

1 **Seismicity Induced by Longwall Coal Mining at the**
2 **Thoresby Colliery, Nottinghamshire, U.K.**

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ABSTRACT

16 *The U.K. has a long history of deep coal mining, and numerous cases of mining-induced seismicity*
17 *have been recorded over the past 50 years. In this study we examine seismicity induced by longwall*
18 *mining at one of the U.K.'s last deep coal mines, the Thoresby Colliery, Nottinghamshire. After public*
19 *reports of felt seismicity in late 2013 a local seismic monitoring network was installed at this site,*
20 *which provided monitoring from February to October 2014. This array recorded 305 seismic events,*
21 *which form the basis of our analysis.*

22 *Event locations were found to closely track the position of the mining face within the Deep Soft Seam,*
23 *with most events occurring up to 300 m ahead of the face position. This indicates that the seismicity is*
24 *being directly induced by the mining, as opposed to being caused by activation of pre-existing tectonic*
25 *features by stress transfer. However, we do not observe correlation between the rate of excavation*
26 *and the rate of seismicity, and only a small portion of the overall deformation is being released as*
27 *seismic energy.*

28 *Event magnitudes do not follow the expected Gutenberg-Richter distribution. Instead, the observed*
29 *magnitude distributions can be reproduced if a Truncated Power Law distribution is used to simulate*
30 *the rupture areas. The best-fit maximum rupture areas correspond to the distances between the Deep*
31 *Soft Seam and the seams that over- and underlie it, which have both previously been excavated. Our*
32 *inference is that the presence of a rubble-filled void (or goaf) where these seams have been removed*
33 *is preventing the growth of larger rupture areas.*

34 *Source mechanism analysis reveals that most events consist of dip-slip motion along near-vertical*
35 *planes that strike parallel to the orientation of the mining face. These mechanisms are consistent with*
36 *the expected deformation that would occur as a longwall panel advances, with the under- and over-*
37 *burdens moving upwards and downwards respectively to fill the void created by mining. This further*
38 *reinforces our conclusion that the events are directly induced by the mining process. Similar*
39 *mechanisms have been observed during longwall mining at other sites.*

40

41 **Keywords**

42 Geomechanics; Earthquake Source Observations; Induced seismicity; seismic anisotropy

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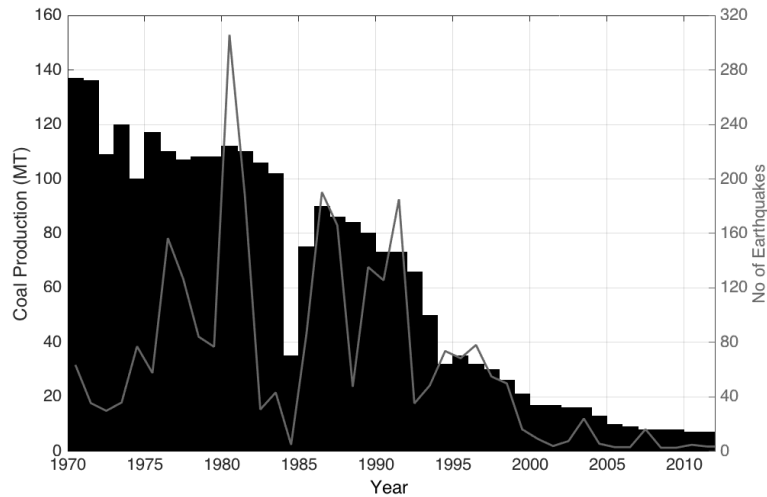
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48 **1. INTRODUCTION**

49 Seismicity induced by coal mining has been a common occurrence in the United Kingdom
50 (e.g., Redmayne, 1988). Indeed, Wilson et al. (2015) estimated that between 20 – 30% of all
51 earthquakes recorded in the UK between 1970 – 2012 were induced by coal mining. From the
52 late 1980s onwards the rate of coal production has declined significantly, as has the rate of
53 associated earthquakes (Figure 1).



54
55 *Figure 1: Deep mined coal production in the UK by year (bars) and the number of induced*
56 *earthquakes per year associated with coal mining (grey line), as categorised by Wilson et al.*
57 *(2015). The drop in both production and induced seismicity in 1984 is associated with the UK*
58 *miner’s strike.*

59 Nevertheless, seismicity associated with deep coal mining still occurs in the UK. Between
60 December 2013 – January 2014, the UK’s national seismometer network detected a series of
61 over 40 earthquakes near to the village of New Ollerton, Nottinghamshire. The largest of
62 these events had a magnitude of $M_L = 1.7$. Given the generally low levels of seismicity in the
63 UK, the village was dubbed the “UK’s Earthquake Capital” (Turvill, 2014). The area has a
64 history of seismic activity relating to coal mining (e.g., Bishop et al., 1993), and it was soon
65 identified that the events were likely to be associated with longwall coal mining at the nearby
66 Thoresby Colliery, which at the time was one of the few remaining deep coal mining sites in
67 the UK.

68 In response to the felt earthquakes, a temporary local monitoring network of surface
69 seismometers was deployed between the 5th February and the 30th October 2014 by the British
70 Geological Survey (BGS). This network recorded a further 300 seismic events. The high
71 quality of the data recorded by the local network permits a detailed study into the nature of
72 seismicity and deformation induced by the longwall mining process.

73

74 *1.1 Longwall Coal Mining at Thoresby*

75 The Thoresby Colliery opened in 1925. Over the history of the site, at least 4 different seams
76 have been mined, including the High Hazels, Top Hard, Deep Soft and Parkgate Seams, in
77 order from shallowest to deepest: see Edwards (1967) for a stratigraphic section showing the
78 positions of these and other seams in the region. The Deep Soft Seam was the last to be
79 developed, with work beginning in 2010: this was the only seam being actively mined during
80 the study period. The colliery closed entirely in mid-2015. This was for economic reasons, i.e.
81 the low price of coal, not because of the induced seismicity.

82 The Deep Soft Seam was mined using standard longwall methods: hydraulic jacks are used to
83 support the roof while a shearing device cuts coal from the face. As the face advances, the
84 jacks are moved forward, allowing the roof to collapse into the cavity that is left behind. The
85 collapsed, brecciated roof material filling this void is known as goaf (e.g. Younger, 2016). At
86 Thoresby, each longwall panel has dimensions of approximately 300 m width, between 1,000
87 – 3,000 m length, and approximately 2.5 m height.

88

89 *1.2 Seismicity Associated with Longwall Coal Mining*

90 Seismicity has often been associated with the longwall mining process (e.g., Cook, 1976;
91 Gibowicz et al., 1990; Bishop et al., 1993; Stec, 2007; Bischoff et al., 2010; Sen et al., 2013).
92 Seismic events associated with coal mining have often been divided into two categories:
93 “mining-tectonic” activity, produced by activation of pre-existing tectonic faults, and
94 “mining-induced” activity, directly associated with the mining excavations (e.g., Stec, 2007).

95 Observed magnitudes have typically ranged from $0.5 < M_L < 3.5$. At some sites event
96 magnitudes have followed the Gutenberg and Richter (1944) distribution (e.g., Bishop et al.,
97 1993; Kwiatek et al., 2011), while in other cases bimodal or other frequency-magnitude
98 distributions have been observed (e.g. Stec, 2007; Hudyma et al., 2008; Bischoff et al., 2010).
99 These non-Gutenberg-Richter distributions have been attributed to the presence of
100 characteristic length scales (the dimensions of the mined panels, for example) that provide a
101 control on rupture dimensions and thereby event magnitudes.

102 Analysis of event focal spheres has revealed a variety of source mechanisms in different
103 settings (e.g., Stec, 2007; Bischoff et al., 2010; Sen et al., 2013) including: non-double-couple
104 events, indicating a volumetric component of deformation usually associated with the roof
105 collapse process; double-couple events showing a direct relationship to mined panels, with
106 vertical fault planes running parallel to the mining face, on which dip-slip motion occurs; and

107 double-couple events that correspond to regional fault orientations and *in situ* tectonic stress
108 conditions.

109 In this paper we follow the processes developed in the aforementioned studies to characterise
110 the seismicity induced by mining at the Thorseby Colliery. We begin by locating events,
111 comparing the event locations to the propagation of the mining faces with time, and
112 seismicity rates with the volume of coal extracted from the mine. We investigate the source
113 characteristics of the events, using spectral analysis combined with event frequency-
114 magnitude distributions to assess the length-scales of structures that have generated the
115 observed events. We use shear-wave splitting analysis to image *in situ* stress orientations at
116 the site, and we calculate focal mechanisms for the events to establish the orientations of fault
117 planes and slip directions generated by the mining process.

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119 **2. EVENT DETECTION AND LOCATION**

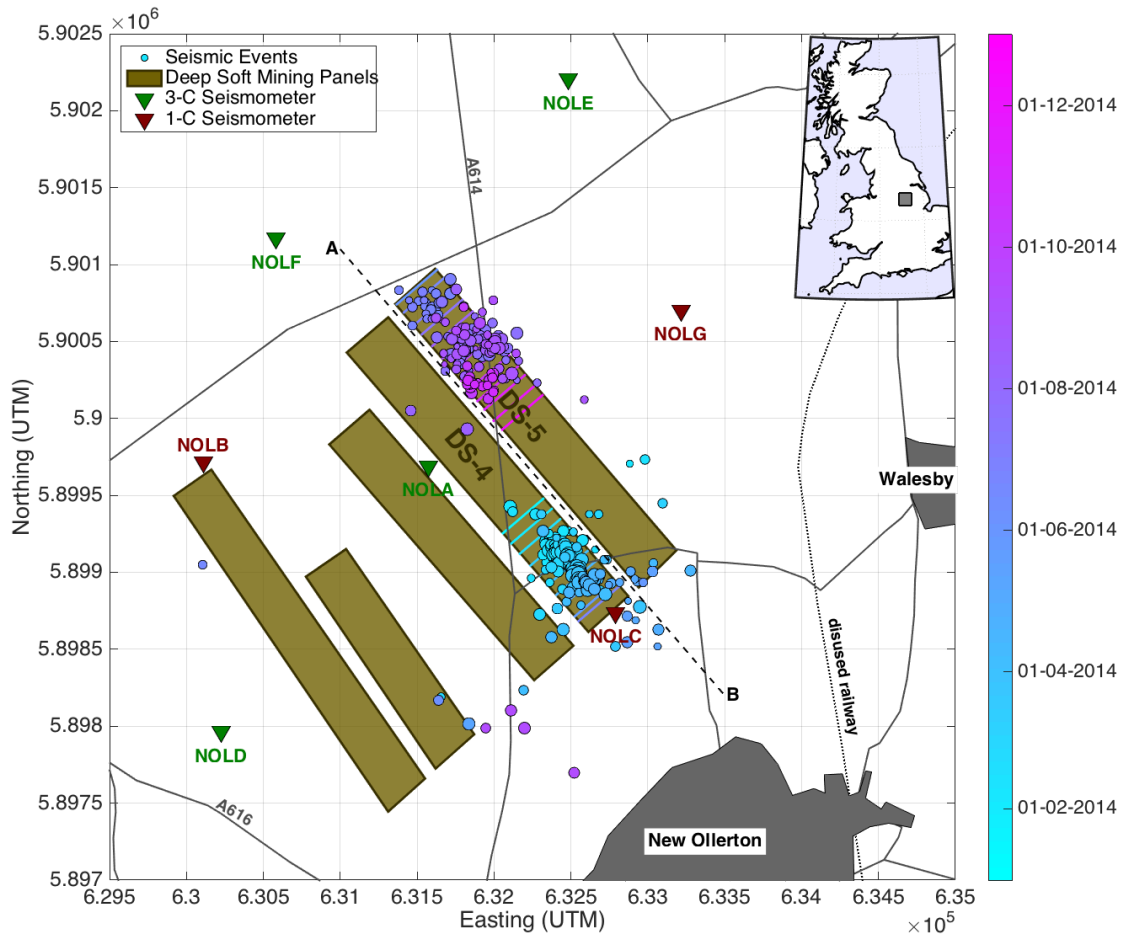
120 *2.1. Monitoring array and event detection*

121 The local surface network deployed to monitor seismicity at the Thoresby Colliery comprised
122 of 4 3-component Guralp 3ESP broadband seismometers (stations NOLA, NOLD, NOLE and
123 NOLF) and 3 vertical-component Geotech Instruments S13J short-period seismometers
124 (NOLB, NOLC and NOLG)). The station positions are shown in Figure 2. Events were
125 detected using the BGS's in-house event detection algorithm, which is based on identification
126 of peaks in running short-time/long-time averages (STA/LTA), as described by Allen (1982).
127 A total of 305 events were identified during the deployment of the local monitoring network.

128 P- and S-wave arrival times were re-picked manually for every event (e.g. Figure 3). For most
129 event-station pairs the P-wave arrival was clear and unambiguous, and so could be accurately
130 picked (83% of station-event pairs where a pick could be manually assigned). Stations
131 NOLB, NOLC and NOLG were single, vertical component stations, so S-wave picks were not
132 made for these stations. For smaller events with lower signal-to-noise ratios, clear S-wave
133 arrivals were sometimes difficult to identify, resulting in a lower number of picks (74% of
134 station-event pairs where a pick could be manually assigned).

135 The velocity model used to locate the events is taken from Bishop et al. (1993), and is listed
136 in Table 1. The arrival time picks were inverted for the best-fitting location that minimises the
137 least-squares residual between modelled and picked arrival times. The search for the best-
138 fitting location was performed using the Neighbourhood Algorithm (Sambridge, 1999), and
139 the modelled travel times were calculated using an Eikonal solver (Podvin and Lecomte,

140 1991). A map of event hypocentres is shown in Figure 2, in which the mining panels and the
 141 position of the mining face with time are also shown.



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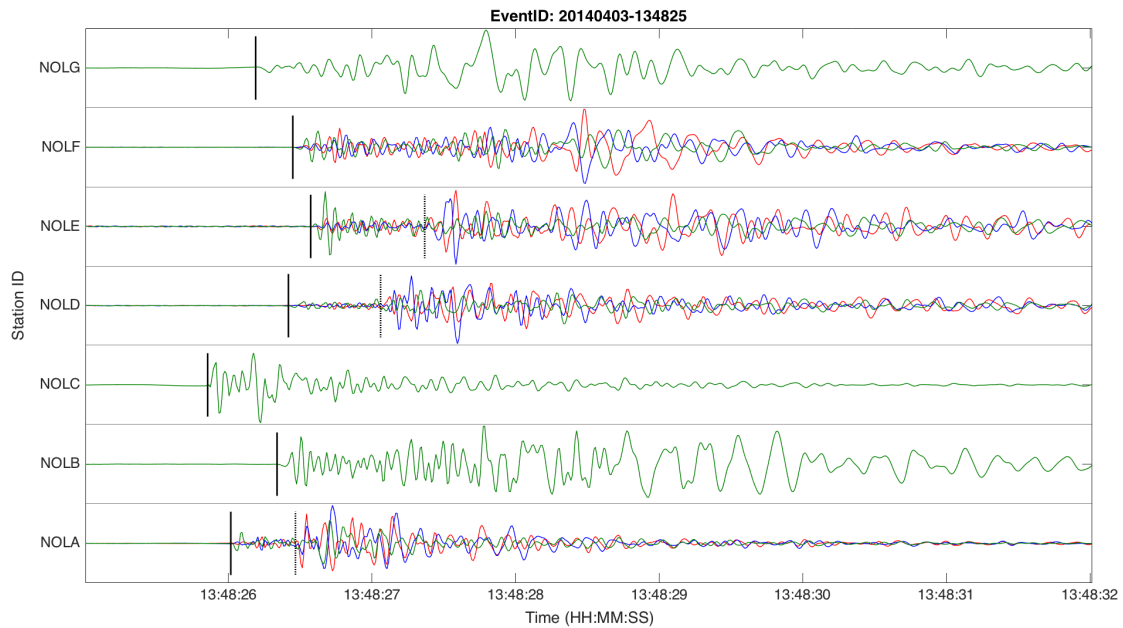
143 *Figure 2: Map of event hypocentres, with events coloured by occurrence date. Also shown are*
 144 *the positions of the monitoring network (triangles) and the mining panels (brown rectangles).*
 145 *Panels DS-4 and DS-5 were active during the monitoring period, and the coloured bars*
 146 *running across these panels show the forward movement of the mining faces with time. The*
 147 *position of the cross-section A – B (Figure 5) is marked by the dashed line.*

148

Layer No.	Depth to Layer Top (m)	V_P (ms^{-1})	V_S (ms^{-1})
1	0	1900	1280
2	60	2750	1540
3	135	3100	1740
4	275	3500	1970
5	1019	4200	2360
6	1351	5250	2920
7	2751	6000	3370

149 *Table 1: 1D, layered, isotropic velocity model used to locate events. Model is based on that used by*
 150 *Bishop et al. (1993).*

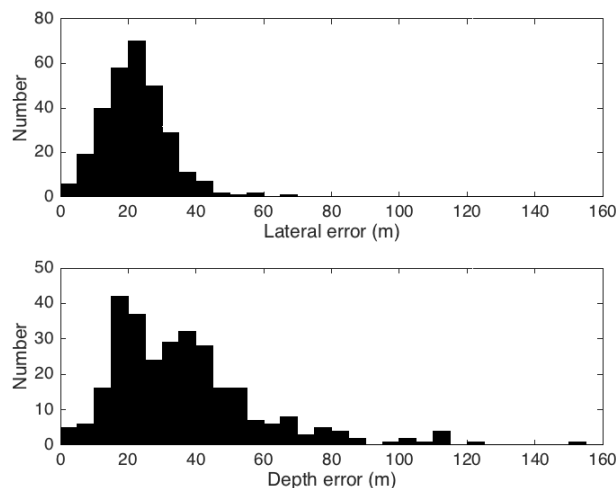
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153 *Figure 3: Recorded waveforms for a larger event ($M_L = 1.3$). The N (red), E (blue) and Z (green)*
 154 *components for each station are overlain. Stations NOLB, NOLC and NOLG are single (Z) component*
 155 *stations. The P- and S-wave picks are marked by the solid and dashed tick marks.*

156 In Figure 4 we show histograms of the event location uncertainties laterally and in depth.
 157 Note that these uncertainties pertain solely to the residuals between picked and modelled
 158 arrival times, and do not account for velocity model uncertainties. The velocity model used is
 159 based on limited site-specific data, relying mainly on regional seismic refraction surveys
 160 (Bishop et al., 1993).



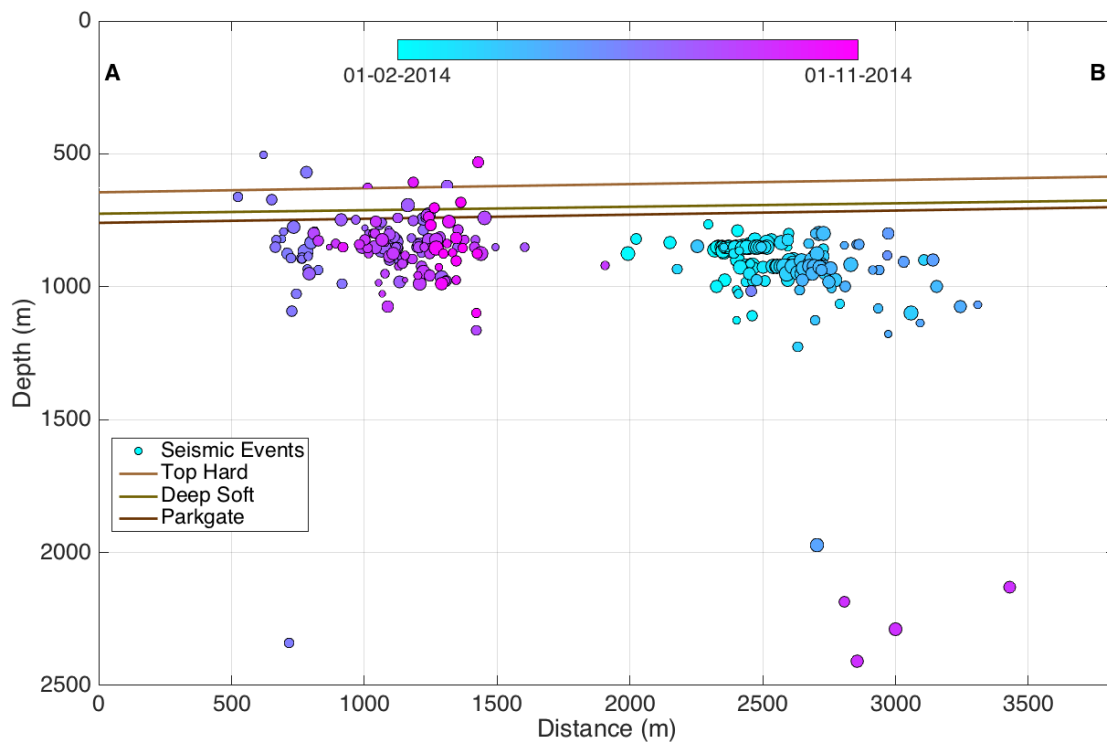
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162 *Figure 4: Histograms showing the lateral and depth uncertainties for the located events.*

163 A brief sensitivity analysis suggested that velocity model uncertainties of up to 10% may
 164 affect depth locations by as much as 150m, while lateral locations are relatively unaffected.
 165 This reflects the geometry of the array, which provides reasonable azimuthal coverage but

166 with surface stations only, such that an uncertain velocity model will primarily affect the
167 event depths.

168 Figure 5 shows a cross-section of event depths relative to the coal seams. We note that, while
169 it appears that the events are located below the seam depths, given the likely velocity model
170 uncertainties, it is not possible to rule out that these events are actually located at the same
171 depths as the Deep Soft Seam being mined.



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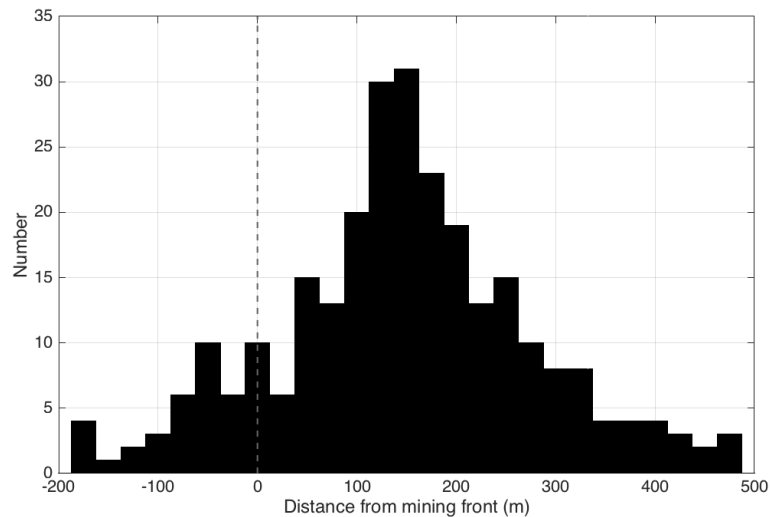
173 *Figure 5: Events depths shown along cross-section A – B (see Figure 2). The positions of the*
174 *Top Hard, Deep Soft and Parkgate Seams are also marked. Note that velocity uncertainties*
175 *mean that the event depths may not be particularly well constrained.*

176 2.2. Event Locations with Respect to Mining Activities

177 The positions of the mining panels, and the progress of the mining face with time, have been
178 provided by the UK Coal Authority in their Mine Abandonment Plans (2017). The position of
179 the mining face with respect to the events can be seen in Figure 2. It is immediately apparent
180 that the event locations are tracking the position of the face as it moves SE along panel DS-4,
181 before switching to DS-5 and again tracking the mining front to the SE. The monitoring
182 period ceases when the events have propagated approximately half-way along the length of
183 panel DS-5.

184 We investigate the position of events in relation to the mining face in greater detail in Figure
185 6, which shows a histogram of event positions relative to the mining face, along an axis
186 parallel to the mining panels. Most events are found to occur ahead of the face, with most

187 events occurring within 300m of the face. This close correlation between events and the
 188 mining face implies that the events are being directly induced by mining activities, as
 189 opposed to the activation of pre-existing tectonic features, in which case we would expect the
 190 events to align along an activated fault. As per the categorisation described by Stec (2007),
 191 we characterise these as mining-induced events.



192

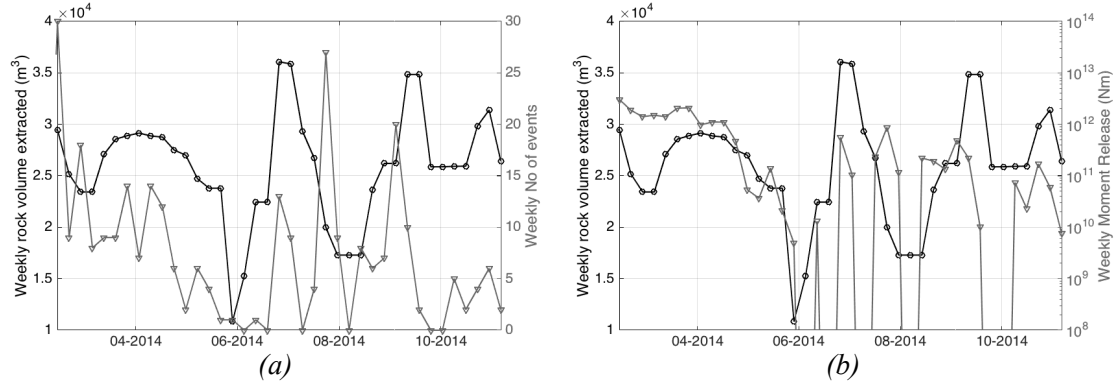
193 *Figure 6: Histograms showing the lateral position of each event relative to the mining face at the time*
 194 *of event occurrence, where a positive distance represents events occurring in advance of the face.*

195 However, we also note small cluster of 5 events that is found at greater depths (>2,000m), to
 196 the SW of the DS-4 panel. 4 of these 5 events occurred within a single 7-hour period.
 197 Establishing the causality of these events is more difficult. It is possible that these events have
 198 been have been triggered by the static transfer of stress changes to greater depths, leading to
 199 fault activation. As per the Stec (2007) categorisation, these may be mining-tectonic events.
 200 However, it is not possible to rule out that these deeper events may in fact have a natural
 201 origin.

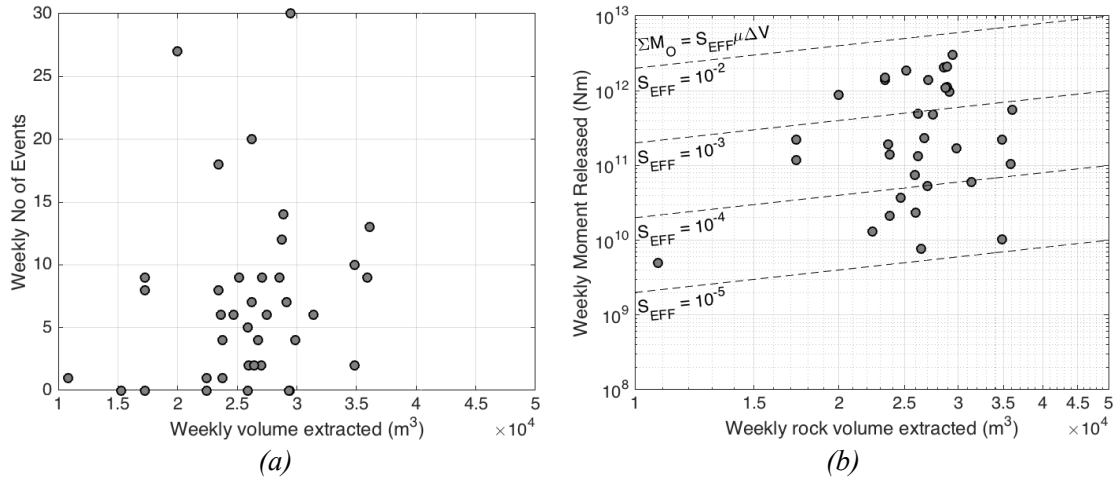
202 **3. CORRELATION BETWEEN SEISMICITY AND MINING RATES?**

203 In Figure 7 we show the volume of rock removed from the mine on a weekly basis (ΔV), the
 204 number of events per week (N_E), and the cumulative seismic moment (ΣM_O) released per
 205 week. The volume of rock removed per week is estimated from the forward progress of the
 206 mining face, multiplied by its dimensions (width and height). To further investigate any
 207 correlation between the extracted volume and seismicity, in Figure 8 we cross-plot these
 208 parameters. From Figure 8 it is apparent that there is little immediate correlation between ΔV
 209 and N_E and ΣM_O on a weekly basis.

210



211 Figure 7: Weekly rock volume extracted (black lines) compared with (a) the weekly number of
 212 recorded events and (b) the weekly cumulative seismic moment released (grey lines).



213 Figure 8: Cross-plots examining potential correlation between weekly rock volume extracted and the
 214 weekly number of recorded events (a) and the weekly cumulative seismic moment released (b). In (b),
 215 the dashed lines show the expected relationship for given values of S_{EFF} .

216 McGarr (1976) posited a linear relationship between ΔV and ΣM_O :

217
$$\Sigma M_O \approx \mu \Delta V, \tag{1}$$

218 where μ is the rock shear modulus. This relationship corresponds to the situation whereby all
 219 of the deformation produced by the volume change is released seismically. In reality, much of
 220 the deformation may occur aseismically. As such, Hallo et al. (2014) proposed a modification
 221 to this relationship via a “seismic efficiency” term, S_{EFF} , which describes the portion of the
 222 overall deformation that is released as seismic energy:

223
$$\Sigma M_O \approx S_{EFF} \mu \Delta V \tag{2}$$

224 In some of the most well-known cases of induced seismicity, values of S_{EFF} have been close
 225 to 1 (e.g. McGarr, 2014). However, these cases represent outliers: during most industrial
 226 operations S_{EFF} is much less than 1 (e.g., Hallo et al., 2014). The dashed lines in Figure 8(b)
 227 show the relationship between ΔV , ΣM_O and S_{EFF} , assuming a generic value of $\mu = 20$ GPa.
 228 We note that the observed moment release rates correspond to values of S_{EFF} between 0.01 to

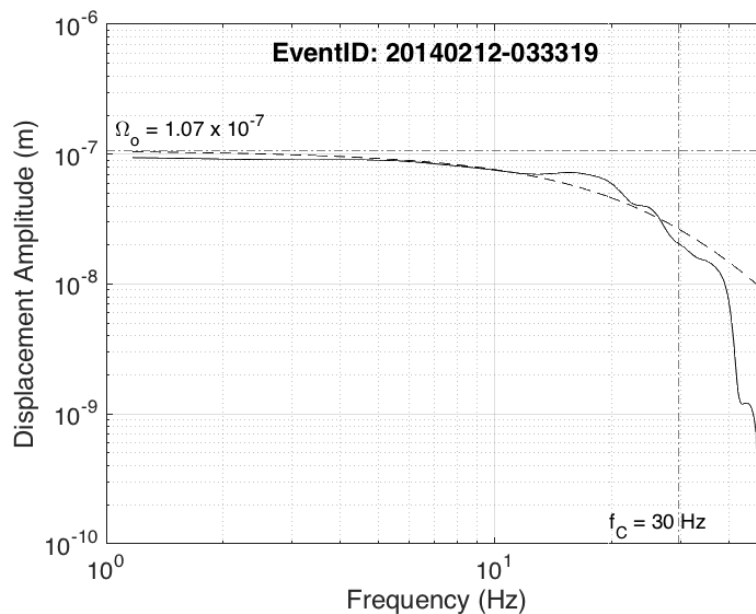
229 0.00001, implying that most of the deformation induced by the mining is released
230 aseismically. This is typical for many cases of seismicity induced by a variety of industrial
231 activities (e.g. Maxwell et al., 2008; Hallo et al., 2014)

232 **4. EVENT MAGNITUDES AND FREQUENCY-MAGNITUDE DISTRIBUTIONS**

233 *4.1 Moment Magnitude Calculation*

234 Local magnitudes for the Thoresby Colliery seismicity have been computed by Butcher et al.
235 (2017), who found that the UK's existing local magnitude scale (Ottmöller and Sargeant,
236 2013) is not appropriate for use when sources and receivers are within a few kilometres of
237 each other. This is because for nearby receivers, the raypath will be predominantly through
238 the softer, more attenuative sedimentary cover, rather than the underlying crystalline crustal
239 rocks, as will be the case for receivers that are more distant to the event. Butcher et al. (2017)
240 have developed an alternative local magnitude scale based on the Thoresby events, which has
241 been recalibrated to ensure consistency between magnitude measurements made on nearby
242 stations and those made using the UK's permanent national monitoring network, the nearest
243 stations of which were some distance from the Thoresby site.

244 However, our aim here is to investigate event magnitude distributions in order to understand
245 the length scales of structures being affected by the mining process. This therefore requires
246 the use of moment magnitudes, since seismic moment can be directly related to rupture
247 dimensions. We compute moment magnitudes by fitting a Brune (1970) source model to the
248 observed S-wave displacement amplitude spectra (Figure 9), following the method described
249 by Stork et al. (2014). The seismic moment is determined from the amplitude of the low-
250 frequency plateau, Ω_0 .



251

252 *Figure 9: Example displacement spectrum used to estimate moment magnitudes. The solid*
253 *line shows the observed spectrum, while the dashed line shows the best-fit Brune (1970)*
254 *source model. The dot-dash lines show the f_C and Ω_0 values for this model.*

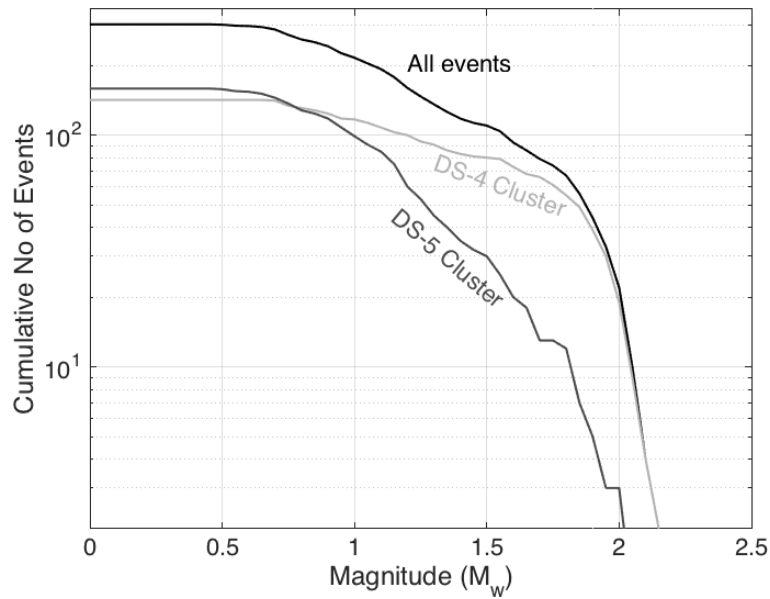
255 Ideally, the measured corner frequency, f_C , of the displacement spectra could be used to
256 determine the rupture length. However, to robustly image the corner frequency, it must be
257 significantly lower than the Nyquist frequency, f_N of the recording system – Stork et al.
258 (2014) recommend that $f_N > 4f_C$ to obtain robust estimates of f_C . The recording systems at
259 Thoresby had sampling rates of 100 Hz, so $f_N = 50$ Hz.

260 We can use generic values for stress drop and rupture lengths to establish the expected corner
261 frequencies for events with $M_W < 1$. Using the relationships between rupture dimensions,
262 seismic moment and stress drop given by Kanamori and Brodsky (2004), assuming a stress
263 drop of 5 MPa and a rupture velocity of 2,000 m/s, the resulting corner frequency $f_C \approx 30$ Hz.
264 Evidently, the $f_N > 4f_C$ criteria is not expected to hold for this particular dataset. However, our
265 observations of event magnitudes, because they are derived from the amplitude spectra at low
266 frequencies, are robust: we therefore use these to make inferences about the length scales of
267 the structures that have generated the observed seismic events.

268

269 *4.2. Frequency-Magnitude Distributions*

270 The observed event magnitude distribution (EMD) is shown in Figure 10. We show the
271 EMDs for the overall dataset, as well as individually for the clusters associated with the DS-4
272 and DS-5 panels. The overall dataset is not well described by the Gutenberg and Richter
273 (1944) distribution $\log_{10} N(M) = a - bM$, where $N(M)$ is the cumulative number of events
274 larger than a given magnitude M , and a and b are constants to be determined. Such a
275 distribution would be represented by a straight line in M vs $\log_{10}(N)$ space. We note that the
276 apparent limit on the largest event size is not an artefact of a short measuring period: while
277 the local array was removed in October 2014, the area continues to be monitored by the BGS
278 National Seismometer Array, which has an estimated detection capability across the UK of
279 magnitude > 2 . Larger events occurring after the study period would therefore be detectable,
280 but no such events have occurred.



281

282 *Figure 10: Observed frequency-magnitude distributions for the full event population (black),*
 283 *as well as for the DS-4 (light grey) and DS-5 (dark grey) clusters individually.*

284 However, fault length and/or earthquake magnitude distributions that are constrained at some
 285 upper limit, leading to a fall-off from the power law relationship at large values, have been
 286 suggested by a number of authors. At the largest scale, Richter (1958) argues that “a physical
 287 upper limit to the largest possible magnitude must be set by the strength of crustal rocks, in
 288 terms of the maximum strain which they are competent to support without yielding”.
 289 Similarly, Pacheco et al. (1992) argue that the rupture dimensions of very large earthquakes
 290 are limited by the thickness of the earth’s seismogenic zone (the portion of the crust that is
 291 capable of undergoing brittle failure). For continental rift zones, Scholz and Contreras (1998)
 292 suggested that the maximum length of normal faults would be limited by the flexural
 293 restoring stress and friction, and found a good match between their model and faults in the
 294 East African Rift and in Nevada. At a much smaller scale, Shapiro et al. (2013) have
 295 suggested these effects will also apply to induced seismicity, with the maximum fault size,
 296 and therefore earthquake magnitude, determined by the dimensions of the volume stimulated
 297 by human activities.

298 To understand the observed EMDs at Thoresby, we consider the statistical distributions of
 299 fault rupture areas that might produce them. Typically, rupture areas are assumed to follow a
 300 self-similar, power law distribution (e.g., Wesnousky et al., 1983; Bonnet et al., 2001). If
 301 stress drops are assumed to be roughly constant (e.g. Abercrombie, 1995) then this power-law
 302 rupture area distribution will result in a power-law distribution of magnitudes, i.e. the
 303 Gutenberg-Richter distribution.

304 A cumulative power law (PL) distribution for rupture area will take the form:

305
$$N(L) = CA^{-\alpha}, \quad (3)$$

306 where $N(A)$ is the number of ruptures with area greater than length A , α is the power law
 307 exponent, and C is a constant. For a PL distribution, there is no upper limit to the maximum
 308 rupture area. Instead, if an upper limit to the rupture area is imposed, for example by the
 309 geometry of the mining panels, then a truncated power law (TPL) distribution results
 310 (Burroughs and Tebbens, 2001; 2002):

311
$$N(A) = C(A^{-\alpha} - A_{MAX}^{-\alpha}), \quad (4)$$

312 where A_{MAX} is the maximum rupture area.

313 To simulate event magnitudes based on rupture area, we use Kanamori and Brodsky (2004):

314
$$M_O = \Delta\sigma A^{3/2}, \quad (5)$$

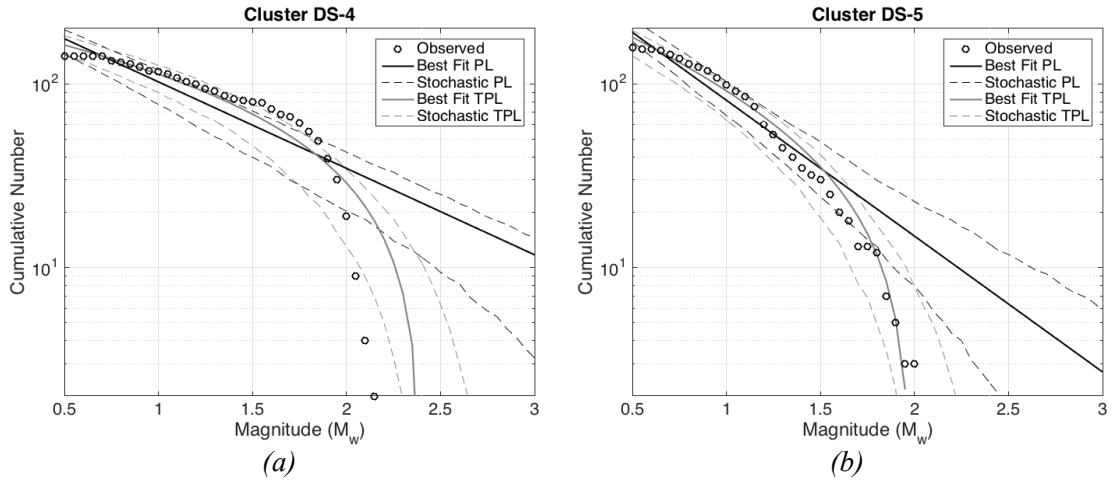
315 where $\Delta\sigma$ is the stress drop. As discussed above, the limitation of a relatively low Nyquist
 316 frequency means that we cannot measure the stress drop directly. Therefore, to estimate the
 317 PL and TPL parameters that best-fit our observations, we initially assume a generic and
 318 arbitrary stress drop of $\Delta\sigma = 5$ MPa.

319 For each of the DS-4 and DS-5 event clusters, we perform a search over the PL and TPL
 320 parameters, finding those that minimise the least-squares misfit between observed and
 321 modelled EMDs. The resulting EMDs are shown as the solid lines in Figure 11, with the PL
 322 and TPL parameters, and the misfit for each of the models, listed in Table 2. The resulting
 323 rupture area distributions are shown in Figure 12.

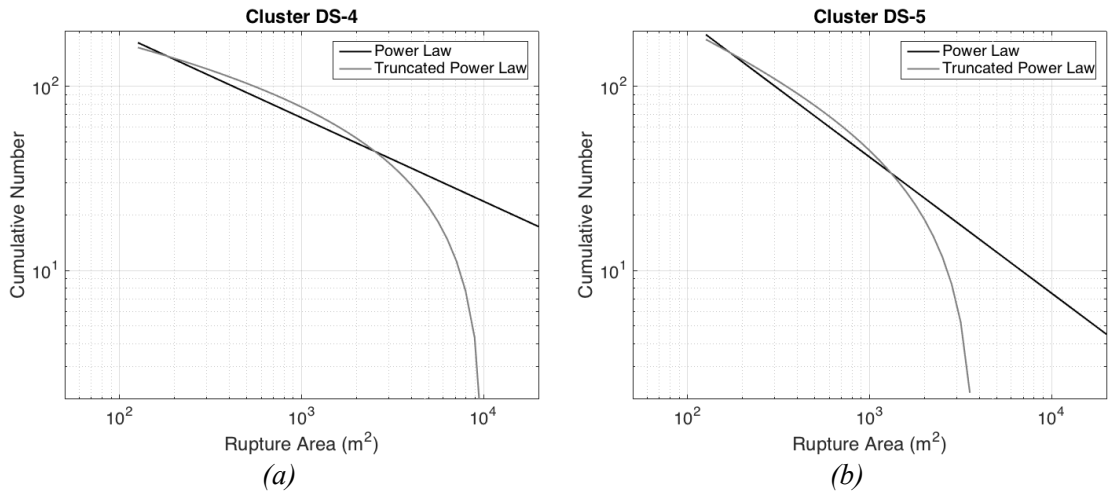
324 Having established the best-fitting PL and TPL distributions with a fixed stress drop value,
 325 we then investigate the impact of a variable range of $\Delta\sigma$. We do this in a stochastic manner,
 326 simulating rupture area distributions based on the PL and TPL parameters, assigning stress
 327 drops randomly from a uniform distribution of $0.1 < \Delta\sigma < 20$ MPa. We repeat this process
 328 over 100 iterations, and in Figure 11 the dashed lines show the range encompassing ± 2
 329 standard deviations around the resulting mean EMD. From Figure 11 we observe that both
 330 event populations are clearly better modelled by a TPL rupture area distribution, even when
 331 stochastic variation in $\Delta\sigma$ is considered.

	<i>Dist. Type</i>	α	C	A_{MAX}	<i>Misfit</i>
DS-4	PL	0.47	1707	NA	5.46
	TPL	0.1	743	10075	1.23
DS-5	PL	0.74	6861	NA	3.05
	TPL	0.38	1536	3870	0.86

332 Table 2: Best fitting power law and truncated power law distributions for each of the DS-4
 333 and DS-5 clusters, and the resulting normalised misfits.



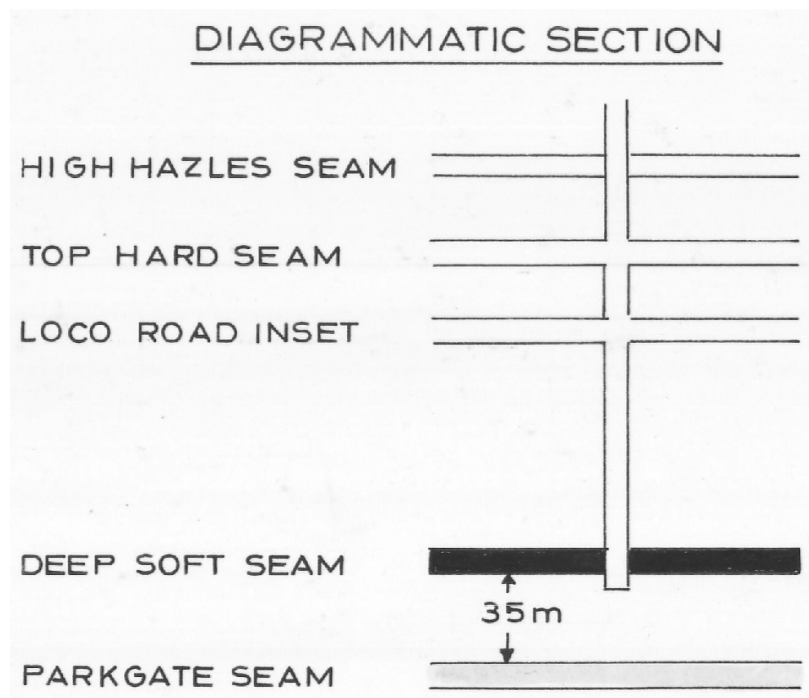
334 Figure 11: Fitting PL (black) and TPL (grey) rupture area distributions to the DS-4 (a) and
 335 DS-5 (b) EMDs. Observed EMDs are shown by black circles. The solid lines show the best
 336 fitting models for a fixed $\Delta\sigma$ value, while the dashed lines show ± 2 standard deviations when
 337 $\Delta\sigma$ is varied stochastically.



338 Figure 12: Best fitting rupture area PL (black) and TPL (grey) distributions for the DS-4 (a)
 339 and DS-5 (b) clusters.

340 Based on these results, it is worth examining whether the best fitting values for A_{MAX}
 341 correspond to any length-scales associated with the mining activities. There are two length
 342 scales in play that might affect rupture dimensions: the width of the mining face
 343 (approximately 300 m); and the separations between (1) the underlying Parkgate Seam, which
 344 is 35 m below the Deep Soft (Figure 13), and (2) the overlying Top Hard Seam, which is
 345 approximately 110 m above the Deep Soft. Both seams have already been mined throughout
 346 our study area. The voids left by the longwall mining of these seams will be filled with goaf,
 347 the rubble and detritus created as the roof collapses behind the mining face. It is difficult to
 348 envisage a mechanism by which ruptures could propagate through such a rubble-filled void.

349 Assuming circular ruptures, areas of 10075 and 3870 m² correspond to rupture radii of 57 and
350 35 m. The larger dimension radius is therefore roughly equivalent to a circular rupture
351 extending from the Deep Soft to the Top Hard. Alternatively, assuming a rectangular rupture,
352 the DS-4 A_{MAX} value could correspond to a rupture with dimensions of approximately 35 ×
353 300 m, equivalent to a rupture extending from the Deep Soft to the Parkgate, across the length
354 of the mined face. In reality, ruptures will not be rectangular nor circular. Nevertheless, the
355 general agreement between the dimensions of the maximum rupture area and these distances
356 leads us to suggest that the presence of the overlying and underlying Top Hard and Parkgate
357 seams is indeed limiting the rupture dimensions. Given the similarities between these
358 dimensions, it is not possible to determine whether one of these features in particular is
359 controlling the maximum rupture area. Indeed, it is likely that all three features: the width of
360 the mining face; the distance to the underlying Parkgate Seam; and the distance to the
361 overlying Top Hard Seam, are all playing a role in limiting the maximum rupture dimensions.



362
363 *Figure 13: Diagrammatic section showing the spacing between the Deep Soft Seam, and the*
364 *underlying Parkgate Seam, which has already been mined out across the study area. Image*
365 *taken from UK Coal Authority Mine Abandonment Plans (2017).*

366

367 5. SEISMIC ANISOTROPY AND SHEAR-WAVE SPLITTING

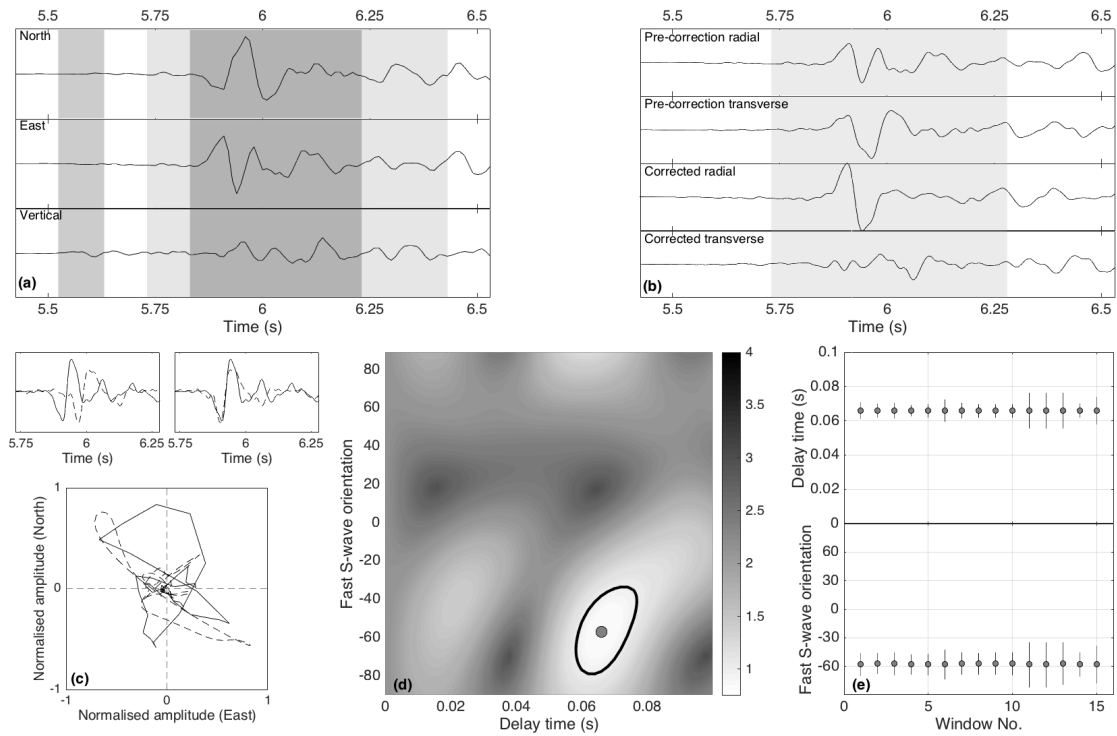
368 Shallow crustal anisotropy can be generated by several mechanisms, including: alignment of
369 macroscopic fracture networks; the preferential alignment of microcracks due to anisotropic

370 stress field (in practice, the microscopic and macroscopic effects often combine, as both
371 larger-scale fracture networks and microcracks are preferentially opened or closed by the
372 same stress field); and by the alignment of sedimentary bedding planes.

373 Shear-wave splitting (SWS), where the velocity of a shear-wave is dependent upon the
374 direction of travel and the polarity of the wave, is an unambiguous indicator of seismic
375 anisotropy, and has been used previously to image stress changes induced by mining activities
376 (Wuestefeld et al., 2011). Shear-waves that propagate near-vertically will not be sensitive to
377 horizontally-layered sedimentary fabrics, which produce Vertically-Transverse-Isotropy
378 (VTI) symmetry systems. Instead, in the absence of other major structural fabrics, the fast
379 shear wave polarisation orientation can be treated as a proxy for the direction of maximum
380 horizontal stress (e.g., Boness and Zoback 2006).

381 We perform SWS measurements on the Thoresby data. Accurate SWS measurements can
382 only be obtained within the “S-wave window” (Crampin and Peacock 2008), because arrivals
383 at an incidence angle greater than $\sim 35^\circ$ from vertical may be disturbed by S-to-P conversions
384 at the free surface. This constraint limits the available data considerably, such that events
385 within the S-wave window are found only on station NOLA, and for only 28 of the recorded
386 events.

387 We perform the SWS measurement using the automated cluster-based approach described by
388 Teanby et al. (2004). Where larger datasets are studied, automated quality assessments such
389 as that described by Wuestefeld et al. (2010) can be used, but in this case, given the small
390 sample size, the quality of measurements were assessed manually. Of the 28 arrivals within
391 the S-wave window at NOLA, 9 provided good-quality, robust results according to the
392 diagnostic criteria specified by Teanby et al. (2004). This is a typical rate-of-return for such
393 studies given the relatively low magnitude (and therefore signal-to-noise) of the events. An
394 example of a robust SWS measurement is provided in Figure 14



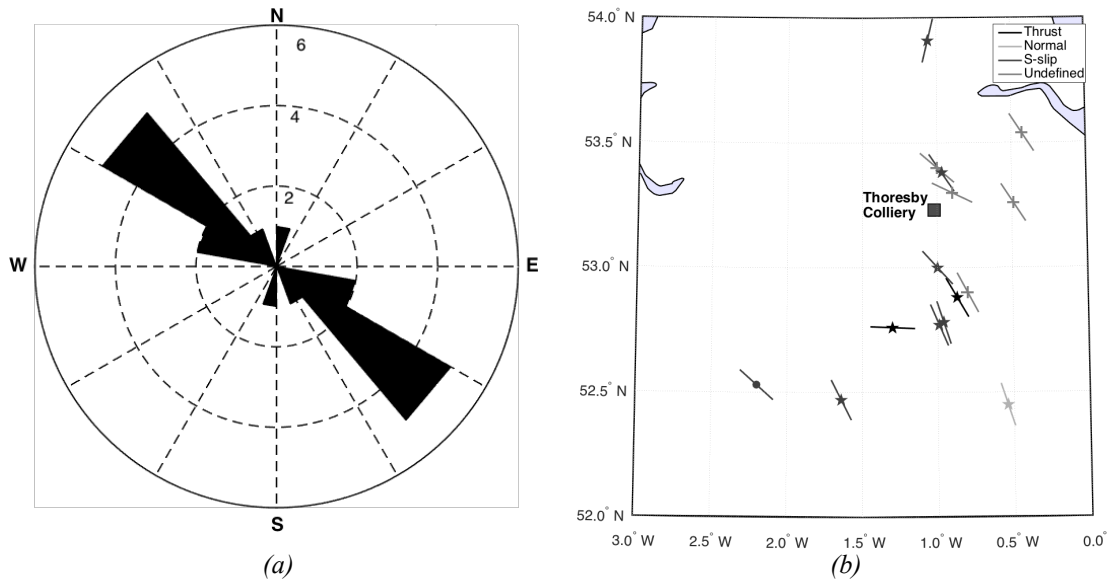
395

396 *Figure 14: Example shear-wave splitting measurement using the method described by Teanby et al.*
 397 *In (a) we plot the N, E and Z components of the recorded waveforms, where P- and S-wave*
 398 *windows are highlighted by the shaded areas. In (b) we plot the radial and transverse components*
 399 *prior to and after the splitting correction, where the aim of the correction is to minimise energy on the*
 400 *transverse component. In (c) we plot the waveform particle motions before (solid lines) and after*
 401 *(dashed lines) correction. In (d) we plot the error surfaces of the correction method as a function of*
 402 *delay time and fast direction normalised such that the 95% confidence interval (highlighted in bold)*
 403 *is 1. In (e) we plot the best-fit delay times and fast directions that result from choosing different S-wave*
 404 *window start and end times (as indicated by the light-grey shaded zone of (a)).*

405 In Figure 15a we show the measured fast directions in the form of an angle histogram. A
 406 dominant fast direction striking NW-SE is clearly observed. The mean fast direction azimuth
 407 is 130° . No temporal variations in SWS fast directions or percentage anisotropy were
 408 observed. The mean delay time was 43 ms, and the mean percentage S-wave anisotropy was
 409 6.8%.

410 In Figure 15b we compare the measured fast S-wave orientations with independent
 411 measurements for S_{Hmax} taken from the World Stress Map database (Heidbach et al., 2008).
 412 These measurements, mainly from borehole breakouts and hydraulic fracturing tests, also
 413 indicate an approximate regional S_{Hmax} strike that is to the NW-SE. We conclude that the
 414 mean measured S-wave fast polarity of 130° can be used as a proxy for S_{Hmax} at this site.

415

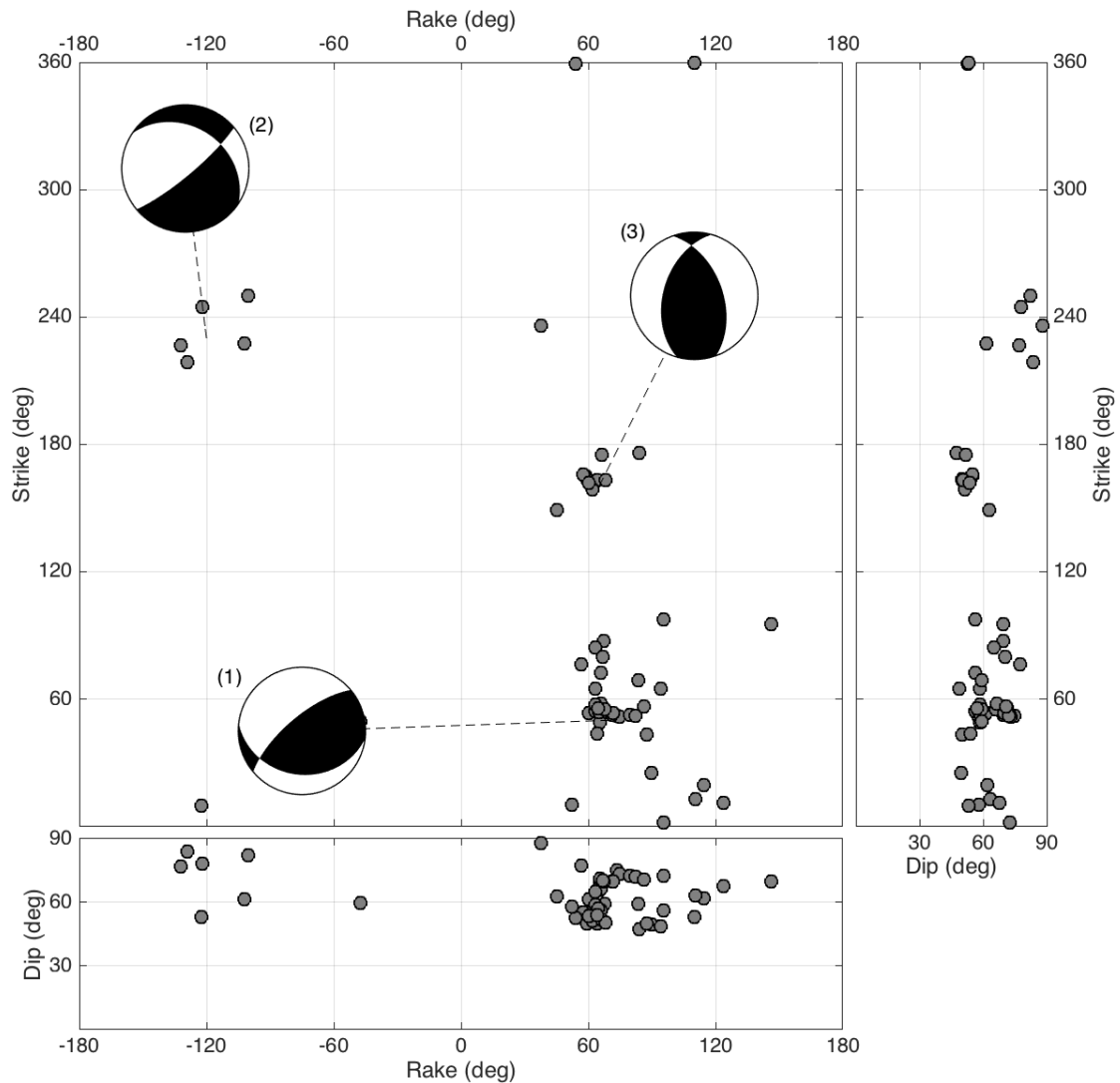


416 Figure 15: SWS and stress anisotropy. In (a) we plot an angle histogram of the measured SWS fast
 417 directions. In (b) we show regional measurements of S_{Hmax} from the World Stress Map database
 418 (Heidbach et al., 2008): '+' symbols represent borehole breakouts, 'o' symbols represent focal
 419 mechanisms, and '★' symbols represent hydraulic fracturing data. The Thoresby site is marked by the
 420 red square. Measurements are coloured by whether they represent a thrust, normal or strike-slip
 421 stress regime (if known).

422 6. SOURCE MECHANISMS

423 We compute event focal mechanisms by inverting the observed P-wave polarities and relative
 424 P-wave, S_H and S_V wave amplitudes for the best fitting double-couple source mechanism. In
 425 doing so, we preclude the possibility of non-double-couple sources in our inversion, as might
 426 be anticipated during mining-induced seismicity. We do this because the monitoring array
 427 consists of only 4 3-C and 3 1-C stations, which limits our ability to robustly constrain non-
 428 double-couple events. However, we note that the recovered mechanisms do a reasonable job
 429 of fitting the observed polarities, i.e. non-double-couple sources do not appear to be necessary
 430 to match the majority of our observations.

431 Of the 305 events, a total of 65 had sufficient signal-to-noise ratios such that P-wave
 432 polarities could be robustly assigned, and produced reliable and consistent source
 433 mechanisms. These strikes, dips and rakes for these events are plotted in Figure 16. We note 3
 434 main clusters of event types, representative source mechanisms for which are also plotted.



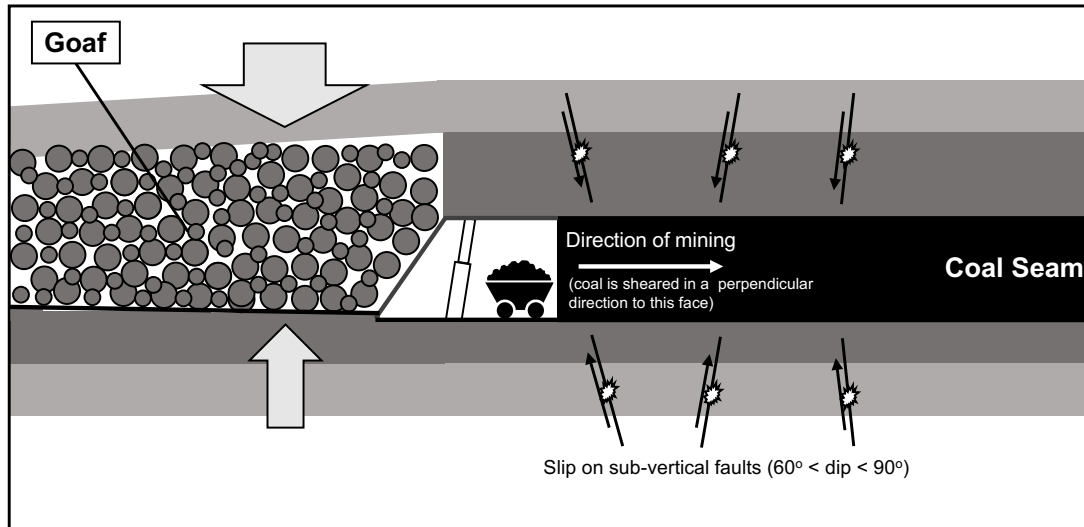
435

436 *Figure 16: Source mechanisms (strike, dip and rake) for each event for which a reliable*
 437 *mechanism could be obtained. Three main clusters of mechanisms can be identified,*
 438 *representative focal spheres for which are shown. These spheres are upper-hemisphere*
 439 *projections where the compressive quadrants are shaded black.*

440 The most common source mechanism type (numbered 1 in Figure 16) consists of events with
 441 strikes of approximately 50° , high angles of dip, and rakes of between $60^\circ - 90^\circ$. This source
 442 mechanism orientation corresponds to near-vertical planes whose strikes match the strike of
 443 the mining face, on which dip-slip movement occurs, with the side of the fault that is towards
 444 the mine moving downwards.

445 A second, less populous source mechanism type (numbered 2 in Figure 16) shows similar
 446 strikes and dips, but with the opposite sense of movement such that the side of the fault
 447 towards the mining face moves upwards. Similar event mechanisms – near-vertical failure
 448 planes striking parallel to the mining face with upward and downward dip-slip motion – were
 449 observed by Bischoff et al. (2010) for longwall mines in the Ruhr Area, Germany, and we

450 share their geomechanical interpretation for these events (Figure 17). As the coal is mined,
 451 the surrounding rock mass will collapse to fill the void. This will result in downward motion
 452 of the overlying rock (as per source mechanism type 1), and upward motion of the underlying
 453 rock (as per source mechanism type 2) along vertical planes that run parallel to the mining
 454 face.



455

456 *Figure 17: Geomechanical interpretation of the observed source mechanisms. As the*
 457 *surrounding rocks move to fill the void created by mining, dip-slip motion occurs on near-*
 458 *vertical slip planes oriented parallel to the mining face. Adapted from Bischoff et al. (2010).*

459 A third type of source mechanism is also observed (numbered 3 in Figure 16), with thrust-
 460 type mechanisms occurring on steeply-dipping planes that strike approximately north-south.
 461 It is possible that they result from the interaction between mining activities and pre-existing
 462 structures in the area, since the N-S orientation of these planes does not match the orientation
 463 of any feature in the mine.

464 Using the source mechanisms for all events, we use the STRESSINVERSE iterative joint
 465 inversion algorithm described by Vavrycuk (2014) to estimate the orientations of principal
 466 stresses and the shape ratio, R (Gephart and Forsyth, 1984):

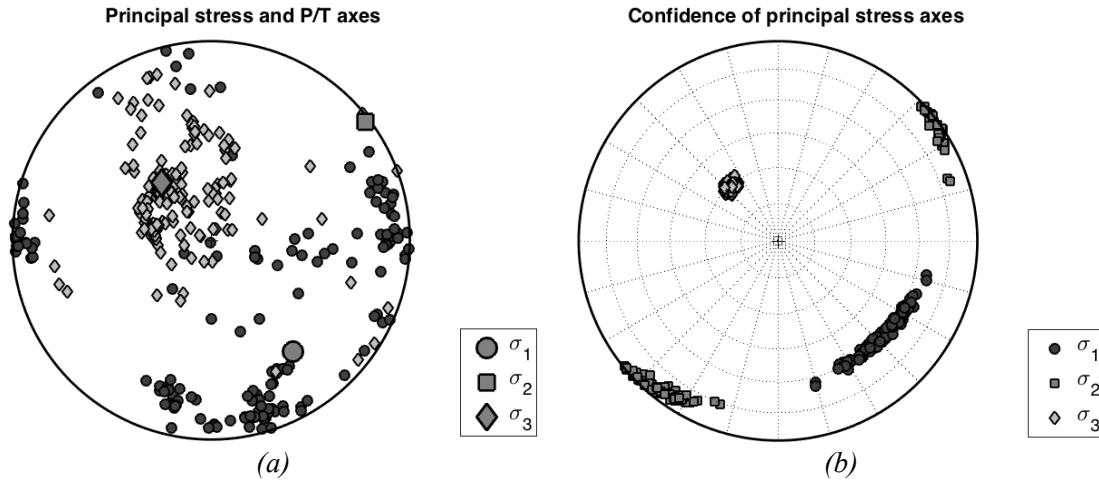
467
$$R = \frac{\sigma_1 - \sigma_2}{\sigma_1 - \sigma_3}, \quad (4)$$

468 where σ_1 , σ_2 , and σ_3 represent the maximum, intermediate and minimum principal stresses.
 469 The results of this inversion are listed in Table 3, and shown in Figure 18. We note that the
 470 resulting maximum horizontal stress is sub-horizontal, with an azimuth of 144° . This is
 471 consistent, within error, with the maximum horizontal stress orientation estimated from SWS
 472 analysis. This implies that, while the orientations of the slip planes are consistent with the
 473 geometry of the mining activities, the resulting deformation is also consistent with the
 474 regional *in situ* stress conditions.

<i>Stress</i>	<i>Azimuth</i>	<i>Plunge (down from horizontal)</i>	<i>Shape Ratio (R)</i>
σ_1	144°	31°	0.17
σ_2	52°	2°	
σ_3	319°	59°	

476 *Table 3: Principal stress orientations and Shape Ratio (R) as inverted from event source mechanisms*

477



478 *Figure 18: Stress tensor inversion results using the STRESSINVERSE algorithm (Vavrycuk,*
479 *2014). In (a) we show a lower hemisphere projection of the P (dark grey \circ) and T (light grey*
480 *\diamond) axes for every event, with the overall estimate for the σ_1 , σ_2 , and σ_3 axes marked by a*
481 *large \circ , \square , and \diamond , respectively. In (b) we show confidence limits for the principle stress*
482 *axes, assuming $\pm 15^\circ$ error in source mechanism orientations.*

483 **6. CONCLUSIONS**

484 In this paper, we characterise the seismicity recorded during longwall mining of the Deep
485 Soft Seam at the Thoresby Colliery, Nottinghamshire, U.K.. A local monitoring network was
486 installed for 8 months, recording 305 events, with the largest event having a local magnitude
487 of $M_L = 1.7$. Event locations are found to track the advance of the mining faces, with most
488 events being located up to 300 m ahead of the face.

489 We conclude that these events are “mining-induced”, i.e. they are directly induced by the
490 mining activity, as opposed to “mining-tectonic” events, which are caused by static stress
491 transfer producing activation of pre-existing tectonic faults. However, comparison between
492 weekly mining rates and the rates of seismic activity do not show strong correlation.
493 Moreover, the amount of deformation released in the form of seismic events is a small

494 percentage of the overall deformation produced by the mining activities (in other words, most
495 of the deformation is released aseismically).

496 Event magnitudes do not follow the expected Gutenberg-Richter distribution. Instead, we find
497 that the observed magnitude distribution can be reproduced by assuming that rupture areas
498 follow a Truncated Power Law distribution, whereby there is a limit to the maximum size of
499 the rupture area. The observed maximum rupture area could correspond to several controlling
500 features around the seam, including the width of the mining face, and the distances to the
501 underlying Parkgate and overlying Top Hard seams, which have already been excavated. Our
502 inference is that the presence of these rubble-filled voids where the excavated seams have
503 been mined out creates a limit to the maximum rupture dimensions.

504 Event source mechanism analysis shows that most events comprise dip-slip motion along
505 near-vertical planes that strike parallel to the orientation of the mining face. This type of
506 deformation is the expected response to the longwall mining process, and has been observed
507 at other longwall mining sites. The observed source mechanisms are also consistent with the
508 orientation of *in situ* regional stresses as inferred from SWS analysis.

509

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514

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