

1 **Prediction of wave ripple characteristics using genetic programming**

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15

16 **Abstract**

17       We integrate published data sets of field and laboratory experiments of wave  
18 ripples and use genetic programming, a machine learning paradigm, in an attempt to  
19 develop a universal equilibrium predictor for ripple wavelength, height, and steepness.  
20 We train our genetic programming algorithm with data selected using a maximum  
21 dissimilarity selection routine. Thanks to this selection algorithm we use less data to train  
22 the genetic programming software, allowing more data to be used as testing (i.e. to  
23 compare our predictor vs. common prediction schemes). Our resulting predictor is  
24 smooth and physically meaningful, different from other machine learning derived results.  
25 Furthermore our predictor incorporates wave orbital ripples that were previously  
26 excluded from empirical prediction schemes, notably ripples in coarse sediment and long  
27 wavelength, low height ripples ('hummocks'). This new predictor shows ripple length to  
28 be a weakly nonlinear function of both bottom orbital excursion and grain size. Ripple  
29 height and steepness are both nonlinear functions of grain size and predicted ripple length  
30 (i.e. bottom orbital excursion and grain size). We test this new prediction scheme against  
31 common (and recent) predictors and the new predictors yield a lower normalized root  
32 mean squared error using the testing data. This study further demonstrates the  
33 applicability of machine learning techniques to successfully develop well performing  
34 predictors if data sets are large in size, extensive in scope, multidimensional, and  
35 nonlinear.

36

37 **1. Introduction**

38           Sufficiently strong water wave propagation over a moveable bed composed of sand  
39 grains results in the development rhythmic bedforms whose crest spacing is of the order  
40 of centimeters to meters while heights are of the order of centimeters. These features are  
41 often termed vortex ripples because of a recirculation cell that develops on the lee side of  
42 the bedform that is subsequently ejected upward during reversals in flow direction.  
43 Accurate prediction of vortex ripple size and shape is crucial for successful determination  
44 of seabed bottom roughness, a first order control on wave attenuation (e.g., Arduin et  
45 al., 2002), as well as sediment transport as suspended load (e.g., Green and Black, 1999;  
46 Bolaños et al., 2012). Furthermore ripple migration is a fundamental mechanism of  
47 bedload transport (e.g., Traykovski et al 1999; Becker et al., 2007), and parameterizations  
48 of bedload flux necessitate an accurate depiction of ripple size and shape.

49           Many predictors of equilibrium ripple geometry have been developed from field  
50 and laboratory datasets (e.g. Clifton, 1976; Nielsen, 1981; Grant and Madsen, 1982;  
51 Wiberg and Harris, 1994; Faraci and Foti, 2002; Styles and Glenn, 2002; Grasmeijer and  
52 Kleinhans, 2004; Soulsby and Whitehouse, 2005; Soulsby et al., 2012; Pedocchi and  
53 García, 2009a; Camenen, 2009). Equilibrium ripple size and shape is frequently broken  
54 down to include 3 subpopulations, a convention developed by Clifton (1976), and  
55 reviewed here in order of increasing hydrodynamic forcing. Orbital ripples are believed  
56 to scale linearly with wave orbital diameter at the seabed and display the largest steepness  
57 (ripple height/wavelength  $\sim 0.15$ ). Suborbital ripples show spacing that depends on wave  
58 orbital diameter and grain size. In even stronger hydrodynamic conditions anorbital  
59 ripples form, whose size is related to grain size alone and whose scaling is irrespective of  
60 wave orbital diameter. Suborbital ripples link the population of anorbital ripples with

61 those of orbital ripples.

62 As noted by Smith and Wiberg (2006), recent field and laboratory work has  
63 challenged the existing typology for wave-generated ripples as a result of the addition of  
64 two new populations (Figure 1). The first are ripples measured in fine sand under strong  
65 hydrodynamic conditions. Field and laboratory campaigns in more energetic conditions  
66 have discovered the presence of long wavelength, low amplitude ripples ('hummocks') in  
67 fine sands that scale with orbital diameter (e.g. Hanes et al., 2001, O'Donoghue et al.,  
68 2006). Predictors are unable to accurately capture this ripple size and shape (e.g. Bolaños  
69 et al., 2012), yet modeling (Chang and Hanes, 2004) and observation (Green and Black,  
70 1999; Cummings et al., 2009) of these bedforms show they eject vortices and are  
71 therefore important for their influence on seabed roughness and sediment transport.  
72 Furthermore at times these long wavelength ripples have superimposed anorbital ripples  
73 (e.g. Southard et al., 1990; Hanes et al., 2001; Williams et al., 2004), another unsolved  
74 problem in wave ripple prediction. Because of these complications, Pedocchi and García  
75 (2009a), who developed a recent well performing predictor, omit long wavelength ripples  
76 from their analysis, but note that these long wavelength 'round crested' ripples are  
77 observed above a critical threshold in  $U/w_s$  (where  $U$  is the maximum orbital velocity at  
78 the bed and  $w_s$  is the sediment fall velocity). Dumas et al. (2005) and Cummings et al.  
79 (2009) also show that the transition from anorbital scale ripples to round crested long  
80 wave orbital scale ripples is a function of orbital velocity (a set value for their given  
81 sediment mixtures).

82 The second new population of ripples are those found in medium to coarse sand  
83 (Traykovski et al., 1999; Arduin et al., 2002; Becker et al., 2007; Masselink et al.,



84 2007; Traykovski, 2007; Cummings et al., 2009; Yamaguchi and Sekiguchi, 2011).  
85 Coarse grained ripples have been observed in shelf environments for several decades  
86 (e.g., Forbes and Boyd, 1987; Leckie et al., 1988 and references therein) but until  
87 recently ripple measurements have not been coupled to the hydrodynamic parameters of  
88 their formation. Recent lab work by Cummings et al., (2009) demonstrated the  
89 persistence of steep ripples with orbital scaling in coarse sand under strong hydrodynamic  
90 conditions.

91         These two new populations of ripples highlight a perennial problem with empirical  
92 predictors; unless equations are built using large, integrated data sets that encompass  
93 many conditions, prediction schemes are difficult to translate to different settings. A non-  
94 empirical approach, such as models based on first principles (e.g., Foti and Blondeaux,  
95 1995; Blondeaux, 2001; Charru and Hinch, 2006), presents different problems: nonlinear,  
96 emergent processes that occur at the ripple scale such as flow separation, vortex ejection,  
97 turbulence, sediment suspension, pattern coarsening, defect creation, migration and  
98 annihilation (Werner and Kocurek, 1999), and the existence of multiple stable  
99 configurations in ripple sizes/shapes at a given hydrodynamic condition (a stability  
100 balloon; Hansen et al., 2001) limit the usefulness of finite-amplitude predictions.  
101 Prediction by numerical models of coupled fluid flow and bed evolution present  
102 promising results but have so far been tested under a narrow range of conditions and  
103 compared to few data sets (Marieu et al., 2008; Chou and Fringer, 2010).

104         If empirical data driven predictors are currently the most broadly applicable tools to  
105 develop field scale predictions, how should they be built? Traditionally the development  
106 of an empirical predictor relies on transforming a single (or several) noisy

107 multidimensional dataset to lower-dimensions and fitting a curve (with a set functional  
108 form) through the resultant point cloud. Here we offer a different solution: a data  
109 integration campaign (the collection of many published datasets) followed by machine  
110 learning (ML), whereby computational optimization techniques are used to find solutions  
111 to multidimensional and nonlinear problems. The suite of techniques encompassed by  
112 ML are essentially identical to empirical data driven techniques used previously except  
113 the trial and optimization of solutions is outsourced to a computer.

114         The most common ML paradigm used in coastal studies is artificial neural networks  
115 (ANN). Recent examples of its use include predictions of alongshore sediment transport  
116 in the surfzone (van Maanen et al., 2010), sand bar behavior (Pape et al., 2010) and  
117 suspended sediment reference concentration under waves (Oehler et al., 2012). Yan et al.  
118 (2008) used an artificial neural network to predict wave ripple geometry (length and  
119 height) based on three input parameters (median grain size, wave period, and the  
120 maximum near bed wave orbital velocity). ANN results give better predictions based on 3  
121 statistical measures (scatter index, correlation coefficient, and mean geometric deviation)  
122 than four common empirical models (Nielsen, 1981; Van Rijn, 1993; Wiberg and Harris,  
123 1994; Grasmeyer and Kleinhans, 2004). Yet the ANN ripple prediction scheme derived  
124 by Yan et al. (2008) was developed and compared to a limited dataset. Furthermore  
125 ANNs are problematic because the highly nonlinear result is difficult to interpret and  
126 does not offer immediate insight into the physical nature of the problem at hand. Decision  
127 or regression trees (e.g., Oehler et al., 2012), another common and well performing ML  
128 technique, is also hampered by the lack of direct physical significance and other  
129 drawbacks such as the lack of smoothness.

130 In this contribution we use genetic programming (GP; Koza, 1992), a population  
131 based optimization technique where the population consists of individual equations (i.e. a  
132 population of individual predictors). The mathematical or logical operations that  
133 constitute each algorithms can be modified at every time step via an ‘evolutionary’  
134 process (such as crossover and mutation) to produce expressions that optimize model-  
135 data fit. Outputs developed by GP can be smooth functions that are easy to examine and  
136 interpret for physical significance. Furthermore, *a priori* determination of the functional  
137 form of the predictor is not required and the final optimized solution can take on any  
138 mathematical form (within user defined limits). Thus far genetic programming has been  
139 applied to a wide range of problems including the prediction of freshwater phytoplankton  
140 dynamics (Whigam and Recknagel, 1999), downscaling of atmospheric model output  
141 (Coulibaly, 2004), determining appropriate parameterization for roughness in vegetated  
142 flows (Baptist et al., 2007), wave forecasting (Kambekar and Deo, 2012) and mapping of  
143 seafloor habitats (Silva and Tseng, 2008).

144 The goal of this study is to demonstrate the applicability of ML techniques  
145 (specifically GP) to research questions in the coastal domain. To accomplish this goal we  
146 compile 27 different field and laboratory data sets of wave ripple prediction (995  
147 individual measurements; Table 1) that span a broad range of conditions and develop a  
148 new wave-ripple predictor that is able to capture the morphology of ripple geometry in a  
149 wide range of forcing conditions, including conditions where long wave orbital ripples  
150 are present. We put our results in the context of existing formulations and theories, and  
151 assess the physical relevance of GP predictors. Our new equilibrium predictor ignores the  
152 effect of ripple orientation, time evolution, heterogeneous sediment, superimposed

153 current, ripple asymmetry, and bio-degradation of ripples. We discuss these limitations in  
154 the discussion section but note here that other existing time dependent ripple prediction  
155 schemes capture one or more (but not all) of these processes (i.e., Soulsby et al., 2012;  
156 Traykovski 2007). Finally, the compilation of published ripple data allows for the  
157 identification of gaps in knowledge and observations that should be pursued in future  
158 research. Future data collection campaigns can be added to this database, allowing for  
159 modifications to the prediction schemes shown below. In this sense the ripple prediction  
160 scheme we demonstrate here is dynamic.

161

## 162 **2. Data**

163 As a result of decades of study, many wave ripple datasets are available in the  
164 scientific literature. Examples of recent wave ripple data integration and compilations are  
165 Soulsby and Whitehouse (2005), Pedocchi and García (2009a) and Camenen (2009).  
166 Here we follow the lead of Pedocchi and García (2009a) and limit our data collection to  
167 studies using sediment with quartz (or near quartz) densities ( $2.65 \text{ g/cm}^3$ ) performed in  
168 large oscillatory tunnels, large wave flumes, wave racetracks and field conditions (i.e. we  
169 omit oscillating trays). Data on rolling-grain ripples, small bedforms that initially appear  
170 when flat beds are subject to oscillatory water motion, are ignored in this study because  
171 they have been experimentally shown to be a transient stage of ripple evolution (Faraci  
172 and Foti, 2001). We use 27 published studies in our dataset. Each measurement contains  
173 wave ripple, hydrodynamic, and sedimentological parameters. The dataset is split 59% /  
174 41% between laboratory and field conditions (Table 1), and laboratory measurements are  
175 obtained from a 49% / 49% / 2% split between oscillatory tunnels, wave flumes, and

176 wave racetracks. Measurement error is different for each data set in our database, a  
177 natural consequence of data integration campaigns that assemble data collected by  
178 different instruments and techniques. We assume that measurements of ripple data  
179 obtained in field settings are at or near equilibrium.

180 Our database can be visualized as a series of histograms showing the parameter  
181 range in our dataset (Figure 2). A majority of ripple measurements in our database occur  
182 at hydrodynamic conditions of  $d_0 < 2$  m,  $U < 0.75$  m/s and sedimentological conditions of  
183  $D_{50} < 0.5$  mm. Another notable attribute is the strong bimodal signature of ripple  
184 steepness centered at values of  $\sim 0.15$  and  $\sim 0.01$ . These clusters represent steep ripples  
185 and ‘hummocky’ ripples, respectively. We base our prediction of wave ripple wavelength  
186  $\lambda$  (m), ripple height  $\eta$  (m), and ripple steepness  $\vartheta$  ( $\eta/\lambda$ ; dimensionless) on four variables:  
187 wave period  $T$  (s), bottom orbital excursion  $d_0$  (m), median grain size  $D_{50}$  (m), and  
188 maximum near bed orbital velocity  $U$  (m/s). A hallmark of field data sets is the irregular  
189 forcing, requiring us to reconcile different measured parameters. Several field datasets  
190 used in the compiled dataset reported hydrodynamic parameters in terms of significant  
191 values ( $U_{sig}$ ,  $d_{0,sig}$ , and  $T_{sig}$ ). We followed the protocol of Pedocchi and García (2009) and  
192 assume  $U = U_{sig}$  (and furthermore  $d_0 = d_{0,sig}$  and  $T = T_{sig}$ ). We acknowledge that the  
193 merging of disparate data sources introduces uncertainty into the data.

194 The hydrodynamic and sedimentological conditions covered by this dataset can be  
195 visualized using 6 projections of the 4 dimensional phase space (Figure 3). Notable  
196 sparseness occurs in this database at strong hydrodynamic conditions, and at median  
197 seabed grain sizes above 0.5 mm. We use  $T$ ,  $d_0$ , and  $U$  as separate independent variables  
198 for input to the GP (though they are related by  $d_0 = UT/\pi$ ) in an attempt to introduce no

199 additional information about which of these parameters is most relevant. As GP is a data  
200 driven technique, the raw hydrodynamic data is given as input and the ML process  
201 determines which hydrodynamic variable(s) is most relevant from a statistical standpoint.  
202 We use  $T$ ,  $d_0$ ,  $D_{50}$ , and  $U$  to predict  $\lambda$ . Predicted  $\lambda$  is incorporated into the suite of  
203 variables (i.e.,  $T$ ,  $d_0$ ,  $D_{50}$ ,  $U$ ) used to predict  $\eta$ . We combining the predictors for  $\lambda$  and  $\eta$   
204 enable the development of a predictor for ripple steepness. Yet we do not enforce the  
205 accurate depiction of steepness in the development of ripple height and length predictors  
206 and imprecision in the  $\lambda$  and  $\eta$  equations may cause imprecision in the prediction of  $\vartheta$ .  
207 However, in some circumstances accurate depiction of height and steepness is required  
208 for the parameterization of relevant processes (e.g. vertical suspended sediment  
209 diffusivity; Nielsen, 1992): therefore we also develop an independent ripple steepness  
210 predictor using the genetic programming technique. The development of a third predictor  
211 also further demonstrates the strengths and weaknesses of GP and ML techniques.  
212 Predicted  $\lambda$  and  $\eta$  are added to the variables ( $T$ ,  $d_0$ ,  $D_{50}$ ,  $U$ ) used to predict  $\vartheta$ . The  
213 development of predictors for  $\lambda$ ,  $\eta$  and  $\vartheta$  without enforcing interoperability relies on users  
214 to decide which predictors are most important for the specific research question.

215         Several published studies measure two superimposed ripple scales (larger orbital  
216 scale ripples and smaller anorbital scale ripples) at a single hydrodynamic condition (e.g.  
217 Hanes et al., 2001; Pedocchi and García 2009b; Cummings et al., 2009). Work by  
218 Cummings et al., (2009) shows that both pattern modes occur as maximum orbital  
219 velocity is increased and the ripple pattern transitions from small scale (anorbital) ripples  
220 to large scale orbital ripples ('hummocks'). Upon further velocity increase, the small  
221 scale ripples are destroyed and only the large scale orbital features remain (Cummings et

222 al., 2009). The threshold of large scale orbital ripple appearance can be estimated from  
223 the work of Pedocchi and García (2009a) who found that large scale features appear at a  
224 threshold value of  $U/w_s \cong 25$ . When both anorbital and large scale ripples are present in  
225 tabulated data (e.g. Hanes et al 2001) we only include large scale ripples: the scaling of  
226 long wave ripples with bottom orbital diameter suggests a physical relationship to small  
227 scale orbital and suborbital ripples. In contrast, anorbital ripples scale with grain size  
228 (similar to current ripples) and the mechanism responsible for their formation may be  
229 different (Wiberg and Harris, 1994). We remove small-scale (anorbital) ripples from our  
230 database if they are present at values of  $U_0/w_s \geq 25$ ; laboratory work by Cummings et al.,  
231 (2009) and Pedocchi and García (2009a) suggests that this regime is dominated by large  
232 scale ripples. The targeted collection of field and laboratory data is needed to refine this  
233 threshold.

234

### 235 **3. Methods**

#### 236 *3.1 Selection of training, validation, and testing data*

237 The database is split into three subsets to be used as training, validation, and  
238 testing. The GP algorithm uses the training dataset to develop and optimize candidate  
239 solutions. The validation dataset is used to evaluate the fitness of GP derived solutions  
240 and define which predictors persist. Testing data is not used or seen by the GP algorithm  
241 and is instead reserved as an independent test of the final predictors (and other published  
242 predictors). In the genetic programming literature there remains no proven 'best practice'  
243 for percentage of training, validation, and testing data, nor a well defined method of  
244 splitting these datasets. This may be because data splitting (e.g., the retention of a testing

245 dataset) is not addressed in the foundational literature of the technique (as noted by  
246 Kushchu, 2002). Yet because our database of ripple measurements contains only sparse  
247 data at energetic hydrodynamic conditions and large grain sizes, the selection and  
248 partitioning of data into these three categories is crucial to develop a well performing  
249 predictor applicable to a range of environments (Bowden et al., 2002). For example,  
250 random division of the data has the potential to produce a significant problem; the  
251 training data is likely to misrepresent the full phase space of the entire dataset (i.e.  
252 exclude coarse grained and/or strong hydrodynamic data).

253 Informed data selection (i.e., selection based on clustering) has been shown to  
254 produce better results with ML predictors than ‘blind’ or random data selection (e.g.,  
255 Bowden et al., 2002; May et al 2010). In this study we select training data through the use  
256 of a maximum dissimilarity algorithm (MDA; e.g., Camus et al., 2011). This algorithm is  
257 not a clustering routine (where cluster centroids are selected to represent a representative  
258 value of the data in the cluster), but instead a selection routine (where a centroid  
259 represents the most dissimilar data point from the previous centroids; Camus et al., 2011).  
260 Though our selection technique is different than the clustering techniques used by  
261 Bowden et al. (2002), our approach leads to a similar result: the use of a minimum of  
262 training data that is able to capture the variance in hydrodynamic and sedimentological  
263 conditions of the entire dataset while leaving more data to be used as validation and  
264 testing.

265 Our implemented version of the maximum dissimilarity algorithm is based on the  
266 description provided in Camus et al. (2011). Selection starts with the normalization of the  
267 data to a value between 0 and 1:



$$X_n = \frac{X - X_{min}}{X_{max} - X_{min}} \quad (1)$$

268 where  $X_n$  is the new normalized data value (between 0 and 1),  $X$  is the original value,  
269  $X_{min}$  and  $X_{max}$  are the minimum and maximum of all values of variable  $X$ , respectively.  
270 After this normalization a single data point, a ‘seed’, is selected as the first centroid.  
271 Since our dataset is typified by sparseness in the coarse grain data, we use the largest  
272 grain size measurements as the first centroid (the ‘seed’). The user selects the number of  
273 centroids and the algorithm then selects the additional centroids through an iterative  
274 process: Each data point in our data set is a 4-dimensional vector (normalized  $T$ ,  $U$ ,  $d_0$ ,  
275  $D_{50}$  space) and is associated with a distance to the nearest centroid. The single data point  
276 with the maximum distance between itself and the nearest centroid is selected as the next  
277 centroid (Camus et al., 2011). This routine continues until the user defined number of  
278 centroids is reached, after which data is denormalized.

279       There remains significant ambiguity in determining the appropriate number of  
280 centroids (or clusters) needed to accurately represent data, especially continuous data  
281 (e.g. May et al 2010). Our dataset on wave ripples is multidimensional and relatively  
282 continuous (i.e. not naturally clustered). Furthermore the dataset is sparse in areas  
283 because of a lack of collected data, while densely populated with measurements in other  
284 regions of phase space (e.g., experimental campaigns at specific hydrodynamic and/or  
285 sedimentological conditions). Since we intend to use selected centroids as representatives  
286 of the entire dataset, selecting too many centroids will likely rob the validation and  
287 testing datasets of poorly represented data (e.g., large  $T$ ,  $U$ ,  $d_0$ ,  $D_{50}$ ) while too few  
288 centroids will leave the testing data with too few data to capture the variability in the  
289 dataset. We use 30 centroids for the prediction of  $\lambda$  and 40 centroids for the prediction of

290  $\eta$ . Centroids used to represent  $\eta$  are also used for analysis of  $\vartheta$ . Centroid locations can be  
291 seen in Figure 4. The use of fewer centroids (10-20) produced too few predictors while  
292 more centroids ( $\sim 100$ ) tended to produce many more nonlinear and potentially overfit  
293 solutions. In addition, the solutions obtained with more centroids were qualitatively  
294 similar to the solutions presented below using only 30-40 centroids. More centroids are  
295 used to predict ripple height because  $\eta$  is more difficult to predict (see also Yan et al.,  
296 2008; Williams et al., 2004). This is likely a result of the nonlinearities associated with  
297 ripple crests protruding into regions of flow with higher velocities: ripple height is likely  
298 more strongly influenced by suspension processes as a result. Data selected as the  
299 centroid locations are used for the training data. The points not selected as centroids (i.e.  
300 not selected as training data) are used for validation and testing data. Data is split  
301 between validation and testing randomly, without using a selection routine. Therefore the  
302 breakdown for the  $\lambda$  datasets is  $\sim 3\%$  training,  $\sim 48\%$  validation,  $\sim 48\%$  testing, while the  $\eta$   
303 (and  $\vartheta$ ) dataset breakdown is  $\sim 5\%$  training,  $\sim 47\%$  validation,  $\sim 47\%$  testing.

304

### 305 *3.2 Genetic programming*

306 We operate on this compiled ripple data using the evolutionary computation  
307 technique of genetic programming (GP), a ML paradigm whereby candidate solutions (in  
308 the form of randomly generated equations) are evaluated and subsequently modified  
309 (Koza, 1992; Poli et al., 2008). The modification of candidate solutions is manifest as  
310 changes in variables and mathematical relationships between variables (i.e. the  
311 mathematical form), hence the description of this style of problem as ‘symbolic  
312 regression’. Variables used in this study to predict wave ripple geometry are  $T$ ,  $U$ ,  $d_0$ ,

313  $D_{50}$ ,  $\lambda$  (for height and steepness prediction),  $\eta$  (for steepness prediction), as well as GP  
314 derived constants. Nondimensional, renormalized input (from 0-1) is not necessary with  
315 GP (as it is with other ML techniques), and input is fed into the algorithm with units.  
316 Only  $D_{50}$  is renormalized in this analysis, and fed into the GP in units of mm (as opposed  
317 to m, but the presentation of all results in this contribution are in meters). Mathematical  
318 operators used in this study are + (addition), - (subtraction),  $\times$  (multiplication),  $\div$   
319 (division),  $\sqrt{\quad}$  (square root), as well as integer powers (e.g.  $x^2, x^3, x^4$ ). Furthermore we  
320 omit logical functions (e.g. if-then-else) because of the lack of smoothness when  
321 incorporating these components.

322 Candidate solutions are evaluated based a 'fitness function', a user defined error  
323 metric that determines how well a given candidate fits the validation data. Mean squared  
324 error (MSE) is used as the fitness function:

$$MSE = \frac{(p - b)^2}{n} \quad (2)$$

325 where MSE is the Mean Squared Error, n is the sample size, p are the predicted values,  
326 and b are the observed values. The correlation coefficient, one of the error metrics used in  
327 previous ripple studies (Yan et al., 2008), was not used as a fitness function because it  
328 tended to develop nonphysical predictors (negative wavelengths and heights under certain  
329 conditions) that matched the shape of the data but did not align well with actual  
330 magnitudes.

331 Equations that minimize mean squared error are retained, while poor performing  
332 solutions are discarded. Retained solutions are combined, rearranged and manipulated in  
333 a probabilistic manner according to evolutionary processes; solutions 'crossover' by  
334 combining elements of other solutions to develop a new solution and 'mutations' develop

335 new mathematical expression to substitute or tack on to a previous solution. As an  
336 example, candidate solutions are commonly encoded in GP software as ‘trees’, and the  
337 modification of candidate solutions (change of variables and/or mathematical expression)  
338 is accomplished through adjustments in tree ‘limbs’ (Figure 5). Through time predictors  
339 gain complexity (i.e. trees grow in size) as they are recombined in a variety of ways,  
340 moving from simple equations (e.g. two variables and one mathematical symbol linking  
341 them) to highly nonlinear, complex expressions (e.g. many variables linked by many  
342 symbols). In this way the growth and adjustment of candidate solutions enables the  
343 searching of an increasingly larger phase space (i.e. variable and symbolic space), and  
344 find optimized solutions to the problem at hand. This search process occurs until a  
345 solution with zero error is found or the routine is terminated.

346 In this study we use a proven symbolic regression/genetic programming software  
347 package developed by Schmidt and Lipson (2009; 2013). This software package,  
348 ‘Eureqa’, modifies the tree-based encoding outlined above by eliminating redundancy  
349 when multiple ‘tree limbs’ are identical. The software output is a suite of solutions with  
350 increasing mathematical ‘complexity’, where complexity is a count of the numbers of  
351 operations and variables used in the candidate solution. Each solution of a given  
352 complexity represents the equation with the least error compared to identically ‘complex’  
353 candidate solutions. Furthermore, to be retained in the solution set, a given solutions must  
354 have less error compared to all previous less-complex solutions. Therefore the suite of  
355 solutions that is developed as output lie along the ‘Pareto front’, a line in complexity-  
356 fitness space that illustrates fitness increases with the increasing complexity of candidate  
357 solutions. Because simple predictors are retained though more complex predictors may fit

358 the data with less error, the user must pick a single solution as the final predictor of  
359 choice.

360

### 361 *3.3 Generalization and overfitting*

362 The lack of a single optimal solution as output from the GP algorithm is likely a  
363 consequence of using noisy data (e.g., field data) and examining a phenomena that may  
364 not have a single solution, but instead a small range of possible solutions (i.e., there may  
365 be multiple stable ripple configurations for a given hydrodynamic/sedimentological  
366 condition, a ‘stability balloon’; Hansen et al., 2001). The determination of an ideal  
367 solution from the GP program was further complicated because there is no stoppage  
368 routine built into the algorithm (e.g., based on fitness) used in this study. We cease the  
369 search after roughly  $10^{10}$  formulas have been evaluated as continued search shows only  
370 marginal increases in predictive power (and this increase occurs only on more complex,  
371 likely overfit, predictors). The solutions were then evaluated to determine the most  
372 appropriate final predictor. Several methods for eliminating overfit solutions exist (e.g.,  
373 Gonçalves et al., 2012). We use several techniques in parallel to determine appropriate  
374 solutions: 1) bias toward shorter, physically reasonable solutions, 2) examining ‘cliffs’ in  
375 the Pareto front, and 3) examination of solution fit.

376 Many of the more complex solutions have lower error with training and validation  
377 data but are physically uninterpretable. Therefore when evaluating the family of solutions  
378 from a given genetic programming iteration we tend to bias our search for the most  
379 universal predictor by preferring compact solutions because they tend to offer more  
380 generalization and are likely less overfit (The minimum description length principle; e.g.,

381 O'Neill et al., 2010). Shorter solutions reappear with repeat initialization of the genetic  
382 programming algorithm, suggesting that these represent the globally optimum solutions  
383 for a given function size. Longer solutions do not tend to reappear, either a result of a  
384 large search space that is not repeated during repeat initializations or the presence of  
385 multiple, equally optimal solutions in the large phase space (i.e. local minima). The  
386 inherent reproducibility of simple, weakly nonlinear solutions suggests their use as  
387 predictors until further data can be used to justify the use of highly nonlinear predictors.

388       Aside from examining the solutions from least complex to most complex,  
389 examining areas along the Pareto front where large gains in prediction are obtained with  
390 small gains in solution complexity is a natural place to observe potential solutions (Figure  
391 6). These areas along the Pareto front are referred to as 'cliffs'. Schmidt and Lipson  
392 (2009) used the last of such 'cliffs' to observe many physically relevant solutions. In this  
393 study final solutions were chosen from the subset of solutions that are 'cliffs' along the  
394 Pareto front

395       Candidate solutions are evaluated by minimizing error functions. Occasionally  
396 candidate solutions are able to minimize the mean squared error but provide unphysical  
397 solutions (e.g. negative ripple wavelengths under some conditions) or generally poor  
398 global performance (e.g. flat, constant predictors). These solutions must be manually  
399 disregarded, as there is as yet no means of excluding them.

400

#### 401 *3.4 Comparison with other predictors*

402       Predictor performance is evaluated, using the independent testing data, with the  
403 Normalized Root Mean Squared Error (NRMSE):

$$NRMSE = \frac{\sqrt{MSE}}{\bar{b}} \quad (3)$$

404 where  $\bar{b}$  is the mean of the observed values. Additionally we report correlation  
 405 coefficient (Pearson's  $r$ ) for each predictor evaluated against the independent testing data.

406 We compare our results to two recently developed and widely used predictors:  
 407 Soulsby and Whitehouse (2005; also reported in Soulsby et al., 2012) and Pedocchi and  
 408 García (2009a). As noted by Soulsby et al., (2012), recent work by Camenen (2009)  
 409 using a large compiled database of ripple measurements found the Soulsby and  
 410 Whitehouse (2005) formulation to be the best overall predictor compared to those  
 411 developed by Grant and Madsen (1982), Wikramanayake and Madsen (1991), Van Rijn  
 412 (1993), Mogridge et al., (1994), Wiberg and Harris (1994), and Grasmeijer and Kleinhans  
 413 (2004). The recent work of Pedocchi and García (2009a), which was not evaluated by  
 414 Camenen (2009), yields good collapse of the data compared to other the predictors  
 415 mentioned above and performs well in field conditions (Bolaños et al., 2012). The  
 416 Soulsby and Whitehouse (2005) predictor for length and steepness ( $\eta/\lambda$ ) is:

$$\frac{\lambda}{A} = \left[ 1 + (1.87 \times 10^{-3}) \frac{A}{D_{50}} \left( 1 - e^{\left\{ - \left( 2.0 \times 10^{-4} \frac{A}{D_{50}} \right)^{1.5} \right\}} \right) \right]^{-1} \quad (4)$$

$$\frac{\eta}{\lambda} = 0.15 \left[ 1 - e^{\left\{ - \left( 5000 \frac{D_{50}}{A} \right)^{3.5} \right\}} \right] \quad (5)$$

417 where  $A$  is the wave orbital amplitude ( $2A=d_0$ ). Combining (4) and (5) yields  $\eta$  alone.

418 The Pedocchi and García (2009a) predictor is:

$$\frac{\lambda}{d_0} = \begin{cases} 0.65 \left[ \left( 0.050 \frac{U}{w_s} \right)^2 + 1 \right]^{-1}, & Re_p \geq 13 \\ 0.65 \left[ \left( 0.040 \frac{U}{w_s} \right)^2 + 1 \right]^{-1}, & 9 \leq Re_p < 13 \\ 0.65 \left[ \left( 0.054 \frac{U}{w_s} \right)^3 + 1 \right]^{-1}, & Re_p < 9 \end{cases} \quad (6)$$

$$\frac{\eta}{d_0} = \begin{cases} 0.1 \left[ \left( 0.055 \frac{U}{w_s} \right)^3 + 1 \right]^{-1}, & Re_p \geq 13 \\ 0.1 \left[ \left( 0.055 \frac{U}{w_s} \right)^4 + 1 \right]^{-1}, & 9 \leq Re_p < 13 \\ 0.1 \left[ \left( 0.055 \frac{U}{w_s} \right)^5 + 1 \right]^{-1}, & Re_p < 9 \end{cases} \quad (7)$$

419

420 where  $w_s$  is evaluated for  $D_{50}$  and  $Re_p$  is a dimensionless particle size (Pedocchi and

421 García, 2009a) evaluated as:

$$Re_p = \frac{\sqrt{gRD_{50}D_{50}}}{\nu} \quad (8)$$

422 where  $g$  is gravity,  $R$  is the submerged specific density of sediment (here taken to be

423 1.65) and  $\nu$  is kinematic viscosity. The three size classes ( $Re_p \geq 13$ ,  $9 \leq Re_p <$

424  $13$  and  $Re_p < 9$ ) correspond to coarse, medium and fine sand respectively and the three

425 separate equations result in slight discontinuities.

426 Lastly we note that we are unable to compare the performance of our GP derived

427 predictor to the ANN model developed by Yan et al., (2008) as we do not know the final

428 optimized ANN equation developed by Yan et al., (2008). In addition, we do not know

429 which data was used as training/validation or testing in the development of the ANN

430 model.



431

## 432 4. Results

### 433 4.1 Ripple wavelength

434 The GP algorithm output is shown in Table 2. This experiment evaluated  $10^{10}$   
435 formulas to develop the Pareto front shown in Figure 6. Cliffs, significant gains in error  
436 for small changes in equation complexity occur along the Pareto front at complexities of  
437 3, 6, and 8 (Figure 6) The first of these cliffs (at complexity 3) is a predictor,  $\lambda = 0.607d_0$ ,  
438 that mimics the basic form of the orbital scale (i.e. weak hydrodynamics) predictor  
439 commonly used today, where ripple wavelength is a linear function of orbital excursion  
440 (e.g.  $\lambda = 0.65d_0$  from Miller and Komar, 1980a;  $\lambda = 0.62d_0$  from Wiberg and Harris, 1994).  
441 Debate surrounds the correct value of the coefficient modifying orbital excursion,  
442 especially in medium to coarse sand (e.g. Becker et al., 2007; Traykovski et al., 1999).  
443 All solutions that are more complex than the solution of complexity 3 demonstrate why  
444 there is debate: the coefficient is likely a function of grain size. We rule out solution 3 as  
445 a viable universal predictor because grain size is a control on ripple length (e.g.,  
446 Cummings et al., 2009). We focus our remaining examination on the solution at  
447 complexity 8.

$$\lambda = \frac{d_0}{1.12 + 2.18(1000D_{50})} \quad (9)$$

448 Figure 7 shows the general behavior of this predictor: increasing wave ripple  
449 spacing with increasing bottom orbital excursion and decreasing wave ripple wavelength  
450 with increasing grain size. Furthermore ripple length is more sensitive to median grain  
451 size at larger orbital diameter. Previous ripple length prediction schemes have focused on  
452 orbital diameter and grain size as they represent the two fundamental length scales in the

453 development of oscillatory bedform. For instance, Soulsby and Whitehouse (2005)  
454 develop an equilibrium predictor where  $A/D_{50}$  is the controlling parameter after  
455 examining the collapse of compiled data with several other variables.

456 Using only the reserved testing data, the NRMSE of the new GP predictor as well  
457 as those developed by Soulsby and Whitehouse (2005) and Pedocchi and García (2009a),  
458 are 0.74, 1.33, and 1.22 respectively, and the correlation coefficient is 0.78, 0.02, and  
459 0.20 respectively. The GP derived predictor performs better than the other predictors  
460 based on the NRMSE and correlation coefficient. Figure 8 shows the performance of  
461 these models in both linear and log-log space. Neither of these previously published  
462 predictors were developed for large scale orbital ripples, and both show predictions that  
463 deviate significantly when observed ripple wavelengths are large. The GP derived  
464 predictor is better able to capture large scale ripples. Both Soulsby and Whitehouse  
465 (2005) and Pedocchi and García (2009a) are able to better capture small scale ‘anorbital’  
466 ripples that deviate significantly from the scaling of (9).

467

#### 468 *4.2 Ripple height*

469 The GP algorithm output is shown in Table 3. This experiment evaluated  $10^{10}$   
470 formulas to develop the Pareto front shown in Figure 9. Cliffs occur along the Pareto  
471 front at complexities of 3, 5, 14, 18 and 36 (Figure 9). Predictor of complexity 3,  $\eta$   
472  $=0.435d_0$ , is qualitatively similar to predictions of ripple height in the orbital regime (i.e.  
473 weak hydrodynamics) presented in Wiberg and Harris (1994), where ripple wavelength is  
474 a function of orbital diameter and ripple steepness ( $\eta/\lambda$ ) is constant, therefore ripple  
475 height is a linear function of bottom orbital diameter. Constant steepness breaks down in

476 stronger hydrodynamic conditions, and this is reflected in the inclusion of grain size and  
477 ripple length in more complex predictors. We have no compelling evidence to use the  
478 most nonlinear but best fit solution (36), nor is there compelling evidence at this time that  
479 ripple height has a such a strongly nonlinear dependence on grain size (Solution 14 and  
480 18). We focus our analysis on solution 5:

$$\eta = 0.313\lambda(1000D_{50}) \quad (10)$$

481 or, replacing  $\lambda$  (which denotes predicted ripple wavelength) with equation 9:

$$\eta = \frac{0.313d_0(1000D_{50})}{1.12 + 2.18(1000D_{50})} \quad (11)$$

482 Figure 10 shows the behavior of this predictor under conditions of various orbital  
483 diameter and grain size. Ripple height increases with increasing grain size and orbital  
484 diameter. As with ripple length, ripple height is more sensitive to changes in grain size  
485 than changes in orbital velocity. Reserved testing data is used as an independent dataset  
486 to compare the GP predictor as well as those developed by Soulsby and Whitehouse  
487 (2005) and Pedocchi and García (2009a): the NRMSE for each predictor is 0.79, 1.02,  
488 and 1.01 respectively, and the correlation coefficient is 0.67, 0.41, and 0.47 respectively.  
489 The GP derived predictor performs better than the other predictors based on the NRMSE  
490 and correlation coefficient. Figure 11 shows the performance of these models in both  
491 linear and log-log space.

492

#### 493 *4.3 Ripple steepness:*

494 Combining the GP predictors for ripple length (9) and height (11), or simply  
495 rearranging (10), yields a predictor for ripple steepness:

$$\vartheta = 0.313(1000D_{50}) \quad (12)$$

496 implying that steepness is a function solely of grain size, which is a gross approximation  
 497 of the variability observed in the data, and to some extent even unphysical. To enhance  
 498 our steepness prediction we produce a GP derived steepness predictor. The GP algorithm  
 499 output is shown in Table 4. This experiment evaluated  $10^{10}$  formulas to develop the  
 500 Pareto front shown in Figure 12. Cliffs, significant gains in error for small changes in  
 501 complexity occur along the Pareto front at complexities 5, 8, 10, and 16. The predictor at  
 502 complexity 5 produces nonphysical results (negative steepness under some conditions) so  
 503 is ruled out. The most nonlinear predictor reported (complexity of 16) shows only small  
 504 decrease in error for increasing equation complexity; we focus our analysis on predictor  
 505 10:

$$\vartheta = \frac{3.42}{22 + \left(\frac{\lambda}{(1000D_{50})}\right)^2} \quad (13)$$

506 by replacing  $\lambda$  (predicted ripple wavelength) with (9), yields:

$$\vartheta = \frac{3.42}{22 + \left(\frac{d_0}{1.12(1000D_{50}) + 2.18(1000D_{50})^2}\right)^2} \quad (14)$$

507 Figure 13 shows the behavior of this predictor under conditions of various orbital  
 508 diameter and grain size. Increasing  $D_{50}$  (for a given  $d_0$ ) results in increasing  $\vartheta$  until a  
 509 saturated value of 0.15 is reached. Increasing  $d_0$  (for a given  $D_{50}$ ) results in decreasing  $\vartheta$ .  
 510 Small grain sizes are very sensitive to changes in  $d_0$ , while large grain sizes are relatively  
 511 insensitive. Figure 14 shows the performance of (14) against the independent testing data  
 512 compared to the linear convolution of GP derived length and height (12), as well as the  
 513 Pedocchi and García (2009a) and Soulsby and Whitehouse (2005) predictors. The

514 NRMSE of these predictors is: 0.36, 0.50, 0.47, and 0.43, respectively, and the  
515 correlation coefficient is 0.70, 0.48, 0.63, and 0.50 respectively. The GP derived predictor  
516 performs better than the other predictors (including the linear convolution of GP derived  
517  $\lambda$  and  $\eta$ ) based on the NRMSE and correlation coefficient.

518

## 519 **5. Discussion**

### 520 *5.1 Predictors derived from genetic programming*

521 The suite of predictors that are produced as output of the genetic programming  
522 show a trend of increasing predictability with increasing complexity. Highly nonlinear  
523 predictors have been avoided in this study because they may be fit to the noise or  
524 variance present in the training dataset (i.e. they are overfit). Yet the more complex  
525 nonlinear predictors can be used as hypothesis for further field and lab studies where  
526 grain size effects are a focus.

527 Dependence on orbital scaling and grain size is not imposed by the authors, it is a  
528 result of the data used to feed the genetic programming software. Aside from the data sets  
529 used in this study, other field observations have shown decreasing ripple height and  
530 increasing ripple length in fine grained sand under strong hydrodynamic forcing (e.g., a  
531 transition from steep to low profile bedforms; Green and Black, 1999; Green et al., 2004;  
532 Trembanis et al., 2004). Pedocchi and García (2009a) and Cummings et al. (2009) note  
533 that  $U$  is the major control on the transition from small ripples (anorbital) to large ripples  
534 (hummocks), yet our GP derived predictors contain only one hydrodynamic parameter,  
535  $d_0$ . Furthermore dependence on  $U$  is not present in any of the candidate predictors (Tables  
536 2,3, and 4). This is likely the result of several factors: First,  $d_0$  and  $U$  are correlated in our

537 database (Figure 3), making  $d_0$  a potential proxy for any dependence on  $U$ . Second, our  
538 database likely contains multiple ripple sizes at similar hydrodynamic conditions,  
539 resulting in the lack of a clear velocity threshold. Third, we focus on developing a  
540 continuous predictor so do not include any logical statements that can accommodate a  
541 threshold.

542 Our results show that ripple height is more difficult to predict than ripple length  
543 (e.g., Yan et al 2008; Williams et al 2004). As mentioned previously, this is likely a  
544 consequence of ripple crests being subject to higher flow velocities and suspension  
545 processes. Yet successful height and steepness determination is important for the  
546 prediction of sediment transport, in particular the reference concentration (e.g., Green and  
547 Black 1999) and sediment diffusivity (e.g., Nielsen 1992, Thorne et al., 2009). Only 2  
548 equations are needed to predict height, length and steepness of ripples, but error in the  
549 two chosen predicted parameters cascades to the third. The basic linear convolution of  
550 predicted  $\lambda$  (9) and predicted  $\eta$  (11) demonstrate this cascading error: the resultant  
551 steepness predictor (12) produces results that are solely dependent on grain size. We  
552 instead offer 3 separate equations in the hope that workers will decide which 2 predictors  
553 are most valuable for a specific research question. Notably, the GP algorithm did have  
554 predicted  $\lambda$  and predicted  $\eta$  available as equation building blocks when determining  
555 ripple steepness but the term ' $\eta/\lambda$ ' did not appear in any candidate solutions (Table 4).  
556 Generating 3 separate predictors that are not self-consistent leads to geometric  
557 inconsistencies, but results in better prediction for work that requires accurate prediction  
558 of height and steepness but does not rely on ripple length measurements.

559 The hydrodynamic and sedimentological limit of the current prediction scheme is

560 represented by the 4-dimensional shape that outlines the point cloud in Figure 3 (Table 1  
561 contains more information regarding the range of the dataset). We excluded conditions  
562 where ripples are not present either as a result of sheet flow conditions (upper plane bed)  
563 or because of insufficient mobility (lower plane bed). Uncertainty in the onset of upper  
564 plane bed exists because of the lack of data at a range of  $D_{50}$  in field-scale conditions  
565 (e.g., Li and Amos, 1999; Trembanis et al., 2004; You and Yin, 2006). Additionally, field  
566 work suggests that upper plane bed conditions may not be flat, but instead typified by  
567 dynamic features that may be similar if not identical to long wave ripples (Green and  
568 Black, 1999; Green et al., 2004; Trembanis et al., 2004). As a result of the ambiguity in  
569 bed state under ‘upper plane bed’ conditions we did not compare our predictor to the  
570 version developed by Camenen (2009), which explicitly includes a sheet flow threshold  
571 (where ripples are destroyed). Furthermore we do not compare the GP derived predictors  
572 with those developed by Williams et al., (2005), who developed separate predictors for  
573 short wavelength and long wavelength ripples. We intentionally did not divide our  
574 dataset (and develop separate predictors) in an attempt to construct a practical prediction  
575 scheme that spans a wide range of conditions. Since this study aims to produce a  
576 continuous predictor of wave ripple geometry, the use of discontinuous functions (logical  
577 statements: e.g. ‘if-then-else’) has not been explored quantitatively in this contribution.  
578 This study does not tackle the issue of time dependent adjustment of bedforms in  
579 unsteady flow (Austin et al., 2007, Soulsby et al., 2012, Traykovski 2007; Davis et al.,  
580 2004; Doucette and O’Donoghue, 2006; Hay 2008), the importance of initial conditions  
581 on final ripple configuration (Traykovski et al., 1999; Hansen et al., 2001), or the explicit  
582 incorporation of emergent ripple parameters (e.g., defect density; Skarke and Trembanis,

583 2011).

584

### 585 *5.2 Open Research Questions: Wave Ripples*

586 Data integration campaigns can highlight gaps in knowledge. The collection of new  
587 ripple datasets will be able to be used as either independent tests of the predictors  
588 developed in this study (if the setting corresponds to an area in figure 3 that is dense with  
589 points) or as new data to train the GP algorithm (if the data correspond to unexplored or  
590 sparse area in figure 3; Bowden et al., 2012). Additional datasets of wave ripple geometry  
591 that include more input parameters (e.g. measures of grain sorting, wave irregularity,  
592 initial conditions, time dependence) are needed if prediction accuracy is to increase.  
593 Furthermore datasets that encompass coarse grained environments (coarse sand and  
594 gravel) and datasets in energetic conditions are still needed. Though coarse grained  
595 conditions reflect a smaller fraction of the seabed than fine grained settings, coarse  
596 grained environments are likely important nursery habitat for fish (Hallenbeck et al.,  
597 2012). Collection of this data will not only help in the determination of ripple  
598 configuration under these specific forcing conditions, but linking these environments  
599 with the present data will allow for the development of a better predictor by defining the  
600 shape of the prediction surface over a greater extent in phase space.

601 Conditions with waves and currents (e.g., Lacy et al., 2007; Khelifa and Ouellet,  
602 2000; Arnott and Southard, 1990) are excluded from this analysis. The collection and  
603 integration of data with waves and currents may lead to a more universal bedform  
604 predictor in the future but more data on ripple geometry under wave and current forcing  
605 is needed for machine learning techniques to be applied successfully. The dataset in this



606 study uses median grain size as the sole sedimentological metric for predicting ripple  
607 geometry. Yet many field settings may not be accurately described by a sharp peaked  
608 unimodal distribution of grain sizes, and therefore prediction of ripples using  $D_{50}$  may  
609 lead to significant error. Foti and Blondeaux (1995) showed that the addition of coarse  
610 sediment can act as stabilizing feature, enhancing ripple length. It is possible that graded  
611 sediment will not conform to the predictive tools outlined above. But what is the effective  
612  $D_{50}$  in graded sediment when predicting ripples? Furthermore variations in grain shape  
613 and bed porosity may also impact the geometry of ripples. More research is needed into  
614 the role of mixed grains in determining equilibrium wave ripple geometry (e.g. Calantoni  
615 et al., 2013).

616 More studies are needed to better constrain thresholds between short ‘anorbital’  
617 ripples and large ‘orbital’ scale ripples (Pedocchi and García, 2009a, 2009b; Cummings  
618 et al., 2009; Maier and Hay, 2009). Experimental work has thus far shown that there  
619 exists no intermediate scale between these two configurations (Dumas et al., 2005;  
620 Cummings et al., 2009). The determination of when large ripples appear and when  
621 superimposed short ripples disappear will allow the pruning of the database in regions  
622 where overlapping ripple scales occur. The decision of which ripple scale to eliminate  
623 when both exist is a function of the research question being studied.

624

### 625 *5.3 Open Research Questions: Data Driven Prediction*

626 In this contribution we demonstrate a selection technique whereby very few data  
627 are used to train the GP algorithm and most data is used as validation and independent  
628 testing. The training data was selected solely from variables representing the forcing

629 conditions. As a result the training data is not representative of the entire population of  
630 ripple configurations as data points that are neighbors in 'forcing space' do not necessarily  
631 have similar ripple geometries. The selected training data is therefore only related to the  
632 range of forcing present in the dataset, not the range of ripple geometry. Therefore we  
633 believe that our sampling strategy does not bias the testing of the predictors (which relies  
634 on ripple geometry) using the reserved, unselected testing data.

635         We define the testing dataset as 'independent' because it was not shown to the GP  
636 algorithm. Additionally we performed experiments by removing several individual  
637 datasets from the composite dataset. The removed data serves as testing data that is not  
638 shown to the selection routine, not shown to the GP algorithm, and additionally not  
639 related to data shown to the selection routine/GP (this is another definition of  
640 'independent'). The resultant predictors (not shown) were quantitatively similar to those  
641 presented in this contribution and similarly performed better than the Pedocchi and  
642 García (2009a) and Soulsby and Whitehouse (2005) predictors using only the smaller  
643 sample of removed datasets as testing. Is it enough for testing data to be unseen by the  
644 ML algorithm, or do entire datasets need to be reserved whole as testing data? More  
645 investigation will resolve this issue.

646         Even though we were able to obtain good results using few centroids, we are  
647 unaware of a technique for quantitatively determining the optimal number of centroids to  
648 capture the variability in the data set while leaving the maximum amount of data for use  
649 as validation/testing. Furthermore many selection and clustering routines are available,  
650 and it is unclear which routine is optimal for a given dataset. It is likely that some of the  
651 answers to these questions lie in statistical science and computer science literature that

652 has not fully percolated into the Earth Sciences.

653 Our observations with the GP software show too few centroids tend to underfit the  
654 data because the GP has too little training data to develop applicable solutions. With few  
655 training data solutions tend to be linear and have low RMSE when compared tot the  
656 validation and testing datasets. Training datasets that are larger than used in this study (>  
657 40 centroids) tend to produce more large (complexity > 30) nonlinear solutions. In  
658 addition, the solutions at complexity less than 30 are similar (if not identical) to the  
659 solutions in this study (using smaller training datasets). The invariance of solutions gives  
660 qualitative justification to the number of centroids used in this study, but we do not offer  
661 a quantitative technique for determining the minimum number of training data needed to  
662 capture dataset variability. Furthermore how is this number linked to the quantity and  
663 quality of the data in the training dataset? Lastly, the stability and final criterion for  
664 selecting a single predictor is subjective and can likely be improved or quantitatively  
665 justified by implementing more sophisticated accounting techniques based on  
666 information such as the Aikake Information Criterion (AIC) or the Bayesian Information  
667 Criterion (BIC).

668

## 669 **6. Conclusion**

670 We develop equilibrium predictors of oscillatory ripple geometry using genetic  
671 programming. Ripple length is a weak nonlinear function grain size and bottom orbital  
672 excursion. Ripple height and steepness are nonlinear functions of grain size and predicted  
673 ripple length (i.e. grain size and bottom orbital excursion). Furthermore these new  
674 predictor encompass a wide range of hydrodynamic and sedimentological conditions not

675 previously included in published prediction schemes. However, the proposed method is  
676 not suitable for practical applications with significant currents present, nor under  
677 conditions that would either be below the threshold of motion or above the threshold of  
678 ripple wash-out. Such conditions should be identified separately by existing methods  
679 (Nielsen, 1992; Lacy et al., 2007; Camenen, 2009; Soulsby et al., 2012)

680 This contribution further demonstrates the viability of developing empirical  
681 predictors through ML techniques. As previously mentioned by Oehler et al., (2012), ML  
682 algorithms could be integrated into future morphodynamics models (model-data fusion  
683 and the development of a ‘hybrid’ model; Krasnapolsky and Fox-Rabinovitz, 2006),  
684 replacing functions with large uncertainty.

685 The data integration campaign (which preceded the implementation of the GP  
686 algorithm) had the side benefit of highlighting the current state of our knowledge on  
687 ripple geometry, potentially motivating targeted data collection campaigns. Newly  
688 collected data can be fed back into the GP software to develop revised predictors.

689

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695

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956

957 FIGURES AND TABLES

958

959 Table 1: Data Summary; Measurement numbers reported are the ripple length

960 measurements used in our study. Measurements with both length and height are less.

Authors	Setting	Measurements	T (s)	U (m/s)	d <sub>0</sub> (m)	D <sub>50</sub> (m)
Boyd et al. 1988	Field	36	3.8-9.8	0.04-0.28	0.05-0.60	0.00011
Cummings et al. 2009	Wave Racetrack	14	4.4-14	0.37-1.22	0.95-4.50	0.00012-0.0008
Delgado Blanco et al. 2004	Wave Flume	17	6.0	0.14-0.74	0.27-1.42	0.00035
Doucette 2000	Field	49	4.7-12.2	0.15-0.52	0.31-1.93	0.00015-0.00053
Doucette 2002	Field	25	2.2-12.2	0.17-0.66	0.31-2.22	0.00035-0.00062
Doucette and O'Donoghue 2006	Osc. Tunnel	32	2.0-12.2	0.29-0.63	0.24-2.00	0.00044
Dumas et al. 2005	Osc. Tunnel	23	7.9-11.0	0.21-1.26	0.51-4.17	0.00011-0.00023
Grasmeijer and Kleinhans 2004	Field	26	4.0-10.5	0.23-0.84	0.58-2.41	0.00024
Hanes et al. 2001	Field	169	7.1-19.7	0.92-1.11	0.47-5.02	0.00012-0.00166
Hume et al. 1999	Field	9	11.0	0.08-0.37	0.30-1.30	0.00040
Inman 1957	Field	59	0.5-15.0	0.06-0.94	0.04-2.74	0.00008-0.00091
Kennedy and Falcon 1965	Wave Flume	10	1.1-2.0	0.12-0.26	0.04-0.13	0.00010-0.00032
Miller and Komar 1980a	Wave Flume	4	3.0-8.0	0.05-0.34	0.14-0.54	0.00017
Miller and Komar 1980b	Field	26	6.0-18.2	0.03-0.41	0.07-2.14	0.00017-0.00029
Mogridge 1972	Osc. Tunnel/ W. Flume	72	1.0-14.0	0.13-0.68	0.05-1.84	0.00036
O'Donoghue and Clubb 2001	Osc. Tunnel	35	2.0-15.0	0.25-0.94	0.16-2.92	0.00018-0.00044
O'Donoghue et al. 2006	Osc. Tunnel	27	3.1-12.5	0.31-0.85	0.42-2.70	0.00022-0.00044
Pedocchi and García 2009b	Osc. Tunnel	22	2.0-18.0	0.20-1.00	0.16-2.86	0.00025
Ribberink and Al-Salem 1994	Osc. Tunnel	25	2.0-10.0	0.30-1.50	0.31-3.82	0.00021
Sleath 1982	Osc. Tunnel	13	2.9-5.1	0.16-0.44	0.17-0.51	0.00020-0.00041
Sleath and Wallbridge 2002	Osc. Tunnel	26	2.8-5.	0.08-0.77	0.12-0.80	0.00020-0.00080
Southard et al. 1990	Osc. Tunnel	63	93.1-19.3	0.16-1.00	0.26-3.56	0.00011-0.00032
Thorne et al. 2002	Wave Flume	14	4.0-6.0	0.26-0.66	0.41-1.05	0.00033
Williams et al. 2000	Wave Flume	9	4.8-5.3	0.19-0.69	0.30-1.10	0.00016-0.00033
Williams et al. 2004	Wave Flume	65	4.0-6.0	0.13-1.02	0.25-1.96	0.00016-0.00035
Xu 2005	Field	13	8.9-14.8	0.11-0.16	0.41-0.76	0.00009
Yamaguchi and Sekiguchi 2011	Wave Flume	111	1.3-5.0	0.18-0.51	0.07-0.55	0.00032-0.00073

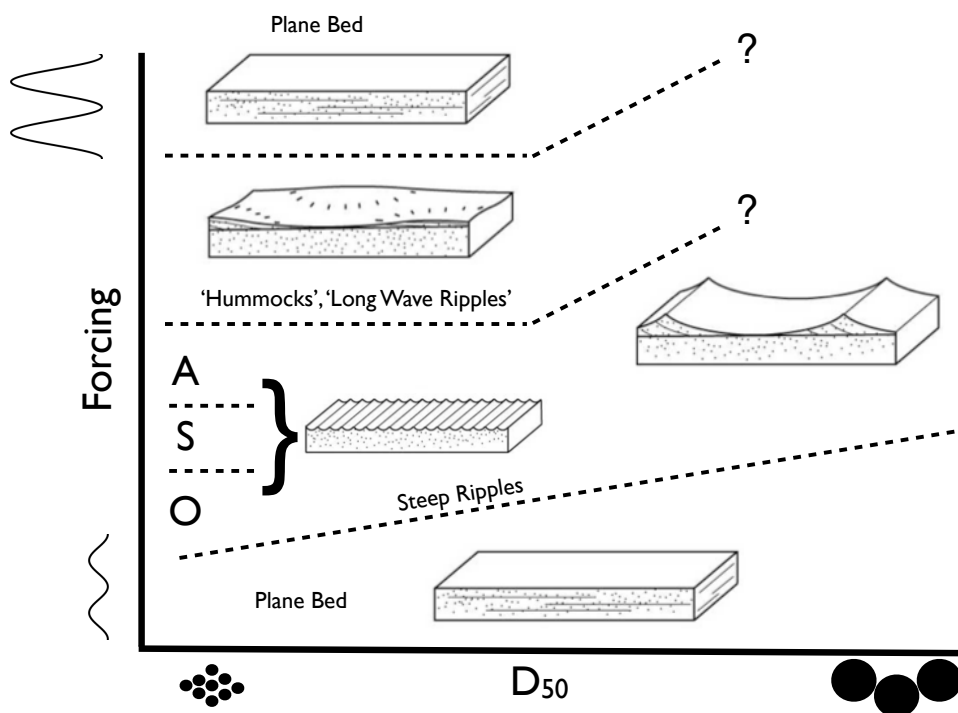
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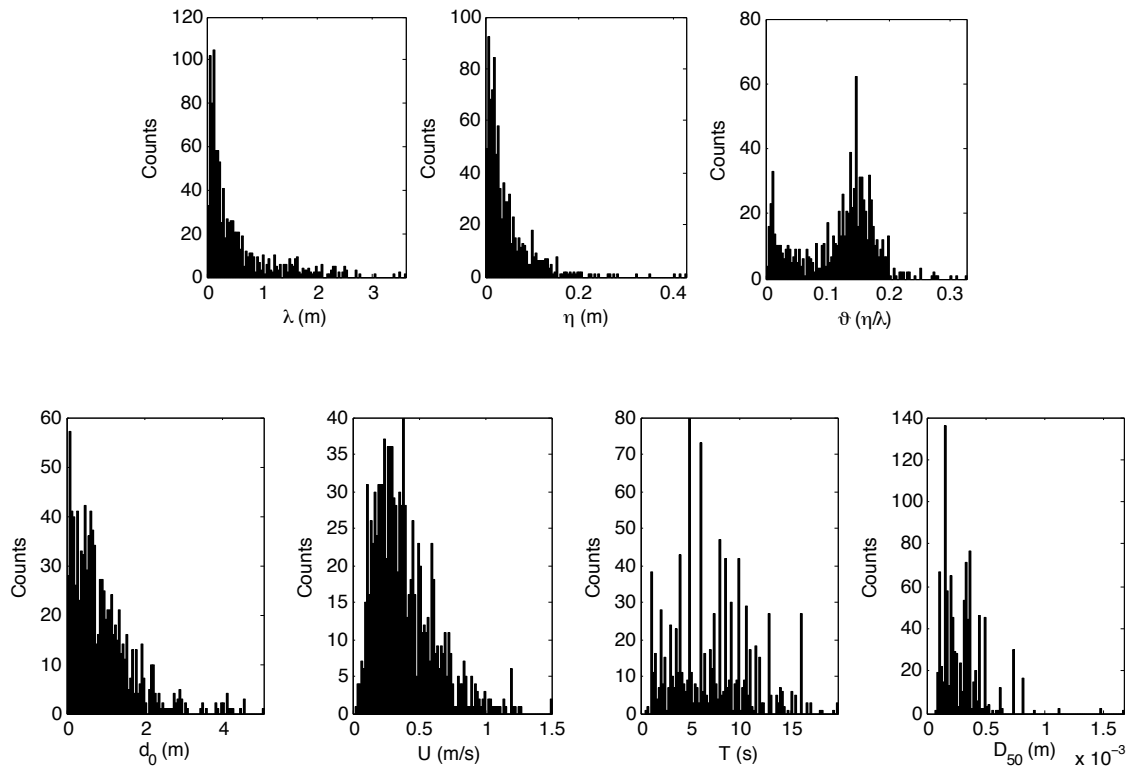
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965 Figure 1: A schematic phase diagram of oscillatory bedforms. O, S, and A represent  
966 orbital, suborbital, and anorbital ripples respectively: smaller steep ripples that occur  
967 under small/moderate hydrodynamic forcing in fine sands. Orbital, suborbital and  
968 Anorbital ripples occur in sequence as hydrodynamic forcing is increased. Recent data  
969 collection campaigns have focused on 1) strong hydrodynamic forcing in fine sands  
970 ('hummocks' or 'long wave ripples') and 2) steep, large ripples in coarse sand. Modified  
971 after Cummings et al., (2009). Question marks denote the unknown threshold for plane  
972 bed in coarse grained environments, and unknown potential for coarse grained  
973 environments to be sculpted into long wavelength 'hummocky' ripples. Additionally it is  
974 unknown if suborbital and anorbital scale ripples exist in coarse grain settings. Lower  
975 plane bed conditions are likely only applicable for laboratory studies where the bed is  
976 artificially flattened (field conditions retain relict or antecedent bed geometry).



977

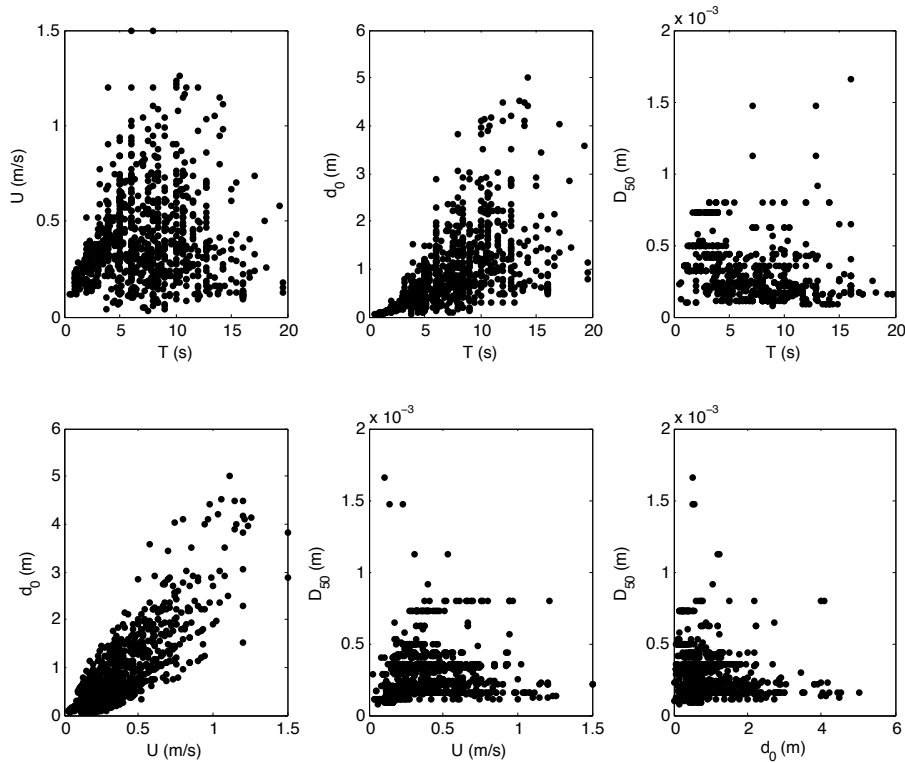
978 Figure 2: Histograms for ripple length (995 measurements), ripple height (872  
979 measurements), ripple steepness (872 measurements), and for hydrodynamic and  
980 sedimentological variables used in this study (includes all 995 data points). Note the  
981 different Y-axis values for each graph.



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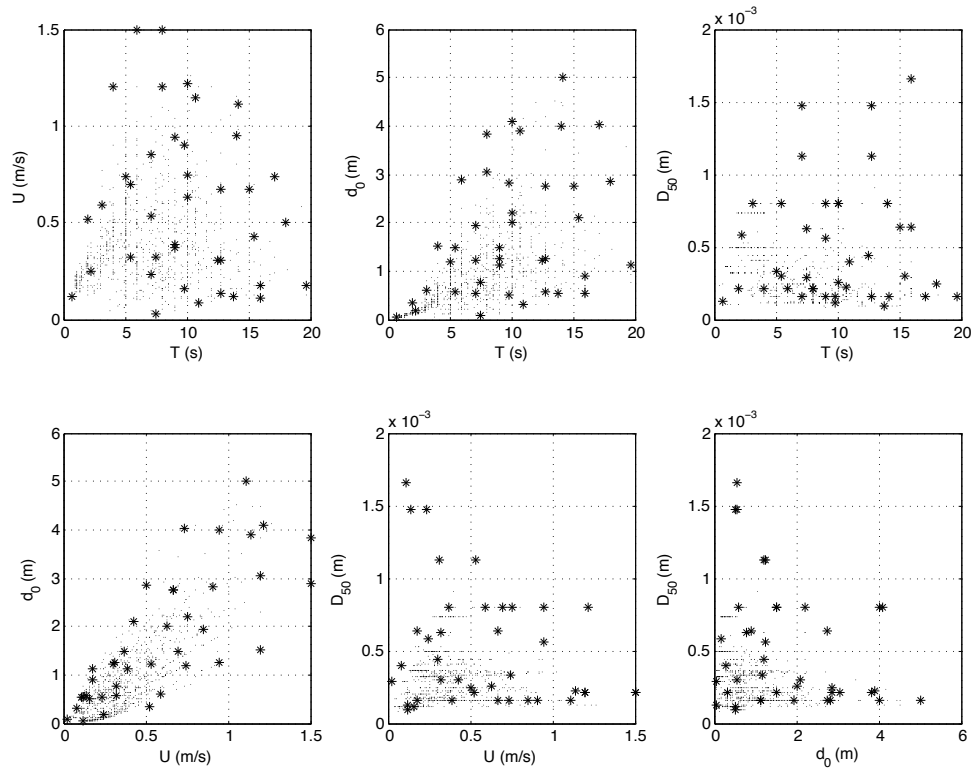
984 Figure 3: Visualization of the range of forcing conditions in the ripple length dataset.  
985 Each plot represents a 2 dimensional projection of the entire data set onto the set of axes  
986 shown. For instance, the first panel with data projected onto the U-T plane shows no  
987 information about  $D_{50}$  or  $d_0$ . Ripple height dataset shows qualitatively similar distribution  
988 and range, but with fewer data points (872 vs. 995).



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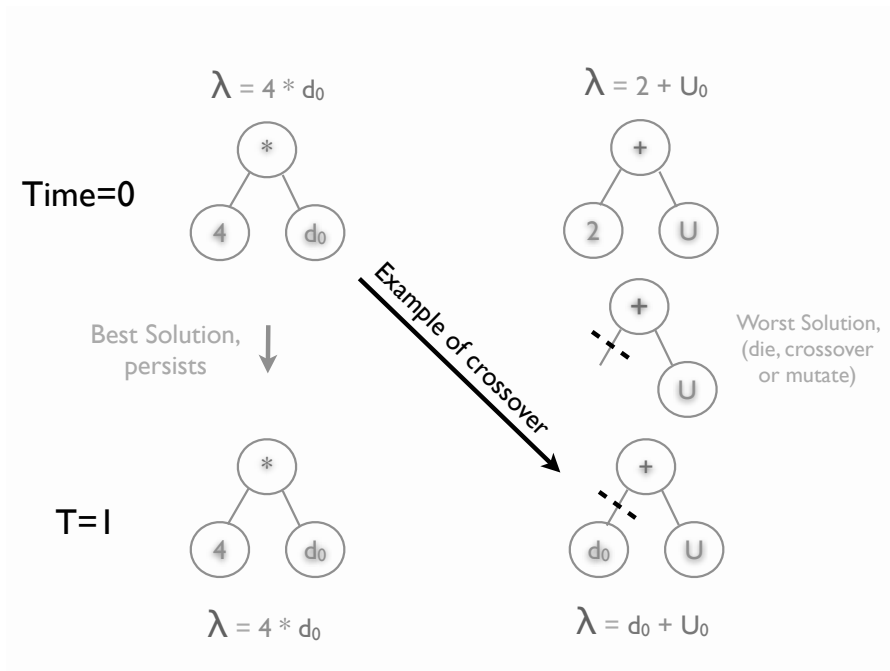
991 Figure 4: Centroid locations in the ripple length dataset, visualized using the projections  
992 shown in Figure 3. Stars denote centroid locations (training data), while points denote  
993 unselected data (validation and testing). Note that centroids are distributed throughout the  
994 dataset. Centroid locations for the ripple height (and steepness) dataset look qualitatively  
995 similar but have more centroids (40 vs. 30) and fewer data points (832 vs. 965).



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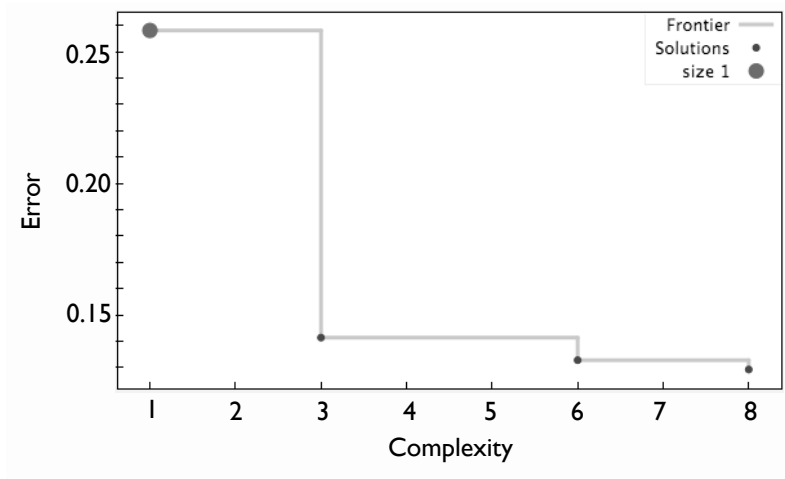
998 Figure 5: Example of the genetic programming process. Potential solutions are encoded  
 999 as a population of ‘trees’. Here a hypothetical population of two solutions is shown. The  
 1000 first solution has a low MSE and therefore persists to the next iteration. The second  
 1001 solution has a high MSE and therefore is subject to removal, mutation, or crossover. Here  
 1002 is an example of ‘crossover’ whereby the old solution is combined with parts of other,  
 1003 better performing solutions to create a new potential solution in the next iteration.



1004  
 1005



1006 Figure 6: Ripple Length Pareto front; Error is expressed as mean squared error of  
 1007 candidate solution versus the validation data set. Complexity is a quantification of the  
 1008 candidate solution length (both mathematical operators and variables).



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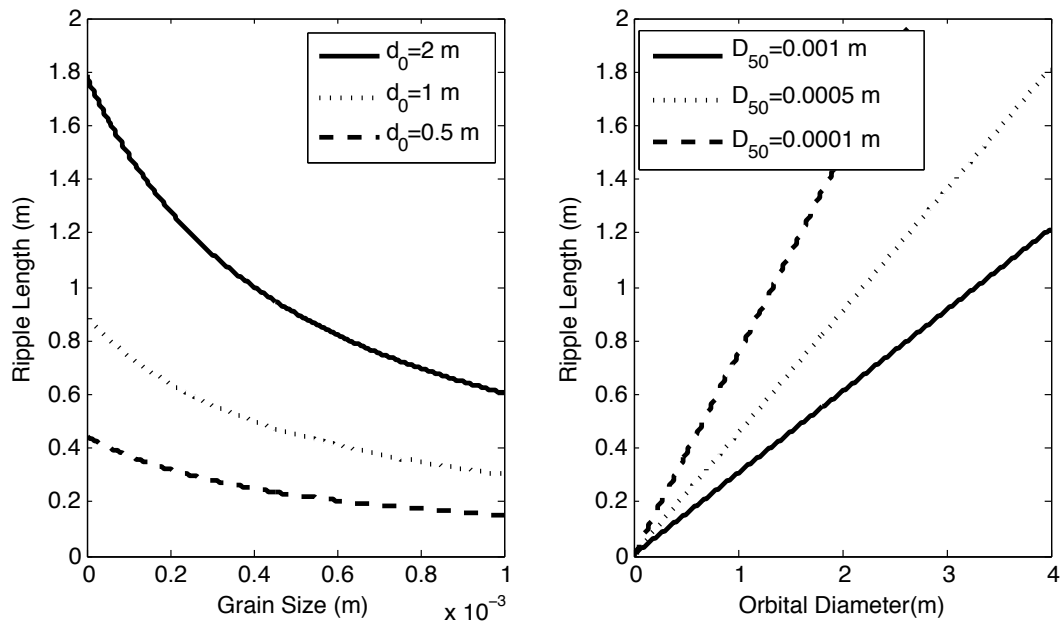
1011 Table 2: Solutions for Ripple Length

Solution	Complexity	MSE
$\lambda = U$	1	0.258
$\lambda = 0.607d_0$	3	0.141
$\lambda = \frac{d_0}{1.39 + (1000D_{50})}$	6	0.133
$\lambda = \frac{d_0}{1.12 + 2.18(1000D_{50})}$	8	0.129

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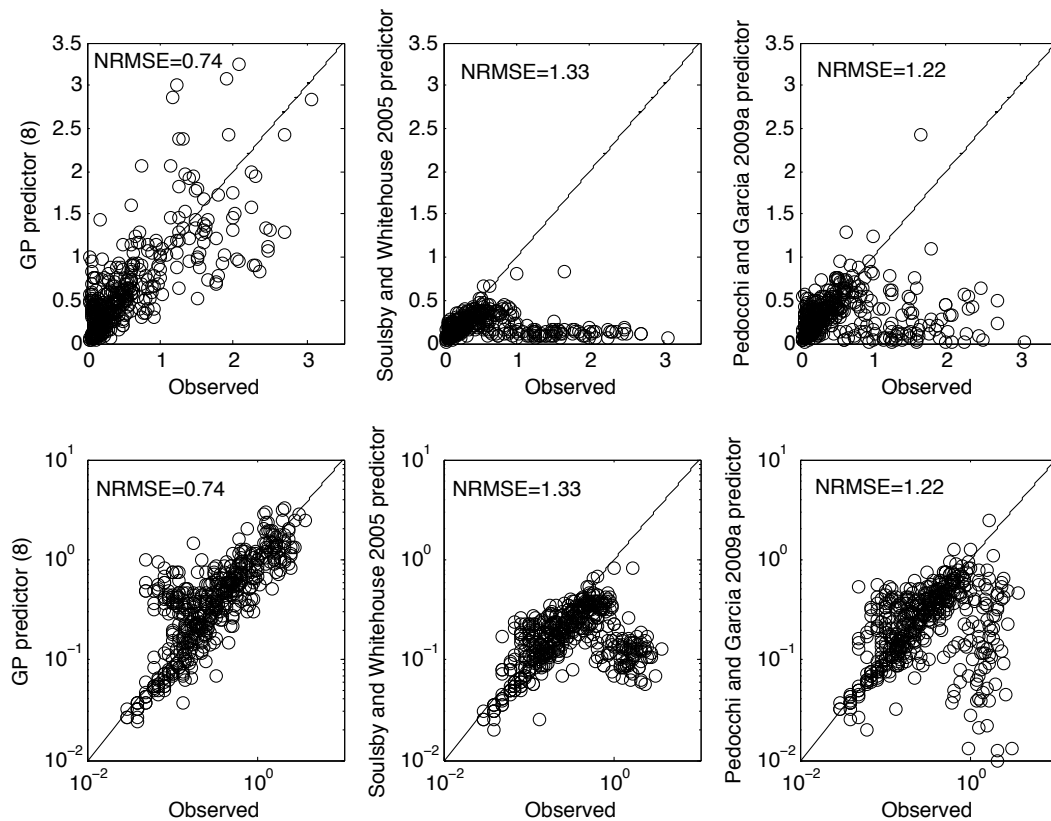
1014 Figure 7: Example behavior of ripple length predictor as a function of grain size for given  
1015 bottom orbital excursions (left panel) and as a function of bottom orbital excursion for  
1016 given grain size (right panel).



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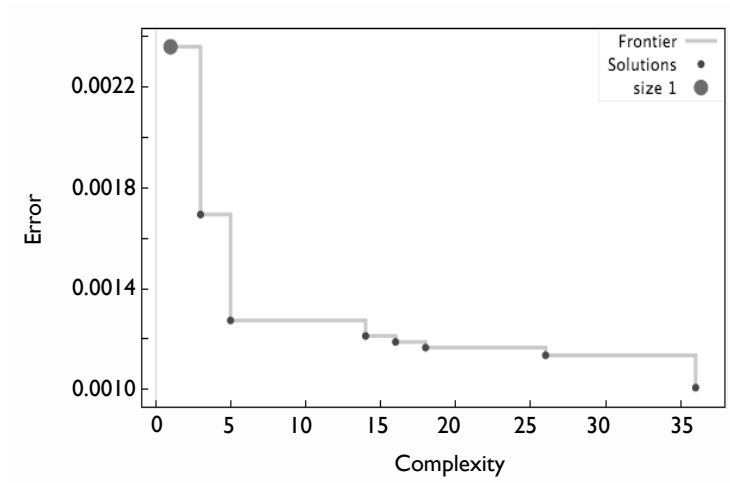
1019 Figure 8: GP predictor of ripple length (8), Soulsby and Whitehouse (2005) predictor (3)  
1020 and Pedocchi and García (2009a) predictor (5) evaluated using only the independent  
1021 testing dataset. Top row shows the predictors in linear space, while bottom row shows  
1022 log-log space.



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1025 Figure 9: Pareto front for ripple height; Error is mean squared error of candidate solution  
 1026 versus the validation data set. Complexity is a quantification of the candidate solution  
 1027 length (both mathematical operators and variables).



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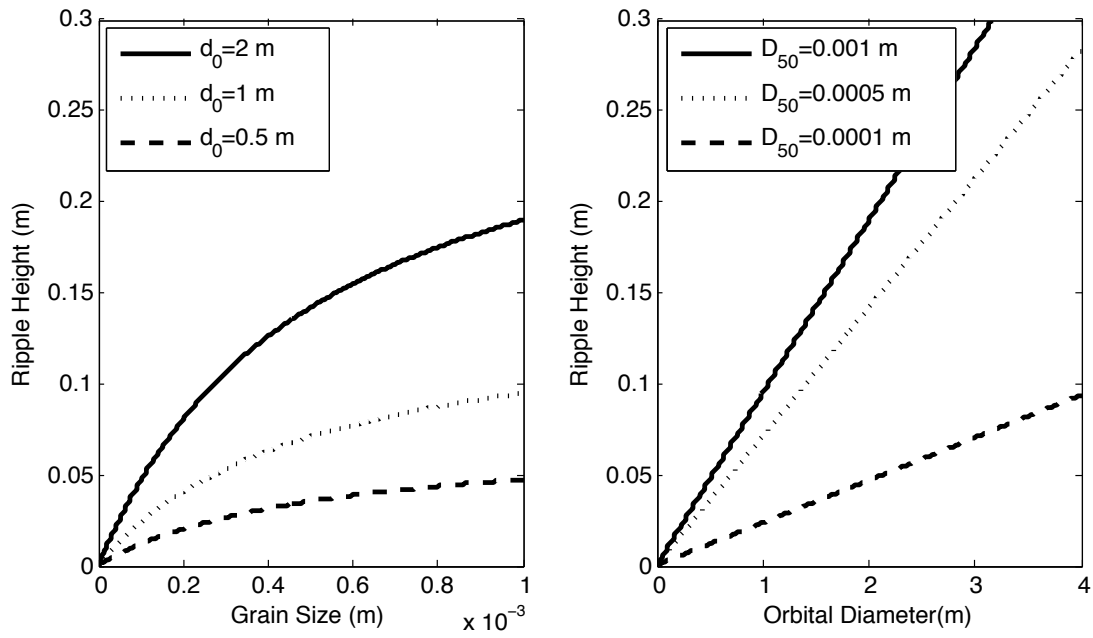
1030 Table 3: Solutions for Ripple Height

Solution	C	MSE
$\eta = 0.435d_0$	3	0.0017
$\eta = 0.313\lambda(1000D_{50})$	5	0.0013
$\eta = \frac{\lambda(1000D_{50})^2}{0.372 + 5.29(1000D_{50})^2}$	14	0.0012
$\eta = \frac{\lambda(1000D_{50})^3}{0.0731 + 5.57(1000D_{50})^3}$	18	0.0012
$\eta = \frac{0.0237\lambda(1000D_{50}) + \lambda(1000D_{50})^3 - 0.308\lambda(1000D_{50})^2}{0.0332 + 4.46(1000D_{50})^3 - 0.321D_{50}}$	36	0.0010

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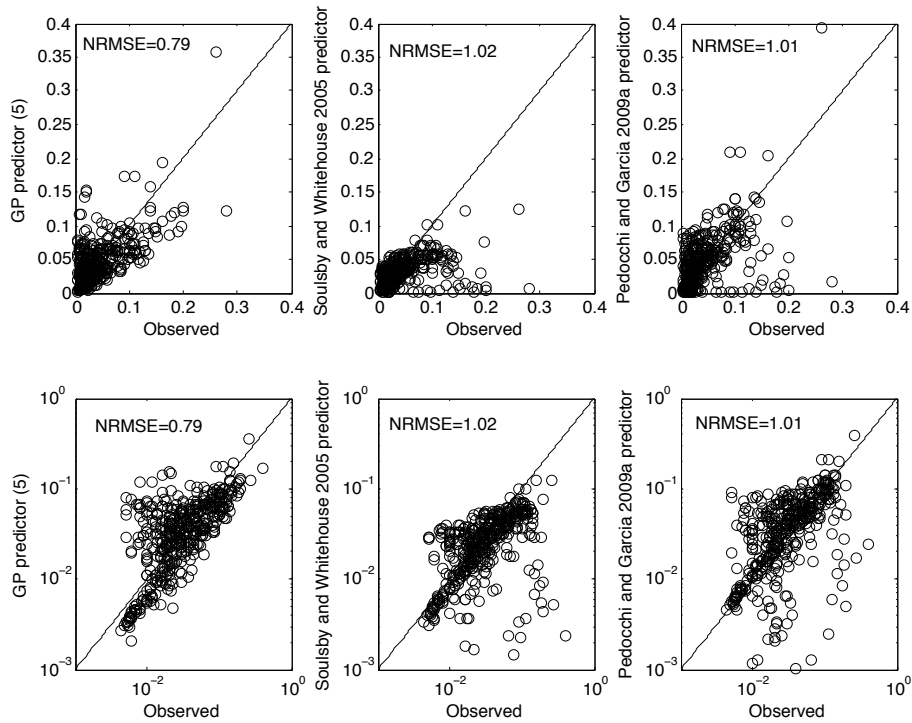
1033 Figure 10: Example behavior of Ripple height predictor as a function of grain size for  
1034 given bottom orbital excursions (left panel) and as a function of bottom orbital excursion  
1035 for given grain size (right panel).



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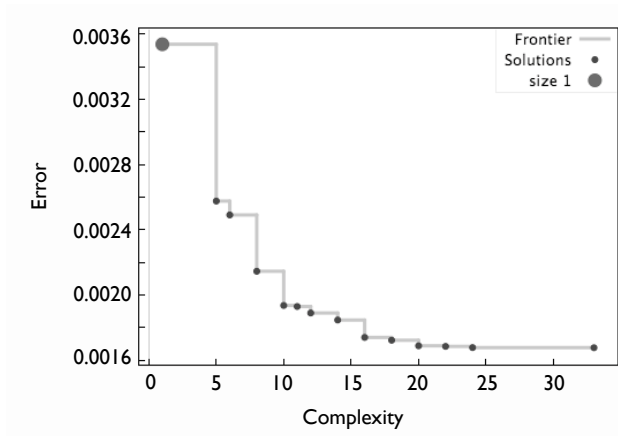
1038 Figure 11: GP predictor of ripple height (10), Soulsby and Whitehouse (2005) predictor  
1039 (3) and (4) and Pedocchi and Garcia (2009a) predictor (6) evaluated using only the  
1040 independent testing dataset. Top row shows the predictors in linear space, while bottom  
1041 row shows log-log space.



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1044 Figure 12: Pareto front for ripple steepness; Error is mean squared error of candidate  
 1045 solution versus the validation data set. Complexity is a quantification of the candidate  
 1046 solution length (both mathematical operators and variables).



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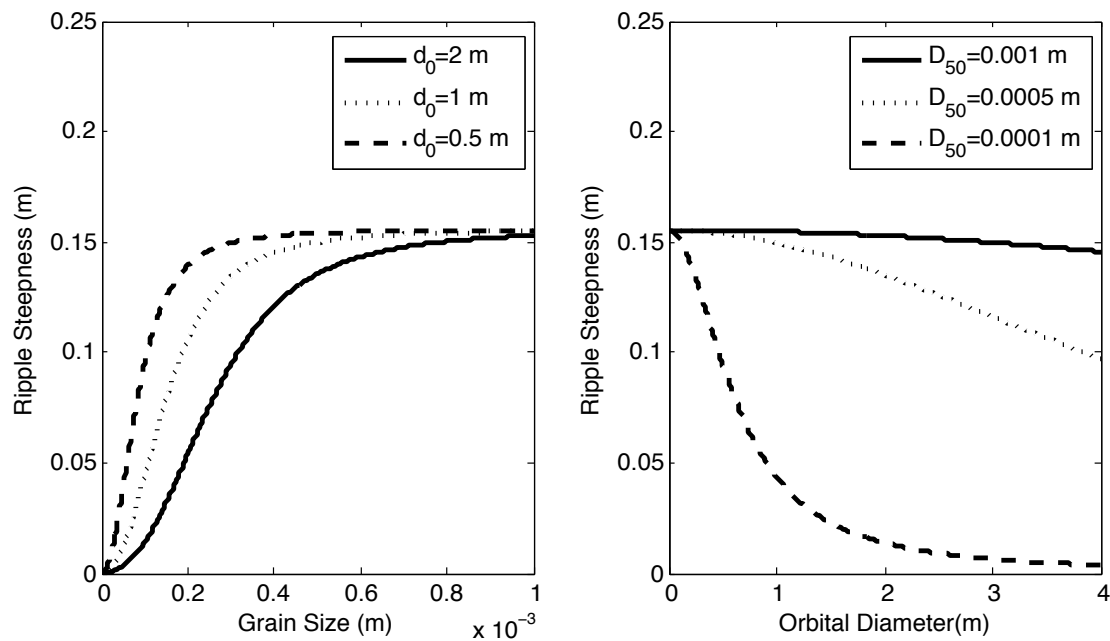
1048 Table 4: Solutions for Ripple Steepness

Solution	C	MSE
$\vartheta = 0.119$	1	0.0035
$\vartheta = 0.154 - 0.0613\lambda$	5	0.0026
$\vartheta = \frac{(1000D_{50})}{\lambda + 6.23(1000D_{50})}$	8	0.0021
$\vartheta = \frac{3.42}{22 + \left(\frac{\lambda}{(1000D_{50})}\right)^2}$	10	0.0019
$\vartheta = \frac{0.447}{2.81 + \left(\lambda^2 + \frac{-0.617\lambda}{(1000D_{50})}\right)^2}$	16	0.0017

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1051 Figure 13: Example behavior of Ripple Steepness predictor as a function of grain size for  
1052 given bottom orbital excursions (left panel) and as a function of bottom orbital excursion  
1053 for given grain size (right panel).

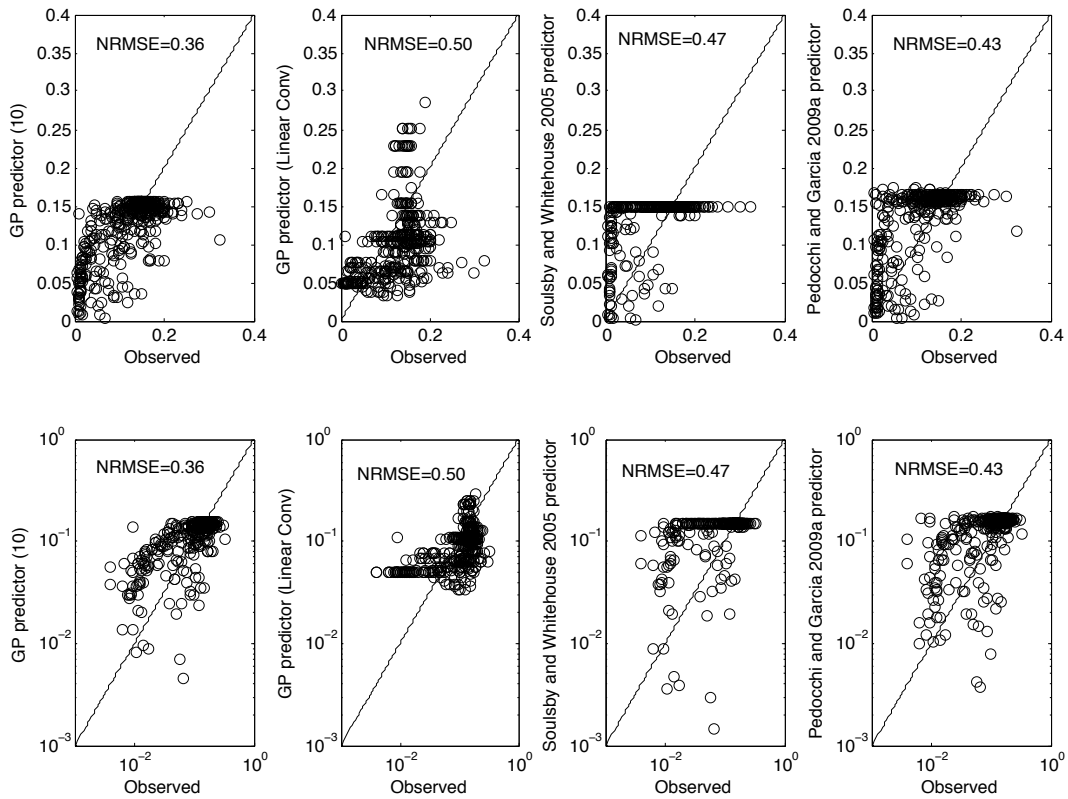


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1056 Figure 14: GP predictor of ripple steepness (13), Predictor based on linear convolution of  
 1057 GP height and length (11), Soulsby and Whitehouse (2005) predictor (4) and Pedocchi  
 1058 and García (2009a) predictor (5) and (6) evaluated using only the independent testing  
 1059 dataset. Top row shows the predictors in linear space, while bottom row shows log-log  
 1060 space.



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