

# Seismic constraints on rock damaging in a failing mountain peak: the Hochvogel, Allgäu

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## Abstract

Large rock slope failures play a pivotal role in long-term landscape evolution and are a major concern in land use planning and hazard aspects. While the failure phase and the time immediately prior to failure are increasingly well studied, the nature of the preparation phase remains enigmatic. This knowledge gap is due, to a large degree, to difficulties associated with instrumenting high mountain terrain and the local nature of classic monitoring methods, which does not allow integral observation of large rock volumes. Here, we analyse data from a small network of up to seven seismic sensors installed during July–October 2018 (with 43 days of data loss) at the summit of the Hochvogel, a 2592 m high Alpine peak. We develop proxy time series indicative of cyclic and progressive changes of the summit. Fundamental frequency analysis, horizontal-to-vertical spectral ratio data and end-member modelling analysis reveal diurnal cycles of increasing and decreasing coupling stiffness of a 126,000 m<sup>3</sup> large, instable rock volume, due to thermal forcing. Relative seismic wave velocity changes also indicate diurnal accumulation and release of stress within the rock mass. At longer time scales, there is a systematic superimposed pattern of stress increases over multiple days and episodic stress release within a few days, expressed in an increased emission of short seismic pulses indicative of rock cracking. We interpret our data to reflect an early stage of stick slip motion of a large rock mass, providing new information on the development of large-scale slope instabilities towards catastrophic failure.

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2                   **mountain peak: the Hochvogel, Allgäu**

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8                   **Key Points:**

- 9                   • Continuous and discrete failure preparation signals of a large-scale slope instabil-  
10                  ity are seismically recorded
- 11                  • Reversible and irreversible mechanisms at the diurnal, multi-day and seasonal scale  
12                  are quantified
- 13                  • We infer an early stage of stick slip motion and thermally forced diurnal stress re-  
14                  lease and rock mass stiffness changes

## Abstract

Large rock slope failures play a pivotal role in long-term landscape evolution and are a major concern in land use planning and hazard aspects. While the failure phase and the time immediately prior to failure are increasingly well studied, the nature of the preparation phase remains enigmatic. This knowledge gap is due, to a large degree, to difficulties associated with instrumenting high mountain terrain and the local nature of classic monitoring methods, which does not allow integral observation of large rock volumes. Here, we analyse data from a small network of up to seven seismic sensors installed during July–October 2018 (with 43 days of data loss) at the summit of the Hochvogel, a 2592 m high Alpine peak. We develop proxy time series indicative of cyclic and progressive changes of the summit. Fundamental frequency analysis, horizontal-to-vertical spectral ratio data and end-member modelling analysis reveal diurnal cycles of increasing and decreasing coupling stiffness of a 126,000 m<sup>3</sup> large, instable rock volume, due to thermal forcing. Relative seismic wave velocity changes also indicate diurnal accumulation and release of stress within the rock mass. At longer time scales, there is a systematic superimposed pattern of stress increases over multiple days and episodic stress release within a few days, expressed in an increased emission of short seismic pulses indicative of rock cracking. We interpret our data to reflect an early stage of stick slip motion of a large rock mass, providing new information on the development of large-scale slope instabilities towards catastrophic failure.

## 1 Introduction

Gravitational mass wasting is the dominant geomorphic process shaping mountain peaks. The release of large rock volumes on steep slopes marks the start of a sediment cascade that evolves to the terrestrial sediment flux into the oceans. Its suddenness, velocity and limited predictability renders gravitational mass wasting hazardous yet hard to constrain in terms of drivers and triggers, and mechanisms and their temporal evolution. Classic approaches to studying mass wasting processes are either *ex post* investigations of the failure mechanism (Frayssines & Hantz, 2006), volume and environmental conditions before and during the event (D’Amato et al., 2016), long-term monitoring, often with remotely sensed imagery, or, increasingly, point measurements from dedicated sensors at sites with known or expected activity (Di Maio et al., 2010; Lévy et al., 2010; Dixon et al., 2018; Collins et al., 2018; Walter et al., 2020). None of these approaches can provide detailed continuous insight into the activity with high temporal resolution, beyond the “point dimension” and throughout the wider process domain, including the downslope pathway of a failing rock mass.

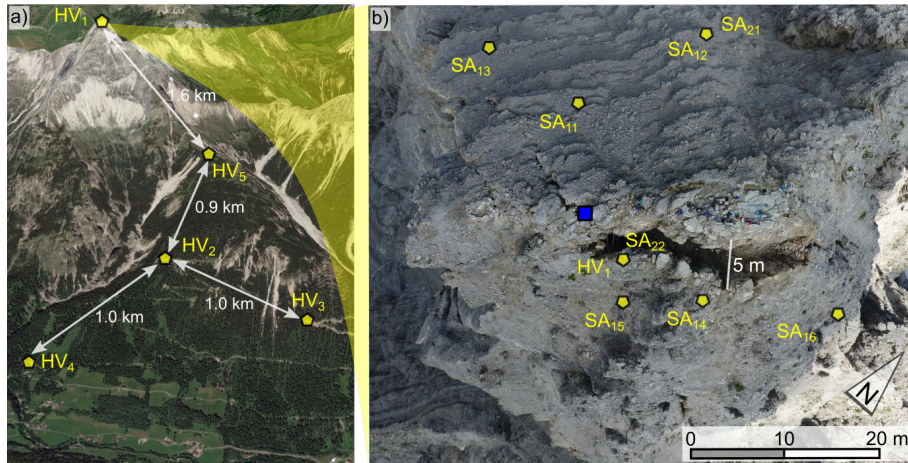
A mass wasting event is preceded by a phase of preparation, during which a set of intrinsic and extrinsic processes drive a rock mass to a state of instability, making it susceptible to trigger mechanisms such as loading with precipitation, seismic ground shaking, wind force fluctuation, and freeze-thaw transitions (see for example Stock et al. (2013) for a summary of trigger mechanisms). While the preparation phase is crucial for hazard anticipation, its mechanisms are little known. It is clear, however, that rock slope failures predominantly develop in the steepest rock flanks of mountains under high topographic stresses (Leith et al., 2014). There, differential stress conditions can be above the micro crack initiation threshold and favour crack coalescence along the weakest trajectories in the rock mass, i.e. the critical path (Einstein et al., 1983). Once a rock mass starts to move, deformation accumulates along that critical path, and shear stress concentrates on a decreasing number of remaining rock bridges, initiating subcritical and critical fracture propagation in these preferential damage zones (Kemeny, 2003). In the final phase of rock slope failure, when most rock bridges have been destroyed, the rock mass often exhibits cyclic accelerations and decelerations referred to as “stick-slip failure” (e.g., Yamada et al., 2016; Schöpa et al., 2018). This behaviour is thought to be

66 controlled by macro-roughness along the sliding plane that inhibits sliding and promotes  
 67 stress concentration along the affected obstacles in the sliding plane and release in episodic  
 68 deformation (Borri-Brunetto et al., 2003; Saettele et al., 2015). In this phase, changes  
 69 in hydrostatic pressure and seismic excitation can initiate significant acceleration of a  
 70 sliding rock mass by enhancing shear stress on remaining obstacles. This is due to tran-  
 71 siently enhanced overall shear stress, reduced total friction of the surrounding sliding plane,  
 72 and the effect of water on rock fracture mechanical strength (Voigtländer et al., 2018).  
 73 As shear planes are mostly inaccessible, insight into spatial and temporal patterns of stress  
 74 concentration and fracture propagation is difficult to achieve with classic approaches, es-  
 75 pecially under natural conditions.

76 In the last decade, networks of seismic sensors have been established as a compli-  
 77 mentary tool to overcome some of the limitations of other measurement and monitor-  
 78 ing approaches. Seismic sensors can record ground motion caused by processes acting  
 79 at or close to the Earths surface. In general, seismic networks can i) detect mass wast-  
 80 ing activity, ii) locate and track the process in space, and iii) infer kinetic and anatomic  
 81 details of a process event. The application fields of environmental seismology (Burtin  
 82 et al., 2014; Larose et al., 2015) comprise, for example, time-resolved investigation of the  
 83 evolution of slope instabilities (Mainsant et al., 2012; Schöpa et al., 2018), detection and  
 84 quantification of event activity from the catchment to the global scale (Lacroix & Helm-  
 85 stetter, 2011; Dammeier et al., 2011; Ekström & Stark, 2013; Fuchs et al., 2018), track-  
 86 ing of mass movements in space (Burtin et al., 2016; Cook et al., 2018; Pérez-Guillén et  
 87 al., 2019; Walter et al., 2020), inversion of seismic signals for event kinetics (Allstadt,  
 88 2013; Ekström & Stark, 2013), and attribution of events to drivers and triggers (Helmstetter  
 89 & Garambois, 2010; Burtin et al., 2013; M. Dietze, Turowski, et al., 2017).

90 In addition to investigating the seismic signals emitted by mass wasting processes,  
 91 it is possible to use random background signals (ambient noise) to survey the material  
 92 properties of landforms (I. Welch & McLamore, 1973). There are several techniques that  
 93 are sensitive to state changes of the near-surface materials in eroding landscapes. A clas-  
 94 sic technique is the horizontal to vertical spectral ratio (HVSR) analysis (Nogoshi & Igarashi,  
 95 1971; Nakamura, 1989). It has been widely used to estimate the near-surface structure  
 96 of sites (for a review see Del Gaudio et al., 2014), though there is controversy about the  
 97 confidence and robustness of quantitative inversions for material properties (see for ex-  
 98 ample discussion in Overduin et al., 2015). Bottelin et al. (2013) used seismic sensors  
 99 to monitor changes in the fundamental frequencies of an unstable rock column in the French  
 100 Alps as a function of temperature and precipitation. Weber, Fäh, et al. (2018) applied  
 101 the same technique to the Matterhorn, identifying cyclic as well as cumulative damage  
 102 evolution. Seismic noise cross correlation (Snieder, 2004; Sens-Schönfelder & Wegler, 2006)  
 103 has the potential to detect subtle, relative changes in the seismic wave velocity, termed  
 104  $dv/v$  throughout this text. The technique can be applied using the individual compo-  
 105 nents of a single seismic station or records from seismic station pairs, given that both  
 106 stations are within the coherence range of the ambient wave field. The seismic wave ve-  
 107 locity is affected by a series of factors, such as rock temperature, ground moisture, ac-  
 108 cumulated rock damage and mechanical stress (Snieder, 2004; Larose et al., 2015; Clements  
 109 & Denolle, 2018; M. Dietze et al., submitted). HVSR, fundamental frequency and  $dv/v$   
 110 are valuable proxies of material state, because they provide information on the average  
 111 state of the medium within the seismic network aperture, not just at a point. Moreover,  
 112 these proxies can be recorded continuously, at arbitrary high resolution and with little  
 113 infrastructural effort, so that the evolution of the material state can be tracked over time  
 114 and changing boundary conditions.

115 Here we study the summit of the Hochvogel, a 2592 m high solitary peak in the South  
 116 German Alps. We exploit the utility of a small, telemetered seismic network on the sum-  
 117 mit, which was operational during a period in which a large rock volume near the sum-  
 118 mit was undergoing failure preparation. We investigate seismically evident patterns of



**Figure 1.** Study area and instrumentation. a) Oblique view from the Austrian side towards the north. The 2592 m high Hochvogel with station HV<sub>1</sub> on the summit is in the background, the town of Hinterhornbach in the foreground. HV<sub>1</sub> to HV<sub>5</sub> are the seismic stations forming the hillslope network to monitor mass wasting activity. b) Close-up top view of the Hochvogel summit with a > 25 m long and 2–4 m wide main fracture in the centre. Yellow polygons depict six seismic sensors of the small aperture summit network, station HV<sub>1</sub> is the same as in a). Blue square depicts data logger and GSM antenna. Figures S1 1–3 provide further details of the summit and the instrumentation.

119 summit deformation, how these are controlled by environmental conditions, and how they  
 120 translate into mechanistic activity.

## 121 2 Study site and instrumentation

122 The Hochvogel is a solitary, 2592 m high peak in the Eastern Allgäu Alps, on the  
 123 border between Germany and Austria. It is formed of folded and thrust upper Tri-  
 124 assic dolomite (Hauptdolomit), which forms competent decimetre to metre thick beds,  
 125 20 ° dipping to the NNW in the summit area. Extensive vertical fracture sets run per-  
 126 pendicular to the general dip, giving rise to an intensely fractured fabric. Formation of  
 127 large fractures in the summit area was first recorded in the late nineteenth century, and  
 128 a progressively accelerating rock slide has developed since the 1950s. Catastrophic fail-  
 129 ure could mobilise up to 260,000 m<sup>3</sup> of rock in several subunits (Krautblatter et al., 2019)  
 130 along a failure plane located on two large, NE–SW trending fracture systems. In 2018,  
 131 the southeastern, main fracture had already opened by 2–4 m, and it continues open-  
 132 ing at a rate of a few mm per month (Krautblatter et al., 2019). The unstable rock mass  
 133 has moved downslope by several metres and developed a severe state of fracturing (Fig.  
 134 SI 1–3).

135 To monitor activity at the Hochvogel summit as well as any mass wasting processes  
 136 in the southeastern mountain slope, we installed a nested network of seismic stations (Fig. 1).  
 137 At the summit we installed a small aperture network (33 m average, 73 m maximum spac-  
 138 ing) consisting of six PE6B one-component 4.5 Hz geophones, three on the less fractured,  
 139 relatively stable German side (SA<sub>11</sub>–SA<sub>13</sub>) and three on the failing Austrian side (SA<sub>14</sub>–  
 140 SA<sub>16</sub>). The sensors were connected by 100 m long cables to a Nanometrics Centaur data  
 141 logger that recorded the ground velocity values at 200 Hz with a gain of 40. These seis-  
 142 mic data were transmitted to the GFZ data server every five minutes. The network was  
 143 operational from 10 July 2018. However, after a lightning strike of the station on 1 Au-

144 gust 2018 no data were recorded for 43 days until 12 September 2018. From this date  
 145 we operated a set of two PE6B 4.5 Hz geophones (SA<sub>21</sub> a 1-component and SA<sub>22</sub> a 3-  
 146 component sensor) and Digos DataCube<sup>3</sup>ext loggers, installed in an aluminium case with  
 147 a 2 cm thick concrete bed. To record slope activity below the summit, we installed five  
 148 seismic stations with spacings between 0.9 and 1.6 km. Station HV<sub>1</sub> contains the same  
 149 setup as SA<sub>22</sub> and was used in this study as additional summit station. Here we only  
 150 consider data from the small aperture array at the summit, for the periods 10 July–1 Au-  
 151 gust and 12 September–12 October.

### 152 3 Data processing

153 All seismic data were processed with the software R v. 3.6.3 (R Development Core  
 154 Team, 2020), using the package *eseis* v. 0.5.0 (M. Dietze, 2018a, 2018b). The Support-  
 155 ing Information (SI) contains dedicated R scripts of all major processing and analysis  
 156 steps. The underlying data is available via GFZ Data Services (M. Dietze & Krautblat-  
 157 ter, 2020). For fundamental frequency analysis we calculated averaged (P. Welch, 1967)  
 158 spectra for non-event periods (Bottelin et al., 2013) of the summit network sensors. For  
 159 this we calculated an STA-LTA ratio (Allen, 1982) and flagged periods with a ratio  $>$   
 160 1.5 as event-contaminated. These periods were excluded from the subsequent step of spec-  
 161 tra calculation. The spectra were calculated for 5 s windows without overlap, which were  
 162 then stacked to non-overlapping hourly spectra. Each of these spectra was subsequently  
 163 normalised. Finally, we executed a 5 sample running average filter in frequency and time.  
 164 Since the PE6B geophone sensors have a natural frequency of 4.5 Hz, we focus our anal-  
 165 ysis on frequencies above this limit, knowing that we most likely miss the lowest funda-  
 166 mental frequency band of the rock mass under investigation (e.g. Lévy et al., 2010; Bot-  
 167 telin et al., 2013).

168 For the HVSR analysis we calculated spectra of station HV<sub>1</sub> and SA<sub>22</sub> for the ver-  
 169 tical and horizontal components within 5 s windows with 50 % overlap. We then calcu-  
 170 lated the ratio of the corresponding horizontal and vertical spectra and stacked these to  
 171 hourly results. To quantify the spectral modes inherent to the HVSR data, we applied  
 172 end-member modelling analysis (EMMA) to the data set. EMMA (E. Dietze & Dietze,  
 173 2019) is a statistical technique, which allows unmixing multimodal data, and describ-  
 174 ing it as a linear combination of loadings (the fundamental underlying spectra distribu-  
 175 tions) and scores (the relative contribution of each loading to each sample). EMMA was  
 176 originally developed for decomposition of high resolution grain-size spectra, but it can  
 177 be extended to other kinds of compositional data (i.e., data that is non-negative and where  
 178 each sample is described by a constant sum). To be able to apply EMMA, we truncated  
 179 the HVSR data set to the most variable frequency range and normalised the spectral ra-  
 180 tios of each hourly time slice between 0 and 1. Since we were interested in identifying  
 181 the position and shape of the spectral modes as well as their relative contributions to  
 182 a temporal sample, but not in the absolute intensity of the HVSR amplitudes, this nor-  
 183 malisation step is valid. We followed the generic deterministic EMMA protocol (using  
 184 the *EMMAgeo* package v. 0.9.8, E. Dietze & Dietze, 2019) and tested between two and  
 185 six end-members for plausibility of the results. These plausibility checks were based on  
 186 model quality (i.e. time-resolved and frequency-resolved variance explained by the model),  
 187 the degree of unmixing of the frequency distributions and the interpretative meaning-  
 188 fulness of the unmixed frequency distributions and their evolution with time.

189 We estimated relative velocity changes from ambient noise cross correlation. First,  
 190 the seismic traces were pre-processed to reduce the impact of erratic, high amplitude sig-  
 191 nals. We trimmed the seismic records to 10 min long segments, downsampled to 50 Hz,  
 192 detrended and filtered in the 4–16 Hz frequency range. We normalised amplitudes by  
 193 spectral whitening, setting all the amplitudes in the Fourier spectrum to 1 and performed  
 194 1-bit normalisation. In the first time period, we performed cross correlations between  
 195 all vertical component sensor pairs. In the second period, since only two stations were



operating, located on different sides of the main fracture, we performed single station cross correlation (Hobiger et al., 2014) on station SA<sub>22</sub> using the three combinations North-Vertical, East-Vertical, and East-North. We stacked the obtained correlation functions to 30 min non-overlapping averages and stored them in a correlation matrix  $C(t_i, \tau)$  where  $t_i$  and  $\tau$  is the lapse time in the correlation process. To infer relative velocity changes  $dv/d$ , we used the stretching technique (Sens-Schönfelder & Wegler, 2006): Each of the 30 min correlation functions were be stretched or compressed depending on their relative velocity change compared to a reference ( $\tau$ ). For the two time periods, we computed the average correlation function ( $\tau$ ) and used it as a reference. We calculated the correlation coefficient  $R(t_i, \varepsilon_j)$  between stretched versions of this reference with each correlation functions stored in  $C(t_i, \tau)$  such as

$$R(t_i, \varepsilon_j) = \int_{\tau_1}^{\tau_2} C(t_i, \tau) \xi(\tau * (1 + \varepsilon_j)) d\tau$$

where  $\varepsilon_j$  indicates stretch-values tested in the time window  $[\tau_1 \tau_2]$ . We used  $\tau_1 = 0.5$  s to avoid unwanted source effects and used a window duration of 3 s. For the second time period, we averaged the measurements between the three combinations by stacking their three corresponding matrices  $R(t_i, \varepsilon_j)$ . Finally, we extract those  $dv(t_i)/v$  values for which  $\varepsilon_j$  yielded the maximum correlation coefficient  $R(t_i, \varepsilon_j)$ , and only kept values where the correlation coefficient was at least 0.5.

In addition to the continuous analysis we isolated discrete seismic events in the data of the summit network. We picked events using a classic STA-LTA ratio routine (Allen, 1982), applied to the 10–40 Hz filtered signal envelopes. The routine was applied with a low detection threshold (on-ratio = 3, off-ratio = 1, STA-window = 0.5 s, LTA-window = 120 s). This yielded many false positives but also guaranteed detection of small events. All potential events were tested against automatic rejection criteria. We required an event to have a minimum duration of 0.2 s as well as a maximum duration of 10 s. We assumed that events shorter than 0.2 s are associated with random sources such as rain drop impacts (M. Dietze, Mohadjer, et al., 2017), and events longer than 10 s are unlikely to be linked to discrete geomorphic activity at the summit. A further rejection criterion imposed that an event had to be detected by at least two stations with a time difference of not more than 0.1 s. This rule enforces that signals must travel across the entire network in a time that corresponds to an apparent seismic wave velocity of at least 730 m/s, a threshold well below typical values for limestone of different origin and degree of fracturing (Assefa et al., 2003; Helmstetter & Garambois, 2010; M. Dietze, Mohadjer, et al., 2017). Signals with longer detection time differences are usually a result of the coincidence of unrelated signals or waves travelling through air (about 330 m/s velocity) before coupling to the ground. All remaining events were then checked manually for plausibility: telltales of genuine geomorphic signals include spindle shaped seismograms, and consistent amplitude and frequency decrease as well as onset shifts with distance between source and sensor (for details see Hibert et al., 2011; Hammer et al., 2015; M. Dietze, Mohadjer, et al., 2017; Vouillamoz et al., 2018).

Meteorological data for the summit were approximated by scaling hourly values of a station near Oberstdorf, 12 km to the west (DWD, 2019). Temperature scaling by -0.6 °C per 100 m altitude increase results in a downward shift of 12 °C of summit temperatures. For precipitation the German Weather Bureau offers hourly gridded data (DWD, 2019). However, for the time of interest there are data gaps. Therefore, we calculated a transfer function between gridded values at the peak and the Oberstdorf data for adjacent periods with data and use this relationship to infer precipitation data also for the peak during relevant data gaps.

## 4 Results

### 4.1 Meteorological conditions

Meteorological dynamics define an important set of boundary conditions for the activity of the summit. The scaled air temperature records (Fig. 2) are dominated by the diurnal cycle, with amplitudes of about 18 ° C in the first instrumentation period (10 July–01 August) and about 22 ° C in the second period (10 September–12 October). Superimposed on these fluctuations was a rising trend from 10 to 23 ° C (interrupted by stable temperatures between 16–24 July) in the first period and a somewhat decreasing trend from 14–10 ° C during the second period. Precipitation events occurred throughout the record time (Fig. 2 e), but lasted only a few hours to a day, with usually less than 6 mm/h peak intensities. Precipitation events did not coincide with changes in any of the seismic data sets.

### 4.2 Continuous seismic observations at the summit

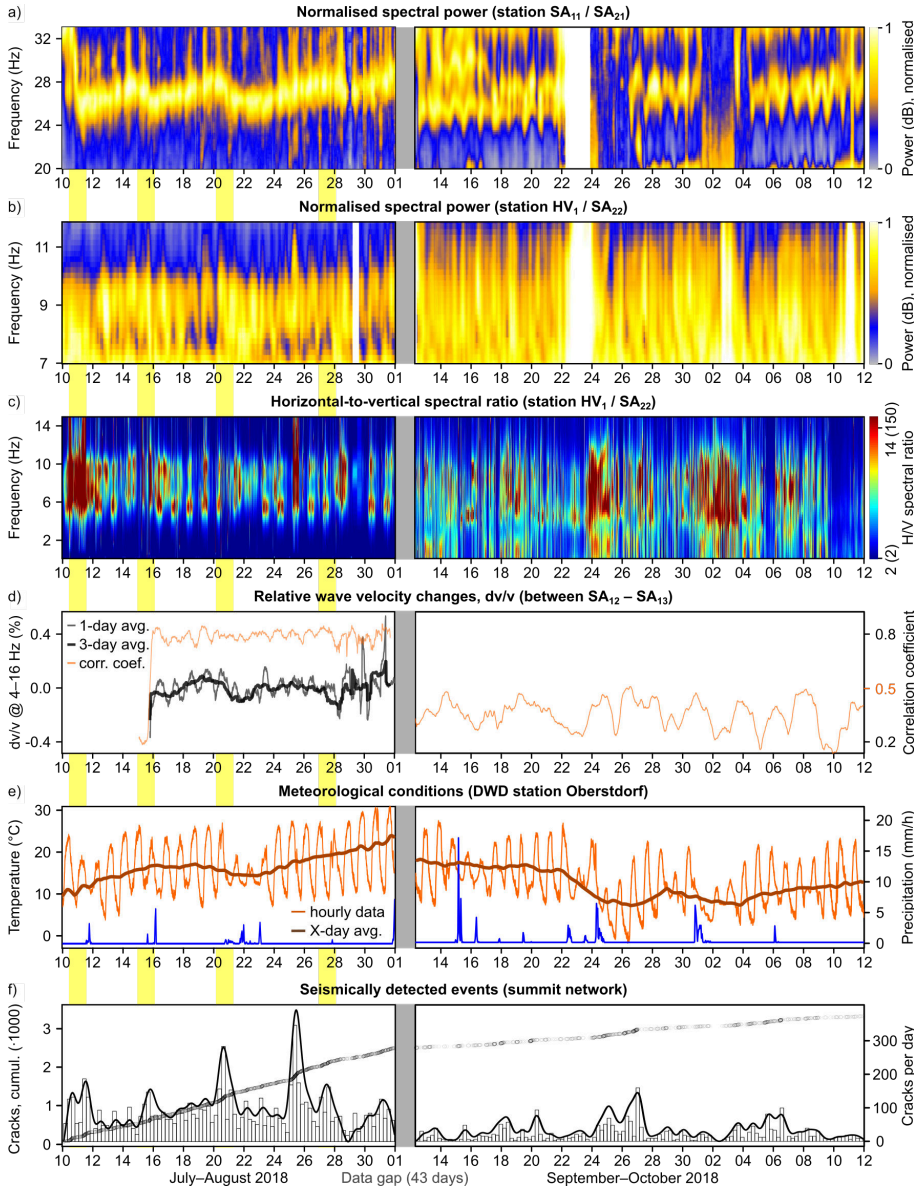
Continuous seismic observations include the results of the fundamental frequency, HVSR and  $dv/v$  analyses, which will be used to infer material property changes of the rock mass. Here, we offer a condensed, synoptic report of the multitude of observed metrics. The SI contain results from all sensors. In general, sensors from the instable mass on the southeast side of the Hochvogel summit (SA<sub>14</sub>, SA<sub>15</sub>, SA<sub>16</sub>, SA<sub>22</sub>) did not record a systematic pattern of fundamental frequency evolution on the vertical component. In contrast, stations on the stable mass, opposite (SA<sub>11</sub>, SA<sub>12</sub>), did reveal consistent vertical component fundamental frequency bands. These were best expressed at station SA<sub>11</sub>, located closest to the summit and the main fracture (see Fig. SI 5 and 2 a). We used the three component stations HV<sub>1</sub> (first period) and SA<sub>22</sub> (second period) for HVSR analysis. Both stations were located on the instable mass adjacent to the main crack. The  $dv/v$  analysis yielded between-station correlation coefficients > 0.5 for the station pair SA<sub>12</sub> and SA<sub>13</sub> while all other pairs did not show meaningful signal coherence. Likewise, the between-components analysis for SA<sub>22</sub> during the second period did not yield correlation coefficients above 0.5 and results were thus discarded.

The two monitoring periods (10 July to 01 August and 12 September to 12 October, separated by 43 days without data) showed contrasting results of the analysed monitoring metrics. The fundamental frequency of the vertical component of stations installed near the top of the stable northern rock mass (SA<sub>21</sub> and SA<sub>22</sub>, Fig. 2 a) ranged between 24 and 28 Hz. There were also higher frequencies present in the data set (see Fig. SI 5) but they are either multiples of the lower frequency band from Fig. 2 a or frequencies corresponding to higher order bending modes (Bottelin et al., 2013). In the first period, the 24–28 Hz band exhibited prominent multi-day cycles, consisting of four to seven days of continuous rise of the fundamental frequency by two to three Hz followed by a sudden drop, almost back to the initial values. Four such cycles are visible in Fig. 2 a, with drops around 10, 15, 21, 29 July. In the second period, the fundamental frequency was dominated by diurnal cycles without clear multi-day effects. On the instable southern rock mass, no fundamental frequency pattern was found in the vertical component (see Fig. SI 5).

The horizontal seismic sensor component (Fig. 2 b) exhibited a fundamental frequency band between 7 and 10 Hz, though weakly expressed, with diurnal cycles throughout both periods. Multi-day patterns were suppressed (first period) or absent (second period). A data gap during the second period was due to event contaminated time slices, which were removed, thus affecting the moving window temporal averaging process.

Diurnal cycles were also visible in the HVSR results (Fig. 2 c), most clearly during the first survey period (10 July to 01 August). The higher resolution of this record in the time dimension allowed a more detailed analysis of the data. This revealed that





**Figure 2.** Seismic properties of the Hochvogel summit network in the summer (left panels) and autumn (right panels) survey period. a) Evolution of the fundamental frequency of vertical component station  $SA_{11}$  ( $SA_{21}$  after reinstallation on 12 September 2018). b) Evolution of the fundamental frequency of the horizontal components of station  $HV_1$  ( $SA_{22}$  after reinstallation on 12 September 2018). c) HVSR evolution for the same stations as in b). Values in the second period are given in parentheses of the legend labels. d) Relative wave velocity change ( $dv/v$ ) between the stations  $SA_{12}$  and  $SA_{13}$  (left panel). Thin grey line depicts 1-day running average smoothed data, Thick black line shows 3-day running average results. Orange line depicts correlation coefficient of the signal pairs. Only  $dv/v$  data above  $r = 0.5$  is shown. e) Meteorological conditions at a station 12 km west of the summit. f) Seismically detected events, shown as cumulative event plot (open circles) and 6-hour histograms of daily crack rates. Kernel density estimate (solid bold line, 6 hour kernel bandwidth) is plotted on top of histogram for better visual interpretation of crack rate evolution.

the HVSR was around 5–7 Hz during night time and around 9–10 Hz during daytime. This diurnal pattern was distorted during episodes when the vertical fundamental frequency drops occurred (Fig. 2 a).

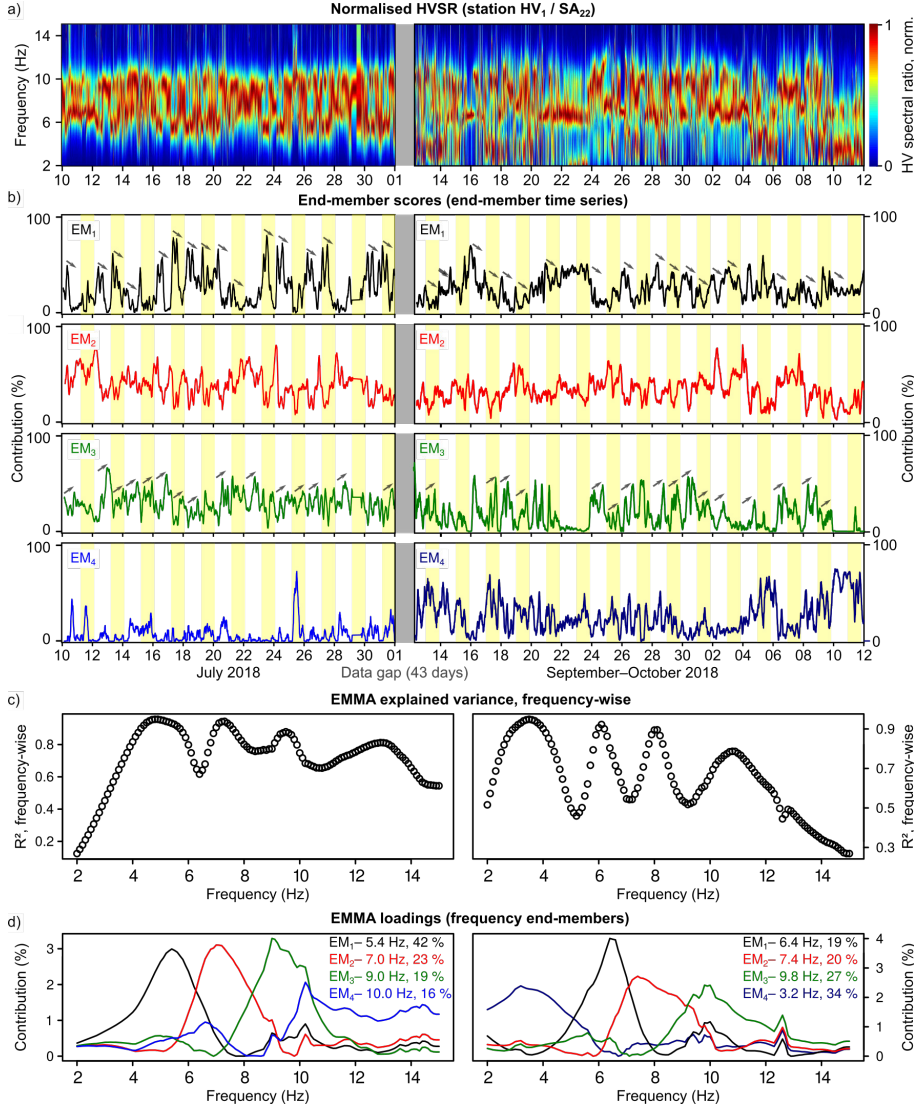
### 4.3 End-member modelling

To better characterise and quantify the underlying frequency patterns and their contribution to the data through time, we performed end-member modelling analysis on the 2–15 Hz band of the normalised HVSR data (Fig. 3 a). Models with two (62.2 % of total variance explained for first period, 55.9 % for second period), three (75.6 %, 70.0 %), five (77.9 %, 70.8 %) and six (80.0 %, 75.3 %) end-members showed mostly lower performance than the model with four end-members (80.6 %, 71.9 %). More importantly, models with other than four end-members did not yield meaningful unmixing results, usually exhibiting multimodal or broad frequency distributions or largely overlapping end-members (see SI for all tested models). The two four-end-member models for the survey periods 10 July–01 August and 12 September–12 October allowed describing the data set in the time domain (i.e., the hourly time windows of HVSR data) with  $R^2 = 0.89$  (first period) and  $R^2 = 0.79$  (second period). The frequency domain (i.e., the model representation of individual HVSR distributions, cf. Fig. 3 c) is described with  $R^2 = 0.73$  (first period) and  $R^2 = 0.65$  (second period, mostly because frequencies above 10 Hz did not contribute much to the data set overall). The shapes of the end-member loadings (Fig. 3 d) agree with the visual inspection of the raw HVSR data (Fig. 2 c). During the first survey period, two recurring frequency modes at 5.4 and 9.0 Hz are represented by EM<sub>1</sub> and EM<sub>3</sub>, while the transitional stages between these two are formed by EM<sub>2</sub> (modal frequency at 7.0 Hz). EM<sub>4</sub> covers a broader frequency band, ranging from 10–15 Hz. During the second survey period, EM<sub>1</sub> to EM<sub>3</sub> had consistently shifted their modal frequencies, to 6.4 Hz, 7.4 Hz and 9.8 Hz, respectively. The EM<sub>4</sub> from the first period was no longer visible, but EM<sub>3</sub> showed a broad shoulder towards higher frequencies. During this second period, EM<sub>4</sub> was in the frequency range 2–5 Hz with a modal frequency of 3.2 Hz. This new low frequency variability in the data is also visible in the normalised (Fig. 3 a) as well as the raw (Fig. 2 c) HVSR data set. In terms of end-member scores (Fig. 3 b), EM<sub>1</sub> provided high contributions to the HVSR data mostly during the early hours of the day, and EM<sub>3</sub> during the late hours. EM<sub>2</sub> takes an intermediate role, between EM<sub>1</sub> and EM<sub>3</sub>. In the first survey period, EM<sub>4</sub> mainly contributed to HVSR data during episodes of enhanced seismic event rates (see below).

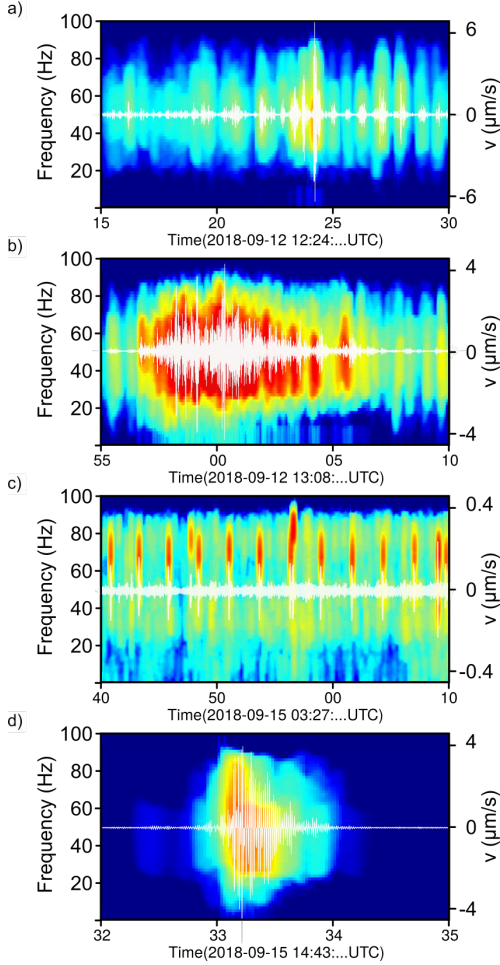
Noise cross correlation analysis yielded meaningful results only for the station pair SA<sub>12</sub>–SA<sub>13</sub> (Fig. 2 d). Even there, the correlation among the stations did not reach  $> 0.8$  before 16 July. From that date on, the 1-day running averaged  $dv/v$  data shows a diurnal pattern superimposed on a longer-term trend. That longer term trend is better visible in the 3-day running average curve: Coincident with the periods of vertical fundamental frequency (Fig. 2 a) and distortions of the bimodal HVSR pattern (Fig. 2 c), the  $dv/v$  trend showed abrupt decreases by about 0.3 % around 15, 21 and 28 July 2018.

### 4.4 Discrete seismic events

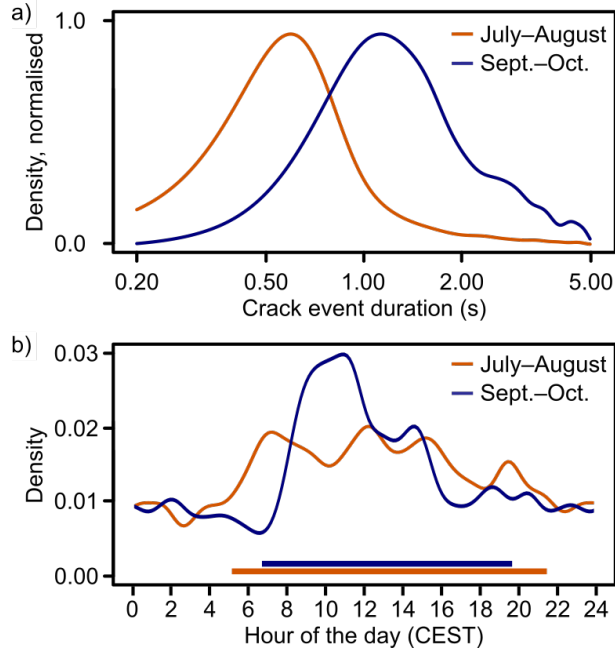
The summit network recorded a series of discrete signals of diverse properties with respect to duration, evolution and spectral content. The STA/LTA event detection algorithm and application of the subsequent rejection criteria yielded a total of 2490 and 845 potential near-surface events for the first and second survey period, respectively. One class of detected events had the hallmarks of rock cracking (e.g., Senfaute et al., 2009): one to two seconds long, erratic, single pulses suddenly emerging from an otherwise silent record, with frequencies predominantly between 30 and 80 Hz. These events showed a distinct evolution during the two survey periods. Between 10 July and 01 August, the average event rate was 100 cracks per day, with episodes of enhanced activity on 10–12, 15–16, 21–22, 24–26 and 27–28 July (Fig. 2 f). In the second period, 12 September to



**Figure 3.** End-member modelling analysis results. a) Normalised HVSR data for both monitoring periods. b) End-member scores (contribution of each end-member to a time slice of a) of all four end-members. Alternating yellow and white bars depict single days. Semi transparent small arrows illustrate days of typical decrease and increase of end-member contribution to the signals with the course of the day. c) Model  $R^2$  in the frequency space. d) End-member loadings (frequency shapes of the four end-members). Legend denotes modal frequency and in parentheses the explained variance of each end-member.



**Figure 4.** Seismic signals due to different sources, as recorded by the Hochvogel summit network (station SA<sub>22</sub>). The plots show spectrograms (scaled between -160 and -100  $10\log_{10}(m^2/s^2)/Hz$ , except for c), which is scaled in the range -180 to -120) of the signals and also seismograms as white line overlays. a) Sequence of steps generated by a person walking towards and past the station along the summit. b) Signature of a small mass wasting event at the summit. c) Seismic signals during a rain event (note order of magnitude lower amplitudes). d) Seismic signature of a stress release event as described in Fig. 5.



**Figure 5.** Properties of discrete stress release signals on the Hochvogel summit. a) Duration of signals. The kernel density estimate (0.1 s kernel bandwidth) depicts the log-normal distribution of signal duration, peaking at 0.58 s during July–August and 1.57 s during September–October. b) Time of occurrence at the diurnal scale. Kernel density estimate (1 h kernel bandwidth) shows temporal clustering of events at different times during the two survey periods. Horizontal bars depict topography corrected sun rise and sunset times.

345 12 October, the average crack rate was lower, 18 per day. Episodes with enhanced crack  
 346 activity occurred also during this period (e.g., 18, 20, 24–28, 03–07 October), but sig-  
 347 nificantly less pronounced compared to the earlier interval. We note that the episodes  
 348 of enhanced activity did not systematically coincide with precipitation events or week-  
 349 end days when mountaineering traffic was likely more frequent. Individual events had  
 350 an average duration of  $0.58^{+0.44}_{-0.02}$  s (median and quartile range) during the first and  $1.57^{+0.95}_{-0.48}$   
 351 s during the second period, both following a log-normal distribution (Fig. 5 a). The sig-  
 352 nals were clustered at the diurnal scale, occurring preferentially at day time (Fig. 5 b).  
 353 That trend is even more pronounced during the second period.

354 Among the cases rejected because of their duration were events as shown in Fig. 4.  
 355 These signals exhibited short, regularly spaced pulses (about 1.3 pulses per second) of  
 356 systematically increasing and and decreasing amplitude (Fig. 4 a), or signals without dis-  
 357 tinct pulses but rather a spindle shaped energy envelope, increasing in amplitude for 4  
 358 s and decreasing for 7 seconds (Fig. 4 b). Likewise, during times of precipitation (Fig. 4 c),  
 359 the sensors recorded up to 0.2 s long single pulses with mainly 60–80 Hz frequency con-  
 360 tent and decreasing recurrence intervals but higher amplitudes as the rain intensity in-  
 361 creased (not shown).

## 362 5 Discussion

### 363 5.1 Individual signals of summit activity

364 According to our seismic data, cyclic activity in the Hochvogel summit occurs on  
 365 diurnal and multi-day scales (Fig. 2). The diurnal cyclicity is expressed in systematic



366 1–2 Hz changes of the vertical and horizontal component’s fundamental frequency, in the  
 367 HVSR bimodal frequency distribution, and in  $dv/v$  fluctuations of about 0.2 %. These  
 368 patterns are clearest in the more intact, northwest facing rock mass, and less so in the  
 369 heavily fragmented rock mass facing towards the southeast. Multi-day cyclicality was promi-  
 370 nent during the first survey period (10 July–01 August), expressed in several seismic ob-  
 371 servables. The vertical component fundamental frequency increased gradually over four  
 372 to six day periods, alternating with two days of decreasing trends. Similarly, a pattern  
 373 of 4–6 days stability with 2 day long interruptions of the bimodal frequency distribution  
 374 pattern was found in the HVSR data. During these interruptions,  $dv/v$  dropped by about  
 375 0.2 %, and the rate of short duration seismic events in the summit area increased by a  
 376 factor of 0.5 to 7. The significance of these coincidence is hard to estimate. However, ten-  
 377 tatively we infer that these temporally connected signals are being driven by the same  
 378 underlying mechanisms. To probe the nature of these drivers, we consider the individ-  
 379 ual seismic proxies (Fig. 2) in more detail.

### 380 *5.1.1 Fundamental frequency*

381 The fundamental frequency with which an object vibrates is at first order defined  
 382 by its geometry and material properties. While the geometry of the Hochvogel summit  
 383 can be considered relatively constant, at least during the time period covered in this study,  
 384 the material properties are subject to changes. These can include thermal changes or vari-  
 385 ations in water content due to the local weather, and also internal mechanical changes  
 386 due to rock mass fracturing, or the failure of rock bridges (Lévy et al., 2010; Moore et  
 387 al., 2010). Similar to the range of values found in the Hochvogel summit, Lévy et al. (2010)  
 388 reported a resonance frequency drop of almost 1 Hz, paired with an increased rate of dis-  
 389 crete seismic emissions in a limestone cliff in the Western Alps, just prior to a 21,000 m<sup>3</sup>  
 390 large failure event. Bottelin et al. (2013) measured and modelled the resonance frequency  
 391 of a 1,000 m<sup>3</sup> limestone slab in the Western Alps, finding systematic, thermally driven  
 392 changes in the fundamental frequency. At the diurnal scale, increasing temperatures can  
 393 cause an increasing fundamental frequency, with a time lag of a few hours due to heat  
 394 dissipation into the rock. Bottelin et al. (2013) attribute this to stiffening of the inter-  
 395 face between the column and the rock mass due to thermal expansion, likely involving  
 396 temporary closure of rock fractures. In contrast, at the annual scale, rising temperatures  
 397 result in a lower fundamental frequency (with a time lag of three months) due to a de-  
 398 creasing Young’s modulus. Following this explanation of empirical and numerical mod-  
 399 elling results, we interpret the diurnal cyclicality of our fundamental frequency data as rep-  
 400 resenting the daily opening and closing of fractures in the rock mass, resulting in increased  
 401 stiffness during warming and decreased stiffness during cooling.

402 In contrast, we did not observe a direct or lagged relation between fundamental fre-  
 403 quency and air temperature on the multi-day scale. This may be due, in part, to the fact  
 404 that the long-term temperature range is smaller than the diurnal range. Note also that  
 405 our periods of observation were significantly shorter than the time lag of three months  
 406 reported by Bottelin et al. (2013). Lévy et al. (2010) interpreted fundamental frequency  
 407 drops as a global parameter characterizing the mass and its coupling to the stable rock  
 408 mass. In the absence of any other likely cause we propose that slow increases in frequency  
 409 were caused by cumulative increases of stress within the rock mass, which were then re-  
 410 leased episodically during a few days when frequency drops back to the values before the  
 411 phase of stress increase. Similar drops of the fundamental frequency have also been re-  
 412 ported to correlate with increased rock mass mobility values (e.g. Burjáněk et al., 2018)  
 413 and changes in the elastic moduli have been reported under stress in laboratory exper-  
 414 iments (Barton, 2007; Draebing & Krautblatter, 2012).

415 The different fundamental frequency windows of the vertical component (24–28 Hz,  
 416 Fig. 2 a) and the horizontal component (8–10 Hz, Fig. 2 b) may represent different modes  
 417 of movement, possibly stamping in the vertical direction and either bending or rotating



418 in the horizontal component. These mode dependent frequency differences were also de-  
 419 scribed and modelled by Bottelin et al. (2013). In view of the complex structural situ-  
 420 ation and under-constrained geometric boundary conditions, numerical modelling of the  
 421 rock mass properties appears not meaningful in the Hochvogel case and we desist from  
 422 deduction of more detailed insights about the quantitative physical meaning of the fun-  
 423 damental frequency values.

### 424 **5.1.2 Horizontal-to-vertical spectral ratio**

425 HVSR data are thought to represent site resonance, solicited predominantly by S-  
 426 waves and to a lesser degree by Rayleigh waves (Del Gaudio et al., 2014). Peaks in the  
 427 ratio reflect the frequency at which the power of the horizontal component of a seismic  
 428 signal significantly exceeds that of the vertical component, thus indicating a site specific  
 429 resonance frequency. The method has been used in steep rock slope landscapes to moni-  
 430 tor changes in system and material states, signaled by changes in the resonance frequency  
 431 (e.g. Burjánek et al., 2012; Weber, Fäh, et al., 2018). These studies predominantly fo-  
 432 cused on seasonal to multi-year time spans to explore reversible and irreversible effects  
 433 contributing to failure preparation. Like Weber, Fäh, et al. (2018) we have found mul-  
 434 tiple discrete fundamental frequency bands in the HVSR data (see SI). Here, we focus  
 435 on the lowermost one, around 8–10 Hz (Fig. 2 c), which was also flagged in the funda-  
 436 mental frequency analysis (Fig. 2 b), indicating the spectral amplification effect due to  
 437 the contribution of the horizontal component.

438 Our measurement intervals are not long enough to compare our data against re-  
 439 ported findings on multi-seasonal time scales. However, we can explore the diurnal and  
 440 the multi-day cyclicity patterns in the data. Our end-member modelling analysis has helped  
 441 to identify and quantify the shape of the underlying frequency distributions and their  
 442 contributions to the resulting normalised HVSR time series (Fig. 3). We interpret EM<sub>1</sub>,  
 443 which predominantly contributes to the data during the night time (see downwards point-  
 444 ing arrows in Fig. 3 b) as indicative of a system with a lower resonance. This is supported  
 445 by the evolution of the fundamental frequency of the horizontal component (Fig. 2 b)  
 446 but at a higher temporal and spectral resolution. The lower resonance frequency may  
 447 be caused, as described above, by thermal shrinkage which results in opening of small  
 448 cracks and an overall decrease in the stiffness of the summit rocks. Vice versa, EM<sub>3</sub> is  
 449 the high resonance mode of the system, predominantly active during day time (green line  
 450 and upwards pointing arrows in Fig. 3 b). EM<sub>2</sub> can be regarded as the intermediate state  
 451 between EM<sub>1</sub> and EM<sub>3</sub>. We interpret EM<sub>4</sub> as noise or residual end-member, account-  
 452 ing for the variability within the data set that is not covered by the three other end-members  
 453 (for details see E. Dietze & Dietze, 2019). During the second survey period, later in the  
 454 year, EM<sub>4</sub> had a different frequency range, around 2–5 Hz, also visible in the raw HVSR  
 455 data (Fig. 2 c). This frequency band had no special weight during the first survey pe-  
 456 riod. Overall, the model quality ( $R^2 = 0.89$  for the first and  $R^2 = 0.79$  for the sec-  
 457 ond survey period) is comparable with results from other studies (cf. E. Dietze & Di-  
 458 etze, 2019), implying that 89 % and 79 % of the system’s variance can be explained by  
 459 just four end-members.

460 Between the two monitoring intervals, end-members EM<sub>1</sub> to EM<sub>3</sub> moved to higher  
 461 frequencies. This is likely due to the seasonally increasing temperature effect: overall higher  
 462 temperatures within the rock result in thermal expansion and a stiffening of the rock mass.  
 463 Likewise, the night time EM<sub>1</sub> contributed systematically less to the data set in the sec-  
 464 ond survey period (decrease from 42 % to 19 %, see Fig. 3 d), while the day time EM<sub>3</sub>  
 465 gained importance (from 19 % to 27 %). We attribute this to the decreasing time for the  
 466 system to swing and remain in the low frequency mode during night time later in the  
 467 year, due to the overall thermal input and heat legacy, and despite the gradually shorter  
 468 day lengths. In the first survey period EM<sub>4</sub> contributions were slightly stronger during  
 469 the episodes of signal disturbance (yellow bars in Fig. 2). However, since we clipped the

470 HVSR data at 15 Hz to isolate the part of the frequency spectrum that is dominated by  
 471 daily variations, interpretation of this trend may be not meaningful. During the second  
 472 monitoring interval, this effect was not found.

473 In summary, we consider the HVSR data (Fig. 2 c) and the estimated end-member  
 474 loadings (Fig. 3 d) and scores (Fig. 3 b) to be a high resolution proxy for reversible ma-  
 475 terial stiffness, as constrained by thermodynamics, predominantly at the diurnal scale,  
 476 but also with an underlying seasonal trend.

### 477 **5.1.3 Seismic wave velocity evolution**

478 Relative changes of the seismic wave velocity are the summed effect of many con-  
 479 tributing factors, and as for the fundamental frequency, these can be exogenic or endo-  
 480 genic in nature. Changes in seismic wave velocity can derive from variations of water sat-  
 481 uration or freezing in pore space and fractures, ambient stress, for example exerted by  
 482 gravity-driven deformation or ice pressure, and reflect changes in fracture extension and  
 483 density (Barton, 2007; Draebing & Krautblatter, 2012; Krautblatter & Draebing, 2013).  
 484 Temperature-driven changes are evident at the diurnal scale, with relatively low  $dv/v$   
 485 values at the start and end of the day and relative maxima in the middle of the day (Fig. 2 d).  
 486 In contrast, the longer-term trends in temperature are not reflected in the  $dv/v$  data set.  
 487 Transitions from positive to negative temperatures, which would cause freezing of wa-  
 488 ter and thus a drastic increase of the seismic wave velocity (James et al., 2019), did not  
 489 affect the  $dv/v$  time series. Elsewhere, the effect of rainwater infiltration is a lowering  
 490 of the relative seismic wave velocity, both for deep groundwater systems (Clements &  
 491 Denolle, 2018) and shallow soil water content in the vadose zone (M. Dietze et al., sub-  
 492 mitted). However, rain (Fig. 2 e) had no easily discernable effects on the  $dv/v$  data from  
 493 the Hochvogel summit network. We remind that we have no direct rain measurements  
 494 from the Hochvogel summit but used the rain gauge data from a station 12 km to the  
 495 west. Under the rapidly changing and filigree weather conditions in the Alps, the rep-  
 496 resentativeness of this rain data set may be limited. However, in the  $dv/v$  data we see  
 497 very few erratic changes that could be attributed to unrecorded rain events.

498 More importantly, we found cyclic patterns of slowly rising and suddenly dropping  
 499  $dv/v$  values, with an amplitude of about 0.2 % (Fig. 2 d, bold black line), coincident with  
 500 the periods of disturbance of the other proxy data. In the absence of meteorological forc-  
 501 ing, we attribute these cycles to reversible, rock-internal material state changes. Increased  
 502 stress has been reported to yield higher  $dv/v$  values (Sens-Schönfelder & Larose, 2010),  
 503 and significant reductions in  $dv/v$  values have been recorded immediately prior to the  
 504 initiation of mass wasting processes (Mainsant et al., 2012). Following these findings, we  
 505 interpret meteorologically unrelated and reversible  $dv/v$  cycles as proxy for stress accu-  
 506 mulation and dissipation within the rock mass.

### 507 **5.1.4 Emission of discrete seismic signals**

508 The different types of short seismic signal pulses recorded by the Hochvogel sum-  
 509 mit network (Fig. 4) have distinct properties, which make them straightforward to at-  
 510 tribute to likely sources. Repeated, regularly spaced pulses with gradually increasing and  
 511 decreasing amplitude (Fig. 4 a) are typical for people approaching and passing a seis-  
 512 mic station (e.g., Vouillamoz et al., 2018). Several such sequences were recorded during  
 513 installation and maintenance visits of the stations, and others registered only during day  
 514 time when hikers frequent the summit. Spindle shaped continuous signals with rapidly  
 515 increasing and less rapidly decreasing amplitudes lasting several seconds (Fig. 4 b) are  
 516 indicative of small mass wasting events down the summit flanks. Similar signals have been  
 517 described in detail by numerous other studies (e.g. Helmstetter & Garambois, 2010; Hi-  
 518 bert et al., 2011; Lacroix & Helmstetter, 2011; Burtin et al., 2016) and their frequency  
 519 spectrum is similar to artificial rock drop experiments (Weber, Fäh, et al., 2018). High-

520 frequency signals unique to independently registered rain events (Fig. 4 c) are evident  
 521 at all sensors but with random, incoherent drop impact signal timings at stations of the  
 522 network. The pauses between the 0.05–0.2 s long pulses decrease after the onset of a rain  
 523 event and become less evenly spaced. The frequency content of the signals is also higher  
 524 (60–80 Hz) than for other recorded signals. We interpret such low amplitude signals as  
 525 rain drop impacts at or very close to the seismic sensors. Finally, sudden pulses of 30–  
 526 80 Hz signals without preceding or subsequent activity (Fig. 4 d) are most likely the re-  
 527 sult of sudden releases of accumulated stress by failing rock bridges or extension of cracks  
 528 in the rock mass. Such signals have been reported under natural (Weber, Faillettaz, et  
 529 al., 2018; Vouillamoz et al., 2018) and laboratory (Murton et al., 2016) conditions. We  
 530 use this type of seismic signal as proxy for stress release, e.g. due to rock cracking.

531 Stress release signals had durations of  $0.58_{-0.02}^{+0.44}$  s (median and quartile range) dur-  
 532 ing the first, and  $1.57_{-0.48}^{+0.95}$  s during the second period (Fig. 5 a). This increased dura-  
 533 tion of pulses may be due to the greater average travel distance of a signal to a sensor:  
 534 only 2 stations instead of up to six during the first period, which results in extension due  
 535 to scattering as waves travel through the rock medium, especially given that during the  
 536 second period the two stations were separated by the main crack through the summit.  
 537 More interestingly, stress release events appear to be linked to daylight conditions or,  
 538 more specifically, potential insolation duration. During July–August, topography cor-  
 539 rected sunrise was between 05:30 and 05:45 local time and sunset between 21:05 and 21:20,  
 540 yielding a duration of potential insolation of between 15:20 and 15:45 hours. In contrast,  
 541 during September–October sunrise was between 06:45 and 07:30 and sunset between 18:45  
 542 and 19:45, separated by 11:15 to 13:00 hours of potential sunlight (see horizontal bars  
 543 in Fig. 5 b). Reflecting these insolation times, stress release events occurred preferen-  
 544 tially between 05:00 and 22:00 local time during the first recording period, with a mul-  
 545 timodal activity distribution. During the second survey period such events clustered be-  
 546 tween 7:00 and 21:00, with a pronounced maximum between 09:00 and 12:00.

547 At the multi-day scale, stress release events were not randomly distributed. Instead,  
 548 we observed discrete bursts of a 150–600 % activity increase above a relatively stable base  
 549 line of 100 cracks per day (median) for the first recording period, and 18 cracks per day  
 550 during the second period. These clusters coincided with the episodes of disturbance in  
 551 the continuous proxy records (Fig. 2). During these episodes, the fundamental frequency  
 552 of stations dropped, the bimodal frequency pattern of HVSR data was interrupted, and  
 553 the local seismic velocity decreased.

554 An exception to the collocation of continuous seismic proxy adjustment and bursts  
 555 of cracking is the prominent spike in stress release activity on 25 July. This could sig-  
 556 nal that there is a hierarchy of stress build up and release cycles within the Hochvogel  
 557 summit, which has not been covered in its entirety by our short surveys. Alternatively,  
 558 stress release in the summit is a process with cyclic and stochastic elements.

559 During the first survey period there were up to seven sensors in operation, with some  
 560 brief, incidental intermissions of data transmission. In contrast, during the second sur-  
 561 vey period only two stations were operating, continuously and without any data gap. The  
 562 average number of stations that detected the same seismic event was 3.5 during the first  
 563 and 2 during the second survey period. Hence, there may be a network bias inherent to  
 564 the time series. However, even if the number of true events were underestimated dur-  
 565 ing the second period, it is unlikely that the general pattern of a significantly reduced  
 566 event rate and overall number during that period is an artifact. Therefore, we conclude  
 567 that there was a stark contrast in stress release activity between the first and the sec-  
 568 ond survey period. A robust interpretation of this difference would only be possible in  
 569 the context of a more extensive data set.

## 5.2 Synthesis

With the interpretation of individual proxy data in hand, we can now investigate underlying system-wide dynamics. On the diurnal time scale, most likely due to thermal forcing, the Hochvogel summit experiences a cyclic increase and decrease of stress, but also strengthening and weakening of the coupling between the strongly fragmented southeastern and the less severely fragmented northwestern parts. Stress evolution is reflected throughout both survey periods and by multiple proxy data: i) the  $dv/v$  data, which records the effects of increased stress as higher  $dv/v$  values, and ii) discrete stress release signals predominantly during daylight conditions. The coupling evolution is evidenced by i) the vertical fundamental frequency of the northwestern summit part, ii) the horizontal fundamental frequency of the southeastern instable summit part, and iii) the HVSR end-member time series. The predominant effect of thermal forcing may be interpreted based on i) the clear diurnal air temperature signal in line with the patterns of all seismic proxy data, ii) the agreement of day lengths and stress release signals at the seasonal scale, and finally iii) the absence of evidence for any other plausible forcing mechanism.

The short sub-diurnal reaction time of the system implies that both stress release and coupling adjustment must be implemented at or very close to the surface, and are thus unlikely to have a significant effect on the wholesale slope instability, especially at many metres depth. Heat as driver for this system is in agreement with the effects that we have measured seismically. That heat can be provided either by direct insolation or as sensible heat of the air circulating around the summit and into the cracks. It would be dissipated by conduction, either directly within the rock or through percolating water, which is able to communicate thermal effects deeper into the rock, and thereby attenuate the expressions of heat input close to the surface. Regardless of the mechanism, heat diffusion into the rock would result in time lags of several hours (e.g., 37 h for 50 cm and 60 h for 1 m, Mulas et al., 2020) and rapid dampening (e.g., Collins et al., 2018) of the thermal input signal. Indeed, the end-member evolution with day time (especially for EM<sub>1</sub> and EM<sub>3</sub>, Fig. 3 b) shows that the low frequency EM<sub>1</sub> usually decreases gradually with the course of the day, whereas the high frequency EM<sub>3</sub> correspondingly increases throughout day time.

The coupling adjustment shows no signs of progressive evolution. Continuously less cracks capable of closing during day time and thus a systematic decrease in the fundamental frequency is not evident from our data during the study interval. We rather see the opposite effect: a systematically rising frequency of the end-member's modal frequencies. However, at least one full year of monitoring would be required to investigate whether the fundamental frequency returns to the values after one annual cycle through all seasons, or whether there is an offset indicative of progressive weakening (cf. Weber, Fäh, et al., 2018).

Stress release signals in the form of short seismic pulses do also argue for an irreversible component in the consequences of the diurnal forcing regime. The thermally driven physical weathering very close to the surface contributes to flaking (Collins et al., 2018) and leads to extended water pathways, which in turn increases infiltration capacity and circulation dynamics, basic foundations for subsequent weathering processes. On top of diurnal cycles, the trends of  $dv/v$  values continuously increasing for four to seven days before dropping for a few days, imply a systematic conditioning of the Hochvogel summit that results in increased accumulated stress that is episodically released and thereby emits an increased amount of short seismic pulses beyond the typical daily rate. The  $dv/v$  proxy, especially as it is based on the 4–16 Hz frequency window, representing wave lengths of tens of metres, integrates mechanical effects significantly beyond just the top few millimetres of the rock mass (Snieder, 2004; Larose et al., 2015). Thus, together with the non-diurnal cyclicality pattern, we see evidence for a process that is disconnected from the

622 set of meteorological drivers and rather points at a mechanism that is active inside and  
 623 driven by endogenic dynamics of the instable rock mass.

624 Stick-slip behaviour has been found in several large mass wasting events prior to  
 625 catastrophic failure (e.g., Yamada et al., 2016; Poli, 2017; Schöpa et al., 2018). Whereas  
 626 this precursor signal appears to bear an exploitable potential for hazard early warning  
 627 strategies during the final part of the preparation stage of an event, little is known about  
 628 the earlier stages of preparation. This is because stick slip motion is overall hard to de-  
 629 tect when the intervals between slip events are long in relation to the monitoring period.  
 630 Seismically, slip events manifest as periods of increased seismic activity, predominantly  
 631 through the emission of short pulses ( $< 2$  s) of 30–80 Hz signals (Poli, 2017; Vouillamoz  
 632 et al., 2018; Weber, Faillettaz, et al., 2018). In the Hochvogel summit, we also evidence  
 633 1–2 day long periods of increased seismic activity exhibiting these pulse-like properties,  
 634 that result in a drop of the  $dv/v$  stress proxy, separated by several days of just background  
 635 activity during which we see the buildup of stress. We interpret this multi-day cyclic pat-  
 636 tern as an early stage of stick-slip evolution of the instable rock mass at the summit. We  
 637 propose that during the stick phase gravity pulls the rock mass down slope, causing elas-  
 638 tic deformation (background seismic pulse emissions), which causes buildup of stress (in-  
 639 crease of  $dv/v$ ) until eventually the rock mass slips and thereby emits more seismic sig-  
 640 nals and the stress within the rock mass is reduced. We anticipate that as failure prepara-  
 641 tion of the Hochvogel rock mass progresses, this frequency of slip events increases, and  
 642 immediately prior to failure, these patterns can turn into a continuous tremor-like sig-  
 643 nal with a frequency of several per second (Schöpa et al., 2018).

644 Throughout the analysed proxy data we identified superimposed cycles of mech-  
 645 anisms that drive the evolution of the slope instability. In addition, the period of increased  
 646 stress release pulses on 25–26 August (Fig. 2 f) points at a further mechanism, which  
 647 may be of stochastic or higher hierarchy nature. The differences in all proxy time series  
 648 between the first and second survey period may suggest also a seasonal component con-  
 649 tributing to the preparation phase of the failure. However, in order to reveal these po-  
 650 tentially additional components, and to fully constrain the progressive evolution towards  
 651 the final stage of the mass movement, we would need a longer, complete data set from  
 652 a constant and sufficiently dense seismic network on the summit – obviously a task of  
 653 future research at this site.

654 We have presented a multi proxy based anatomy of a major future slope failure dur-  
 655 ing a key stage of its preparation phase, which is inherently difficult to survey. We pro-  
 656 pose that it is possible to generalise the Hochvogel observations to also describe the fate  
 657 of other high mountain peaks prone to failure. Most seismic proxies can be measured  
 658 and interpreted in a similar way as for the case presented here. However, the outcomes  
 659 of this study have also shown that it will be key to design a system that is reliable and  
 660 capable of operating under the hostile environmental conditions encountered at this soli-  
 661 tary peak. Such a system should ideally also probe the surface expressions of the evo-  
 662 lution the rock mass undergoes, for example by dedicated direct measurements of crack  
 663 opening, reference point position tracking and surface as well as below-surface temper-  
 664 ature and moisture logs. That way, the seismically sensed effects can be linked more di-  
 665 rectly to other first order physical quantities and thus allow for a distributed, continu-  
 666 ous and effective survey of a hitherto enigmatic yet essential phase in the evolution of  
 667 landscapes dominated by steep slopes.

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