

Geostatistical Earth Modeling of Cyclic Depositional Facies and Diagenesis

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July 2, 2018

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13 **Acknowledgements**

14 The authors would like to thank Carl Jacquemyn and Claire Veillard for providing
15 practical examples and contributing ideas to this paper, and the department of
16 Earth Science and Engineering at Imperial College for a scholarship to Thomas Le
17 Blévec. Olivier Dubrule would like to thank Total for seconding him at Imperial
18 College London.

19 **Abstract**

20 In siliciclastic and carbonate reservoirs, depositional facies are often described as be-
21 ing organized in cyclic successions that are overprinted by diagenesis. Most reservoir
22 modeling workflows are not able to reproduce stochastically such patterns. Herein,
23 a novel geostatistical method is developed to model depositional facies architectures
24 that are rhythmic and cyclic, together with superimposed diagenetic facies.

25 The method uses truncated Pluri-Gaussian random functions constrained by
26 transiograms. Cyclicity is defined as an asymmetric ordering between facies, and
27 its direction is given by a three-dimensional vector, called shift. This method is
28 illustrated on two case studies. Outcrop data of the Triassic Latemar carbonate
29 platform, northern Italy, are used to model shallowing-upward facies cycles in the
30 vertical direction. A satellite image of the modern Bermuda platform interior is used
31 to model facies cycles in the windward-to-leeward lateral direction.

32 As depositional facies architectures are modelled using two Gaussian random
33 functions, a third Gaussian random function is added to model diagenesis. Thereby,
34 depositional and diagenetic facies can exhibit spatial asymmetric relations. The
35 method is applied in two regions of the Latemar carbonate platform that expe-
36 rience two different types of diagenetic transformation: syn-depositional dolomite
37 formation, and post-depositional fracture-related diagenesis. The method can also
38 incorporate proportion curves to model non-stationary facies proportions. This is
39 illustrated in Cretaceous shallow-marine sandstones and mudstones, Book Cliffs,
40 Utah, for which cyclic facies and diagenetic patterns are constrained by embedded
41 transition probabilities.

42 **Introduction**

43 In reservoir modeling applications, an important step is the representation of three-
44 dimensional facies architecture and the quantification of associated uncertainty. The
45 geomodeling community routinely uses geostatistical methods to reach this goal

46 (Koltermann and Gorelick, 1996; Alabert and Modot, 1992; Pyrcz and Deutsch,
47 2014). However, the commonly-used geostatistical approaches have some signif-
48 icant limitations. For instance, geostatistical models often show the same facies
49 successions in the upward as in the downward direction, which does not allow the
50 representation of classic geological features such as facies cyclicity or certain types
51 of syn-depositional diagenesis.

52 **Modeling Cyclicity and Rhythmicity**

53 Depositional facies in vertical successions exhibit extensive cyclicity and rhythmic-
54 ity (Strasser, 1988; Goldhammer et al., 1990; Wilkinson et al., 1997; Lindsay et al.,
55 2006; Burgess, 2016). These features are defined respectively as facies ordering (Gin-
56 gerich, 1969; Hattori, 1976) and repetition of facies at intervals of constant thickness
57 (De Boer and Wonders, 1984; House, 1985). Their origin is attributed to various
58 controls, including relative sea level oscillations (e.g., Grotzinger, 1986), local tec-
59 tonic activity (e.g., Cisne, 1986) and autogenic mechanisms. These different origins
60 may lead to cycles and rhythms of differing lateral extent and stacking patterns,
61 which should be reproduced by the modeling method. For example, facies cycles
62 are commonly interpreted at reservoir scale with reference to sequence stratigraphic
63 models, implying that they are laterally extensive (e.g., Goldhammer et al., 1990),
64 although over such distances, some facies cycles are documented to pinch out (e.g.,
65 Egenhoff et al., 1999). In order to represent these diverse facies cycles and rhythms,
66 reservoir-wide deterministic correlations may not be appropriate.

67 Diverse facies distributions are modelled by geostatistical methods, but their
68 current implementation cannot generate facies cycles and rhythms simultaneously.
69 For example, cyclicity quantification is possible with Markov Chain analysis (Gin-
70 gerich, 1969; Hattori, 1976), but the method is originally limited to one dimension.
71 It was later improved by Carle and Fogg (1996) who model cyclic three dimensional
72 Earth models thanks to asymmetric transiograms. However, the method does not
73 incorporate rhythmicity, because the transiogram models are not flexible enough to

74 incorporate the characteristic periodic oscillations (Jones and Ma, 2001; Dubrule,
75 2017), called hole-effects. Facies cyclicity and rhythmicity could theoretically be
76 modelled by Multi-Point Statistics (Strebelle, 2002), but it is challenging to include
77 those patterns in the required three dimensional training image.

78 A geostatistical method has been developed recently to model simultaneously
79 facies cyclicity and rhythmicity (Le Blévec et al., 2018), thus improving the real-
80 ism of facies Earth models. The method is based on Pluri-Gaussian Simulations
81 (Armstrong et al., 2011), constrained by facies transiograms. The facies asymmetric
82 ordering (or cyclicity) is defined by two Gaussian random functions spatially shifted
83 from each other, and rhythmicity of the facies succession is modelled by defining
84 new hole-effect covariance models (Le Blévec et al., 2018).

85 So far, this method has only been used to model cyclicity in the vertical direction,
86 although cyclicity can also be observed in lateral directions. Stratigraphic forward
87 models can produce asymmetry between facies in lateral directions (Burgess et al.,
88 2001), and such lateral facies asymmetry is also explicit within Walther’s Law (e.g.,
89 Middleton, 1973). This could possibly be modelled with the shifted Pluri-Gaussian
90 method (Le Blévec et al., 2018) by defining a spatial shift with a lateral component.

91 **Modeling Diagenesis**

92 Reservoir quality is not only affected by depositional facies cyclicity. Rock proper-
93 ties of carbonate (e.g., Bartok et al., 1981; Moore and Wade, 2013) and siliciclastic
94 (e.g., Taylor et al., 2010) deposits are also influenced by diagenesis. Diagenetic pro-
95 cesses give rise to depositional and diagenetic facies with a variety of geometrical
96 relationships, which should be captured by the modeling method. Early diagenesis
97 tends to closely follow the texture and stratal configuration of depositional facies
98 (e.g., Ginsburg, 1957; Egenhoff et al., 1999; Peterhänsel and Egenhoff, 2008; Rameil,
99 2008) while late diagenesis either follow depositional features, or other structures
100 such as faults, fractures, and karsts, thus resulting in diagenetic bodies and trends
101 that cut across depositional facies geometries (e.g., Sharp et al., 2010; Vandeginste

102 [et al., 2013](#); [Jacquemyn et al., 2014](#); [Beckert et al., 2015](#)). It is therefore highly desir-
103 able that reservoir modeling methods are flexible enough to embrace these different
104 possibilities.

105 In many geostatistical studies, diagenesis is modelled as porosity or permeability
106 variations ([Wang et al., 1998](#); [Pontiggia et al., 2010](#)). This is a useful approach,
107 but it cannot be applied to the representation of distinct diagenetic geobodies or
108 of different diagenetic phases within a depositional facies. Therefore, some authors
109 model diagenesis as a facies random field that is superimposed on the depositional
110 facies field ([Renard et al., 2008](#); [Doligez et al., 2011](#); [Barbier et al., 2012](#); [Carrera](#)
111 [et al., 2018](#)). These authors use a version of truncated Pluri-Gaussian Simulations
112 (Bi-PGS) developed by [Renard et al. \(2008\)](#), which models two facies fields with dif-
113 ferent Gaussian random functions. The depositional and diagenetic facies fields can
114 be either independent of or correlated to each other, which allows to model depo-
115 sitional and diagenetic facies geometries that are either discordant or conformable.
116 However, this method does not generate distributions of diagenetic facies that are
117 asymmetric such as occurring preferentially towards the top or the base of deposi-
118 tional facies bodies.

119 The algorithm of [Renard et al. \(2008\)](#) to model diagenesis is thus extended
120 here by including a shift between depositional and diagenetic facies fields, which
121 allows diagenetic facies to overprint depositional facies preferentially at their top or
122 at their base. These relationships are constrained by cross-transiograms between
123 the two facies fields, and the method is also combined with the advancements of
124 [Le Blévec et al. \(2018\)](#), so that diagenesis is modelled in the context of depositional
125 facies cyclicity and rhythmicity.

126 **Aims**

127 Therefore, this paper presents a new geostatistical facies modeling method that is
128 able to represent facies cyclicity and rhythmicity, together with diagenetic facies
129 bodies, which have either conformable or non-conformable geometries. The struc-

130 ture of the paper is outlined below. First, we illustrate the concepts of cyclicality and
131 rhythmicity and show that these concepts are captured by transiograms. Model-
132 ing of cycles and rhythms is then illustrated using: (1) synthetic facies successions;
133 (2) facies successions from the outcropping Triassic Latemar carbonate platform
134 of Northern Italy; and (3) lateral facies relationships on the interior of the mod-
135 ern Bermuda carbonate platform. Then, diagenesis is modelled by adding another
136 Gaussian random function to the method. Three examples are modelled to illustrate
137 the flexibility of the method: (1) syn-depositional diagenesis below hardgrounds in
138 facies cycles of the Latemar carbonate platform; (2) early diagenetic development
139 of concretions in shallow-marine, siliciclastic facies cycles in the Cretaceous Black-
140 hawk Formation, Book Cliffs outcrops (Utah), in which the facies proportions are
141 non-stationary; and (3) post-depositional diagenesis caused by localised movement
142 of hydrothermal fluids through faults and fractures in the Latemar carbonate plat-
143 form.

144 **Quantifying Cyclicality and Rhythmicity with Tran-** 145 **siograms**

146 **Defining Cyclicality and Rhythmicity**

147 Although facies cyclicality and rhythmicity are commonly interpreted in sedimentary
148 sequences, these concepts have different meanings to different authors. Formal,
149 quantitative definitions of cyclicality and rhythmicity are needed for facies modeling,
150 as a facies succession can be more or less ordered or exhibit more or less variability in
151 facies thickness. Cyclicality is defined as facies ordering in a given direction ([Gingerich,](#)
152 [1969](#); [Hattori, 1976](#); [Le Blévec et al., 2018](#)), usually vertically (Fig. 1). The ordering
153 considered here is asymmetric, which means that it differs going upwards from going
154 downwards. For instance, in vertical shallow-marine carbonate and siliciclastic suc-
155 cessions, facies cycles tend to be shallowing-upward ([Strasser, 1988](#); [Goldhammer](#)

156 et al., 1990; Lindsay et al., 2006), which is equivalent to deepening-downward.

157 Another commonly observed feature is that the same facies tends to appear re-
158 peatedly at intervals of constant thickness (e.g., Goldhammer et al., 1993; Lindsay
159 et al., 2006), which defines rhythmicity (De Boer and Wonders, 1984; House, 1985;
160 Le Blévec et al., 2018). If cyclicity and rhythmicity are both present, it implies
161 that the facies cycles have low variability in thickness. For instance, the vertical
162 succession in Figure 1d is cyclic and rhythmic because the facies are fully ordered
163 and have constant thickness intervals between them. The succession illustrated in
164 Figure 1a has non-ordered transitions between facies and also contains two facies
165 cycles. The succession in Figure 1b also contains two facies cycles, and the blue
166 facies is rhythmic, because intervals between occurrences of this facies have similar
167 thickness. Figure 1c shows a cyclic and non-rhythmic facies succession, and the suc-
168 cession in Figure 1d is cyclic and rhythmic, because the facies are fully ordered and
169 the blue facies is separated by intervals of constant thickness. For three-dimensional
170 Earth models to be geologically realistic, facies cyclicity and rhythmicity must be
171 properly modelled.

172 **The Transiogram: a Tool to Capture Cyclicity and Rhythmic-** 173 **ity**

174 Standard geostatistical simulation approaches quantify geologic patterns by comput-
175 ing experimental variograms, modeling them mathematically and then ensuring that
176 the variogram models are reproduced in the final simulation (Pyrcz and Deutsch,
177 2014). However, Carle and Fogg (1996) show that variograms are not able to quan-
178 tify asymmetric cycles, and promote the use of the transiogram instead.

179 The transiogram gives the probability $t_{AB}(h)$ of finding a facies B at a vector
180 h from a given facies A (Carle and Fogg, 1996; Le Blévec et al., 2018). If the two
181 facies A and B are identical, the transiogram is referred to as an auto-transiogram,
182 otherwise it is referred to as a cross-transiogram. Auto-transiograms and cross-
183 transiograms are calculated experimentally and gathered in a transiogram matrix

184 (Fig. 2). As with variograms, the direction h is usually vertical, but it can also
185 have a lateral component if calculated along other directions. For sedimentary fa-
186 cies, transiograms are commonly different in opposite directions (e.g., upward and
187 downward) (Carle and Fogg, 1996).

188 Transiograms have specific properties, which are described in detail by Carle and
189 Fogg (1996). One property is that at long distances h , $t_{AB}(h)$ tends towards the
190 proportion of facies B . For example, Figure 2b-e shows that the transiograms tend
191 towards the value of 0.5, which is the proportion of facies 1, and 0.25, which is the
192 proportion of facies 2. Also, the tangent at the origin $t'_{AA}(0)$ of the auto-transiogram
193 $t_{AA}(h)$ defines the mean length of facies A , denoted as \overline{L}_A (Carle and Fogg, 1996),
194 as shown in Figure 2b, e (\overline{L}_1 and \overline{L}_2).

195 Figure 2c, and 2d also show that cyclicity is captured by the behavior at the
196 origin of the cross-transiograms (Le Blévec et al., 2018). $t'_{12}(0)$ is high while $t'_{21}(0)$ is
197 low, which means that facies 2 tends to overlie facies 1, while facies 1 does not tend
198 to overlie facies 2. Consequently facies 3 overlies facies 2, creating facies cycles with
199 facies 1 at the base, facies 2 in the center and facies 3 at the top. This cyclicity is
200 observed in the corresponding succession (Fig. 2a), which shows that facies 1 almost
201 always transitions upward to facies 2 (except on one occasion when it transitions
202 directly to facies 3), and facies 2 transitions upward to facies 3.

203 Rhythmicity is characterized by the oscillations of the transiograms or variograms
204 (Jones and Ma, 2001; Le Blévec et al., 2018), as shown in Figure 2. The average
205 distance separating two repetitions of a facies is given by the first maximum of the
206 corresponding auto-transiogram, as this is associated with the highest probability
207 of finding the same facies (Fig. 2b, e, $\overline{L}_c = 0.4$ m). It also corresponds to the
208 first minimum of the cross-transiograms (Fig. 2c, d), which is associated with the
209 lowest probability of finding two different facies. In this case, because there is also
210 cyclicity, this length \overline{L}_c corresponds to the average thickness of the facies cycle
211 and is approximately the sum of the mean thicknesses of all facies present in a
212 cycle. This also explains why the auto-transiogram of facies 2 shows the same

213 rhythmicity (Fig. 2e) as the auto-transiogram of facies 1 (Fig. 2b). Rhythmicity can
214 be visually verified by examining the corresponding succession (Fig. 2a), which shows
215 that facies cycles indeed exhibit low thickness variations (thickness of approximately
216 0.4 m). Therefore, transiograms appear to be better suited than variograms to the
217 quantification of cyclicity and rhythmicity.

218 Modeling Cyclicity and Rhythmicity with Shifted Pluri- 219 Gaussian Simulations

220 Principle of Truncated Gaussian Simulations

221 The truncated Gaussian approach for facies modeling was first developed by [Math-](#)
222 [eron et al. \(1988\)](#) and is explained in detail by [Armstrong et al. \(2011\)](#). It has
223 two steps: (1) first, the simulation of a continuous Gaussian random function, and
224 then (2) the truncation of this continuous function into facies with the help of a
225 truncation rule.

226 A Gaussian random function defines at every location (x, y, z) (usually in a
227 grid) a Gaussian random variable. The Gaussian random function is controlled by
228 its covariance model ([Chiles and Delfiner, 2012](#)). In this paper, as explained in
229 the Appendix, Gaussian cosine covariances (with frequency parameter b) are used
230 with scale factor noted r_z (Eq. A.4a) in the vertical direction. In lateral directions
231 Gaussian covariances are used, with scale factors noted r_x and r_y for each principal
232 direction (Eq. A.4a). Scale factors control the average length scale of the Gaussian
233 random functions in the corresponding direction and b their periodicity ([Le Blévec](#)
234 [et al., 2018](#)). Figure 3b (red curve) shows an example of a realization of a Gaussian
235 random function Z_1 along a vertical succession (i.e., on a one-dimensional grid).

236 The truncation rule defines the number of facies, their proportions, and their
237 contacts. For instance, Figure 3a shows a truncation rule with three facies, with
238 a small area for facies 3 defined by the threshold q_2 . This results in a smaller

239 proportion of facies 3 in the corresponding vertical succession (Fig. 3b). As shown
240 by Figure 3b, when Z_1 is higher than q_1 , facies 2 is allocated, and when it reaches
241 q_2 , facies 3 is defined.

242 By using only one Gaussian random function, modeling is limited because each
243 facies can only transition into one or, at most two other facies. For instance, in
244 Figure 3 facies 1 and 3 can only transition into facies 2, while facies 2 can transition
245 into both facies 1 and 2 upward or downward. Therefore, cyclicity cannot be mod-
246 elled because there is no asymmetry associated to the simulation. [Armstrong et al.](#)
247 (2011) extend the method to Pluri-Gaussian Simulations, and it was then modified
248 by [Le Blévec et al. \(2018\)](#) to model cyclicity.

249 The Shifted Pluri-Gaussian Simulations Approach

250 Here we summarize the modeling method developed in [Le Blévec et al. \(2018\)](#). The
251 method is based on Pluri-Gaussian Simulations (PGS) ([Armstrong et al., 2011](#)),
252 which generalizes Truncated Gaussian Simulations by using several Gaussian ran-
253 dom functions instead of just one. An example is given in Figure 4b, which shows
254 a Pluri-Gaussian Simulation using two Gaussian random functions Z_1 and Z_2 . The
255 variations of each Gaussian random function are controlled by their respective co-
256 variance model (Eqs. [A.4a](#) and [A.4b](#)). The truncation rule applied to them is two
257 dimensional (Fig. 4a) and defines in this example three facies separated by two
258 thresholds q_1 and q_2 , with all three facies in contact with each other. For instance,
259 the defined facies is yellow if q_1 is smaller than Z_1 and q_2 is smaller than Z_2 . The cor-
260 responding facies succession (Fig. 4b) shows no specific cyclicity, because all facies
261 can transition into each other randomly.

262 In order to model cyclicity, [Le Blévec et al. \(2018\)](#) introduced a spatial shift
263 α between the two Gaussian random functions. More specifically, the Gaussian
264 random functions are correlated (or anti-correlated) to each other by a correlation
265 coefficient β , then shifted by a vector noted α (Eq. [A.3](#)), which gives the direction
266 of the cyclicity. This is illustrated in Figure 4c, in which the Gaussian random

267 functions are anti-correlated ($\beta < 0$), with a small shift α oriented upward. This
268 results, after truncation, into a highly cyclic facies succession (Fig. 4c). The upward
269 cycle from facies 1 to facies 2 then to facies 3 is repeated almost everywhere because
270 Z_2 tends to cross its threshold q_2 (from facies 2 to facies 3) just after Z_1 crosses its
271 threshold q_1 (from facies 1 to facies 2), as if the truncation rule (Fig. 4a) had an
272 anti-clockwise motion in the upward direction. The cyclicity of the succession shown
273 in Figure 4c is confirmed by its corresponding transiograms (Fig. 2) as explained
274 previously.

275 Modeling Vertical Facies Cyclicity and Rhythmicity in the 276 Latemar Carbonate Platform

277 Dataset

278 The Triassic Latemar carbonate platform (northern Italy) is renowned for its cyclic-
279 ity (Goldhammer et al., 1990; Hinnov and Goldhammer, 1991) and is thus well suited
280 for analysis to the new method described above. Using the original data of Peter-
281 hänsel and Egenhoff (2008), part of the Upper Cyclic Facies interval has previously
282 been modelled by Le Blévec et al. (2018) with a simplified, three-fold classification
283 of depositional facies that is modified from Egenhoff et al. (1999). Here, the same
284 interval is modelled in the Cimon Latemar outcrop with the full four-fold classifi-
285 cation of depositional facies of Egenhoff et al. (1999): subtidal (e_1), intertidal (e_2),
286 supratidal (e_3) and subaerial exposure facies (e_4). Diagenetic overprinting is at first
287 not considered in the model described here, but models of the Latemar platform
288 presented later include diagenetic facies. Although depositional facies are here de-
289 nominated as environments of deposition, their interpretation is directly based on
290 application of the Dunham classification to observations in thin sections (Egenhoff
291 et al., 1999). Therefore, it is possible that they transition laterally with each other
292 several times at the same stratigraphic level, in a mosaic-like fashion, as shown by
293 the interpreted cross section of Peterhänsel and Egenhoff (2008). The measured
294 sections of Peterhänsel and Egenhoff (2008) are presented in Figure 5.

295 As discussed by Egenhoff et al. (1999) and Peterhänsel and Egenhoff (2008), the
296 facies tend to be organized in shallowing-upward facies cycles that comprise, from
297 base to top, facies 1, facies 2, facies 3, facies 4. This interpreted organization is
298 supported by logs in Figure 5. For example, the subtidal facies e_1 tends to overlie
299 the subaerial exposure facies e_4 , which defines the base of a cycle, and is generally
300 overlain by intertidal facies e_2 . However, many cycles are incomplete and lack one or
301 more facies (Fig. 5). There is also a high number of alternations between intertidal
302 facies e_2 and subaerial exposure facies e_4 (e.g., in log N17, Fig. 5). Therefore, the
303 cyclicity of the facies succession is not perfect and this imperfect pattern should be
304 reproduced statistically in the model. It is also noteworthy that subtidal facies e_1
305 and supratidal facies e_3 are never in contact (Fig. 5).

306 Model

307 The first step is to define an appropriate truncation rule based on the observed
308 contacts between facies and their cyclicity. As the typical cycle is $e_1 \rightarrow e_2 \rightarrow e_3 \rightarrow$
309 e_4 , these facies should be arranged clockwise (or counter-clockwise) in the truncation
310 rule. Moreover, as observed (Fig. 5), subtidal facies should not be in contact with
311 supratidal facies. Figure 6 shows a truncation rule satisfying these constraints. The
312 thresholds q_1 , q_2 and q_3 are computed according to the proportions of the different
313 facies (an example of how to compute the thresholds from the proportions is given
314 in the Appendix, using Eqs. A.7, A.8 and A.9).

315 The next step is to find the parameters of the model (β_{12} , α_{12} , r_1, r_2 , b_1, b_2)
316 from the experimental transiograms computed from the logs. The results are shown
317 in Figure 7 (grey points). The parameters of the method are found so that they
318 generate theoretical transiograms that provide a good match to the experimental
319 transiograms. The computation of a theoretical transiogram from the parameters
320 of the method is explained in the Appendix (Eqs. A.10, A.11 and A.12). A trial-
321 and-error test is performed on the parameters, and the ones which give the best fit
322 between experimental and theoretical transiograms are chosen.

323 It is important to note that transiograms are inter-dependent and cannot be fitted
324 individually. For instance, the first maximum of the auto-transiograms and first
325 minimum of the cross-transiograms are related to the cycle thickness, as explained
326 earlier (Fig. 2). Thus, one parameter such as the shift α_{12} controls the behavior of
327 several theoretical transiograms (see [Le Blévec et al., 2017](#), for details). Because
328 of these relationships between transiograms, it is usually not possible to obtain a
329 perfect fit between experimental and theoretical transiograms, and a compromise
330 should be made based on which feature (or combination of features) is considered
331 by the user to be more important.

332 The theoretical transiograms after fitting are shown in Figure 7 (black curves).
333 The tangents at the origin of the different auto-transiograms and cross-transiograms
334 are matched, which means that the different facies thicknesses and the contacts
335 between them are well constrained. Therefore, the fit between experimental and
336 theoretical transiograms is satisfactory. Subtidal facies e_1 and supratidal facies e_3
337 are not in contact because $t_{e_1e_3}(h)$ and $t_{e_3e_1}(h)$ both have a tangent at the origin
338 with a very low value, which comes from the truncation rule (Fig. 6). The only
339 significant mismatch is for the transiogram $t_{e_4e_1}(h)$, for which the tangent at the
340 origin of the theoretical transiogram is not high enough (Fig. 7). This means that
341 in the model, facies e_1 has less tendency to overlie facies e_4 than in the dataset.

342 Some rhythmic facies patterns are also captured, such as the one observed in the
343 transiogram $t_{e_1e_1}(h)$ (Fig. 7).

344 The scale factors in the lateral direction r_x and r_y are chosen by visual comparison
345 of the resulting facies models with the outcrop cross section of [Peterhänsel and](#)
346 [Egenhoff \(2008\)](#). The higher their values, the larger the extent of the facies. As the
347 facies are quite laterally extensive, the scale factors are of the order of the size of
348 the final Earth model of depositional facies.

349 Simulation

350 The Earth model for depositional facies is now built using the above fitted parame-
351 ters. The Gaussian random functions are simulated in the grid described below, and
352 then truncated into facies. The simulations are also conditioned to the measured
353 sections so that the facies observed in the measured sections are reproduced in the
354 model realizations. The algorithms to perform these steps are explained in [Le Blévec](#)
355 [et al. \(2018\)](#).

356 The number of grid cells in each direction (East, North, Z) is (100, 10, 182), and
357 the grid dimensions are (1000 m, 250 m, 9.1 m) \sim (0.62 mi, 820 ft, 29 ft). Hence the
358 size of each cell is (10 m, 25 m, 5 cm) \sim (33 ft, 82 ft, 16 ft). The number of cells is
359 a compromise between the desired computational speed of the simulation and the
360 level of details at which the heterogeneities are represented. Here, a high resolution
361 is chosen in the vertical direction, because most of the transitions between facies are
362 vertical. The simulation is fast and several equiprobable realizations are obtained in
363 a two or three minutes with a standard desktop PC. Two realizations are shown in
364 [Figure 8](#), together with the original measured sections of [Peterhänsel and Egenhoff](#)
365 [\(2008\)](#), reproduced in both realizations.

366 Incomplete facies cyclicity, as observed in the measured sections ([Fig. 5](#)) is visible
367 in the realizations ([Fig. 8](#)). For instance, subaerial exposure facies e_4 are not only
368 overlain by subtidal facies e_1 at the base of each cycle, but also by intertidal facies e_2
369 or supratidal facies e_3 . Subtidal facies e_1 and supratidal facies e_3 are not in contact,
370 as defined by the truncation rule ([Fig. 6](#)). Laterally, facies transition randomly
371 into each other because no lateral transition constraint is given. This aspect of the
372 Earth model realizations can be improved by using conceptual knowledge of the
373 platform-interior facies architecture, leading to Earth models that exhibit lateral
374 facies cyclicity or non-stationarity, as illustrated below.

375 For model validation, the transiograms are computed in three realizations of
376 the simulation and shown in [Figure 7](#) (thin grey curves). Most transiograms of
377 the realizations are a good fit to the experimental and theoretical transiograms,

378 which shows that the Earth models are geologically realistic. Some mismatches also
379 appear, for instance in $t_{e_2e_2}(h)$, for which it seems that the realizations have a higher
380 plateau than the model. However, these statistical variations are not systematic and
381 are common with stochastic simulations (Chiles and Delfiner, 2012).

382 **Extension to Lateral Cyclicity**

383 Lateral facies cyclicity can be observed in modern environments or generated by
384 forward stratigraphic modeling (Burgess et al., 2001). Tidal-flat and reef islands
385 deposits in modern shallow-water carbonate environments can exhibit lateral di-
386 rectionality, induced by currents in the water column, which results in lateral and
387 vertical facies cyclicity (e.g., Burgess et al., 2001). The method developed here mod-
388 els such vertical and lateral facies cyclicity by adding a lateral component to the shift
389 α between the Gaussian random functions. This procedure is demonstrated using a
390 satellite image of reef islands in the interior of the modern Bermuda platform, which
391 was first described by Verrill (1907) (Fig. 9a). The reef island deposits show a lateral
392 facies asymmetry, with a typical facies cycle comprising reef, backreef, and lagoonal
393 facies (after Jordan Jr, 1973). Although there are no data describing the vertical
394 facies succession, we assume that Walther’s law (Middleton, 1973) is followed, such
395 that the lateral facies transitions are equivalent to the vertical facies transitions.
396 This equivalence is modeled by incorporating the lateral component into the shift
397 vector.

398 One unconditional (no vertical sections are matched) realization of an Earth
399 model for facies distribution is shown in Figure 9c, along with the model truncation
400 rule (Fig. 9b). The three modelled facies are in contact, and lateral facies cyclicity
401 similar to that observed in the satellite image is generated. The vertical cyclicity is
402 such that reef facies overlie backreef facies (Fig. 9c). The combination of lateral and
403 vertical facies cyclicity results in an overall eastward progradation of reef islands.
404 Therefore, the shift controls the movement over time of the facies belts and bodies.
405 For instance, if the shift was oriented to the west, then this would be the direction

406 of progradation. If the shift was purely vertical, there would only be aggradation.
407 Due to the lateral component of the shift, Walther's Law is respected in the model.

408 **Modeling Diagenesis with Shifted Pluri-Gaussian Sim-** 409 **ulations**

410 Siliciclastic and, particularly, carbonate reservoirs are widely documented to un-
411 dergo extensive diagenetic modification during deposition and subsequent burial,
412 which modifies their petrophysical properties (Moore and Wade, 2013). Therefore,
413 it is important to provide a flexible modeling method for diagenetic overprinting
414 of depositional facies that can mimic patterns resulting from multiple diagenetic
415 events, in order to capture the impact on hydrocarbon recovery. Diagenesis can
416 follow the original depositional fabric in some cases, but can also be templated by
417 faults and fractures and thus cross-cut depositional facies. A novel method able to
418 model these two end members, based on the Shifted Pluri-Gaussian Simulations is
419 presented. By adding a third Gaussian function that controls diagenetic facies, the
420 method co-simulates a depositional facies field and a diagenetic facies field. The
421 shifts between the three Gaussian random functions allows the user to model asym-
422 metric relations between diagenetic and depositional facies.

423 **Modeling Syn-Depositional Diagenesis: Revisiting the Latemar** 424 **Carbonate Platform**

425 **Syn-Depositional Diagenesis in the Latemar Platform**

426 Previously we modelled the depositional facies of the Latemar carbonate platform
427 using the measured sections of Peterhänsel and Egenhoff (2008) as input data. How-
428 ever, the studies of Egenhoff et al. (1999) and Peterhänsel and Egenhoff (2008) also
429 show that diagenesis affect these facies. Tepee structures, dolomitization and caliche
430 crusts suggest an early diagenetic overprinting.

431 The measured sections of [Peterhänsel and Egenhoff \(2008\)](#) are again chosen as
432 data for the model. These sections show two diagenetic facies: completely dolomi-
433 tized crusts and partial dolomitization, which overprint different depositional facies
434 (Fig. 10). The dolomitic crust diagenetic facies only occurs in conjunction with sub-
435 aerial exposure depositional facies, while the partially dolomitized diagenetic facies
436 occurs in conjunction with intertidal and (marginally) supratidal depositional facies.
437 This observation from vertical measured sections is supported by the interpreted lat-
438 eral correlations of [Peterhänsel and Egenhoff \(2008\)](#), in which the dolomitic crust
439 diagenetic facies transitions laterally only into subaerial exposure depositional fa-
440 cies. Table 1 shows the proportions of each diagenetic facies within each depositional
441 facies.

442 Model

443 In the Earth model realizations shown in Figure 8, depositional facies were mod-
444 elled using two Gaussian random functions. If diagenetic facies were included in
445 the corresponding two dimensional truncation rule (Fig. 6), they would necessarily
446 have geometrical properties similar to those of depositional facies. Adding a third
447 Gaussian random function as a third dimension in the truncation rule gives a greater
448 flexibility to represent diagenetic facies geometries and their relations with deposi-
449 tional facies. Moreover, diagenesis can then be modelled as a superimposed facies
450 field that overprints the depositional facies as explained in [Renard et al. \(2008\)](#).

451 A three dimensional truncation rule for the Latemar platform is thus defined in
452 Figure 11. The truncation rule for the depositional facies is the same as that shown
453 in Figure 6. The third Gaussian random function defines two diagenetic facies:
454 dolomitic crust d_1 (which overprints subaerial exposure depositional facies e_4) and
455 partial dolomite d_2 (which overprints intertidal and supratidal depositional facies e_2
456 and e_3). Depositional facies e_1 is not affected by diagenesis. The thresholds q con-
457 trolling the proportions of diagenetic facies within depositional facies are computed
458 from Table 1, as explained in the Appendix (Eq. A.7). For example, diagenetic

459 facies d_2 is more abundant in depositional facies e_2 than in depositional facies e_3 ,
460 and so its area is larger in the truncation rule (Fig. 11).

461 Once the truncation rule is chosen, the experimental transiograms of diagenetic
462 facies are fitted with the parameters of the method, as described previously. Cross-
463 transiograms between depositional and diagenetic facies are fitted first, because they
464 are controlled by a smaller number of parameters: α_{13} , β_{13} , α_{23} and β_{23} (Eq. A.3).
465 These parameters define the relations of the first two Gaussian random functions
466 Z_1 and Z_2 with the third Gaussian random function Z_3 and thus control relations
467 between depositional facies and diagenetic facies. These cross-transiograms have
468 different properties from usual cross-transiograms because they relate to two super-
469 imposed facies fields, for which facies can both be present at the same location.
470 Therefore, their value at a distance $h = 0$ is not 0 but the probability of finding
471 the two facies types at the same location (Table 1). The fit between theoretical
472 transiograms (black curves, Eq. A.11) and experimental transiograms (grey points)
473 of depositional facies and diagenetic facies is shown in Figure 12.

474 For most transiograms, the experimental curve at the first distance step is com-
475 monly higher than the theoretical curve (Fig. 12). This is due to the small number
476 of data points, because there are few occurrences of diagenetic facies in the mea-
477 sured sections (Fig. 10), thus causing the transiograms to be statistically unreliable.
478 However, it is worth noting that the theoretical transiograms generally show reason-
479 able behaviors at the origin. For instance, the tangent at the origin of transiogram
480 $t_{d_1e_1}(h)$ has a high value (Fig. 12), which shows that subtidal depositional facies
481 tends to overlie dolomitic crust diagenetic facies, as observed in the measured sec-
482 tions (Fig. 10). This spatial relationship supports the facies cyclicity of the model,
483 because the dolomitic crust diagenetic facies is present in the subaerial exposure
484 depositional facies, which are themselves overlain by subtidal depositional facies.
485 Similarly the high value of the tangent at the origin of transiogram $t_{d_2e_4}(h)$ shows
486 that subaerial exposure depositional facies tends to occur above partial dolomite
487 diagenetic facies, which is also observed in the measured sections (Fig. 10). The

488 transiograms thus confirm that the method is able to capture asymmetry between
489 depositional and diagenetic facies, so that diagenetic facies are ordered with respect
490 to the depositional facies.

491 As stated above, the cross-transiograms between depositional facies and diage-
492 netic facies are not equal to zero for a zero distance. For instance transiogram $t_{d_1e_4}(h)$
493 starts at a value close to 1 (Fig. 12) because the dolomitic crust diagenetic facies d_1
494 is only present in the subaerial exposure diagenetic facies e_4 . The cross-transiogram
495 then decreases abruptly, which suggests that units of the subaerial exposure dia-
496 genetic facies are thin, which is consistent with the measured sections (Fig. 10).
497 Finally, rhythmicity, although not very pronounced, is captured in transiograms
498 $t_{d_2e_1}(h)$ and $t_{d_2e_2}(h)$ (Fig. 12). This suggests that partial dolomite diagenetic facies
499 d_2 is separated from subtidal depositional facies e_1 by a nearly constant thickness
500 of 0.3 m and that partial dolomite diagenetic facies d_2 is separated from intertidal
501 depositional facies e_2 by a nearly constant thickness of 1 m (i.e., the first maxima
502 of transiograms $t_{d_2e_1}(h)$ and $t_{d_2e_2}(h)$; Figure 12).

503 The auto- and cross-transiograms of the diagenetic facies themselves are now
504 fitted using the same procedure. The parameters controlling these transiograms
505 are the parameters of the third covariance r_3 and b_3 (Eq. A.5), and all the other
506 parameters mentioned above, which are left unchanged. They control the spatial
507 properties of Z_3 and thus the geometries of diagenetic facies. Figure 13 shows the fit
508 between experimental and theoretical transiograms. The method is able to capture
509 the asymmetry of cross-transiograms between the two diagenetic facies as $t_{d_2d_1}(h)$
510 and $t_{d_1d_2}(h)$, showing that the dolomitic crust diagenetic facies d_1 tends to overlie
511 the partial dolomite diagenetic facies d_2 , and the modelled transiograms are able to
512 match exactly this behavior at the origin (Fig. 13). Theoretical auto-transiograms
513 $t_{d_1d_1}(h)$ and $t_{d_2d_2}(h)$ also exhibit the correct behavior at the origin, which confirms
514 that the mean thicknesses of these diagenetic facies are well constrained (Fig. 13).

515 This section has shown the value of the method for capturing complex tran-
516 siograms between depositional facies and diagenetic facies. Shifts α_{13} and α_{23} play

517 an important role, which emphasizes the value of incorporating asymmetry in the
518 modeling of syn-depositional diagenetic patterns.

519 **Simulation**

520 The simulation is performed as for previously described models (e.g., Figure 8),
521 with the added third Gaussian random function. Two realizations of the Earth
522 model showing diagenetic facies superimposed on depositional facies are shown in
523 Figure 14. Both realizations honor the data along the measured sections (e.g., long
524 measured section N22; Figure 14), but differ away from them (e.g., in the volume
525 above measured section N22; Figure 14).

526 To verify that the resulting simulations honor the data statistics, transiograms
527 are computed on three realizations (thin grey curves in Figure 12 and Figure 13).
528 The simulated transiograms match the experimental transiograms quite well, even
529 better than the theoretical transiograms. For instance, transiograms $t_{d_2e_2}(h)$ for the
530 realizations reproduce the complex hole-effect observed in the data (Fig. 12). Sim-
531 ilarly, transiogram $t_{d_2d_2}(h)$ of the realizations follows the experimental transiogram
532 more closely than the theoretical transiogram (Fig. 13). This could be due to the
533 conditioning of the simulation, which provides significant constraints on the Earth
534 models.

535 **Modeling syn-depositional diagenesis in non-stationary shallow-** 536 **marine deposits, Book Cliffs, Utah**

537 The Upper Cretaceous Spring Canyon Member of the Blackhawk Formation, which
538 is exposed in the Book Cliffs (Utah), consists of shallow-marine, wave-dominated
539 shoreface sandstones that contain overprinting diagenetic features such as carbonate-
540 cemented concretions and leached zones (whitecaps) (Van Wagoner et al., 1990;
541 Kamola and Huntoon, 1995; Hampson and Storms, 2003; Taylor et al., 2004). Due
542 to their large lateral extent, the deposits display non-stationary facies proportions
543 from proximal to distal. Herein we model the outcrop dataset to show the flexibility

544 of the method and highlight the use of embedded transition probabilities in a non-
545 stationary context.

546 **Dataset**

547 The nine measured sections reported by [Taylor et al. \(2004\)](#) are used here, and the
548 facies classification is simplified into three depositional facies: distal lower shoreface
549 heteroliths and offshore mudstones (E_1), proximal lower shoreface sandstones (E_2)
550 and foreshore and upper shoreface sandstones (E_3). There are also two diagenetic
551 facies: carbonate cement D_1 and leached sandstones ("whitecaps") D_2 , in which
552 carbonate material has been removed via syn-depositional diagenesis. [Table 2](#) shows
553 the proportion of each diagenetic facies within the different depositional facies, based
554 on measured sections with this facies classification ([Fig. 15](#)).

555 No cyclicity is observed between depositional facies. Facies proportions in each
556 measured section (represented by pie charts in [Figure 15](#)) show that from west
557 (proximal) to east (distal), the proportion of depositional facies E_3 decreases while
558 that of depositional facies E_1 increases. Vertical facies proportion curves show that
559 depositional facies E_3 is only present at the top of the Spring Canyon Member in the
560 area sampled by the measured sections. Diagenetic facies are also non-stationary
561 because their distribution is controlled by the distribution of depositional facies
562 ([Table 2](#)).

563 **Modeling Non-Stationary Facies Proportions**

564 As described above, facies proportions vary systematically over the dataset to be
565 modelled ([Fig. 15](#)). Therefore, the final Earth model should account for these vari-
566 ations. This is achieved by estimating the proportions of each facies in each cell of
567 the Earth model. ([Armstrong et al., 2011](#); [D'Or et al., 2017](#)). First, the proportions
568 of each facies are computed at each horizontal level from all the measured sections
569 to give vertical facies proportion curves ([Fig. 15](#)). The vertical proportion curves are
570 then smoothed with a moving average algorithm to remove random variations, as

571 described in [White et al. \(2003\)](#). Then, the proportions of each facies are computed
572 at each vertical measured section (pie charts of [Figure 15](#)). Finally, at each grid
573 cell intersecting a measured section, the proportion of each facies is calculated by
574 averaging the proportion given by the vertical proportion curve with the proportion
575 of the facies at the measured section.

576 A procedure to interpolate these facies proportions between measured sections is
577 then required. Here this is achieved by lateral simple kriging interpolation ([Chiles
578 and Delfiner, 2012](#)) using a Gaussian covariance with a large scale factor and a
579 mean chosen as the global proportion of each facies. Once the proportions of each
580 facies have been calculated for every cell of the model, they are transformed into
581 thresholds for the Gaussian random functions according to the same procedure used
582 for the models presented earlier (Appendix, [Eq. A.8](#)).

583 **The Model**

584 The truncation rule can be inferred from the observation of facies contacts in the
585 measured sections ([Fig. 15](#)). Because of the facies distribution's non-stationarity,
586 the truncation rule is different in every cell and depends on the cell's facies pro-
587 portion. Therefore, a general truncation rule is first defined in [Figure 16](#), which is
588 then adapted to the local facies proportions in the different cells of the Earth model
589 ([Fig. 16](#)). The foreshore and upper shoreface sandstone depositional facies (E_3) and
590 distal lower shoreface heteroliths and offshore mudstone depositional facies (E_1) are
591 not in contact, because of a limited presence of foreshore and upper shoreface sand-
592 stones (which occur only five times in the measured sections) and non stationarity
593 ([Fig. 15](#)). However, there is no reason why these facies should not be in contact
594 away from the measured sections, and the global truncation rule is thus defined to
595 allow this contact relationship ([Fig. 16](#)). The carbonate cement diagenetic facies
596 (D_1) and leached sandstone diagenetic facies (D_2) are respectively present in the
597 proximal lower shoreface sandstone depositional facies (E_2) and both the proximal
598 lower shoreface sandstone depositional facies and the foreshore and upper shoreface

599 sandstone depositional facies (E_2 , E_3) (Fig. 16).

600 Transiograms are not fitted here because their behavior is strongly influenced
 601 by non stationarity, especially at long distances (Armstrong et al., 2011). However,
 602 embedded transition probabilities (Krumbein and Dacey, 1969) are not much af-
 603 fected by non stationarity because they just measure facies juxtapositions. They
 604 can be deduced from the parameters of the model by taking the derivative at the
 605 origin of the transiograms (Eq. A.13). Thus, they are compared to the experimen-
 606 tal embedded transitions computed from the measured sections in order to infer the
 607 parameters α_{12} and β_{12} . The experimental (red) and model (blue) embedded matrix
 608 for the three depositional facies after fitting is

$$R_{logs/model} = \begin{bmatrix} E_1 & E_2 & E_3 \\ 0 & 1.0/0.63 & 0.0/0.36 \\ 0.72/0.79 & 0 & 0.28/0.23 \\ 0.0/0.15 & 1.0/0.85 & 0 \end{bmatrix}. \quad (1)$$

609 The matrix shows that foreshore and upper shoreface sandstones (E_3) and distal
 610 lower shoreface heteroliths and offshore mudstones (E_1) are not in contact in the
 611 measured sections because their embedded probability is zero, while in the model
 612 they can be in contact ($r_{31}=0.15$, $r_{13} = 0.36$) according to the truncation rule
 613 (Fig. 16). The embedded transitions from proximal lower shoreface sandstones (E_2)
 614 to the other depositional facies are similar in the model and in the measured sections.

615 In order to constrain the vertical component of the scale factors r_1 and r_2 , the
 616 thicknesses of the depositional facies are computed in the measured sections and
 617 matched with the theoretical thicknesses, which are obtained from the derivative at
 618 the origin of the auto-transiograms (Carle and Fogg, 1996). The resulting theoretical
 619 thicknesses for the three depositional facies E_1 , E_2 , and E_3 are respectively 1.3 m,
 620 0.8 m, and 0.5 m, while the experimental thicknesses computed from the measured
 621 sections are 1.4 m, 0.8 m, and 0.6 m, which is a good match.

622 Embedded transition probabilities between the diagenetic facies are not shown

623 because they are simply not in contact with each other. The vertical scale factor r_3 is
 624 chosen to be similar to r_1 and r_2 because diagenetic facies have a similar thickness to
 625 depositional facies. Lateral components of the scale factors r_1 , r_2 and r_3 are chosen
 626 by visual comparison of the resulting Earth model realizations and the correlation
 627 panel between measured sections of Taylor et al. (2004). The depositional facies
 628 have a large lateral extent, of the same order as the west-to-east lateral extent of
 629 the Earth model.

630 Simulation

631 The number of cells in the grid in each direction is 100 (west-to-east), 20 (north-to-
 632 south), 566 (height) and the dimensions of the grid are 20 km (\sim 12.4 mi) (west-to-
 633 east), 5 km (\sim 3.1 mi) (north-to-south), 56 m (\sim 184 ft) (height). The simulations
 634 are conditioned to the measured sections with the procedure outlined in Le Blévec
 635 et al. (2018).

636 Two realizations of the resulting Earth model are shown in Figure 17. It is clear
 637 that the realizations are non-stationary as, for instance, the proportion of foreshore
 638 and upper shoreface sandstone depositional facies (E_3) decreases towards the west.
 639 Leached sandstone diagenetic facies (D_2) also exhibit a decreasing proportion to-
 640 wards the west, because they are constrained by the presence of foreshore and upper
 641 shoreface sandstone depositional facies (E_3) (Table 2).

642 As a post-validation step, embedded transition probabilities are computed in
 643 three resulting realizations and averaged, to give the embedded matrix of transition
 644 probabilities

$$R_{simu} = \begin{bmatrix} E_1 & E_2 & E_3 \\ 0 & 0.75 & 0.25 \\ 0.79 & 0 & 0.21 \\ 0.06 & 0.94 & 0 \end{bmatrix}. \quad (2)$$

645 This matrix matches the embedded matrix computed from the measured sections
 646 (Eq. 1), although foreshore and upper shoreface sandstone depositional facies (E_3)

647 and distal lower shoreface heteroliths and offshore mudstone depositional facies (E_1)
648 are in contact, as discussed above.

649 **Towards Modeling Post-depositional Hydrothermal Diagenesis**

650 Post-depositional hydrothermal diagenesis is commonly observed at outcrop ([Jacquemyn et al., 2014](#);
651 [Vandeginste et al., 2013](#); [Beckert et al., 2015](#)) and interpreted in the
652 subsurface ([Davies and Smith Jr, 2006](#); [Smith Jr, 2006](#)). Hydrothermal diagenesis
653 produces diagenetic bodies that are discordant with strata and instead follow faults,
654 fractures and other structures. Depositional facies may differ in their permeability,
655 such that hydrothermal fluids can also flow laterally away from faults and fractures
656 along relatively permeable facies belts and bodies, thus creating so-called "Christ-
657 mas tree" geometries ([Beckert et al., 2015](#)). For example, outcrops of the Latemar
658 carbonate platform in Valsorda valley exhibit such fracture-related hydrothermal di-
659 agenesis, which generate dolomite that is distributed along and nearby to fractures
660 (Fig. 18, after [Jacquemyn et al. \(2014\)](#)).

661 This type of dolomitization can be represented in our method thanks to the third
662 Gaussian random function covariance $\rho_3(h)$, which can have a different anisotropy
663 from that of the two other Gaussian random functions covariances (Eq. A.4c). An
664 unconditional realization of such a model is shown in Figure 19. The third Gaussian
665 random function is modelled independently from the two other Gaussian random
666 functions ($\beta_{13} = \beta_{23} = 0$), so that the geometries of diagenetic dolomite bodies cut
667 across depositional facies geometries. The truncation rule controls the extent of
668 dolomite within each depositional facies (Fig. 19). Depositional facies E_3 contains
669 more dolomite than depositional facies E_2 , because the volume of dolomitized facies
670 D_{E_3} in the truncation rule is larger than that of dolomitized facies D_{E_2} . On the contrary,
671 depositional facies E_1 is not affected by diagenesis.

672 The Earth model realization (Fig. 19) shows that depositional facies tend to be
673 organized in shallowing-upward asymmetric cycles and diagenetic dolomite bodies
674 cut across them. The dolomite diagenetic facies (D_{E_2} and D_{E_3}) is more abundant

675 in depositional facies E_3 than in depositional facies E_2 , and is not present at all in
676 depositional facies E_1 , as constrained by the truncation rule.

677 Conclusion and Recommendations

678 The new method proposed in this paper models depositional and diagenetic facies
679 fields with cyclic and rhythmic patterns. The method is based on a novel Pluri-
680 Gaussian approach, using three dimensional truncation rules and Gaussian random
681 functions shifted from each other. Qualitative information and concepts are used to
682 construct the truncation rule, and the other parameters of the method are defined by
683 fitting the experimental auto- and cross- transiograms. The resulting models show
684 that a combination of lateral and vertical facies cyclicity can be used to generate
685 aggradational and progradational facies geometries.

686 In addition, we model depositional facies overprinted by cross-cutting or con-
687 formable diagenesis. This is possible because the three Gaussian random functions
688 are spatially shifted from each other, and depositional and diagenetic facies are
689 ordered according to the cross-transiograms.

690 The method has also shown its capability to model non-stationary facies propor-
691 tions, which is a predominant feature in datasets that contain pronounced proximal-
692 to-distal or axial-to-marginal facies trends. In such cases, it is not appropriate to
693 use transiograms to constrain the parameters of the method. Instead, it is sug-
694 gested to use embedded transition probabilities, because non stationarity does not
695 significantly impact facies juxtapositions.

696 The method significantly improves the capability of geostatistical Earth models
697 to represent geologically realistic facies architectures, and thus can lead to more
698 realistic geostatistical reservoir models and more accurate hydrocarbon production
699 forecasts.

700 A Appendix: Shifted Pluri-Gaussian Model

701 The model developed in this paper is an extension of that developed by [Le Blévec](#)
702 [et al. \(2018\)](#). Three Gaussian random functions Z_1, Z_2, Z_3 are correlated and shifted
703 relative to each other and truncated into facies according to a truncation rule (e.g.,
704 [Figure 11](#)). The first two Gaussian random functions control depositional facies while
705 the third Gaussian random function controls diagenetic facies. A shifted version of
706 the linear model of co-regionalization ([Wackernagel, 2003](#)) is used

$$\begin{cases} Z_1(x) = Y_1(x), \\ Z_2(x) = \beta_{12} Y_1(x + \alpha_{12}) + \sqrt{1 - \beta_{12}^2} Y_2(x), \\ Z_3(x) = \beta_{13} Y_1(x + \alpha_{13}) + \beta_{23} Y_2(x + \alpha_{23}) + \sqrt{1 - \beta_{13}^2 - \beta_{23}^2} Y_3(x), \end{cases} \quad (\text{A.3})$$

707 where $-1 < \beta_{ij} < 1$ are the correlations coefficients between $Y_i(x + \alpha_{ij})$ and $Z_j(x)$,
708 α_{ij} being the shifts, and Y_1, Y_2, Y_3 are uncorrelated Gaussian random functions with
709 respective covariances in three dimensions

$$\rho_1(h_x, h_y, h_z) = \exp\left(-\frac{h_x^2}{r_{1x}^2} - \frac{h_y^2}{r_{1y}^2} - \frac{h_z^2}{r_{1z}^2}\right) \cos(b_1 h_z), \quad (\text{A.4a})$$

710

$$\rho_2(h_x, h_y, h_z) = \exp\left(-\frac{h_x^2}{r_{2x}^2} - \frac{h_y^2}{r_{2y}^2} - \frac{h_z^2}{r_{2z}^2}\right) \cos(b_2 h_z), \quad (\text{A.4b})$$

711

$$\rho_3(h_x, h_y, h_z) = \exp\left(-\frac{h_x^2}{r_{3x}^2} - \frac{h_y^2}{r_{3y}^2} - \frac{h_z^2}{r_{3z}^2}\right) \cos(b_3 h_z), \quad (\text{A.4c})$$

712 with $r_i = (r_{ix}, r_{iy}, r_{iz})$ the scale factors in three dimensions and b_i the frequencies of
713 the cosine functions. Therefore, the auto-covariances of the three Gaussian random
714 functions Z_1, Z_2, Z_3 are respectively

$$\begin{cases} \rho_{Z_1}(h) = \rho_1(h), \\ \rho_{Z_2}(h) = \beta_{12}^2 \rho_1(h) + (1 - \beta_{12}^2) \rho_2(h), \\ \rho_{Z_3}(h) = \beta_{13}^2 \rho_1(h) + \beta_{23}^2 \rho_2(h) + (1 - \beta_{13}^2 - \beta_{23}^2) \rho_3(h), \end{cases} \quad (\text{A.5})$$

715 and the cross-covariances between them

$$\begin{cases} \rho_{Z_1 Z_2}(h) = \beta_{12} \rho_1(h + \alpha_{12}), \\ \rho_{Z_1 Z_3}(h) = \beta_{13} \rho_1(h + \alpha_{13}), \\ \rho_{Z_2 Z_3}(h) = \beta_{12} \beta_{13} \rho_1(h + \alpha_{13} - \alpha_{12}) + \beta_{23} \sqrt{1 - \beta_{12}^2} \rho_2(h + \alpha_{23}). \end{cases} \quad (\text{A.6})$$

716 These covariances are used to derive the thresholds of the Gaussian random functions
 717 from the proportions of the different facies. For instance, let us determine the
 718 threshold q_{d_1} of the third Gaussian random function Z_1 that controls the proportion
 719 of the facies d_1 (Fig. 11)

$$p_{d_1} = Pr[Z_1(x) > q_1, Z_2(x) < q_3, Z_3(x) > q_{d_1}], \quad (\text{A.7})$$

720 which can be re-written by integration of the multi-variate Gaussian density $G_\Sigma(u, v, w)$

$$p_{d_1} = \int_{q_1}^{\infty} \int_{-\infty}^{q_2} \int_{q_{d_1}}^{\infty} G_\Sigma(u, v, w) du dv dw, \quad (\text{A.8})$$

721 with Σ the covariance matrix

$$\Sigma = \begin{bmatrix} 1 & \rho_{Z_1 Z_2}(0) & \rho_{Z_1 Z_3}(0) \\ \rho_{Z_1 Z_2}(0) & 1 & \rho_{Z_2 Z_3}(0) \\ \rho_{Z_1 Z_3}(0) & \rho_{Z_2 Z_3}(0) & 1 \end{bmatrix}. \quad (\text{A.9})$$

722 Equation A.8 is then solved numerically with the algorithm of Genz (1992). The
 723 same methodology is applied to compute theoretical transiograms (Fig. 7). For
 724 instance, let us examine the transiogram between facies e_1 and e_2 (Fig. 11)

$$t_{e_1 e_2}(h) = \frac{Pr[Z_1(x) < q_1, Z_2(x) < q_2, Z_1(x+h) < q_1, Z_2(x+h) > q_2]}{p_{e_1}}, \quad (\text{A.10})$$

725 which can be re-written by integration of Gaussian multi-variate density

$$t_{e1e2}(h) = \frac{1}{p_{e1}} \int_{-\infty}^{q_1} \int_{-\infty}^{q_2} \int_{-\infty}^{q_1} \int_{q_2}^{\infty} G_{\Sigma(h)}(u, v, w, y) du dv dw dy, \quad (\text{A.11})$$

726 with $\Sigma(h)$ the Gaussian covariance matrix

$$\Sigma(h) = \begin{bmatrix} 1 & \rho_{Z_1 Z_2}(0) & \rho_{Z_1}(h) & \rho_{Z_1 Z_2}(h) \\ \rho_{Z_1 Z_2}(0) & 1 & \rho_{Z_2 Z_1}(h) & \rho_{Z_2}(h) \\ \rho_{Z_1}(h) & \rho_{Z_2 Z_1}(h) & 1 & \rho_{Z_1 Z_2}(0) \\ \rho_{Z_1 Z_2}(h) & \rho_{Z_2}(h) & \rho_{Z_1 Z_2}(0) & 1 \end{bmatrix}. \quad (\text{A.12})$$

727 Equation A.11 is then solved numerically with the algorithm of [Genz \(1992\)](#) and the
 728 same methodology is applied for the other transiograms. The embedded transition
 729 probabilities are computed from the transiograms as follows

$$r_{ij} = -\frac{t'_{ij}(0)}{t'_{ii}(0)}. \quad (\text{A.13})$$

730

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892 Table and Figure Captions

893 Table 1: Proportions of diagenetic facies overprinted on depositional facies in the
894 Latemar carbonate platform, taken from measured sections (Fig. 10).

895 Table 2: Proportions of diagenetic facies overprinted on depositional facies in the
896 Spring Canyon Member of the Blackhawk Formation, taken from measured sections
897 (Fig. 15).

898 Figure 1: Four synthetic facies successions: (a) non rhythmic with two cycles;
899 (b) rhythmic (blue facies) with two cycles; (c) cyclic and non rhythmic; and (d)
900 cyclic and rhythmic. Modified from [Le Blévec et al. \(2018\)](#).

901 Figure 2: Cyclic and rhythmic facies succession (a) with associated transiogram
902 matrix between facies 1 and 2 (b-e). \overline{L}_c is the mean thickness of a facies cycle, and
903 \overline{L}_1 and \overline{L}_2 are the mean thicknesses of facies 1 and 2. Proportion of facies 1 is 0.5
904 and proportion of facies 2 is 0.25.

905 Figure 3: Facies succession (b) modelled with Truncated Gaussian Simulations
906 according to the truncation rule (a) and parameters $r_1 = 0.1$ m , $b_1 = 0$ m⁻¹,
907 $(p_1, p_2, p_3) = (0.4, 0.4, 0.2)$ (Eq. A.4a).

908 Figure 4: Comparison between conventional Pluri-Gaussian Simulation (PGS)
909 (b) and shifted PGS (c) with the same truncation rule (a). For (b), the parameters
910 are $r_1 = r_2 = 0.6$ m, $b_1 = 15$ m⁻¹, $b_2 = 30$ m⁻¹ (Eqs. A.4a, A.4b), and facies
911 proportions $(p_1, p_2, p_3) = (0.5, 0.25, 0.25)$ and for (c), the same parameters are ap-
912 plied together with the shift $\alpha_{12} = 0.04$ m and correlation coefficient $\beta_{12} = -0.7$
913 (Eq. A.3).

914 Figure 5: Measured sections through part of the Upper Cyclic Facies interval in
915 the Cimon Latemar outcrop, Latemar platform. Figure modified from [Peterhänsel
916 and Egenhoff \(2008\)](#).

917 Figure 6: Truncation rule used for modeling depositional facies in the Latemar
918 platform dataset (Fig. 5).

919 Figure 7: Experimental transiograms (grey points) in the upward vertical direc-
920 tion of depositional facies computed from the measured sections shown in Figure 5,

921 theoretical transiograms fitted to these points (black line), and transiograms com-
922 puted in three realizations of the depositional facies Earth model (thin grey lines).
923 The parameters used for the theoretical transiograms are $r_1 = (800, 800, 0.3)$ m,
924 $r_2 = (800, 800, 1.2)$ m, $b_1 = 0 \text{ m}^{-1}$, $b_2 = 5 \text{ m}^{-1}$, $\beta_{12} = 0.67$, $\alpha_{12} = 0.1$ m (Eqs. A.4a,
925 A.4b).

926 Figure 8: Two realizations of an Earth model for depositional facies in the Cimon
927 region of the Latemar carbonate platform conditioned by four measured sections
928 (Fig. 5) with modeling parameters explained in Figures 7.

929 Figure 9: Three dimensional unconditional realization from a satellite image
930 of Bermuda carbonate platform interior. (a) satellite image (with latitudinal and
931 longitudinal position) showing three types of facies based on visual interpretation:
932 blue represents the lagoon, light green the backreef, and dark green the reef; (b)
933 truncation rule; and (c) 3D Earth model of facies distributions. The parameters
934 of the simulation are $r_1 = r_2 = (20, 100, 0.4)$ m, $\alpha_{12} = (0.1, 5)$ m, $(p_1, p_2, p_3) =$
935 $(0.15, 0.15, 0.7)$ (Eqs. A.4a, A.4b).

936 Figure 10: Depositional facies and diagenetic facies in the measured sections
937 through part of the Upper Cyclic Facies in Cimon Latemar outcrop, Latemar car-
938 bonate platform (Fig. 5). Measured sections are adapted from Peterhänsel and
939 Egenhoff (2008).

940 Figure 11: Three dimensional truncation rule used for modeling the depositional
941 facies and diagenetic facies in the Latemar platform dataset (Fig. 10, Table 1).

942 Figure 12: Experimental vertical cross-transiograms between depositional facies
943 and diagenetic facies (grey points) from measured sections shown in Figure 10, theo-
944 retical cross-transiograms fitted to these points (black lines), and cross-transiograms
945 computed in three realizations of a resulting Earth model (thin grey lines). The pa-
946 rameters defining the theoretical transiograms are the same as those for Figure 7,
947 with in addition $\beta_{13} = -0.8$, $\beta_{23} = -0.5$, $\alpha_{13} = -0.1$ m, $\alpha_{23} = 0.1$ m (Eq. A.3).

948 Figure 13: Experimental transiograms between diagenetic facies (grey points),
949 theoretical transiograms fitted to these points (black lines), and transiograms com-

950 puted in three realizations of a resulting Earth model (thin grey lines). The param-
951 eters defining the theoretical transiograms are the same as those for Figure 7 and
952 12, with in addition $r_3 = (800, 800, 0.3)$ m and $b_3 = 0$ m⁻¹ (Eq. A.4c).

953 Figure 14: Two realisations of an Earth model for depositional facies and di-
954 agenetic facies in the Cimon Latemar region of the Latemar carbonate platform,
955 conditioned by four measured sections with modeling parameters noted in Figure 7,
956 12 and 13.

957 Figure 15: measured sections through the Spring Canyon Member, Blackhawk
958 Formation in outcrops of the Book Cliffs, as reported by Taylor et al. (2004), with
959 simplified classification of depositional facies and diagenetic facies, corresponding fa-
960 cies vertical proportion curves, and pie charts of facies proportions in each measured
961 section.

962 Figure 16: Global truncation rule and two examples of local truncation rules
963 for modeling the Spring Canyon Member, Blackhawk Formation in outcrops of the
964 Book Cliffs. The facies E are depositional facies and D are diagenetic facies.

965 Figure 17: Two realizations of an Earth model for depositional facies and di-
966 agenetic facies in the Spring Canyon Member, Blackhawk Formation, conditioned
967 by nine measured sections with modeling parameters $r_1 = (0.6, 3000, 3000)$ m, $r_2 =$
968 $(0.7, 3000, 3000)$ m, $r_3 = (1, 1500, 1500)$ m, $b_1 = b_2 = 0$ m⁻¹, $\alpha_{12} = \alpha_{13} = \alpha_{23} = 0$
969 m, $\beta_{12} = \beta_{13} = \beta_{23} = 0$ (Eqs. A.3, A.4a, A.4b, A.4c).

970 Figure 18: Uninterpreted (a) and (c); and interpreted (b) and (d) photographs of
971 post-depositional hydrothermal dolomite associated with fractures in the Valsorda
972 valley outcrops of the Latemar carbonate platform. Hydrothermal dolomite confined
973 to the fracture area is shown in red, and hydrothermal dolomite expanding in the
974 host rock is shown in yellow. Modified from Jacquemyn et al. (2014).

975 Figure 19: (a) Truncation rule and (b) resulting unconditional realization of
976 Earth model of depositional facies (cf. Figure 8) overprinted by post-depositional
977 hydrothermal dolomite diagenetic facies. The parameters used for the simulation
978 are $r_1 = r_2 = (60, 60, 0.2)$ m, $r_3 = (5, 5, 5)$ m, $b_1 = b_2 = b_3 = 0$ m⁻¹, $\beta_{12} = 0.99$,

979 $\alpha_{12} = 0.1$ m, $\beta_{13} = 0.8$, $\beta_{23} = 0$, $\alpha_{13} = \alpha_{23} = 0$ m (Eqs. [A.3](#), [A.4a](#), [A.4b](#), [A.4c](#)).

Table 1:

	Dolomitic crust	Partial dolomite
Subtidal	0	0
Intertidal	0	0.10
Supratidal	0	0.02
Exposure	0.32	0

Table 2:

	Carbonate concretion	White caps
Distal mudstones	0	0
Shoreface sandstones	0.21	0.03
Foreshore sandstones	0.59	0.4

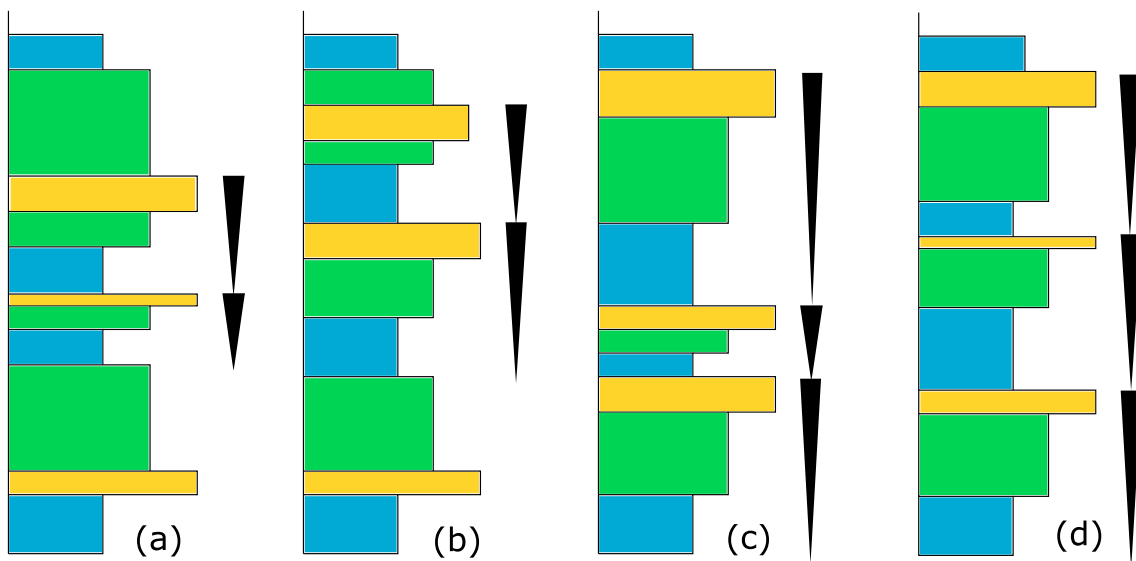


Figure 1:

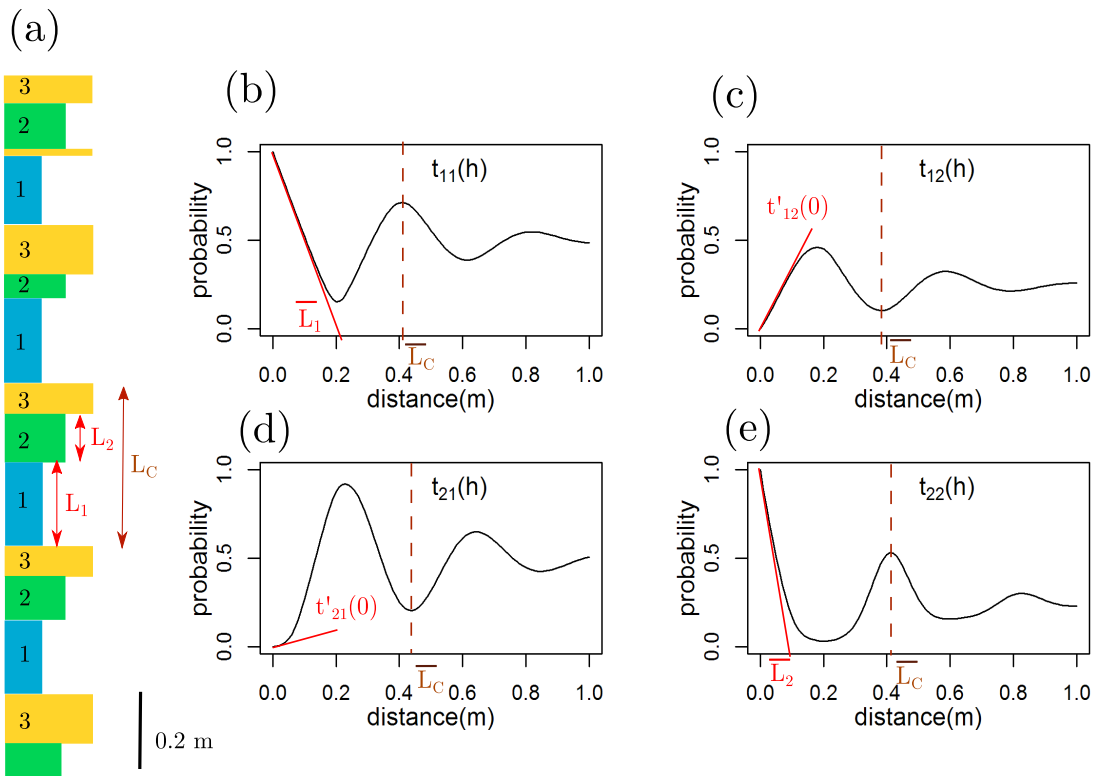


Figure 2:

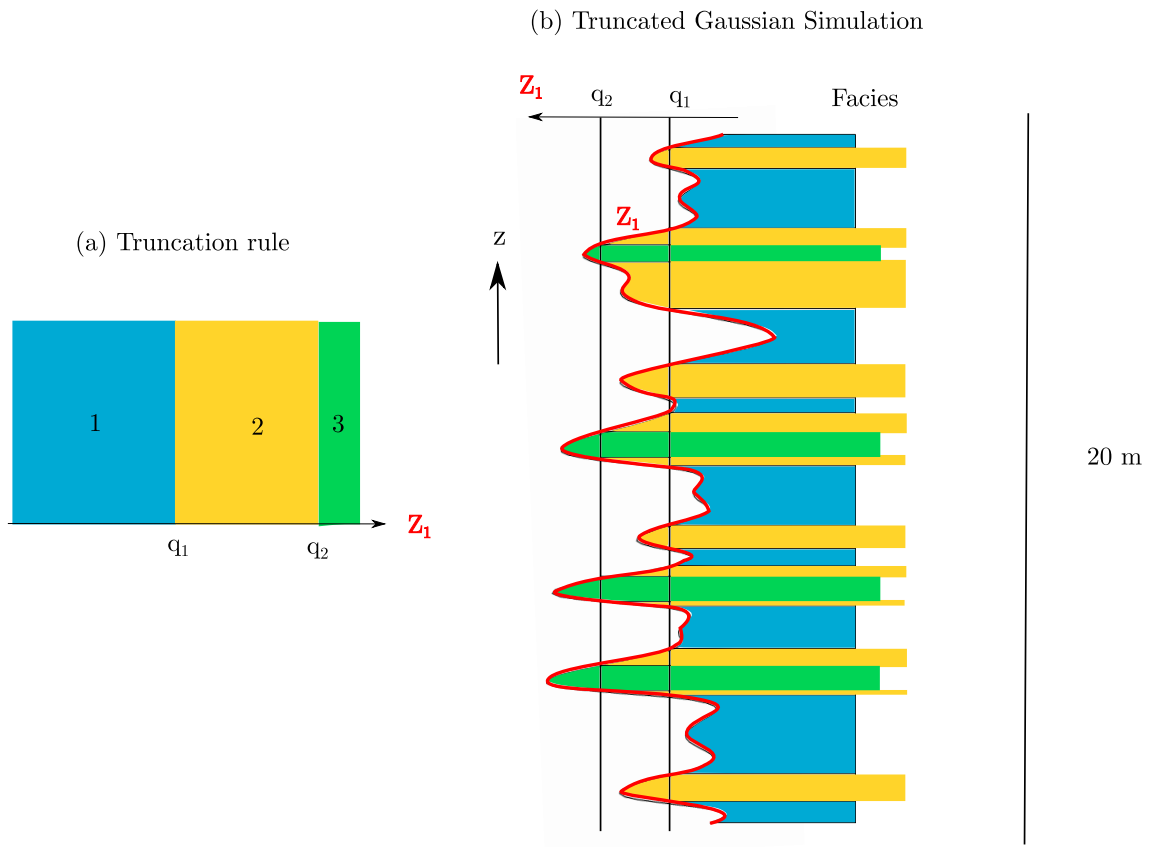


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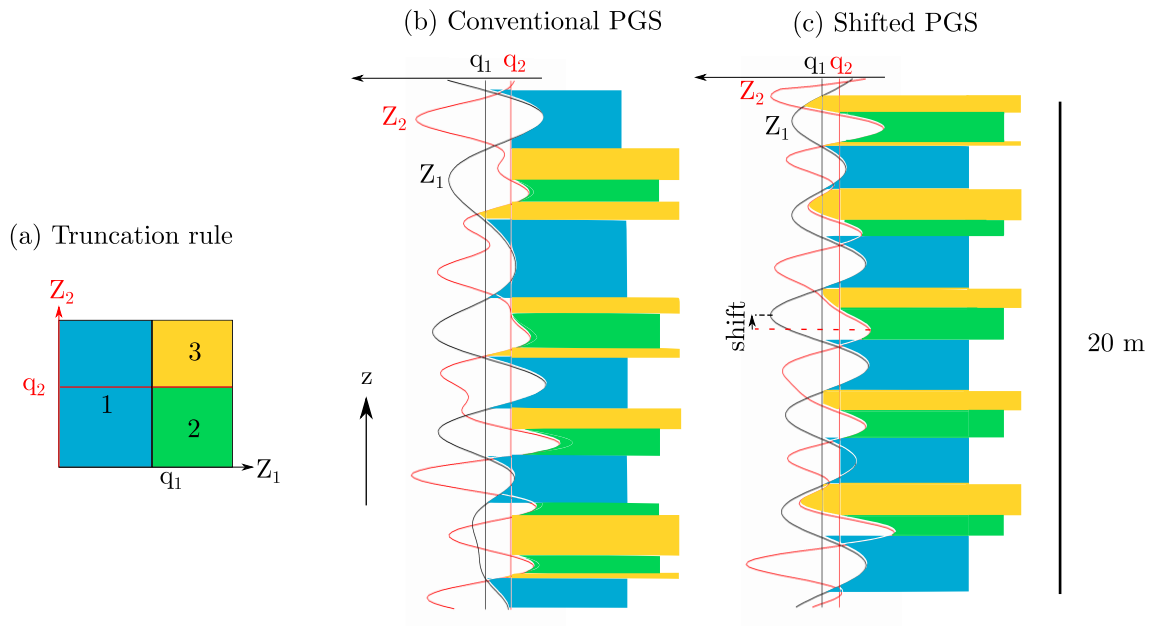


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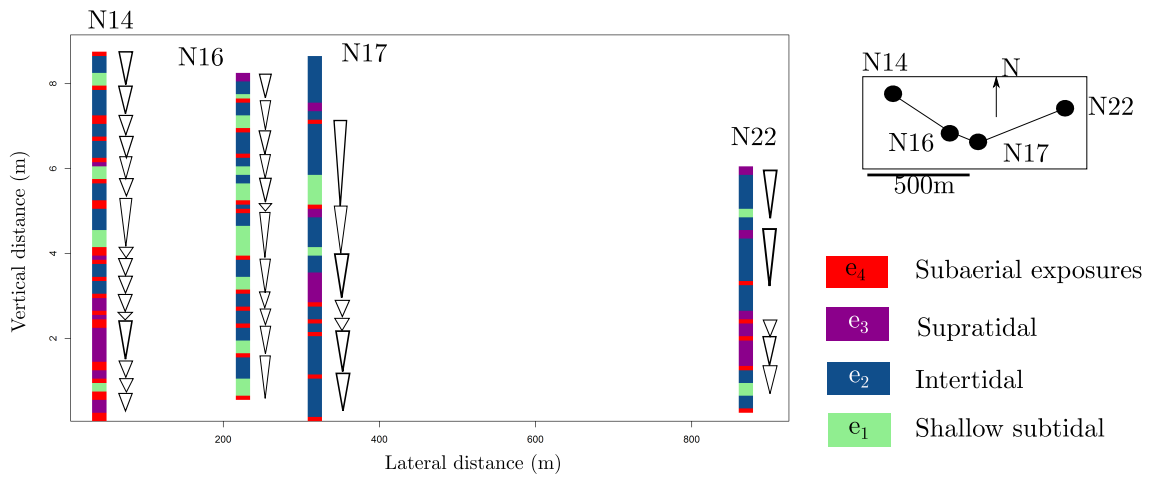


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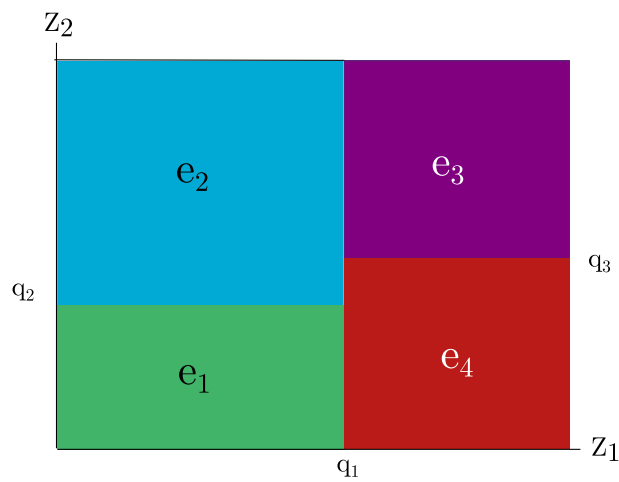


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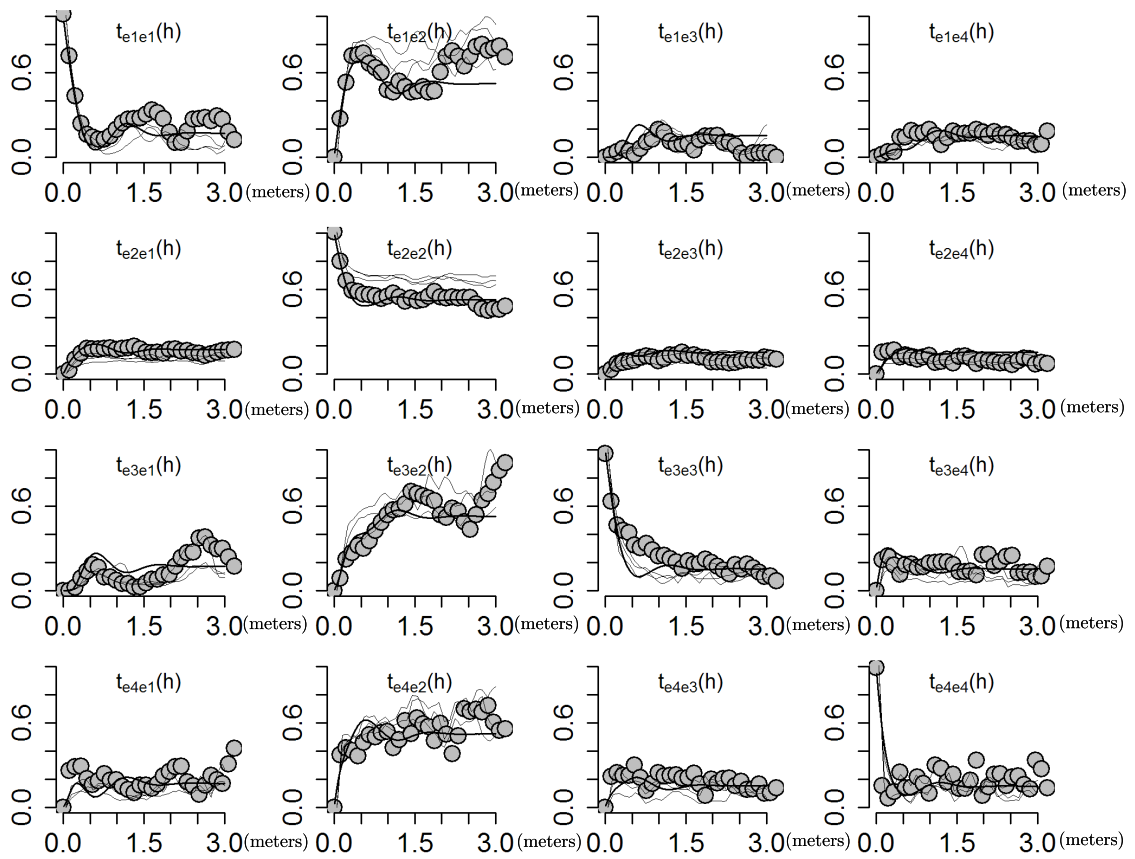


Figure 7:

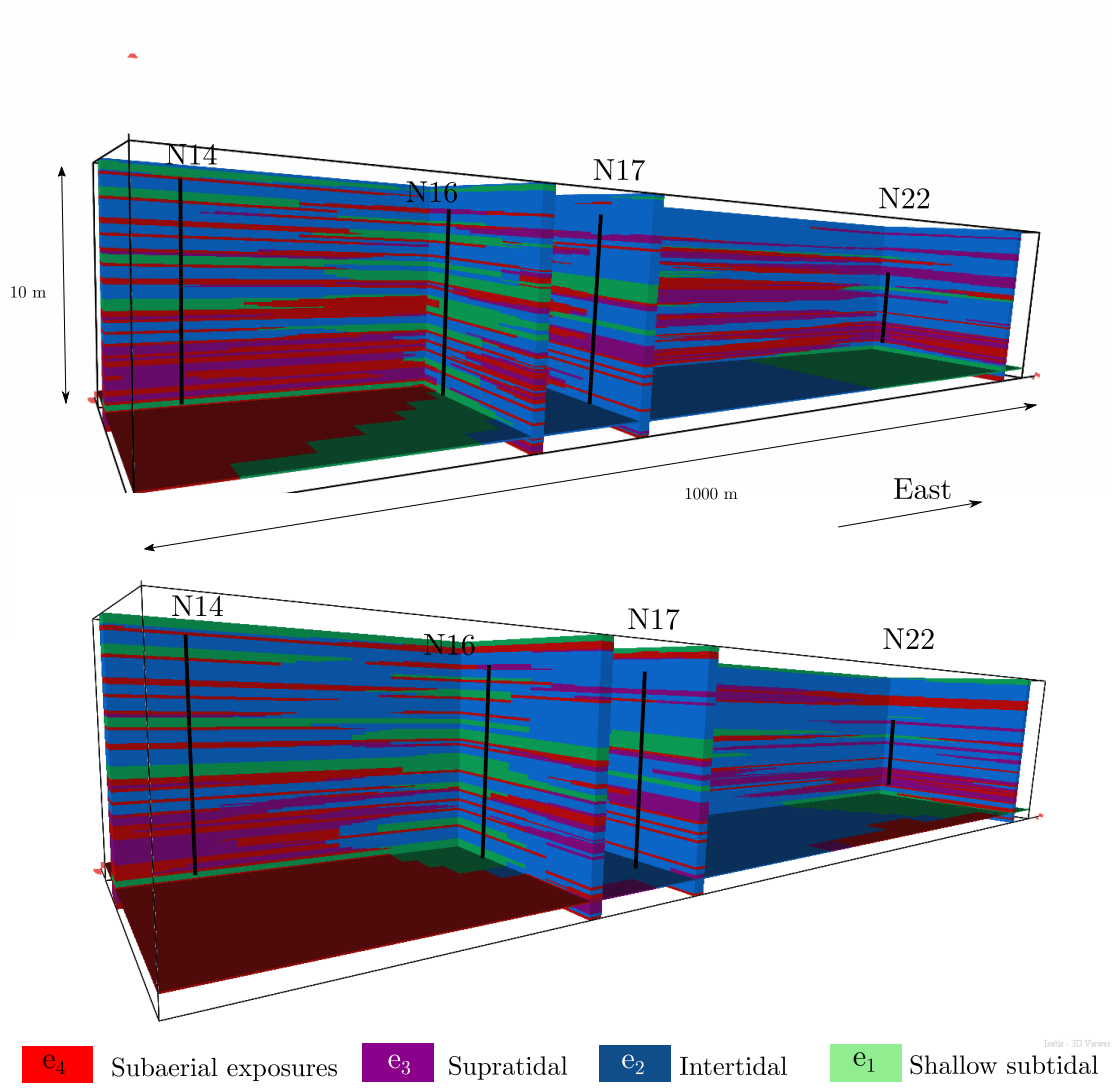


Figure 8:

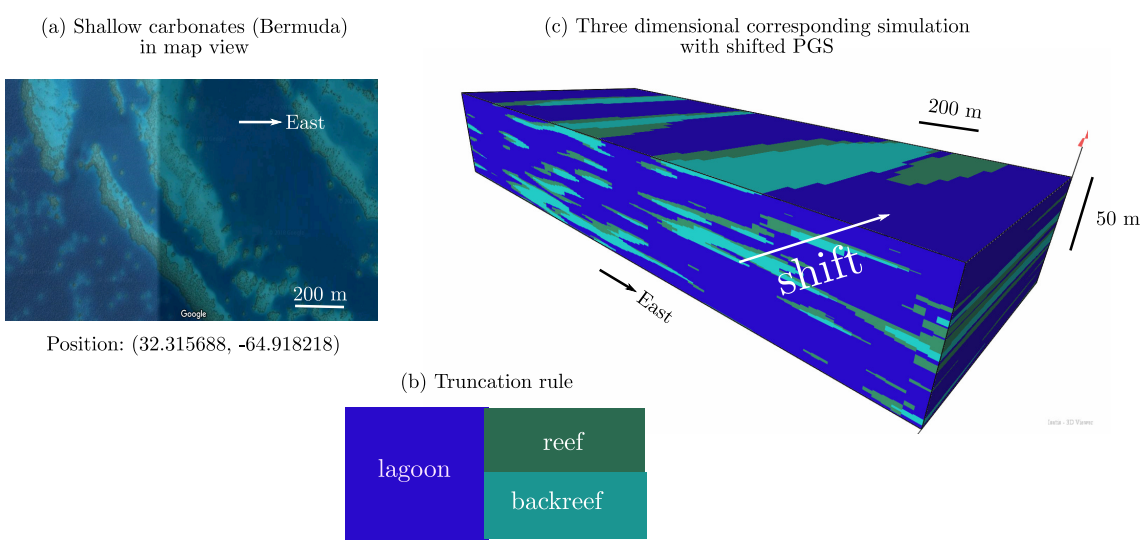


Figure 9:

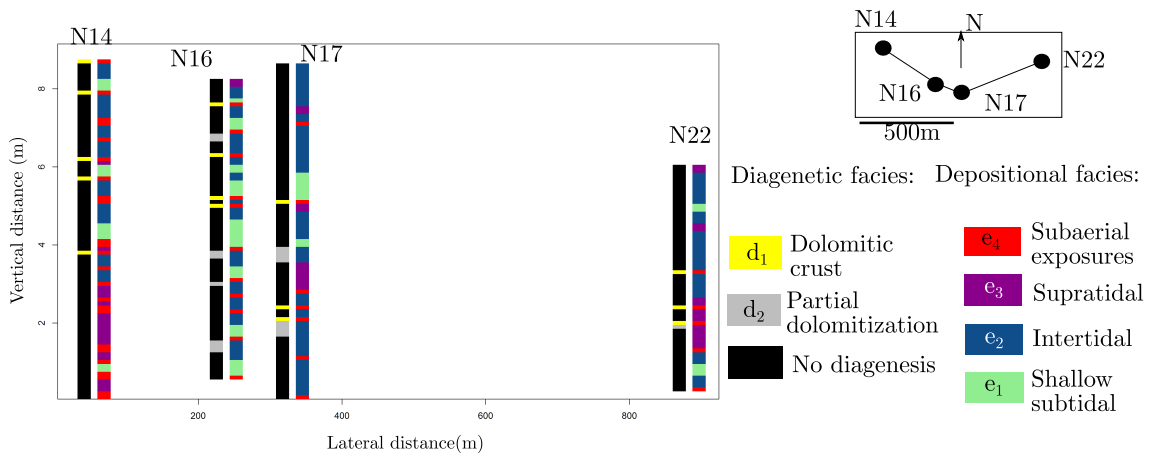


Figure 10:

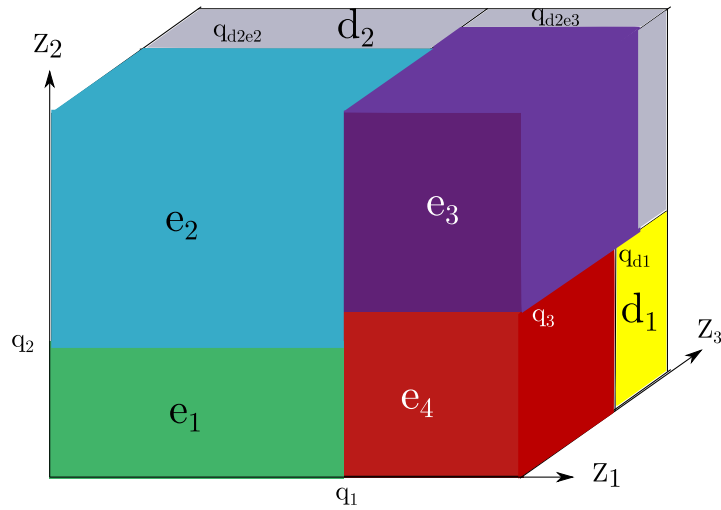


Figure 11:

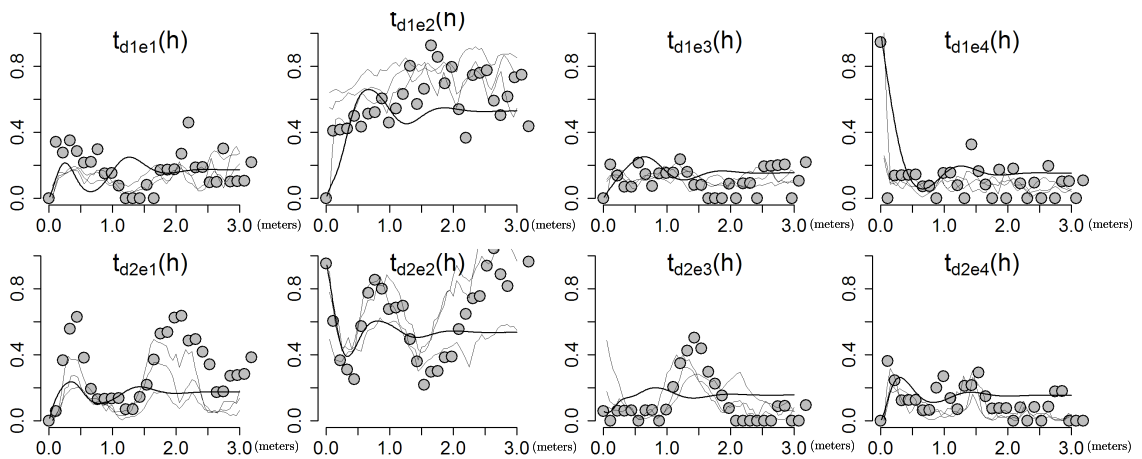


Figure 12:

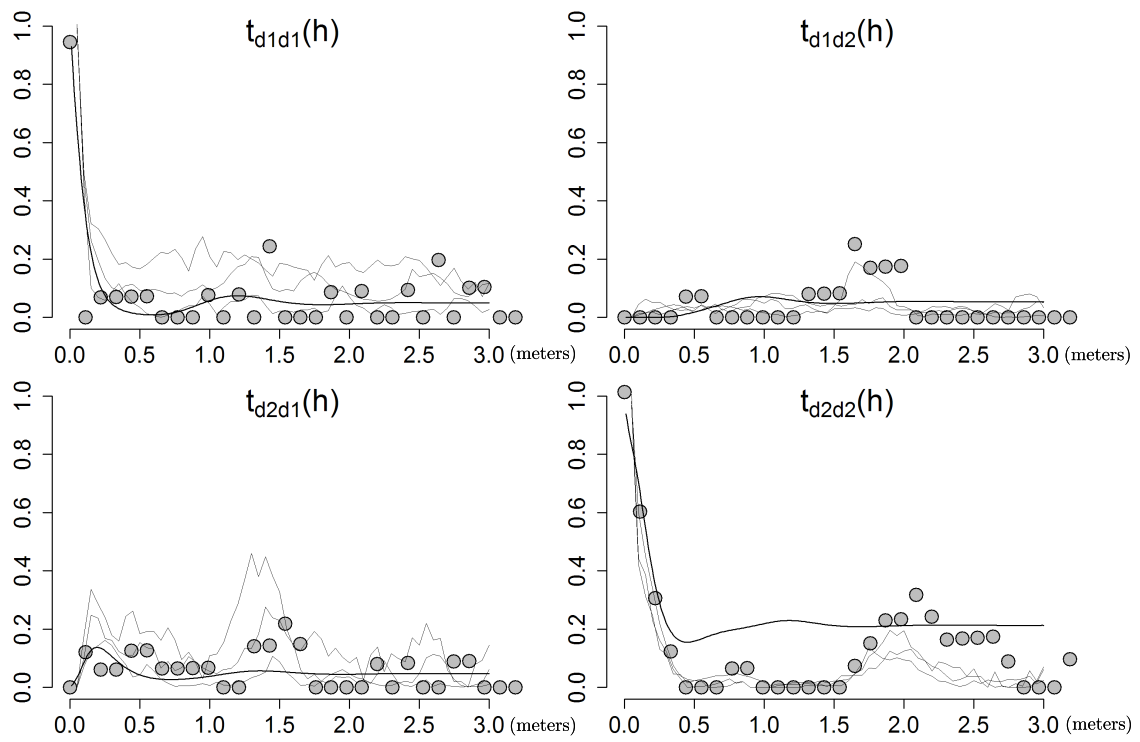


Figure 13:

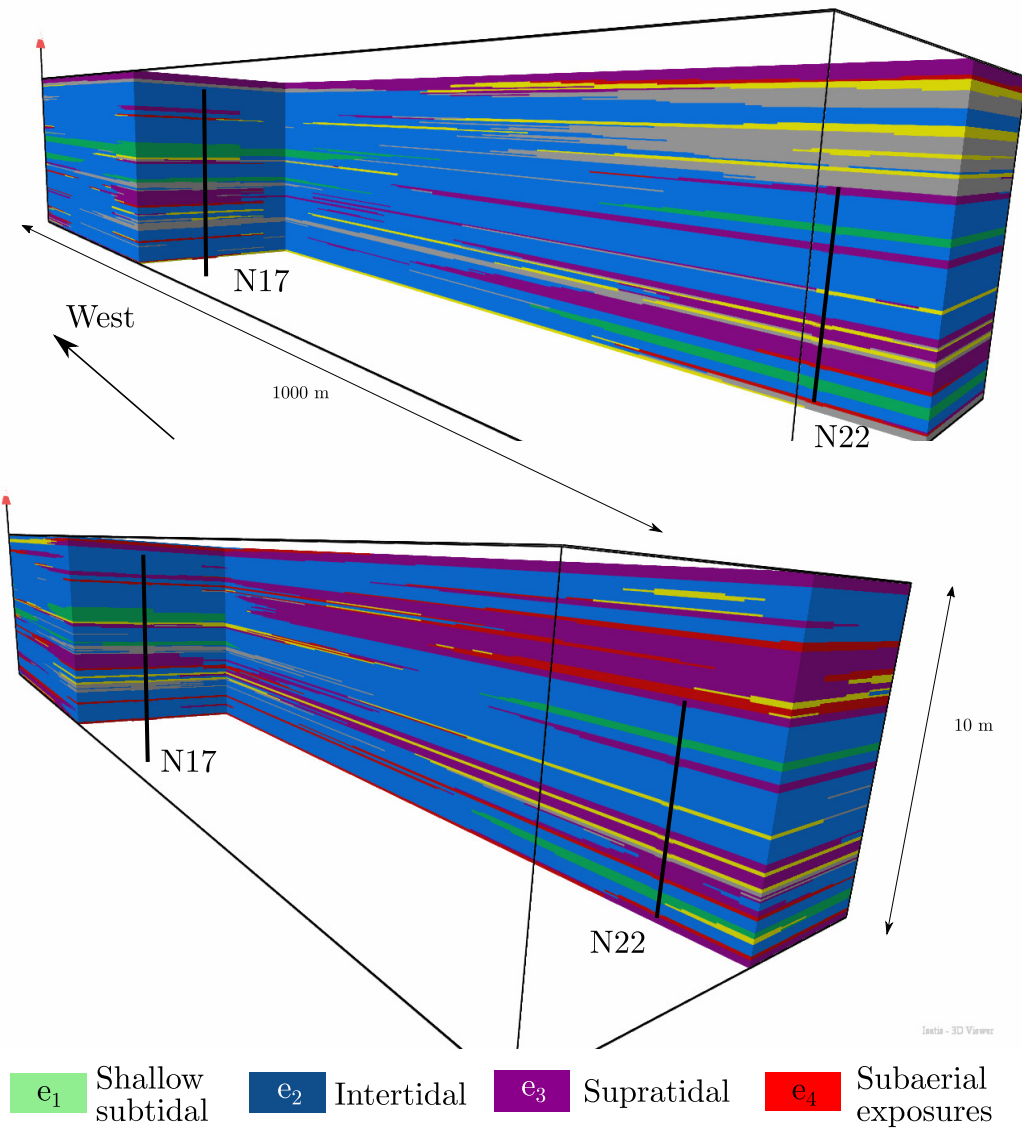


Figure 14:

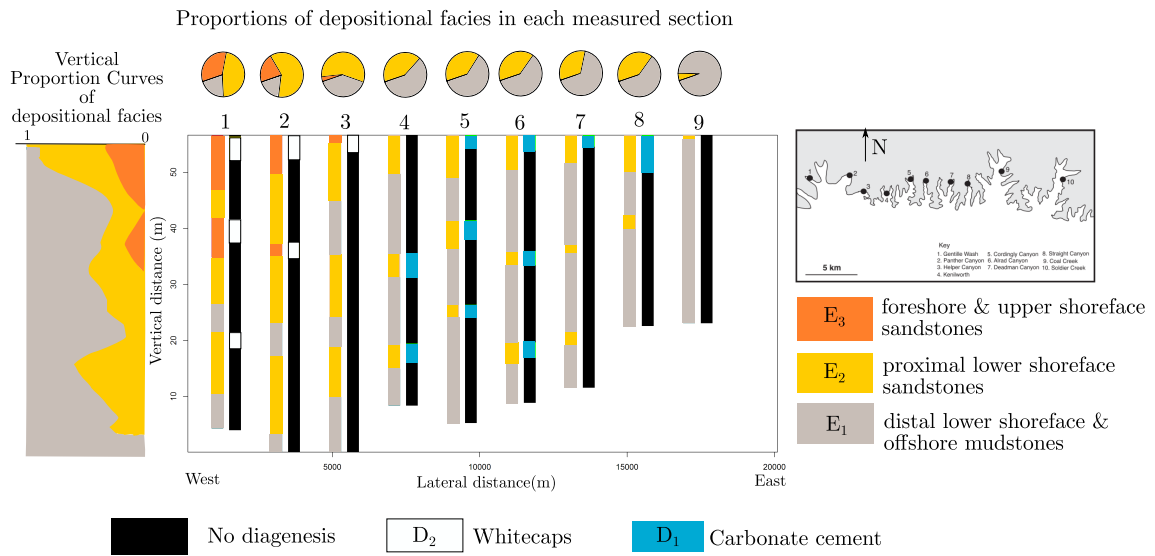


Figure 15:

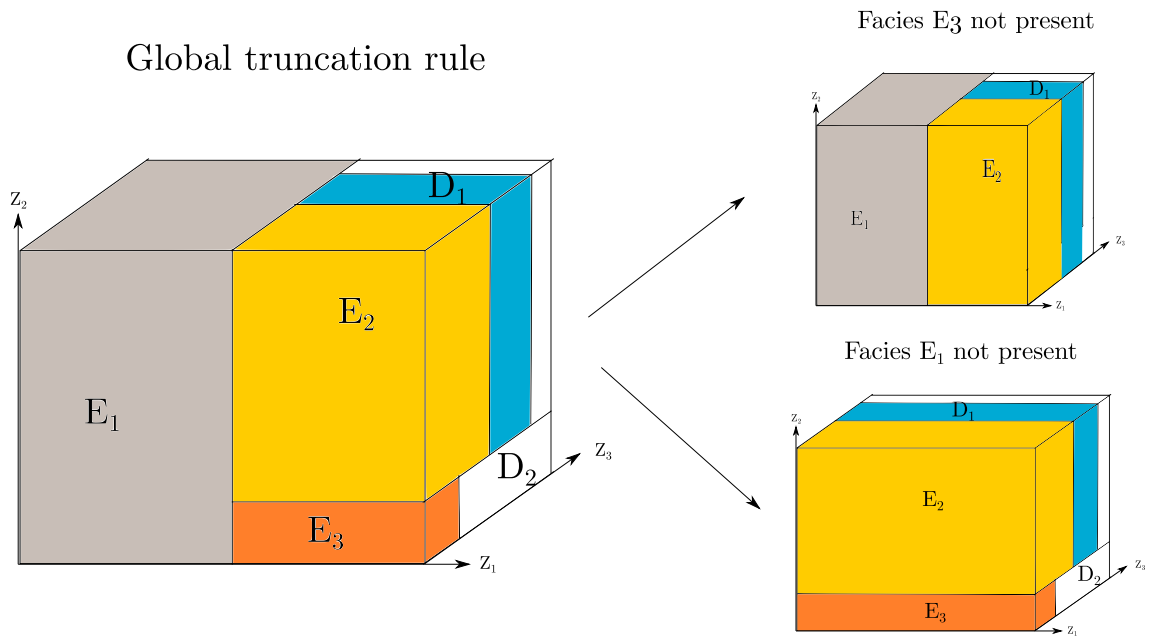


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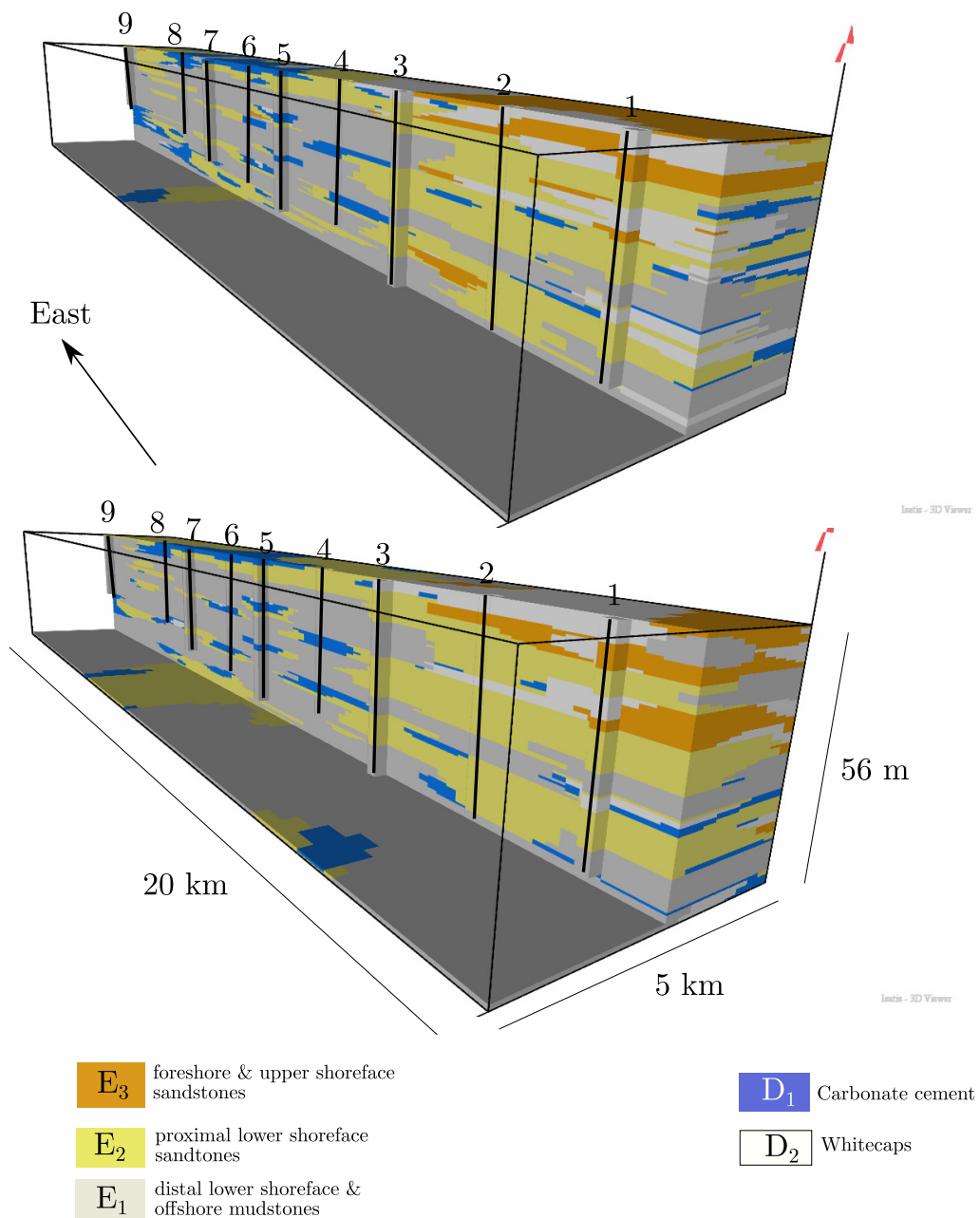


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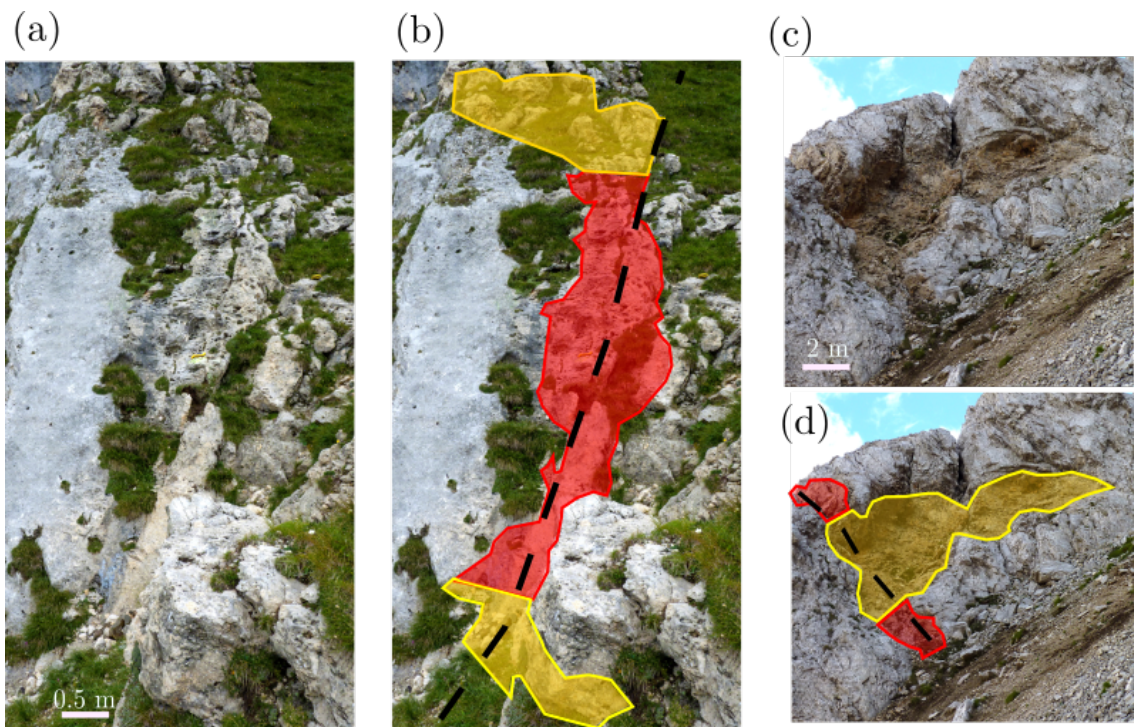


Figure 18:

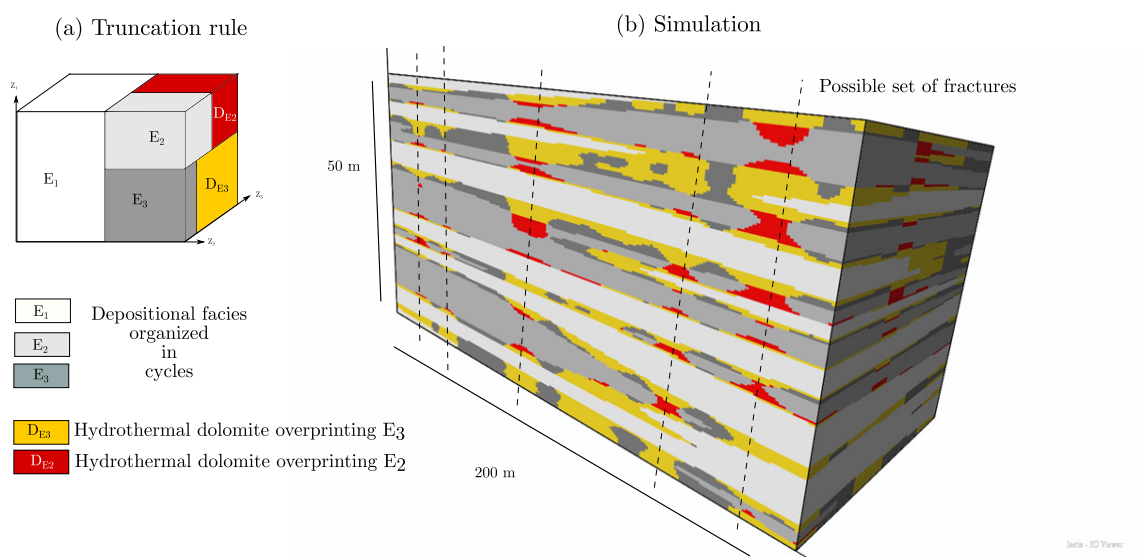


Figure 19: