

1 **Joint inversion of co-seismic and early post-seismic slip to**
2 **optimize the information content in geodetic data: Application**
3 **to the 2009 M_w 6.3 L'Aquila earthquake, Central Italy**

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8 **Key Points:**

- 9 • Simultaneous inversion of geodetic data with different spatio-temporal resolution
10 maximizes the information content
11 • Incorporating early afterslip deformation when solving for the co-seismic processes
12 (as with InSAR data) overestimate inferred co-seismic models
13 • Estimations of the post-seismic processes neglecting early afterslip deformation
14 may largely underestimate the total afterslip amplitude

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Abstract

When analyzing the rupture of a large earthquake, geodetic data are often critical. These data are generally characterized by either a good temporal or a good spatial resolution, but rarely both. As a consequence, many studies analyze the co-seismic rupture with data also including one or more days of early post-seismic deformation. Here, we propose to invert simultaneously for the co- and post-seismic slip with the condition that the sum of the two models remains compatible with data covering the two slip episodes. We validate the benefits of our approach with a toy model and an application to the 2009 M_w 6.3 L'Aquila earthquake, using a Bayesian approach and accounting for epistemic uncertainties. For the L'Aquila earthquake, we find that if early post-seismic deformation is acknowledged as co-seismic signal, co-seismic slip models may overestimate the peak amplitude while long-term post-seismic models may largely underestimate the total post-seismic slip amplitude. This example illustrates how the proposed approach could improve our comprehension of the seismic cycle, of the fault frictional properties, and how the co-seismic rupture, afterslip and aftershocks relate to one another.

1 Introduction

The occurrence of earthquakes and seismic sequences is mainly controlled by the spatial and temporal evolution of crustal stresses. The co-seismic stress change and the redistribution of stress following an earthquake thus both play an important role in the seismic cycle and the mechanical behavior of faults, including the generation of new seismic sequences. To understand both co-seismic and post-seismic processes, and their relationship, is thus a crucial step to propose realistic earthquakes scenario and reliable hazards estimates.

While earthquakes can last for a few seconds to minutes, their post-seismic stress relaxation can last months to years. Post-seismic relaxation is generally modeled by several interacting mechanisms, such as localized shear on the fault (a.k.a. afterslip) [e.g. *Marone et al.*, 1991; *Freed*, 2007; *Johnson et al.*, 2012], visco-elastic deformation in the lower crust or mantle [e.g. *Nur and Mavko*, 1974; *Pollitz et al.*, 1998; *Freed and Burgmann*, 2004] or poroelastic rebound [e.g. *Peltzer et al.*, 1998; *Jonsson et al.*, 2003]. The interactions between co-seismic stress changes, aftershocks and post-seismic deformation are still poorly understood [e.g. *Perfettini and Avouac*, 2007]. Slip on the fault may be governed by two brittle deformation modes following rate and state friction laws [*Rice and L. Ruina*, 1983]: seismic rupture may occur in velocity weakening area, whereas afterslip may develop in the velocity strengthening zone [e.g. *Marone et al.*, 1991]. In contrast, *Helmstetter and Shaw* [2009] also show that afterslip processes may be primarily driven by stress heterogeneities, independently of the rate and state friction behavior. Aftershocks may be triggered by co-seismic stress changes, without direct relation with post-seismic deformation [*Dieterich*, 1994]. Or, aftershocks may also be primarily triggered by the post-seismic reloading due to afterslip [e.g. *Perfettini and Avouac*, 2004; *Hsu et al.*, 2006; *Peng and Zhao*, 2009; *Ross et al.*, 2017]. The variability of these theories emphasizes the need to refine our comprehension and description of the co-seismic and post-seismic phases and their transition.

Our understanding of the co-seismic processes mainly derives from modeling of seismic, geodetic and tsunami data, and our understanding of post-seismic behaviors is mainly improved with the modeling of geodetic observations [e.g. *Burgmann et al.*, 1997; *Wang et al.*, 2012; *Perfettini and Avouac*, 2014; *Gualandi et al.*, 2017] or simulation [e.g. *Smith and Sandwell*, 2004; *Barbot and Fialko*, 2010; *Cubas et al.*, 2015]. The observations thus remain a cornerstone to identify and characterize the co- and post-seismic processes. GNSS time series are commonly used and can provide a good temporal resolution. But the spatial resolution of such observation is usually limited. In contrast, synthetic aperture radar interferometry (InSAR) can provide extensive spatial coverage but with a limited

66 temporal resolution. Indeed, while earthquakes last for a few seconds, very often satellites
 67 have a revisit time of more than a few days. If earthquakes do not nucleate just before
 68 the visit of a satellite, which is generally the case, the measured deformation is the co-
 69 seismic signal plus a fraction of the post-seismic deformation. As a consequence, most
 70 earthquakes models based on geodetic observations are biased by unwanted deformation
 71 signal. In practice, used interferograms or campaign GNSS offsets generally cover time
 72 periods extending at least a few days before and after the mainshock. Pre-earthquake sig-
 73 nals, when evidenced, are usually related to small slip episodes at depth near the hypocen-
 74 ter. The associated surface deformation signals are usually hard to detect and neglected
 75 in co-seismic studies. The post-seismic deformation happening on the first few days after
 76 the mainshock is usually detectable in the geodetic data but incorporated in source es-
 77 timation problems as if it was part of the co-seismic signal [e.g. *Elliott et al.*, 2013; *Lin*
 78 *et al.*, 2013; *Cheloni et al.*, 2014; *Bletery et al.*, 2016; *He et al.*, 2017; *Salman et al.*, 2017;
 79 *Barnhart et al.*, 2018], with the justification that it is comparatively small. Similarly, post-
 80 seismic models generally do not account for observations related to the early post-seismic
 81 deformation because they are often contaminated by co-seismic signal [e.g. *D'Agostino*
 82 *et al.*, 2012; *Cheloni et al.*, 2014]. What we name here the early post-seismic phase cor-
 83 responds to the overlooked part of the post-seismic deformation, and can last for a few
 84 hours after the mainshock in the best case, or a few days in most studies. Yet, the largest
 85 post-seismic deformation rate is expected during the first few days after the mainshock,
 86 considering that its main trend is to decrease exponentially with time after an earthquake.

87 The early post-seismic processes remain largely unexplored, because of the limited
 88 temporal and spatial resolution of geodetic data. Neglecting the early post-seismic signal
 89 may also affect our understanding of both co-seismic and post-seismic processes. And this
 90 effect is probably persisting if seismic data (i.e. purely co-seismic) are added to the in-
 91 verse problem, since geodetic data tend to have a stronger control on the inferred distribu-
 92 tion of slip, at least in the first 10 km below the Earth surface [e.g. *Delouis et al.*, 2002].

93 The recent advent of high-frequency GNSS has allowed to record the strictly co-
 94 seismic signal (10 to 30 seconds after the earthquake time occurrence) without any con-
 95 tamination by early post-seismic deformation. Well instrumented earthquakes are thus now
 96 characterized by at least two geodetic datasets, one being strictly co-seismic and the other
 97 which also includes some days of early afterslip. In this study, we use an original inver-
 98 sion methodology to jointly infer co-seismic and early post-seismic slip models, taking
 99 advantage of the complementary spatial and temporal resolutions of different geodetic ob-
 100 servations. We first validate the approach through a toy model, and then analyze and il-
 101 lustrate the benefits of our methodology with a real event. We consider the 2009 M_w 6.3
 102 L'Aquila earthquake, Central Italy, which has been intensively studied but whose very
 103 early post-seismic phase has not been imaged. The choice of the L'Aquila event is also
 104 motivated by the large density of near field observations and the overall quality of the in-
 105 strumentation. Additionally, this event ruptured a relatively well known and simple fault
 106 geometry, in an area where crustal properties have been investigated in detail: this will
 107 ensure the forward physics and its uncertainties can be estimated. In this work we inves-
 108 tigate the impact of accounting for early afterslip on co-seismic models. We explore the
 109 impact of uncertainties in the slip imagery with a probabilistic approach and account for
 110 uncertainties in the physics of our problem. The results will allow us to investigate the re-
 111 lationship between co-seismic rupture, early afterslip processes, longer term afterslip and
 112 the distribution of aftershocks.

113 **2 Inversion Framework**

114 **2.1 Dual time inversion of co-seismic and early post-seismic data**

115 For a number of earthquakes, we have the opportunity to use two static datasets:
 116 one which is strictly co-seismic ("co") and the other which contains co-seismic and early

117 post-seismic signal ("co+post"). To infer the strictly co-seismic and early post-seismic
 118 slip distributions, one approach could be to invert separately for the two datasets, and as-
 119 sume that the strictly post-seismic ("post") solution is the difference between the "co" and
 120 "co+post" models. However, in this case, the model "co" would be constrained by fewer
 121 observations (only few GNSS offsets), most of the co-seismic information being in the
 122 "co+post" dataset (dense map of InSAR offsets). An alternative approach is to assume
 123 that the "co+post" slip model is the sum of the "co" and "post" slip distributions. We then
 124 have:

$$\begin{cases} \mathbf{d}_{\text{co}} = \mathbf{G}_{\text{co}}^{\text{co}} \cdot \mathbf{m}_{\text{co}} \\ \mathbf{d}_{\text{co+post}} = \mathbf{G}_{\text{co+post}}^{\text{co}} \cdot \mathbf{m}_{\text{co}} + \mathbf{G}_{\text{co+post}}^{\text{post}} \cdot \mathbf{m}_{\text{post}} \end{cases} \quad (1)$$

125 where matrices of the Green's functions $\mathbf{G}_{\text{data}}^{\text{model}}$ have been calculated for the corresponding
 126 dataset and model. For instance, $\mathbf{G}_{\text{co+post}}^{\text{co}}$ is the matrix of the Green's functions calculated
 127 from the model "co" for the data "co+post". The Eq. 1 can also be represented in the fol-
 128 lowing matrix form:

$$\begin{pmatrix} \mathbf{d}_{\text{co}} \\ \mathbf{d}_{\text{co+post}} \end{pmatrix} = \begin{pmatrix} \mathbf{G}_{\text{co}}^{\text{co}} & 0 \\ \mathbf{G}_{\text{co+post}}^{\text{co}} & \mathbf{G}_{\text{co+post}}^{\text{post}} \end{pmatrix} \cdot \begin{pmatrix} \mathbf{m}_{\text{co}} \\ \mathbf{m}_{\text{post}} \end{pmatrix}. \quad (2)$$

129 The redesigned Green's functions matrix is now composed of 3 sub-matrices. As we fo-
 130 cus on the early post-seismic phase, we can make the assumption that both $\mathbf{G}_{\text{co+post}}^{\text{co}}$ and
 131 $\mathbf{G}_{\text{co+post}}^{\text{post}}$ matrices are identical because we suppose both co-seismic and early post-seismic
 132 deformations are elastic. We can thus write

$$\begin{pmatrix} \mathbf{d}_{\text{co}} \\ \mathbf{d}_{\text{co+post}} \end{pmatrix} = \begin{pmatrix} \mathbf{G}_{\text{co}} & 0 \\ \mathbf{G}_{\text{co+post}} & \mathbf{G}_{\text{co+post}} \end{pmatrix} \cdot \begin{pmatrix} \mathbf{m}_{\text{co}} \\ \mathbf{m}_{\text{post}} \end{pmatrix}. \quad (3)$$

133 If strictly post-seismic observations are available, we could also incorporate these data into
 134 our equation to help constrain the "post" model:

$$\begin{pmatrix} \mathbf{d}_{\text{co}} \\ \mathbf{d}_{\text{co+post}} \\ \mathbf{d}_{\text{post}} \end{pmatrix} = \begin{pmatrix} \mathbf{G}_{\text{co}}^{\text{co}} & 0 \\ \mathbf{G}_{\text{co+post}}^{\text{co}} & \mathbf{G}_{\text{co+post}}^{\text{post}} \\ 0 & \mathbf{G}_{\text{post}}^{\text{post}} \end{pmatrix} \cdot \begin{pmatrix} \mathbf{m}_{\text{co}} \\ \mathbf{m}_{\text{post}} \end{pmatrix}, \quad (4)$$

135 with $\mathbf{G}_{\text{post}}^{\text{post}}$ reflecting the response of the Earth for the strictly post-seismic data. The "post"
 136 dataset then corresponds to the same post-seismic time window as what is covered by the
 137 "co+post" dataset.

138 The off-diagonal terms of the redesigned Green's function matrix allow us to make
 139 use of the "co+post" dataset to constrain both "co" and "post" models. In the following,
 140 we refer to this approach as Combined Time Windows (CTW) approach. The CTW ap-
 141 proach can be generalized to cover various intervals of post-seismic deformation. Indeed,
 142 while for many earthquakes strictly co-seismic data are now available, non-strictly co-
 143 seismic datasets usually cover variable time intervals. If, for instance, two intervals of
 144 post-seismic deformation contaminate the co-seismic signal, with only one of these in-
 145 tervals observed independently, our equation 3 can be adapted as

$$\begin{pmatrix} \mathbf{d}_{\text{co}} \\ \mathbf{d}_{\text{co+post1}} \\ \mathbf{d}_{\text{co+post2}} \\ \mathbf{d}_{\text{post2}} \end{pmatrix} = \begin{pmatrix} \mathbf{G}_{\text{co}}^{\text{co}} & 0 & 0 \\ \mathbf{G}_{\text{co+post1}}^{\text{co}} & \mathbf{G}_{\text{co+post1}}^{\text{post1}} & 0 \\ \mathbf{G}_{\text{co+post2}}^{\text{co}} & \mathbf{G}_{\text{co+post2}}^{\text{post1}} & \mathbf{G}_{\text{co+post2}}^{\text{post2}} \\ 0 & 0 & \mathbf{G}_{\text{post2}}^{\text{post2}} \end{pmatrix} \cdot \begin{pmatrix} \mathbf{m}_{\text{co}} \\ \mathbf{m}_{\text{post1}} \\ \mathbf{m}_{\text{post2}} \end{pmatrix}, \quad (5)$$

146 with $\mathbf{d}_{\text{post2}}$ reflecting the surface displacement for the time interval between times 1 and
 147 2, and $\mathbf{G}_{\text{post2}}^{\text{post2}}$ and $\mathbf{m}_{\text{post2}}$ associated Green's functions and slip model. Indeed, this ap-
 148 proach could be used to investigate as many time windows of post-seismic deformation
 149 as needed.

150 To refine co-seismic models and investigate early post-seismic deformation of the
 151 L’Aquila earthquake, we follow here the approach described by Eqs 1 and 3. We do not
 152 incorporate any information on the strictly post-seismic phase to investigate the very sim-
 153 ple case where only co-seismic data (contaminated or not by early post-seismic deforma-
 154 tion) are available.

155 2.2 Accounting for Epistemic Uncertainties

156 When imaging a slip distribution on a fault, the physics of the forward model is
 157 usually assumed of minimum complexity to simplify the computation and also often be-
 158 cause we don’t know well the Earth interior. For instance, Earth interior is frequently ap-
 159 proximated as an elastic and homogeneous environment and the causative fault geometry
 160 is usually reduced to a flat rectangular plane. The uncertainties related to our approxima-
 161 tions of the physics of the Earth affect the inferred source models [Ragon *et al.*, 2018].
 162 As the early post-seismic slip is of limited amplitude, it may be particularly impacted by
 163 uncertainties of the forward model. We thus account for epistemic uncertainties following
 164 the approach developed by Duputel *et al.* [2014] for the Earth elastic properties and Ragon
 165 *et al.* [2018] for the fault geometry. The epistemic uncertainties are calculated from the
 166 sensitivity of the Green’s Functions and are included in a covariance matrix \mathbf{C}_p .

167 2.3 Bayesian approach

168 Our inverse problem solves for both co-seismic and early post-seismic slip param-
 169 eters, the later being of limited amplitude. While the co-seismic parameters will be rea-
 170 sonably well constrained, multiple early post-seismic models will probably be realistic.
 171 To get a robust image of the early post-seismic phase, we thus solve our problem with
 172 a Bayesian sampling approach which relies on the AITar package, which is a rewrite of
 173 the code CATMIP [Minson *et al.*, 2013]. AITar combines the Metropolis algorithm with a
 174 tempering process to realize an iterative sampling of the solution space of the source mod-
 175 els. A large number of samples are tested in parallel at each transitional step. Addition-
 176 ally, a resampling is performed at the end of each step to replace less probable models.
 177 The probability of each sample to be selected depends on its ability to fit the observations
 178 \mathbf{d}_{obs} within the uncertainties $\mathbf{C}_\chi = \mathbf{C}_d + \mathbf{C}_p$, with \mathbf{C}_d the observational errors and \mathbf{C}_p the
 179 epistemic uncertainties.

180 The ability of each model parameter to solve the source problem is evaluated through
 181 repeated updates of the probability density functions (PDFs)

$$f(\mathbf{m}, \beta_i) \propto p(\mathbf{m}) \cdot \exp[-\beta_i \cdot \chi(\mathbf{m})], \quad (6)$$

182 where \mathbf{m} is the current sample, $p(\mathbf{m})$ is the prior information on this sample, i corre-
 183 sponds to each iteration and β evolves dynamically from 0 to 1 to ensure an exhaustive
 184 exploration of the solution space [Minson *et al.*, 2013]. $\chi(\mathbf{m})$ is the misfit function:

$$\chi(\mathbf{m}) = \frac{1}{2} [\mathbf{d}_{\text{obs}} - \mathbf{G} \cdot \mathbf{m}]^T \cdot \mathbf{C}_\chi^{-1} \cdot [\mathbf{d}_{\text{obs}} - \mathbf{G} \cdot \mathbf{m}]. \quad (7)$$

185 The use of AITar with the CTW approach allows us to specify prior information on each
 186 model, and thus to ensure the non-negativity of both co-seismic and post-seismic slip
 187 models (or of any time window model).

188 3 Application to a simplified 2D model

189 To ensure that our methodology allows to reliably infer the slip distribution of dif-
 190 ferent time windows, we first analyze a synthetic 2D case where the slip is imaged either

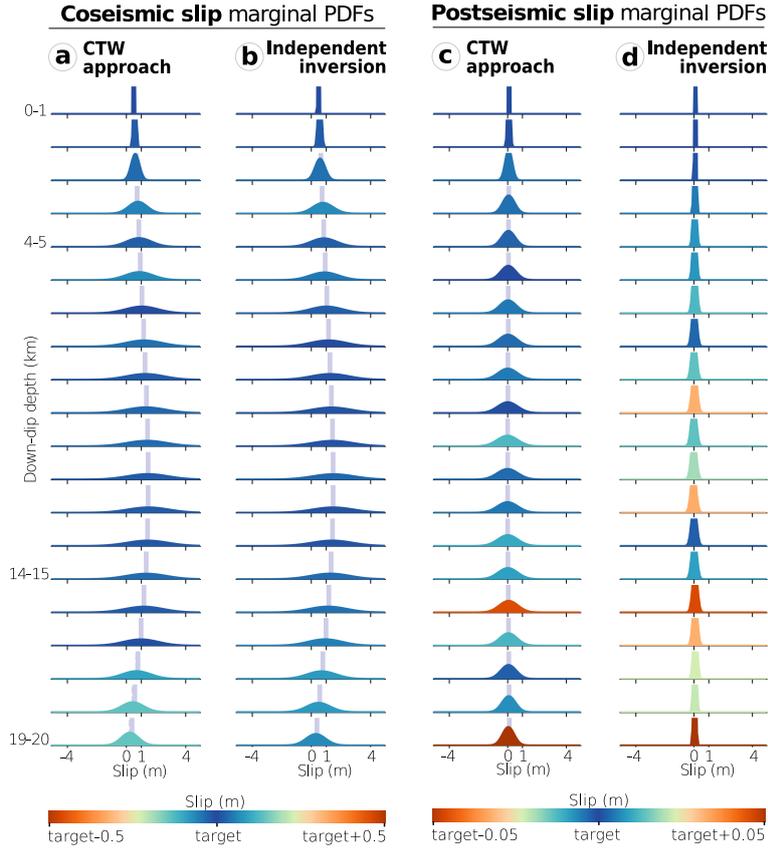


Figure 1. Co-seismic and post-seismic slip inferred for the simplified case of a fault that extends infinitely along strike. The co- and post-seismic slip models inferred from the CTW approach are shown in (a) and (c), and can be compared to the slip inferred from the inversion of co-seismic data only (b) and the post-seismic slip distribution (d) resulting from the difference between slip inferred from the inversion of co+post data and slip of (b). The fault is discretized along dip in 20 subfaults, for which are represented the target parameters as gray vertical lines. For each subfault, the posterior PDFs of co-seismic (a and b) and post-seismic (c and d) slip is colored according to the offset between the target parameter and the posterior mean, with a colorscale saturated at 50 cm for the co-seismic slip and at 5 cm for the post-seismic slip. The target slip is well inferred if the PDF of a particular parameter is colored in dark blue, while it is not if the PDF is colored in red.

191 independently or with the CTW approach. For this simple case, we assume two time win-
 192 dows named for the purpose of simplicity co-seismic and post-seismic.

193 3.1 Forward Model

194 We assume a fault extending infinitely along strike and which is 20 km wide along
 195 dip. The fault is discretized along dip into sub-faults of 1 km width and is dipping 55° .
 196 We assume the co-seismic slip on this fault to be purely dip-slip and to vary gradually
 197 with depth between 0 m and 1.5 m, with maximum slip between 9 and 14 km depth. We
 198 also assume that there is post-seismic slip on the same fault, with a similar location and
 199 direction and an amplitude equal to a tenth of the co-seismic slip amplitude. We compute
 200 the corresponding "co" and "co+post" synthetic observations using the expressions of sur-
 201 face displacement in an homogeneous elastic half-space [Segall, 2010]. These synthetic
 202 observations are computed for 100 data points at the surface, spaced every kilometer. A

203 correlated Gaussian noise of 5 mm is added to the synthetic data to simulate measurement
 204 errors. Note that, for this toy model, the number of "co" data is the same as the number
 205 of "co+post" observations.

206 Using these 100 synthetic observation points, we then estimate the depth distribu-
 207 tion of slip still assuming a homogeneous elastic half-space. We use a uniform prior dis-
 208 tribution $p(\mathbf{m}) = \mathcal{U}(-0.5 \text{ m}, 5 \text{ m})$ for the dip slip component (uniform implies that all
 209 values are considered equally likely with no a priori knowledge), a zero-mean Gaussian
 210 prior $p(\mathbf{m}) = \mathcal{N}(-0.1 \text{ m}, 0.1 \text{ m})$ on the strike-slip component and include 5 mm of obser-
 211 vational uncertainty in \mathbf{C}_d . We do not account for epistemic uncertainties as our forward
 212 model is the replicate of the one used to generate the data. We first solve for the "co" and
 213 "post" slip following the CTW approach (Figures 1a and 1c). Then, we run independent
 214 inversions, one to solve for the "co" slip and the other one to infer the "co+post" slip, the
 215 post-seismic solution being the difference between "co" and "co+post" models (Figures 1b
 216 and 1d).

217 3.2 Results

218 Both independent and CTW inversion approaches allow to correctly infer the "co"
 219 slip, as the median of the PDFs is very close to the target model (Figures 1a and 1b). As
 220 expected from the inversion of surface data, the resolution is very good on shallow parts
 221 of the fault but quickly decreases with depth. The posterior uncertainty on deepest param-
 222 eters is slightly decreased because the lower tip of the fault acts as an additional constrain.
 223 In contrast, the inversion methodology has a larger impact on the inferred "post" slip dis-
 224 tributions. When jointly inverting "co" and "co+post" observations, the correct "post" slip
 225 is well estimated at almost all subfaults (Figure 1c). When solving the two slip stages sep-
 226 arately, the mean of the models is not as good at estimating the target model (Figure 1d).
 227 The reduced posterior uncertainty of the "post" model for the independent inversion is an
 228 artifact resulting from the subtraction of two gaussian-shaped curves.

229 In summary, the two inversion approaches allow to reliably infer the "co" slip dis-
 230 tributions, probably because its signal is dominating in the observations. But the CTW
 231 approach provides a more robust estimation of the "post" slip distribution. In this 2D case,
 232 co-seismic and co+post signals have been observed by the same number of stations. How-
 233 ever, for most earthquakes, the number of "co" data points available (usually GNSS) will
 234 be very limited compared to the quantity of "co+post" observations (usually InSAR). We
 235 thus expect that if performing independent inversions for a real event, the inferred "co"
 236 slip distribution will be less reliable than in the case of a CTW inversion, where the whole
 237 "co+post" dataset is used to guide the choice of co-seismic parameters. We now compare
 238 these two approaches on a real earthquake.

239 4 Application to the 2009 M_w 6.3 L'Aquila earthquake, Central Italy

240 The L'Aquila earthquake nucleated within the Apennines orogenic system (Fig-
 241 ure 2), where the current seismic activity results from the ongoing extensional tectonics
 242 of the area. The mainshock nucleated on the Paganica fault [Figure 2, *Atzori et al.*, 2009;
 243 *Falucci et al.*, 2009; *Chiaraluca et al.*, 2011; *Vittori et al.*, 2011; *Lavecchia et al.*, 2012;
 244 *Cheloni et al.*, 2014], southwest of the city of L'Aquila, and has been followed by at least
 245 4 aftershocks of $M_w > 5$ [*Scognamiglio et al.*, 2009; *Chiarabba et al.*, 2009; *Pondrelli*
 246 *et al.*, 2010]. Although the L'Aquila earthquake has been intensively studied, most co-
 247 and post-seismic models have considered the first days of post-seismic deformation as if
 248 they were part of the co-seismic phase [e.g. *Anzidei et al.*, 2009; *Atzori et al.*, 2009; *Che-*
 249 *loni et al.*, 2010; *Trasatti et al.*, 2011; *Cirella et al.*, 2012; *D'Agostino et al.*, 2012; *Cheloni*
 250 *et al.*, 2014; *Balestra and Delouis*, 2015; *Volpe et al.*, 2015]. To avoid the contamination of
 251 co-seismic signal by early afterslip, *Yano et al.* [2014] proposed to explore independently
 252 the early post-seismic deformation, yet with datasets covering different time intervals (1

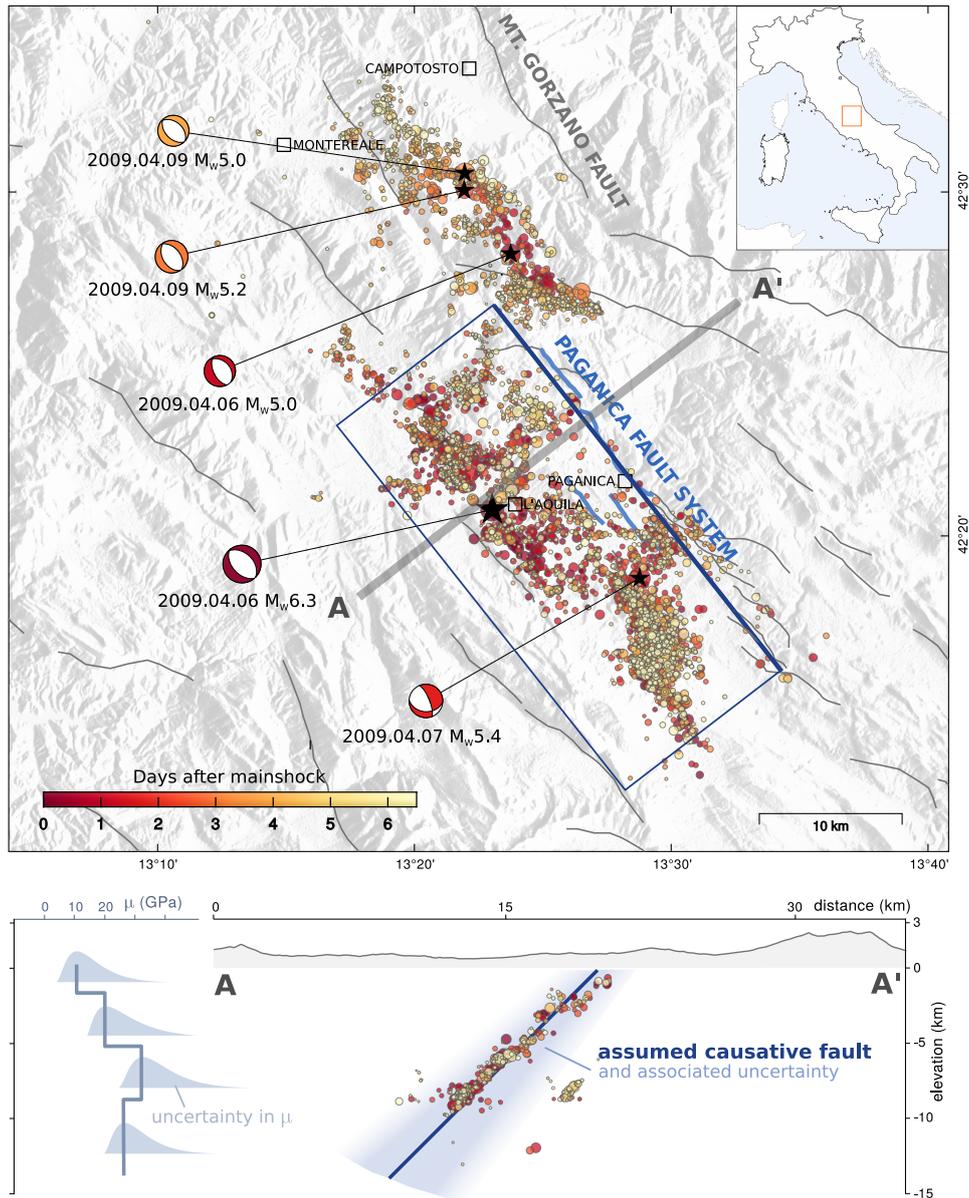


Figure 2. Seismotectonic framework of the area involved in the 2009 seismic sequence (top) and assumed forward model and associated uncertainties (bottom). In the map, coloured circles are the aftershocks from 2009 April 6 at 01:32 UTC to 2009 April 12, from the catalog of *Valoroso et al.* [2013]. The aftershocks are coloured from their occurrence time after the mainshock. Beach balls are the focal mechanisms of the mainshock and four main aftershocks, with their respective epicenters located by black stars. Solid gray lines are the major seismogenic faults of the area [*Boncio et al.*, 2004a; *Lavecchia et al.*, 2012]. The observed co-seismic surface rupture is indicated with continuous blue lines [*Boncio et al.*, 2010]. Our assumed fault geometry is shown with a dark blue rectangle. In the elevation profile (bottom), uncertainty in the fault geometry is illustrated in blue. The assumed elastic modulus μ and associated uncertainties are also illustrated for the 12 first kilometers below the Earth surface.

254 ment is recorded up to several months after the mainshock [e.g. *D'Agostino et al.*, 2012;
 255 *Gualandi et al.*, 2014; *Cheloni et al.*, 2014; *Albano et al.*, 2015], yet most studies of long-
 256 term post-seismic signal did not analyze the first few days of post-seismic deformation.

257 4.1 Data, Forward Model and Prior Information

258 From a geodetic perspective, this event has been particularly well documented. We
 259 can distinguish two main static datasets: one which is strictly co-seismic ("co", using con-
 260 tinuous GNSS data), and the other which also includes some days of post-seismic slip
 261 ("co+post", using cGNSS and InSAR). Two SAR images were acquired 6 days after the
 262 mainshock rupture, making the L'Aquila earthquake a perfect case study for our proposed
 263 approach. The "co" dataset corresponds to surface displacements measured between the
 264 earthquake time occurrence (t_0) and 25-30 s after t_0 , and includes the static offsets of 41
 265 (including high-rates) GPS stations surrounding the earthquake area processed by *Aval-
 266 lone et al.* [2011]. The "co+post" dataset covers the co-seismic phase plus 6 days of post-
 267 seismic slip, documented by 40 static GPS offsets and 2 InSAR frames: an ascending
 268 COSMO-SkyMed frame and a descending Envisat frame (Tab. S1). The observations and
 269 their processing are detailed in Supplementary Material [Section S1, *Rosen et al.*, 2004;
 270 *Lohman and Simons*, 2005; *Jolivet et al.*, 2012].

271 The Paganica fault is generally thought to be responsible for the co-seismic rupture
 272 of the L'Aquila earthquake, and also for most of its post-seismic deformation [*D'Agostino
 273 et al.*, 2012; *Cheloni et al.*, 2014; *Yano et al.*, 2014] along with the northernmost Cam-
 274 potosto fault [Figure 2, *Gualandi et al.*, 2014]. Although the distribution of relocalized
 275 aftershocks and surface rupture suggest that the Paganica fault system is possibly seg-
 276 mented [*Boncio et al.*, 2010; *Lavecchia et al.*, 2012] and/or curved at depth [*Chiaralu-
 277 ce et al.*, 2011; *Lavecchia et al.*, 2012; *Valoroso et al.*, 2013], its geometry remains poorly
 278 constrained below the surface. The variability of published morphologies for the causative
 279 fault [*Lavecchia et al.*, 2012] suggests that even with a large amount of observations and a
 280 great seismotectonic knowledge of the area, it is not possible to determine a unique fault
 281 geometry. This is why we approximate the Paganica fault geometry as a planar surface.
 282 We determine strike and position from the trace of the co-seismic surface rupture [*EMER-
 283 GEO Working Group*, 2010; *Boncio et al.*, 2010] and formerly identified seismogenic faults
 284 [e.g. *Boncio et al.*, 2004b]. We select the dip and width based on aftershocks relocations
 285 and focal mechanisms [e.g. *Chiaralu-
 286 ce et al.*, 2011; *Chiaralu-
 287 ce*, 2012; *Valoroso et al.*,
 288 2013]. Hence, our preferred geometry extends over 25 km south of coordinates (13.386°
 289 E, 42.445° N) with a strike of N142°. We set fault dip at 54° and width at 18 km, such
 290 that the fault is reaching the ground surface. This geometry is in agreement with already
 291 proposed causative structures [e.g. *Lavecchia et al.*, 2012]. The fault is divided into 154
 292 subfaults of 1.8 km length and 1.6 km width. As our fault geometry does not reflect the
 293 reality and is poorly constrained, we account for its uncertainties [*Ragon et al.*, 2018] and
 assume a standard deviation on the fault dip of 5° and on the fault position of 1.5 km,

294 We perform the static slip inversion assuming a 1-D layered elastic structure de-
 295 rived from the CIA velocity model [*Herrmann et al.*, 2011], and calculate Green's func-
 296 tions with the EDKS software [*Zhu and Rivera*, 2002]. We precompute Green's functions
 297 at depths intervals of 500 m down to 15 km depth and every 5 km below. Laterally, the
 298 Green's functions are computed every kilometer to reach the maximum epicentral distance
 299 of 100 km. Then, we interpolate and sum pre-computed Green's functions given our fault
 300 geometry and data locations. The strong variability in published elastic models for the
 301 central Italy [*Herrmann et al.*, 2011] can have a strong influence on co-seismic slip esti-
 302 mates [e.g. *Trasatti et al.*, 2011; *Volpe et al.*, 2012; *Gallovic et al.*, 2015]. We thus account
 303 for the uncertainties in our Earth model [*Duputel et al.*, 2014] assuming a standard devia-
 304 tion on shear modulus of 4 % at depths greater than 15 km and 13 % above. These values

305 have been chosen a priori considering the variability between layered models and the hori-
 306 zontal variability of 3D crustal models for several depth intervals [Magnoni *et al.*, 2014].

307 We perform our static slip inversion as previously detailed in Section 2. We specify
 308 prior distributions for each model parameter: a zero-mean Gaussian prior $p(\mathbf{m}) = \mathcal{N}(-10$
 309 cm, 10 cm) on the strike-slip component (we assume that, on average, the slip direction is
 310 along dip) and we consider each possible value of dip-slip displacement equally likely if
 311 it does not exceed 20 cm of reverse slip and 5 m of normal slip: $p(\mathbf{m}) = \mathcal{U}(-20$ cm, 500
 312 cm).

313 4.2 Co-seismic and early post-seismic slip models

314 We will start by analyzing models inferred by the independent approach as applied
 315 to the two datasets. The first model is inferred from "co" data (model COgps) and the
 316 second is estimated from the "co+post" dataset (model COPOST). The "co+post" dataset
 317 is similar to what has been used in several previous studies to infer the co-seismic slip
 318 [Cirella *et al.*, 2012; D'Agostino *et al.*, 2012; Cheloni *et al.*, 2014; Yano *et al.*, 2014; Volpe
 319 *et al.*, 2015]. The results of these two inversions will then be compared to those of the
 320 CTW inversion. For the sake of comparison, these inversions are performed without ac-
 321 counting for epistemic uncertainties. This refinement will only be added in a final inver-
 322 sion.

323 For each approach, the results are a set of 300,000 models corresponding to the
 324 most plausible samples of the full solution space whose interpretation can provide nu-
 325 merous information: posterior uncertainty of the parameters, trade-off between parameters
 326 of the model, entropy of our model, etc. As we are tied to the need of presenting our re-
 327 sults with 2D figures when the exploration is done on a parameter space of tenth of di-
 328 mensions, we choose to represent our results in 3 different ways. In the main manuscript,
 329 the first representation illustrates the relations between neighboring subfaults and the vari-
 330 ability of most probable parameters. To do so, we divide our models into 25 families, and
 331 represent the median model of each family in different pixels in each subfault (e.g. Fig-
 332 ures 3a-d, 5a-b). A selected model will be added to the first family if it is equal to the
 333 median of the 300,000 models within a tolerance of 50 cm (for each co-seismic parame-
 334 ter) or 25 cm (for post-seismic parameters). The other families of models are built iter-
 335 atively. If the selected model does not fall into the first family, it is used as model of ref-
 336 erence to define the next family. When 24 families have been created, orphaned samples
 337 are added to the 25th family (more information in Figure S2). A second representation
 338 shows the posterior Probability Density Functions which, for a particular parameter, will
 339 inform on the amount of slip uncertainty associated to each subfault (e.g. Figures 3e-g,
 340 5c-e). Our third representation shows the median models of the 300,000 inferred samples
 341 in map view (e.g. Figures a and b in S3, S5, S9).

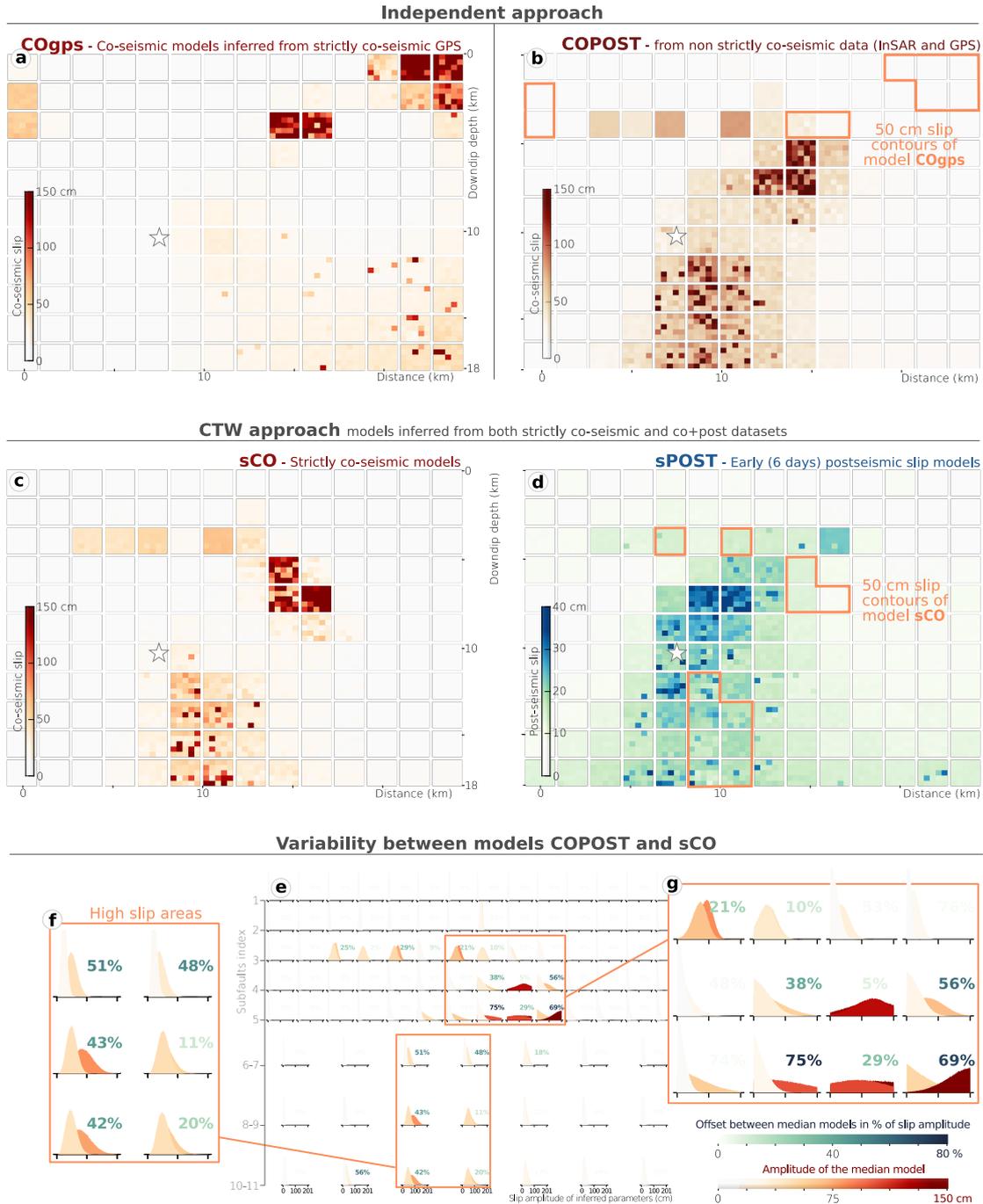


Figure 3. Comparison between finite-fault models inferred with the independent or the CTW approach. (a) Strictly co-seismic (30 s after the mainshock) slip model, named COgps, inferred from the strictly co-seismic dataset (GPS only). (b) Non-strictly co-seismic model COPOST inferred from the co-seismic dataset contaminated with some post-seismic deformation. (c) and (d) Strictly co-seismic sCO and early postseismic sPOST (6 days after the mainshock) slip models inferred jointly with the CTW approach. (a) to (d) illustrate the slip amplitude of the median models of 25 families of inferred models (more information in the text and Figure S2). Each subfault (large square) is divided into 25 pixels colored from the slip amplitude of the corresponding median model.

Figure 3. (Previous page.) (e) Comparison between the posterior Probability Density Functions of models COPOST (b) and sCO (c), colored from the amplitude of their median model. In the last four rows, the PDFs show the repartition of parameters for patches covering 2 subfaults along strike and 2 subfaults along dip (i.e. patches two times bigger than for the first four rows). The COPOST model PDFs are in the background while the sCO PDFs are in the foreground. The offset between the median models is shown as percentage with a different color scale. Two high slip areas are illustrated in detail: the highest slip patch (g) and the deep slip patch (f).

342 **4.2.1 Approaches assuming independent datasets**

343 When solving for the model COgps, we find that most of the slip is concentrated
 344 in the shallow parts of the fault (Figures 3a, S3a,c,e). Slip amplitudes reach 230 cm in
 345 the first two kilometers below the Earth surface. These values largely contradict field
 346 observations hardly reporting more than 15 cm of surface offset [*Faluccci et al.*, 2009;
 347 *Vittori et al.*, 2011]. This contradiction probably derives from our limited set of observa-
 348 tions, with only 4 GPS stations documenting the rupture in the near field (Figure 4a). The
 349 COgps model is thus largely under-determined and unlikely to represent a reliable im-
 350 age of the co-seismic deformation. In contrast, the COPOST slip model is inferred from
 351 a more populated dataset extending over a large part of the Paganica fault (Figures 3b,
 352 S3b,d,f). The patch of highest slip amplitude, reaching more than 150 cm, is well con-
 353 strained and located between 5 to 7 km depth (Figure 3b). Up to 100 cm slip is also in-
 354 ferred below the epicenter. The scalar seismic moment of model COPOST, calculated with
 355 $\mu = 3.5 \cdot 10^{10}$, is $M_0 = 4.9 \pm 0.67 \cdot 10^{25}$ dyne.m. This value corresponds to a M_w 6.4 earth-
 356 quake rather than a M_w 6.3. The comparison between observations and predicted surface
 357 displacement is shown in Figure 4 for the GPS datasets and in Figure S4 for the interfer-
 358 ograms. As expected, the COgps model well explain the "co" dataset (Figure 4b), but its
 359 predictions hardly fit the interferograms of the "co+post" dataset (Table S2). In contrast,
 360 the predicted surface displacement of the COPOST model well approaches the "co+post"
 361 observations (Figures 4b and S4), with limited residuals (Tab. S2).

362 **4.2.2 Dual time approach, without epistemic uncertainties**

363 With the CTW approach, we infer two slip models: the strictly co-seismic model
 364 sCO (see Figures 3c, S5a,c,e and S6 for an animated compilation of probable models)
 365 and the model sPOST which reflects the 6 days displacement following the mainshock
 366 (Figures 3d and S5b,d,f). The model sCO, exploiting information from both the "co" and
 367 "co+post" datasets, is in agreement with the main characteristics of the COPOST model
 368 (Figure 3b): the location and amplitude of the maximum slip patch are comparable, and
 369 a large amount of slip is also inferred at depth, up to 75 cm on average and exceeding
 370 150 cm for some models (Figure 3c). However, unlike the COPOST model, the two main
 371 slip patches of the sCO model are delimited by an unruptured area (Figure 3c). Overall,
 372 the two models differ on average by 44% and by up to 75% for some subfaults character-
 373 ized by high slip amplitudes (Figures 3e-g, S7), mainly because of the variability of the
 374 amount of slip inferred below 5 km depth. For the model sCO, $M_0 = 3.50 \pm 0.63 \cdot 10^{25}$
 375 dyne.m, corresponding to the moment magnitude value (GCMT) of 6.3.

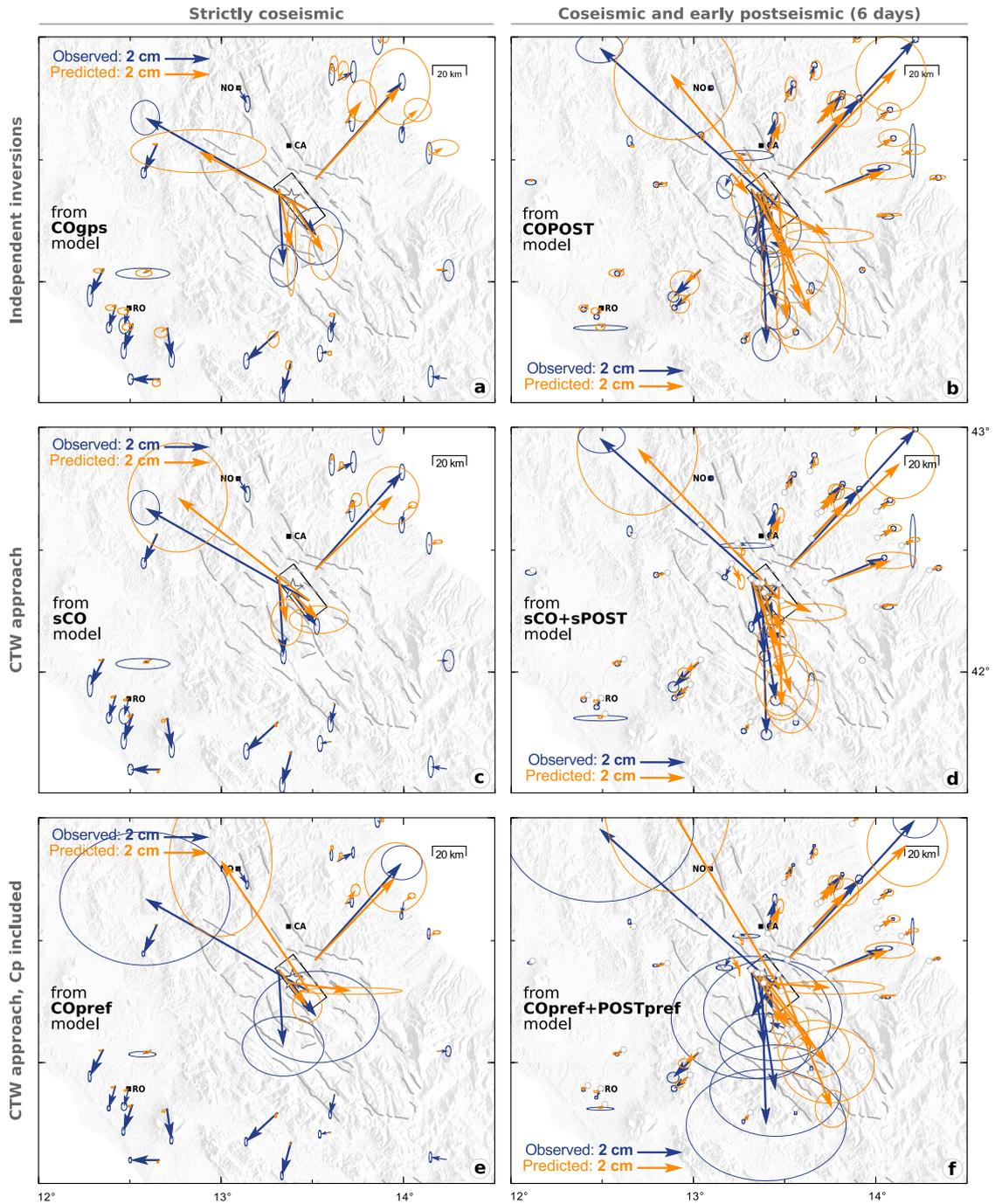


Figure 4. Comparison of horizontal surface displacement at GPS stations. Strictly co-seismic displacement is shown on the left while "co+post" displacement is shown on the right. Observed surface displacement is in blue with 95% confidence ellipses. Predictions are in orange with 95% confidence ellipses. In the top and middle rows, observational confidence ellipses (in blue) include only data errors. (a) and (b) The predictions have been calculated independently: using "co" data (a) and the "co+post" dataset (b). In (c) and (d), predictions are derived from the CTW approach. (e) and (f) show the predictions for a similar inversion setup, except epistemic uncertainties have been added to the data errors, enlarging the confidence ellipses. Our fault geometry is shown with a black rectangular box. The cities of Norcia (NO), Campotosto (CA) and Roma (RO) are indicated with black squares. Major seismogenic faults are shown in gray solid lines and the epicenter is the white star.

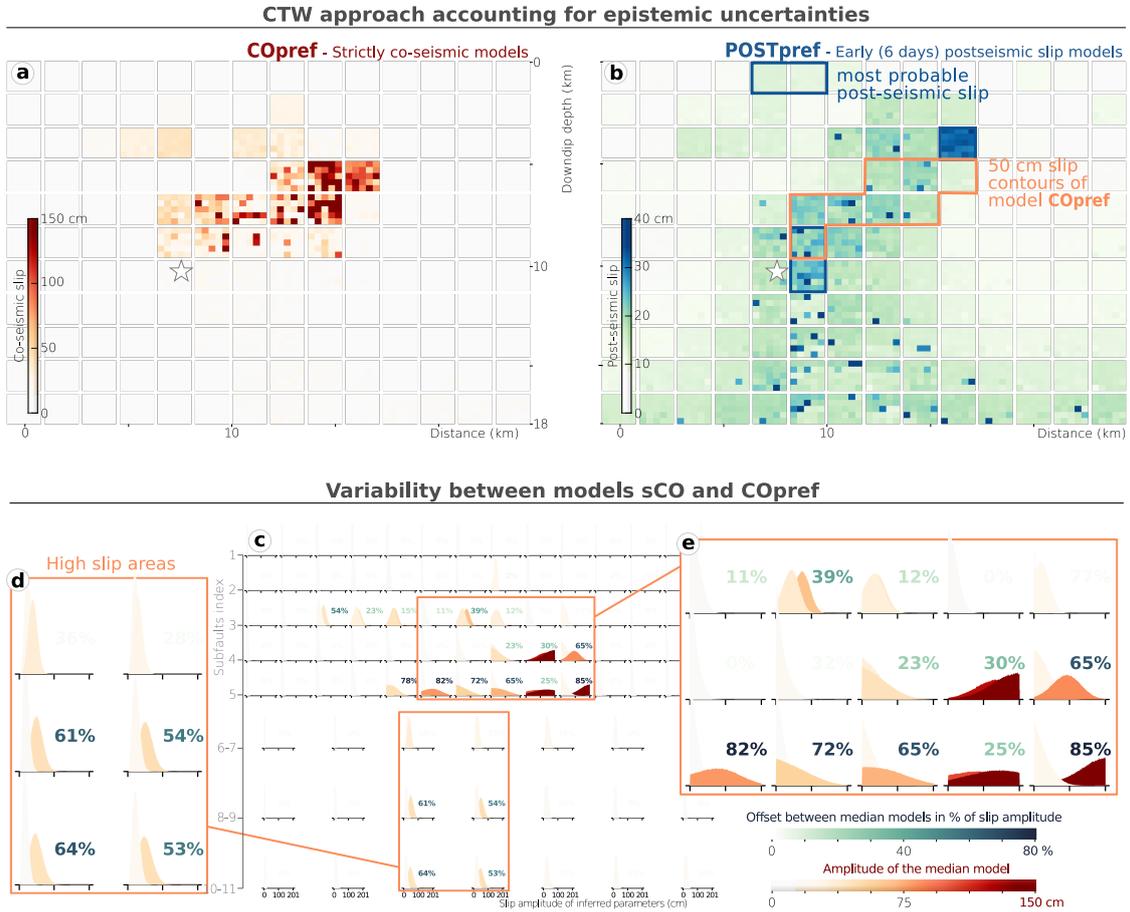


Figure 5. (a) and (b) Strictly co-seismic COpref and early postseismic POSTpref (6 days after the main-shock) preferred slip models, inferred with the CTW approach and accounting for epistemic uncertainties. (c) Comparison between the posterior Probability Density Functions of models COpref (a) and sCO (Figure 3c), colored from the amplitude of their median model. In the last four rows, the PDFs show the repartition of parameters for patches covering 2 subfaults along strike and 2 subfaults along dip (i.e. patches two times bigger than for the first four rows). The COpref model PDFs are in the foreground while the sCO PDFs are in the background. The offset between the median models is shown as percentage with a different color scale. Two high slip areas are illustrated in detail: the highest slip patch (e) and the deep slip patch (d).

376 With the CTW approach, we also find that a large portion of the fault slipped during
 377 a 6 days time window after the mainshock (Figure 3d), with maximum amplitude of
 378 30 cm in the dip-slip direction (Figure S5b). The largest post-seismic slip (> 45 cm) is
 379 located between the co-seismically ruptured patches (Figure 3c), and is well constrained
 380 with only 15 cm of posterior uncertainty (Figure S5f). Overall, post-seismic slip (30 cm
 381 and below) tends to locate around the highest co-seismic slip patch and the epicenter, but
 382 also overlaps the deepest co-seismic slip patch. Yet, below 10 km depth, the posterior un-
 383 certainty can reach 100% of the median slip amplitude, meaning that it is difficult to inter-
 384 pret anything at that level of detail (Figures S5d,f). The seismic moment of model sPOST
 385 is $M_0 = 1.58 \pm 0.63 \cdot 10^{25}$ dyne.m. The predicted surface displacements fit well the ob-
 386 servations (Figures 4c,d and Figure S8) with residuals similar to the ones of the COPOST
 387 model (Tab. S2).

388 As expected, the areas of largest post-seismic slip in the sPOST model correspond
 389 to the locations of largest divergence between COPOST and sCO models (3b-g). In sum-
 390 mary, usual approaches using independent datasets do not allow us to infer reliable im-
 391 ages of the strictly co-seismic and early post-seismic phases. Whereas the "co+post" slip
 392 model is reliable, the "co" model is not robust enough to retrieve the early afterslip from
 393 the subtraction of these two slip distributions. Additionally, the scalar seismic moment of
 394 model "co+post" corresponds to a moment magnitude greater than the GCMT M_w of 6.3.
 395 In contrast, the CTW approach allows us to infer robust estimates of both co-seismic and
 396 post-seismic slip, to exploit all the information collected within our geodetic observations,
 397 and to correctly estimate the seismic moment. However, the reliability of these models can
 398 be questioned as they do not account for uncertainties in the forward model.

399 **4.2.3 Dual time approach, accounting for epistemic uncertainties**

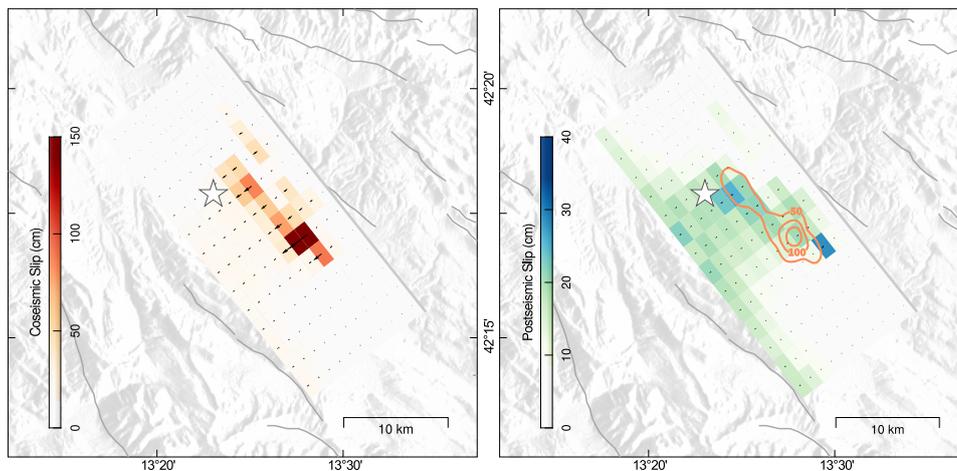


Figure 6. Our preferred slip models of the L'Aquila earthquake, inferred with the CTW approach and accounting for epistemic uncertainties. (a) and (b) show the dip-slip amplitude and rake of the average model in map view, the epicenter being the white star. In (b), orange lines also show the 50 cm co-seismic slip contours.

400 Here, we present the results of the CTW approach, as in the previous section, but
 401 accounting for epistemic uncertainties. We will name the resulting models COpref and
 402 POSTpref since they correspond to our preferred approach providing the most complete
 403 and objective evaluation of the problem (see Figures 5, 6, S9 and S11 for an animated
 404 compilation of probable co-seismic models). The distribution of the co-seismic slip differs
 405 by 42% on average and by up to 88% locally from the models inferred without accounting

406 for uncertainties (Figures 5c, S10). The co-seismic slip is now limited (on average) to a
 407 single 10 km long patch located between 5 and 10 km depth, reaching more than 150 cm
 408 amplitude to the south (right-hand side of Figure 5a). The corresponding scalar seismic
 409 moment $M_0 = 3.03 \pm 0.64 \cdot 10^{25}$ dyne.m is slightly lower than what was estimated for the
 410 model sCO but is still very close ($M_w = 6.28 \pm 0.06$) to the (GCMT) value of M_w 6.3.

411 Compared to the COpref model, the main characteristics of the POSTpref model
 412 are not strongly affected by the inclusion of uncertainties. Overall, post-seismic slip (20
 413 cm and below) occurs mostly below the co-seismic high slip patch, where almost no (less
 414 than 20 cm) co-seismic slip is imaged. Subfaults with the largest post-seismic slip (more
 415 than 40 cm) tend to be located around or on the edges of the co-seismic high slip patch
 416 (Figure 5b). The presence of large post-seismic slip below 10 km depth is unlikely as the
 417 posterior uncertainty reaches 150% of the median slip (Figure S19f). Thus, only 3 narrow
 418 zones most probably slipped post-seismically (see Figure 5b, and a comparison of median
 419 and maximum a posteriori models in Figure S12). M_0 is similar to model sPOST with a
 420 value of $1.60 \pm 0.63 \cdot 10^{25}$ dyne.m. The addition of epistemic uncertainties has increased
 421 the residuals between observations and predictions (see Tab. S2 and Figure S13). This
 422 behavior was expected as the inclusion of C_p allows the inversion to tolerate for larger
 423 misfits at data points where the forward model predictions are less reliable [Ragon *et al.*,
 424 2018].

425 In summary, the CTW approach shows that if early post-seismic is not acknowl-
 426 edged as post-seismic signal, co-seismic models may be biased by more than 40% on av-
 427 erage and of up to 75% locally. But we also learn from these different tests that adding
 428 more data into the problem is not sufficient, and epistemic uncertainties remain critical for
 429 the inference of a reliable model. Altogether, our results emphasize the need to account
 430 for two types of bias in the slip models: the contamination of co-seismic observations by
 431 some early post-seismic signal, and not acknowledging for the uncertainties associated to
 432 the forward problem.

433 5 Discussion

434 5.1 Discussion of the CTW approach

435 Observations of co-seismic or post-seismic processes are often contaminated by
 436 other sources of deformation (mainly post-seismic or co-seismic, respectively) and are
 437 widely used, when non-contaminated data are rare and scarcely distributed. Optimizing
 438 the use of the information content in each dataset is thus critical to improve the robust-
 439 ness of both co-seismic and post-seismic slip models. A first approach would be to ac-
 440 count for potential uncertainties in the co-seismic model due to early afterslip in the form
 441 of a covariance matrix, as already proposed in *Bletery et al.* [2016]. While this approach
 442 helps inferring more reliable co-seismic models at a low computational cost, it does not
 443 allow us to estimate the early afterslip and needs a prior evaluation of the amount of af-
 444 terslip considered as co-seismic signal. Another strategy would be to jointly infer "co"
 445 and "co+post" data as if they were strictly co-seismic, and to select models that better
 446 explain the "co" observations, as in [Chlieh *et al.*, 2007]. In this case, the computational
 447 cost is increased because several models have to be tested. Additionally, with these ap-
 448 proaches the early post-seismic slip is not estimated. In contrast, the CTW approach we
 449 use in this study allows us to discriminate co-seismic from early post-seismic slip and to
 450 reliably estimate corresponding slip models. Our approach takes advantage of the InSAR
 451 data that recorded both co- and post-seismic deformation to help constrain both strictly
 452 co- and early post-seismic models.

453 Our results on the L'Aquila event show that the early afterslip, here corresponding
 454 to 6 days after the co-seismic rupture, can reach a fourth to a third of the amplitude of
 455 the co-seismic slip. If the early afterslip is acknowledged as co-seismic signal, co-seismic

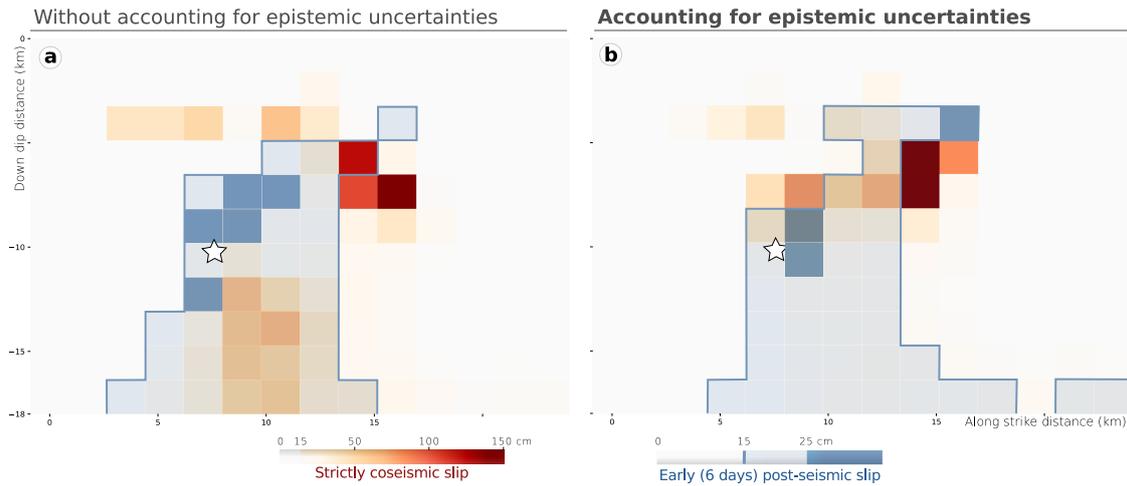


Figure 7. Effect of epistemic uncertainties (C_p) on the distribution of strictly co-seismic slip and after-slip. The slip models have been inferred accounting for epistemic uncertainties (b) or not (a). The strictly co-seismic slip median model is in light gray to dark orange colorscale. The subfaults that slipped of more than 15 to 25 cm up to 6 days after the mainshock, according to our median model, are in transparent light to medium blue. The afterslip does not overlap the co-seismic slip when C_p is accounted for (b), whereas the two slip distributions overlap at depth when no C_p is included (a).

456 models of the L'Aquila event are biased. The impact of early afterslip on the co-seismic
 457 models is particularly large in the case of the L'Aquila event and questions the generic
 458 nature of this result. Overall, early afterslip remains poorly studied but has been shown
 459 to range from 0.6% to more than 8% of the co-seismic peak slip in the first 3-4 hours
 460 following an earthquake [respectively for the 2009 great Tohoku-Oki earthquake and the
 461 2012 M_w 7.6 Nicoya earthquake, *Munekane, 2012; Malservisi et al., 2015*]. Thus, that the
 462 post-seismic deformation ongoing 6 days after the mainshock reaches up to 20 % of the
 463 co-seismic slip of the L'Aquila earthquake might not be an extreme case.

464 Our tests also demonstrate that models are largely impacted by the introduction of
 465 epistemic uncertainties (Figure 7). This impact could mean the assumed fault and Earth
 466 properties are not realistic enough to capture the real seismic rupture, and/or that small
 467 variations of the fault geometry (slight curvature, roughness) or of the Earth model (3D
 468 heterogeneities) largely affect our slip models. The influence of epistemic uncertainties is
 469 greater on the co-seismic model, as expected from the fact that these uncertainties scale
 470 with the amount of slip [*Duputel et al., 2014; Ragon et al., 2018*]. Accounting for uncer-
 471 tainties of the forward model allowed us to exclude the possibility of deep slip for the co-
 472 seismic models, but not totally for the post-seismic models probably because of the much
 473 lower slip amplitudes. Additionally, accounting for C_p prevented the most probable co-
 474 and post-seismic slips to overlap in deeper parts of the fault. The inclusion of epistemic
 475 uncertainties acts like a smoothing constraint on the slip distribution, but with a smoothing
 476 factor being controlled by the inaccuracies of the forward problem.

477 5.2 Non-unicity of co-seismic and afterslip models of the L'Aquila earthquake

478 Our results on the L'Aquila event indicate that the strictly co-seismic slip is con-
 479 centrated in a thin horizontal band located between 5 and 7 km depth and reaching more
 480 than 150 cm in amplitude at its southern end, with no large slip amplitudes inferred be-
 481 low 8 km depth (Figures 6 and 8). The highest amplitude is reached at about 6 km depth
 482 south west of the epicenter, a rupture area also imaged in the co-seismic models of *Gua-*

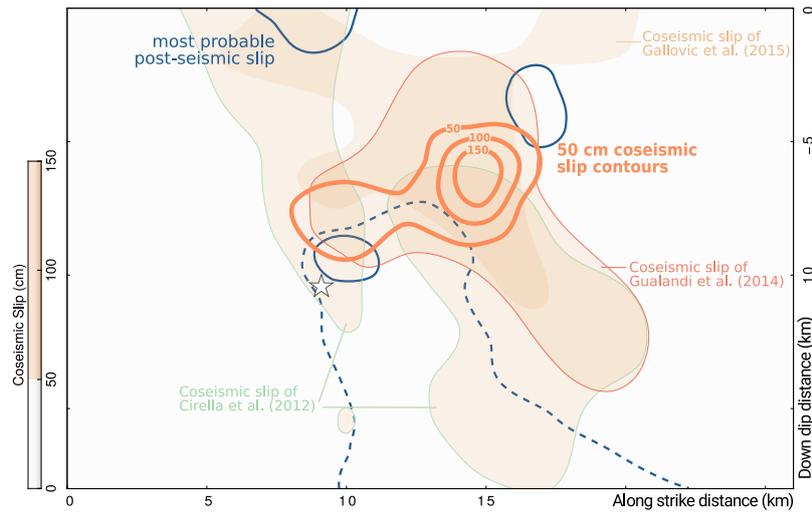


Figure 8. Comparison between the slip distributions inferred with the CTW approach and co-seismic slip distributions of other studies. The strictly co-seismic slip of *Gualandi et al.* [2014] inferred from GPS only, the strictly co-seismic slip of *Gallovic et al.* [2015] inferred from accelerometric and high rate GPS data, the co-seismic slip of *Cirella et al.* [2012] inferred from GPS, InSAR and strong motion, are projected in our fault plane in transparent light orange when slip exceeds 50 cm. The 50 cm contours of our strictly co-seismic slip distribution and the 15 cm contours of our afterslip inferred accounting for epistemic uncertainties are in bold lines, respectively orange and dark blue. The epicenter is the white star.

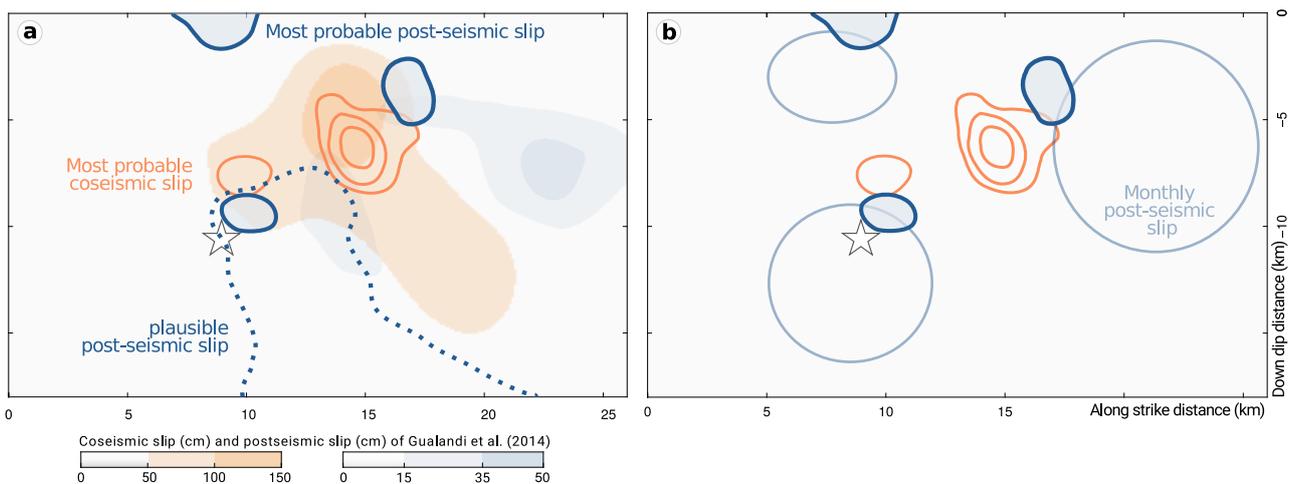


Figure 9. Comparison between our most probable strictly co- and post-seismic slip distribution 6 days after the mainshock and the post-seismic slip up to 306 days after the mainshock. (a) Our most probable slip distributions are represented with bold orange and dark blue lines, respectively for co-seismic (50 cm contours) and post-seismic slip (slip > 10 cm). The area of afterslip delimited with a dotted blue line is considered as less plausible as inferred with large uncertainties. The co-seismic slip and afterslip 306 days after the mainshock inferred by *Gualandi et al.* [2014] are plotted with the same color codes but as color swaths. (b) Our results are compared to the area that slipped post-seismically during about 6 months (176 and 194 days respectively) after the mainshock as modeled by *D'Agostino et al.* [2012] and *Cheloni et al.* [2014]. The epicenter is the white star.

483 *landi et al.* [2014], *Gallovic et al.* [2015] and *Cirella et al.* [2012] (inferred respectively
 484 from GPS only, from accelerometric and high rate GPS data, and from GPS, InSAR and
 485 strong motion, see Figure 8) and most of other authors [*Atzori et al.*, 2009; *Trasatti et al.*,
 486 2011; *D'Agostino et al.*, 2012; *Serpelloni et al.*, 2012; *Cheloni et al.*, 2014; *Balestra and*
 487 *Delouis*, 2015; *Volpe et al.*, 2015]. It is the only recurrent pattern we can notice between
 488 the 4 slip models of Figure 8. Indeed, while we do not image any shallow slip, other pub-
 489 lished slip models do with up to 1.5 m in amplitude [Figure 8, *Cirella et al.*, 2012; *Volpe*
 490 *et al.*, 2015]. At greater depths, most authors infer large slip amplitudes while our pre-
 491 ferred model shows no slip below 8 km depth.

492 The imaged patches of post-seismic slip (>15 cm) are located around our co-seismic
 493 slip, near its epicenter and southern end. Interestingly, our inferred post-seismic slip is
 494 also located near areas that ruptured co-seismically as inferred by other studies (Figure 8).
 495 The post-seismic slip that occurred several days to months after the mainshock is char-
 496 acterized by 3 wide slip areas, located SW of the main co-seismic slip patch, above the
 497 epicenter close to the surface, and around the epicenter [*D'Agostino et al.*, 2012; *Che-*
 498 *loni et al.*, 2014; *Gualandi et al.*, 2014]. Most of these post-seismic models acknowledge
 499 the first days of post-seismic signal as a co-seismic deformation. While we infer likely
 500 afterslip in similar locations, the afterslip patches are limited to narrower areas near the
 501 co-seismic rupture (Figure 9). Most of these longer-term post-seismic models cover time
 502 periods ranging from 6 days to 9 months after the mainshock, they overlook a large part
 503 of the early post-seismic deformation. Thus, the peak amplitude of the early afterslip is
 504 up to 3 times larger than what was imaged for several months by *D'Agostino et al.* [2012]
 505 and *Cheloni et al.* [2014].

506 Our results show that the amplitude and distribution of long-term afterslip may be
 507 largely underestimated (here by a factor of 3) if the deformation occurring the first few
 508 hours to days after the mainshock is not accounted in the post-seismic budget. Thus, over-
 509 looking the early part of the postseismic phase measured in geodetic data may not only
 510 bias the estimates of the coseismic slip, but also our estimates of the postseismic phase.

511 **5.3 Fault frictional properties and relationship between afterslip and aftershocks**

512 The comparison between our early post-seismic model and images of longer-term
 513 post-seismic slip suggest that afterslip may nucleate preferably around the co-seismic rup-
 514 ture in the days following the mainshock (Figure 9b). Afterwards, the afterslip propagates
 515 and extends, both along-dip and laterally, away from the co-seismic slip [*D'Agostino et al.*,
 516 2012; *Cheloni et al.*, 2014; *Gualandi et al.*, 2014]. This behavior agrees with models ex-
 517 plaining afterslip as a result of rate dependent friction behavior. Indeed, in these models
 518 the afterslip relaxes the stress increment induced in velocity-strengthening area by the
 519 co-seismic rupture [*Marone et al.*, 1991; *Perfettini and Avouac*, 2004]. The post-seismic
 520 sliding thus nucleates close to the mainshock asperity and propagates with time outward
 521 from the rupture zone. That early afterslip relates to the stress changes induced by the
 522 co-seismic rupture has also been modeled for other events [e.g. the M_w 8.0 Tokachi-oki,
 523 M_w 7.6 Chi-Chi, and the M_w 6.0 Parkfield, *Miyazaki et al.*, 2004; *Chan and Stein*, 2009;
 524 *Wang et al.*, 2012, respectively].

525 Additionally, our results show that early afterslip nucleate within narrow areas (1-2
 526 km wide), and does not happen everywhere around the co-seismic rupture. Areas sliding
 527 aseismically just after the mainshock are thus limited in size around the co-seismic rup-
 528 ture, suggesting that frictional properties vary at a small-scale around the rupture zone. It
 529 may also suggest that the regions adjacent to co-seismic rupture are potentially unstable
 530 (i.e. are steady-state velocity weakening). This interpretation agrees with the results of
 531 *Gualandi et al.* [2014] suggesting the longer-term afterslip regions, that are also located
 532 farther away from the co-seismic ruptured zone, are characterized by a transition between
 533 velocity weakening and velocity strengthening behavior. This implies that co-seismic rup-

534 ture occurs and triggers early afterslip in velocity weakening regions; while afterslip prop-
 535 agates away from the ruptured zone in fault regions that progressively become stable with
 536 the distance to the mainshock.

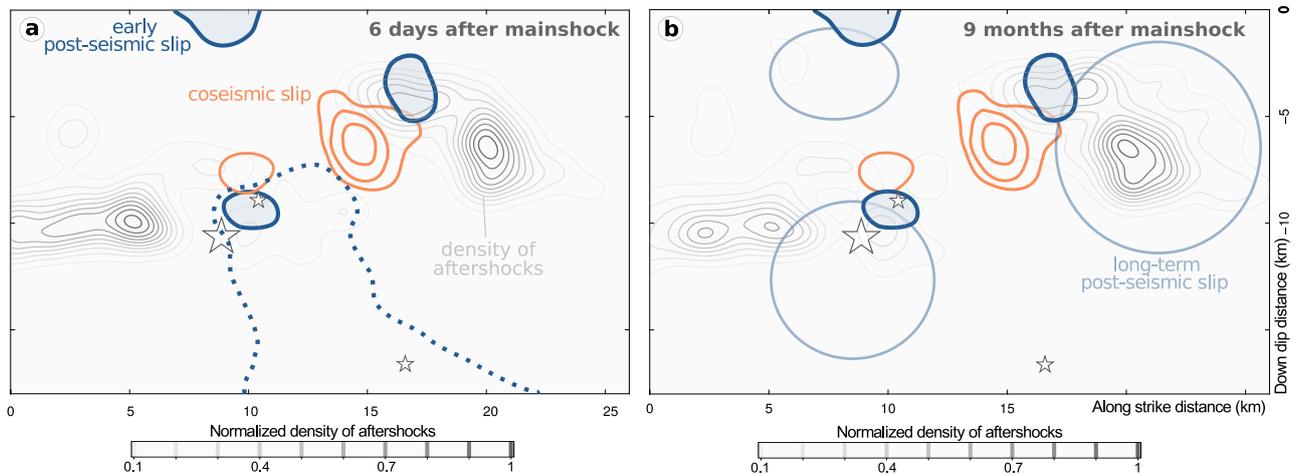


Figure 10. Distribution of most probable co-seismic slip and afterslip models, and the normalized density of aftershocks that occurred (a) within 6 days after the mainshock or (b) up to 9 months after the mainshock [$M_C = 0.88$, catalog of *Valoroso et al.*, 2013]. The strictly co-seismic 50 cm slip contours are in orange, while the contours of most probable afterslip (slip > 10 cm) are in blue. The area delimited by dotted blue lines has plausibly hosted some afterslip, but not as probably as the other regions. The areas that slipped post-seismically during about 6 months after the mainshock as modeled by *D’Agostino et al.* [2012], *Cheloni et al.* [2014] and *Gualandi et al.* [2014] are the blue circles. The density of aftershocks located within 3 km of the fault (to account for potential uncertainty of the fault geometry) is calculated with a kernel density estimation method [*Parzen*, 1962] with a smoothing factor of 0.6. The cumulated number of aftershocks of $M_C = 0.88$ is of ~ 6000 6 days after the mainshock and 8 times larger 9 months after the mainshock (Figure S14). The epicenter and aftershocks of $M_w \geq 4.4$ are the white stars.

537 In Figure 10, we compare the slip distributions imaged for the mainshock and 6
 538 days after, with the distribution of aftershocks detected over 6 days and 9 months after the
 539 mainshock [*Valoroso et al.*, 2013]. As for many earthquakes, aftershocks are distributed
 540 mainly at the ends of the fault [*Das and Henry*, 2003] with few events located near the
 541 co-seismic rupture. Six days after the mainshock (Figure 10a), our results show no clear
 542 correlation between the location of early afterslip and aftershocks. Months after the main-
 543 shock, the areas with a high density of aftershocks are similar to 6 days after the main-
 544 shock [as suggested by *Henry and Das*, 2001, whereas the cumulated number of afters-
 545 hocks is 8 times larger, see Figure S14] and the post-seismic slip has extended farther
 546 away from the co-seismic rupture. This is why we can observe a spatial correlation be-
 547 tween some areas of long-term post-seismic slip and aftershocks [*D’Agostino et al.*, 2012;
 548 *Cheloni et al.*, 2014]. The spatial correlation is particularly striking for the southern afterslip
 549 patch, for which few early aftershocks are located within the early afterslip area
 550 (Figure 10a) while the aftershock cluster overlies the monthly afterslip that propagated
 551 outward from the co-seismically ruptured zone (Figure 10b).

552 From our results, we can thus draw only one conclusion: there is no correlation be-
 553 tween the area of large (>15 cm) early afterslip and the location of aftershocks for the
 554 first few days after the mainshock. This conclusion contradicts the observations made for
 555 some other earthquakes although mainly at longer time scales [e.g. *Hsu et al.*, 2006; *Per-
 556 fectini and Avouac*, 2007; *Wang et al.*, 2012; *Ross et al.*, 2017, for time periods spanning

557 respectively 11 months, 3.5 years, 5 days and 2.5 months]. Our results could also suggest
 558 that, for some parts of the fault, aftershocks nucleation precedes aseismic slip that occur
 559 months after the mainshock; aftershocks could thus be partly explained by stress changes
 560 due to the co-seismic rupture. But these aftershocks could also be triggered by early af-
 561 terslip with an amplitude so low that it is not inferred by our model. The absence of clear
 562 correlation between early afterslip and aftershocks may also be related to the presence of
 563 high pressure fluids in the seismogenic zone of the L'Aquila event, and of Central Italy
 564 in general, with the widespread emissions of CO₂ rich fluids for deep origin [*Chiodini*
 565 *et al.*, 2000; *Frezzotti et al.*, 2009; *Chiodini et al.*, 2011]. Already, *Miller et al.* [2004] and
 566 *Antonoli et al.* [2005] proposed that the aftershocks and spatio-temporal migration of
 567 the seismicity of the 1997 Umbria-Marche seismic sequence (80 km NE of the L'Aquila
 568 event) were driven by the co-seismically induced fluid pressure migration. Similarly, the
 569 increase in seismicity rate of the L'Aquila earthquake and the occurrence of some after-
 570 shocks may have been driven by fluid flows [*Luccio et al.*, 2010; *Terakawa et al.*, 2010;
 571 *Malagnini et al.*, 2012]. High pressure fluids have been observed before the co-seismic
 572 rupture, and may have impacted the nucleation phase of the L'Aquila earthquake [*Lucente*
 573 *et al.*, 2010]. Finally, *Malagnini et al.* [2012] show that the strength of the Campotosto
 574 fault, just north of the main rupture (see Figure 2), has been controlled by fluid migration
 575 for at least 6 days after the mainshock, a time window corresponding to our study of early
 576 afterslip. The perturbations in pore fluid pressure induced by the co-seismic rupture may
 577 have triggered the first aftershocks of the L'Aquila earthquake. Fluid migration may have
 578 prevented aftershocks and early afterslip to affect the same areas of the fault, especially if
 579 the increase in fluid pressure first produced aseismic slip, followed by triggered seismicity
 580 around the pore pressure front [*Miller et al.*, 2004]. Finally, if early aftershocks were trig-
 581 gered by changes in fluid pressure, it may justify the possibility that some of these early
 582 aftershocks nucleated before the occurrence of long-term afterslip in similar regions of the
 583 fault.

584 6 Conclusion

585 In this study, we use a simple and efficient approach to account for the differences
 586 in temporal resolution of various geodetic datasets. A redesign of the Green's Functions
 587 matrix allows us to optimize the use of the information content of datasets covering differ-
 588 ent time periods. With this approach, we image simultaneously the strictly co-seismic slip
 589 and the early afterslip (6 days after the mainshock) of the 2009 M_w6.3 L'Aquila earth-
 590 quake using two datasets: one covers the two slip episodes (e.g. InSAR) while the other
 591 records the co-seismic signal only (e.g. continuous GNSS). We show that when the two
 592 phases are inverted independently, as is usually the case, the estimated slip distributions
 593 are not reliable because strictly co-seismic observations are usually of poor spatial reso-
 594 lution. Additionally, overlooking the early post-seismic deformation results in models that
 595 overestimates the co-seismic slip, and underestimates the total post-seismic slip budget.
 596 In contrast, our approach allows us to accurately estimate both co-seismic and early post-
 597 seismic slip models.

598 Our results show that neglecting the contribution of the early post-seismic defor-
 599 mation will likely bias estimates of the co-seismic and/or the post-seismic slip. For our
 600 test case of the L'Aquila earthquake, the peak co-seismic slip is likely 30% greater when
 601 early post-seismic signal is recorded as co-seismic deformation. The long-term afterslip
 602 estimates are underestimated by a factor 3 when the first 6 days of post-seismic defor-
 603 mation are not acknowledged. Our investigation of the L'Aquila event also stressed the
 604 strong influence of uncertainties in the forward model, mainly stemming from our imper-
 605 fect knowledge of the fault geometry and the Earth structure, on the imaged slip distribu-
 606 tions. These uncertainties alone are sufficient to cause contradictory interpretations on the
 607 slip history on the fault (e.g. with the existence of shallow or dip slip).

608 Our preferred slip model for the L'Aquila earthquake tends to be simpler than many
 609 previous models, with one thin horizontal band of slip located around 7km depth, reach-
 610 ing 150cm in amplitude near its southern end. Our model thus excludes the possibility
 611 of major shallow or deep co-seismic slip patches (less than a few km or deeper than 10).
 612 The early post-seismic slip (6 days after the mainshock) was limited to the same inter-
 613 mediate depth range (7 km +/- 3 km), initiating on the edges of the co-seismic slip, with
 614 possibly some overlap. Some afterslip may also have occurred at greater depths. A com-
 615 parison with longer term afterslip models suggest that the early afterslip patches might
 616 have simply expanded over time from their initial position. Aftershocks are more spatially
 617 distributed (7 km +/- 5 km) but still concentrated at intermediate depth. Several studies
 618 suggest that aftershocks might be driven by afterslip [e.g., *Perfettini and Avouac*, 2007;
 619 *Hsu et al.*, 2006; *Sladen et al.*, 2010; *Ross et al.*, 2017; *Perfettini et al.*, 2018] but here af-
 620 tershocks are only partially overlapping. This result suggests that post-seismic reloading
 621 may be influenced by fluids as advocated in several previous studies [e.g., *Luccio et al.*,
 622 2010; *Terakawa et al.*, 2010; *Malagnini et al.*, 2012; *Guglielmi et al.*, 2015; *Scuderi and*
 623 *Collettini*, 2016].

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 632 livet was used to build inputs for the Bayesian algorithm. Discussed slip models of *Gua-*
 633 *landi et al.* [2014], *Gallovic et al.* [2015] and *Cirella et al.* [2012] have been retrieved from
 634 SRCMOD [<http://equake-rc.info/SRCMOD/>], *Mai and Thingbaijam*, 2014]. Figures were
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