

# **Structure of the North Anatolian Fault Zone imaged via teleseismic scattering tomography**

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## **SUMMARY**

Information on fault zone structure is essential for our understanding of earthquake mechanics, continental deformation and our understanding of seismic hazard. We use the scattered seismic wavefield to study the subsurface structure of the North-Anatolian Fault Zone (NAFZ) in the region of the 1999 İzmit and Düzce rupture using data from an 18-month dense deployment of seismometers with a nominal station spacing of 7 km. Using the forward and backscattered energy following the direct *P*-wave arrival from teleseismic earthquakes, we apply a scattered wave inversion approach and are able to resolve fine-scale changes in lithospheric structure on scales of 10 km or less in an area of about 130 km by 100 km across the NAFZ. We find several crustal interfaces that are laterally incoherent beneath the surface strands of the NAFZ and evidence for contrasting crustal structures either side of the NAFZ, consistent with the presence of juxtaposed crustal blocks and ancient suture zones. Although the two strands of the NAFZ in the study region strike roughly east-west, we detect strong variations in structure both north-south, across boundaries of the major blocks, and east-west, parallel to the strike of the NAFZ. The shallow NAFZ is coincident with features detected on the crust-mantle boundary (Moho) and deeper into the mantle. We show that a dense passive network of seismometers is able to capture information from the scattered seismic wavefield and using a tomographic approach, resolving the fine scale structure of crust and lithospheric mantle even in geologically complex

regions. Our results show that major shear zones exist as narrow structures beneath the NAFZ throughout the crust and into the lithospheric mantle, suggesting a strong coupling of strain across these depths.

**Key words:** Seismology, Teleseismic Scattering, Tomography, North Anatolian Fault Zone.

## 1 INTRODUCTION

The North Anatolian Fault Zone (NAFZ) is one of the longest continuous continental strike slip fault systems on Earth. The NAFZ is a 1500 km long, right-lateral strike slip fault running the length of northern Anatolia, separating the Eurasian plate to the north from the deforming Anatolian region in the south accommodating the relative motion and deformation between these tectonic domains. Together with the East Anatolian Fault, the NAFZ accommodates the westward motion of Anatolia (McKenzie, 1972; Reilinger et al., 2006; Şengör et al., 2005; Barka, 1992) driven by kinematic extrusion of a rigid Anatolian block, or by the gradient of gravitational potential energy from the Anatolian plateau to the Hellenic Trench (England et al., 2016).

While the deformation at the surface is localized on faults (Hussain et al., 2016; Bürgmann & Dresen, 2008), the distribution of deformation throughout the crust and into the mantle remains unclear (Bürgmann & Dresen, 2008; Vauchez et al., 2012; Moore & Parsons, 2015). The structure and dynamics of fault zones such as the NAFZ are essential for our understanding of continental deformation and seismic hazard.

The NAFZ ruptured in a series of  $M \geq 6.7$  earthquakes during the 20th century from east to west (Stein et al., 1997) interpreted as stress transfer along the strike of the NAFZ from one earthquake bringing the next segment closer to failure. The two most recent events in the current series occurred in 1999 with epicentres in İzmit ( $M=7.6$ ) and Düzce ( $M=7.2$ ) (Barka et al., 2002; Gülen, 2002) with the fault rupture extending into the Sea of Marmara and the next anticipated event in the series posing a pronounced risk to the city of Istanbul.

To better understand the structure of the NAFZ, especially in the middle and lower crust and

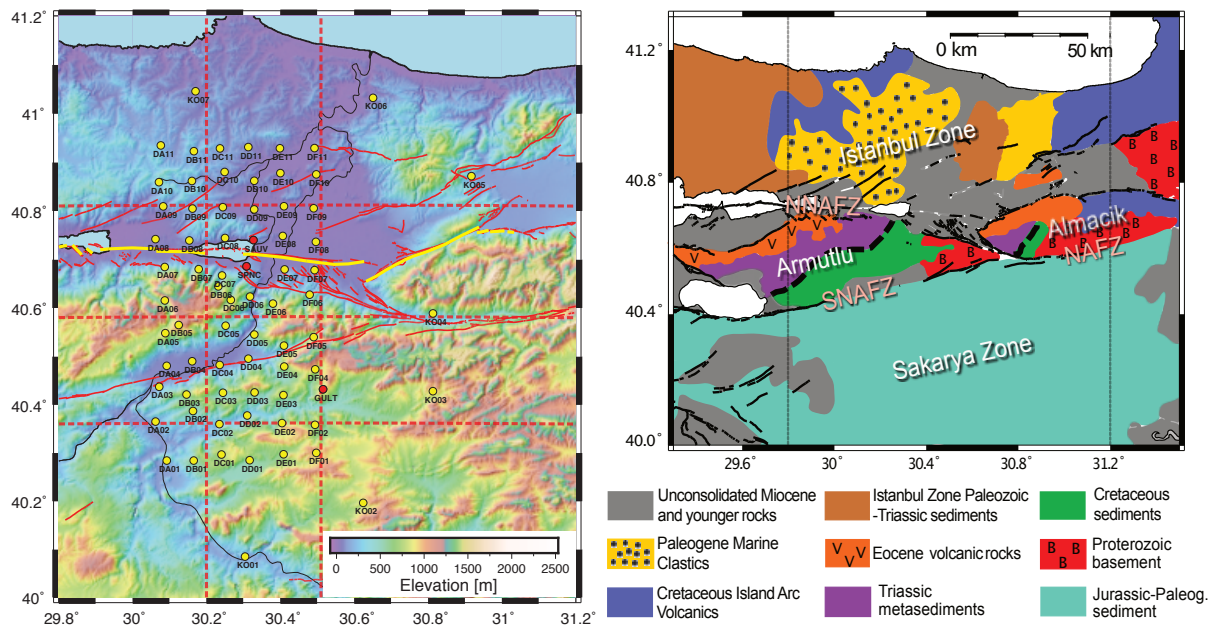
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22 into the upper mantle, we exploit the scattered seismic wavefield following the *P*-wave arrivals  
23 of teleseismic events (Frederiksen & Revenaugh, 2004), allowing us to resolve the fine scale  
24 structure of the lithosphere in the study region using data from the 18-month DANA deploy-  
25 ment (DANA, 2012) deployed across the NAFZ in the region of the 1999 ruptures (Fig. 1a). The  
26 *P*-wave coda contains energy from *P*-to-*P* and *P*-to-*S* scattering at small-scale heterogeneities  
27 along the ray-paths. Structure can be recovered from the scattered seismic energy through mi-  
28 gration approaches ranging from common-conversion-point or common-scattering-point stacking  
29 (e.g. Dueker & Sheehan, 1997) to full depth migration (e.g. Ryberg & Weber, 2000). Here we are  
30 using a tomographic waveform approach based on linear inverse theory of the scattered wavefield  
31 (Ji & Nataf, 1998; Frederiksen & Revenaugh, 2004).

32 We find that the two strands of the NAFZ evident in the shallow structure coincide with main  
33 interfaces and interface terminations throughout the crust and into the upper mantle indicating that  
34 the fault zone structure may extend to depths of at least  $\sim 75$  km in this region. We find evidence  
35 for small-scale variation of structure in the vicinity of the strands that might indicate the detection  
36 of heterogeneity related to past deformation along the present day fault.

## 37 **2 TECTONICS AND PREVIOUS GEOPHYSICAL SURVEYS**

38 The study region (Fig. 1b) is an amalgam of continental and subduction-related oceanic fragments  
39 that remain after the closing of the Tethyan Ocean in the late Tertiary (e.g. Okay et al., 2008;  
40 Robertson & Ustaömer, 2004) and the NAFZ diverges into a northern and southern strand west  
41 of about  $30.65^\circ\text{E}$ . There are three main tectonic blocks sampled by the DANA network: (i) The  
42 Istanbul-Zonguldak Zone (IZ) to the north of the northern strand of the NAFZ, (ii) the Sakarya  
43 zone (SZ) to the south of the southern strand and (iii) the Armutlu and Almacık blocks (AA)  
44 between the two strands. The northern (NNAFZ) and southern strands (SNAFZ) in the study re-  
45 gion run roughly along the Armutlu-Istanbul and Sakarya-Armutlu block boundaries, respectively  
46 (Emre et al., 2018). Slip on the northern and southern NAFZ strands has been estimated to be  
47 approximately 16-25 mm/yr and 5-19 mm/yr, respectively (Stein et al., 1997; Flerit et al., 2003;  
48 Meade et al., 2002). The northern branch of the NAFZ ruptured in the 1999 İzmit earthquake (Tibi



**Figure 1.** Study region. (a) Topographical map of the study region using SRTM data (Farr et al., 2007). Stations are indicated by yellow circles (permanent stations in red). Mapped faults (red lines) provided by Emre et al. (2018) and rupture of the 1999 Izmit and Düzce earthquakes (yellow) provided by Gülen (2002). Dashed north-south and east-west lines indicate location of depth profiles shown in Fig. 6 and 7 and are approximate locations of depth profiles provided by Kahraman et al. (2015). (b) Geological map of the region outlining the three main tectonic blocks and geological areas. After Taylor et al. (2019).

et al., 2001; Barka et al., 2002) and still shows active slip at the surface (Hussain et al., 2016) although current seismicity is not focussing beneath either fault strand (Altuncu Poyraz et al., 2015). In general, the NAFZ seems to follow the Intra-Pontide Suture (IPS) between the Istanbul and Sakarya Zone (e.g. Okay et al., 2008) with the complication of the Almacik and Armutlu blocks lying between these two continental fragments in the study area, with a distinctly different near-surface velocity structure (Taylor et al., 2019). The structure of the Almacik and Armutlu blocks are potentially the cause for the splay (e.g. Akbayram et al., 2016).

The study region has been sampled by two seismic refraction experiments crossing all major tectonic blocks, finding crustal thicknesses of  $32 \pm 2$  km at about  $30.10^\circ$  longitude in the east (Bekler & Gürbüz, 2008) and  $\approx 38$  km at about  $29.25^\circ$  longitude in the west of our study region (Karahan et al., 2001). The refraction data show evidence for a 5 km thick upper crustal layer with  $P$ -wave velocities ( $V_P$ ) of 5.6 to 6.1 km/s and lower crustal velocities of  $V_P = 6.7 - 7.2$

61 km/s (Bekler & Gürbüz, 2008). The upper mantle between 29.5° and 30.0° E is relatively slow  
62 ( $V_P = 7.6$  km/s) (Bekler & Gürbüz, 2008), while Karahan et al. (2001) find higher mantle velocities  
63 further east (29.0° to 29.5° E) with velocities of  $V_P \approx 8.1$  km/s. There is evidence for seismic  
64 discontinuities at crustal depths of  $\sim 17$  km and  $\sim 24$  km (Karahan et al., 2001).

65 A study using local earthquake waveforms (Horasan et al., 2002) finds a Moho depth at 32 km  
66 in the Marmara region. They find crustal discontinuities at 4 km and 17 km depth with upper  
67 mantle velocities of 8.00 km/s and 4.6 km/s for  $V_P$  and  $V_S$ , respectively and mantle densities of  
68  $3.35$  g/cm<sup>3</sup>. Upper crustal velocities are found to be 5.8 km/s and 3.4 km/s for  $V_P$  and  $V_S$ , lower  
69 crustal velocities of 6.2 km/s and 3.6 km/s and near-surface velocities of 3.5 km/s and 2.2 km/s  
70 (Horasan et al., 2002).

71 *P*-wave receiver functions (PRFs) east of the Marmara sea indicate a deepening of the Moho  
72 from west (29 to 32 km) to east (34 to 35 km) (Zor et al., 2003; Vanacore et al., 2013). The average  
73 crustal  $V_P/V_S$  in our study region is  $\sim 1.75$  (Vanacore et al., 2013). PRFs of the DANA dataset  
74 (Kahraman et al., 2015) find crustal thickness and  $V_P/V_S$  variation in both EW and NS directions  
75 with the crust deepening from 36.5 km ( $V_P/V_S = 1.73$ ) to 40 km ( $V_P/V_S = 1.73$ ) in the IZ, a  
76 constant crustal thickness of  $\sim 37$  km ( $V_P/V_S = 1.69$  to  $1.70$ ) in the AA, and a slight thinning from  
77  $\sim 35$  km ( $V_P/V_S = 1.73$ ) in the west to  $\sim 34$  km ( $V_P/V_S = 1.85$ ) in the east of the SZ (Fig. 1).  
78 Combining data from several permanent stations and temporary station deployments, including  
79 DANA data, Jenkins et al. (2020) determined Moho depths across the Sea of Marmara region  
80 finding thick crust of up to 41 km in the IZ, with a shallower Moho (32-34 km) in the AA and  
81 SZ with a sharp step in Moho depth across the NNAFZ. The transition also shows complex Moho  
82 structure around the NNAFZ. Additionally, Jenkins et al. (2020) find east-west variation with a  
83 general deepening of the Moho towards the east.

84 There is evidence in previous studies for strong crustal heterogeneity on scales of less than  
85 10 km with sharp truncations of sub-horizontal interfaces coinciding with the surface locations of  
86 the northern and southern NAFZ strands. The northern strand seems to penetrate deeper into the  
87 crust and may extend into the upper mantle (Kahraman et al., 2015). Similar Moho structure has  
88 also been imaged using ambient noise auto-correlations (Taylor et al., 2016).

89 Using *P*-wave transfer functions and a grid-search inversion approach Frederiksen et al. (2015)  
90 detected a sharp change of crustal thickness across the northern NAFZ which is believed to follow  
91 the trace of the Intra-Pontide suture in this location and a change of the  $V_P/V_S$  ratio across the  
92 southern branch indicating a change in basement composition. The IZ shows thick crust but low  
93 topography indicating that it is in isostatic disequilibrium or underlain by thicker lithosphere, a  
94 result supported by Jenkins et al. (2020). The transfer functions also provide evidence for thick  
95 sediments in Sakarya and Pamukova basins in agreement with ambient noise analysis (Taylor et al.,  
96 2019).

97 Teleseismic *P*-wave and *S*-wave tomography using the DANA dataset (Papaleo et al., 2017,  
98 2018) shows narrow low-velocity zones originating at the surface trace of both strands of the NAFZ  
99 and extending throughout the crust. These low velocities were interpreted as zones of localised  
100 shear related to the active faults which likely extend into the upper mantle, although it appears to  
101 widen with depth as lateral resolution decreases (Papaleo et al., 2017, 2018).

102 Magnetotelluric (MT) data show differences in the crustal conductivity from south to north  
103 across the NAFZ (Tank et al., 2005) with a high resistivity ( $\geq 1000 \Omega\text{m}$ ) crustal basement in the  
104 IZ to the north and a less resistive crustal basement ( $500 \Omega\text{m}$ ) in the SZ. The MT data resolve a  
105 narrow conductive zone (30 to  $50 \Omega\text{m}$ ) within the AA that extends into the upper mantle that has  
106 been attributed to partial melts or pore fluid flow from the upper mantle beneath the NAFZ.

### 107 3 METHOD

108 We apply the teleseismic scattering tomography approach by Frederiksen & Revenaugh (2004)  
109 to the DANA dataset to resolve the small-scale structure beneath the array. The scattered seis-  
110 mic wavefield is more sensitive to short-wavelength variations in material properties than is the  
111 path-integrated sensitivity of transmitted phases such as used in e.g. seismic traveltime tomogra-  
112 phy. The *P*-to-*p* and *P*-to-*s* scattered energy in the coda of teleseismic *P*-waves travelling along  
113 different paths to the main arrival can uniquely determine Earth structure if the sampling of the  
114 seismic wavefield is dense enough to avoid spatial aliasing. In the tomographic approach some  
115 aliasing can be accepted without introducing issues with non-uniqueness of the solution due to

116 the regularization of the problem. Several approaches to use the scattered coda energy to image  
 117 the subsurface have been developed, forming a continuous spectrum of method complexity. The  
 118 common approach of receiver function analysis uses stacked records of *P*-to-*s* (or *S*-to-*p*) scattered  
 119 (converted) energy (Vinnik, 1977; Langston, 1979) which may be binned according to their com-  
 120 mon conversion point to improve signal-to-noise ratio (Dueker & Sheehan, 1997) and mapped to  
 121 depth. The method assumes a 1D stratified seismic structure which is often violated in practice  
 122 (Rondenay, 2009). Lateral variation of structure leads to diffraction of the seismic wavefield and  
 123 diffraction stacking, a backprojection of the diffracted energy along its traveltime hyperbola, can  
 124 be used to image small-scale perturbations of the structure at depth. These methods are widely  
 125 used in controlled-source type applications (Yilmaz, 2001), and are commonly described as mi-  
 126 gration techniques (Rondenay, 2009) but implementation requires dense spatial sampling of the  
 127 seismic wavefield. General improvements and densification of recent passive seismic deployments  
 128 make the application of more complex methods, such as traveltime stacking of the scattered wave-  
 129 field (Revenaugh, 1995) or the application of inversion or backprojection operators in a 2D or 3D  
 130 model space (Bostock & Rondenay, 1999) possible and allow higher resolution of detail. For a full  
 131 review of these methods see Rondenay (2009).

132

133 For a more complete treatment of the scattering problem, the scattering image problem can be  
 134 formulated as a tomographic inversion (Ji & Nataf, 1998). Using a waveform inversion, Frederik-  
 135 sen & Revenaugh (2004) have developed a linear tomographic inversion of the scattered seismic  
 136 wavefield which we apply here. A full description is given in Frederiksen & Revenaugh (2004)  
 137 and we outline only the main points of this approach here.

138 In the standard derivation of seismic scattering in the Born approximation, assuming single  
 139 scattering, the scattering properties are represented as perturbations in elastic parameters ( $\delta\lambda$ ,  $\delta\mu$ ,  
 140  $\delta\rho$ ) to a background model ( $\lambda_0$ ,  $\mu_0$ ,  $\rho_0$ ). The seismic equation of motion in an isotropic medium is  
 141 given by:

$$\rho\ddot{u}_i = (\lambda\nabla \cdot \mathbf{u})_{,i} + [\mu(u_{i,j} + u_{j,i})]_{,j} \quad (1)$$



142 which can be expanded, using small perturbations to the elastic properties ( $\delta\lambda, \delta\mu, \delta\rho$ ) (Frederiksen  
143 & Revenaugh, 2004), to:

$$\begin{aligned} \rho_0 \ddot{u}_i - (\lambda_0 + \mu_0)(\nabla \cdot \mathbf{u})_{,i} = \\ - \delta\rho \ddot{u}_i + (\delta\lambda + \delta\mu)(\nabla \cdot \mathbf{u})_{,i} + \delta\mu \nabla^2 u_i \\ + (\delta\lambda)_{,i}(\nabla \cdot \mathbf{u}) + (\delta\mu)_{,j}(u_{i,j} + u_{j,i}) \end{aligned} \quad (2)$$

144 The wavefield can then be divided into a primary (background) and scattered component ( $\mathbf{u} =$   
145  $\mathbf{u}^0 + \delta\mathbf{u}$ ) with the unperturbed wavefield satisfying the unperturbed wave equation

$$\rho_0 \ddot{u}_i^0 = (\lambda_0 + \mu_0)(\nabla \cdot \mathbf{u}^0)_{,i} + \mu_0 u_{i,jj}^0 \quad (3)$$

146 Assuming that the scattered wavefield is much weaker than the unperturbed wavefield this gives  
147 the first-order Born approximation by discarding higher-order terms:

$$\rho_0 \delta \ddot{u}_i - (\lambda_0 + \mu_0)(\nabla \cdot \delta \mathbf{u})_{,i} - \mu_0 \nabla^2 \delta u_i = Q_i \quad (4)$$

148 with  $Q_i$  being a term of the unperturbed wavefield and the perturbed model parameters which is  
149 given as equation 13.22 in Aki & Richards (2002).

$$Q_i = -\delta\rho \ddot{u}_i^0 + (\delta\lambda + \delta\mu)(\nabla \cdot \mathbf{u}^0)_{,i} + \delta\mu \nabla^2 u_i^0 + (\delta\lambda)_{,i}(\nabla \cdot \mathbf{u}^0) + (\delta\mu)_{,j}[u_{i,j}^0 + u_{j,i}^0] \quad (5)$$

150 with  $\mathbf{u}^0$  being a solution for the unperturbed medium.

151

152 Assuming Rayleigh scattering, where the wavelength of the incident wavefield is much larger  
153 than the scale of the heterogeneity, the scattering problem reduces to a point scatterer and the full  
154 scattered wavefield is approximated by that of an array of point scatterers. Following Wu & Aki  
155 (1985), it is possible to derive expressions for the equivalent point source in Rayleigh scattering.  
156 These expressions also contain the directivity of the radiation of the scattered wavefield, and are  
157 provided as equations (7) to (10) in Frederiksen & Revenaugh (2004). This gives us the ability

158 to use both the amplitude and radiation pattern of scattering from small-scale heterogeneities in  
 159 solving the forward problem of the waveform inversion.

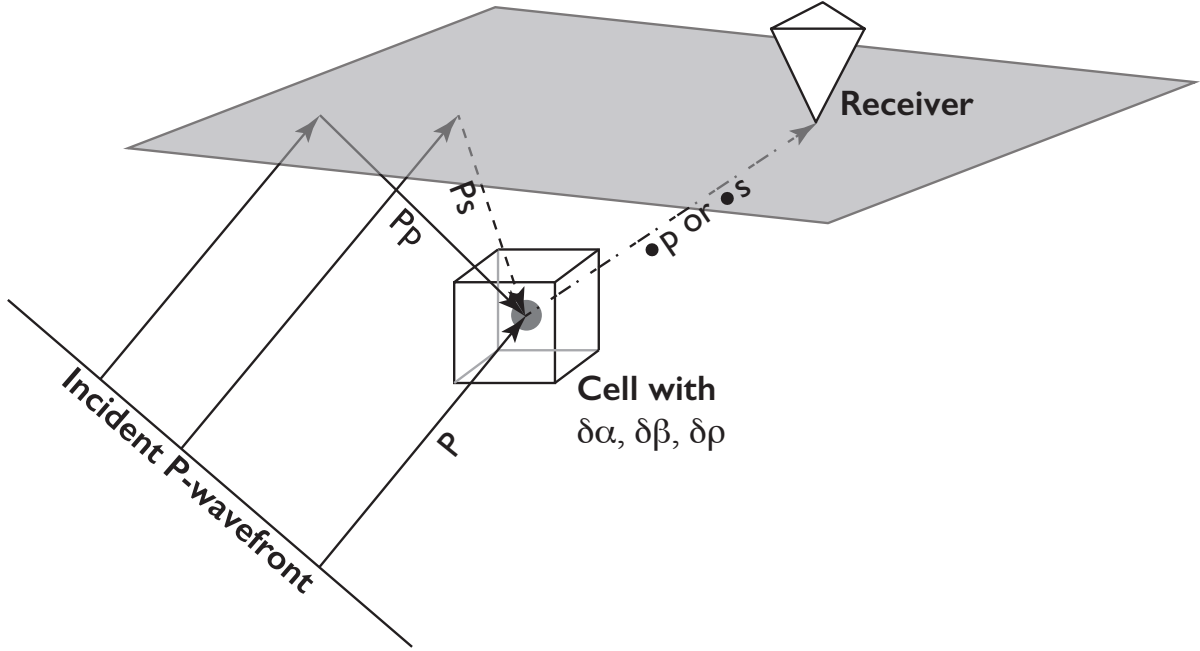
160 We assume the incident  $P$ -wave to be planar (Fig. 2) with a known slowness vector, a condi-  
 161 tion well met for teleseismic records. The scattered wavefield is derived from the seismic obser-  
 162 vations by computing the 3-component receiver functions relative to the first arriving  $P$  wave. The  
 163 considered input seismic wavefield includes the direct incident  $P$  and the free-surface reflections  
 164 ( $Pp$  and  $Ps$ ), producing forward and backscattering in the volume, respectively. The forward and  
 165 backscattering of the input wavefield produces six possible scattered phases (where  $\bullet$  indicates the  
 166 scattering event along the raypath) at small-scale elastic heterogeneities:  $P\bullet p$ ,  $P\bullet s$ ,  $Pp\bullet p$ ,  $Pp\bullet s$ ,  
 167  $Ps\bullet p$  and  $Ps\bullet s$ . In the forward modelling, we consider every possible combination of perturbed  
 168 parameter ( $V_P$ ,  $V_S$ ,  $\rho$ ), incident wave (forward scattering  $P$  and backscattering free surface reflec-  
 169 tion  $Pp$  and  $Ps$ ) and station location. The inclusion of the free surface backscattered energy as well  
 170 as the forward scattered direct wave increases the resolution of the study volume and allows us to  
 171 resolve a 3D perturbation model, here represented as a regular grid of perturbed cells. We use ray  
 172 tracing in a 1D velocity model to determine traveltimes to and from the scattering heterogeneity  
 173 and for incidence and refraction angles. We use equations (7) to (10) of Frederiksen & Revenaugh  
 174 (2004) including a geometrical spreading factor for a layered medium to determine the amplitudes  
 175 of the scattered energy in an elastic velocity model.

176 The Born approximation prescribes that single scattered waves propagate in the unperturbed  
 177 medium and do not interact with heterogeneities again. Therefore, the scattered wavefields from  
 178 individual heterogeneities are independent, which allows the simple summation of the contribu-  
 179 tions of individual heterogeneities to calculate the complete scattered wavefield  $T$

$$T = \sum_{i=1}^N \sum_{j=1}^3 t_{ij} \quad (6)$$

180 with  $t_{ij}$  representing the time series representing the scattering contribution of the  $j$ th perturbed  
 181 parameter of the  $i$ th scatterer (Frederiksen & Revenaugh, 2004).

182 The medium beneath the array is parameterised into a 3D grid of cells with each cell potentially  
 183 containing a perturbation of elastic parameters. The perturbation for all cells can be collapsed in



**Figure 2.** Sketch of the scattered phases included in the full waveform inversion. An incoming planar P-wave wavefront interacts with a cell with a parameter disturbance ( $\delta\alpha, \delta\beta, \delta\rho$ ) either from the direct wavefront (forward scattered) or the back-scattered wave from the free-surface reflection. The wavetype can convert upon scattering from P to S.

184 an  $M$ -element vector  $\mathbf{m}$  with dimension  $M = \# \text{ cells in } [x, y, z] \times \text{properties } [\Delta\alpha, \Delta\beta, \Delta\rho]$ .  
 185 Summing over all contributing elements we obtain the  $N$ -element vector  $\mathbf{d}^i$  with the number of  
 186 points depending on  $N = \text{samples} \times \text{stations} \times \text{components} \times \text{events}$ . The dependence of the full  
 187 scattered wavefield on arbitrary model  $\mathbf{m}$  is then described as

$$\mathbf{d} = \mathbf{A}\mathbf{m} \quad (7)$$

188  $\mathbf{A}$  is an  $N \times M$  matrix describing the sensitivity of each data point to each model parameter, i.e.  
 189 each column of  $\mathbf{A}$  represents a differential seismogram for a perturbation of a single parameter in  
 190 a single cell of the perturbed model. Equation (7) is linear and can therefore be solved using linear  
 191 inverse theory. To pose this problem as a damped inversion the inverse problem is formulated as a  
 192 minimization:

$$\min \left\| \begin{bmatrix} \mathbf{A} \\ \lambda \mathbf{I} \end{bmatrix} \mathbf{m} - \begin{bmatrix} \mathbf{d} \\ 0 \end{bmatrix} \right\|_2 \quad (8)$$

193 with  $\mathbf{I}$  being an  $M \times M$  identity matrix and  $\lambda$  a weighting factor, representing uniform damping. We  
 194 use the LSQR method (Paige & Saunders, 1982) to solve for the material properties in  $\mathbf{m}$ . For the  
 195 inversion of real data it has been found that regularisation by smoothing is preferable to damping  
 196 (Frederiksen & Revenaugh, 2004) as it provides results with higher coherence. Using LSQR, the  
 197 model is smoothed by posing  $\mathbf{m} = \mathbf{S}\mathbf{x}$  with  $\mathbf{S}$  being a matrix containing a Gaussian smoother. We  
 198 use  $\mathbf{B} = \mathbf{A}\mathbf{S}$  and the minimization

$$\min (\|\mathbf{B}\mathbf{x} - \mathbf{d}\|^2 + \lambda^2 \|\mathbf{x}\|^2) \quad (9)$$

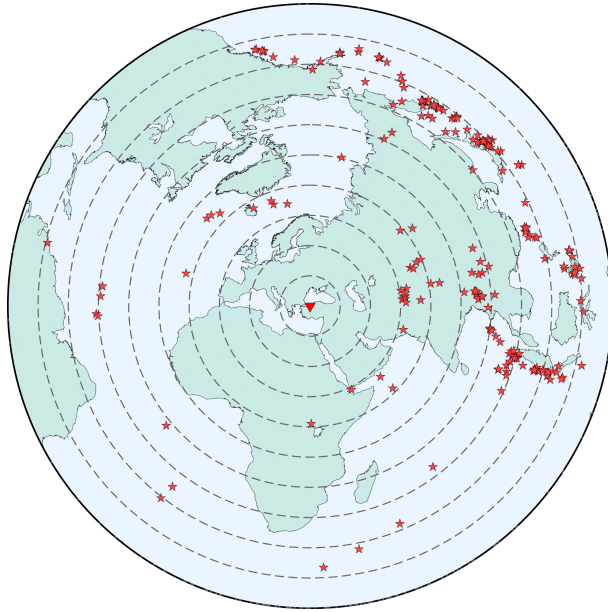
199 in which we solve for  $\mathbf{x}$  rather than  $\mathbf{m}$  (Van der Lee & Nolet, 1997; Frederiksen & Revenaugh,  
 200 2004). For all recovery tests and real data inversions, we apply a moving Gaussian smoother with  
 201 a standard deviation of one model element in the horizontal directions, but we do not smooth in the  
 202 vertical direction. No smoothing is applied beyond three standard deviations. This choice biases  
 203 the recovered model towards lateral coherence, making recovered lateral changes more coherent  
 204 in our study region where we expect strong lateral changes across the NAFZ.

205 The model space is parameterised as a regular grid with 5 km horizontal grid spacing and 2 km  
 206 vertical grid spacing with 30 (0-29) cells in horizontal directions and 60 (0-59) in vertical direction.  
 207 Each cell is treated as a point scatterer with vertical and horizontal locations at depths  $2 \cdot j$  km ( $j =$   
 208  $0, \dots, 59$ ) and longitude/latitude location of  $5 \cdot k$  ( $k = 0, \dots, 29$ ), respectively. The maximum grid  
 209 size is controlled by the maximum memory required to invert the dataset (see below). We tested the  
 210 method with doubled lateral and vertical grid spacing and do not find noticeable differences in the  
 211 general structure of the solutions except for obvious impacts on the maximum possible resolution  
 212 of the solutions.

## 4 DATA

We use passive seismic data from stations of the Dense Array for Northern Anatolia (DANA) that were installed across the NAFZ in the region of the 1999 İzmit and Düzce ruptures (DANA, 2012). DANA was deployed between May 2012 and October 2013 and stations were arranged in a quasi-rectangular region of 35 km by 70 km with a nominal station spacing of 7 km (Fig. 1). Stations were aligned along seven north-south oriented lines (labelled A to F) and 11 east-west lines (labelled 01 to 11). Seven additional stations were installed in an eastern semi-circle with a radius of about 60 km. Three permanent stations (SPNC, SAUV, GULT) of Boğaziçi University and Kandilli Observatory and Earthquake Research Institute/National Earthquake Monitoring Center (BU-KOERI/NEMC) located within the DANA network grid were included in the analysis. Stations were equipped mainly with Güralp CMG-6TD and CMG-3T medium broadband and broadband three-component instruments (full information on the network can be found in DANA (2012)). Data were sampled at 50 Hz.

We use earthquakes within the deployment period with  $m_b > 5.5$  from the catalogue of the National Earthquake Information Centre (NEIC) and angular distances of  $30^\circ$  to  $90^\circ$ . For the permanent stations we add events from 2009 onwards (in total 47 additional events contributing typically a single 3-component seismogram (ZRT) to the dataset). Low frequency noise was suppressed by applying a 2-way, 2-pole high-pass filter with cut-off frequency of 0.1 Hz. We calculate 3-component receiver functions (RFs) with a maximum frequency of 1.2 Hz using the time domain iterative deconvolution approach by Ligorría & Ammon (1999) deconvolving the Z component from the vertical, radial and transverse components. The calculated receiver functions were visually inspected to select events following these criteria: (1) transverse RFs show lower or comparable amplitudes than radial RFs, (2) the direct  $P$ -wave arrival is close to the predicted travel time for a 1D Earth model and (3) no evidence for large amplitude ringing. The pre-processing used to obtain the receiver functions is similar to the method used by Kahraman et al. (2015), but applied to all three components (vertical, radial, transverse) of the traces in our analysis. To remove the first arrival, which does not contain any additional structural information, we mute the first 2.5 s of each trace following the theoretical  $P$ -wave arrival.



**Figure 3.** Earthquake locations relative to the center of the DANA array. Dashed circles give distance from DANA center in  $10^\circ$  steps. Earthquakes with  $m_b > 5.5$  occurring during the deployment and since 2009 for the permanent stations in an epicentral distance of  $30^\circ$  to  $90^\circ$  were used in the analysis.

241 In total, we use 1396 traces from 176 events in our analysis. The distribution of sources is  
 242 shown in Figure 3. Traces were cut and tapered to 100 s and downsampled to 5 Hz (from the orig-  
 243 inal 50 Hz sampling). Despite the downsampling, the matrix to invert is very large which limits  
 244 the achievable resolution and model depth. Typical storage requirements for the matrix inversion  
 245 using the sparse storage method are  $\approx 338$  Gb for a model space dimension  $(x \times y \times z)$   $145 \times 145$   
 246  $\times 118$  km<sup>3</sup> with an element size of  $5 \times 5 \times 2$  km<sup>3</sup> and 1396, 100 s long traces, sampled at 5 Hz.  
 247 We are able to invert the full dataset without reverting to inverting subsets of data and stacking  
 248 the resulting images (Frederiksen & Revenaugh, 2004; Zhang & Frederiksen, 2013) leading to  
 249 improved image quality of our results.

250

## 251 5 RECOVERY TESTS

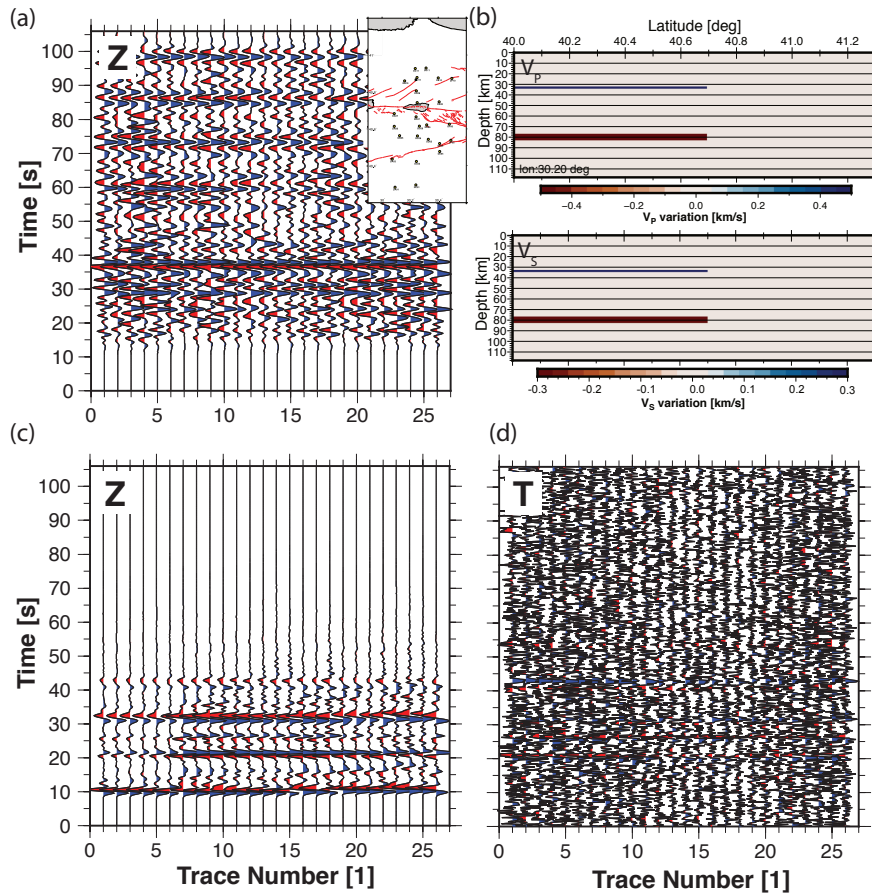
252 We tested several background velocity models for data inversion and synthetic data recovery in-  
 253 cluding models by Karahan et al. (2001), Bekler & Gürbüz (2008), and Horasan et al. (2002) and  
 254 models including constant velocity and linear vertical gradients. The background models are used

Depth (km)	$V_P$ (km/s)	$V_S$ (km/s)	$\rho$ (kg/m <sup>3</sup> )
0 - 2	3.6	1.967	2.376
2- 14	5.900	3.225	2.814
14 - 26	6.500	3.552	2.955
26 - 40	7.000	3.691	2.975
⋮	8.055	4.347	3.326

**Table 1.** 1D velocity model used in inversion. Depth,  $P$ -wave,  $S$ -wave and density ( $\rho$ ) following (Karahan et al., 2001)

for raytracing to determine traveltimes of the incident and scattered wavefield. While timing of arrivals changes slightly for all realistic velocity models, the overall recovered structure in our tests does not depend significantly on the choice of background model, although depths of interfaces change due to changes in the traveltimes. We chose to use the model by Karahan et al. (2001) for all inversions presented here (Table 1, Fig. 4). This velocity model is derived from seismic experiments in the study area and has been used in previous studies using this dataset (Kahraman et al., 2015; Altuncu Poyraz et al., 2015).

Figure 5 shows the result of an inversion of the full (1396 traces) noisy synthetic data generated through the perturbation model shown in Fig. 4b). A sub-set of the synthetic traces used in this inversion, i.e. the stations recording event 20123211812 and used in the data inversion, are shown in Fig. 4a) and d). Synthetic data were generated using ray tracing through the background velocity model with the addition of the scattered wavefield (i.e. the summation of all contributions of the single scatterers in the model). We use a 0.25 s wide Gaussian wavelet as the source time function. Synthetic tests use the source-receiver combinations for each event in the dataset, therefore recreating the same resolution as the recorded dataset. For comparison we show the recorded data in Fig. 4a) with the first arrival muted. Scattered phases can be seen coherently across the traces. The synthetic traces (Fig. 4c) show similar structure although clearly are not able to capture the full complexity of the data due to the simplicity of the model (Fig. 4b). Noise is added to the synthetic data through a random number generator (Marsaglia & Bray, 1964) using 10% RMS

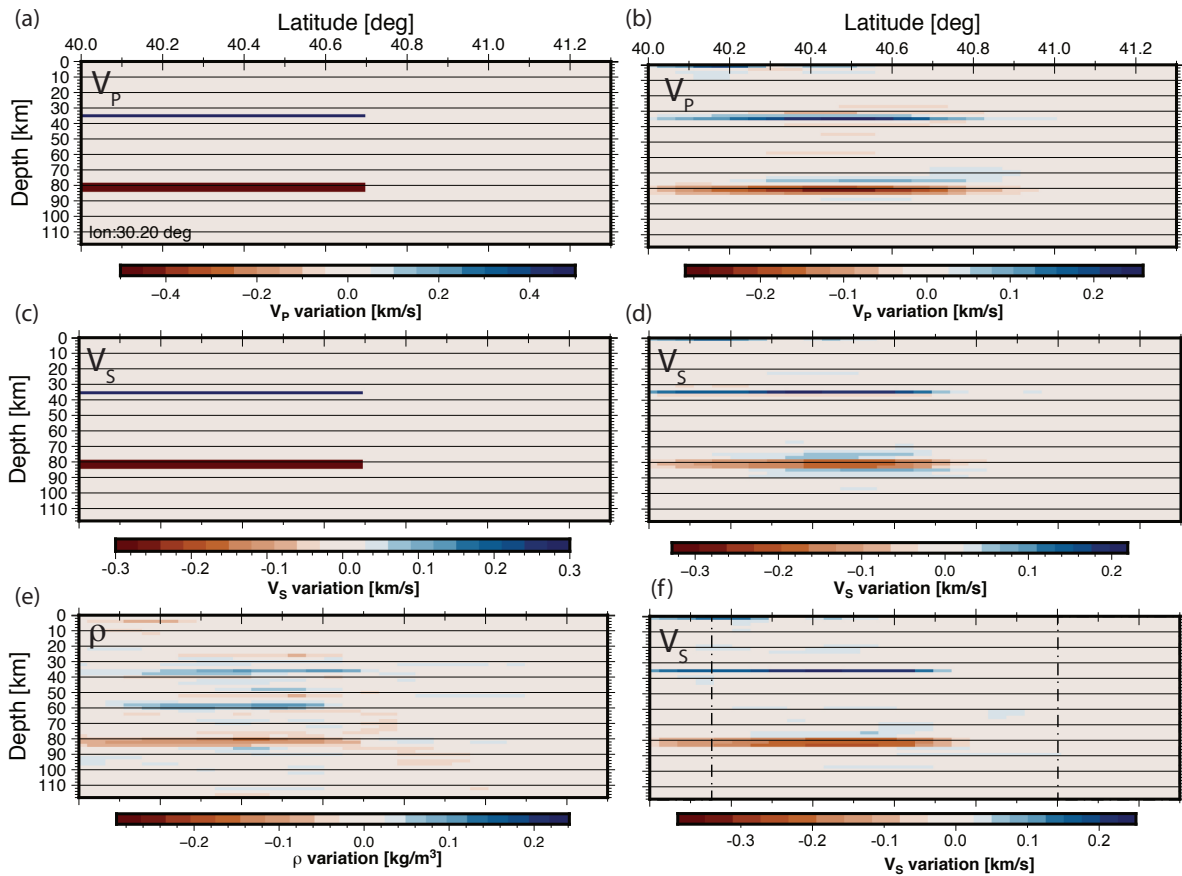


**Figure 4.** Example data and synthetics. (a) Recorded data example of vertical component of event 20123211812 (insert shows station configuration). Data are cut to 100 s and tapered. The direct P-arrival is suppressed. (b) North-South slices through perturbation model to calculate synthetics. (c) Noiseless synthetics (vertical) through perturbation model shown in b), including all scattering events from direct wave and free surface reflections. Synthetics represent the event-station configuration of the quality-controlled receiver functions for event 20123211812 shown in a). Although the perturbation is not continuous throughout the model, the scattered wavefield can be recorded across the network with distinctive moveout allowing the localisation of the perturbation. (d) Transverse component of the synthetic data calculated for model shown in (b) with added noise. We add Gaussian noise with a 10% standard deviation relative to the maximum signal amplitude to the synthetic data.

274 amplitude variation Gaussian noise compared to the direct wave amplitude to produce this noisy  
 275 synthetic dataset (Fig. 4d).

276 The synthetic model is parameterised with 5 km cell spacing horizontally and 2 km vertically.  
 277 The model contains a  $V_P = +0.5$  km/s and  $V_S = +0.3$  km/s anomaly for a single depth element  
 278 (2 km) starting at 34 km depth and a  $V_P = -0.5$  km/s and  $V_S = -0.3$  km/s anomaly with thickness





**Figure 5.** Recovery tests for model shown in Fig. 4 containing heterogeneous  $V_P$  and  $V_S$  structure. Slices were taken from the 3D model at longitude  $30.2^\circ$ . (a)  $V_P$  anomaly of input model containing terminating velocity anomalies at 32 km (negative) and 76 km (positive). Anomaly strength is  $\pm 0.5$  km/s. (b) Recovered model after scattering inversion. Due to damping the velocity recovered is smaller than the input model but locations are well resolved. Horizontal smoothing of a single lateral element (5 km) is applied. (c) as (a) but for  $V_S$ . Velocity anomaly is  $\pm 0.3$  km/s. (d) As (b) but for  $V_S$ . (e) Recovered low amplitude density ( $\rho$ ) anomaly after inversion. Input model does not contain density variations so recovered anomaly represents cross-talk between the different components. (f) Wiener filtered longitude slice of recovered model to fit recovered depth of anomaly as a zero-phase wavelet. Black dashed lines in f indicate the perceived limits of the well-resolved region.

279 of 6 km starting at 78 km (Fig. 5 a and c). No density variation was added to the model. The  
 280 anomalous layers extend across the model in longitude but terminate 80 km into the model from  
 281 the south (approximately at  $40.7^\circ\text{N}$ ), leaving the part of the model that corresponds to the Istanbul  
 282 zone free of a velocity anomaly. The inverted model in Fig. 5 b and d shows the recovery of  
 283 the input model. Because the inversion uses the Born approximation, which generates signals

284 from localized perturbation, the recovered model will be a band-pass filtered version of the input.  
 285 We apply a Wiener optimum filter to minimize the effects of the inversion process, mainly to  
 286 reduce sidelobes to aid interpretation. The optimization filter, as described for example by Gubbins  
 287 (2004), is obtained by minimising the residual between the desired output  $g_t$  (Fig. 5a,c) and the  
 288 signal obtained by convolution of the filter  $f_t^0$  with the actual output  $x_t$  (Fig. 5b,d)

$$(e_t)^2 = (g_t - x_t * f_t^0)^2 \quad (10)$$

289 The effect of the inversion and the filter terms acting on a single trace of the synthetic model  
 290 are shown in Supplemental Figure S1. Although the input model in this test does not contain any  
 291 density ( $\rho$ ) heterogeneity, Fig. 5e shows that the inverted model for the density structure is affected  
 292 by cross-talk between the different parameters (more examples given in Supplemental Figure S2).  
 293 However, relative amplitudes  $\Delta\rho$  in this model are small and the effect is most prominent in  
 294 areas with velocity anomalies. Tests with models including  $\Delta\rho$  show that density structure can be  
 295 resolved. Complete input and output models for this recovery test and further recovery tests are  
 296 shown in the Supplemental Figures S2 to S7.

297 These tests show that the recovery of velocity and density anomalies is variable within the  
 298 model volume due to the relative sampling of the model volume by the ray configuration of the  
 299 dataset. Peripheral regions are generally less well resolved than the center of the volume (Fig. 5).  
 300 Within the central zone we do not observe strong depth or amplitude variations of the recovered  
 301 model, adding confidence to our interpretation. Areas of the model space that are not well resolved  
 302 are masked in all following figures (and supplemental material) and the approximate limits of the  
 303 well-resolved volume are shown in the N-S profiles (dot-dashed vertical lines in Fig. 5f), to which  
 304 we limit our interpretation.

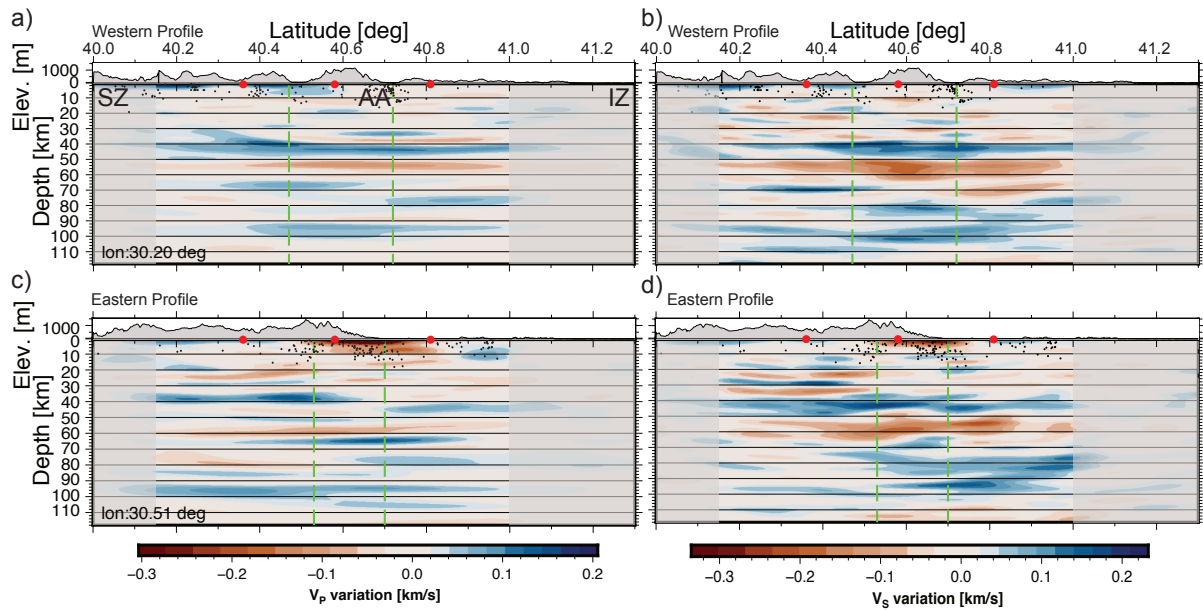
305 Changing the depth of the inverted model space between 48 km and 118 km (in 20 km steps)  
 306 does not lead to strong changes in the inverted model. A comparison between a 48 km and 118 km  
 307 deep model containing the same structure for  $V_P$  and  $V_S$  is shown in supplemental Figure S3 a)  
 308 and b). This holds even when synthetic traces were generated including structure below the in-  
 309 verted volume (Fig. S3c) showing that heterogeneities underneath the volume are not erroneously

310 mapped into the model volume. We show models down to depths of 118 km (60 nodes with 2 km  
311 spacing) in a trade-off between achievable resolution, model size and required computer memory.  
312 The horizontal smoothing leads to some smearing of energy in horizontal directions. Nonetheless,  
313 Fig. 5 shows that terminating discontinuities can be accurately located within 1 to 2 horizontal el-  
314 ements (i.e. 5 to 10 km) in the central region of the model space. We also performed recovery tests  
315 using other structural models including velocity and density heterogeneities to better understand  
316 the performance of the method (see Supplemental Material).

## 317 **6 RESULTS**

318 The results of the tomographic scattering inversion of the DANA dataset are shown in Fig. 6 and  
319 Fig. 7. Slices in Fig. 6 and 7 were extracted from the three-dimensional inversion volume along  
320 North-South (Fig. 6) and East-West (Fig. 7) profiles at locations shown in Fig. 1. The locations  
321 of the profiles were chosen to be in similar locations to those shown in Fig. 6 of Kahraman et al.  
322 (2015) (for an equivalent display to Kahraman et al. (2015) see Supplemental Figure S8, S10 and  
323 S12). Full solutions are presented in the form of animated GIFs in Figures S14 and S15 of the  
324 Supplemental Materials. The model is filtered with the Wiener optimization filter as discussed  
325 above.

326 Generally, the *S*-wave images show greater amplitude and are better constrained due to the  
327 combination of upgoing *Ps* conversions as well as the free-surface multiples, while the *P*-wave  
328 inversion solely relies on the free-surface multiples. The *S*-wave tomographic images also seem to  
329 show more fine scale structure likely related to the shorter wavelength. The density ( $\Delta\rho$ ) profiles  
330 show some of the major structure and are shown in the supplemental Figs. S12 and S14 but suffer  
331 from cross-talk as shown earlier. As the interpretation of the  $\Delta\rho$  profiles is more difficult and there  
332 is no independent constraint on the density structure we do not discuss this parameter further in  
333 the text.



**Figure 6.** North-south oriented slices along  $30.2^\circ$  (top) and  $30.51^\circ$  (bottom) for the inversion for  $V_P$  (left) and  $V_S$  (right) structure. Profile locations are indicated in Fig.1. Approximate locations of the southern and northern branches of the NAFZ are shown as dashed green lines. Areas with limited resolution as determined from the recovery tests (Fig. 5) are masked in grey. Red dots indicate locations of EW slices shown in Fig.7. Black circles indicate local seismicity as determined by Altuncu Poyraz et al. (2015) within a  $\pm 5$  km corridor projected onto the profile. The top panel in each subpanel shows SRTM topography along the profile (Farr et al., 2007).

### 334 6.1 Western Profile

335 Profiles for  $V_P$  (Fig. 6 a) and  $V_S$  (Fig. 6 b) have been extracted along a longitude of  $30.20^\circ$  E.  
 336 Areas with limited resolution as determined from the recovery tests (Fig. 5) have been masked  
 337 in this profile in transparent grey. The  $V_P$  profile (Fig. 6a) is dominated by a velocity increase at  
 338  $\sim 40$  km depth for most of the profile, which we associate with the Moho. Depths to interfaces are  
 339 measured at the top of the interface. The Moho velocity increase bifurcates south of  $\sim 40.4^\circ$  with a  
 340 shallower velocity increase located at  $\sim 32$  km depth deepening to 40 km at  $\sim 40.4^\circ$ N. The anomaly  
 341 also seems to fade, i.e. showing less of a velocity anomaly, south of about  $40.3^\circ$ N. The point of  
 342 bifurcation coincides with the surface expression of the southern strand of the NAFZ. A similar  
 343 Moho signal is observed in the  $S$ -wave anomaly at  $\sim 40$  km, shallowing to about 38 km within the  
 344 Armutlu block, which shows a thickening of this interface. The  $S$ -wave anomaly does not show

345 the same shallow branch observed in the P-waves but shows lower amplitudes south of  $\sim 40.4^\circ\text{N}$ ,  
 346 i.e. south of the southern NAFZ strand.

347 Observed crustal structure includes a weak high  $V_P$  anomaly at  $\sim 18$  km in the Armutlu block  
 348 with weak complex  $V_S$  structure in the Sakarya zone. Complex structure starting at  $\sim 32$  km depth  
 349 (positive and negative anomalies) in  $V_P$  and  $V_S$  can be seen in the vicinity of the northern strand  
 350 ( $40.7^\circ\text{N}$ ) just overlying the Moho. The  $V_P$  model shows less structure in the upper crust except a  
 351 fast anomaly to depths of  $\sim 5$  km around the southern branch and a slow (also seen in  $V_S$ ) overlying  
 352 fast anomaly between  $\sim 10$  km and  $\sim 20$  km depth in the Armutlu block.

353 The high velocity anomaly at 40 km depth is underlain by a strong low  $V_P$  and  $V_S$  anomaly  
 354 at depths of  $\sim 50$  km. This anomaly shows lower amplitudes in the Sakarya Zone with the change  
 355 coinciding with the surface expression of the southern NAFZ strand. We also identify a velocity  
 356 increase in  $V_P$  and  $V_S$  at  $\sim 64$  km and  $\sim 66$  km depth, respectively, around  $40.4^\circ\text{N}$  (southern  
 357 strand) and  $\sim 74$  km in  $V_P$  beneath the Istanbul zone (with a termination at the northern strand).  
 358 The  $V_S$  anomaly shows a low velocity anomaly at  $\sim 68$  km depth just north of the northern strand  
 359 changing to a high velocity anomaly at  $\sim 74$  km depth north of  $40.9^\circ\text{N}$ .

360 At greater depths we observe a fast anomaly in  $V_S$  at  $\sim 78$  km depth and a fast anomaly in  $V_P$   
 361 and  $V_S$  at  $\sim 92$  km but showing depth variation in  $V_S$ . The 78 km anomaly seems to merge with  
 362 the deeper anomaly in the Istanbul zone.

## 363 6.2 Eastern Profile

364 The eastern North-South profile at  $30.51^\circ$  E (Figs. 6c and 6d) shows more structure, especially in  
 365 the crust, than the western profile despite the close proximity of the two profiles.

366 We observe a strong, fast  $V_P$  anomaly at a depth of  $\sim 34$  km terminating halfway through  
 367 the Armutlu block and re-emerging at a depth of  $\sim 42$  km just north of the northern strand in  
 368 the Istanbul zone. In  $V_S$  we observe a more continuous structure with a high velocity anomaly at  
 369  $\sim 36$  km depth in the south stepping to  $\sim 42$  km at  $\sim 40.7^\circ\text{N}$  coinciding with the northern strand.  
 370 The  $V_P$  anomaly is weak in the Armutlu block on this profile and seems to terminate at  $40.6^\circ$  E,

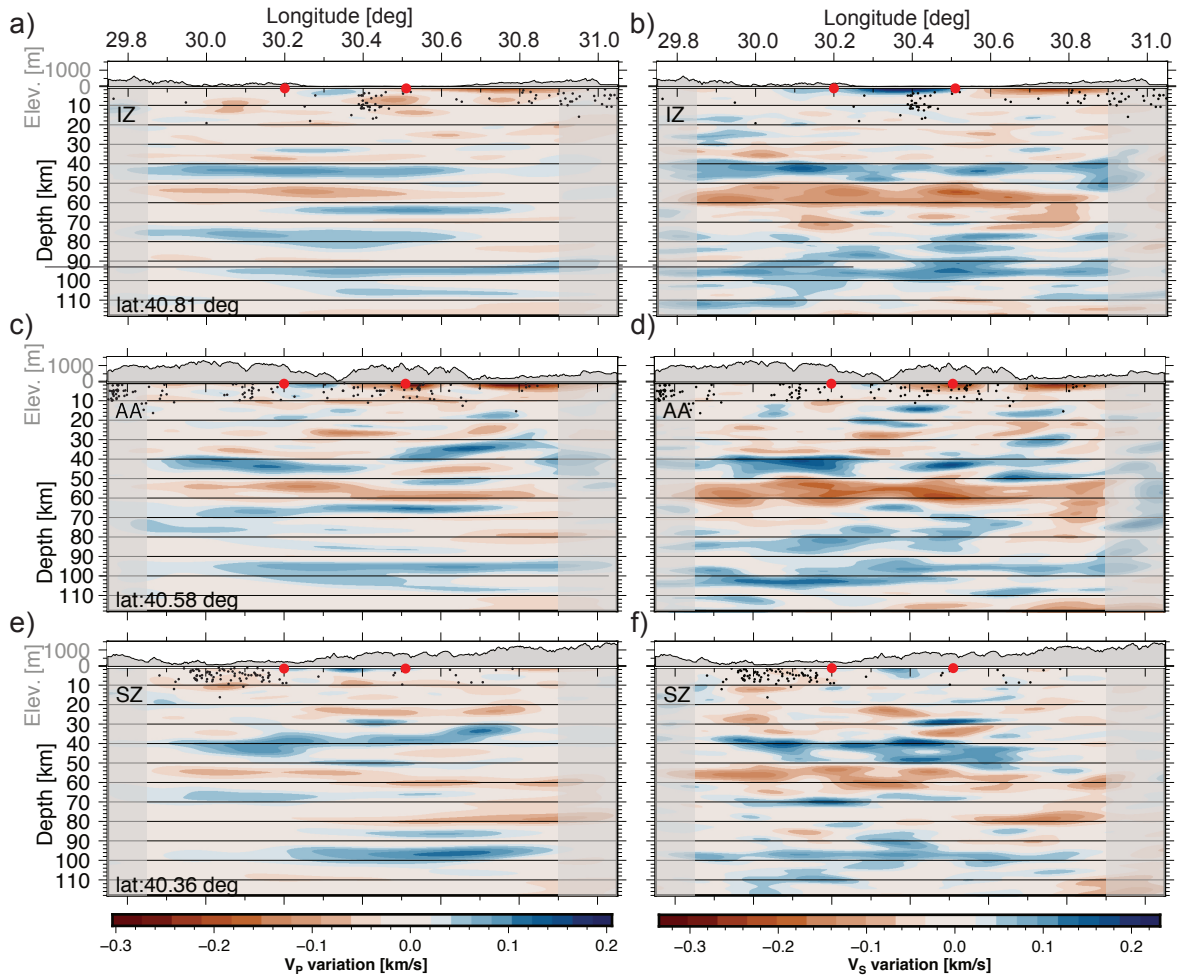
371 while the  $V_S$  anomaly is more continuous, but also weakens in this region. The amplitude variation  
 372 of these anomalies cannot be explained by the limitations of the sampling (see Fig. 5).

373 Especially striking in this profile is the complex  $V_S$  structure in the Sakarya Zone down to  
 374 depths of about 30 km manifesting as series of fast and slow anomalies between  $\sim 10$  km and  
 375 32 km (see supplemental Figure S7). The  $V_P$  structure is similar but weaker than  $V_S$ . The structure  
 376 terminates abruptly at the southern strand with little crustal structure in the Armutlu block. The  
 377 Adapazarı basin (centred at about  $40.7^\circ$  N) is representing as a low velocity anomaly between  
 378  $40.6^\circ$ N and  $40.7^\circ$ N to depths of about 6 km ( $V_S$ ).

379 Similar to the western profile we identify a slow anomaly in both  $V_P$  and  $V_S$  at depths of  
 380  $\sim 56$  km and  $\sim 52$  km, respectively. The  $V_S$  anomaly seems to show more complexity. We identify  
 381 a slow anomaly at  $\sim 76$  km depth in the Sakarya zone in  $V_P$  which cannot be identified in  $V_S$ . This  
 382 anomaly seems to terminate at the southern branch. Fast anomalies are detected at  $\sim 92$  km in  $V_P$   
 383 and  $V_S$  across the profile with shallower fast anomalies for  $V_P$  and  $V_S$  at  $\sim 76$  km depth beneath  
 384 the Istanbul zone and the Armutlu block. In  $V_P$  there is evidence of this interface splitting into a  
 385 deeper interface deepening to  $\sim 102$  km across the southern strand.

### 386 **6.3 Sakarya Zone**

387 The West-East profile for  $V_P$  and  $V_S$  (Figs. 7e and 7f, respectively) has been extracted along  
 388  $40.36^\circ$ N and is fully located within the Sakarya zone. The Sakarya zone is the southernmost  
 389 tectonic block in the study region. The inverted scattering tomography model shows a positive  
 390 anomaly at depths of  $\sim 38$  km. In  $V_P$  this interface moves to  $\sim 32$  km around  $30.6^\circ$ E. This anomaly  
 391 seems rather complex and might be discontinuous. We also identify a laterally limited fast anomaly  
 392 at  $\sim 30$  km between  $30.4^\circ$ E and  $30.6^\circ$ E. Deeper a slow anomaly at about 54 km depth can be seen  
 393 that shows a slight step down to about 60 km ( $V_P$ ) at about  $30.4^\circ$ E and seems complex in  $V_S$ . The  
 394 western part of the profile shows a fast anomaly at  $\sim 68$  km, with a slow anomaly at  $\sim 78$  km in the  
 395 east. A fast anomaly at  $\sim 98$  km depth (94 km in  $V_P$ ) is identified which terminates at  $30.2^\circ$ E in  
 396  $V_P$ .



**Figure 7.** As Fig.6 showing west-east oriented slices through the inverted  $V_P$  (left) and  $V_S$  (right) structure. Slices are located in the Istanbul zone (IZ - a,b) at longitude  $40.81^\circ$ , the Armutlu-Almacik block (AA - c,d) at  $40.58^\circ$  and Sakarya zone (SZ - e,f) at  $40.36^\circ$ . Areas with limited resolution as determined from the recovery tests (Fig. 5) masked in grey.

#### 397 6.4 Armutlu Block

398 In contrast to the Sakarya Zone, the Armutlu Block (Figs. 7c and 7d for  $V_P$  and  $V_S$ , respectively)  
 399 shows more structure down to depths of 40 km. A fast anomaly at  $\sim 40$  km terminates around  
 400  $30.4^\circ\text{E}$  and appears as shallow as 30 to 32 km further east in  $V_P$ .  $V_S$  also shows the termination  
 401 but a less pronounced step. The step around  $30.6^\circ\text{E}$  seems to coincide with the profile moving  
 402 from the Armutlu block to the Almacik mountains. West of  $\sim 30.4^\circ$ , this interface is underlain  
 403 by a slow anomaly at  $\sim 50$  km showing a step to  $\sim 58$  km at  $30.4^\circ\text{E}$  in  $V_P$ . Overall  $V_S$  seems  
 404 more complex. We identify several small scale fast and slow anomalies in the crust, the strongest

405 at  $\sim 14$  km around  $30.4^\circ\text{E}$  in  $V_S$ . Slow anomalies shallower than 40 km are indicated between  
406  $30.5^\circ\text{E}$  and  $30.9^\circ\text{E}$ .

407 A fast anomaly at  $\sim 94$  km stretches across most of the profile in  $V_P$ , with comparable but more  
408 complex structure in  $V_S$ . The  $V_S$  section also shows more localised structures at depths greater than  
409 80 km.

## 410 **6.5 Istanbul Zone**

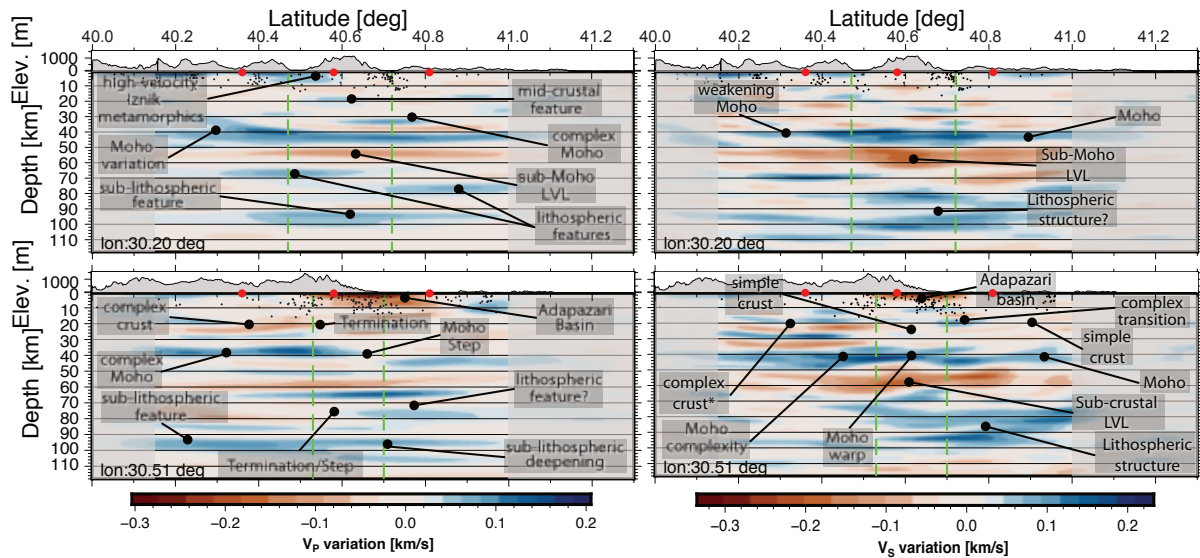
411 The Istanbul zone (Figs. 7a and 7b for  $V_P$  and  $V_S$ , respectively) shows very little structure down  
412 to depths of about 40 - 42 km where a strong fast anomaly can be detected in  $V_P$  and  $V_S$ . This  
413 fast anomaly seems to terminate around  $30.6^\circ\text{E}$  for  $V_P$  but remains continuous for  $V_S$ . A slow  
414 anomaly is visible in  $V_S$  at depths less than 10 km between  $30.6^\circ\text{E}$  and  $30.9^\circ\text{E}$  and a fast anomaly  
415 between  $30.2^\circ\text{E}$  and  $30.5^\circ\text{E}$ .

416 The strong Moho signal is underlain by a slow anomaly around 52 km depth again terminating  
417 around  $30.6^\circ\text{E}$  for  $V_P$ .  $V_P$  shows a fast anomaly at  $\sim 74$  km depth, which like the Moho signal in  
418 this block, terminates at about  $30.6^\circ\text{E}$ ; the corresponding structure in  $V_S$  is weaker and discontin-  
419 uous. A strong fast anomaly at  $\sim 92$  km depth can be seen in  $V_P$  and  $V_S$ , but again the  $V_S$  structure  
420 is complex.

## 421 **7 DISCUSSION**

422 The scattering tomography results show changes in the structure over distances of 5 km. These  
423 changes can be related to the different structure of the tectonic blocks and manifest in the North-  
424 South profiles. Nonetheless, we observe structural changes in East-West direction where more  
425 continuous structure might be expected, at least for the Istanbul and Sakarya zones. The profile  
426 across the Armutlu block follows the strike of the NAFZ east of about  $30.7^\circ\text{E}$  and moves from the  
427 Armutlu block to the Almacık Mountains just north of the NAFZ (Fig. 1). The depth slices through  
428 the model shown in Fig. 6 and 7 show strong changes between the two north-south trending profiles  
429 despite their close proximity. Interpreted NS sections are shown in Fig. 8. To highlight the most  
430 coherent part of the model we stack the depth profiles in longitude and divide these at  $30.4^\circ\text{E}$  to





**Figure 8.** Interpreted NS cross sections for  $V_P$  (left) and  $V_S$  (right). Shown are North-South oriented slices as in Fig. 6.

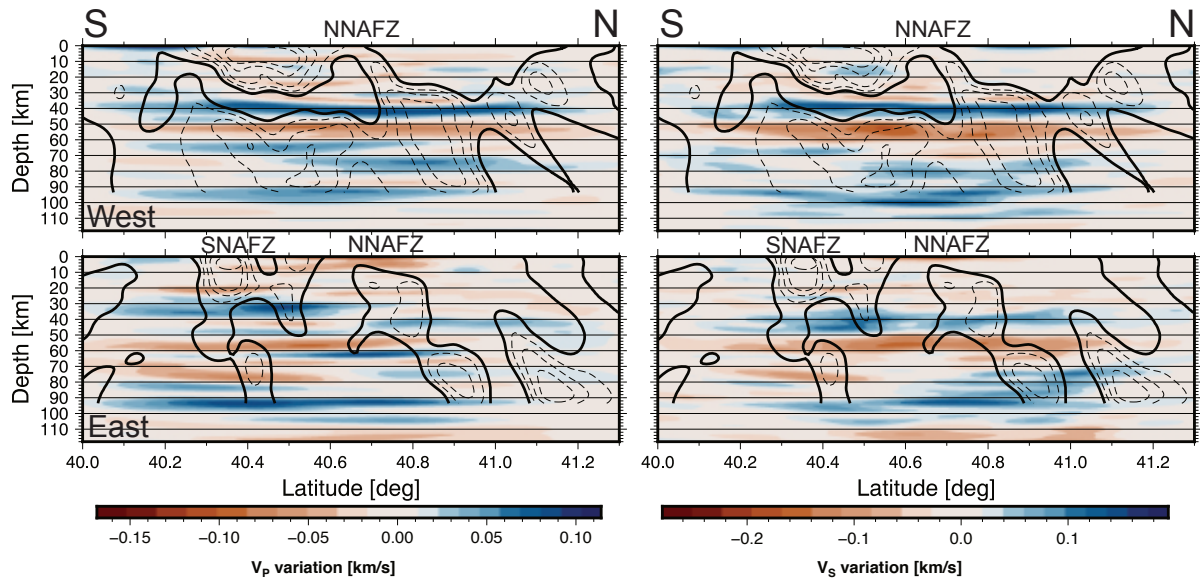
431 show western and eastern stacks in Fig. 9 for both  $V_P$  and  $V_S$ .

432 Comparing the individual slices and the stacked velocity-depth profiles shows that many features  
 433 are coherent along stretches of the profile, but can change on short scale-lengths.

### 434 7.1 Mohorovičić discontinuity

435 In the west the Mohorovičić discontinuity (the Moho) is visible in both  $V_P$  and  $V_S$  as a dominant  
 436 fast velocity at depths of  $\sim 40$  km with variations in  $V_P$  in the south and in  $V_S$  in the Armutlu  
 437 block. In the east the Moho is shallower in the south (34 km ( $V_P$ ), 38 km ( $V_S$ )) but shows a step  
 438 to larger depths (42 km) between  $40.6^\circ$  N to  $40.7^\circ$  N at  $30.51^\circ$ E. The deepening of the Moho  
 439 coincides with the low-velocity region in the teleseismic  $S$ -wave tomography model by Papaleo  
 440 et al. (2018) which has been interpreted as the location of the shear zone of the northern branch of  
 441 the North Anatolian Fault Zone (Fig. 9). The Moho at  $30.2^\circ$ E is overlain by a slow anomaly in  $V_P$   
 442 between  $40.7^\circ$ N and  $40.9^\circ$ N coinciding with the slow anomaly in the travelt ime tomography. The  
 443 corresponding anomaly in  $V_S$ , however extends further south to  $40.5^\circ$ N and coincides across that  
 444 range with both high and low velocity in the travelt ime tomography.

445 In the east the Moho seems much weaker and discontinuous across all three tectonic blocks.



**Figure 9.** North-South depth slices of the model averaged in longitude (a) for  $dV_P$  and (b)  $dV_S$ . Top figure for each velocity variation is for the western section (west of  $30.4^\circ$  E) of the model with the bottom the eastern part (east of  $30.4^\circ$  E). Contour lines show the  $V_S$  tomography of Papaleo et al. (2018) with the 0% contour line shown as solid and velocity reductions in 0.5% increments shown as dashed. Teleseismic traveltime tomography Papaleo et al. (2018) contours are shown at  $30.2^\circ$  E and  $30.4^\circ$  E

446 The strongest change in Moho depth can be identified around  $40.8^\circ$ N in the eastern profile where  
 447 we observe a step from 32 km to 40 km coinciding with the tomographically inferred edges of the  
 448 southern and northern branches of the NAFZ and the surface expression of the northern branch.  
 449 For the northern branch the fault zone seems to extend into the mantle.

450 While there are strong north-south changes in the profiles in Figs. 6 to 9 we also observe strong  
 451 east-west changes, e.g. in the Sakarya zone with a complex Moho structure around  $30.5^\circ$ E and the  
 452 weakening of the Moho east of  $30.4^\circ$ E or the change in the Armutlu block at  $\sim 30.5^\circ$ E, which  
 453 might be related to the step-over structure of the NAFZ related to the differential movement of the  
 454 Armutlu and Almacik blocks and the trend of the suture zones between the tectonic blocks. We  
 455 also observe a pronounced change of the Moho depth between the Armutlu block and the Almacik  
 456 mountains at around  $30.6^\circ$  E (Suppl. Fig. S14, S15), which is more pronounced in  $V_P$  but also  
 457 detectable in  $V_S$ , indicating strong contrasts in crustal structure between these two blocks.

## 458 7.2 Crustal structure

459 We find evidence for strong crustal structure variation along some of the profiles. The most striking  
460 structure is the apparent strong layering south of the southern branch in the Sakarya zone visible  
461 in the north south profiles. This structure is much stronger in the eastern profile for both  $V_P$  and  
462  $V_S$  but is dominant in  $V_S$ . The crustal heterogeneity is clearly limited by the surface location  
463 of the southern branch and forward models indicate that it consists of a series of high and low  
464 velocity anomalies (e.g. Fig. S7) perhaps related to emplacement of magmatic sills during the  
465 Tethys closure as have been seen in the eastern Marmara region (Karabulut et al., 2003). The crust  
466 in the Armutlu block on the other hand is relatively homogeneous and the heterogeneity is clearly  
467 truncated by the southern strand of the NAFZ.

468 Overlying the Moho in the area of the northern strand along the eastern profiles we detect  
469 small-scale, complex Moho structure. Modeling indicates that it could be related to a small-scale  
470 heterogeneity with limited extent approximating a point scatterer perhaps related to the material  
471 property changes in the fault zone. The lateral smoothing inherent to our inversions leads to a lack  
472 of resolution in this case.

473 We detect evidence for the Adapazari basin as low velocity anomalies between  $40.6^\circ\text{N}$  and  
474  $40.8^\circ\text{N}$  in the eastern profiles. Our method does not allow the necessary depth resolution at these  
475 depths for conclusion on the depth of the basin. The high velocity Iznik metamorphics (Taylor  
476 et al., 2019) can be detected between  $40.4^\circ\text{N}$  and  $40.6^\circ\text{N}$ .

477 Areas in the proximity of the surface expressions of the northern and southern strand show  
478 more heterogeneous structures than areas further away, perhaps related to increased damage around  
479 the fault zone (Ben-Zion & Sammis, 2003). However, the consistency between near surface ve-  
480 locity variation from our inversions and mapped surface structure provides confidence that we can  
481 interpret scattered signals from the fault zone as reflecting real structure. We detect a few localized  
482 crustal heterogeneities in the Sakarya zone and Armutlu block. There is evidence for a more con-  
483 tinuous low velocity anomaly at  $\sim 10$  km and  $\sim 25$  km depth in the Armutlu block and the Sakarya  
484 zone. These are clearer in the  $V_P$  models. The Istanbul Zone shows very little structure compared

485 to the neighbouring tectonic units, which could be related to the absence of metamorphism and the  
486 lack of major deformation (Okay, 1989).

### 487 **7.3 Sub-crustal structure**

488 Below the Moho we identify a dominant low velocity layer at depths between  $\sim 50$  and  $\sim 60$  km in  
489 both north-south profiles in  $V_P$  and  $V_S$ . The low velocity layer weakens but remains observable in  
490 the area of a prominent low velocity anomaly in recent tomographic models in this region (Papaleo  
491 et al., 2017, 2018) that is linked to the surface expression of the northern strand. This low velocity  
492 anomaly in the scattering tomography seems weaker in the Istanbul Zone. The weakening is more  
493 pronounced in  $V_P$  than  $V_S$ . The interface to the anomaly is slightly deeper (52 km) in the stacked  
494 eastern profile but also shows changes in the area of the reduced seismic velocities from traveltime  
495 tomography. The continuity of this structure beneath all tectonic blocks, although with possible  
496 depth and structural variations, indicates that it is related to lithospheric structure post-dating the  
497 amalgamation of northern Anatolia and the development of the suture zones. It is similar to a signal  
498 detected by Kahraman et al. (2015) and can be related to anisotropy beneath the Moho related to  
499 a lithospheric shear or decoupling zone active during continent amalgamation.

500 The fast anomalies detected at depths larger than 60 km shows changes in depth and struc-  
501 ture in the vicinity of the surface locations of the NAFZ branches although slightly offset to the  
502 north possibly linking mantle structure to dipping fault zones. The shear zone structure might  
503 penetrate deep into the lithospheric mantle showing a dip towards the north in agreement with  
504 previous results (Kahraman et al., 2015; Papaleo et al., 2017, 2018). The tectonic implications of  
505 this northerly dip remain unclear.

506 There is little evidence for a dominant low velocity anomaly at depths in our model that can be  
507 interpreted as a Lithosphere-Asthenosphere boundary (LAB). The lower part of the models seems  
508 dominated by high velocity anomalies, although there is weak evidence for a low velocity anomaly  
509 between 110 and 120 km depth at the boundary of our model. Therefore we cannot confirm a de-  
510 tection of the Lithosphere-Asthenosphere boundary in our models. Lower resolution deeper models

511 of the study region do not show evidence for a LAB in our images at greater depth. The LAB might  
512 be too gradational to show up as signal in the *P*-wave coda and to be imaged using our method.

#### 513 **7.4 Shear zones**

514 In the scattering tomography we see the strongest evidence for the NAFZ shear zone in the  
515 abrupt changes of crustal and sub-crustal structures. We see crustal structures that terminate on  
516 or near both fault strands, most clearly in the changes of the crustal structure transitioning from  
517 the Sakarya Zone to the Armutlu block (i.e. the SNAFZ) which we can trace to the Moho (e.g. Fig.  
518 6c,d). In general the AA shows almost no heterogeneity in the crust. Fig 6 a,b) shows evidence for  
519 heterogeneity at crustal depths coinciding with the surface expression of the NNAFZ.

520 The Moho step detected in the eastern profiles (e.g. Fig 6c, d)) seems to coincide with the  
521 surface expression of the NNAFZ and might indicate a localized shear zone extending deeper than  
522 the Moho and into the mantle. Some interfaces in the lithospheric mantle (e.g. Fig 6a,c) also show  
523 terminations coinciding with the NNAFZ indicating clear sub-Moho structure related to the shear  
524 zone.

525 Other continental continental transform faults such as the San Andreas Fault system (SAF), the  
526 Alpine Fault (AF) and the Dead Sea Transform (DST) show similar structures (e.g Weber et al.,  
527 2004; Mohsen et al., 2005; Ford et al., 2014; Stern & McBride, 1998) indicating localized shear  
528 throughout the crust as we find in this study. The SAF in southern California (Yan & Clayton,  
529 2007) and the DST along the Aravia fault (Mohsen et al., 2005) seem to offset the Moho in close  
530 proximity to the surface expression of the fault similar to the eastern profiles across the NAFZ.  
531 There is evidence that the SAF also offsets the LAB indicating localised shear even in the upper  
532 mantle (Ford et al., 2014). Although we do not detect clear arrivals from the LAB, mantle features  
533 down to depths of 90 to 100 km where the LAB has been detected (Kind et al., 2015) could indicate  
534 that the shear beneath the NAFZ is localized in the upper mantle similar to the SAF.

535 Due to the intra-Pontide suture zone located within the study area and the juxtaposition of  
536 the tectonic blocks described above it is difficult to separate the potential effects of a suture zone  
537 from the shear zone in this locality. The NAFZ seems to exploit a crust weakened by the presence

538 of sutures. Nonetheless, our results provide first evidence that the southern branch might extend  
539 throughout the crust with some evidence that mantle structures are also disrupted (e.g. termina-  
540 tions of high velocity anomalies at depths between 90 and 100 km). We also see evidence in the  
541 crust indicating small-scale heterogeneity coinciding with the location of the southern and north-  
542 ern strands. This could be related to scattering heterogeneity related to the shear zone along the  
543 northern branch.

## 544 **8 CONCLUSION**

545 We have used data from a dense deployment of seismometers to analyse the scattered seismic  
546 wavefield following teleseismic *P*-wave arrivals. Extending the analysis to a tomographic inver-  
547 sion we detect crustal and mantle heterogeneities that can be linked to the tectonics of the region  
548 around the North Anatolian Fault Zone. The high-resolution images from the scattering tomog-  
549 raphy down to depths of 120 km allow unprecedented insight into the small-scale structure of a  
550 major continental strike-slip fault. We show complex structure in crust and lithospheric mantle that  
551 can be linked to modern active tectonic processes as well as the structure of the crustal terranes  
552 that form the region.

553 Our images show that crustal structures contrast across the southern and northern strands of  
554 the NAFZ with stark differences between the Sakarya zone, the Armutlu block and the Istanbul  
555 zone. The terminations of crustal structure are sharp within the resolution of our approach. We  
556 detect a step in Moho depth coinciding with the surface location of the northern NAFZ strand and  
557 possibly changes in Moho structure in the region of the southern strand. We detect changes of  
558 mantle structure correlated to the fault zone indicating that the NAFZ extends into the mantle to  
559 depths of at least 75 km.

560 We show that scattering tomography in conjunction with dense recordings of the seismic wave-  
561 field is able to provide deeper insight into crustal and mantle structure and the fine-scale structure  
562 around fault zones. Strain localisation across the NAFZ seem to be narrow across the crust and into  
563 the mantle. This indicates that at least the upper part of the mantle lithosphere shows localized de-  
564 formation which also extends throughout the crustal layer. The NAFZ likely exploits weaknesses

565 due to old sutures in this region following the northwards subduction of the Tethys during the  
566 amalgamation of Anatolia.

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