basin  
138 is mid-Pliocene (Fig. 3A-c). In the Cilicia Basin northern margin, Messinian salts pinch out where  
139 Pliocene rocks overlay an erosional contact with Miocene rocks, attesting to pre-Pliocene uplift.  
140 Uplifted Miocene rocks in the Mut Basin delineate a flexural monocline with no Miocene or  
141 younger surface-reaching thrusts (Fig. 3A-d). These regional-scale structures result in structural  
142 highs that bound basins or basin sectors and compartmentalize the Cyprus forearc at distances  
143 of ~40-50 km (Fig. 3A, a to f). A basement high and the Kyrenia Range bound the Messaoria  
144 Basin (a to b, ~40 km), and a deep-rooted thrust system in the center of the Cilicia Basin (Fig.  
145 3A-c) set two sub-basins with similar length (~50 km). Basement highs discriminating sectors  
146 within the Mut Basin (d, e, f in Fig. 3A) also appear at similar distances. These observations are  
147 consistent with strain accommodation lead by accretion in the Central Cyprus subduction  
148 margin.

149 Basin infill is regionally continuous until the Messinian and deposited exclusively in  
150 seaward sectors of the Central Cyprus forearc thereafter (Fig. 3B). After terrestrial  
151 sedimentation, pre-Messinian Miocene neritic limestones were deposited atop pre-Miocene  
152 basement (Cosentino et al., 2012). These shallow-water rocks are continuous from the  
153 Messaoria Basin, where the pre-Messinian basin thins to the south, to Central Turkey. Since the  
154 Messinian, rocks deposited seaward off the present Turkish coast and have basin depocenters  
155 occurring at northward locations at younger ages (Fig. 3B). This evidence suggests protracted,  
156 large-wavelength subsidence of a wide forearc basin prior to the Messinian, followed by  
157 younger surface uplift of the modern Central Taurides with concomitant, counteracting



158 subsidence in the Cilicia Basin.

## 159 **4 Thermo-mechanical models of accretion**

160 Upper-plate strain and morphology at accretionary margins is often described using the critical  
161 wedge theory (e.g., Davis et al., 1983), which defines the geometry of the orogenic wedge as a  
162 function of the mechanical properties of the accreting wedge. In its strict, brittle form, critical  
163 Coulomb wedge theory does not include the ductile properties of these systems. Research  
164 inclusive of the visco-plastic attributes show the influence of thermal or rheological variations  
165 or that of sediment load and/or competence in the strain distribution and deformation patterns  
166 within the accretionary wedge (e.g., Fillon et al., 2013; Fuller et al., 2006; Mannu et al., 2016;  
167 Royden, 1996; Vanderhaeghe et al., 2003; Willett and Schlunegger, 2010; Williams et al., 1994).

### 168 **4.1 Model set-up and strategy**

169 We use a 2D kinematic-dynamic model with standard rheological and thermal parameters to  
170 explore feasible mechanisms leading to the present structure of the Anatolian subduction  
171 margin (Fig. 4). Models are visco-plastic and similar to those in other studies exploring the  
172 evolution of accretionary settings over millions of years (e.g., Fillon et al., 2013; Fuller et al.,  
173 2006; Mannu et al., 2016; Royden, 1996; Vanderhaeghe et al., 2003; Willett and Schlunegger,  
174 2010; Williams et al., 1994). Albeit important for shorter-term processes, elasticity is only  
175 considered for the flexural isostatic response to vertical loads and included as an elastic  
176 foundation that is calculated by assuming two elastic plates that remain in contact (Fuller,  
177 2006). Our models consist of coupled mechanical and thermal domains. The domain where  
178 mechanical laws apply represents the crust of a deforming subduction zone where sedimentary  
179 accretion of incoming sediments is driven by tangential velocities at its base. Tangential  
180 velocities decrease toward, and become zero, at the “S” point, which represents the contact  
181 point between the subducting slab and continental Moho (Fig. 4). The thermal domain covers  
182 the whole model, including the mechanical domain (Fig. S1). Models simulate the growth of an

183 accretionary wedge at a rate determined by the accretionary flux, as defined by the thickness of  
184 accreting material and convergence velocity.

185 Our simulations aim at being consistent with the time evolution of the Anatolian subduction  
186 margin and parameter values are chosen to match plate-scale and upper crustal observations  
187 along the transects in Central Cyprus (Figs. 2 & 3) at the end of model run. Models simulate 25  
188 Ma of subduction in a transect of 550 km, with an accretionary thickness  $h = 3$  km and a  
189 convergence velocity  $v_c = 35$  mm/yr (Fig. 4), i.e. an accretionary flux of  $105 \text{ km}^2/\text{My}$ , and have a  
190 sedimentation rate of  $Sed_r = 0.5$  mm/yr that occurs as the wedge grows until sediments fill  
191 depressions to capacity. Models have constant accretionary thickness and convergence velocity  
192 that are estimates derived from the extrapolation of present-day values over 25 Ma run time  
193 and are lower and higher than present, respectively. Present-day sedimentary thicknesses in  
194 the East Mediterranean Sea, ranging from 10 km to 15 km (e.g., Makris and Stobbe, 1984), are  
195 the largest in the history of the margin, given the narrower confinement of the modern  
196 Mediterranean and the presence of the Nile. Similarly, present-day convergence velocities of  $9.3$   
197  $\pm 0.3$  mm/yr (Reilinger et al., 2006) are slower than in the recent past due to the underthrusting  
198 of the Eratosthenes Seamount below south Cyprus. In general, results scale such that the  
199 accretionary flux and run time trade-off at nearly one to one.

200 The subducting lithosphere is 50 My old at the left side of the model and has a thickness of  
201 70 km that remains constant during model run time. Since thicknesses in the mechanical  
202 domain change as material is accreted, we chose an initial thickness of 30 km that leads to end-  
203 model crustal thicknesses of 45 km near the “S” point, thereby matching the crustal thickness  
204 below the Central Taurides. The rest of the overriding lithosphere is 80 km thick. To let the  
205 thermal structure equilibrate, the thermal model runs for 20 My before the crustal model  
206 onsets. Cohesion and internal friction angles control the mechanical strengths in our model (Fig.  
207 4). Cohesion,  $c$ , is set to 1000 Pa, a value higher than expected for the crust, to maintain model  
208 stability. Lower values do not affect the outcomes (Fuller, 2006). The internal friction angle of  
209 the crustal material,  $\varphi$ , is set to  $27^\circ$  and the friction angle between the subducting and

210 overriding plate,  $\varphi_b$ , to  $8^\circ$ . Friction values are set low to include the effect of fluid pressures not  
211 explicitly taken into account and imply fluid pressure ratios within the range of those at  
212 accretionary wedges (Fuller, 2006, and references therein). Other parameter values are not  
213 specific to the Anatolian margin nor to our numerical models and are described in detail in  
214 Fuller (2006) and Cassola (2013).

## 215 **4.2 Model results**

216 Model results portray the evolution over millions of years of accreting subduction wedges with  
217 forearc highs, and show the typical morphologic elements of these systems (Fig. 5 & Video S2).  
218 North of the seaward migrating trench, the trench-slope wedge bounds a wide topographic  
219 depression that grows continuously as accommodation space is created by the landward  
220 increasing depth of the subduction slab. Steady infill of the forearc depression suppresses strain  
221 rates and deformation of the underlying wedge, providing the stability to maintain an  
222 undeformed sedimentary basin in a negative-alpha setting (Fig. 5B). Temperature increases  
223 under the basin during its growth leading to a viscosity drop in the lower crust and ductile  
224 strain that ultimately results in uplift of the forearc high and subsidence at seaward locations  
225 (Fig. 5C). At the end of the model run, subduction wedge accretion has led to elevated strain  
226 rates and widespread deformation from the trench to the forearc high. Subduction wedge  
227 accretion also results in wedge topography and wedge thickness continuously increasing  
228 landward until the forearc high, where topographic height is maximum.

### 229 **4.2.1 Sedimentation and forearc high uplift**

230 Sedimentation rates in the forearc basin fundamentally affect vertical motions in the forearc  
231 high at advanced stages of wedge evolution. To explore this effect, we plot how models with  
232 different sedimentation rates (0 to 1.3 mm/yr) change the time evolution of relative vertical  
233 motion between the highest and lowest point of the forearc basement (Fig. 6). At early stages of  
234 wedge evolution (15 Ma model run time), variations in sedimentation rate have minimum effect

235 in forearc topography with relative vertical motions that have constant rates of  $\sim 0.04$ - $0.05$   
236 mm/yr. Thereon, sedimentation rates above and below a threshold value result in two different  
237 trends in the rate of change of relative vertical motions (Fig. 6, Inset A); whereas lower, below-  
238 threshold sedimentation rates ( $< 0.3$  mm/yr) lead to convex-up trends in the time evolution of  
239 relative vertical motions, concave-up trends occur with higher, above-threshold sedimentation  
240 rates ( $0.3$  mm/yr and higher). Below-threshold sedimentation rates result in rates of relative  
241 vertical motion of  $\sim 0.2$ - $0.4$  mm/yr (black and blue in Fig. 6). Above-threshold sedimentation  
242 rates lead to relative vertical motion rates ranging from  $\sim 0.72$  mm/yr to  $\sim 1.75$  mm/yr as  
243 sedimentation rate changes from  $0.3$  mm/yr to  $1.3$  mm/yr, respectively (orange to magenta in  
244 Fig. 6).

245 We also track the horizontal distance between the highest and lowest point of the forearc  
246 basement during the last stages of model run, i.e. the period of differentiation in vertical  
247 motions as led by sedimentation (Fig. 6, Inset B). For all simulations, the horizontal distance  
248 between the highest point on the forearc high and the lowest point on its depocenter is short  
249 ( $< 70$  km) with regards to its associated vertical motion (up to  $> 12$  km). Such horizontal distance  
250 is also controlled by a threshold value in sedimentation rate (Fig. 6, Inset B); whereas models  
251 with lower sedimentation rates lead to horizontal distances of  $\sim 20$  km that are consistent  
252 throughout the model run, those with higher sedimentation rates show horizontal distances  
253 between  $40$  km and  $70$  km that vary during the model run. For the latter simulations, once  
254 vertical motions onset, horizontal distances for the tracked points increase suddenly for a  
255 period of  $\sim 3$  Ma and decrease thereon for the rest of the model run (Fig. 6, Inset B). These  
256 changes in horizontal distance reflect the steepening in time of the transitional area between  
257 uplifting and subsiding areas.

258 The aforementioned non-linear relation between sedimentation rates and forearc high  
259 growth controlled by a threshold value suggest the activation of an external forcing that  
260 contributes to vertical motions in the interior of mature wedges. Below, we evaluate whether  
261 such external forcing is forearc basin sedimentation leading to thermally-activated deformation

262 and viscous flow at the base of the crust. This evaluation also provides information on the  
263 sensitivity to variations of sedimentation rate regarding this sediment “blanketing” effect.

#### 264 ***4.2.2 Sediment blanketing controls on temperature and viscosity***

265 We derive the evolution in time of temperature and viscosity (Fig. 7) using all model grid  
266 elements located in the mechanical domain within a 200 km box centred in the S-point (boxes of  
267 Fig. 5 insets). We plot single values of median temperature (Fig. 6A) and median viscosity (Fig.  
268 6B) for 250 model snapshots (every  $10^5$  yr) for simulations with sedimentation rates ranging  
269 from 0.1 mm/yr to 0.9 mm/yr. To assure a representative, non-parametric calculus, we use the  
270 median for both temperature field and viscosity distribution. For the former, we compute  
271 median temperature values for each grid point, thereby avoiding potential complications  
272 related to multi-modal distributions in temperature fields of geodynamic models. For the latter,  
273 we compute the median viscosity of highly-strained regions (shear zones). To avoid the  
274 tendency of average values to misrepresent power-law distributions, we calculate the strain  
275 rates for all grid elements and use the median viscosity of the 10% with larger strain rates.

276 The amount of wedge-top sediments control the vertical gradient of temperature and the  
277 evolution of viscosity in the underlying wedge (Fig. 7). Models reveal a transition at  $\sim 15$  Ma,  
278 when the same wedge temperature is recorded for all models, regardless of their sedimentation  
279 rate. Models with higher sedimentation rates decrease the overall wedge temperature before  
280 the transition and increase it after the transition, with respect to lower sedimentation rate  
281 models (Fig. 7A). As sedimentation rate increases, temperature variations become not linear  
282 and are twice as large in the advance stages of wedge evolution, with models varying  
283 sedimentation rate from 0.1 mm/yr to 0.9 mm/yr resulting in a temperature change of  $-10^\circ\text{C}$  in  
284 before the transition and  $+30^\circ\text{C}$  degrees after it. This change in wedge temperature with  
285 increased sedimentation rates and the aforementioned transition are also observed in the  
286 evolution of viscosity in shear zones (Fig. 7B). Whereas lower sedimentation rate models have  
287 the lowest values of viscosity in shear zones for early times of model run, higher sedimentation

288 rate models record the lowest values of viscosity at advanced stages (Fig. 7B).

289 Sedimentation rates control the growth of a forearc high in mature accretionary wedges,  
290 and the rate of associated topographic growth, ultimately leading to regional vertical motions of  
291 opposite sense and short wavelength (Fig. 6). The correlation of increased sedimentation rates  
292 with larger relative vertical motions in the forearc high as well as with increased temperatures  
293 and lower viscosity below it (Fig. 7) strongly support that sediments have a “thermal  
294 blanketing” effect that induces viscous flow in the lower crust and the uplift of the forearc high.  
295 Thermal resistivity of sediments leads to “thermal blanketing” and increasingly higher  
296 temperatures in the underlying wedge as the basin grows, and ultimately results in the  
297 progressive change in the deformation mechanism from Coulomb friction to nonlinear viscous  
298 at the base of the crust (Figs. 5B & 7). Lowering strength results in viscous flow at the base of  
299 the orogen and, given the compressional state of the wedge, it shortens horizontally and  
300 thickens, uplifting the forearc high (Fig. 5C). Thus, the forearc basin “thermal blanket” promotes  
301 deep-seated deformation that, in the context of accretion, propels the uplift of the forearc high  
302 while subsidence continues in seaward regions that are unaffected by viscous flow (Figs. 5, 6 &  
303 7).

## 304 **5 Discussion**

305 Deep-seated flow (Pavlis and Bruhn, 1983) at the base of an orogenic wedge (Platt, 1986;  
306 Willett et al., 1993) provides a simple general framework to explain the dynamic formation of  
307 forearc highs. Crustal thickening by protracted wedge accretion increases the depth of burial  
308 and the temperature of the lower crust (Willett et al., 1993). Synorogenic sedimentation filling a  
309 forearc basin similarly contributes to wedge thickening, and also raises lower crustal  
310 temperatures by increasing thermal resistance and, if sediment conductivity is low, the  
311 geothermal gradient through the basin (Fuller et al., 2006).

## 312 **5.1 Dynamic growth of forearc highs**

313 The fundamental conditions required by the model presented here are quite simple. Accretion  
314 and syn-accretion sedimentation result in a progressive increase of crustal thickness in a  
315 subduction margin forearc. The increased thickness increases thermal resistance and Moho  
316 depth, thereby increasing lower crustal temperatures and changing deformation from frictional  
317 to viscous. Under protracted accretion and shortening, ductile strain in the lower crust switches  
318 vertical tectonic motions in the overlying wedge, from forearc basin subsidence to uplift,  
319 forming a new forearc high directly under the former basin (Fig. 8). We expect this process to  
320 take place in any accretionary system as it matures and increases in size, as for example  
321 documented on the Cascadia margin (Fuller et al., 2006; McNeill et al., 2000).

322 A viscous-flow controlled, structurally internal forearc high as we propose will form at a  
323 location determined by the geometry of the slab, as integral parts of accretionary wedges that  
324 uplift in a dynamic, non-linear manner. In this context, forearc highs become more probable as  
325 the accretion system matures and form at a time dictated by accretionary flux, wedge  
326 temperature, and wedge viscosity. We note, however, that there are other mechanisms for the  
327 formation of a forearc high, including forced mechanical accretion against areas of relatively  
328 larger strength (Byrne et al., 1993), changes in wedge taper or stress state (Willett &  
329 Schlunegger, 2010), or other processes leading to deep-seated ductility (Pavlis and Bruhn,  
330 1983), such as the presence of fluids.

331 Competing, dynamic effects control the uplift of the forearc high. Synorogenic sedimentation  
332 increases the thermal resistance and thus temperature of the underlying crust. If sediments  
333 have a low thermal conductivity, this effect is even more pronounced. While a weaker lower  
334 crust tends to decrease the wedge taper, it also facilitates ductile strain that, when sustained by  
335 accretion, results in increasing the wedge taper. Therefore, the taper geometry of the internal  
336 sectors of the wedge depends not only on convergence velocity (Willett et al., 1993) but also on  
337 its interplay with forearc basin sedimentation. The area undergoing lower crustal flow  
338 generates uplift in regions immediately above it, while trenchward regions not affected by

339 lower crustal flow continue to subside through sediment loading. The transitional area from  
340 forearc high to trenchward subsidence is regarded as ductile-to-frictional decollement in  
341 Williams et al., (1994) and as decoupled-to-coupled flow in Royden (1996).

342 Sedimentation controls the lag time between wedge growth and thermal activation of lower  
343 crustal flow. Whereas no relevant topographic growth occurs in the internal sectors of the  
344 wedge during its early evolution, regardless of the amount of sediments in the forearc, a  
345 threshold in the amount of sediment in the forearc controls the growth of the forearc high at  
346 advanced stages of wedge (Figs. 6, 7 & 8). This implies that the lower crustal flow is controlled  
347 by the thickness of the crust, which depends in turn on the sedimentation rates in the forearc  
348 basin. In other words, the sedimentary blanketing effect controls the lower crustal flow in so  
349 that the drop in viscosity does not occur until sediment thickness reaches a threshold.  
350 Otherwise, crustal thickening alone would eventually lead to viscous deformation, albeit  
351 occurring later and possibly elsewhere in the mechanical system.

352 Sedimentation and sedimentation rate have other effects that control forearc high uplift.  
353 Sediment infill of the forearc topographic depression reduces the surface angle of the forearc  
354 wedge to zero and stabilizes the wedge underneath, resulting in a broad wedge where active  
355 deformation is confined to an outer wedge removed from any material-controlled backstop  
356 (Fuller et al., 2006). In addition, isostatic basin subsidence by sediment loading of the forearc  
357 persists seaward and landward of the uplifting forearc high. This leads to the apparent paradox  
358 that regional subsidence in the forearc basin can control uplift in the forearc high for cases  
359 where sedimentation rate outpaces accommodation space, i.e. when subsidence controls the  
360 amount of sediment entering the system.

## 361 **5.2 Growth of the Anatolian margin**

362 Our simulations are consistent with SCAP formation as a dynamic, thermo-viscous forearc high  
363 led by forearc sedimentation and accretion along Central Cyprus (Fig. 8). Models reproduce the  
364 growth of the SCAP, including the surface uplift of Central Taurus and coeval subsidence in the



365 Cilicia Basin, resulting in the monoclinical flexure of Late Miocene rocks at the plateau margin  
366 scale (Fernández-Blanco et al., 2019) (Figs. 3 & 8). This is compatible with surface uplift onset  
367 between 8 and 5.45 Ma in the plateau margin hinterland (Cosentino et al., 2012), and a well-  
368 developed orogenic rain shadow by 5 Ma (Meijers et al., 2018) as well as the concomitant, short-  
369 wavelength vertical tectonic motions described for S Turkey (Fernández-Blanco et al., 2019;  
370 Walsh-Kennedy et al., 2014). Models are also coherent with strong observational evidence like  
371 the undisrupted and copious sedimentary record in the SCAP offshore thorough the latest  
372 Messinian - Recent times (Aksu et al., 2014, 2005; Walsh-Kennedy et al., 2014). Therefore, the  
373 mechanism of dynamic, thermo-viscous forearc high growth provides a physical support for  
374 models of SCAP growth by contraction and crustal thickening (Fernández-Blanco, 2014;  
375 Fernández-Blanco et al., 2019; Meijers et al., 2018).

376 Accelerated uplift rates during the uplift of the Central Taurides forearc can be inferred from  
377 the elevations of Miocene to Pleistocene marine rocks. These rocks show that the onset of uplift  
378 initiated at ~8 Ma (Cosentino et al., 2012) and accelerated to ~0.75 mm/yr over the last ~1.6 Ma  
379 (Schildgen et al., 2012). Our models show the occurrence of an equivalent accelerated uplift as a  
380 natural consequence of the non-linear uplift associated with thermal weakening of the lower  
381 crust (Fig. 5B,-C, 6, 7, 8B). Most observations of uplift rates are consistent with our models, but  
382 one study has suggested that 1200-1500 m of topographic growth occurred within the last ~450  
383 ka (Öğretmen et al., 2018). The spatial location of the marine deposits analysed by this study is  
384 relatively local, but if applicable to the entire margin would be difficult to reconcile with our  
385 models.

386 Primary alternative models for uplift of the Taurides forearc high suggest shallow slab break-  
387 off and multiple-phase surface uplift (e.g., Schildgen et al., 2014). However, slab break-off models  
388 do not inherently predict extended periods of uplift or accelerated uplift, so the earlier uplift  
389 during the Miocene has been attributed to structural thickening, similar to what we predict here  
390 (Schildgen et al., 2014) with the high rates of uplift associated with slab break-off in the last ~1.6  
391 Ma (Schildgen et al., 2012) or younger times (Öğretmen et al., 2018). Slab break-off is an

392 independent process to lower crustal viscous flow and it is possible that both have occurred and  
393 contribute to the uplift of S Turkey. However, our models show that there are combinations of  
394 parameters that predict growth rates, timing and accelerations consistent with the observations,  
395 with no need for an additional mechanism such as slab break-off.

396 Schildgen et al. (2012) argued that uplift rates exhibited a rapid increase in rate initiating at  
397 about ~1.6 Ma, on the basis of analysis of knickpoints in rivers draining the Central Taurides.  
398 However, the river knickpoints are not particularly uniform in their elevation or along-channel  
399 distribution, and could be consistent with a gradual increase in uplift rate from Late Miocene to  
400 the present, particularly if other complicating factors such as tilting and river capture are taken  
401 into account. Similarly, these studies interpret the depositional age or subaerial exposure of  
402 marine rocks at individual sites in terms of uplift age and uplift spatiotemporal pattern, i.e.  
403 temporal phases of uplift at margin scale. However, the uplifted marine sediment data cannot  
404 resolve sharp changes in uplift rate and only show that rates averaged over the Pliocene and  
405 Late Miocene were lower than rates over the Pleistocene to modern.

406 Our models show that compression-driven wedge growth with sedimentation can not only  
407 provide an uplift mechanism for the southern margin of the Central Anatolian Plateau but also  
408 reproduce first-order upper plate strain and the complex geometry and patterns of vertical  
409 motion in space and time that characterize the southern Anatolian margin (Figs. 3 & 8). These  
410 simulations agree with the regional frame and the geological and geophysical observables in the  
411 Anatolian margin along the Central Cyprus subduction zone, including in and under the area  
412 undergoing maximum uplift. The Kyrenia Range trench-slope break divides the active frictional  
413 deformation in the seaward areas, resulting in wedge top basin of Messaoria, from the negative-  
414 alpha Cilicia Basin, and from areas farther landward where thermally-activated viscosity in the  
415 deeper sectors of the wedge resulted in the uplift of the modern Central Taurides (Figs. 3, 5 &  
416 8). Our simulations could reproduce the dynamic growth of wider plateau-like terrains, but the  
417 mechanism presented here cannot be responsible for the topography of the entire Central  
418 Anatolian Plateau, given the thin crust in the plateau interior (e.g., Abgarmi et al., 2017). This

419 mechanism is also compatible with disruption of the former forearc basin by uplift of the  
420 forearc high in Cascadia (McNeill et al., 2000), and suggests similar processes in the Alaskan and  
421 Nankai accretionary margins (Pavlis and Bruhn, 1983).

## 422 **6 Conclusions**

423 Integration and interpretation of geophysical and geological evidence along the Anatolian  
424 subduction margin from the Central Cyprus trench to the SCAP suggests that lithospheric and  
425 crustal thicknesses, as well as Miocene and younger regional-scale structures and derived  
426 tectono-stratigraphic features formed by accretionary subduction. In this context, thermally-  
427 activated viscous flow of the lower crust is a physical mechanism of forearc high growth.  
428 Thermo-mechanical models of this process show that this single mechanism can explain much  
429 of the complex space and time pattern of vertical motions in the Anatolian subduction margin,  
430 with no need for an additional mechanism such as slab breakoff. We conclude that the plateau  
431 margin in South Turkey, and areas with a similar sequence of vertical motions in the interior of  
432 other accreting subduction wedges, grow as dynamic, thermo-viscous forearc highs.

433

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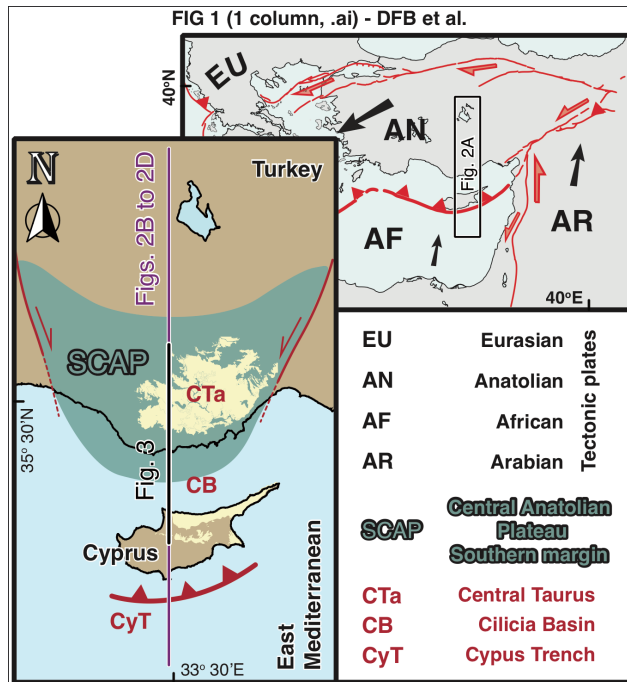
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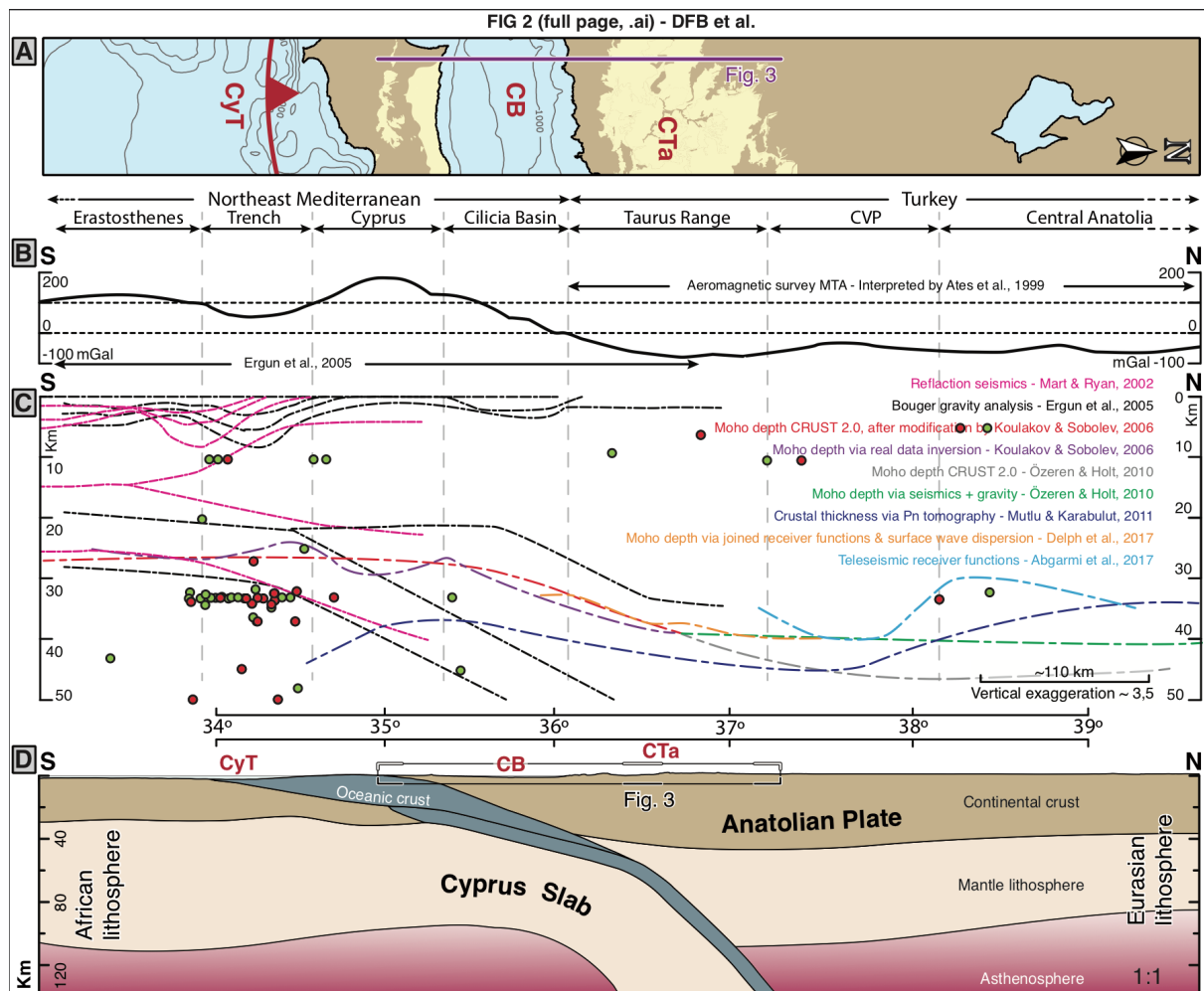
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## Figures

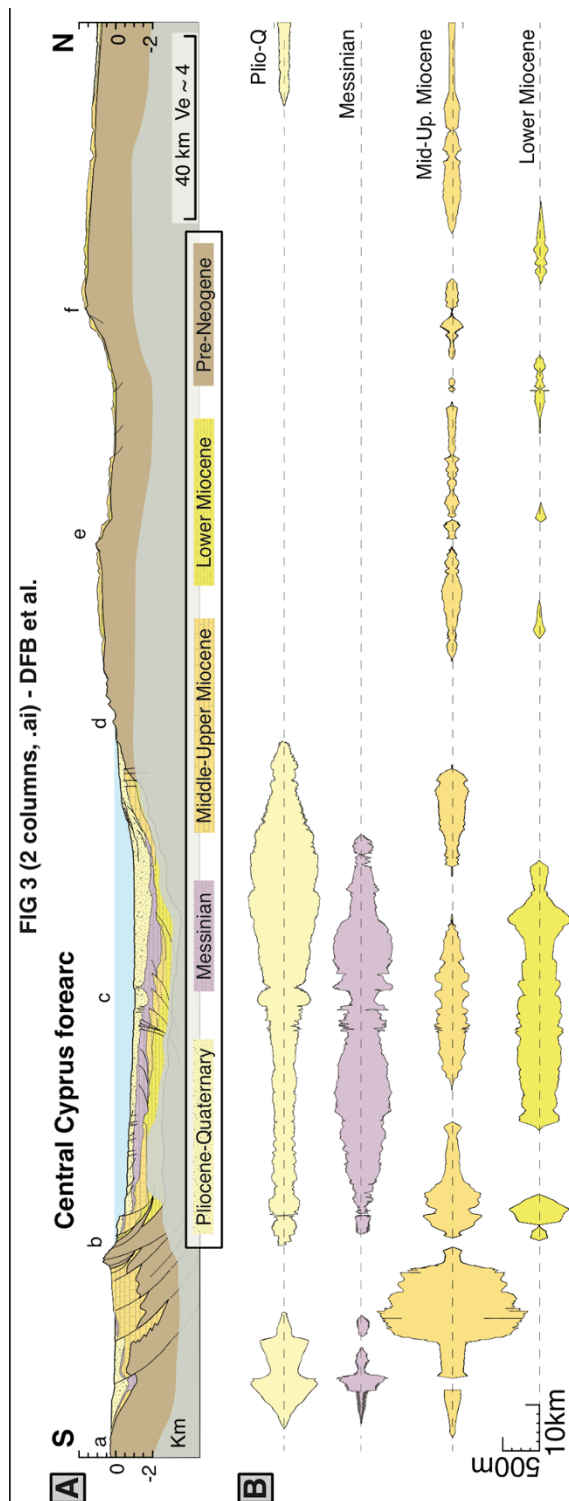


**Figure 1.** Central Cyprus subduction zone and its surrounding tectonic frame, with the location of transects in figures 2 and 3. Neogene rocks on Central Taurides and Cyprus are in yellow. EU, AN, AF and AR = Eurasian, Anatolian, African and Arabian tectonic plates; SCAP = Southern margin of the Central Anatolian Plateau; CTA = Central Taurides; CB = Cilicia Basin; CyT = Cyprus trench.



**Figure 2. Plate-scale transect.** (A) Map view of a  $2^\circ$ -longitude wide ( $32^\circ30' \text{ E}$  to  $34^\circ30' \text{ E}$ ) swath running  $\sim 650 \text{ km}$  along latitude, as a reference for data along transects in panels (B) and (C). (B) Values along the section of interest derived from the two major gravimetric studies in the area (Ates et al., 1999; Ergun et al., 2005). (C) Published geophysical data, including the interpretation of the offshore section C in Ergun et al. (2005), and that of the seismic study performed by Mart & Ryan (2002). The plot also includes several cross-sectional values retrieved from maps of Moho depth models derived from different geophysical approaches, including Pn tomography and receiver functions (Abgarmi et al., 2017; Delph et al., 2017; Koulakov and Sobolev, 2006; Mutlu and Karabulut 2011; Özeren and Holt, 2010). The circles are focal epicenters with  $M_w > 5$  recorded in a longitudinal area from  $32^\circ30' \text{ E}$  to  $34^\circ \text{ E}$  in red and  $31^\circ30' \text{ E}$  to  $34^\circ30' \text{ E}$  in green. (D) Lithospheric-scale transect along the Central Cyprus subduction (for  $\sim 650 \text{ km}$  at  $33^\circ30' \text{ E}$ ) at real scale derived from the interpretation of the data shown in (B) and (C).





**Figure 3. Uppermost crustal transect. (A)** Geologic transect along the Central Cyprus forearc (for ~300 km at 33°30' E), exaggerated ~4 times in the vertical. Letters “a” to “f” show the approximate location of structural highs bounding basinal sectors with similar lengths. See main text for data used and interpretation. **(B)** Thicknesses of main stratigraphical units derived from the transect and their age.

FIG 4 (1.5 column, .ai) - DFB et al.

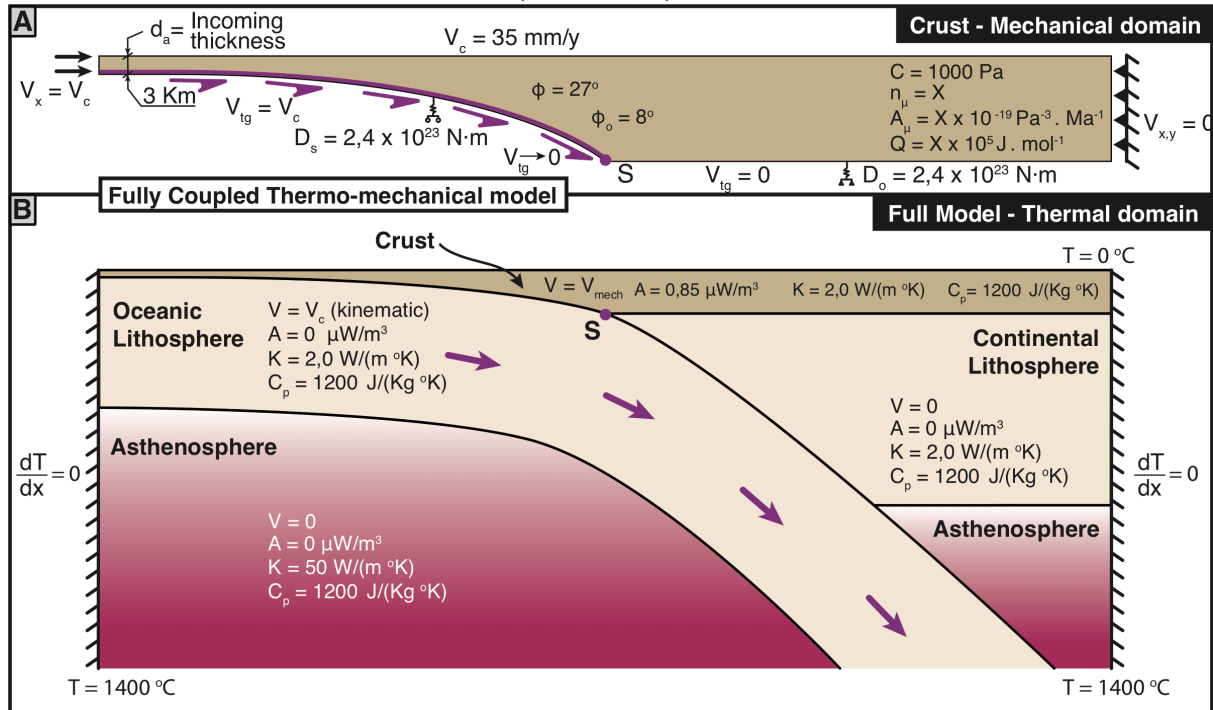
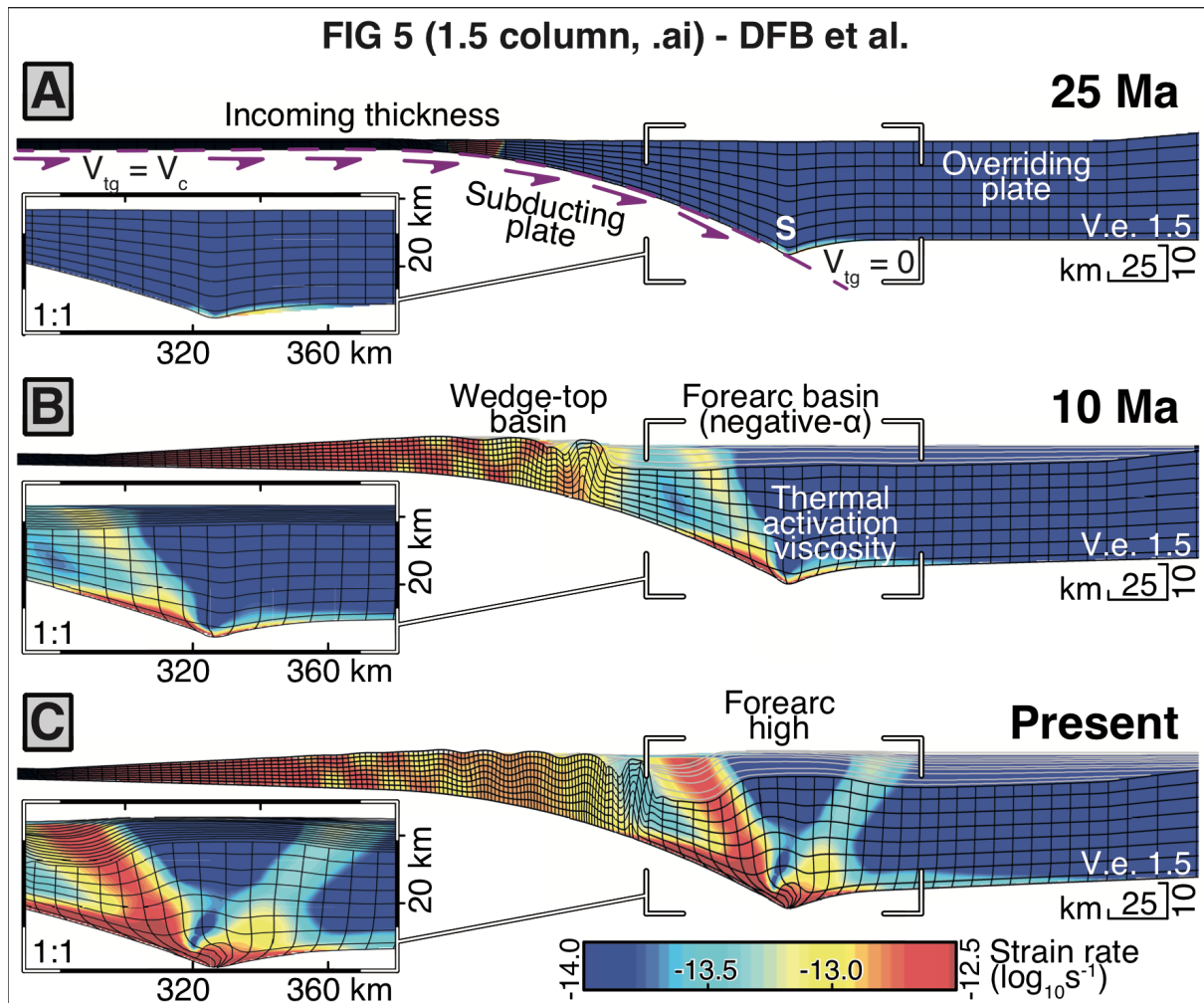
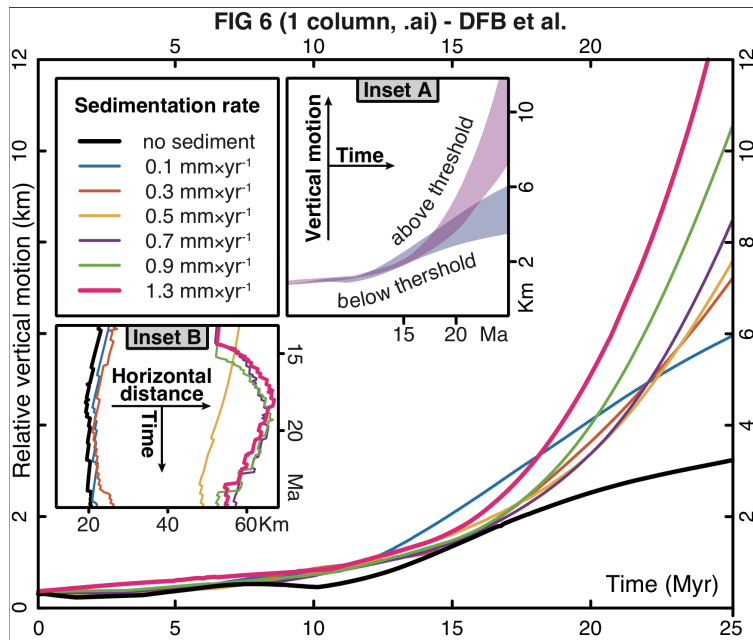


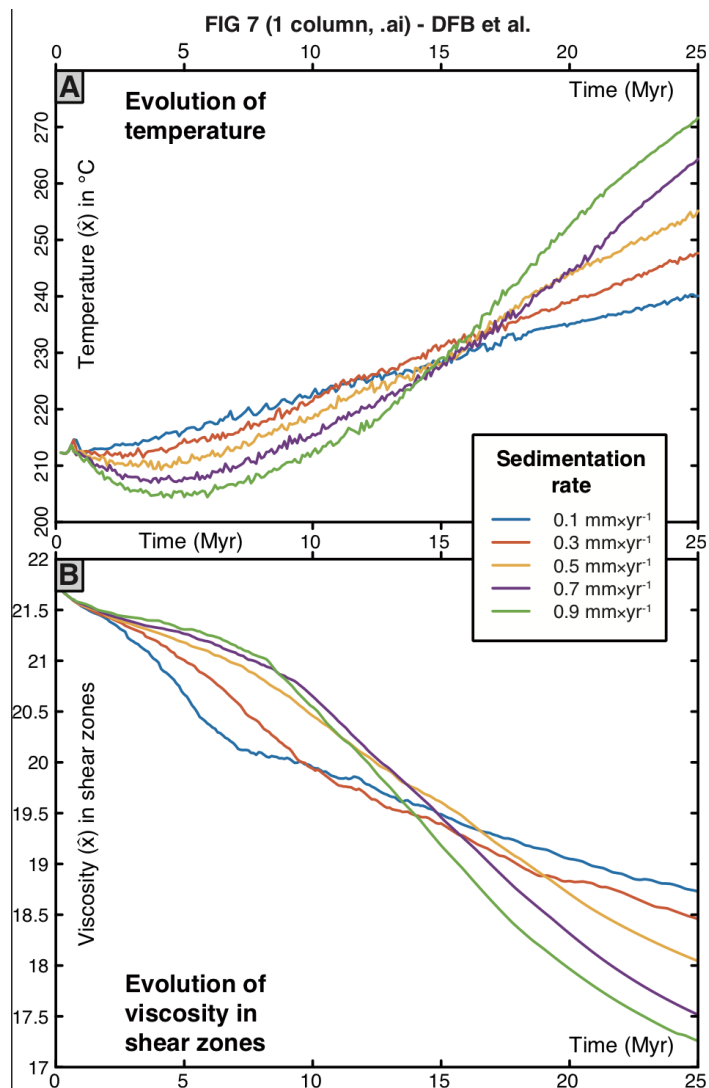
Figure 4. Model setup, with an indication of mechanical and thermal parameters.



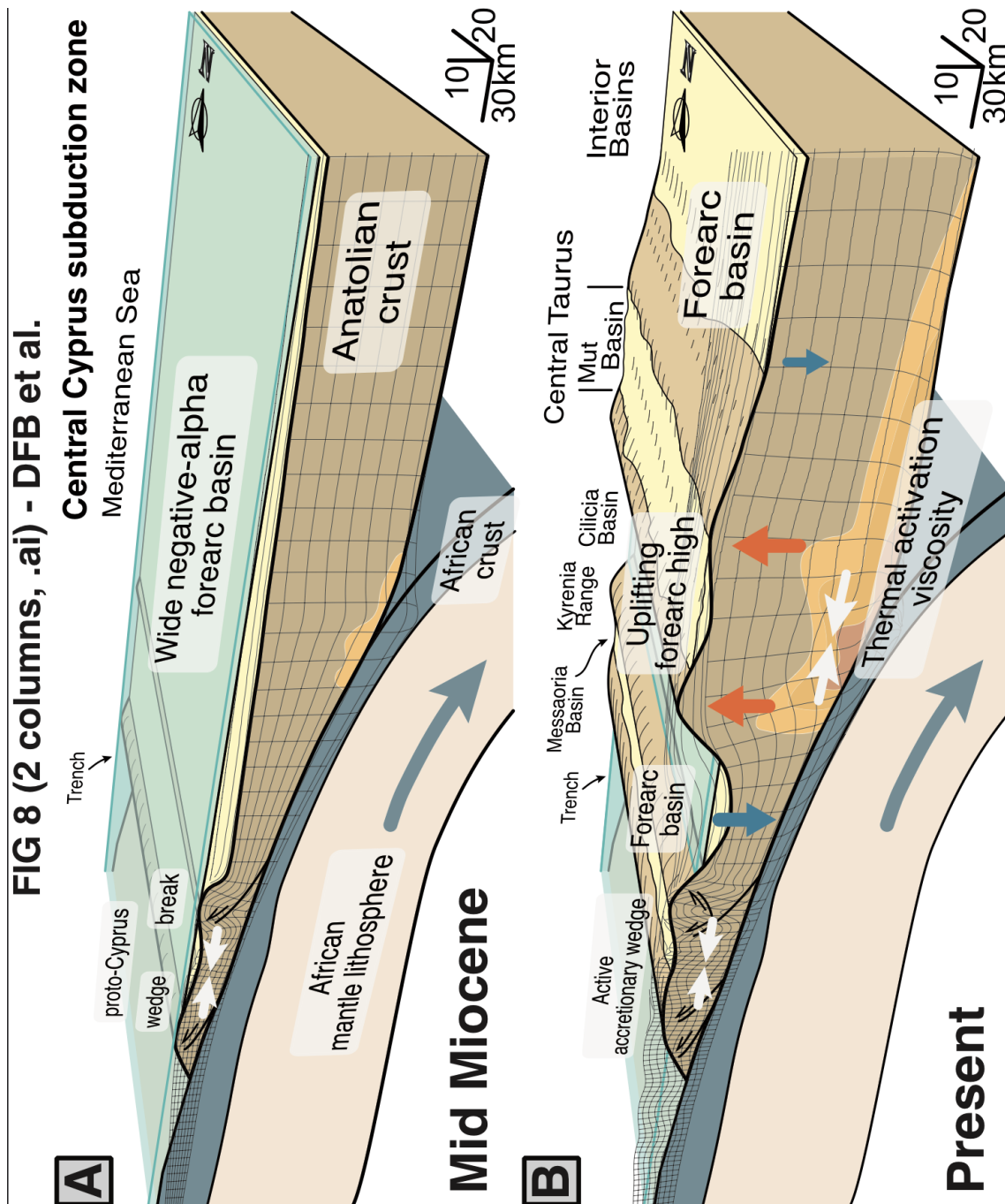
**Figure 5.** Three time-steps (initiation, ~10 Ma and present) in the numerical thermo-mechanical model of viscous-plastic deformation for a convergent wedge undergoing basal traction to simulate subduction (see DR3). Insets are zoom-ins of the area where the forearc high develops at the model scale. The color shows the second invariant of strain rate tensor, i.e. strain rate. The cumulative strain is shown by the Lagrangian mesh. Individual lines on top of the basement are isochrones (synthetic stratigraphy) that reflect the overall geometric relationships expected for strata. The upper panel represents the moment of subduction initiation at 25 Ma. The middle panel at 10 Ma shows a wide forearc (negative- $\alpha$ ) basin, increase in strain rate below it and wedge deformation reaching the future location of the forearc high. The bottom panel at present shows a forearc high developed by a wide shear zone and onlaps in the synthetic stratigraphy above large strain rates at the lower crust. The model does not include erosion, so sediments draping the forearc high are still present, but would be expected to erode quickly, were this process included. See video Video S1 in Supplementary Material.



**Figure 6.** Sedimentation rate influence on the time evolution of vertical motions in a forearc high. Each colored line shows the height difference between the highest and lowest basement elevation in a 100 km strip centered around the S point and above it. A total of 250 height-difference data points (one every  $10^5$  yr) compound each line of sedimentation rate, from a simulation with no sedimentation rate (black) to rates of 1.3 mm/yr (magenta). Inset A is a schematic representation of vertical motions in time for the groups below and above the threshold in sedimentation rates. Inset B is the evolution in time of the horizontal distance between the highest and the lowest point in the basement surface for the last 12 Ma of model run.



**Figure 7.** Evolution of median temperature (A) and median viscosity in shear zones (B) for models with different sedimentation rate, calculated every 100 kyr (250 snapshots). Different line colors represent sedimentation rates of 0.1 mm/yr (blue), 0.3 mm/yr (red), 0.5 mm/yr (yellow), 0.7 mm/yr (purple), and 0.9 mm/yr (green).



**Figure 8.** Box model representation of the mechanism of thermo-viscous forearc high uplift in two time-steps. Boxes show the evolution and forearc elements for a generic subduction wedge with forearc high and for the Central Cyprus margin. Time steps are interpreted as representative of the Central Cyprus margin. Integrated deformation is shown using the Lagrangian mesh of the model. Note that the model is only two-dimensional.