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# 1 The Ediacaran Grenville dykes (SE Canada) reveal the weakest

# 2 sustained palaeomagnetic field on record.

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# 13 Abstract

Long-term variations of the geomagnetic field, observed in the palaeomagnetic record, have 14 the potential to shed much light on the evolution of Earth's deep interior. With a geomagnetic 15 field characterised by anomalous directions and ultra-low intensities, the Ediacaran period 16 17 (635-541 Ma) is a time of special interest. Steep and shallow directions, leading to virtual geomagnetic poles (VGPs), separated by angles of up to 90° and very close in age could have 18 19 recorded a geomagnetic field switching between axial and equatorial dipole-dominated 20 states. Alternatively, the field may simply have been highly nondipolar and subject to rapid reversals. Palaeointensity determinations of units that record the anomalous directions could 21 potentially help to discriminate between morphologies but the spatial and temporal 22

distribution of palaeomagnetic data require improvement. Here we present new 23 palaeointensities from 11 sites from the western end of the Grenville Dyke swarm that 24 recorded directionally anomalous geomagnetic fields around ~585 Ma. Palaeointensities, 25 obtained through microwave Thellier, Shaw and pseudo-Thellier methods, show field 26 27 strength values of 2.9 $\pm$ 2.2  $\mu$ T and corresponding virtual dipole moments of 0.3-1.7 x10<sup>22</sup> Am<sup>2</sup>. These field strengths are an order of magnitude weaker than the present-day field. The most 28 extreme palaeointensity values of 1.4-2.1 µT are half as strong as seen in previous studies of 29 30 the Ediacaran field and as low as Mars' recently measured crustal field intensity, giving a new lower bound for the Earth. VGPs grouping in two distinct clusters with almost identical angular 31 dispersions of VGPs ( $S_B$  = 18.5° and 18.9°) may argue for the presence of an equatorial dipole. 32 33 In contrast, the palaeointensities associated with the steep and shallow components are indistinguishable. This observation, together with the overall very large VGP dispersion may 34 35 rather support that the Grenville Dykes have recorded enhanced secular variation linked to a 36 highly unstable, multipolar and reversing field.

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### 38 **1. Introduction**

Extreme climatic changes in the late Neoproterozoic (Evans and Raub, 2011) and tectonic activity between the final breakup of Rodinia and the assembly of Gondwana have made the Ediacaran period (542-635 Ma) a prime target for palaeomagnetic and geodynamic research. Especially the individual drifting histories of Laurentia and Baltica remain mysterious (Li et al., 2013) although both their positions are relatively well constrained around the beginning and the end of the Ediacaran at ~615 Ma and from ~550 Ma (e.g McCausland et al., 2007; Meert, 2014). Several sets of palaeomagnetic data are available to constrain the palaeogeography

for Laurentia (e.g. Bono and Tarduno, 2014; McCausland et al., 2007; McCausland and 46 Hodych, 1998) and Baltica (e.g. Bazhenov et al., 2016; Meert, 2014), but the resulting 47 palaeopoles lead to ambiguous apparent polar wander paths (AWPW). Datasets from 48 Laurentia and Baltica show both steep and shallow primary directional components that are 49 50 extremely close in age. For Laurentia, both components were found in studies of the Sept-Îles 51 intrusion (Bono and Tarduno, 2014; Tanczyk et al., 1987), the Grenville Dykes (Murthy, 1971, 52 Halls et al., 2015) and the Catoctin Basalts (Meert et al., 1994). These components usually 53 show dual-polarities and are often supported by field tests. The limitations imposed by the nature of a geocentric axial dipole field lead to discussions about which of the two 54 components of each study could be less reliable and would not be used for the calculation of 55 the APWP (e.g. Bono and Tarduno, 2014; McCausland et al., 2007; Robert et al., 2017). 56 Depending on which of the seemingly primary components are rejected, high plate motion 57 58 velocities are required to allow the latitudinal changes in the resulting APWPs which lead to 59 different theories explaining the geodynamic processes in the Ediacaran, including inertial interchange true polar wander (Kirschvink et al., 1997). In more recent studies, the puzzling 60 61 APWPs have been attributed to anomalous directional behaviour of the geomagnetic field itself. Specifically, the steep and shallow directions have been interpreted as being caused by 62 the geomagnetic field flipping between an axial and an equatorial dipole field configuration 63 64 (Abrajevitch and Van der Voo, 2010) or as showing an intermediary state of an axial dipole 65 field during continuous reversals (Halls et al., 2015). Numerical geodynamo simulations have suggested that the occurrence of a protracted equatorial dipole field may be possible during 66 the reversal of an axial dipole field (Aubert and Wicht, 2004; Gissinger et al., 2012) with 67 significant differences between the intensities of the axial and equatorial field states. Studies 68 69 of magnetostratigraphic data from the late Ediacaran and the early-mid Cambrian (Bazhenov

et al., 2016), resulting in reversal frequencies of more than 20 reversals/ Ma, claim that the 70 field might have been in a hyperactive state at that time. Due to the correlation between 71 reversal frequency and dipole field strength (Kulakov et al., 2019; Tauxe et al., 2013) the 72 dipole field strength in the Ediacaran is expected to be low. First palaeointensity studies of 73 74 Ediacaran rocks from Canada (Bono et al., 2019) and Ukraine (Shcherbakova et al., 2020) show 75 ultra-low virtual dipole moments (VDM) that are an order of magnitude weaker than dipole 76 moments in the Phanerozoic. To date, similarly low dipole moments have only been found in 77 the Devonian (Hawkins et al., 2019; Shcherbakova et al., 2017) and the Jurassic (Kulakov et al., 2019; Tauxe et al., 2013). Information about reversals and palaeointensities are critical to 78 79 delineate the anomalous field behaviour in the Ediacaran. However, no estimates for the strength of the geomagnetic field exist for the early-mid Ediacaran before 580 Ma. Here we 80 report multi-method palaeointensity measurements performed on mid-Ediacaran age units 81 82 showing both steep and shallow directions to look for differences that could help explain the 83 directional observations of that time.

## 84 2. Materials and methods

#### 85 2.1. Sample material:

The samples from dykes within the Grenville province used in this study were those collected for the study of Halls et al. (2015) and were taken preferentially from chilled margins. Several dykes were sampled at multiple sites over distances of up to 150 km. Analysis of geochemical composition helped with longitudinal correlation of the dykes. Samples were selected for intensity determination if sister-samples from the same site showed a stable component of ChRM in the previous study, which was the case for 99 samples from 15 sites within nine dykes (Figure 1). Following the naming convention used in Halls et al. (2015), the five dykes

analysed in this study were: Coniston dyke (sites GD02, GD33), French River dyke (site GD23),
Key River dyke (sites GD10, GD15, GD16, GD19, GD37) and Sand Bay dyke (site GD29).
Additionally, four sites from other dykes in the area associated with the dyke swarm (sites
GD07, GD25, GD26, GD30) were also analysed. U-Pb ages exist for Sand Bay dyke (585.2±0.8
Ma), Augusta Lake dyke (584±0.6 Ma), Key River dyke (587.3±0.7 Ma and French River dyke
(598.0±1.4 Ma).

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101 Distances between sites of the same dyke were generally larger than 2km and the sites were 102 therefore regarded as separate spot readings of the field strength for the purpose of this study. Differences in chemical composition between the individual dykes suggest single 103 104 intrusion events at different development stages of the magma chamber (Halls et al., 2015). 105 The directional results in Halls et al. (2015) show four components resulting in antipodal steep 106 and shallow directions close to directions of older studies of the Grenville dykes (Hyodo and 107 Dunlop, 1993; Murthy, 1971). Positive contact (Halls et al., 2015) and reversal tests (Hyodo and Dunlop, 1993) support the primary nature of the steep direction. This comprises 108 109 components B and E, that can be seen in sites of the Key River dyke (Halls et al., 2015) which are close to the directions recorded by the Mattawa dyke (585.9 Ma, Hyodo and Dunlop, 110 111 1993). The shallow component from the almost antipolar components C and D can also be 112 observed continuously in sites of the same dyke and component D differs significantly from 113 the regional overprint recorded in the close-by Whitestone anorthosite (Ueno et al., 1975). 114 Because of an inconclusive contact test, it cannot be excluded that this component was remagnetised (Halls et al., 2015; Robert et al., 2017). In this study, we nevertheless analysed 115 samples from sites that recorded both steep and shallow components. 116

#### 117 **2.2 Methods**:

All measurements were carried out at University of Liverpool's Geomagnetism Laboratory. At 118 119 least one specimen per site was selected for rock magnetic measurements. Measurements of 120 IRM, hysteresis parameters and thermomagnetic curves were done on a Variable Field 121 Translation Balance (VFTB) using both crushed specimens and cylinder specimens with a 122 diameter of 5 mm. Additionally, temperature dependent susceptibilities were measured on an AGICO MFK-1A Kappabridge using crushed specimens heated in air to 700°C. To monitor 123 124 thermochemical alterations in more detail, the same measurements were also performed in 125 incremental heating and cooling cycles in temperature steps of 100°C between 150°C and 650°C. Scanning electron microscope (SEM) images were taken with a Hitachi TM3000 126 tabletop microscope and analysed with Quantax70. 127

The magnetic viscosity of 59 specimens was determined by measuring their remanent magnetisation after being stored for three weeks in different orientations without shielding from the ambient field (Prévot, 1981). For this experiment, the specimens were stored on a rack in the laboratory with their z axis parallel to the field lines of the geomagnetic field for the first duration and antiparallel to the field lines for the second duration. The remanent magnetisations were measured on an AGICO JR6 spinner magnetometer.

134 Four different methods were used to acquire palaeointensity data:

Microwave Thellier type palaeointensity experiments were performed on the high frequency (14GHz) microwave system and SQUID magnetometer (Hill and Shaw, 1999). This instrument uses cylinder specimens with a 5mm diameter. These were drilled from the available 1" cylinders, which allowed for 5 to 15 sub-specimens to be drilled from one standard cylinder specimen. Specimens are subjected to microwave treatment with incrementally increasing

power and/or duration in a range between 20W.s and 500 W.s. The succession of zero-field (Z) and in-field (I) steps for the experiments followed the IZZI protocol (Tauxe and Staudigel, 2004). A pTRM-check was done after every other I-Z-pair, resulting in 4-5 pTRM-checks per specimen. To detect multidomain effects, the laboratory field was applied in angles of 45°-90° to the ChRM of the specimen (Yu and Tauxe, 2005). The laboratory field itself was varied between 3 and 30  $\mu$ T between experiments to detect possible field dependencies of the results.

147 Thermal Thellier experiments were undertaken using the IZZI protocol with pTRM checks in 148 vacuum, using an MMTD80 oven with a bias field of 10  $\mu$ T in the temperature range of 100°C-149 550°C. Magnetisations of the specimens were measured on an AGICO JR6 spinner 150 magnetometer after being 'cleaned' with a 2 mT AF step before every measurement using an 151 AGICO LDA5 AF demagnetiser.

The double-heating Shaw method of (Tsunakawa and Shaw, 1994) used AF demagnetisations and ARM acquisitions carried out on a 2G RAPID superconducting rock magnetometer system in incremental steps of 2-10 mT up to a peak AF-field of 100 mT. ARMs were given at peak AFfield in bias fields of 57.9  $\mu$ T or 81.2  $\mu$ T. The specimens were given a full TRM twice by heating them to 610° or 650°C in bias fields of 5  $\mu$ T or 10  $\mu$ T to be able to apply an ARM correction (Rolph and Shaw, 1985).

The pseudo-Thellier method (Paterson et al., 2016; Tauxe et al., 1995) was performed on the automatic 2G RAPID system with the same AF levels as used in the Shaw method. This was done to have results that avoided heating the specimen at all during palaeointensity determination but its results were only used to confirm the results from the other methods. For the purpose of this experiment, bias fields of 11.4  $\mu$ T and 81.2  $\mu$ T were used for the ARM

acquisition and, in some cases, Shaw and pseudo-Thellier experiments were run as a
 combined experiment, where the ARM acquisition steps were performed between the NRMO
 and ARMO steps of the Shaw experiments.

Different sets of selection criteria have been used to assess the quality of the palaeointensityexperiments using the different methods.

For all Thellier-type experiments, the selection criteria followed the Standard Palaeointensity Definition (SPD, (Paterson et al., 2014) and were modified from the SELCRIT2 criteria (Biggin et al., 2007). The criteria of N=4, FRAC  $\ge$  0.35,  $\beta \le$  0.1,  $q \ge 1$ , MAD  $\le$  15°,  $\alpha \le$  15°, DRAT  $\le$  10% and CDRAT  $\le$  10% were used together with a curvature factor if the best-fit line (Paterson, 2011) of  $|\mathbf{k}'| \le$  0.48.

The selection criteria used for the Shaw method were similar to those set out by Yamamoto et al., (2010),: number of consecutive data points N  $\geq$  5, correlation coefficients of the linear parts of the NRM-TRM1\* diagram and the TRM1-TRM2\* diagram of  $r^2_N \geq 0.995$  and  $r^2_T \geq$ 0.995, the fraction of used NRM f  $\geq$  0.2 and the slope of the linear part of the TRM1-TRM2\* diagram 0.95  $\leq$  slope<sub>T</sub>  $\leq$  1.05. In addition, the selected part of the NRM must appear convergent on the orthographic plot of the NRM demagnetisation with  $\alpha$  and MAD values  $\leq$ 10°.

The SPD selection criteria used for pseudo-Thellier experiments were slightly relaxed from Paterson et al., (2016) and applied to a convergent part of NRM in the orthographic plot:  $N \ge$ 6, f  $\ge$  0.3, ß  $\le$  0.1, r<sup>2</sup>corr  $\ge$  0.990, f<sub>resid</sub>  $\le$  0.15,  $\alpha \le$  10°, MAD  $\le$  10°, |k'|  $\le$  0.27 and 0.85  $\le$  |b<sub>AA</sub>|  $\le$  1.15.

Results that passed all criteria were classified as 'class A' result. We found that by relaxing a single criterion to a less strict value, we could dramatically increase success rates without severely affecting the overall result quality. Results that required a single criterion to be relaxed were classified as 'class B' results. Limits for relaxed selection criteria were: FRAC  $\geq$ 0.25, DRAT  $\leq$  15% or CDRAT  $\leq$  15% for Thellier-type experiments and  $r^2_N \geq$  0.990,  $r^2_T \geq$  0.990 or 0.90  $\leq$  slope<sub>T</sub>  $\leq$  1.10 for Shaw experiments. If more than one criterion would have to be relaxed for a result to pass, the estimate was rejected.

191 **3 Results** 

192 Thermomagnetic curves and susceptibility versus temperature curves measured on samples 193 from different dykes showed quite diverse behaviour (supplementary figure 1). Most specimens showed one or two Curie temperatures between 550°C and 580°C. Specimens 194 from the French River and Key River dykes as well as specimens from sites GD25 and GD26 195 also had noticeable Hopkinson peaks, indicating fine magnetite. Exceptions were 196 197 measurements on specimens from site GD01 with a Curie temperature of 470°C and from 198 sites GD07 and GD30, where susceptibilities were about two orders of magnitude weaker. With the exception of these specimens and those from Coniston dyke (GD01) and site GD33 199 200 of Augusta Lake dyke, thermomagnetic curves were mostly reversible. The susceptibility curves of specimens from Coniston (GD01, 14), French River (GD23) and Sand Bay (GD29) 201 dykes showed a small 'toe' with another Curie temperature up to 610-620°C, possibly 202 203 suggesting the presence of a small fraction of maghemite or titanohematite. Backscattered 204 electron imaging (figure 2a-c) shows mostly relatively fresh coarse (~30-200 μm) (Ti-)magnetite grains with sharp edges as well as small dendritic (< ~40 µm) TM grains (e.g. figure 205 2a,b), indicating rapid cooling of the rocks. EDS elemental composition mapping of larger 206

207 grains shows structures of full and partial exsolutions with lamellae of titanium rich ilmenite 208 and titanium poor magnetite (figure 2c, supplementary figure 2). The preservation of high 209 temperature textures and lack of low-temperature oxidation structures suggest that the 210 grains are carriers of a primary TRM. For some samples it cannot be ruled out that the oxy-211 exsolution continued to temperatures lower than the Curie temperatures, but this has shown 212 to only have small effects on the resulting palaeointensities (Shcherbakov et al., 2019).

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Hysteresis parameters, shown in Figure 2e) and f), fall in two clusters on the Day et al., (1977) plot with one large cluster in the "PSD" area and another cluster with higher Bcr/Bc ratios. Most specimens plot close to the TMO line (Wang and Van der Voo, 2004). This includes results from the French River dyke, which showed a substantially higher TiO<sub>2</sub> content in the geochemical analysis of the sites in Halls et al. (2015). In contrast, specimens from the low-Titype Coniston dyke (Halls et al., 2015) plot further away from the TMO and closer to the TM60 line.

Fractions of viscous remanent magnetisations (VRM) of specimens from the different sites, 221 acquired after 3 weeks in the ambient field, did not show any differences at the dyke level. 222 223 The average viscosity index (v) for all sites (figure 2 d) was 14.4% (median = 9.6%, standard 224 deviation = 14.9%). This mean viscosity index includes v values of up to 80% from sites GD07 225 and GD30. Specimens from those sites were basically non-magnetic and the viscosity index was mostly calculated from noise. Specimens with v values of >30% were not used for 226 intensity determination. We note that low palaeointensities will naturally produce weak 227 228 NRMs which will then be more radically affected by a VRM acquired in the strong present-day field. 229

Results of palaeointensity experiments are summarised in table 1. The total number of Shaw 230 231 experiments only accounts for specimens that were retained for the full routine. The experiment was stopped after the first demagnetisation series (NRMO) for 16 specimens 232 whose ChRM could not be isolated. Microwave demagnetisation experiments were 233 performed on 29 sister specimens to determine their response to microwave treatments. 234 Specimens from the same core as specimens that did not respond well to the microwave 235 236 treatment during demagnetisation, or were too weak to be demagnetised at all, were not 237 used for intensity experiments with the microwave Thellier method. In total, palaeointensity measurement was attempted on 189 specimens of the Grenville dykes. Of these, 128 (68%) 238 239 used microwave IZZI, 20 (11%) used thermal IZZI, and 41 (21%) used Shaw and/or pseudo-Thellier experiments (Table 1). 240

Specimens universally showed non-ideal behaviour during thermal Thellier experiments.
These mostly resulted in chaotic Arai and orthographic plots (figure 3e) and the experiment
was stopped at 550°C. Strong 'zig-zagging' and high β values suggest strong MD behaviour
over the full temperature range. In contrast, pTRM checks passed the criteria up to the highest
temperature steps.

The majority of Arai plots from all Thellier experiments (75%) and of passing Thellier results (94%) showed two slopes as illustrated in figure 3. Similar to the results of other studies yielding extremely low palaeointensities (Hawkins et al., 2019; Shcherbakova et al., 2020), the sharp bend between the high and low temperature/power slopes usually coincided with the junction between ChRM and secondary components in the associated Zijderveld plots (figure 3 a,c), but this is was not clearly visible for all results (figure 3 b). Of the accepted microwave Thellier type results, 10 passed as class A and 15 as class B result. For all class B

results, either DRAT or CDRAT (but not both) had to be relaxed. About 10% of the results that 253 254 were not accepted failed only one of the selection criteria. All other non-accepted results failed more than one selection criterion (e.g. figure 3e). The most commonly failed criteria 255 were  $\beta$  (57% of results), FRAC (56%) and DRAT/CDRAT (48/52%). The high failure rate due to 256  $\beta$  and DRAT/CDRAT shows that MD effects and mineral alterations during heatings were both 257 primary issues in the results of thermal and microwave experiments. Of the low 258 temperature/power slopes, 6 results (5%) would pass the selection criteria, but it is clearly 259 260 visible in the orthographic plot that the selected fraction is not the ChRM (figure 3d). One outlier intensity from a high-power slope of 26.7 µT was rejected as unreliable even though 261 262 the result passed the selection criteria (GD19-43B1). Here, sister specimens from the same 1" core gave inconsistent intensities ranging from 3.6 to 140  $\mu$ T. 263

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265 Of the 12 accepted Shaw results, 7 passed as class A and 5 as class B results. For all 7 class B 266 results, the correlation coefficient of the best fit line in the NRM-TRM1\* plot r<sup>2</sup><sub>N</sub> was relaxed 267 to  $r_N^2 \ge 0.990$  (figure 4a). Results not passing  $r_N^2$  and/or slope was the main reason for unsuccessful experiments. Specimens showed a wide range of coercivities with a mean MDF 268 = 19.2 ±9.6 mT. Secondary magnetisations that accounted for NRM fractions of up to 80% 269 were usually removed by AF fields of less than 40 mT. Similar to the AF demagnetisation 270 271 experiments in Halls et al. (2015), that successfully generated reliable directions and passed 272 field tests, the remaining fraction of primary NRM was large enough in the Shaw experiments 273 as well so that only two specimens failed the FRAC criterion.

From the 7 pseudo-Thellier and the 21 Shaw/pseudo-Thellier experiments 5 pseudo-Thellier
results passed selection criteria and were accepted. Due to the high uncertainty of pseudo-

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Thellier intensities associated with the used calibration factor (Paterson et al., 2016), the 277 278 resulting palaeointensities were not considered in the calculation of the mean palaeointensities. However, since thermal alteration of the specimens during experiments 279 280 was one of the main issues in the other methods, the pseudo-Thellier results can be useful to 281 confirm that the results from the other methods are not biased by such alterations. Results from successful pseudo-Thellier experiments showed good agreement with the site mean 282 values calculated from microwave and Shaw results (Table 1) with the mean values of sites 283 GD33, GD29 and GD26 being within the 25% uncertainty range of the associated pseudo-284 Thellier results. For sites GD26 and GD37 that show a larger difference (>1 µT) between 285 intensities from pseudo-Thellier experiments and intensities from other methods, the 286 287 palaeointensities from pseudo-Thellier experiments are lower than the palaeointensities from the other methods. Most of the unsuccessful pseudo-Thellier experiments failed due to 288 289 a high curvature of the selected line segment in the demag-demag plot (figure 4b), suggesting that NRM and ARM (de)magnetisations behave differently in these specimens and a 290 palaeointensity calculated from these would be unreliable. 291

292 Combining class A and class B results, 25 results of the microwave Thellier experiments (20% 293 success rate) and 12 results of the Shaw experiments (35%) passed the used selection criteria. 294 No result of the thermal Thellier experiments passed selection criteria, leading to an overall 295 success rate of 20% for Shaw and Thellier-type experiments. Palaeointensity estimates were 296 averaged for each site as separate spot readings of the field strength (Table 1). The resulting 297 intensity values range from 1.4  $\mu$ T to 7.6  $\mu$ T, and yielded virtual dipole moments (VDM) from 298 0.3 to 1.7 \*10<sup>22</sup> Am<sup>2</sup>. Five pseudo-Thellier results passed selection criteria as well and gave intensity values between 1.1  $\mu$ T and 4.1  $\mu$ T. All intensity results with critical values are listed in supplementary tables 1-3 and measurement data of all intensity experiments are available on the MagIC database (www2.earthref.org/MagIC).

302 4 Discussion

The geomagnetic field in the Ediacaran is characterised by ambiguous field directions that impede the construction of reliable APW paths. Recently, it has been stated that the time averaged dipole field strength is uniquely low in the Ediacaran (Bono et al., 2019) which is supported by the only other palaeointensity study of Ediacaran rocks to date (Shcherbakova et al., 2020). We expand the intensity data from these studies with the whole rock palaeointensities of the Grenville dykes with virtual dipole moments between 0.3 \*10<sup>22</sup>Am<sup>2</sup> and 1.7 \*10<sup>22</sup>Am<sup>2</sup>.

310 Due to the low number of available samples and non-ideal behaviour during thermal Thellier experiments, microwave experiments were preferred as the small sample size allowed for a 311 312 high number of experiments with sub-specimens. To account for a possible method-bias of the MW results, a number of thermal Thellier, Shaw and pseudo Thellier experiments were 313 314 conducted as well. All methods suffered from low success rates that made it hard to compare 315 the different methods on the site level. Where such a comparison was possible, the results from different methods differed by less than 1 µT, similar to the overall intensity variability 316 on the dyke level. From this, a method related systematic bias seems unlikely. Furthermore, 317 318 results from Thellier and Shaw experiments are also supported by the results of the pseudo-319 Thellier experiments.

To confirm that the intensity results are not biased by the choice of selection criteria, other commonly used sets of selection criteria have been used to re-analyse all Thellier-type results.

Accepted results from the different sets of criteria have been compared (supplementary table 4). All sets show only small differences in the range of ~1  $\mu$ T and agree well with the results of the Shaw experiments. Supplementary table 4 also shows that results from combining class A and class B results only differ from pure class A results by 0.2  $\mu$ T and are well inside the standard deviation of class A results, while the number of accepted results almost doubled.

327 The two-slope behaviour, seen in most Arai plots of Thellier experiments in this study, has been recognised in other studies of basalts with low palaeointensities as well (Hawkins et al., 328 329 2019; Kodama et al., 2019). Similar behaviour is often connected to MD behaviour (Riisager and Riisager, 2001; Smirnov et al., 2017) and/or alteration of magnetic minerals during 330 heating steps of the experiments (Kissel and Laj, 2004) as well as to instabilities of aged 331 thermoremanence (Shaar and Tauxe, 2015). Thermochemical alteration during the heatings 332 333 was very common in the experiments of this study and results that were affected by alteration, as seen from pTRM checks, were not accepted. Curved or zig-zagging Arai plots, 334 335 caused by MD behaviour and/or instability of TRM have been excluded as well if the curvature 336 or β criteria were not met. Hawkins et al. (2019) argued that the two-slope behaviour without a corresponding directional change in their study was attributed to strong thermal or 337 338 thermoviscous overprints of similar direction. This is also the case in our study, but here a 339 directional change is often observable at the bend between the two slopes in the Arai plots as well. Therefore, the slope of the high unblocking temperature ranges, that also carried the 340 341 ChRM in Halls et al. (2015), have been selected in the Arai plots to calculate palaeointensities. 342 The overall reliability of palaeointensity results was quantified by the application of QPI criteria (Biggin and Paterson, 2014) to each of the sites (see supplementary table 5). 343

The AGE criterion was met by all sites in the dykes dated in Halls et al. (2015). Sites GD33 from Augusta Lake dyke and site GD29 from Sand bay dyke were the only sites to meet the STAT criterion. The site GD23 from French River dyke has more than the minimum of 5 individual intensity estimates, but the estimate had a dispersion (standard deviation/mean)  $\leq$  25% (Paterson et al., 2010), which is often the case for ultra-low palaeointensity results (e.g. Hawkins et al., 2019; Shcherbakova et al., 2020). All other sides yielded less than the 5 successful palaeointensity results needed to meet the STAT criterion.

351 Mostly fresh looking titanomagnetite grains with high temperature oxy-exsolution structures and rapidly cooled dendritic magnetite grains without signs of low-temperature oxidations as 352 seen from SEM suggest that the magnetisation is a primary TRM. However, concerns can be 353 raised about the validity of the magnetisation itself due to the anomalous nature of the 354 355 palaeodirections (Halls et al., 2015). This is especially the case for sites associated with the shallow directional component where the possibility of a remagnetisation event cannot be 356 357 completely dismissed due to inconsistent results of a baked contact test (Robert et al., 2017). Halls et al. (2015) argued that these directions are also primary because they are almost 358 antipodal and can be followed continuously along the dyke. Palaeointensity estimates, 359 360 irrespective of the associated directions were similar and extremely low and the angular 361 dispersions of VGPs show almost identical behaviour of steep and shallow components. On the basis of the combined microscopic and palaeomagnetic evidence, we chose to award the 362 363 TRM criterion to all sites.

The use of pTRM checks and the application of the IZZI protocol with β and k' criteria as well
as the double heating checks of the Shaw experiments enabled the exclusion of all estimates
that could be significantly biased by thermochemical alteration or MD behaviour during the

367 experiments. Therefore, all sites passed both the ALT and MD criteria. Following the standard palaeointensity definitions 1.1 (Paterson et al., 2014), the angle between the laboratory field 368 and the last pTRM check of the Thellier experiments was calculated as  $\gamma = 3.2^{\circ} \pm 1.2^{\circ}$ , showing 369 that the results were not majorly influenced by anisotropy effects. Systematic bias of 370 estimates due to non-linear TRM behaviour was avoided by using different laboratory TRM 371 fields between 3  $\mu$ T and 30  $\mu$ T and ARM fields between 10  $\mu$ T and 81.2  $\mu$ T. The similarity of 372 373 results between the Shaw method with cooling times of ~1h and the microwave Thellier method with cooling times in the order of a fraction of a minute exclude a cooling rate bias. 374 375 Therefore, the ACN criterion was given to all sites.

Four different methods were used to determine palaeointensities. To pass the TECH criterion, 376 the palaeointensity estimate of a unit has to comprise results from at least two different 377 378 methods and this was the case for 4 sites (supplementary table 5). The criterion LITH was only awarded to site GD37 since it combines results from both the dyke and baked host rocks. It 379 380 was not met by any of the other sites as all intensity estimates came from the same lithology. Results of all intensity experiments with critical values are available in supplementary tables 381 1-3. In addition, all intensity measurement data is available on the MagIC database, which 382 awards the MAG criterion to all units. 383

Summing up the QPI criteria results in scores of 6-8, showing that the palaeointensity estimates in this study are of high quality. This allows the palaeointensities associated with the steep B and shallow C+D directions found in Halls et al. (2015) to be compared. Averaging site mean palaeointensities of all sites with shallow directions results in 3.7 ±2.3  $\mu$ T whereas the average palaeointensity of sites with steep directions gives 5.0 ± 0.5  $\mu$ T. Assuming a highlatitude Laurentia around ~590 Ma, the palaeointensities of the shallow component – if 390 recording an equatorial dipole field state – might be expected to be much lower than the 391 intensities of the steep component. The average intensities of sites with the shallow 392 component are indeed weaker, but with the higher values being within one standard deviation of the weaker values, the difference of < 2  $\mu$ T is negligible. The similarity of the 393 394 ultra-low palaeointensities of the two components suggests that the Grenville Dykes have recorded a highly unstable field as proposed in Halls et al. (2015). In contrast, the high and 395 396 almost identical values for angular dispersions (supplementary figure 3) of two distinct groups 397 of VGPs around the mean VGPs of the steep ( $S_B$ =18.5°) and the shallow component ( $S_B$ =18.9°) look consistent with the existence of an equatorial dipole. However, some caution is advised 398 399 when taking the dispersions at face value. Previous studies of VGP scatter required a minimum of N = 9 sites (Doubrovine et al., 2019; Veikkolainen and Pesonen, 2014). Due to 400 401 the exclusion of lower quality directions from sites with  $n \le 4$  or  $k \le 30$  in the calculations, this 402 requirement was only met by the shallow component (N = 11), but not by the steep (N = 7). 403 The dispersion values seem reasonable for the time period with comparable values for 404 Laurentia showing a wide range between S<sub>B</sub> = 13.5° (Skinner Cove volcanics, McCausland and 405 Hodych, 1998; Veikkolainen and Pesonen, 2014) and S = ~26° (Sept-Îles intrusion, Bono et al., 2019). If the two groups of VGPs were interpreted as one group, showing a transitional field, 406 then the resulting VGP dispersion would be  $S_B = 32.4^\circ$  (variable cutoff = 63.4°) at low latitude. 407 408 This would be an extremely high value but the current lack of constraints on the Ediacaran 409 field means that it is not implausible.

With the exception of sites GD14 and GD25, where single high (7.6 and 12.6  $\mu$ T) microwave results lead to high VDMs of 1.7 \*10<sup>22</sup>Am<sup>2</sup>, the site mean VDMs of 0.3 - 0.9 \*10<sup>22</sup>Am<sup>2</sup> are comparable to, but lower than, the results of the Sept Iles (Bono et al., 2019) and the Volyn Traps (Shcherbakova et al., 2020) with corresponding VDMs of 0.5 and 1.0 \*10<sup>22</sup>Am<sup>2</sup>. They

therefore define a new lower boundary for the field strength in the Ediacaran. The 414 palaeointensities, coming from dykes with ages that span ~15 Ma, suggest a sustained 415 geomagnetic field with these extremely low intensities. These results are an order of 416 magnitude weaker than the strength of the present-day field and the average 417 418 palaeointensities of site GD26 and sites from the French River dyke and the Augusta Lake dyke 419 with values between 1.4 and 2.1 µT are even as low as recent measurements of Mars' crustal field made by the InSight lander of  $\sim 2 \mu T$  (Johnson et al., 2020). Palaeointensities this low 420 421 have been reported for Earth before, but were generally not attributed to a sustained field. The PINT database (Biggin et al., 2015) contains 6 site mean estimates with  $H_{pal} \le 5 \mu T$  or VDM 422  $\leq$  0.5 \*10<sup>22</sup>Am<sup>2</sup>, N > 1 and reliable experiment types (excluding single-heating Shaw and total 423 TRM experiments) that can be roughly divided into two groups. The first group comprises 424 entries with mid-Miocene or younger ages that show the ultra-low intensity values in single 425 426 basalt flows from Iceland (Lawley, 1970) and the Canary Islands (Brown et al., 2009; Leonhardt 427 and Soffel, 2002) that are all connected to the short-term drop of dipole moments during polarity transitions interrupting a much stronger sustained field. The palaeointensities with 428 Mesoproterozoic to Archaean ages from the second group are either extremely low due to 429 fractions of CRM (Yoshihara and Hamano, 2004), would not satisfy any modern sets of 430 selection criteria (Ueno, 1995), or are only seen as spot reading of the field in a single dyke 431 432 (Smirnov and Tarduno, 2005) in a dyke swarm showing a weak, but overall stronger field (Halls et al., 2004). Non-Ediacaran weak sustained fields as in the Jurassic (e.g. ~2.8 \*10<sup>22</sup>Am<sup>2</sup>, Tauxe 433 et al., 2013) or the Devonian (e.g. ~1.1 \*10<sup>22</sup>Am<sup>2</sup>, Hawkins et al., 2019) are all stronger. 434 However, a field strength behaviour, similar to the one in the Ediacaran, can be observed in 435 436 the Upper Devonian around ~370 Ma, where site mean palaeointensities as low as 2.4  $\mu$ T (0.4 437 \*10<sup>22</sup>Am<sup>2</sup>) have been reported in the weak time averaged field as well (Hawkins et al., 2019).

438 The ultra-low Ediacaran intensities are consistent with the predicted weak field state of the geodynamo before the onset of inner core growth (Bono et al., 2019; Driscoll, 2016; Landeau 439 et al., 2017). Under this scenario, the dynamo was operating marginally, powered by thermal 440 441 convection due to heat loss at the core-mantle boundary alone and the field was diminished. Subsequently, the inner core nucleated providing additional convective power from the 442 release of light elements and latent heat of crystallisation. We cannot, however, rule out an 443 444 entirely different cause of a massively diminished dipole moment in the Ediacaran, perhaps 445 related to a reconfigured convective pattern in the core perhaps related to an unusual coremantle heat flow pattern. That the measured palaeointensities are similar in magnitude to 446 447 ground measurements of the crustal field of Mars, a planet which is suspected to have been without a core dynamo for 4.1 billion years, raises profound questions concerning 448 measurement limits and the history of the geodynamo. Is it even possible to measure 449 450 geomagnetic-field-derived palaeointensities that are any lower than the values presented 451 here, given the possibility of contamination by a (palaeo-) crustal field? Can we therefore be certain that the geodynamo did not die out altogether in the Ediacaran? Such questions are 452 453 unfortunately beyond the scope of the present study to answer but motivate urgent future investigations that may yield transformative insights into our planet's history. 454

455

### 456 **5 Conclusions**

We report new high quality palaeointensities from eleven sites of the western end of the Grenville dyke swarm. Success rates of palaeointensity experiments were substantially lowered by thermochemical alteration during laboratory heating as well as MD behaviour of the studied rocks. The resulting palaeointensity values range between  $1.4 - 7.6 \mu T (0.33 - 1.76)$ 

\*10<sup>22</sup>Am<sup>2</sup>). Most of these estimates agree well with results from other palaeointensity studies 461 from this time period (Bono et al., 2019; Shcherbakova et al., 2020) but several 462 palaeointensities are substantially lower with values being similar to values of Mars' crustal 463 field (Johnson et al., 2020). This opens up questions about what geodynamo regimes could 464 result in such weak fields or, to take it one step further, if such weak fields could even suggest 465 an inactive geodynamo. The behaviour of VGPs with almost identical angular dispersion 466 467 around two clusters argues for the presence of an equatorial dipole field. In contrast, the 468 consistency of presented ultra-low intensities along dykes for both directional components supports the idea of an unstable and/or transitional field in the Ediacaran (Bono et al., 2019; 469 470 Halls et al., 2015). The presented data are another argument for the Ediacaran field behaving strangely but at this point we are still struggling to define it. However, data for the early to 471 mid-Ediacaran is still scarce and a better coverage of palaeointensity and 472 473 magnetostratigraphic data would be immensely useful to better characterise the 474 geomagnetic field in the Ediacaran.

475

#### 476 Acknowledgements

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482

## 483 **Data availability:**

- 484 Data has been uploaded to a private repository in the MagIC database and will be accessible
- 485 after publication. Currently, data from this study are available from the corresponding author
- 486 on reasonable request.

487

488 Figures:

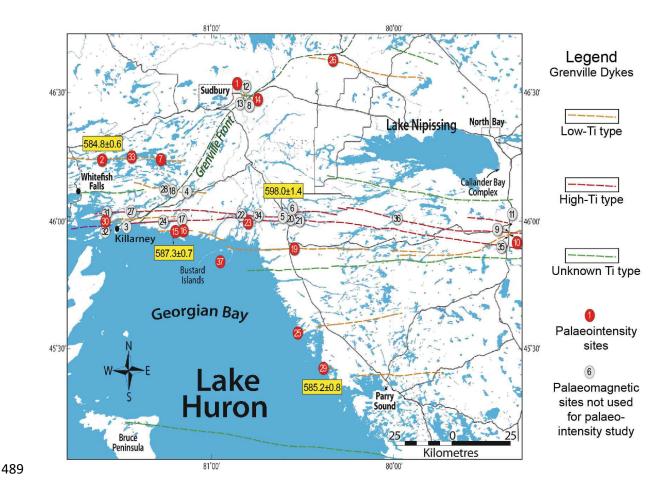


Figure 1: Map of southeast Ontario, showing the Grenville dykes and locations of sites studied
in Halls et al. (2015) and sites selected for this palaeointensity study; from Halls et al. (2015)

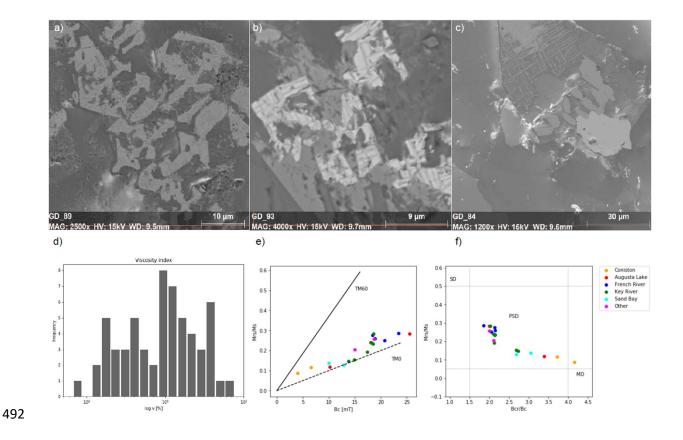
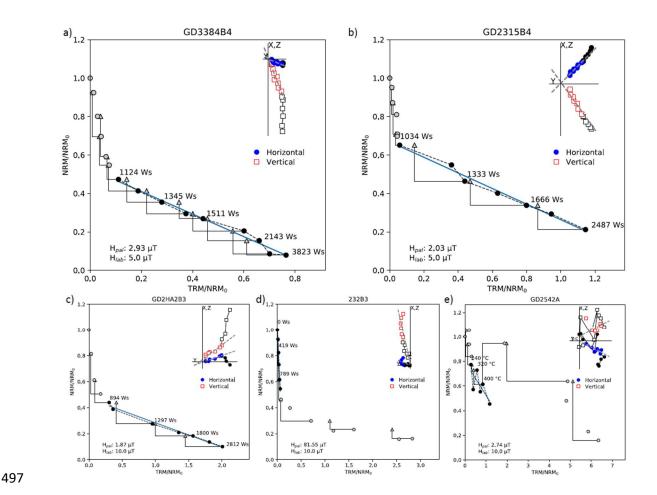


Figure 2: SEM backscatter images: (a) GD23, (b) GD29, examples of dendritic TM grains, (c)
GD02, coarse Tm grain, showing magnetite-ilmenite exsolution. (d) frequencies of viscosity
indexes of all sites, (e) Hysteresis parameters and (f) Day plot of representative samples of
each site.



498 Figure 3: Arai plots of Thellier experiment results. a)-d) show microwave experiments, e)
499 shows a result of a thermal Thellier experiment. For more details, readers are referred to
500 section 3 of the text.

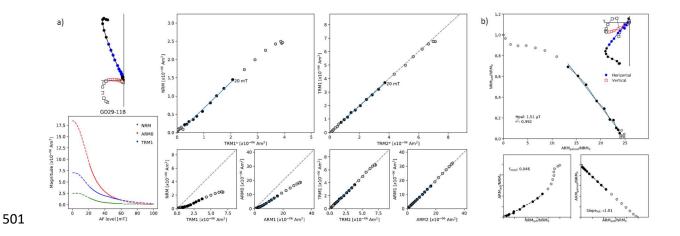


Figure 4: a) Representative example of a successful Shaw experiment: left column shows the
orthogonal plot of the demagnetisation of NRM (top) and absolute values of magnetisation

- 504 during AF demagnetisations of NRM, ARMO and TRM1 (bottom). middle and right columns
- show results from first and second laboratory heating, respectively; b) example of a successful

506 pseudo-Thellier experiment: pseudo-Arai plot with orthogonal projection, demag-demag plot

507 (NRM vs ARM) and ARM-ARM plot (acquisition vs. demagnetisation).

508

# 509 **Tables:**

		Age	Directions (Halls et al 2015)			Intensities [μT]				pThel [μT]	VDM [10 <sup>22</sup> Am <sup>2</sup> ]				
Dyke	Site	[Ma]	N/n	Dec	Inc	a95	k	n/n <sub>Pl</sub>	Methods	PI	Std	PI	VDM	Std	QPI
Augusta	GD02	584	9/8	120.7	23.1	6.7	69	12/2	MW+S	2.1	0.04	-	0.52	0.01	7
Lake	GD33	±0.6	8/6	135.6	31.2	8.8	60	17/5	MW	2.6	0.62	1.9	0.60	0.14	7
Coniston	GD01		5/3	115.5	17.1	32.4	16	4/0	-	-	-	-	-	-	-
	GD14		13/3	95.5	23.8	26.3	13	6/1	MW	7.6	-	-	1.76	-	6
French		598													_
River	GD23	±1.4	6/3	132.7	32.6	8.7	202	24/13	MW+S	1.8	0.67	-	0.42	0.15	7
Кеу	GD10		10/9	137.7	75.8	6.7	61	14/0	-	-	-	-	-	-	-
River	GD15		5/5	140.9	66.8	6.7	130	12/1	MW	5.6	-	-	0.88	-	6
	GD16		6/6	163.8	76.7	14.4	22	10/1	MW	4.3	-	1.3	0.67	-	6
	GD19	587	9/5	100.6	56.2	11.9	42	11/0	-	-	-	-	-	-	-
	GD37	±0.7	8/7	143.1	71.7	5.1	141	6/3	S	5.3	1.57	3.1	0.78	0.23	7
		585													
Sand Bay	GD29	±0.8	9/9	297	40.1	5.2	97	41/7	MW+S	3.7	0.69	4.1	0.80	0.15	8
Other	GD07		7/6	134.8	1.2	11.2	37	1/0	-	-	-	-	-	-	-
Grenville	GD25		9/9	136.5	9.6	8.8	36	14/2	MW+S	6.8	5.82	-	1.74	1.41	7
Dykes	GD26		9/6	130.1	23.8	10.1	45	7/2	S	1.4	0.40	1.5	0.33	0.10	6
	GD30		10/6	139	67.3	8.4	64	3/0	-	-	-	-	-	-	-

510

511	Table 1: Summary of palaeointensity results of the Grenville dykes: Dyke/Site: dyke/site name, Ages
512	and directional information from Halls et al. (2015), $N/n_{int}$ : total number of specimens/ number of
513	successful results, Methods: type of method that contributed successful results (MW: microwave
514	Thellier, TH: thermal Thellier, S: Shaw), PI: palaeointensity results in $\mu T$ , Std: standard deviation of
515	palaeointensity results (*: standard error of single specimen result instead of multiple specimen

- result), pThel: palaeointensity results in  $\mu$ T from pseudo-Thellier experiments. Values are shown as comparison to the results from the heating methods but are not used in the calculation of the mean palaeointensities. VDM: virtual dipole moment in  $10^{22}$ Am<sup>2</sup>, Std: standard deviation of VDM results, QPI values of sites.
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