

16 **Abstract**

17 The lapse-rate feedback is the dominant driver of stronger warming in the Arctic than the
18 Antarctic in simulations with increased CO₂. While Antarctic surface elevation has been
19 implicated in promoting a weaker Antarctic lapse-rate feedback, the mechanisms in which
20 elevation impacts the lapse-rate feedback are still unclear. Here we suggest that weaker
21 Antarctic warming under CO₂ forcing stems from shallower, less intense climatological
22 inversions due to limited atmospheric heat transport above the ice sheet elevation and
23 elevation-induced katabatic winds. In slab ocean model experiments with flattened Antarctic
24 topography, stronger climatological inversions support a stronger lapse-rate feedback and
25 annual-mean Antarctic warming comparable to the Arctic under CO₂ doubling. Unlike the
26 Arctic, seasonality in warming over flat Antarctica is mainly driven by a negative shortwave
27 cloud feedback which exclusively dampens summer warming, with a smaller contribution from
28 the winter-enhanced lapse-rate feedback.

29 **Plain Language Summary**

30 Models project stronger surface warming in the Arctic than the Antarctic under climate change.
31 A climate feedback in which more warming occurs near the surface than at higher altitudes in
32 the atmosphere promotes this increased warming in the Arctic. Antarctica's surface elevation is
33 thought to weaken this feedback in comparison to the Arctic, but how this occurs is unclear.
34 Here we show that Antarctic elevation weakens surface warming by changing the base-state
35 vertical temperature structure. When Antarctic topography is flattened in model experiments,
36 Antarctica experiences more warming under climate change, resembling Arctic warming.
37 Similarly to the Arctic, flat Antarctica warms most during the winter, but this seasonality is
38 driven by different climate feedbacks in the Arctic versus Antarctic. These results indicate the
39 importance of base-state temperatures for warming under climate change, and suggest that
40 strong polar amplification is possible without local sea-ice loss.

41 **I Introduction**

42 The Arctic has warmed about twice the global average in recent decades in a pattern
43 known as Arctic amplification (Serreze et al., 2009; Screen and Simmonds, 2010a). In contrast,
44 Antarctic amplification is not observed in the same timeframe, and Antarctic warming is
45 dwarfed by Arctic warming in 21st century projections (Marshall et al., 2015; Smith et al., 2019).
46 While Antarctic warming is delayed in part by Southern Ocean upwelling and associated heat
47 uptake (Collins et al., 2013; Armour et al., 2016), global climate models also project weaker
48 equilibrium warming for the Antarctic than the Arctic (e.g., Danabasoglu and Gent, 2009). In an
49 evaluation of models in the fifth phase of the Coupled Model Intercomparison Project (CMIP5),

50 Goosse et al. (2018) find that the lapse-rate feedback is the single greatest factor contributing
51 to this hemispheric asymmetry in polar warming.

52 In a warming climate, a positive polar lapse-rate feedback results from stable
53 temperature inversions which contribute to stronger warming near the surface than aloft,
54 leading to inefficient longwave emission to space. In the Arctic, the ice-albedo feedback
55 promotes surface warming and thus contributes to a more-positive lapse-rate feedback (Feldl et
56 al., 2017; Graversen et al., 2014). Dai et al. (2019) have more recently argued that Arctic sea-
57 ice loss enables the lapse-rate feedback via increased turbulent heat fluxes and upward
58 longwave radiation, enhancing lower-tropospheric warming over newly opened ocean. These
59 results align with evidence that sea-ice loss promotes seasonality in Arctic warming by
60 enhancing winter heat transfer from the ocean to the atmosphere and strengthening winter
61 longwave radiative feedbacks (e.g., Screen and Simmonds, 2010b; Bintanja and van der Linden,
62 2013). Given the proposed dependence of Arctic lower-tropospheric warming and lapse-rate
63 feedback on sea-ice loss, a weaker Antarctic lapse-rate feedback may be driven by the
64 persistence of the Antarctic ice sheet.

65 Using fully-coupled Community Earth System Model (CESM) experiments, Salzmann
66 (2017) instead finds that Antarctic elevation drives the weaker Antarctic lapse-rate feedback. As
67 the sign and magnitude of the lapse-rate feedback depends on base-state static stability (Cronin
68 and Jansen, 2015; Payne et al., 2015), we expect Antarctic elevation to control the lapse-rate
69 feedback through impacts on base-state inversions. Confinement of the radiatively-active
70 atmospheric column over Antarctica to a shallower layer than in the Arctic may additionally
71 weaken the lapse-rate feedback.

72 With the mechanisms linking Antarctic elevation to a weaker lapse-rate feedback still
73 unclear, we investigate hemispheric asymmetry in the polar lapse-rate feedback using slab ocean
74 CESM experiments with present-day and flattened Antarctic topography under preindustrial
75 and doubled CO₂. We analyze Antarctic elevation impacts on climatological inversions, the
76 lapse-rate feedback, and polar amplification, in addition to investigating the seasonality of
77 warming in the flat Antarctic compared to the Arctic. These experiments provide insight into
78 hemispheric polar warming differences, with an ultimate goal of better understanding the
79 mechanisms behind polar amplification.

80 **2 Data and Methods**

81 2.1 Model Experiments

82 To investigate hemispheric asymmetry in the lapse-rate feedback, we use CESM (Hurrell
83 et al., 2013) version 1.2.2 with the Community Atmosphere Model version 4 (CAM4; Neale et
84 al., 2013). CAM4 has a horizontal resolution of 0.9° × 1.25° and 26 vertical levels, with vertical
85 grid spacing further described in Richter et al. (2014). For all experiments, CAM4 is coupled to
86 a slab ocean forced with a spatially heterogeneous monthly climatology of ocean heat flux
87 convergence derived from a fully-coupled preindustrial control simulation (Bitz et al., 2012).
88 Differing from Salzmann (2017) in the use of a slab ocean rather than fully-coupled model, these
89 experiments isolate the role of Antarctic elevation in determining the equilibrium climate
90 response to CO₂ forcing and exclude potential effects of ocean heat uptake changes on the
91 lapse-rate feedback (e.g., Po-Chedley et al., 2018; Singh et al., 2018). CESM also uses the
92 Community Land Model version 4 (CLM4; Oleson et al., 2010) and the Los Alamos Sea Ice
93 Model version 4 (CICE4; Hunke and Lipscomb, 2008).

94 We perform CESM experiments with present-day topography, referred to as the
95 control Antarctic, or with the elevation of Antarctica flattened to 0 m above sea level, referred
96 to as the flat Antarctic. Branching from preindustrial runs for the control and flat Antarctic,
97 CO₂ is abruptly doubled from 285 to 570 ppm. We run each simulation for 50 years, and
98 calculate climatologies over the final 30 years. Net top-of-atmosphere (TOA) energy imbalances
99 of -0.09 W/m² (control, preindustrial), -0.06 W/m² (control, doubled CO₂), -0.14 W/m² (flat,
100 preindustrial), and -0.09 W/m² (flat, doubled CO₂) indicate near equilibrium for all experiments
101 in this period.

102 2.2 Radiative Feedbacks

103 We calculate the lapse-rate, Planck, surface-albedo, water-vapor, and cloud radiative
104 feedbacks using the radiative kernel method (Shell et al., 2008; Soden et al., 2008). This method
105 calculates radiative feedbacks as the product of (a) the change in radiative flux per unit change
106 in a given climate variable, termed the radiative kernel, and (b) the modeled change in this
107 climate variable normalized by the surface temperature change. While the Planck feedback is
108 defined by propagating the surface temperature change through the entire troposphere, the
109 lapse-rate feedback calculates the effect of departures from this vertically uniform temperature
110 change. To calculate cloud feedbacks, the kernel method is used to determine the effect of
111 noncloud variables (temperature, water vapor, and surface albedo) on the change in cloud
112 radiative forcing (Δ CRF), and this cloud masking effect is subtracted from the total Δ CRF (Shell
113 et al., 2008; Soden et al., 2008). We calculate the residual term by subtracting the kernel-
114 estimated TOA radiation change from the modeled TOA radiation change, normalizing by the
115 surface temperature change.

116 Following Goosse et al. (2018), we also calculate feedback contributions to polar
 117 warming (ΔTS) by dividing the energetic contribution of each feedback ($\lambda_i \Delta TS$), CO₂ forcing (F),
 118 change in atmospheric heat transport convergence (ΔAHT), and residual term (R ΔTS) by the
 119 magnitude of the Planck response in either the control Antarctic (Figure 3b) or control Arctic
 120 (Figure 3c) ($\lambda_{p,c}$), where $\lambda'_p = \lambda_p - \lambda_{p,c}$ is the difference between the local Planck response
 121 (λ_p) and $\lambda_{p,c}$:

$$122 \quad \Delta TS = -F/\lambda_{p,c} - \lambda'_p \Delta TS/\lambda_{p,c} - \sum_i \lambda_i \Delta TS/\lambda_{p,c} - \Delta AHT/\lambda_{p,c} - R\Delta TS/\lambda_{p,c}. \quad (1)$$

123 For this study, changes in climate variables are obtained from the doubled CO₂ minus
 124 preindustrial simulations, and radiative kernels are taken from Shell et al. (2008), calculated with
 125 an offline radiative transfer version of the Community Atmospheric Model, version 3 (CAM3).
 126 Feedbacks are calculated with respect to local rather than global surface temperature change
 127 (Armour et al., 2013; Feldl & Roe, 2013), and Arctic and Antarctic regional averages are defined
 128 for 70 to 90 °N and °S, respectively.

129 For the flat Antarctica experiment, we apply zonally averaged kernels from the Arctic
 130 between 60 and 90°N to the flat Antarctic between 60 and 90°S, shifted by six months. As the
 131 standard kernels do not exist below the Antarctic surface elevation, this allows for feedback
 132 calculation with the assumption that climate variable effects on TOA radiation for the flat
 133 Antarctic are similar to the Arctic. We note that radiative kernels are affected by climatological
 134 characteristics such as cloud patterns and land-ocean distributions, which differ between
 135 hemispheres. However, the flipped kernel method is supported by similar climatological winter
 136 temperature profiles in the flat Antarctic and control Arctic, with similar inversion depth and
 137 intensity (Section 3.1) and representative profiles (Figure S1a). We note that summer
 138 temperatures are colder in the flat Antarctic than the Arctic (Figure S1b) likely due to the

139 absence of leads over Antarctica. Nevertheless, both summer and winter demonstrate strong
 140 agreement between the modeled and kernel-estimated longwave (LW) TOA radiation change
 141 in the flat Antarctic (Figure S2). Annual LW and shortwave (SW) kernel-estimated changes in
 142 TOA radiation approximate the modeled TOA radiation change for the flat Antarctic almost as
 143 well as the control Antarctic experiments (Figure S3), further supporting this flipped kernel
 144 method. An alternative estimation of the SW feedbacks using the approximate partial radiative
 145 perturbation (APRP) technique (Taylor et al., 2007) produces control and flat Antarctic
 146 feedbacks that are qualitatively similar to those calculated using the kernel method (Figure S9).

147 2.3 Atmospheric Heat Transport

148 We calculate AHT across each latitude using the poleward integral of the difference
 149 between TOA and surface energy fluxes (and seasonal atmospheric energy storage for seasonal
 150 averages) from monthly model output (e.g., Kay et al., 2012; Donohoe et al., 2020a). We
 151 calculate latent heat transport as the poleward integral of evaporation minus precipitation at a
 152 given latitude, multiplied by the latent heat of vaporization, and dry static AHT as the residual
 153 between total AHT and latent AHT. As in Kay et al. (2012), we estimate polar AHT
 154 convergence by dividing the change in AHT at the edge of polar regions by the surface area of
 155 these regions. The change in AHT is then normalized by the Planck feedback for the AHT
 156 warming contribution and by the polar surface temperature change for the AHT feedback.

157 We also determine the vertical structure of AHT due to stationary eddies (SE) and
 158 transient eddies (TE) from CESM monthly output at 60°N and 60°S. For the vertical structure
 159 of TE transport, we neglect the potential energy term in Donohoe et al. (2020), which is small
 160 at high latitudes. Vertical structures of SE and TE are calculated as:

$$161 \quad SE = [\bar{V}^* \overline{MSE}^*] \quad (2)$$

162 and

$$163 \quad TE = [c_p(\overline{VT} - \overline{V}\overline{T}) + L(\overline{VQ} - \overline{V}\overline{Q})], \quad (3)$$

164 where V is meridional velocity, MSE is moist static energy, c_p is the specific heat of air, L is the
 165 latent heat of vaporization of water, T is atmospheric temperature, and Q is specific humidity.
 166 Square brackets indicate zonal averages, overbars indicate monthly means, and asterisks denote
 167 departures from the zonal mean.

168 For comparison with CESM, we include vertical profiles of AHT and temperature from
 169 the ERA-Interim (Dee et al., 2011) and NCEP-NCAR (Kalnay et al. 1996) reanalyses, using 6-
 170 hourly fields for 1979-2018 to calculate SE and TE following Donohoe et al. (2020).

171 2.4 Inversion Depth and Intensity

172 Following Zhang et al. (2011), we define surface-based inversion depth as the
 173 geopotential thickness between the surface pressure and the first pressure level above which
 174 temperature decreases with height, and inversion intensity as the difference in temperature
 175 between these two levels.

176 3 Results

177 3.1 Climatological polar inversion asymmetry

178 Antarctic surface elevation drives modeled differences in base-state surface inversion
 179 depth and intensity between the poles. Using the CESM preindustrial control simulation during
 180 Arctic winter, Figure 1a shows a schematic of radiative-advective equilibrium, which controls
 181 base-state inversions in polar regions (Cronin and Jansen, 2015; Payne et al., 2015). In this
 182 framework, LW cooling and weak surface solar absorption promote cold near-surface
 183 temperatures, while poleward AHT (shown at 60°N) maintains warmer temperatures aloft. As
 184 is also seen in the NCEP and ERA-Interim reanalyses (Figure S4), AHT supporting these

185 inversions maximizes in the lower troposphere (around 900 hPa), near the tropospheric
186 temperature maximum. We propose that surface elevations above this level of maximum AHT
187 promote shallower, radiatively-driven winter inversions in the Antarctic (Figure 1b).

188 This preindustrial temperature profile at 90°S demonstrates CESM's ability to resolve
189 shallow surface inversions (here 50 hPa deep) matching radiosonde observations over the
190 Antarctic Plateau, although the model underestimates the intensity of observed mean
191 wintertime inversions which exceed 20 K at South Pole Station (Hudson & Brandt, 2005).
192 Weaker winter inversions over coastal slopes in CESM (Figure 2a; Figure S5d) align well with
193 coastal Antarctic radiosonde observations, which indicate typical depths shallower than 300 m
194 and intensities less than 5 K (Zhang et al., 2011).

195 Considering CESM preindustrial winter inversion depth over the entire control
196 Antarctic ice sheet, the deepest inversions are found where surface elevations are below the
197 level of maximum poleward AHT (black contour, Figure 2a). Shallower winter inversions exist
198 over the Antarctic Plateau, and inversion depth generally increases with increasing surface
199 pressure over Antarctica (Figure 2b). Points with shallow inversions for a given surface pressure
200 are located on (grey points, Figure 2b) or at the bottom of (boxed points, Figure 2b; stippling,
201 Figure 2a) steep slopes (Figure S5a), where inversions may be disturbed by mixing due to
202 katabatic winds (Vihma et al., 2011).

203 In contrast to Antarctica, most points in the Arctic are at higher surface pressures,
204 allowing AHT to support deep inversions (Figure 2c, S5c). Inversions are shallower over the
205 Greenland ice sheet, particularly over steep slopes, and in regions with low sea-ice fraction.
206 While intense, shallow inversions (Figure S5d) exist at high Antarctic elevations, many areas
207 show weaker inversion intensity than the Arctic (Figure S5f), especially over steep Antarctic

208 slopes. These inversions in CESM are consistent with radiosonde observations of generally
209 shallower, weaker inversions in the Antarctic compared to the Arctic (Zhang et al., 2011).

210 In comparison to the control Antarctic, inversions in the flat Antarctic preindustrial
211 experiment are intensified and deepened, particularly over steep slopes (Figures S5b,e, S6a,b).
212 An average inversion depth of 583 m and intensity of 7.5 K over the control Antarctic increases
213 to 846 m and 11.9 K over the flat Antarctic, more comparable with the average inversion depth
214 of 854 m and intensity of 9.0 K in the Arctic. We next investigate the extent to which these
215 stronger, deeper climatological inversions in the flat Antarctic simulation, which more closely
216 resemble Arctic inversions, may promote an Arctic-like lapse-rate feedback and stronger
217 Antarctic amplification.

218 3.2 Polar amplification and lapse-rate feedback asymmetry

219 In contrast to control Antarctic warming (5.4 K) under doubled CO₂, flat Antarctic
220 warming (6.3 K) is more comparable to control Arctic warming (6.7 K; Figure 3a) in the annual
221 mean. Applying Eq. (1), we find that the largest contributor to stronger amplification in the flat
222 Antarctic is the lapse-rate feedback (Figure 3b). A more-positive water-vapor feedback and
223 ΔAHT also support enhanced flat Antarctic warming, while more-negative Planck and cloud
224 feedbacks and a less-positive albedo feedback oppose flat Antarctic warming. The LW residual
225 term is similar for the control and flat Antarctic, although the SW residual term is slightly larger
226 in the flat Antarctic (Figure S7a). Flat Antarctica experiences lapse-rate, water-vapor, and Planck
227 warming contributions nearly equivalent to the Arctic, while a larger ΔAHT and residual term
228 in the flat Antarctic partly balance more-positive albedo and cloud feedbacks in the Arctic
229 (Figure 3c).

230 Figures S7b-g show the response to CO₂ doubling for flat minus control Antarctic
231 topography experiments for various climate variables relevant for feedbacks (left) and the
232 difference for the radiative kernels (right), where T is atmospheric temperature, TS is surface
233 temperature, q is specific humidity, and α is albedo. The strengthened lapse-rate feedback in the
234 flat Antarctic is supported both by more surface-trapped warming in response to CO₂ doubling
235 (more-negative $\Delta(T-TS)$), and by a more-negative T kernel due to deepening and moistening the
236 atmospheric column (Figure S7b,c). While it is difficult to quantify exactly how much of the
237 increased lapse-rate feedback in the flat versus control Antarctic stems from the deeper
238 troposphere alone, we find that only 26% of this increase can be obtained from simply
239 extending the control Antarctic $\Delta(T-TS)$ from the lowest elevations above the ice sheet to all
240 pressure levels below the ice sheet and multiplying by the flat Antarctic temperature kernel.
241 This suggests that the increased lapse-rate feedback over flat Antarctica is strongly driven by
242 enhanced surface-trapped warming, rather than simply a deeper atmospheric column.

243 Surface-trapped warming enhancement for the flat Antarctic appears largest over the
244 Transantarctic Mountains and coastal slopes, where steep slopes promote shallow, weak
245 inversions in the control topography experiment (Figure S7b). Particularly for these regions,
246 increased preindustrial inversion depth and intensity in the flat experiment support stronger
247 low-level warming (Figure S6) and an enhanced lapse-rate feedback.

248 Both the stronger water-vapor feedback and weaker surface-albedo feedback in the flat
249 Antarctic are supported by deepening and moistening the atmospheric column. This amplifies
250 the greenhouse effect of water vapor by increasing the column-integrated specific humidity and
251 strengthening the water-vapor kernel (Figure S7d,e). The deeper atmospheric column over flat
252 Antarctica also dampens surface albedo impacts on TOA radiation, and this weaker (less-

253 negative) albedo kernel (Figure S7g) drives a weakened albedo feedback in the flat Antarctic.
254 Since feedbacks are normalized by local surface temperature change, the more-negative flat
255 Antarctic Planck feedback shown in Figure S7a can be attributed to a stronger atmospheric
256 temperature kernel due to a warmer, deeper emitting column.

257 The more-negative cloud feedback in the flat versus control Antarctic is explained by
258 more-negative SW cloud forcing (Figure S8a; Figure S1 Iaa). The flat Antarctic experiences a
259 larger increase in cloud cover and cloud water path due to increased low-level liquid-bearing
260 clouds under CO₂ doubling (Figure S8d,e), supporting a stronger negative SW cloud feedback.
261 Preindustrial cloud water path (not shown) is also larger for the flat Antarctic, likely due to an
262 extended lower atmosphere, enhanced water-vapor transport, and increased stability. Areas of
263 stronger sea-ice loss (Figure S8c) for the flat Antarctic also correspond to increased cloud
264 water path and SW forcing.

265 Increased preindustrial AHT toward the Antarctic balances stronger cooling to space
266 over flat Antarctica (Figure S10a), consistent with Singh et al. (2016) and Salzmann (2017).
267 Under doubled CO₂, AHT to the flat Antarctic also increases more, largely due to enhanced
268 southward latent AHT (Figure S10b). Δ AHT and feedback calculations for these slab ocean
269 experiments using the kernel method produce similar results to Salzmann (2017), who employ
270 partial radiative perturbation (PRP) feedback computations for the transient response to CO₂
271 doubling in fully-coupled runs: flat Antarctica experiences stronger lapse-rate and water-vapor
272 feedbacks and increased poleward Δ AHT, opposed by more-negative Planck and cloud
273 feedbacks and a less-positive albedo feedback.

274 3.3 Radiative feedback contributions to seasonality in Arctic and Antarctic amplification

275 Enhanced warming in the flat Antarctic compared with the control topography
276 experiment occurs predominantly during the winter season (Figure 4a-c), with similar summer
277 warming for flat and control experiments. This produces a greater difference between winter
278 and summer warming (4.0 K) in the flat Antarctic compared to the elevated Antarctic (2.1 K),
279 more comparable with control Arctic warming seasonality (5.3 K). To investigate what causes
280 enhanced seasonality in warming in the flat Antarctic experiment, we compare winter and
281 summer climate feedbacks. Feedbacks shown are normalized by the annual-mean local warming,
282 although we find similar results for normalizing by seasonal warming. Figures 4d-i highlight
283 seasonal lapse-rate and SW cloud feedbacks, with seasonality for all feedbacks shown in Figure
284 S11.

285 In the Arctic, the lapse-rate feedback strongly promotes greater warming in winter than
286 summer (Figure 4f), consistent with Pithan and Mauritsen (2014). This enhanced winter lapse-
287 rate feedback is driven by stronger base-state inversions in winter compared to summer, when
288 sea-ice melting keeps surface temperatures near the freezing point. While strong lapse-rate
289 feedback seasonality occurs over Southern Ocean sea ice, seasonality in this feedback is weaker
290 over Antarctica for both the control and flat experiments. Flat Antarctica experiences stronger
291 seasonality in warming than control Antarctica due in part to strengthened winter inversions,
292 but the primary feedback enhancing flat Antarctic warming seasonality is the SW cloud
293 feedback.

294 Due to polar night during winter, enhanced SW cloud cooling in the flat Antarctic
295 occurs exclusively during summer (Figure 4h). While the lapse-rate and water-vapor feedbacks
296 promote greater warming for the flat Antarctic throughout the year, greater warming
297 seasonality over flat Antarctica results largely from the SW cloud feedback damping only

298 summertime warming. About two thirds of the difference in SW cloud feedback seasonality
299 between control and flat Antarctica arises from cloud radiative forcing, with the other third
300 coming from differences in the cloud masking term. A weaker summer surface albedo feedback
301 over flat Antarctica also contributes slightly to stronger winter versus summer warming, with
302 stronger albedo feedback seasonality changes over sea ice (Figure S1 I d).

303 **4 Conclusions**

304 With a goal of understanding the mechanisms driving lapse-rate feedback differences
305 between the Arctic and Antarctic, we compare CESM slab ocean experiments with control and
306 flattened Antarctic topography under preindustrial and doubled CO₂ forcing. We find
307 climatological differences in CESM preindustrial Arctic and Antarctic inversions, supported by
308 radiosonde observations. Limited poleward atmospheric heat transport above ice sheet
309 elevations, in addition to mixing due to katabatic winds on steep slopes, drives shallower,
310 weaker inversions for the Antarctic than the Arctic. Combined with weaker emission from the
311 relatively shallow and dry Antarctic atmospheric column, these weaker Antarctic inversions
312 prohibit the strong positive lapse-rate feedback seen in the Arctic. In contrast, over the flat
313 Antarctic, stronger climatological inversions support a lapse-rate feedback that is comparable to
314 that in the Arctic, with nearly equivalent degrees of polar amplification in each hemisphere in
315 these slab ocean experiments. While the seasonality of warming in the flat Antarctic is also
316 more comparable to the Arctic, this seasonality is supported by different climate feedbacks at
317 each pole: lapse-rate feedback seasonality contributes strongly to enhanced winter warming in
318 the Arctic, while a negative SW cloud feedback that only applies to non-winter months
319 contributes more to flat Antarctic seasonality.

320 As noted by Salzmann (2017), enhanced Antarctic amplification in flat Antarctica
321 experiments suggests that reduced Antarctic surface elevation due to mass loss would
322 accelerate Antarctic amplification under climate change. In the context of previous studies
323 considering the relative roles of sea ice and the lapse-rate feedback for polar amplification (e.g.,
324 Graversen et al., 2014; Dai et al., 2019), investigating Antarctic amplification can also provide
325 insight into mechanisms supporting Arctic amplification. Flat Antarctic experiments here and in
326 Salzmann (2017) demonstrate that strong, Arctic-like polar amplification is possible without
327 local sea-ice loss, although non-local sea-ice loss in the Southern Ocean likely contributes to
328 Antarctic amplification through changes in AHT. Our experiments additionally indicate that the
329 strongest seasonality in the lapse-rate feedback occurs in regions with sea ice: even over the
330 flattened Antarctic continent, lapse-rate feedback seasonality pales in comparison to lapse-rate
331 feedback seasonality over the Southern Ocean and Arctic. In line with Cronin and Jansen (2015)
332 and Payne et al. (2015), enhanced flat Antarctic warming supports the dependence of the lapse-
333 rate feedback on base-state inversions, which themselves depend on sea-ice concentration and
334 surface albedo in the Arctic. Further work to disentangle the lapse-rate feedback from sea-ice
335 effects may clarify how they will change separately and in tandem under climate change.

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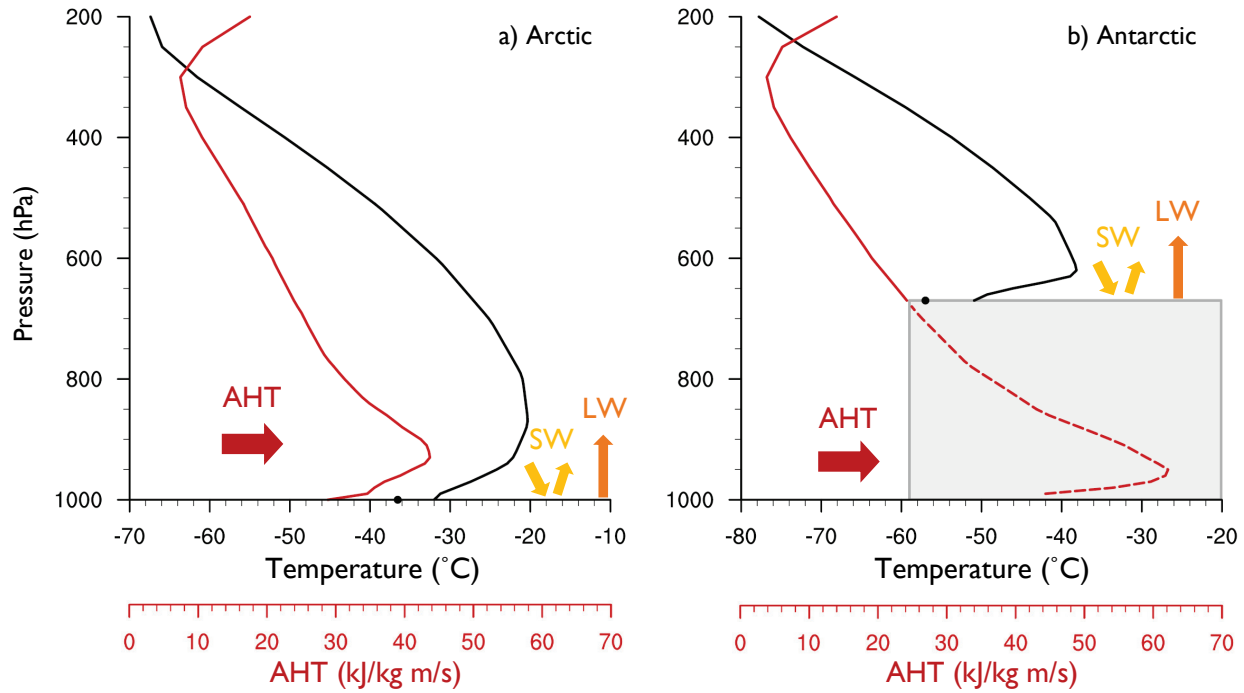


Figure 1. Radiative-advective equilibrium schematic using the CESM preindustrial control topography temperature and AHT. Vertical profiles of winter poleward AHT due to transient and stationary eddies at (a) 60°N and (b) 60°S (red; kJ/kg m/s) and winter temperature at (a) 90°N and (b) 90°S (black; dot for surface temperature; °C).

a) Surface inversion depth (m)

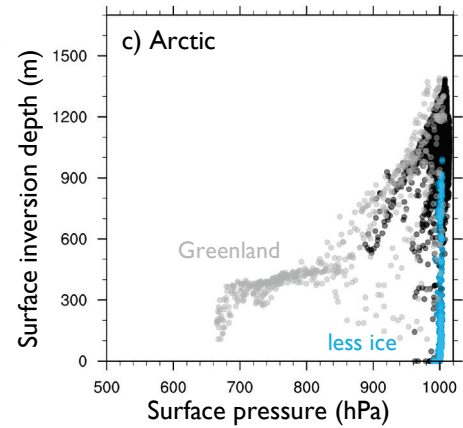
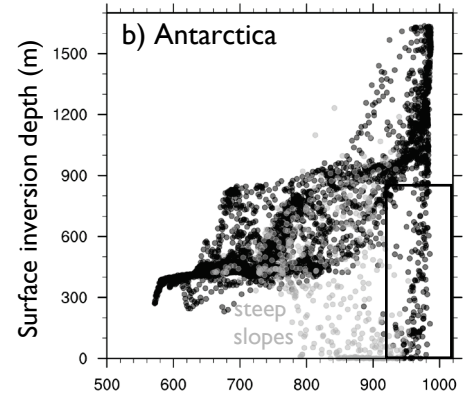
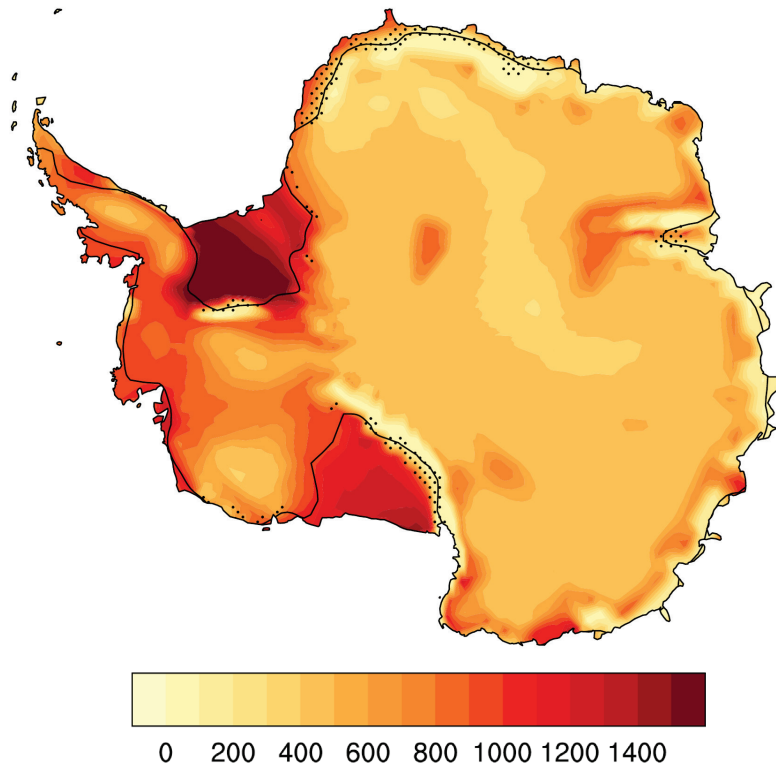


Figure 2. Winter surface-based inversion depth (m) for (a) Antarctica, (b) Antarctica from 70 to 90°S, and (c) the Arctic from 70 to 90°N from the CESM preindustrial control topography experiment. Black contour in (a) shows where the winter surface pressure equals 950 hPa: the level of the maximum AHT due to transient and stationary eddies during winter at 60°S. Grey points in (b) indicate surface slopes greater than .007. Boxed points in (b), with surface pressure greater than 920 hPa and inversion depth less than 850 m, correspond to stippled regions in (a). In (c), blue points indicate winter sea ice fraction less than .97, and grey points are located over Greenland.

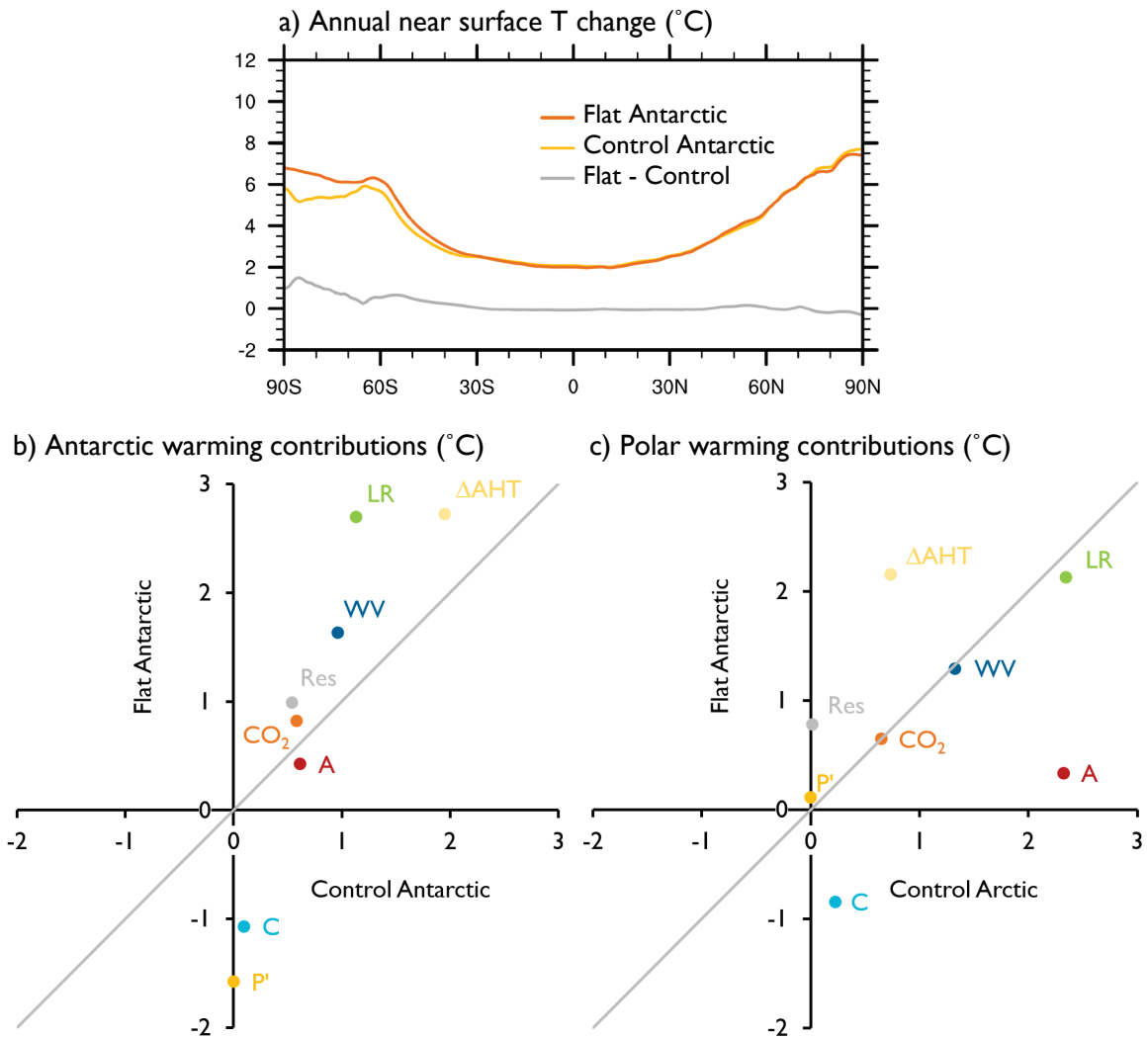


Figure 3. (a) Zonal- and annual-mean near-surface temperature change (°C) under CO₂ doubling in the control (yellow) and flat (orange) Antarctic experiments and their difference (grey); (b,c) Contributions of each feedback and atmospheric forcing to warming (°C) for the flat Antarctic compared to (b) the control Antarctic and (c) the control Arctic for the lapse rate (LR), surface albedo (A), water vapor (WV), and cloud (C) feedbacks, the variation in the Planck response from its value in (b) the control Antarctic and (c) the control Arctic (P'), CO₂ forcing (CO₂), change in AHT convergence (ΔAHT), and residual term (Res).

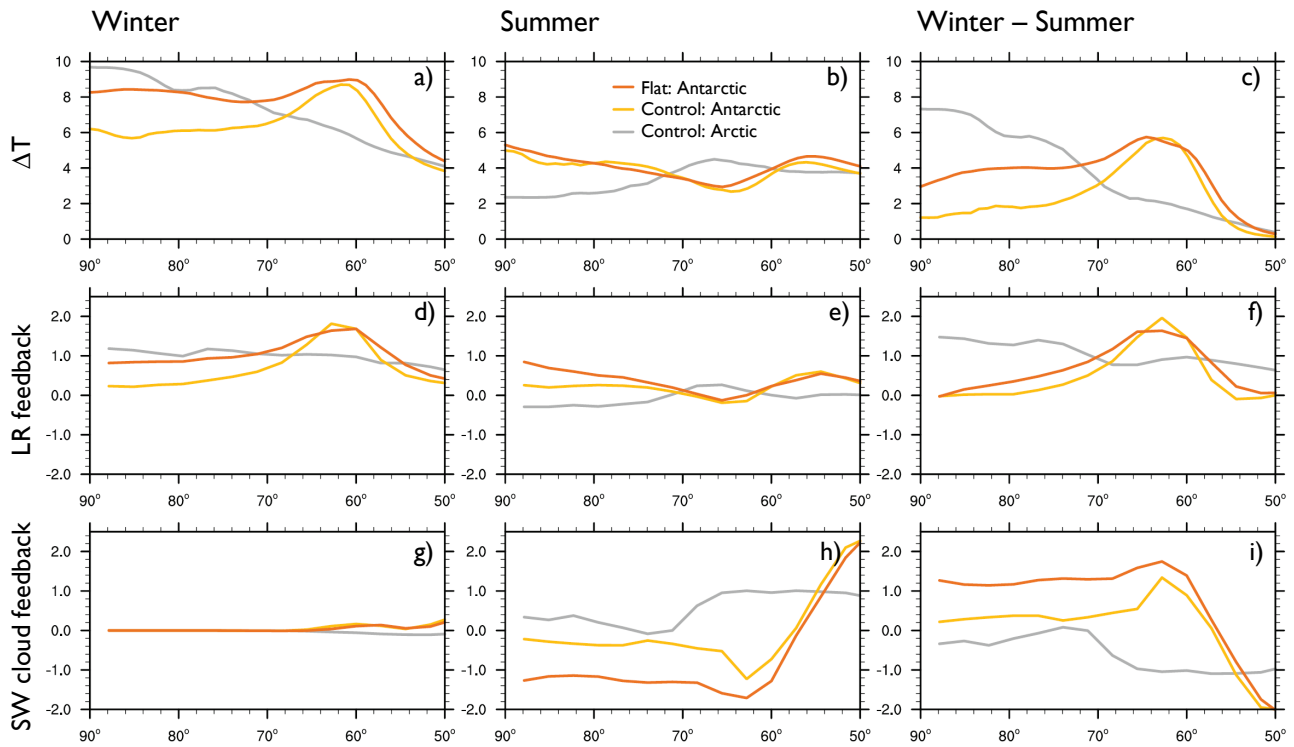


Figure 4. Zonal-mean (a-c) near surface temperature change ($^{\circ}\text{C}$), (d-f) lapse rate feedback ($\text{W}/\text{m}^2/\text{K}$), and (g-i) shortwave (SW) cloud feedback ($\text{W}/\text{m}^2/\text{K}$) under CO_2 doubling in the control Arctic (grey) and control (yellow) and flat (orange) Antarctic for (a,d,g) winter, (b,e,h) summer, and (c,f,i) winter minus summer.