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1 **Implications of impact-energy dependent erosional efficiency on bedrock river sediment dynamics and**
2 **form: 1. Reach-scale dynamics and the effective flood**

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10 **Key Points:**

- 11 • Incorporation of experimentally-constrained, impact-energy-dependent rock resistance
12 into a model for impact-driven erosion evaluates its influence on bedrock incision rates
- 13 • Impact-energy-dependent rock resistance strongly amplifies grain-size controls on
14 bedrock detachment per unit impact
- 15 • Differences in detachment efficiency with increasing impact energy emphasizes the
16 relative contribution of rarer, high-magnitude floods

17

Abstract

Bedrock river incision reflects the cumulative geomorphic work performed across a distribution of flood magnitudes and frequencies. Mechanistic models of bedrock incision by bedload impacts typically assume that bedrock resistance to erosion is constant with respect to particle impact energy. However, recent impact experiments demonstrate that rock resistance to erosion decreases systematically as impacts become more energetic. Here we incorporate experimentally constrained, impact-energy-dependent rock-resistance coefficient into the widely used saltation-abrasion model and evaluate how this modification alters predicted incision across a range of grain sizes and transport conditions. We find that allowing resistance to scale with impact energy strongly amplifies the grain-size dependence of detachment, such that coarse grains remove substantially more bedrock rock per impact. Applying the model across a full discharge distribution to evaluate long-term incision reveals that while impact-energy-dependent rock resistance does not meaningfully alter the effective floods, it shifts cumulative incision toward rarer, high-magnitude events. These results demonstrate that accounting for impact-driven differences in erosional efficiency fundamentally alters how geomorphic work is distributed across floods, increasing the contribution of extremes to long-term bedrock erosion rates.

Plain Language Summary

Rivers carve into bedrock when moving sediment grains strike the channel bed. How fast a river incises depends not only on how large floods are, but also on how often they occur and how energetic sediment impacts are during those floods. Many models assume that bedrock resistance to impacts is constant, meaning low- and high-energy collisions are equally efficient at breaking rock. However, recent experiments show that more energetic impacts erode bedrock more effectively. We incorporate this behavior into a widely used bedrock erosion model to test how it alters predicted incision across different grain sizes and flow conditions. As impact energy increases, coarse grains remove substantially more rock per collision. We find that this effect that is stronger than changes to the efficiency of erosion caused by increasing flow intensity alone. When we account for the full range of floods a river experiences, from frequent small events to rare large ones, we find that energy-dependent erosion increases the relative importance of extreme floods. Our results suggest that very large floods may contribute more to long-term bedrock incision than previously expected when erosion efficiency increases with sediment impact energy.

1 Introduction

River incision into bedrock drives landscape evolution, with both channel form and incision rate reflecting the combined influence of tectonics (Kirby & Whipple, 2012; Wobus et al., 2006), climate (Ferrier, Huppert, et al., 2013; Ferrier, Perron, et al., 2013), hydrology (Forte

56 et al., 2022; Lague et al., 2005), and geology (Duvall et al., 2004; Mitchell & Yanites, 2021).
57 Because bedrock incision is slow and often concentrated in rare, high-magnitude events (e.g.
58 Hartshorn, 2002), it is difficult to measure or model directly. Most models treat incision as a
59 detachment-limited process, neglecting downstream sediment transport (e.g. Howard & Kerby,
60 1983; Seidl & Dietrich, 1992; Tucker & Bras, 2000). In this framework, erosion rate scales with
61 flow intensity – commonly expressed as unit stream power or total or excess boundary shear
62 stress. Unit stream power is typically calculated as the product of drainage area, A (km^2), a
63 proxy for water discharge, Q (m^3/s), and channel slope, S (m/m). This formulation, which relates
64 downstream energy dissipation to bedrock detachment, serves as the basis of the stream
65 power incision model (SPIM). Traditionally, SPIM assumes a constant effective discharge,
66 neglecting the full distribution of events that drive incision (Howard, 1994; Lague et al., 2000;
67 Tucker & Slingerland, 1994; Willgoose et al., 1991). More recent work has incorporated variable
68 discharge, demonstrating that when erosion thresholds are significant, they strongly influence
69 landscape evolution by controlling the fraction of the discharge distribution that contributes to
70 bedrock incision (DiBiase & Whipple, 2011; Lague et al., 2005; Snyder et al., 2003; Tucker &
71 Bras, 2000).

72 These thresholds represent the minimum flow intensity required to initiate measurable
73 bedrock erosion. They capture the idea that not all flows can detach bedrock. Only flows
74 exceeding this threshold, set by rock strength, sediment transport, or sediment caliber or cover,
75 can drive incision. River sediments play a central role in this process, both shielding the bed
76 through surface armoring and promoting erosion via abrasion (Finnegan et al., 2007, 2014;
77 Johnson et al., 2009; Johnson & Whipple, 2010). Most stream-power-based models, however,
78 treat rivers as effectively sediment-free, implicitly folding sediment effects into the erodibility
79 coefficient, K ($\text{m}^{1-2m}\text{t}^{-1}$), which describes the relative efficiency with which bedrock is detached
80 per unit stream power, which also summarize the combined influences of rock strength, uplift
81 rate, and climate (Han et al., 2014; Hancock et al., 2011; Leonard & Whipple, 2021; Small et al.,
82 2015). While some studies have explicitly incorporated erosion thresholds into SPIM by
83 modifying K (e.g. DiBiase & Whipple, 2011; Scherler et al., 2017), these efforts have been
84 agnostic to what sets these erosion thresholds. More recent work has begun to incorporate the
85 influence of sediment transport and sediment supply on long profile evolution and bedrock
86 channel form more directly (e.g. Gasparini et al., 2007; Yanites, 2018), but even in these
87 treatments, the physical mechanisms driving erosion are simplified.

88 In contrast, mechanistic bedrock erosion models provide a more detailed representation
89 of the processes driving bedrock erosion. The saltation-abrasion model (SAM) predicts bedrock
90 erosion rates as a function of particle impacts, adopting an energy-based approach that
91 explicitly links river erosion to sediment transport capacity and particle size (Sklar & Dietrich,
92 2004). The model integrates over many impacts using an impact rate, I_r ($\#/ \text{m}^2/\text{s}$), derived from
93 bedload flux equations. The instantaneous erosion rate, E_i (ms^{-1}), is the product of this impact

94 rate, the average volume of rock eroded by an individual impact, V_{imp} (m^3), and the fraction of
 95 bedrock that is exposed to impacts, F_e (-), such that $E_i = I_r V_{imp} F_e$.

96 While the saltation-abrasion model explicitly models the kinetic energy delivered to the
 97 bedrock riverbed, KE_{imp} (J), by solving for particle trajectories under some imposed transport
 98 capacity, it does not directly model the bedrock detachment process. Rather, like the
 99 erodibility term, K , in the SPIM, the model employs a rock-resistance coefficient, k_v , which
 100 modulates the relationship between particle impact energy and the volume of material
 101 removed per unit impact, V_{imp} , such that $V_{imp} \propto KE_{imp}/k_v$. As such, larger values of k_v imply
 102 higher rock resistance, or less efficient erosion, for a given impact energy. The original SAM
 103 implementation treats the rock resistance coefficient as a constant with a value of $k_v \sim 10^6$
 104 (Sklar & Dietrich, 2004). Importantly, this assumption of a constant k_v posits that particles
 105 impacts across all impact energies are equally efficient at detaching bedrock.

106 More recent work has shown that the rock resistance coefficient, k_v , can vary with
 107 lithology (Auel et al., 2017b), across field settings (Turowski et al., 2013), and with particle
 108 impact energy (Masteller et al., 2024). Specifically, Masteller et al., (2024) conducted impact
 109 experiments that identified a systematic increase in erosional efficiency, or decreasing rock
 110 resistance k_v , as impact energy increased. This observed increase in erosional efficiency ($1/k_v$)
 111 at high impact energies is related to an observed decrease in the coefficient of restitution, or
 112 rebound, of the impacting particles. Masteller et al., (2024) hypothesized that this observed
 113 increase in erosional efficiency was indicative of enhanced energy dissipation for high energy
 114 impacts, due to increasing crack extension lengths and likelihood of crack intersection. This
 115 hypothesized enhanced crack growth ultimately results in more efficient bedrock detachment
 116 for high energy collisions, as observed by their study.

117 This newly identified, impact energy-dependent behavior of k_v , as observed by
 118 Masteller et al., (2024), has several implications as particle impact energies can vary in both
 119 space and time and are also closely related to flow strength. Particle impact energy, KE_{imp} , is a
 120 function of both the mass of the impacting particle and the particle's velocity as it collides with
 121 the underlying bedrock. During larger floods, peak transport capacity, defined as the
 122 dimensionless shear stress exceeding the threshold for motion, τ^*/τ_c^* , increases with increasing
 123 flow depths. As transport capacities increase, flows can mobilize larger grains and accelerate
 124 saltating particles, increasing both the mass and velocity components of impact energy (Auel et
 125 al., 2017b, 2017a; Larimer et al., 2021). Extreme floods therefore deliver disproportionately
 126 energetic collisions through the combined effects of both heavier particles impacting the
 127 riverbed as well as higher particle impact velocities overall. As such, rare, high-magnitude
 128 floods may contribute more to long-term bedrock incision than expected based on transport
 129 capacity alone, elevating their relative importance to driving bedrock incision.

130 Landscape-scale applications of the saltation-abrasion framework have provided insights
 131 into how variable discharge distributions may translate into long-term incision rates (Sklar &

132 Dietrich, 2008; Sklar & Dietrich, 2006). By upscaling the SAM, Sklar and Dietrich (2006) linked
133 variable discharge, sediment transport, and long-term bedrock incision. This upscaled
134 application demonstrated that moderately high-flow events, rather than the largest floods,
135 dominate incision. Sklar and Dietrich (2006) reasoned that this relative dominance of moderate
136 floods arises because low flows fail to mobilize sediment and very high flows are too rare to
137 maintain sufficient bedrock exposure to drive continued abrasion. However, this previous
138 study assumed a constant rock resistance, ultimately weighting the relative influence of all
139 impact energies as having an equal efficiency of detaching bedrock material.

140 Long-term bedrock incision ultimately reflects the cumulative geomorphic work
141 performed across the full range of discharges, or the integrated instantaneous erosion rates
142 across flood events with a range of magnitudes and frequencies. Within this framework, the
143 effective flood emerges as the discharge that contributes most to cumulative incision over time
144 (Wolman & Miller, 1960). If erosional efficiency increases systematically with impact energy,
145 then the distribution of geomorphic work across the full range of the discharge distribution may
146 shift toward more extreme events, potentially altering the magnitude of the effective flood
147 itself. Motivated by this possibility, we extend the SAM to incorporate impact-energy-
148 dependent erosional efficiency and evaluate how this modification alters the partitioning of
149 geomorphic work across the discharge distribution. Specifically, we test whether increasing
150 erosional efficiency at high impact energies shifts the effective flood toward rarer, more
151 extreme events, or whether long-term bedrock incision rates continue to favor moderately
152 frequent floods. By combining experimentally constrained, energy-based scaling for rock
153 resistance with a variable discharge framework, we quantify how particle impact energy, flow
154 magnitude, and sediment properties together govern equilibrium bedrock incision at the reach
155 scale.

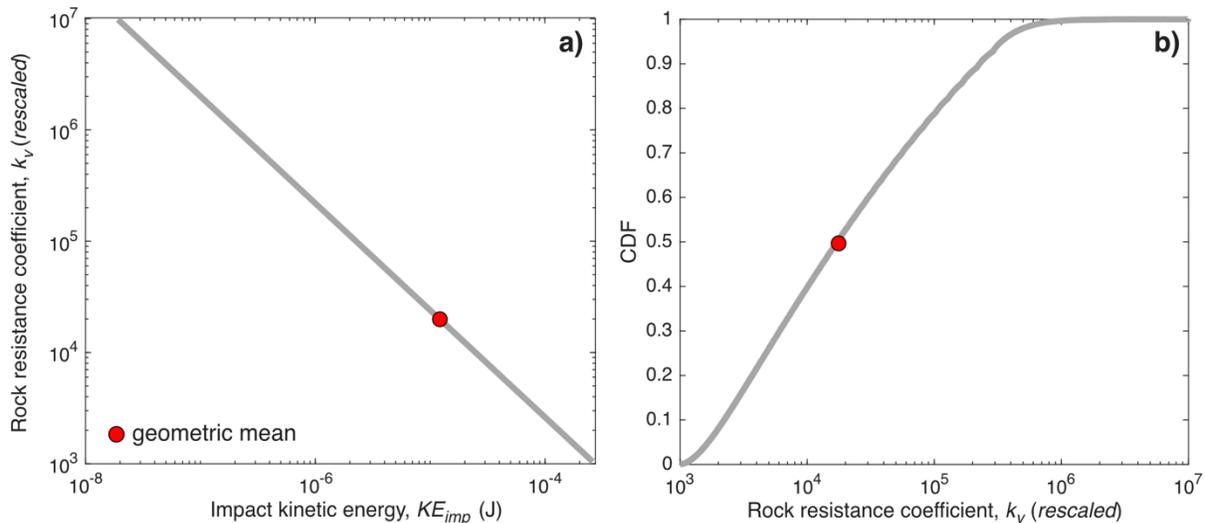
156 **2 Materials and Methods**

157 *2.1 Incorporation of energy-dependent rock resistance coefficient*

158 To evaluate the relative influence of a range of particle impact velocities for bedrock
159 incision, we incorporate an impact energy-dependent rock resistance into the upscaled
160 saltation-abrasion model following Sklar and Dietrich (2006). We define k_v as a function of
161 impact kinetic energy, KE_{imp} , based on the empirical data from Masteller et al. (2024).
162 Masteller et al. (2024) showed from repeated impact experiments that rock-resistance
163 coefficients decrease with increasing impact energy as a greater fraction of the particle impact
164 energy is transferred to bedrock, reducing rock resistance and driving more efficient
165 detachment. To incorporate this observation into the SAM, we use the experimental data from
166 Masteller et al., (2024) to specify the form of the relationship between KE_{imp} and k_v . This is
167 necessary because it is not possible to directly constrain the fraction of energy translated to
168 bedrock during collisions in the current framework of the SAM. Using these data, the impact-
169 energy-dependent rock resistance coefficient can be calculated as $k_v = 839.4KE_{imp}^{-1.21}$.

170 While this empirical scaling provides a translation between impact energy and erosional
 171 resistance as described by k_v , this scaling was derived using brick as a bedrock surrogate rather
 172 than natural rock. Masteller et al., (2024) used bricks in their experiments because they are
 173 relatively homogenous, readily erodible, and consistent with previous impact experiments
 174 (Litwin Miller & Jerolmack, 2021). They also used a steel impactor. Due to the use of these
 175 materials rather than natural rock, experimentally derived rock resistance coefficients ranged
 176 from $k_v = 10^2$ - 10^4 , orders of magnitude less than the commonly used value of $k_v = 10^6$, but
 177 consistent with values of $k_v = 10^3$, for soft rocks (Auel et al., 2017b; Turowski et al., 2013). To
 178 more directly reflect the range of values reported for natural rocks, and to offer a more direct
 179 comparison to Sklar and Dietrich (2006), we rescale the impact-energy dependent rock
 180 resistance coefficient to span $k_v = 10^3$ - 10^7 , based on compiled values from Auel et al. (2017b).

181 To ensure that relative k_v values are comparable across all models explored, we apply a
 182 global rescaling procedure in which impact energies are first sampled over a representative
 183 range of grain sizes, $D_{50} = 0.002$ - 0.1 m, and flow depths, $h = 0.06$ - 4.5 m. The resulting impact
 184 energies span over 5 orders of magnitude and define global minimum and maximum bounds on
 185 k_v (Fig. 1a). For each model run, calculated impact energies are then mapped onto this fixed
 186 global range, ensuring that a given impact energy corresponds to the same k_v across all model
 187 experiments. We compare simulations with impact-energy-dependent, variable k_v to control
 188 runs in which rock resistance is held constant for all impacts. The constant k_v selected in this
 189 case is the geometric mean of the global range of scaled rock resistance values, $k_v = 1.988 \times 10^4$
 190 (Fig. 1B).



191
 192 **Figure 1.** A) Globally scaled range of k_v values as a function of impact kinetic energy, KE_{imp} .
 193 Geometric mean k_v used in constant treatment of rock resistance is denoted as a red dot. B)
 194 CDF plot of globally scaled k_v values with geometric mean again denoted as a red dot.

195
 196 **2.2 Constant discharge case**

197 As a first step, we isolate the influence of impact-energy-dependent k_v under a constant
198 transport capacity, τ^*/τ_c^* , (or discharge, Q) condition. This simplified case allows us to examine
199 how incorporating an impact-energy-dependent rock resistance alters the volume of rock
200 removed per impact, V_{imp} , for a range of flood magnitudes. In the saltation-abrasion model, k_v
201 directly scales the volume removed per bedload impact, with V_{imp} proportional to impact
202 kinetic energy, KE_{imp} .

203 To evaluate how an energy-dependent k_v modifies erosion across a range of hydraulic
204 and sediment conditions, we calculate instantaneous bedrock erosion rates, E_i , for all
205 combinations of $D_{50} = 0.002$ - 0.1 m and $\tau^*/\tau_c^* = 1$ - 10 . The selected grain size range reflects
206 moderate-sized gravel, encompassing the median grain size ($D_{50}=0.06$ m) reported for the
207 South Fork Eel River, the reference site used in Sklar and Dietrich (2006, 2008). The range of
208 selected transport capacities that encapsulates the range of incision rates modeled for the
209 same reference site by Sklar and Dietrich (2006), where the bankfull transport capacity
210 associated with the effective flood for the reference site is $\tau^*/\tau_c^*=1.2$. All other parameters,
211 including reach-scale channel morphology properties and bedrock properties are consistent
212 with Sklar and Dietrich (2006, 2008; Table S1). This constant-discharge analysis provides a
213 baseline for understanding how impact-energy-dependent rock resistance modifies bedrock
214 detachment rates under fixed hydraulic forcing.

215

216 *2.3 Variable discharge case*

217 We extend the constant-discharge framework to incorporate the full magnitude–
218 frequency distribution of discharge, again guided by the South Fork Eel River reference site
219 (Sklar and Dietrich, 2006, 2008). To isolate the role of sediment grain size in mediating the
220 interaction between discharge variability and impact-energy-dependent erosional efficiency, all
221 channel geometry, slope, and hydraulic scaling relationships are held constant (Table S1).

222 We used a 24-year record of mean daily discharge from the South Fork Eel River (USGS
223 gage 11475500, near Branscomb), covering the same period analyzed by Sklar and Dietrich
224 (2006). Discharges were converted to flow depth using empirical power-law relationships (Table
225 S1). Exceedance probabilities for each flood size were then computed as one minus the
226 cumulative distribution function, and the resulting flow-duration curve was resampled and
227 interpolated to smooth noise and increase point density in the distribution tails where
228 observations are sparse. Recurrence intervals for each discharge were calculated using the
229 flow-duration curve method.

230 Bedrock incision was modeled using a modified saltation–abrasion framework following
231 Sklar and Dietrich (2006), comparing a constant rock resistance coefficient with an impact-
232 energy-dependent formulation of k_v . Median grain size was varied from $D_{50} = 0.02$ - 0.1 m,
233 restricted compared to the constant discharge case to better reflect field-measured grain size
234 values along the South Fork Eel River (Masteller et al., under review). Channel slope was fixed

235 at $S = 0.0053$ m/m, consistent with field measurements. For each combination of grain size and
 236 flow depth, we calculated the dimensionless transport capacity (τ^*/τ_c^*) using standard
 237 rectangular-channel formulations. Sediment impact velocity was calculated as a function of
 238 transport capacity, grain size, and fall velocity, and the corresponding impact kinetic energy,
 239 KE_{imp} , was then computed for each flow magnitude. For the impact-energy-dependent model
 240 runs, the rock resistance coefficient, k_v , was then rescaled as a function of KE_{imp} . Bedrock
 241 incision was calculated as the product of the impact rate, I_r , the volume removed per impact,
 242 V_{imp} , and the fraction of the bedrock riverbed exposed to particle impacts, F_e , yielding modeled
 243 incision rates across the full discharge distribution for each grain size scenario.

244 To quantify the cumulative contribution of each flood magnitude to total bedrock
 245 incision, we multiplied the modeled instantaneous incision rates by the exceedance probability
 246 associated with each flow. This probability-weighted incision rate represents the geomorphic
 247 work performed by each discharge, integrating the relative magnitude of incision associated
 248 with each event with the frequency of occurrence for each flood size (e.g. Wolman & Miller,
 249 1960). Framing the results in terms of geomorphic work captures both the sustained
 250 contribution of frequent flows and the relative influence of rare, high-magnitude events. From
 251 this distribution of geomorphic work across the full range of flood flows, we then identified the
 252 effective discharge, defined as the flood magnitude that contributes most to long-term erosion.
 253 We additionally constrained the relative contributions of floods within specified recurrence
 254 intervals to better isolate the influence of less frequent, higher-magnitude events.

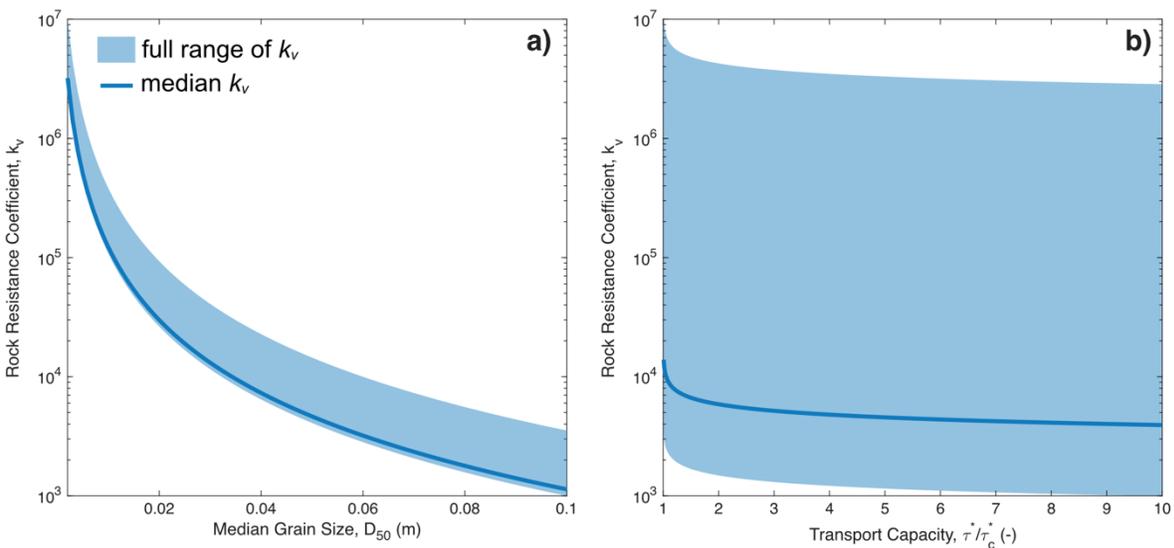
255 **3 Results**

256 *3.1 Constant discharge case*

257 Calculated impact energies span many orders of magnitude, ranging from $KE_{imp} =$
 258 $1.89 \times 10^{-8} - 2.71 \times 10^{-4}$ J. From the saltation-abrasion model, particle impact energies increase
 259 with increasing grain size with the form $KE_{imp} \propto D_{50}^4$, reflecting the combined effects of
 260 increasing particle mass and impact velocity of more massive sediment. This strong nonlinear
 261 scaling between particle size and impact energy produces a heavy-tailed distribution of impact
 262 energies across the parameter space explored (Fig. S1). Because the rock-resistance coefficient
 263 scales inversely with impact energy, the corresponding k_v values decrease sharply with
 264 increasing grain size. Prior to rescaling to values representative of natural materials, k_v exhibits
 265 a strong negative dependence on grain size, with $k_v \propto D_{50}^{-4.84}$. This dependence arises directly
 266 from combining the mechanistic scaling between impact energy and grain size specified by the
 267 SAM with the empirically derived scaling relationship of rock resistance with impact energy
 268 from Masteller et al., (2024). Following rescaling, the exponent of the scaling relationship is
 269 reduced to $k_v \propto D_{50}^{-2.03}$, reflecting some compression of the dynamic range of k_v when
 270 rescaled to values representative of natural materials rather than a change in the underlying
 271 dependence of rock resistance on impact energy.

272 For fixed grain size, the saltation-abrasion model predicts only a weak dependence of k_v
 273 on transport capacity (see Sklar and Dietrich, 2004 for governing equations). In the SAM, impact
 274 velocity scales with excess transport capacity, such that impact energy increases as $KE_{imp} \propto$
 275 $(\tau^*/\tau_c^*)^{0.36}$. Incorporating this relationship into the impact-energy-dependent scaling for rock
 276 resistance implies a comparatively weak sensitivity of k_v to differences transport capacity.
 277 Consistent with this expectation, $k_v \propto (\tau^*/\tau_c^*)^{-0.44}$ for unscaled values and $k_v \propto (\tau^*/\tau_c^*)^{-0.33}$
 278 after global rescaling. Thus, while increasing transport capacity reduces k_v , reducing rock
 279 resistance, the magnitude of this effect is modest compared to the much stronger grain-size
 280 sensitivity observed across the explored D_{50} range.

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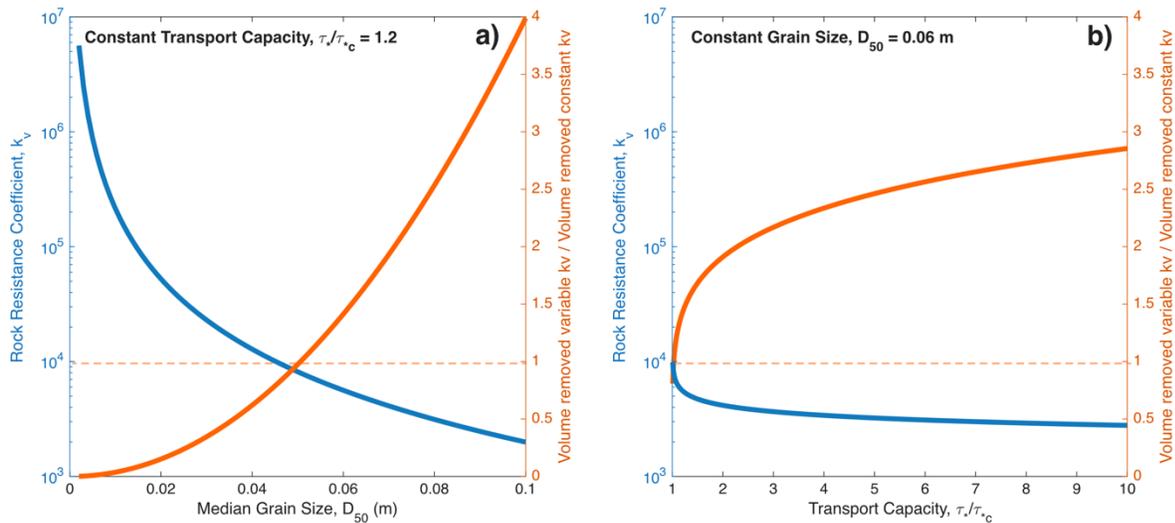
283 **Figure 2.** A) Impact-energy-dependent rock resistance coefficient, k_v , as a function of median
 284 grain size. B) Impact-energy-dependent rock resistance coefficient, k_v , as a function of
 285 transport capacity. Shading highlights range from minimum to maximum k_v , dark blue line
 286 indicates median k_v .

287

288 To place these scaling relationships in context, we evaluate how k_v varies along each
 289 axis of the model parameter space within the constant-discharge experiments (Fig. 2). When
 290 grain size is varied, the entire envelope of k_v values, spanning the full range of modeled
 291 transport capacities, decreases dramatically with increasing sediment caliber (Fig. 2A). For the
 292 finest grains considered, k_v ranges from approximately 10^6 to 10^7 . In contrast, for the coarsest
 293 modeled grain sizes, k_v decreases to roughly 10^3 to 10^4 . Median k_v values consistently trends
 294 toward the lower end of the envelope, indicating that typical impact conditions favor lower
 295 rock resistance and more efficient bedrock detachment. This nearly four-order-of-magnitude
 296 reduction in k_v implies that more massive grains generate greater impact energies and

297 therefore remove bedrock material far more effectively across the entire transport-capacity
 298 range explored.

299 By comparison, across the full span of modeled transport capacities ($\tau^*/\tau_c^* = 1-10$), the
 300 range of rock-resistance coefficients remains consistently wide (Fig. 2B). At any fixed transport
 301 capacity, k_v ranges from approximately 10^3 to 10^7 , reflecting the dominant role of sediment
 302 grain size in setting particle impact energies, and by extension, rock resistance. Again, median
 303 k_v values lie closer to the lower bound of the envelope. Although increasing transport capacity
 304 slightly lowers both the minimum and maximum k_v values, the overall spread in rock resistance
 305 coefficients still spans multiple orders of magnitude. Transport capacity modifies impact
 306 velocity, adjusting impact kinetic energy, but grain size more strongly controls particle mass and
 307 therefore exerts the primary influence on impact energy. As a result, larger grains maintain
 308 substantially lower effective rock-resistance coefficients and drive more efficient bedrock
 309 erosion regardless of hydraulic forcing.



310
 311 **Figure 3.** A) Rock resistance coefficient, k_v (blue), and the ratio of volume removed per impact
 312 using an impact-energy-dependent k_v relative to the constant- k_v case (red), as a function of
 313 median grain size. B) Rock resistance coefficient, k_v (blue), and the ratio of volume removed
 314 per impact using an impact-energy-dependent k_v relative to the constant- k_v case (red), as a
 315 function of transport capacity.

316
 317 These differences in the rock resistance propagate directly to the volume of bedrock
 318 detached per impact, V_{imp} (Fig. 3). At constant transport capacity ($\tau^*/\tau_c^* = 1.2$), increasing grain
 319 size reduces k_v by almost four orders of magnitude (Fig. 3A). This reduction produces up to a
 320 four-fold increase in the volume removed per impact relative to the constant- k_v formulation.
 321 Because $V_{imp} \propto D^4/k_v$, the decline in k_v with increasing sediment caliber compounds the
 322 intrinsic scaling between grain size and impact energy within the SAM. Consequently, the

323 impact-energy-dependent rock resistance predicts progressively larger volumes removed per
324 impact for coarse grains, with departures from the constant- k_v case increasing systematically
325 beyond grain sizes of ~ 5 cm. For grains finer than ~ 5 cm, V_{imp} for an impact-energy-dependent
326 k_v is reduced relative to the constant- k_v case. However, this reduction largely reflects the
327 specific value selected for the constant- k_v reference (Fig. 1), around which the global rescaling
328 is centered. As a result, the ratio of volume removed for impact-energy-dependent rock
329 resistance relative to a constant k_v fall below one at low impact energies. However, the overall
330 pattern remains that increasing grain size enhances the efficiency of bedrock detachment, with
331 the largest grains removing the greatest volume of bedrock per impact relative to a constant
332 rock resistance assumption.

333 At constant grain size ($D_{50} = 0.06$ m), increasing transport capacity reduces k_v by only
334 about a factor of five – substantially smaller than the multiple order-of-magnitude reduction
335 observed with increasing grain size (Fig. 3B). When coupled with the increase in impact velocity
336 at higher τ^*/τ_c^* , this decrease in rock resistance produces up to a three-fold increase in V_{imp}
337 relative to a constant- k_v case. Thus, although hydraulic forcing enhances detachment
338 efficiency, its influence on the volume removed per impact remains secondary to the control
339 exerted by sediment caliber. For a constant grain size, the largest relative increases in V_{imp}
340 occur at transport capacities less than ~ 2 . In this regime, small increases in excess shear stress
341 produce relatively large increases in impact velocity and therefore impact kinetic energy. At
342 higher transport capacities, the response of impact velocity to additional hydraulic forcing
343 diminishes, and the relative enhancement in erosional efficiency and associated increases in
344 V_{imp} become progressively smaller. This pattern indicates that, for a fixed grain size, moderate
345 flows are most effective at amplifying detachment efficiency through increases in impact kinetic
346 energy.

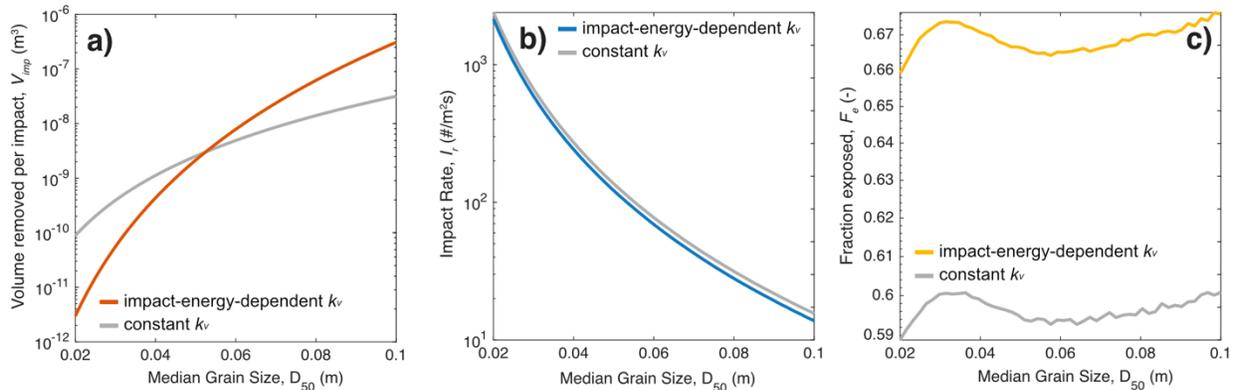
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348 3.2 Variable discharge case

349 3.2.1 The effective flood and components of the saltation-abrasion model

350 For the variable discharge case, we first focus on the effective flood, defined as the
351 discharge that maximizes geomorphic work when magnitude-frequency trade-offs are
352 accounted for across the discharge distribution. We first explore the individual components of
353 the saltation-abrasion model evaluated for the effective flood (Fig. 4).

354



355
 356 **Figure 4.** A) Volume removed per unit impact, V_{imp} (m^3/s), at the effective flood for impact-
 357 energy-dependent rock resistance (red) and constant rock resistance (gray), B) Impact rate, I_r
 358 ($\#/m^2/s$), at the effective flood for impact-energy-dependent rock resistance (blue) and
 359 constant rock resistance (gray), C) Fraction of bedrock exposed, F_e (-), at the effective flood for
 360 impact-energy-dependent rock resistance (yellow) and constant rock resistance (gray).

361

362 The most pronounced differences between the constant and impact-energy-dependent
 363 treatments of rock resistance emerge in the volume removed per impact, V_{imp} . Because
 364 k_v enters directly into the detachment expression, allowing rock resistance to scale with impact
 365 energy dramatically expands the range of V_{imp} at the effective flood relative to the constant- k_v
 366 case (Fig. 4a). In both formulations, V_{imp} increases with grain size, reflecting the strong scaling
 367 between particle mass and impact energy. However, this effect is amplified when the rock
 368 resistance coefficient is framed as a function of impact energy, such that k_v decreases
 369 dramatically when impacts are more energetic. Across the explored grain-size range, V_{imp} at
 370 the effective flood increases by approximately six orders of magnitude under the energy-
 371 dependent formulation of k_v , compared to only 2–3 orders of magnitude when k_v is held
 372 constant. These strong differences in V_{imp} and bedrock detachment efficiency per unit impact
 373 highlight the importance of grain size in determining the effective flood, reflecting the strong
 374 control of particle mass.

375 By comparison, the impact rate, I_r , and fraction of exposed bedrock, F_e , at the effective
 376 flood exhibit much smaller differences between constant and impact-energy-dependent
 377 treatments of rock resistance. Within the saltation-abrasion framework, both I_r and F_e vary as
 378 functions of excess transport capacity, $(1 - \tau^*/\tau_c^*)$, such that their values at the effective flood
 379 reflect the hydraulic transport stage that maximizes geomorphic work. Differences in I_r and F_e
 380 between constant and energy-dependent k_v formulations therefore appear to arise indirectly
 381 from changes in detachment efficiency encapsulated by V_{imp} rather than from independent
 382 modifications to transport mechanics.

383 Impact rate decreases systematically with increasing grain size under both constant and
 384 variable rock resistance (Fig. 4b). This decrease reflects the reduced mobility of coarse

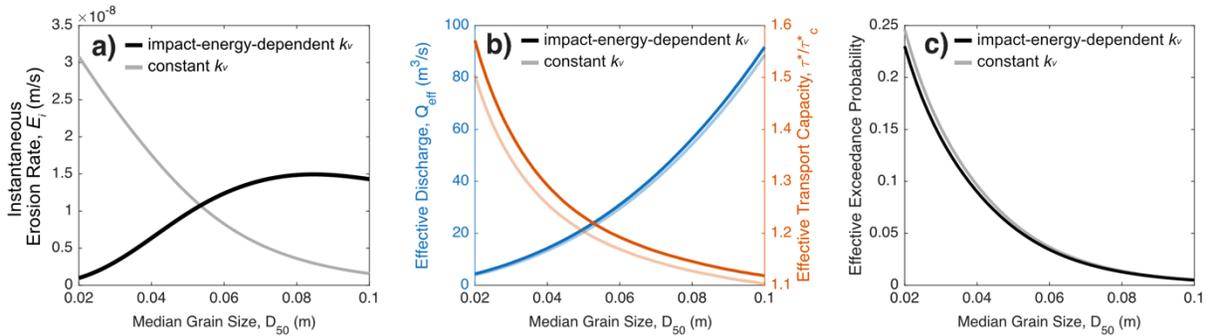
385 sediment across the modeled discharge distribution. As grains become more massive and the
386 volume removed per impact increases, fewer collisions are required to maximize geomorphic
387 work. The decline in I_r also reflects a narrowing of the accessible transport-capacity range for
388 coarse grains, as converting discharge to dimensionless shear stress compresses the effective
389 transport capacity domain with increasing grain size. Incorporation of an impact-energy-
390 dependent rock resistance produces a modest but systematic reduction in impact rate at the
391 effective flood, which remains consistently lower, at approximately 90% of the constant- k_v
392 values, across the modeled grain-size range.

393 The fraction of exposed bedrock, F_e , shows a similarly modest variation across grain
394 sizes. Because F_e reflects the balance between sediment supply and transport capacity, it
395 responds directly to changes in discharge and grain size. At the effective flood, F_e follows a non-
396 monotonic pattern in both constant and energy-dependent resistance formulations (Fig. 4c).
397 The shape of the F_e curves at the effective flood reflects how the transport state that maximizes
398 geomorphic work evolves with grain size. In both constant and energy-dependent treatments of
399 rock resistance, F_e increases from the fine grains to moderate grain sizes, dips slightly at these
400 intermediate sizes, and then rises again toward the coarsest grains.

401 This structure emerges from shifts in the balance among imposed sediment supply
402 (Table S1), declining impact frequency, and increasing detachment per impact. Fine, highly
403 mobile grains generate large transport capacities and moderate exposure. As grain size
404 increases, reduced mobility lowers transport capacity relative to supply, slightly decreasing
405 exposure. At coarse grain sizes, however, the increase in V_{imp} shifts the effective flood toward
406 higher transport capacities where sediment transport more consistently exceeds supply,
407 producing an increase in F_e . The same overall pattern persists in the impact-energy-dependent
408 k_v case, though the curve is systematically elevated relative to the constant- k_v formulation.
409 This offset, an approximately 10-15% increase in bed exposure, reflects a shift of the effective
410 flood toward more intense sediment transport conditions, where impact rates are reduced and
411 bed exposure is enhanced.

412 Nevertheless, the magnitude of variation in both I_r and F_e remains small compared to
413 the orders-of-magnitude divergence observed in V_{imp} . While bedrock riverbed exposure and
414 impact rate adjust in response to the transport state, they do not appear to control the primary
415 differences in the effective flood between constant and impact-energy-dependent treatments
416 of rock resistance. Instead, the dominant control on bedrock incision rates and the effective
417 flood arise from impact-energy-dependent changes in the efficiency of bedrock detachment per
418 unit impact, V_{imp} . These component-level patterns help to clarify how the effective flood
419 adjusts with grain size under each formulation of rock resistance. Impact rate and bed exposure
420 respond to shifts in transport capacity, but their variations remain comparatively small. In
421 contrast, the volume removed per unit impact varies by orders of magnitude as grain size
422 increases, an effect that is strongly enhanced when rock resistance scales with impact energy.

423 The combined effect of these terms ultimately determines the instantaneous erosion rate, E_i at
 424 the effective flood.
 425



426
 427 **Figure 5) A)** Instantaneous bedrock erosion rate, E_i (m/s), at the effective flood for impact-
 428 energy-dependent rock resistance (black) and constant rock resistance (gray), B) Effective
 429 discharge, Q_{eff} (m^3/s ; left axis), and effective transport capacity, τ^*/τ_c^* (right axis) for impact-
 430 energy-dependent rock resistance (dark) and constant rock resistance (light), C) Effective flood
 431 exceedance probability for impact-energy-dependent rock resistance (black) and constant rock
 432 resistance (gray).

433
 434 Instantaneous bedrock erosion rates at the effective flood exhibit fundamentally
 435 different trends depending on whether rock resistance is held constant or allowed to scale with
 436 sediment impact energy (Fig. 5a). Under the constant- k_v formulation, bedrock erosion rates at
 437 the effective flood decrease systematically with increasing grain size. Although larger grains can
 438 detach more material per impact (Fig. 4a), their reduced mobility lowers the impact rate
 439 sufficiently to drive an overall decline in erosion rate. Given this, across the explored range of
 440 D_{50} , effective-flood erosion rates decrease by nearly an order of magnitude under the
 441 assumption of a constant- k_v .

442 In contrast, when k_v scales inversely with increasing impact energy, instantaneous
 443 erosion rates at the effective flood increase with increasing grain size (Fig. 5a). At the coarsest
 444 grain sizes, instantaneous bedrock erosion rates plateau, before slightly declining at the
 445 coarsest grain sizes because the imposed power-law scaling of k_v limits further gains in
 446 erosional efficiency. Importantly, this increase in bedrock erosion rate with increasing grain size
 447 at the effective flood is not driven by large differences in bedload impact rate (Fig. 4B). Rather,
 448 the impact rate associated with the effective flood varies only slightly between the constant
 449 and energy-dependent formulations, exhibiting similar grain-size dependence in both cases.
 450 Indeed, the impact rate is slightly lower under the impact-energy-dependent rock resistance
 451 formulation across much of the grain-size range. The divergence in bedrock erosion rate at the
 452 effective therefore does not arise from increases transport capacity or enhanced particle
 453 impact rates. Instead, it reflects changes in detachment efficiency described by the rock

454 resistance coefficient. These results reinforce that incorporating impact-energy-dependent rock
455 resistance fundamentally alters the grain-size sensitivity of effective-flood incision, with
456 differences arising primarily from the scaling of particle mass and detachment efficiency rather
457 than from changes in transport capacity or impact frequency.

458 Across the full range of grain sizes examined, the discharge associated with the effective
459 flood increases monotonically, spanning $Q_{eff} \approx 4\text{--}95 \text{ m}^3/\text{s}$ (Fig. 4b). This systematic increase
460 reflects the greater flow depths required to mobilize coarser grains and achieve transport
461 stages sufficient for abrasion. In dimensional terms, progressively larger floods dominate long-
462 term incision as grain size increases. However, when expressed as dimensionless transport
463 capacity (τ^*/τ_c^*), the effective flood shifts toward lower transport stages with increasing grain
464 size (Fig. 4b). This apparent paradox arises from the nondimensionalization by grain diameter in
465 the Shields formulation (Shields, 1936), where boundary shear stress is normalized by grain
466 size. For a given discharge, increasing D_{50} reduces τ^* , effectively compressing the accessible
467 range of transport stages across the flow distribution. As a result, although larger discharges are
468 required to mobilize coarser sediment, the maximum in frequency-weighted geomorphic work
469 occurs at more moderate τ^*/τ_c^* values because extreme transport stages become increasingly
470 rare. These dynamics are reflected in systematic shifts in the exceedance probability of the
471 effective flood (Fig. 5c). Effective events become progressively less frequent with increasing
472 grain size, declining from ~ 0.25 at $D_{50} = 0.02 \text{ m}$ to < 0.01 at $D_{50} = 0.10 \text{ m}$. This monotonic
473 decrease reinforces the magnitude-frequency tradeoffs between large and small floods, with
474 coarser sediment shifting the effective flood toward larger, rarer floods capable of moving
475 larger particles.

476 This behavior is consistent across both constant- and impact-energy-dependent
477 formulations of rock resistance. Although the energy-dependent case shifts effective work
478 toward higher discharges at a given grain size (Fig. 5b), the differences in Q_{eff} and associated
479 differences in exceedance probabilities remain very small ($< 10\%$ across all grain sizes). Thus,
480 while individual components of the saltation-abrasion model, specifically the volume removed
481 per unit impact, are strongly influenced by introducing an impact-energy-dependent rock
482 resistance coefficient (Fig. 4), integrating incision across the full flood distribution damps those
483 differences. Weighting incision by flood frequency appears to offset much of the enhanced
484 detachment efficiency at high impact energies, such that the effective flood remains broadly
485 similar between formulations for any given grain size.

486

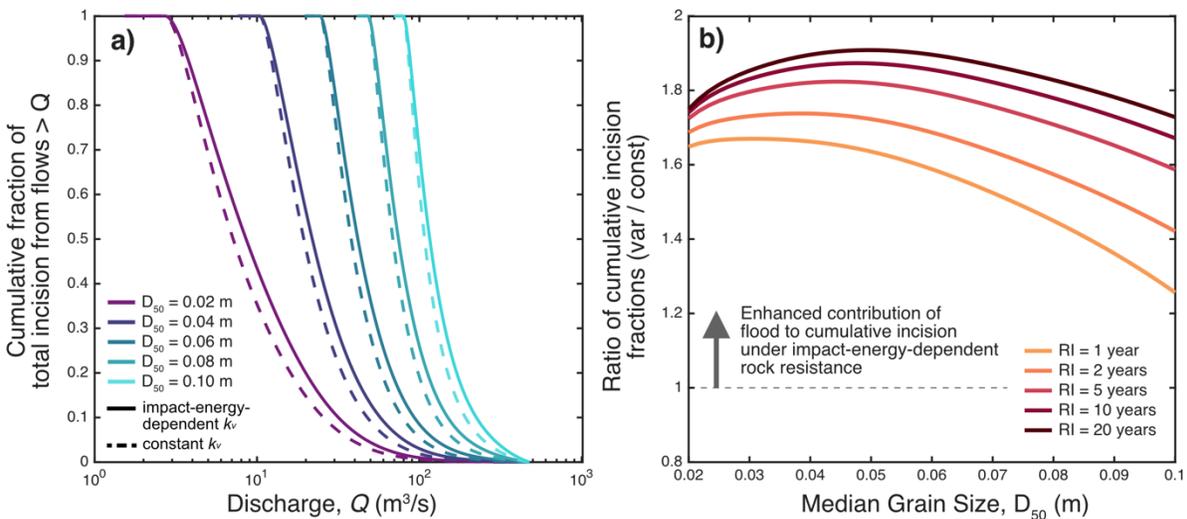
487 *3.2.2 Contribution of less frequent, high-magnitude floods*

488 To directly evaluate the role of more infrequent, high-magnitude events on total
489 bedrock erosion, we examine how cumulative incision is partitioned across the discharge
490 distribution and how this partitioning changes under impact-energy-dependent rock resistance.
491 Although shifts in effective discharge and its exceedance probability remain small across grain

492 sizes (Fig. 5), impact energy increases nonlinearly with particle mass and velocity. As a result,
 493 larger floods are more likely to have enhanced erosional efficiency under impact-energy-
 494 dependent k_v , shifting their overall contribution to cumulative bedrock incision.

495 We first calculate the cumulative fraction of total incision as a function of discharge for a
 496 subset of modeled grain sizes under both constant and impact-energy-dependent rock
 497 resistance (Fig. 6a). For all grain sizes and both treatments of rock resistance, cumulative
 498 incision increases most rapidly across moderate discharges, reflecting the combined influence
 499 of event magnitude and frequency. As grain size increases, the range of discharges that
 500 meaningfully contribute to cumulative incision both shifts to larger flows and narrows,
 501 reflecting a reduction in the range of flows that can transport sediment and erode bedrock via
 502 abrasion. For fine sediment, incision is distributed across a broader range of discharges, and
 503 cumulative incision curves decline gradually as progressively larger floods are included in the
 504 cumulative fraction calculation. In contrast, coarse sediment restricts abrasion to a much
 505 narrower band of high discharges, rendering low-to-moderate flows largely ineffective. Across
 506 all grain sizes, incorporating impact-energy-dependent rock resistance shifts a greater fraction
 507 of total incision toward the upper tail of the discharge distribution, indicating that a larger
 508 portion of total erosion can be attributed to high-magnitude floods. In contrast to the small
 509 shifts in the effective flood (Fig. 5), the cumulative distributions reveal a systematic
 510 redistribution of incision toward the upper tail of the discharge distribution under impact-
 511 energy-dependent rock resistance.

512



513

514 **Figure 6.** A) Cumulative fraction of total bedrock incision as a function of discharge, Q , for a
 515 range of grain sizes for impact-energy-dependent rock resistance (black) and constant rock
 516 resistance (gray). B) Ratio of cumulative incision predicted by the impact-energy-dependent
 517 formulation to that predicted by the constant-resistance case, for a range of recurrence
 518 intervals, as a function of grain size.

519

520 This redistribution is more clearly expressed when cumulative incision is explored at
521 specified flood recurrence intervals (Fig. 6b). For recurrence intervals ranging from 1-20 years,
522 we take the ratio of the cumulative fraction of incision of the variable k_v case relative to the
523 constant k_v case. Across all modeled grain sizes, floods with recurrence intervals exceeding 1
524 year contribute more to total incision under impact-energy-dependent rock resistance
525 compared to the assumption of a constant rock resistance, indicating a systematic increase in
526 the relative importance of larger, less frequent floods. It is worth noting that the magnitude of
527 this increase is not monotonic with sediment caliber. It peaks at intermediate grain sizes, where
528 sediment mobility and impact-energy scaling combine to maximize sensitivity to hydraulic
529 forcing. In this range, floods with 10–20 year recurrence intervals contribute approximately
530 1.8–1.9 times more to total bedrock incision than in the constant- k_v case - an increase of
531 roughly 80–90%. Even more frequent floods with 1–5 year recurrence intervals also show a
532 pronounced contribution to total bedrock incision, exceeding ~60% in this same grain-size
533 window.

534 For finer grains, incision is already distributed across a broad range of discharges, so
535 increasing detachment efficiency at high flows produces a smaller shift in the contribution of
536 high magnitude flows relative to moderate grain sizes. Instead, cumulative erosion remains
537 dominated by smaller, more frequent floods, limiting the relative change in the contribution of
538 rare events. In contrast, for the coarsest grains, incision is already concentrated within high-
539 magnitude floods due to sediment transport constraints. Although impact-energy-dependent
540 rock resistance further increases the incision attributed to these events, the proportional
541 change is smaller than at intermediate grain sizes because the erosion budget is already
542 strongly weighted toward extreme flows. As a result, sensitivity declines at both ends of the
543 grain-size spectrum, producing a pronounced peak in response at intermediate sediment
544 calibers.

545 These results demonstrate that introducing while impact-energy-dependent rock
546 resistance does not substantially alter the effective flood itself, it does redistribute the
547 cumulative incision budget toward less frequent, higher-magnitude events. This redistribution
548 is maximized at intermediate grain sizes where hydraulic forcing and detachment efficiency
549 interact most effectively.

550

551 **4 Discussion**

552 *4.1 Controls on the effective flood under impact-energy-dependent rock resistance*

553 A primary motivation for introducing impact-energy-dependent rock resistance is that
554 particle impact energy increases nonlinearly with both grain size and transport stage (Fig. 2), so
555 erosional efficiency should not be uniform across floods or sediment conditions. Within the
556 saltation-abrasion framework, the effective flood emerges as the discharge that maximizes

557 frequency-weighted geomorphic work across the flow distribution (Wolman & Miller, 1960;
558 Sklar & Dietrich, 2006). Our results show that allowing k_v to decrease with impact energy
559 fundamentally changes the grain-size sensitivity of the detachment component of the saltation-
560 abrasion model (Fig. 3; 4a). However, these shifts in erosional efficiency per unit impact only
561 slightly shift the effective flood itself for any given grain size (Fig. 4b-c). Importantly, this result
562 reveals that while the physics of individual impacts are strongly altered by incorporation of an
563 impact-energy-dependent rock resistance term, the flood event that maximizes geomorphic
564 work remains anchored by the magnitude-frequency structure of the hydrograph.

565 The dominant mechanism by which impact-energy-dependent rock resistance affects
566 bedrock incision is through the volume removed per impact, V_{imp} , because k_v enters the
567 detachment expression directly (Sklar & Dietrich, 2004). Under the impact-energy-dependent
568 formulation, V_{imp} at the effective flood spans a much wider dynamic range with grain size than
569 in the constant- k_v case (Fig. 4a), reflecting the combined scaling of impact energy with particle
570 mass and velocity. By contrast, the impact rate I_r and the exposed fraction F_e vary less between
571 variable and constant- k_v formulations (Fig. 4b–c) because they are primarily governed by
572 transport capacity and sediment cover mechanics in SAM rather than by detachment physics
573 (Sklar & Dietrich, 2004; Sklar & Dietrich, 2006).

574 Despite the strong sensitivity of V_{imp} to impact energy, the effective flood remains
575 broadly similar because the discharge that maximizes long-term incision is determined by the
576 balance between event magnitude and event frequency. Increasing detachment efficiency at
577 high impact energies does not inherently shift dominance to the largest floods, as these events
578 are ultimately de-emphasized by their low exceedance probabilities (Wolman & Miller, 1960;
579 Sklar & Dietrich, 2006). Consequently, the peak in frequency-weighted geomorphic work
580 reflects intermediate-sized flows for both variable and constant- k_v . This behavior is consistent
581 with prior applications of the saltation-abrasion model, which show that incision rates are
582 maximized where transport stage is sufficiently high to generate frequent energetic impacts,
583 yet not so rare that cumulative contribution is limited by infrequent occurrence or by evolving
584 sediment cover dynamics (Sklar & Dietrich, 2006; Sklar & Dietrich, 2008).

585 Importantly, this partitioning assumes that sediment supply and cover respond
586 deterministically to hydraulics, as prescribed in the canonical SAM framework. In natural
587 channels, however, stochastic sediment inputs from hillslopes, episodic pulses of coarse
588 material, and transient storage can partially decouple F_e from instantaneous hydraulic forcing
589 (DeLisle et al., 2022; DeLisle & Yanites, 2023; Golly et al., 2017; Turowski, 2009). Variability in
590 supply relative to transport capacity can shift the bed between detachment- and transport-
591 limited states, altering both the threshold for erosion and the channel's ability to convey
592 sediment (Finnegan et al., 2007; Johnson, 2016; Johnson & Whipple, 2010). Under conditions of
593 excess supply, persistent or patchy cover may suppress exposure even at high shear stress,
594 limiting the delivery of high-energy impacts. Conversely, sediment-starved conditions may

595 increase exposure but reduce impact frequency, lowering the cumulative work delivered to the
596 bed. Because impact energy and V_{imp} scale strongly with transport stage, such stochastic
597 modulation of cover and transport efficiency has the potential to reorganize when and how
598 impact energy is transmitted to the bed, potentially shifting erosion thresholds and
599 redistributing geomorphic work across the hydrograph (DeLisle & Yanites, 2023; Turowski,
600 2009). Thus, while our results isolate the detachment consequences of energy-dependent
601 resistance, natural variability in sediment supply introduces an additional degree of freedom
602 that could amplify or damp the predicted grain-size sensitivity on bedrock incision rates.

603 *4.2 The shifting contribution of less frequent, high-magnitude floods*

604 Notably, our analysis of less frequent, higher-magnitude floods demonstrates that while
605 the effective flood appears to remain relatively constant, this does not imply invariance in the
606 partitioning of overall incision across all flood events. Instead, our recurrence-interval analysis
607 reveals a systematic redistribution of cumulative incision toward the upper tail of the discharge
608 distribution under impact-energy-dependent rock resistance (Fig. 6). This result reinforces that
609 while the effective flood represents a summary of the geomorphic work, identifying the
610 discharge at which frequency-weighted incision is maximized, it does not necessarily describe
611 how incision is distributed across the rest of the discharge distribution. The recurrence-interval
612 incision portioning analysis that we present here, by contrast, integrates incision across
613 portions of the flow distribution and revealing changes in the overall shape of that distribution.
614 Our analysis reveals that under impact-energy-dependent resistance, a greater fraction of total
615 incision is accomplished by less frequent, higher-magnitude events.

616 This behavior parallels results from variable-discharge stream-power models in which
617 erosion thresholds determine which portions of the hydrograph contribute to long-term
618 incision (Tucker & Bras, 2000; Snyder et al., 2003; Lague et al., 2005; DiBiase & Whipple, 2011).
619 In those frameworks, increasing the incision threshold may not shift the discharge that
620 maximizes geomorphic work, yet it alters the relative contribution of the distribution tails by
621 concentrating incision into flows that exceed that threshold. Here, impact-energy-dependent
622 resistance operates through a different mechanism, enhancing the efficiency of bedrock
623 detachment rather than directly modifying entrainment threshold, but produces a similar
624 outcome by increasing the erosive effectiveness of high-energy impacts during larger floods.

625 These findings have broader implications for landscape-scale modeling. First, they
626 provide a process-based pathway for explicitly incorporating bedload impact energy, and by
627 extension, sediment-transport driven differences in bedrock detachment efficiency, into
628 models of long-term incision. In SPIM-like formulations, thresholds are often implemented
629 phenomenologically by introducing a critical shear stress or modifying erodibility without
630 explicitly specifying what sets this threshold or how it evolves with sediment transport
631 conditions. In contrast, the saltation–abrasion framework defines the relevant erosion
632 threshold mechanistically as the boundary shear stress required to mobilize sediment, and ties

633 incision directly to sediment entrainment, impact delivery, and bed cover. Incorporation of an
634 impact-energy-dependent resistance extends this framework by demonstrating that even
635 when the entrainment threshold is fixed, the efficiency of impacts above this threshold can vary
636 strongly with grain size and discharge. In this treatment, bedrock erodibility becomes an
637 emergent property of coupling hydraulics and sediment transport.

638 Importantly, because the erosion threshold is explicitly linked to sediment mobilization,
639 this framework also allows variability in grain-scale motion thresholds to be incorporated
640 directly into long-term incision models. Critical Shields stress is known to vary with channel
641 slope (Lamb et al., 2008; Prancevic & Lamb, 2015), flood magnitude and intermittency
642 (Masteller et al., 2019; Masteller et al., 2025; Ockelford et al., 2019), reworking of riverbed
643 structure (Hodge et al., 2013; Masteller & Finnegan, 2017; Ockelford & Haynes, 2013; Whitfield
644 et al., 2025), and shifts in sediment supply (Hassan et al., 2020; Johnson, 2016; Yager et al.,
645 2012), all of which influence sediment mobility. These process-driven variations in mobility
646 thresholds influence both the frequency and intensity of particle impacts during floods of a
647 given size. By explicitly coupling these threshold dynamics to impact-energy-dependent
648 resistance, the saltation–abrasion framework provides a mechanistic route for linking
649 hydrologic variability, sediment caliber, and bedrock erosion efficiency across the full flood
650 distribution.

651 Our results reinforce that the effective flood alone is not sufficient to capture the
652 relative contribution of rare, large magnitude flows on overall bedrock incision rates. We find
653 that the effective discharge only shifts slightly between constant and impact-energy-dependent
654 formulations of rock resistance. However, exploration of low-recurrence interval floods reveals
655 a redistribution of cumulative incision toward rarer floods, particularly at intermediate grain
656 sizes. The effective flood metric obscures this shift because it does not capture changes in the
657 skewness or tail weighting of incision across the hydrograph. This distinction is critical for
658 interpreting landscape response to hydroclimatic shifts that disproportionately alter the tails of
659 flood distributions. Even if the nominal effective discharge remains stable, increases in flood
660 intermittency or intensity may elevate the geomorphic influence of rare events. Likewise, shifts
661 in sediment caliber or mobility may also alter the sensitivity of bedrock incision to extreme
662 flows. These results highlight the need for long-term incision models to further explore both
663 grain-scale erosion thresholds and impact-energy-dependent shifts in erosional efficiency, and
664 to account for how these processes redistribute geomorphic work across the full flood
665 distribution.

666

667 **5 Conclusions and future work**

668 In this contribution, we introduced an impact-energy-dependent rock resistance
669 coefficient into the saltation-abrasion model, which calculates bedrock erosion rates as a
670 function of sediment transport rates. We found that incorporating this variable rock resistance

671 alters bedrock incision predictions primarily by changing efficiency of bedrock detachment per
672 unit impact at the grain scale. Under this formulation, rock resistance decreases as impact
673 energies increase, amplifying the grain-size dependence of the volume removed per impact. At
674 a constant-discharge experiments, this produces a strong reduction in rock resistance with
675 increasing grain size and a corresponding systematic increase in the volume of material
676 detached per impact relative to a constant resistance assumption.

677 When the full discharge distribution is considered, these component-level shifts
678 translate into two reach-scale outcomes. First, the effective flood remains broadly similar for a
679 given grain size, indicating that the discharge that maximizes frequency-weighted geomorphic
680 work is still largely set by magnitude-frequency tradeoffs within the discharge distribution.
681 Second, despite minimal changes in the effective flood, the cumulative incision budget shifts
682 toward less frequent, higher-magnitude events under impact-energy-dependent rock
683 resistance, increasing the relative contribution of the upper tail of the flood distribution. This
684 effect is maximized for intermediate grain sizes where impacts become substantially more
685 energetic while sediment remains sufficiently mobile to sustain abrasion.

686 Two limitations are worth emphasizing when interpreting these results into long-term
687 models for bedrock incision or in field application. First, the impact-energy-dependent scaling
688 of the rock resistance coefficient, is experimentally constrained but derived using surrogate
689 materials and a specific impactor configuration, and then rescaled to span a range reflecting
690 natural rocks (Auel et al., 2017b; Turowski et al., 2013; Masteller et al., 2024). While the relative
691 trends are mechanistically-motivated, extending the framework across lithologies and
692 transport-driven distributions in impact trajectories will require additional calibration and
693 testing. Second, the modeling isolates grain size while holding channel geometry and slope
694 fixed to attribute differences in incision partitioning to the mechanics of bedrock detachment.
695 In natural channels, changes in sediment caliber and discharge regime have strong feedbacks
696 with channel width, slope, and bed state, which in turn alter transport stage, cover, and the
697 reach's capacity to translate impacts into incision (Finnegan et al., 2017; Gasparini et al., 2007;
698 Yanites, 2018) (Gasparini et al., 2007; Yanites, 2018). Incorporating these feedbacks is a natural
699 next step for assessing how impact-energy-dependent resistance influences long-profile form
700 and transient channel adjustment.

701 Motivated by these results, future experimental work should aim to better constrain the
702 relationship between particle impact energy and rock resistance across representative natural
703 lithologies, sediment shapes, and mixed-size populations. Flume experiments should also
704 quantify how impact angle, impactor material, and bed condition alter the fraction of kinetic
705 energy transmitted to fracture growth. To better characterize the consequences of these
706 impact-energy-dependent differences in rock resistance on channel form, modeling should
707 could explore how channel width, slope, and bed cover to evolve with variable sediment supply
708 and discharge regime. By coupling grain-scale processes with evolution of channel form, this

709 work could evaluate whether the observed redistribution of incision toward rarer floods has an
710 extractable signature on equilibrium channel morphology.

711

712 **Acknowledgments**

713

714 **Open Research**

715 All code developed as part of this study is available in a Zenodo repository (Masteller,
716 2026). *Note to editors and reviewers: The Zenodo repository is currently accessible as a draft*
717 *with the sharing link below, however, the data are not yet formally published with a DOI. The*
718 *formally published data will be cited here and linked with a DOI following review. This delay is to*
719 *enable edits to the dataset if substantive methodological changes are suggested during the*
720 *review process resulting in material changes to the data.*

721 <https://zenodo.org/records/18851708?preview=1&token=eyJhbGciOiJIUzUxMiJ9.eyJpZCI6IjdhNzViNzY0LWM0M0MwQTNNGNlOS1iOGUwLTg0ZTVkYzQyNzQ1ZCIsImRhdGEiOiJ9LCJyYW5kb20iOiJzYjRkYzE3ZDE5ZGM2Y2IxZTUxM2M5NjZmZDVjNjY4YiJ9.CLBZ4WzEV5f5cKVL0A1QLN3NS1Q-LvmZjhOGIPVFC7W0IbHilqJi3skbcvo0PSSlqjcpUnBhSbY0pMGeKj0Bfw>

725

726 **Conflict of Interest Disclosure**

727 The authors declare there are no conflicts of interest for this manuscript.

728

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Supporting Information for

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Implications of impact-energy dependent erosional efficiency on bedrock river sediment dynamics and form: 1. Reach-scale dynamics and the effective flood

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Contents of this file

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Table S1

Figure S1

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Introduction

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- This supplementary file contains one table defining the model parameter values used in this study and their values (Table S1). It also contains one figure, Figure S1, demonstrating the distribution of impact kinetic energies for the constant discharge case explored in the main manuscript.

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941 **Table S1.**942 List of model parameter values for reference field site, reference values from Sklar and Dietrich,
943 2006.

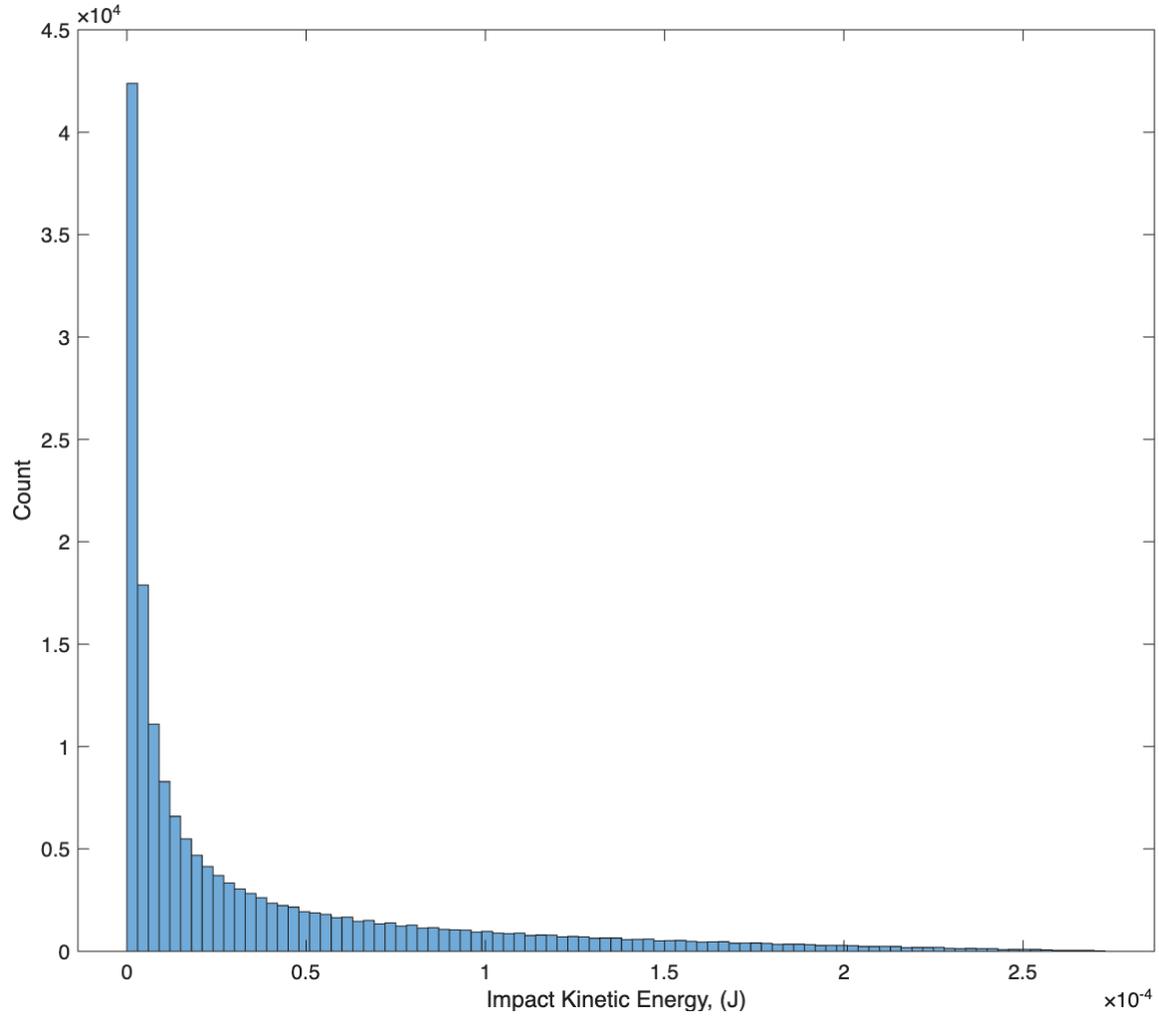
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Parameter name	Parameter symbol	Units	Value
Drainage area	A	km ²	112
Slope	S	m/m	0.005
Discharge-drainage area coefficient	c		0.347
Discharge-drainage area exponent	h		1
Channel width	w	m	18
Depth-discharge scaling coefficient	k		0.3147
Depth-discharge scaling exponent	f		0.4321
acceleration due to gravity	g	m/s ²	9.81
Density of water	ρ	kg/m ³	1000
density of sediment	ρ_s	kg/m ³	2650
Submerged sediment density	R_b		1.65
Dynamic viscosity	μ	Pa s	0.00
critical Shields stress	τ_c^*		0.03
Rock tensile strength	σ_T	Pa	7000000.00
Young's modulus	Y	Pa	50000000000.00
Sediment supply rate	Q_s	kg/s	42.60
Median rock resistance coefficient	k_v		19880.00
Minimum rock resistance coefficient	$k_{v\min}$		1000.00
Maximum rock resistance coefficient	$k_{v\max}$		10000000.00
Median grain size, reference site	D_{50}	m	0.06

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949 **Figure S1.** Distribution of Impact Kinetic Energy, KE_{imp} (J), for the constant discharge case

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