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1 Evolution of normal fault displacement and length as the  
2 continental lithosphere stretches

3 Sophie Pan<sup>1\*</sup>, Rebecca E. Bell<sup>1</sup>, Christopher A-L. Jackson<sup>1</sup>, John Naliboff<sup>2</sup>

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5 <sup>1</sup>Basins Research Group (BRG), Earth Science and Engineering, Imperial College, Prince  
6 Consort Road, London, SW7 2BP

7 <sup>2</sup>Department of Earth and Environmental Science, New Mexico Institute of Mining and  
8 Technology, NM, USA

9

10 \*Corresponding author: Sophie Pan ([sophie.pan16@imperial.ac.uk](mailto:sophie.pan16@imperial.ac.uk))

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13 **HIGHLIGHTS**

- 14 • 7.2 Myrs after rift initiation, fault trace lengths were established to their near-full lengths
- 15 • At 28.5 Myrs, fault *segment* throw profiles were symmetric, but overall fault *systems* were  
16 asymmetric
- 17 • In the last 67% of rifting, antithetic faults located in stress shadows remained inactive and  
18 under-displaced
- 19 • We suggest faults grow in alternating phases of fault lengthening and displacement  
20 accumulation
- 21 • Our model suggests the frequently observed D-L scatter reflects fault growth and maturity

22

23 **ABSTRACT**

24 Continental rifting is accommodated by the development of normal fault arrays. Fault growth  
25 patterns control their related seismic hazards, as well as influencing the tectonostratigraphic  
26 evolution, resource extraction and CO<sub>2</sub> storage potential of rifts. Our understanding of fault  
27 evolution is largely derived by observing the final geometry and displacement (D)-length (L)  
28 characteristics of active and inactive fault systems, and by making subsequent inferences on their  
29 kinematics. We rarely consider how these properties change through time, and how the growth of  
30 individual fault *systems* relates to the temporal evolution of their host *arrays*. Here we use 3D  
31 seismic reflection and borehole data from the Exmouth Plateau, NW Shelf, Australia to  
32 determine the growth of rift-related, crustal-scale fault systems and arrays over geological  
33 timescales (>10<sup>6</sup> Ma). The excellent-quality seismic data allows us to reconstruct the entire  
34 Jurassic-to-Early Cretaceous fault array over a large area (~1200 km<sup>2</sup>). We find that fault trace  
35 lengths were established early, within the first ~7.2 Myr (8%) of rifting, and that along-strike  
36 migration of throw maxima towards the centre of individual fault systems occurred after ~28.5  
37 Myr (33%) of rifting. We propose that D and L may scale linearly, but increase via alternating  
38 phases of fault lengthening and displacement accumulation. Growth trajectories produce  
39 inflections in D-L space, reflecting times when fault lengths and/or displacement saturate a given  
40 rock volume, possibly controlled by crustal thickness. At the array-scale, faults located in stress  
41 shadows become inactive and appear under-displaced relative to adjacent larger faults, onto  
42 which strain localises as rifting proceeds. This implies that the scatter frequently observed in D-L  
43 plots can simply reflect fault growth and array maturity. We show that by studying complete rift-  
44 related normal arrays rather than individual faults, we can better understand how faults grow and  
45 more generally how continental lithosphere deforms as it stretches.

46        **1. INTRODUCTION**

47        The formation of extensional basins is controlled by the development of normal faults. Large,  
48 basin-bounding fault *systems* typically grow by the initiation, propagation, interaction, and  
49 linkage of smaller fault *segments* (e.g. Peacock and Sanderson, 1991; Anders and Schlische,  
50 1994; Trudgill and Cartwright, 1994; Dawers and Anders, 1995; McLeod et al., 2000). A  
51 kinematically linked network of fault segments and systems comprise a fault *array*. Fault arrays  
52 evolve in response to co- and interseismic stress feedbacks between the constituent segments and  
53 systems; this can lead to the temporal progression from numerous, short, low-displacement  
54 segments to a few, long, high-displacement systems ('strain localisation'; Cowie, 1998;  
55 Gawthorpe and Leeder, 2000) (Figure 1).

56        As strain localisation is often associated with an increase in basin subsidence rate, the  
57 temporal and spatial evolution of fault arrays are reflected in changes in rift physiography, and  
58 the location and rate of generation of accommodation (e.g. Dawers and Underhill, 2000, McLeod  
59 et al., 2000). The way in which fault arrays grow therefore strongly influences the size, shape,  
60 and distribution of sedimentary depocenters and drainage catchments, which in turn control the  
61 location of various energy resources, groundwater reservoirs, and potential CO<sub>2</sub> storage sites  
62 (e.g. Gawthorpe and Leeder, 2000).

63        Mechanical interactions between adjacent faults profoundly affects the stress-dependent  
64 nature of earthquakes in potentially hazardous, seismically active regions (e.g. Harris, 1998;  
65 Stein, 1999; Allmendinger et al., 2000; Dolan et al., 2007; Nicol et al., 2010; Mildon et al.,  
66 2019). An improved understanding of both long-term fault displacement patterns from ancient  
67 fault arrays imaged in seismic reflection data, together with short-term earthquake slip rates most  
68 commonly derived from regions of active extension, would allow us to characterise fault growth

69 and interactions over multiple timescales, from a few seconds up to tens-to-hundreds of millions  
70 of years (Allmendinger et al., 2000; Nicol et al., 2010; Nicol et al., 2006; Scholz and Gupta,  
71 2010; Nicol et al., 2019). However, our current understanding of fault growth largely stems from  
72 detailed studies of individual fault *systems* (e.g. Walsh et al., 2002; Nicol et al., 2005; Jackson  
73 and Rotevatn, 2013; Fossen and Rotevatn, 2016; Nicol et al., 2016; Childs et al., 2017; Jackson  
74 et al., 2017), rather than rift-scale fault *arrays*, thereby often overlooking the important role fault  
75 interactions and stress feedbacks may play in controlling the early stages of fault and rift  
76 development.

77 Historically, fault growth models have been proposed based on field and subsurface  
78 observations of the final geometry of ancient faults, often in the absence of kinematic constraints  
79 on how the faults evolved. The key geometric relationship that has been used is the relationship  
80 between fault length (L) and maximum fault displacement (D):

81

$$82 \quad D_{\max} = cL^n$$

83

84 Where  $c$  represents the scaling D/L ratio and  $n$  represents an exponent value argued to vary from  
85 2 (Watterson, 1986; Walsh and Watterson, 1988) to 1 (Cowie and Scholz, 1992; Dawers et al.,  
86 1993; Schlische et al., 1996), with a value of 1 indicating a linear scaling law (i.e. self-similarity)  
87 (Fig 2) (see also Marrett and Allmendinger, 1991). The exponent  $n$  remains unresolved due to  
88 the significant amount of scatter in global D-L profiles (Fig. 2) (Cowie and Scholz, 1992;  
89 Dawers et al., 1993; Scholz et al., 1993, Cartwright et al., 1995; Nicol et al., 1996; Cladouhos  
90 and Marrett, 1996; Schlische et al., 1996), which may reflect the natural geological variations  
91 between study areas (e.g. host rock mechanical properties), and uncertainties in data quality and

92 sampling approaches (Gillespie et al., 1992; Walsh and Watterson, 1988; Kim & Sanderson,  
93 2005; Walsh et al., 2017). However, it is possible that variations in  $n$  reflect fundamentally  
94 differing styles of fault growth. For example, current conceptual models suggest that faults may  
95 grow by either the: (i) ‘propagating model’, in which fault displacement and length  
96 simultaneously increase, producing a linear growth trajectory (Fig 2; e.g. Walsh and Watterson,  
97 1988; Cartwright, 1995; Walsh et al., 2003); or (ii) the ‘constant length model’, in which faults  
98 reach their near-full lengths rapidly before significant displacement accumulation, producing an  
99 initially shallow gradient trajectory on D-L plots as the fault lengthens, followed by a near  
100 vertical trajectory as the fault accumulates displacement (Fig 2; e.g. Walsh et al., 2002; Meyer et  
101 al., 2002; Nicol et al., 2005; Nicol et al., 2016; Childs et al., 2017; Jackson et al., 2017; Rotevatn  
102 et al., 2019). The two fault models may represent end-members (Fossen 2016; Childs et al.,  
103 2017), leading to the proposition of a hybrid growth model (Rotevatn et al., 2019), in which  
104 faults initially lengthen via propagation and linkage for 20-30% of the fault’s lifespan, before  
105 later accumulating displacement for the final 70-80%. However, the implications of the  
106 constant/hybrid model in D-L space are not yet clear.

107         Our current understanding of fault segment, system and array growth is largely geometric  
108 and based on observations from modern extensional basins or ancient basins imaged in seismic  
109 reflection data. There are very few studies that provide a *temporal* framework for the growth of  
110 normal fault segments and systems (e.g. Tvedt et al., 2016; Meyer et al., 2002). This is mainly  
111 due to the lack of age-constrained growth strata to record the timing of fault activity (e.g. Dawers  
112 et al., 1993; Schlische et al., 1996; Peacock and Sanderson, 1991). There are even fewer attempts  
113 to determine the evolution of entire fault arrays at the rift scale (e.g. Claringbould et al., 2017).  
114 We suspect this may reflect the time-consuming effort required to manually interpret many

115 seismically imaged normal faults, and to subsequently extract their related geometric properties  
116 (i.e. displacement, length). As a result, most studies focus on the growth of large, basin-bounding  
117 structures (i.e. major fault systems; e.g. McLeod et al., 2000), whilst ignoring smaller, intra-basin  
118 ones (i.e. fault segments; e.g. Tvedt et al., 2002).

119         Here we use 3D seismic reflection and borehole data from the Exmouth Plateau, NW  
120 Shelf, Australia to determine the growth of rift-related, crustal-scale fault systems and arrays  
121 over geological timescales ( $>10^6$  Ma). This study area was subject to Late Triassic-to-Early  
122 Cretaceous extension, forming extremely well-defined and well-imaged normal faults, half-  
123 grabens, and grabens (Exon and Willcox., 1978; Mutter and Larson, 1989). The large extent of  
124 the study area (1200 km<sup>2</sup>) is of a comparable size to areas of active continental extension  
125 containing large, seismogenic fault arrays, including East Africa (Baker et al., 1972; Albaric et  
126 al., 2010; Poggi et al., 2017), Central Italy (Luccio et al., 2010; Collettini and Barchi, 2002;  
127 Roberts and Michetti, 2004; Cowie et al., 2013), Central Greece (King et al., 1985; Bell et al.,  
128 2009; Taylor et al., 2011), SE Russia (Logatchev and Zorin, 1992; Mats, 1993), western Turkey  
129 (Seyitoglu, 1996), and the western USA (Hamilton, 1987; Thatcher et al., 1999). This makes our  
130 study area an excellent ancient analogue with which to better understand the geological evolution  
131 of areas of active continental extension.

132         To the best of our knowledge, this study is the first to quantify the temporal evolution of a  
133 complete rift-scale fault array using geometric observations of all seismically resolvable faults.  
134 We first define the present-day structure of the fault array (i.e. the final product of rifting).  
135 Second, we use depocentre mapping (via the use of time-thickness maps known as isochrons)  
136 and fault displacement backstripping to reconstruct the geometry of the entire fault array at an  
137 earlier stage of rifting. This allows us to determine the growth trajectory of individual fault

138 segments and systems, which we then compare to global D-L datasets (Figure 2). Based on our  
139 results, we propose a model for the growth of fault arrays. Our model links our current  
140 knowledge of the styles of segment- and system-scale growth (i.e. propagating vs constant length  
141 models) to the development of the larger fault array. Our model provides an explanation for the  
142 long-recognised scatter observed in global D-L datasets and sheds light on how continental  
143 lithosphere stretches during early rifting (e.g. Cowie and Schultz, 1992; Cartwright et al., 1995;  
144 Kim and Sanderson, 2005; Rotevatn et al., 2019).

145

## 146 **2. GEOLOGICAL HISTORY**

147 The Thebe dataset is located on the Exmouth Plateau (Exon and Willcox 1978) portion of the  
148 Northern Carnarvon Basin, NW Shelf of Australia (Geoscience Australia 2014) (Fig. 3a).  
149 Formation of the Northern Carnarvon Basin initiated during the Carboniferous to Permian in  
150 response to the breakup of Pangea, with episodic rifting continuing until the Late Cretaceous  
151 (Stagg and Colwell, 1994; Driscoll and Karner, 1998). Late Triassic rifting produced NE-  
152 trending sub-basins (Fig 3a), until a phase of regional uplift in the Callovian (Barber 1988; Pryer  
153 et al., 2002; McCormack and McClay, 2013; Metcalfe, 2013). Extension renewed in the Late  
154 Jurassic as the Argo block separated from Gondwana, culminating with seafloor spreading and  
155 the opening of the Argo Abyssal Plain during the Hauterivian (Tindale, 1998; Heine and Muller,  
156 2005; Gibbons et al., 2012; Metcalfe, 2013). The period of Late Jurassic-to-Early Cretaceous  
157 rifting was also characterised by the formation of NE-SW-striking faults, locally reactivating  
158 Late Triassic faults (e.g. Jitmahantakul and McClay, 2013; Magee et al., 2016).

159 The stratigraphic framework of the Exmouth Plateau consists of a thick, Upper Triassic pre-  
160 rift succession of the fluvio-deltaic and marginal marine deposits of the Mungaroo Formation



161 (Fig 4) (Tindale et al., 1998; Driscoll and Karner, 1988; Longley et al., 2002). From the Late  
162 Triassic to Late Jurassic, a syn-rift succession comprised of shallow marine claystone and  
163 limestone of the Brigadier Formation (Late Triassic) and marine claystones of the Athol and  
164 Murat formations (Early to Middle Jurassic), were deposited within fault-bound, rift-related  
165 depocentres (e.g. Stagg and Colwell, 1994; Tindale et al., 1998; Stagg et al., 2004; Longley et  
166 al., 2002). During the Late Jurassic to Early Cretaceous, uplift resulted in a decrease in water  
167 depth and an influx of interbedded shale and fluvial deltaic sands of the Barrow Group (Fig 4)  
168 (Boyd et al., 1992). In association with the subsequent subsidence, a major transgression resulted  
169 in the deposition of a thick succession of post-rift marine deposits of the Muderong Shale,  
170 Windalia Radiolarite and Gearle Siltstone, which are collectively known as the Winning Group  
171 (Hocking, 1987; Driscoll and Karner, 1998) (Fig. 4).

172

### 173 **3. DATASET AND METHODOLOGY**

#### 174 **3.1. Seismic reflection and well data**

175 Our database comprises a pre-stack time-migrated 3D seismic reflection survey (HEX07B)  
176 located in the WA-346-P permit area, and two wells (Thebe-1 and Thebe-2) (Fig. 3). The seismic  
177 data images to depths of ~4.5 s TWT (~3.7 km) and covers 45 km x 39 km, providing a total  
178 areal extent of 1200 km<sup>2</sup>. Crosslines are spaced at 25 m and inlines at 12.5 m. The dominant  
179 frequency is 45 Hz at the depth of the studied fault array at Thebe-1. From time-depth plots  
180 derived from well checkshot data, we estimate the average velocity to be 2480 m/s within the  
181 Jurassic syn-rift units; when combined with frequency data, this results in an estimated spatial  
182 vertical resolution of c. 14 m ( $\lambda/4$ ). Overall, the data quality is excellent throughout and enables a  
183 high-resolution analysis of fault structure and evolution.

184 The Thebe-1 and Thebe-2 wells are drilled in the footwalls of tilted fault blocks (Fig 3b) and  
185 collected a standard well-log suite (e.g. gamma ray, density, sonic, checkshot and  
186 biostratigraphic data). We choose to present our measurements of fault throw in time (ms TWT)  
187 rather than depth, given we are primarily interested in the relative changes of throw rather than  
188 absolute values (cf. Tvedt et al., 2016; Jackson et al., 2017), and we do not have a high-  
189 resolution 3D velocity model to constrain along-strike changes in velocity. To estimate depth in  
190 metres from observations in ms TWT, we have used checkshot data from wells that define a  
191 polynomial relationship between time and depth. We have not attempted to decompact our throw  
192 values as the loss of throw is likely relatively low (< 15%) in mixed sand-shale growth sequences  
193 that characterise much of the syn- to post-rift succession of the Exmouth Plateau (cf. Taylor et al.  
194 2008); we also lack information on sediment compaction parameters from the two wells. For  
195 better comparison with global D-L compilations, we converted our throw data to displacement  
196 data by assuming an average fault dip of 60°, as this is an average for the faults in this area.

197

### 198 **3.2. Structural framework**

199 We mapped ten horizons throughout the seismic reflection volume (labelled H1 – H10; see  
200 ages in Fig 4) to delineate the present basin structure and then unravel the fault evolution. Fault  
201 interpretation was conducted throughout the seismic volume, aided by overlaying the dip  
202 attribute onto mapped horizons (i.e. the deviation of a seismic reflection from a horizontal plane)  
203 (Chopra and Marfurt, 2007). We measured the tip-to-tip lengths of all seismically resolvable,  
204 individual fault segments at the base of the syn-rift, given this structural level captured most of  
205 the rift-related strain (see Fig 1 for definitions). In D-L plots, we show the scaling properties of  
206 individual segments, as well as fault systems (F1-F3). We classify faults systems as *major* (F1-

207 F3) if their throw is >150 ms (TWT) and *minor* if their throw is <150 ms TWT (labelled on Fig  
208 3b).

209 Fault throw was constrained by measuring the vertical distance between horizon cut-offs in  
210 the footwall and hanging wall on seismic profiles trending perpendicular to local fault strike (cf.  
211 McLeod et al., 2000; Tvedt et al., 2013). Values are taken every 100 - 400 m along the fault  
212 depending on the structural complexity and throw variability along the fault. In the presence of  
213 erosional footwall fault scarp degradation and fault-related ductile (continuous) deformation, the  
214 horizon cut-offs are projected towards the fault plane along the regional structural trend (e.g.  
215 Long and Imber, 2010; Whipp et al., 2014, Bell et al., 2014; Wilson et al., 2015; Barrett et al.  
216 submitted).

217

### 218 **3.3. Constraining rift evolution**

219 To constrain the fault array evolution, we use isochron (time-thickness) mapping and throw  
220 backstripping techniques, which are kinematically constrained by our mapped age-constrained  
221 horizons (H2, H3, H4, H8, H9 and H10; Fig 4). The ages of three syn-rift horizons are not  
222 constrained by the wells (H5, H6 and H7); their ages are therefore inferred based on an  
223 assumption of constant sedimentation rates between bounding, age-constrained horizons (Fig 4).  
224 We define H2 (Top Mungaroo; 208.5 Ma) as the base syn-rift (i.e. the start of rifting), consistent  
225 with other publications (e.g. Stagg et al., 2004; Tindale et al., 1998), given we observe no across-  
226 fault sediment thickness changes between H1 and H2 (Figure 5). Fault activity ceases at H9 (Top  
227 Muderong; 113 Ma), based on observations of no across-fault changes in thickness above H9,  
228 which we thus define as top syn-rift (i.e. the end of rifting). Our ages for top and base syn-rift  
229 give a total rift duration of ~85.5 Myr (Fig 4). We produced six isochron maps within the syn-rift

230 interval (SU1–6; Figs 4 and 6), which we use to constrain patterns of fault-controlled subsidence  
231 and thus fault growth (Thorsen, 1963; McLeod et al., 2000; Jackson and Rotevatn, 2013).

232 We use fault throw backstripping to constrain the style and patterns of fault growth (Childs et  
233 al., 1993; Rowan et al., 1998; Dutton and Trudgill, 2009, Jackson et al., 2017). Throw  
234 backstripping involves the subtraction of fault throw at multiple stratigraphic levels to determine  
235 the throw accumulated during deposition of the backstripped unit (see Chapman and Meneilly,  
236 1991). Faults can be backstripped using the ‘original method’ or the ‘modified method’. The  
237 ‘original method’ directly subtracts throw across a shallower horizon from that of the deeper  
238 horizon (Chapman and Meneilly, 1991; Petersen et al., 1992), making no assumptions about fault  
239 growth style. The ‘modified method’ subtracts the maximum throw along the horizon from the  
240 entire fault surface. This method implicitly assumes that the studied fault grew in accordance  
241 with the propagating fault model (Rowan et al., 1998; Jackson et al., 2017). In this study we use  
242 the original method to avoid any *a priori* assumption of the style of fault growth. We then  
243 compare our throw backstripping results with depocenter lengths revealed by isochrons; this  
244 allows us to assess fault length at different points in the rift history (cf. Jackson and Rotevatn,  
245 2013; Jackson et al., 2017).

246 The Exmouth Plateau was sediment-starved (i.e. underfilled) during Jurassic extension  
247 (Marshall and Long, 2013; Gartrell et al., 2016). Because of this, accommodation associated with  
248 the very earliest stage of extension may have been unfilled, meaning: (i) we cannot determine if  
249 any observed fault lengthening occurred by segment linkage or tip propagation (infilling  
250 sediments do not exist to record this); and (ii) our measured trace-lengths may thus have been  
251 established earlier (i.e. more quickly) than estimated, such that our estimates should be seen as  
252 maximum possible values (Jackson et al., 2017; Lathrop et al. submitted).

253

## 254 4. RESULTS

### 255 4.1. Structural framework

#### 256 4.1.1. *Time-structure maps*

257 We first describe the present-day structure of the study area (Fig 6 and Fig 7). We mapped  
258 150 faults over an area of 45 x 39 km that displace the 208.5 Myr base syn-rift horizon (H2, Top  
259 Mungaroo); this surface records the total cumulative rift-related strain of the fault array. The  
260 base syn-rift surface also contains the largest number of faults. Faults are strongly segmented,  
261 being either hard- and soft-linked across relays (Fig 6a). Fault trace lengths at the structural level  
262 of H2 (Top Mungaroo) range from c. 300 m to >15000 m, with the largest faults extending  
263 beyond the study area (Fig 6).

264 A fault map and rose diagram (Fig 7) reveal the strike and dip direction of the syn-rift fault  
265 array at the structural level of H2. Major NE-SW-striking faults define the fault array, with a  
266 minor set of E- dipping, N-S striking faults (Fig 7). The main NE-SW-striking population  
267 comprises 86 E- dipping faults and 57 W- dipping faults, although major faults (i.e. F1-F3)  
268 exclusively dip westwards, whereas the remaining antithetic faults dip eastward (Figure 7). The  
269 antithetic E- dipping faults are spaced ~2 km, whereas fault spacing for the major, W- dipping  
270 fault systems is 12-16 km (labelled on Fig 3b).

271 Time-structure maps reveal a reduction in faulting in the shallower structural levels as faults  
272 die-out upwards, with only the major fault systems (F1-F3) present at the top syn-rift horizon  
273 (H9; Top Muderong) (Fig 5, Fig 6f). These major fault systems extend upward above H10,  
274 where they link to normal faults contained in a widespread polygonal fault array (Fig 5). A series  
275 of small en-echelon faults occur at the structural level of H6 (~143 Myr) and H7 (~141 Myr),

276 whereas at deeper levels there is a continuous N-S-striking fault. This observation suggests that  
277 faults splay vertically upwards as well as horizontally (i.e. at lateral fault tips) (Fig 6c and d).  
278 Cross-cutting relationships suggest that the N-S-striking faults formed before the NE-SW-  
279 striking fault population (Fig 5, Fig 6a).

280

#### 281 **4.1.2. Throw distribution**

282 Maximum throw on the syn-rift fault array varies from 10 to 650 ms TWT (10 – 670 m)  
283 (Figure 8a). Throw varies along-strike on major fault systems F1-F3, with maximum throw (up  
284 to 650 ms TWT) typically occurring near their centres. Throw gently decreases to zero at the  
285 tips of fault systems (e.g. F2-a, Fig 8a). Individual segments forming the fault systems may show  
286 relatively steep throw gradients towards their tips (e.g. F2-b; see Fig 1 for definitions). Strain is  
287 distributed on several splay faults at the ends of fault systems (Figure 9a for location). Overall,  
288 the aggregate throw of major faults produce a broadly coherent, symmetric throw shaped profile  
289 with a gradual decrease in throw towards the tips (Figure 8a).

290 In detail, however, throw minima exist along many major fault systems (e.g. F2-a, F2-b, F3-  
291 b), with these minima tending to correlate with areas where fault systems change strike (F2, F3;  
292 Fig 8a). The relay zone between F3-a and F3-b shows a decrease in throw when fault segments  
293 overlap (Fig 8a). These two observations suggest major fault systems were previously  
294 segmented, at least at this structural level (Figure 8a).

295 In contrast to the major fault systems, minor fault systems typically have a uniform throw  
296 value along strike (ranging from 20-80 ms TWT) but exhibit an abrupt decrease in throw towards  
297 their tips (Figure 8a). Local throw minima on minor faults also correspond to abrupt changes in  
298 the strike, again suggesting a history of fault linkage via segment linkage (Figure 6a).

299

## 300 **4.2. Rift evolution**

301 In order to determine the growth of the fault array, we now quantify the activity of individual  
302 fault systems through the use of isochron maps and throw backstripping.

303

### 304 **4.2.1. Isochron analysis**

305 Within the earliest resolvable time period, SU1 (from 208.5 to 201.3 Ma; i.e. capturing fault  
306 activity during the first 7.2 Myrs of rifting), we observe across-fault thickening (~100 ms TWT)  
307 on all major faults (F1-F3; Fig 9a). These fault-controlled depocenters span the entire present-  
308 day fault trace length, indicating that major faults had already reached their near-final length  
309 during the first 7.2 Myr of the ~85.5 Myr rift history (i.e. only 8% of the total rift duration).

310 Thick (up to 200 ms TWT or 225 m), mound-shaped features are present in the hanging walls of  
311 F1, F3-a and F3-b, which we interpret as the deposition of fault scarp-derived talus (Fig 9a)  
312 (Bilal et al., 2019; Barrett et al., 2020). Fault scarp-derived deposits provide further evidence that  
313 the faults were active during this period. These deposits also suggest that fault slip rates outpaced  
314 sediment accumulation rates, and that at-surface relief formed which was then eroded.

315 Thickness variations in SU2 (Figure 9b) show that fault-controlled depocenters persisted  
316 adjacent to the major fault systems F1-F3 (from 201.3 to 163.5 Mya; i.e. capturing fault activity  
317 during the subsequent 37.5 Myrs of rifting). SU2 depocenters are wider (normal to fault strike)  
318 than those in SU1, and show a greater thickness increase up to 3.5 km away from the fault,  
319 indicating increased hanging wall flexure (Fig 9b). SU2 is thickest in the middle of the  
320 depocenters adjacent to F3-a and F3-b (<170 ms TWT) instead of immediately adjacent to the  
321 bounding faults, which we attribute to the underlying scarp-derived sedimentation in SU1

322 leading to an apparent reduction in the thickness of deposits in SU2 (Fig 9b, Fig 5). Although  
323 there is an overall thickness increase across both F3-a and F3-b (~120 ms TWT), high thickness  
324 variations (up to ~220 ms TWT) are distinct across F3-a and F3-b, separately indicating that the  
325 fault segments were soft-linked and separated by a relay zone (Fig 9b). Thickness variations also  
326 suggest that F2-a and F2-c segments were active during the SU2 time interval, but a lack of  
327 thickness changes in association with F2-b suggest that the F2 system was not hard-linked at this  
328 time (Figure 9b).

329         Within SU3 (from 180 to 143 Mya; i.e. capturing fault activity during the subsequent 37  
330 Myrs of rifting) we observe a thickness increase of 100 ms TWT across *all* segments of F2. This  
331 is in contrast to activity within SU2, where no thickness changes occurred across the faults  
332 central segment (F2-b), indicating that by 37 Myrs after the onset of rifting, F2 had hard-linked  
333 to form a through-going system (Fig 9c). Additionally, thickening across both F3-a and F3-b  
334 segments suggest that F3 was now acting as a single, hard-linked fault system that was  
335 principally growing via the accumulation of displacement at the expense of lengthening (Figure  
336 9c). In summary, F2 and F3 had grown to the maximum lengths after only 33% of the total rift  
337 duration, achieving this by segment linkage and relay breaching. This style of growth resulted in  
338 the development of abandoned splays in their hanging wall.

339         SU4 (from 143 to 141 Mya; i.e. capturing fault activity during the subsequent 2 Myrs of  
340 rifting) shows a significant thickness increase across F2, indicating this major structure was  
341 active at this time. SU5 (from 141 to 138 Mya; i.e. capturing fault activity during the subsequent  
342 3 Myrs of rifting) shows no thickness variations across F1 and F2, suggesting that these  
343 structures were inactive by this time (Fig 9e). F3-a also became inactive as evidenced by  
344 constant across-fault thickness of SU5. In contrast, the southern part of the fault system remained



345 active (F3-c and F3-d), defining a relatively wide (~4 km across strike) depocenter. SU6 (from  
346 138 to 123 Mya; i.e. capturing fault activity during the subsequent 15 Myrs of rifting) shows that  
347 nearly all the northern faults were inactive during this time, as only the southern segments of F2  
348 and F3 are associated with thickness variations. However, as these variations appear broad and  
349 regional, occurring on a length-scale larger than individual faults (Figure 9f), and given that the  
350 timing of deposition of SU6 coincides with a period of widespread regional subsidence along the  
351 margin (Hocking, 1987; Driscoll and Karner, 1998), we infer that F3-c, F3-d, F2-b, F2-c may be  
352 less active than the spatially related thickness changes suggest.

353

#### 354 **4.2.2. Fault throw backstripping**

355 Isochrons qualitatively reveal that fault lengths were established relatively quickly, as fault-  
356 controlled depocenters nearly span across their entire present-day fault trace lengths within the  
357 earliest resolvable slip increments (Fig 9a, 9b). This style of growth is consistent with the  
358 constant-length model of fault growth, hence the original rather than modified method is  
359 considered more appropriate to backstrip the fault array (e.g. Jackson et al. 2017; see  
360 methodology section). H3 is not present across the entire fault array due to footwall erosion (see  
361 Figure 7a), therefore backstripping is undertaken to the next resolvable increment H4 (180 Ma)  
362 to view an intermediate stage of array evolution (i.e. at 33% of the total rift duration). As a result,  
363 we can gain a view of rift-related strain after 33% (Fig 8b) and 100% of the rift duration (Fig 8a).  
364 Our results show that by 33% of rift history, major faults F1, F2, and F3 had accumulated  
365 maximum throws of throws of 261, 181 and 365 ms TWT (c. 261, 183, 372 m) respectively (Fig  
366 8b); this is approximately half of the present-day throw.

367        Throw accumulation was relatively uniform (averaging ~200 ms TWT) along F1, with a  
368 relatively abrupt decrease in throw occurring near a throw minima located at 16.3 km along the  
369 fault (Fig 8b). This minima corresponds to a small hanging wall splay and a subtle change in  
370 fault strike, thus we infer it defines a relay that was breached before this backstripped interval  
371 (i.e. before 33% of the rift duration; Fig 8b, Fig 3b). F1 continues north beyond the seismic  
372 dataset, so we cannot tell whether the maximum throw on this structure is located further north;  
373 however, extrapolation of the throw gradient at the northern end of F1 suggests this is not the  
374 case (Figure 8b). Our results indicate that during the latter part of rifting (33-100% of the rift  
375 duration), F1 accumulated throw near its centre without appreciable lengthening. Our present-  
376 day throw-length plot shows a broadly symmetric, bell-shaped profile, with no strong indication  
377 of paleo-segmentation (i.e. displacement minima; Fig 8a).

378        During the initial stage of rifting, F2 exhibited a similar throw distribution as F1, in that  
379 throw (average of 130 ms TWT) was fairly evenly distributed along its constituent segments (F2-  
380 a, F2-b and F2-d; Fig 8b). The throw minima between F2-a and F2-b suggests that these  
381 segments were not geometrically linked at this time. Instead, F2-b develops alongside two small  
382 hangingwall splays (pink and brown in Fig 8b). At this time, a small splay (purple in Fig. 8b; see  
383 Figure 3 for corresponding map) appears to be kinematically interacting with fault segments F2-c  
384 and F2-d, given they together define a throw profile that shows a gradual southward increase in  
385 throw (Figure 8b). However, the purple splay and F2-d appear to have been abandoned during  
386 the latter phase of rifting (33-100% of rift duration); instead, F2-c gained significant throw (from  
387 ~80 to 330 ms TWT) and hard-linked with F2 as rifting continues (Figure 8). During the latter  
388 phase of rifting, F2-b also accumulated a significant amount of displacement (up to 550 ms

389 TWT) to produce the relatively symmetric, bell-shaped throw profile that presently characterises  
390 F2 (Fig 8a).

391 Backstripping of F3 reveals that during the initial stage of rifting, a series of throw  
392 maxima were located towards the northern portions of segments F3-a, F3-b and F3-c. Whereas  
393 F3 still displays a generally symmetric throw profile, its constituent segments are defined by a  
394 rather irregular, asymmetric throw profile (Fig 8b). Throw minima (e.g. between F3-a and F3-b)  
395 indicate relay zones that was breached prior to this time-step (Fig 8b). By comparing the throw  
396 distribution after c. 33% of the rift history with that presently observed (Figure 8a), our results  
397 show that during the last 67% of rift history, the asymmetric and irregular profile of fault system  
398 F3 evolved to have a symmetric, bell-shaped profile, with the throw maxima migrating towards  
399 the present centre of the fault.

400 Backstripping of minor faults show that fault lengths and throw distribution remained  
401 relatively similar during the latter 67% of the rift history (Fig 8). A number of minor faults  
402 appeared to have died-out before the deposition of H4 (180 Ma; Top Murat Siltstone); however,  
403 ductile or so-called *continuous* deformation, make offset difficult to determine at H4. Our  
404 interpretation that the minor faults became inactive <28.5 Myr into the rift event is supported by  
405 isochron data, which shows that post-SU2 units do not thicken across them, as well as seismic  
406 cross-section evidence that show minor fault tipping-out below H4.

407

## 408 **5. DISCUSSION**

### 409 **Timescales of fault growth**

410 Our results clearly demonstrate that many of the fault's established their near-final length  
411 relatively early in the ca. 85.5 Myr rift history. Throw backstripping quantitatively shows that for

412 many of the faults, their lengths were essentially fixed after only ca. 33% of the total rift duration  
413 (Fig 8); isochron analysis suggest that fault length establishment may have occurred after as little  
414 as 8% of the total rift duration (Fig 9a). This duration is similar to that documented by Walsh et  
415 al., (2002), Jackson et al., (2017), and Rotevatn et al., (2019), who show that faults typically  
416 establish their near-final lengths within c. 30% of their total growth history. Our results show  
417 that this early stage of fault lengthening could be even shorter (i.e. 8%). Whereas current models  
418 of fault growth typically focus on the development of the largest fault *systems* within a given  
419 region, we note that some of the minor, antithetic faults located *between* the larger fault systems,  
420 are relatively long (i.e. >20 km) with low displacements (<30 m). Although we cannot directly  
421 establish if these structures were lengthening and/or accumulating displacement immediately  
422 prior to their death, observations of their long trace lengths, paleo-segmentation on time-structure  
423 maps (Fig 6), thickness variations from isochrons (Fig 9) and throw minima from throw  
424 backstripping (Fig 8), suggest that prior to becoming inactive, the majority of growth was  
425 occupied by fault lengthening via segment linkage. We suspect that these long, low-displacement  
426 faults may be ‘fossilised’ in their early stages, providing a snapshot of fault development before  
427 localisation and displacement accrual – we discuss this later in the context of our updated model.

428 The last c. 67% (and possibly up to c. 92%) of the 85.5 Myr rift history was characterised by  
429 strain localisation and displacement accumulation on the major fault systems (F1-F3). Strain  
430 localisation onto larger faults that are ultimately active for longer than smaller faults is also seen  
431 in the Inner Moray Firth, North Sea (Nicol et al., 1997; Walsh et al., 2003), and the Timor Sea,  
432 NW Shelf of Australia (Meyer et al., 2002), and in numerical (e.g. Cowie et al., 1995; Cowie,  
433 1998; Gupta et al., 1998; Naliboff et al., 2020) and physical models (e.g. Ackermann et al., 2002;  
434 Mansfield and Cartwright, 2001). Our results demonstrate that the pattern of displacement

435 accumulation and/or localisation itself is not straightforward, due to complex fault interactions  
436 that result in relay breaching and associated splay abandonment (e.g. Fig 9). Isochrons and throw  
437 backstripping reveal a distributed fault array that is initially composed of fault systems with  
438 multiple displacement maxima and minima. Detailed backstripping analyses show that in the  
439 early stages of fault development, fault segments exhibited profiles that were broadly symmetric  
440 (i.e. throw was greatest at the fault centres and decreases towards their tips). At the end of rifting  
441 (captured by present day geometry), throw profiles of fault systems are also symmetric.  
442 However, at an intermediate stage of rifting (33% of total rift duration; Figure 8b), while fault  
443 *segments* are symmetric, the overall fault *system* in which they are contained exhibited an  
444 irregular, asymmetric throw profile. At this stage of rifting, we observe multiple examples of  
445 splay abandonment (e.g. F2-c) related to strain localisation onto more optimally positioned fault  
446 segments that lie outside of stress shadows (e.g. F2-b). The along-strike migration of maximum  
447 slip at this stage may be important when trying to understand multi-fault ruptures that may occur  
448 on previously unrecognised faults, producing larger-than-expected earthquakes (e.g. Wesnousky,  
449 1986; Leeder et al., 1991; Goldsworthy and Jackson, 2000; Gupta and Scholz, 2000; Hamling et  
450 al., 2017). Our results suggest that interactions between faults can be better understood through  
451 changes in fault geometry as an array evolves; this may allow for better predictability of  
452 earthquake recurrences in seismic hazard analysis, where a somewhat random Poisson-based  
453 process may be used (e.g. Reiter, 1990; Main, 1996; Mildon et al., 2019; Sgambato et al., 2020).

454

#### 455 **Fault growth trajectory using throw-length relationships**

456 Figure 10 shows our new fault throw-length data for all 150 faults mapped in the Thebe  
457 fault array, relative to previously collected D-L data. Note that, in this plot, we have converted

458 our throw data in ms TWT to displacement in metres, using an average dip of  $60^\circ$  and borehole-  
459 derived time-depth relationships (see section 3.1). On a log-log scale, our data approximately  
460 follow a linear trend (i.e.  $n = 1$ ) across almost three orders of magnitude; this is similar to other  
461 global D-L scaling relationships for normal faults (e.g. Watterson, 1986; Walsh and Watterson,  
462 1988; Marrett and Allmendinger, 1991; Cowie and Scholz, 1992) (Fig 10). Our data are coloured  
463 by dip direction; blue-coloured faults are broadly east-dipping and red-coloured faults are  
464 broadly west-dipping (see Fig 7). We observe that, whereas all faults abide by a broadly similar  
465 linear scaling relationship ( $n = 1$ ), the major west-dipping (red-coloured) faults are presently  
466 somewhat ‘over-displaced’ ( $c \sim 0.1$ ) relative to the smaller, antithetic, east-dipping faults ( $c \sim$   
467  $0.01$ ) (Fig 10). Most of the over-displaced faults are abandoned splays, left behind in the hanging  
468 walls of fault systems after relay breaching and segment linkage. These faults exhibit very short  
469 trace lengths but high throws. In contrast to over-displacement, under-displaced faults may be  
470 indicative of interpreting faults as *systems* instead of *segments*. For example, one fault segment  
471 F3-a exhibits a displacement-length ratio of 0.06, but the entire fault system in which it is  
472 contained (i.e. F3) is 46 km long and has a throw of 656 ms TWT (c. 858 m in displacement; Fig  
473 8a). This gives a lower D-L ratio of 0.02 (Fig 10), demonstrating that D-L (and T-L) ratios are  
474 higher for individual fault segments than systems, although both still plot within the scatter  
475 observed in global range D-L data (Fig 10). Similar observations are noted by Peacock and  
476 Sanderson (1991) and Roberts and Michetti (2004), highlighting how mixing the analysis of fault  
477 segments (e.g. Meyer et al., 2002; Roberts et al., 2004) and/or fault systems (e.g. McLeod et al.,  
478 2000) may explain the range of scatter already seen in D-L profiles.

479         Fault length and displacement measurements for the present-day fault array (Fig 8a),  
480 coupled with the reconstructed fault array at c. 33% of the total rift duration (Fig 8b), allow us to

481 define how D-L relationships have changed over the last c. 67% (57 Myr) of rift history. Our  
482 time-constrained displacement-length relationships reveal either: i) a vertical growth line,  
483 indicating that the fault accumulated displacement without lengthening (i.e. the constant fault  
484 growth model; Walsh et al., 2003; Nicol et al., 2005; Jackson and Rotevatn, 2013), or; ii) a static  
485 point, indicating that the faults was inactive during the last c. 67% of rifting (Figure 10 inset).  
486 Typically, faults that were inactive were the east-dipping antithetic faults that did not slip post-  
487 180 Ma and thus remained ‘under-displaced’ (Fig 10). Because all the major and minor fault  
488 systems dip westward and eastward, respectively, the structures onto which strain was eventually  
489 localised must have been set early in rift history (i.e. as early as the first c. 8% of the total rift  
490 duration). This may be attributed to positive stress feedback processes occurring during early  
491 array development, whereby shear stress reduction or ‘stress shadow zones’ inhibit fault growth  
492 and the nucleation of new cracks (e.g. Hu and Evans, 1989; Ackermann and Schlische, 1997;  
493 Cowie, 1998), and optimally positioned faults form and slip (i.e. faults with across-strike spacing  
494 of 12-14 km). Early establishment of the near-final fault lengths reinforces the suggestion that  
495 lateral tips of faults are effectively pinned by stress interactions with neighbouring faults (e.g.  
496 Burgmann et al., 1994; Willemsse et al., 1996; Gupta and Scholz, 2000; Contreras, 2000). We  
497 propose that distributed faulting, principally characterized by fault lengthening, occurs until a  
498 given rock volume is ‘strain saturated’, as predicted by numerical and physical models (e.g.  
499 Cowie and Scholz 1992b; Cowie et al., 1995; Wu and Pollard, 1995; Ackermann et al., 2001).

500 Our observation that minor (E- dipping) faults either die early (where blue coloured faults  
501 are static points; Fig 10), or grow and still appear relatively under-displaced (where blue faults  
502 have a vertical trajectory that is relatively under-displaced; Fig 10) relative to larger faults onto  
503 which strain is subsequently localised (with higher displacement of red faults; Fig 10), has only

504 been revealed from studying a complete fault array. Considering our data spans nearly three  
505 orders of magnitude, and given that the 150 faults formed in broadly similar rock types in the  
506 same stress regime, this implies that some of the scatter seen in D-L plots may simply be related  
507 to fault maturity and not only from limitations inherent to data collection (e.g. Kim and  
508 Sanderson, 2005) and/or differences in rock type and stress regime (e.g. Cowie and Scholz,  
509 1992; Gillespie et al., 1992; Scholz et al., 1993).

510 Our findings illustrate the importance of collecting data from entire fault arrays  
511 developed over large areas if we hope to determine how normal faults grow. Datasets that allow  
512 inspection of only a small part of the rift (e.g. small-scale field-based studies), or analysis of  
513 faults lacking growth strata, could produce D-L relationships that appear anomalous (i.e.  
514 relatively under- or over-displaced) simply due to fault maturity (i.e. the datasets sample only the  
515 very young or old faults). Taking the trendline of datasets may produce an exponent value  $n$  that  
516 may not be representative of geometrical fault scaling relationships. We suspect that this may be  
517 why the true value of  $n$  is contested (e.g. Watterson, 1988; Marrett and Allmendinger, 1991;  
518 Gillespie et al., 1992; Cowie and Scholz, 1992; Clark and Cox, 1995; Dawers and Anders, 1995;  
519 Soliva et al., 2005; Xu et al., 2006; Torabi et al., 2011). By studying complete fault arrays rather  
520 than individual fault systems, we can better understand how faults grow and how continental  
521 extension proceeds.

522

### 523 **A model of strain accommodation as the continental lithosphere stretches: incorporating** 524 **segment- and system-scale growth**

525 Although our study demonstrates that during the last c. 67% of rifting faults either grew  
526 by displacement accumulation or became inactive, the magnitude of displacement accumulation



527 (i.e. the vertical trajectory on D-L profiles) does not extend outside of the overall scatter  
528 observed in global D-L plots (where  $c \sim 0.001$ , Figure 10). This is important, given that recently  
529 proposed fault models (constant-length and hybrid) commonly present growth trajectories only  
530 very schematically on D-L plots (e.g. Childs et al., 2017; Rotevatn et al., 2019; Fig 2). Due to the  
531 lack of scale on each axes, the constant-length and hybrid models also lack geometric constraint,  
532 which D-L profiles provide (i.e. faults cannot, for example, exhibit trace lengths of 10 km with 1  
533 m of displacement, as this would plot well below the lower trendline of D-L profiles). However,  
534 current literature has frequently superimposed the schematic model trajectories onto global D-L  
535 data, in turn implying that faults that grow (or grew) via the constant-length and hybrid model  
536 apparently lie well outside of the known D-L scaling relationships. We suspect that many  
537 observations from ancient (inactive) and modern (active) extensional settings that endorse the  
538 constant-length model (e.g. Meyer et al., 2002; Tvedt et al., 2016; Jackson et al., 2017; Rotevatn  
539 et al., 2019) in fact still lie within the bounds of our proposed geometrical constraint, as our  
540 results from offshore Australia show (Figure 10).

541 We propose a new model for how fault systems grow within an array, where faults  
542 ultimately scale by a linear relationship in log-log space ( $n = 1$ ), but grow by alternating periods  
543 of lengthening and localisation, thereby producing a more step-like growth trajectory in D-L  
544 space than previously proposed (Fig 11). First, faults initiate as numerous, isolated segments  
545 (Stage 1 in Fig 11; e.g. Gawthorpe and Leeder, 1993). Fault lengthening via segment linkage  
546 occurs until a small rock volume (e.g. a mechanical layer of rock, which could at the largest scale  
547 represent the seismogenic crust) is saturated with faults with relatively little displacement. These  
548 early faults appear ‘under-displaced’ when viewed on D-L plots (2; Fig 11). Once the rock  
549 volume is sufficiently saturated with fault trace lengths defining the overall pattern of the fault

550 network, optimally positioned faults (i.e. those that receive positive stress feedback from  
551 adjacent structures) then accrue displacement, whereas faults in stress shadow zones remain  
552 inactive and ultimately die (3; Fig 11). This phase of fault localisation and displacement  
553 accumulation produces a vertical growth trajectory on D-L plots, where active faults move  
554 vertically up until the upper scaling limit (approximately where  $c = 0.1$ ). Fault that are not  
555 optimally positioned become inactive faults and die along the vertical trajectory, leaving a  
556 vertical spread of points. At this point, stress feedback mechanisms cause the fault array to be  
557 fully mechanically interacting, essentially inhibiting further fault interaction and growth. To  
558 further accommodate strain, the considered rock volume and the overall scale *increases* (from  
559 stage 3 to 4; fig 11). During this time the largest fault systems can mechanically interact (e.g.  
560 Ackermann et al 2001), acting again as isolated segments during the ‘initiation’ stage (4; Fig 11).  
561 These faults then lengthen, predominantly by segment linkage (5; Fig 11), and then strain  
562 localises onto optimally spaced zones, producing more widely spaced major faults (i.e. red fault  
563 segment; Fig 11). Faults located in stress shadow zones become inactive, appearing relatively  
564 under-displaced in D-L space (i.e. the green segment; Fig 11).

565 Our model includes aspects of the isolated and constant-length fault models. This is  
566 important because in the absence of data that truly constrains the temporal evolution of fault  
567 arrays and their constituent segments and systems, faults may appear to follow a sloping, linear  
568 trajectory, thereby endorsing the isolated model. The hybrid and constant-length models are also  
569 supported given that the final increment of fault growth, which in our study accounts for at least  
570 c. 67% of the rift history, is defined solely by displacement accumulation. Our fault growth  
571 model builds on this premise, however, by considering that faults systems interact as part of an  
572 array and are not kinematically or mechanically isolated (Fig 11).

573 Our suggested model inevitably leads to a lot of scatter in D-L profiles from the onset of  
574 extension, with this outcome indicating that the natural variability in D-L relationships may be  
575 an even bigger consequence of fault growth than previously suggested (e.g. Peacock and  
576 Sanderson 1991; Cartwright et al., 1995; Peacock and Sanderson, 1996). We suggest that the  
577 ‘inflection point’ between fault lengthening and displacement accumulation, bounded at D-L  
578 scaling trendlines (c, Fig 11), is controlled by the physical (e.g. mechanical strength, effective  
579 elastic thickness) characteristics of the region undergoing strain and, more specially, relates to  
580 layer thickness (cf. Ackermann et al., 2002) and areal extent, as this determines a ‘saturation’  
581 point of fault lengths and displacements. Studies proposing a linear correlation between layer  
582 thickness and fault spacing are consistent with our updated model of fault growth, as this this  
583 relationship is largely driven by stress feedback mechanisms (Soliva and Benedicto, 2005; Soliva  
584 et al., 2006). Due to many natural mechanical heterogeneities in the subsurface, the ‘inflection’  
585 transition may occur at any point in scale and account for the more subtle changes in growth  
586 trajectory path variability that may be lost in log-log space (Rotevatn et al., 2019). Further  
587 decrypting the ‘natural’ variability of global D-L plots by unravelling fault growth trajectories  
588 may allow us to better predict certain structural and mechanical properties, such as crustal  
589 thickness or mechanically confined layers in data-poor regions.

590 From our model, we are able to characterise the scaling properties and consider the  
591 growth stage of active and ancient, natural fault arrays. Fault populations that sit near the lower  
592 D-L boundary, where  $c = 0.01$ , exhibit distributed faulting, are seen in the Central Afar and on  
593 the Asal Ghoubbet faults in East Africa ( $c \sim 0.012$  in Manghietti et al., 2015). In contrast, fault  
594 populations that lie closer to the upper D/L boundary ( $c = 0.1$ ) may be characterised by more  
595 localised faulting, similar to that seen in fault populations imaged in seismic reflection data (e.g.

596 northern North Sea; McLeod et al., 2000, NW Shelf of Australia; Black et al., 2017). 3D seismic  
597 reflection surveys will often image now-inactive, ancient rifts that have proceeded through a  
598 phase of strain localisation. Therefore, typical fault system-scale studies may only capture the  
599 most mature faults that have accrued the largest amounts of displacement, potentially leading to  
600 an erroneous endorsement of the constant-length model. Conversely, studies from active rifts  
601 may be biased toward fault populations with lower  $c$  values, as immature faults in their early  
602 phase of development may be more readily captured by these datasets.

603         Resolving whether length establishment occurs within, for example, a few tens of  
604 thousands to a few million years is difficult with seismic reflection data from ancient rifts, given  
605 that the limited spatial and temporal resolution of seismic data at depth; it may thus be difficult  
606 to learn more about the fault lengthening stage from solely the kinematic analysis of seismic  
607 reflection data that image ancient basins. We suggest that future studies should focus on active,  
608 over-filled, shallowly buried (and thus seismically well-imaged) rifts (e.g. the Whakatane rift,  
609 New Zealand; Taylor et al., 2004), but 3D seismic data rather than grids of 2D data will really be  
610 required. Additionally, numerical and physical models may be able to test our hypothesis by  
611 constraining fault growth through global scaling relationships to investigate the inflection point  
612 of fault growth trajectories observed in D-L space.

613

## 614 **6. CONCLUSIONS**

- 615         • The timescales over which a fault array evolves have previously not been well  
616         constrained, due to a lack of dynamic, kinematic data.

- 617 ● To fill this gap, we studied an extensive fault array (1200 km<sup>2</sup>) in the Exmouth Plateau,  
618 NW Australia and characterised 150 seismically resolvable faults using displacement-  
619 length relationships.
- 620 ● Isochron maps reveal that within the earliest resolvable slip increment of 7.2 Myrs (out of  
621 an overall duration of 85.5 Myrs) the majority of fault lengths have already been  
622 established.
- 623 ● Fault displacement backstripping to 28.5 Myrs (33% of total rift duration) shows that the  
624 fault array was characterised with distributed faulting. Backstripping reveals along-strike  
625 slip migration through time to produce symmetric fault systems. Throw minima suggest  
626 previous segmentation and relay breaching.
- 627 ● Our results allow us to view the last growth trajectory (representing the last 67% of total  
628 rift duration) on D-L plots. Faults undergo varying amounts of displacement accrual (i.e.  
629 the constant length model). Less (to no) displacement is accrued on faults in stress  
630 shadow zones.
- 631 ● We propose that fault length and displacement scale linearly ( $n = 1$ ) but grow via  
632 alternating phases of lengthening and displacement accrual. Our model is largely based  
633 on stress feedback mechanisms driven by layer/crustal thickness, which produce  
634 ‘inflection’ points in D-L space.
- 635 ● Inactive faults located in stress shadow zones contribute towards the large amount of  
636 global displacement-length scatter due to differences in fault maturity. As the scatter  
637 may be inherent to fault growth, D-L relationships may help us to differentiate fault  
638 characteristics in ancient and modern rifts.

639

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648 **FIGURE CAPTIONS**

649 **Figure 1.** Fault array evolution, modified after Gawthorpe and Leeder, (2000). Faults grow from  
650 an initiation stage (A) with isolated, short (20-10 km) and low-displacement fault segments into  
651 a period of interaction and linkage (B), and finally strain localises onto a few long (>10 km),  
652 high displacement ‘through-going’ fault zones (C). Typically, studies investigating fault growth  
653 measure the geometry of fault systems when they have reached (C).

654 **Figure 2.** Global displacement-length (D-L) data of normal faults. D-L relationships are  
655 presented in log-log space, therefore the line of best fit results in the scaling exponent,  $n$ . Faults  
656 may grow via the isolated where increases in displacement and length occur sympathetically,  
657 producing the green growth trajectory; or the constant-length model where the final length is  
658 established early in slip history, prior to accumulation of significant displacement, producing the  
659 blue trajectory.

660 **Figure 3.** The location of our study area: (a) Location of the Thebe field seismic dataset within  
661 the Exmouth Plateau, North Carnarvon Basin, North West Shelf of Australia. Sub-basin  
662 boundaries are taken from Geoscience Australia. The main structural NE- trends in Fig 3b are

663 from Bilal et al. 2019; and (b) The main structural elements of the Thebe dataset. Major fault  
664 systems are labelled in colour, correlating to their respective throw-profiles in Figure 8. Minor  
665 faults are coloured in light grey. Wells Thebe-1 and Thebe-2 are shown.

666 **Figure 4.** Horizons and their respective formation ages. H1 and H10 are pre-rift and post-rift  
667 reflectors, whereas H2 – H9 define the syn-kinematic unit of the study area. H4, H5 and H6 do  
668 not contain biostratigraphic ages, therefore age is inferred through the assumption of constant  
669 sedimentation. Our different methodologies to understand rift evolution are resolved to different  
670 stages in time, which has been demonstrated on the right.

671 **Figure 5.** Cross section of inline 1524 (location shown on Figure 3b): (a) Without interpretation;  
672 and (b) With interpretation and annotation highlighting the main structural features of the overall  
673 study area.

674 **Figure 6.** Time-surface maps (ms TWT) showing the present day structure of the dataset. Faults  
675 at each successive interval are mapped using a variance attribute overlay for higher confidence.

676 **Figure 7.** Dip direction of the fault array. The rose diagram shows the average strike azimuth  
677 and dip direction per fault segment. The rose diagram is split into two halves, whereby the blue  
678 coloured faults are east dipping, and the red coloured faults are west dipping.

679 **Figure 8.** Throw length profiles (segments coloured by Fig 3b) and their map view throw  
680 distribution, for (a) The present day fault array, i.e. the final product of rift history; and (b) The  
681 reconstructed, backstripped fault array, where faults are backstripped to 28.5 Myrs (H4; Top  
682 Murat).

683 **Figure 9.** Isochron maps showing thickness changes between syn-depositional horizons. The  
684 syn-rift unit is divided into six, labelled as SU1-SU6.

685 **Figure 10.** Our results plotted against global displacement-length data in grey. Our data is  
686 coloured by the dip map (see Figure 7) whereby blue-coloured faults are east dipping, and red-  
687 coloured faults are red dipping. Our results were measured in throw, but we have converted our  
688 values into displacement by assuming an average dip of 60 (see supplementary table for raw  
689 values). The inset shows the last 67% of rift history, as we connect our present-day values of  
690 fault geometry with their respective backstripped values.

691 **Figure 11.** Our new proposed fault growth model, where the fault trajectory shows growth via  
692 alternating phases of fault lengthening (producing a near horizontal trajectory) and fault  
693 displacement (producing a near vertical trajectory). Stage 1, 2 and 3 are examples of initiation,  
694 linkage and localisation over a smaller region. Faults initiate as small numerous fault segments  
695 (1), then lengthen via segment linkage until the rock volume is sufficiently saturated, producing  
696 relatively under-displaced D-L values (2), displacement accrues on optimally located and spaced  
697 faults (3). Once the smaller region is saturated by fault lengths and displacement, the considered  
698 region increases, allowing the fault array to undergo initiation (4), lengthening (5) and  
699 localisation (6) again at a larger scale.

700

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## 702 **DATA AVAILABILITY STATEMENT**

703 The data that support the findings of this study are available from the corresponding author upon  
704 reasonable request.

705

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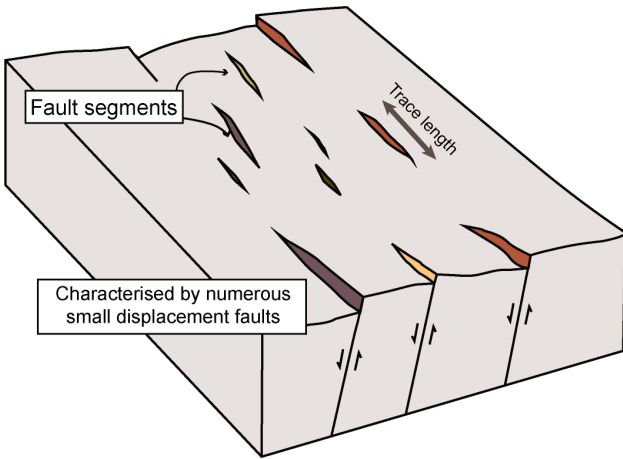
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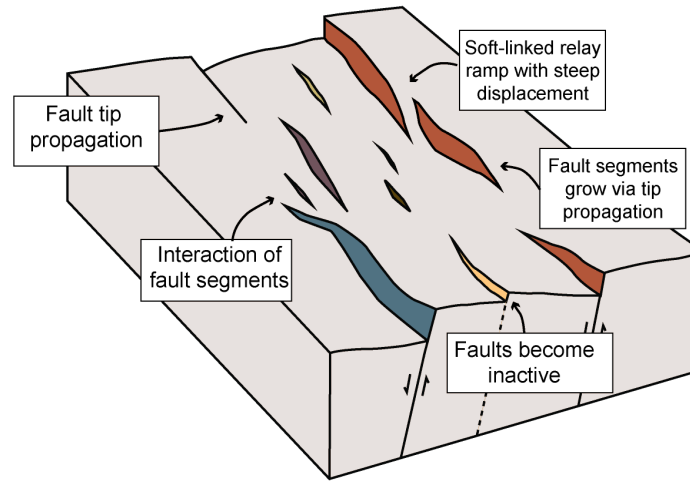
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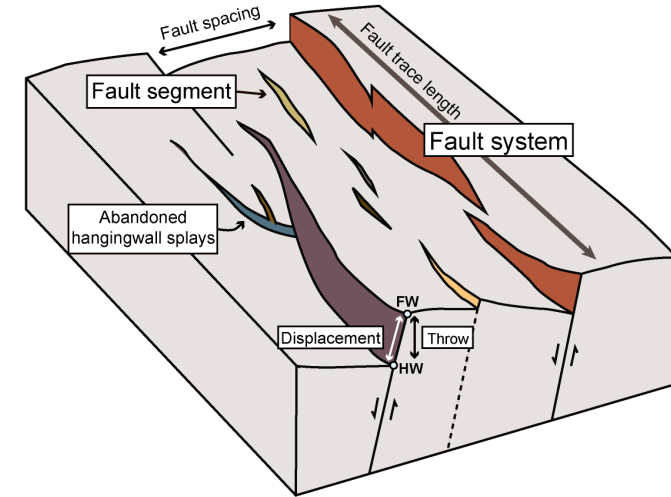
### A) Initiation



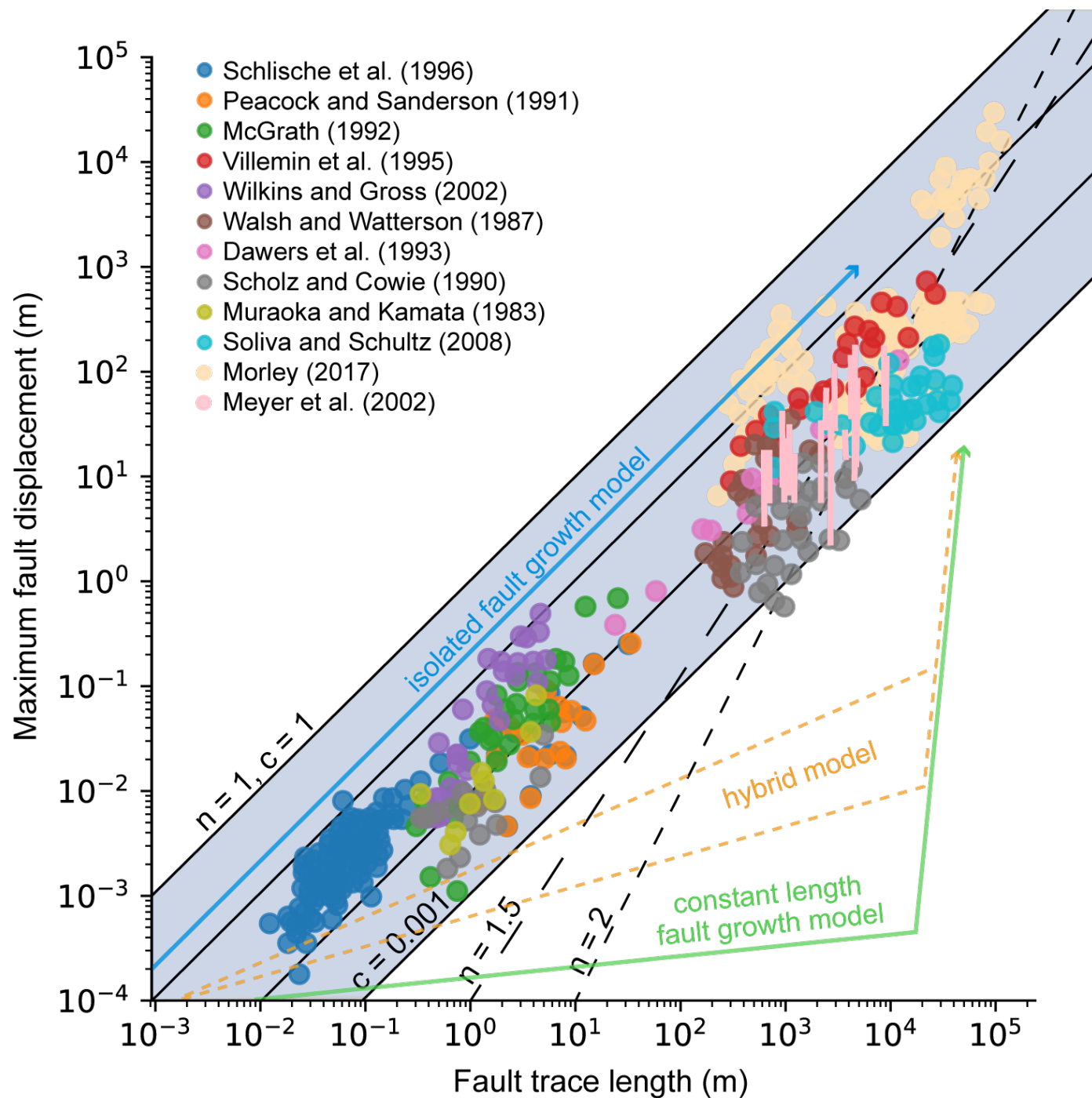
### B) Interaction and linkage



### C) Through-going fault zones

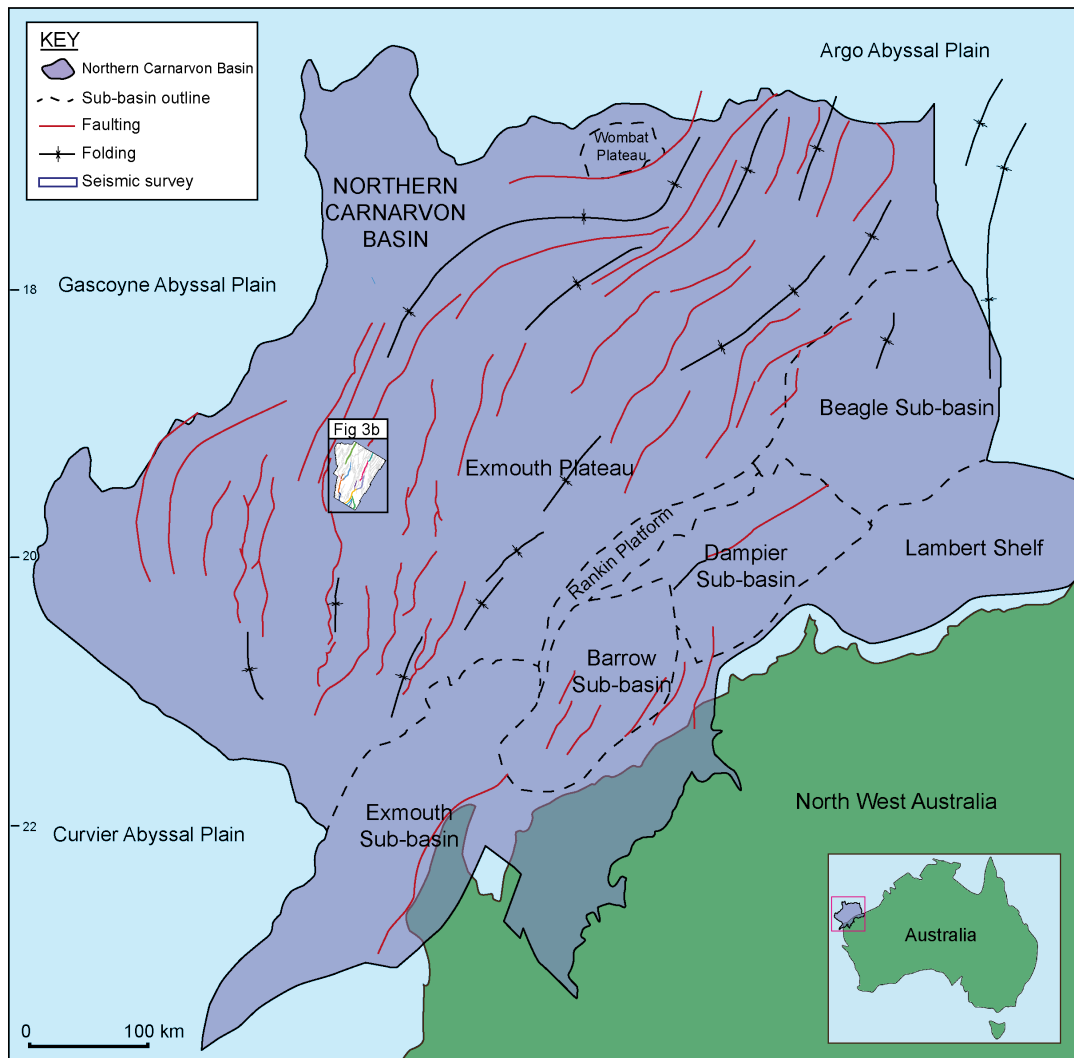


**Figure 1.** Fault array evolution, modified after Gawthorpe and Leeder (2000). Faults grow from an initiation stage (A) with isolated, short (20-10 km) and low-displacement fault segments into a period of interaction and linkage (B), and finally strain localises onto a few long (>10 km), high displacement 'through-going' fault zones (C). Typically, studies investigating fault growth measure the geometry of fault systems when they have reached (C).

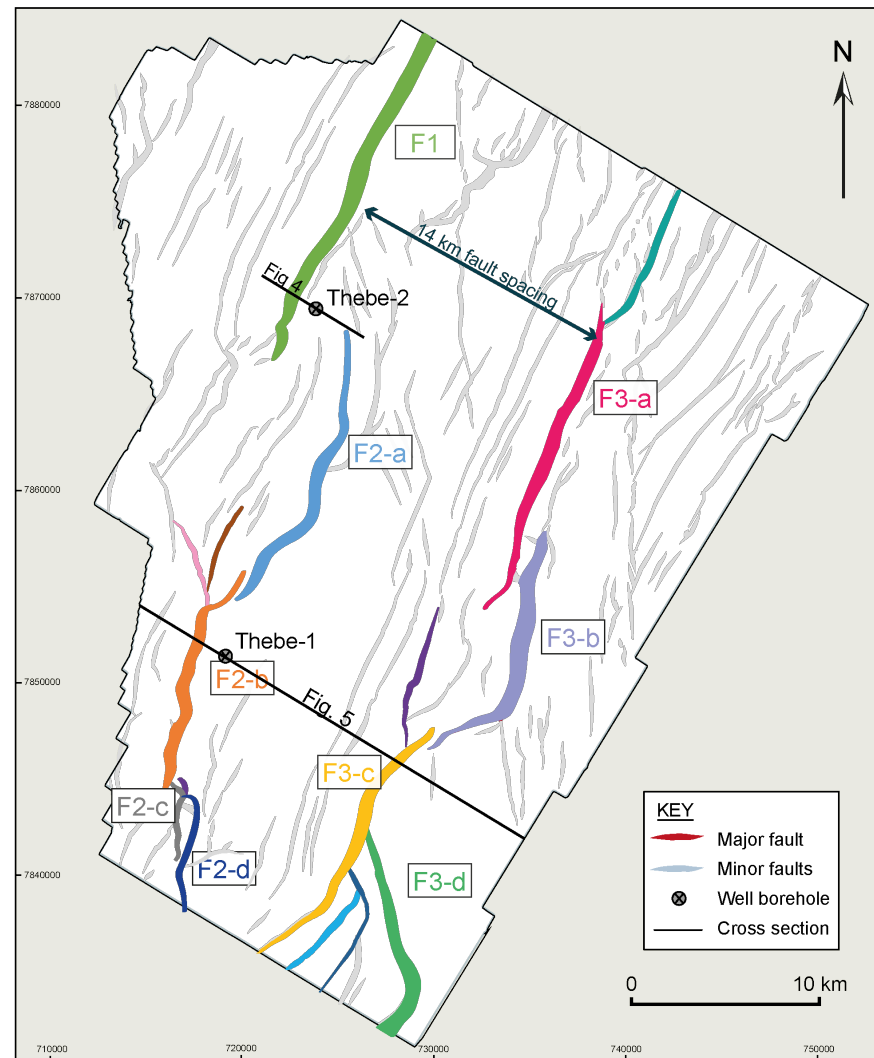


**Figure 2.** Global displacement-length (D-L) data of normal faults. D-L relationships are presented in log-log space, therefore the line of best fit results in the scaling exponent,  $n$ . Faults may grow via the isolated model where increases in displacement and length occur sympathetically, producing the blue growth trajectory; or the constant-length model where the final length is established early in slip history, prior to accumulation of significant displacement, producing the green trajectory.

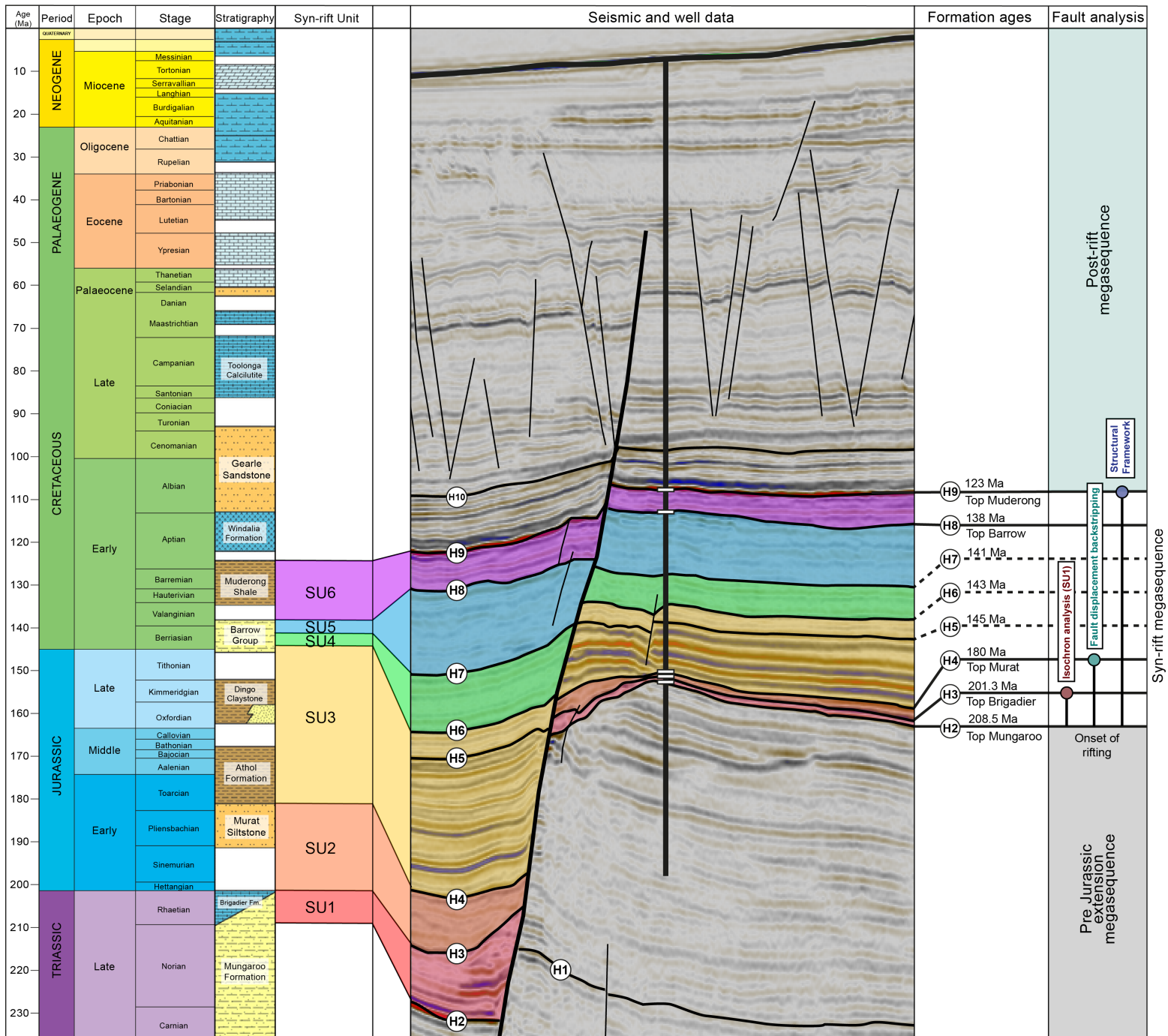
a) Location of Thebe dataset, Exmouth Plateau, NW Shelf of Australia



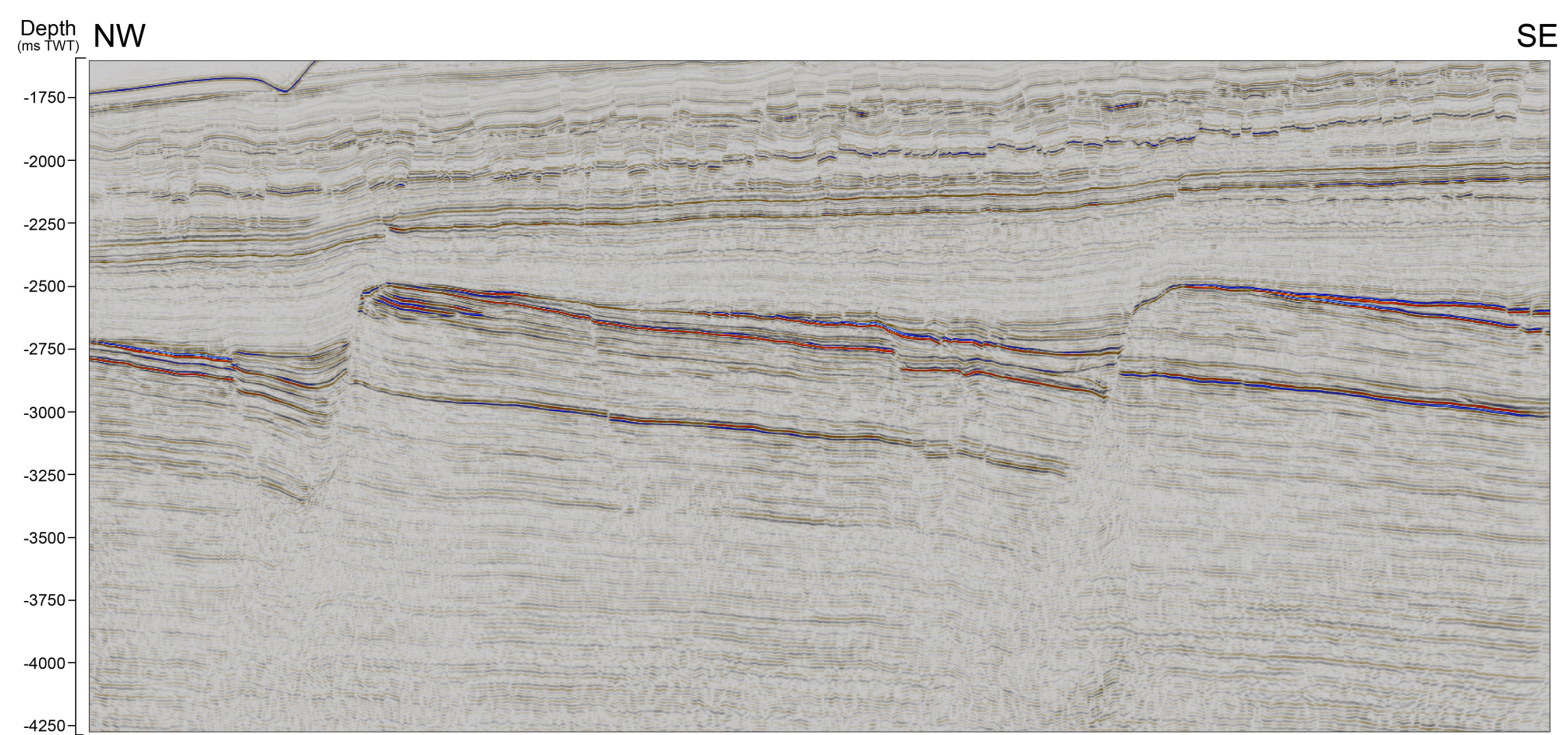
b) Main structural elements of Thebe dataset



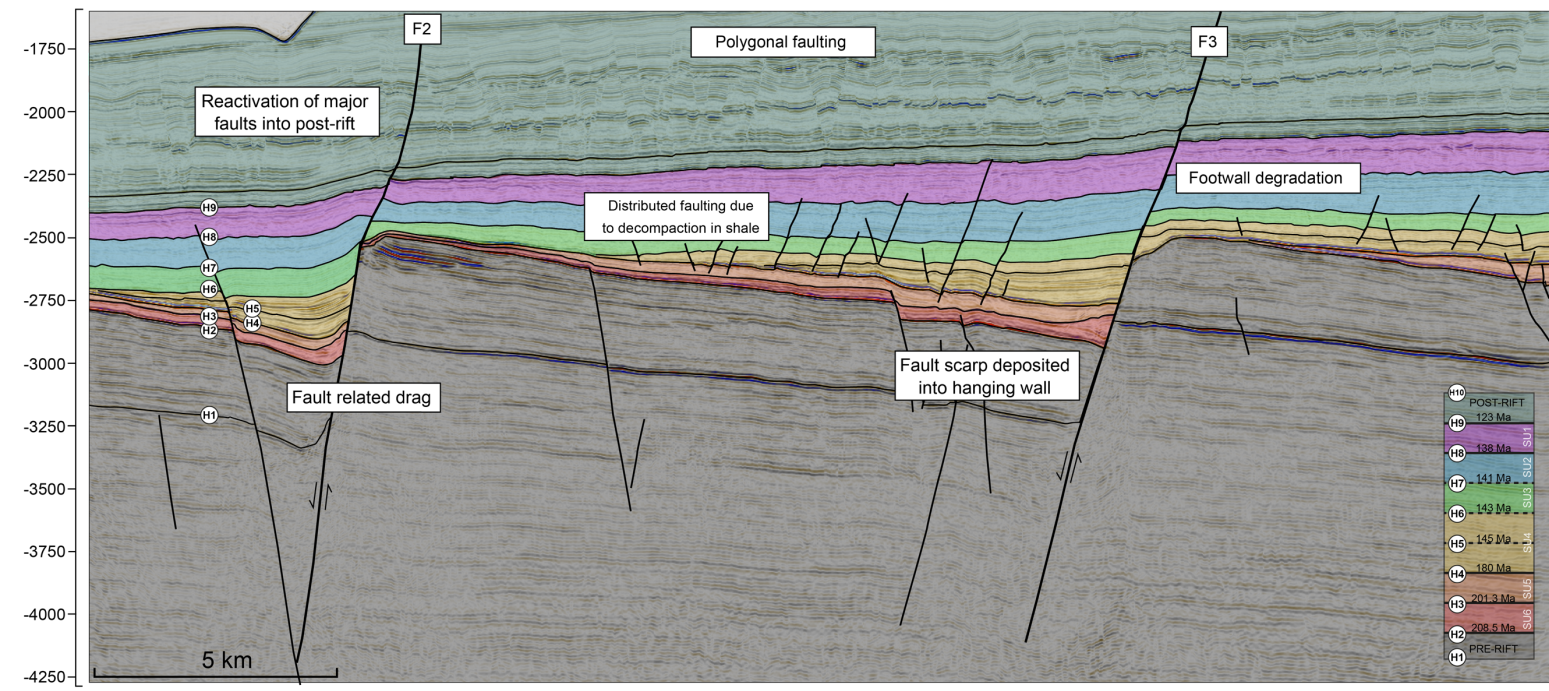
**Figure 3.** The location of our study area: (a) Location of the Thebe field seismic dataset within the Exmouth Plateau, North Carnarvon Basin, North West Shelf of Australia. Sub-basin boundaries are taken from Geoscience Australia. The main structural NE-trends in Fig 3b are from Bilal et al., (2019); and (b) The main structural elements of the Thebe dataset. Major fault systems are labelled in colour, correlating to their respective throw-profiles in Figure 8. Minor faults are coloured in light grey. Wells Thebe-1 and Thebe-2 are shown.



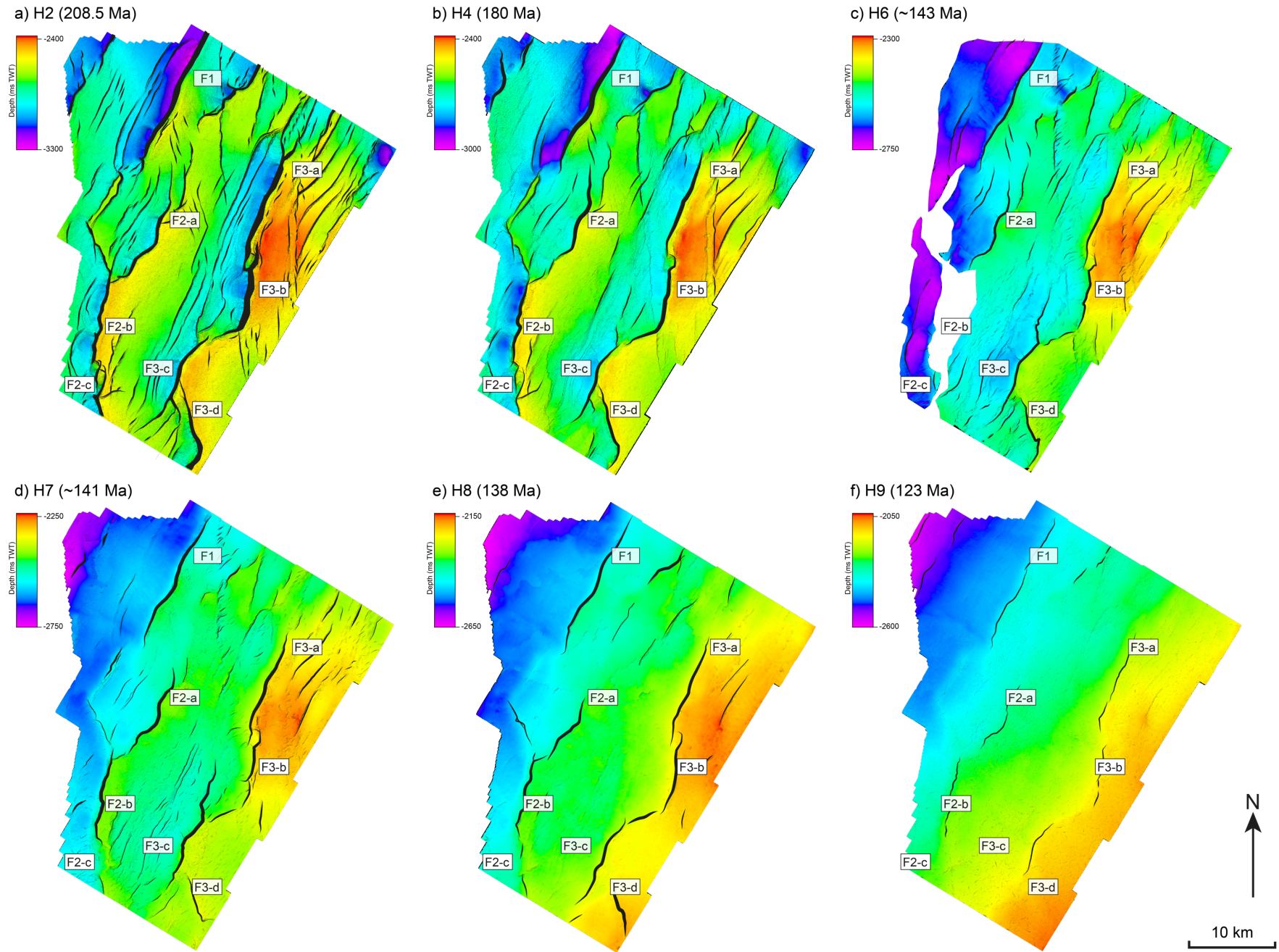
**Figure 4.** Horizons and their respective formation ages. H1 and H10 are pre-rift and post-rift reflectors, whereas H2 – H9 define the syn-kinematic unit of the study area. H4, H5 and H6 do not contain biostratigraphic ages, therefore age is inferred through the assumption of constant sedimentation. Various methodologies used to understand rift evolution are resolved to different stages in time, which has been demonstrated on the right.



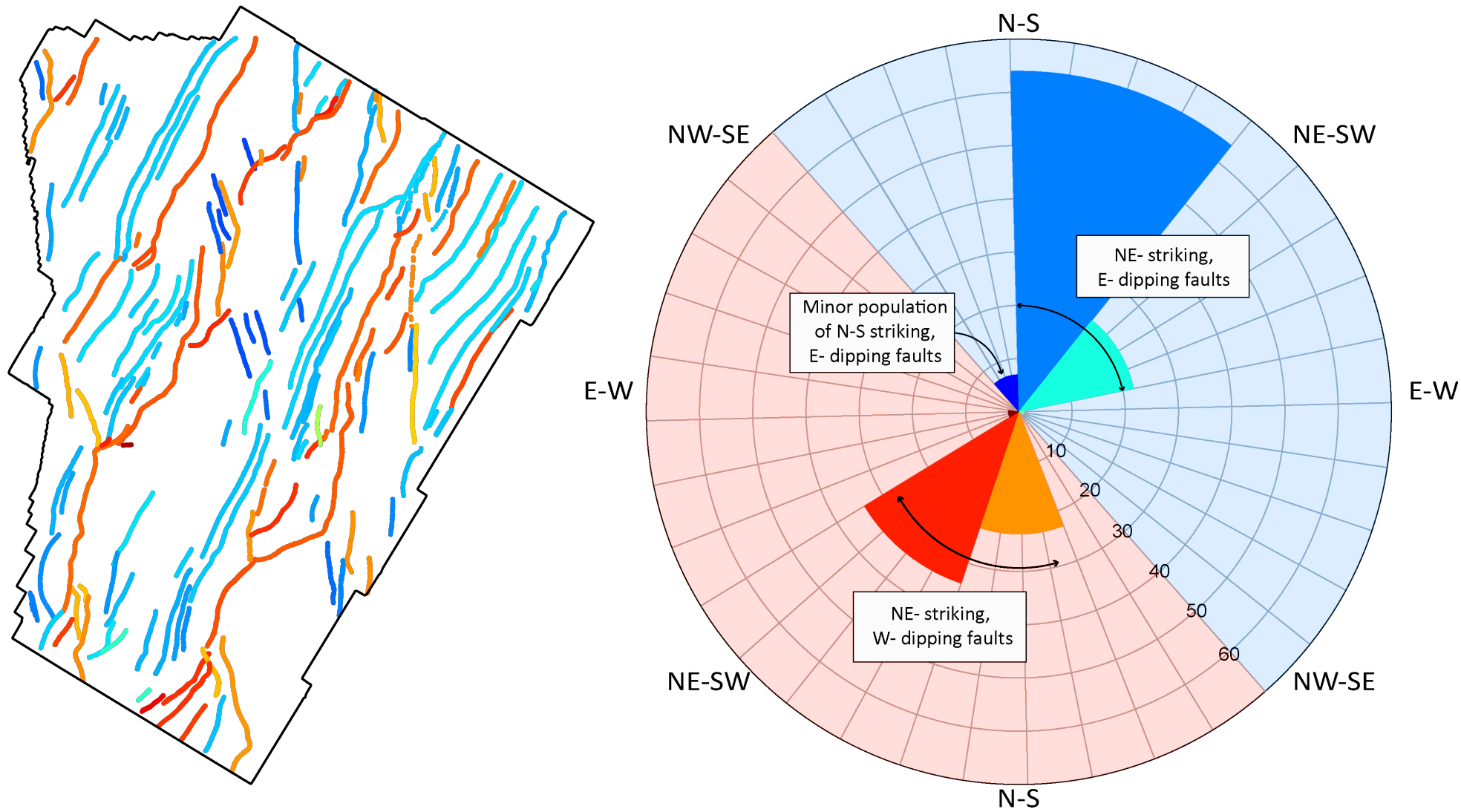
**Figure 5.** Cross section of inline 1524 - location shown on Figure 3b: (a) Without interpretation; and (b) With interpretation and annotation highlighting the main structural features of the overall study area.



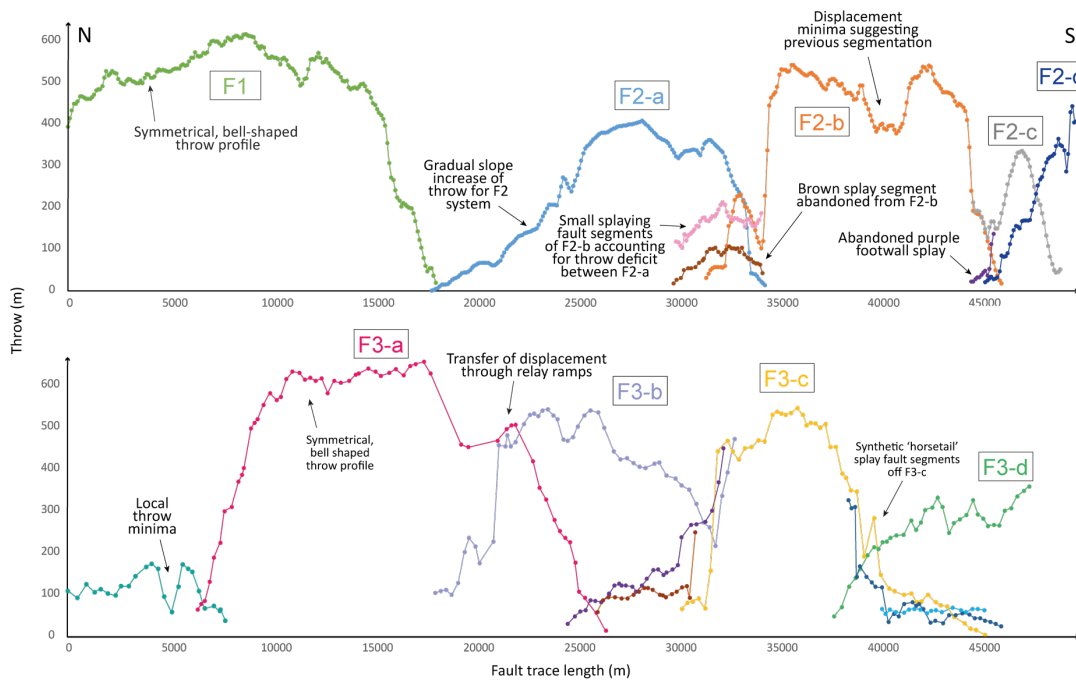
**Figure 6.** Time-surface maps (ms TWT) showing the present day structure of the dataset. Faults at each successive interval are mapped using a variance attribute overlay for higher confidence.



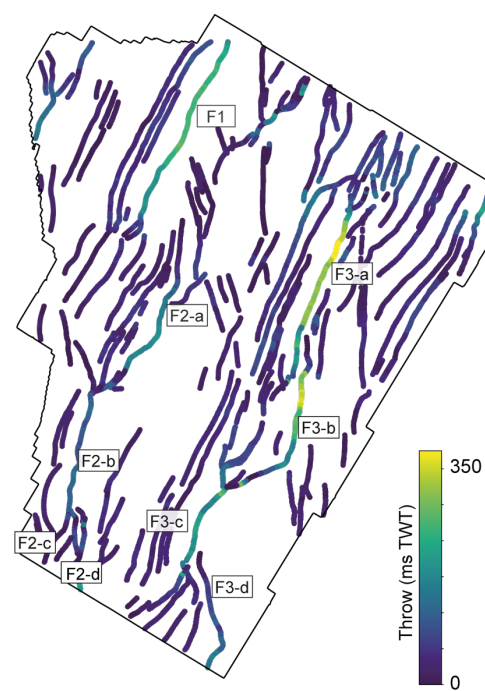
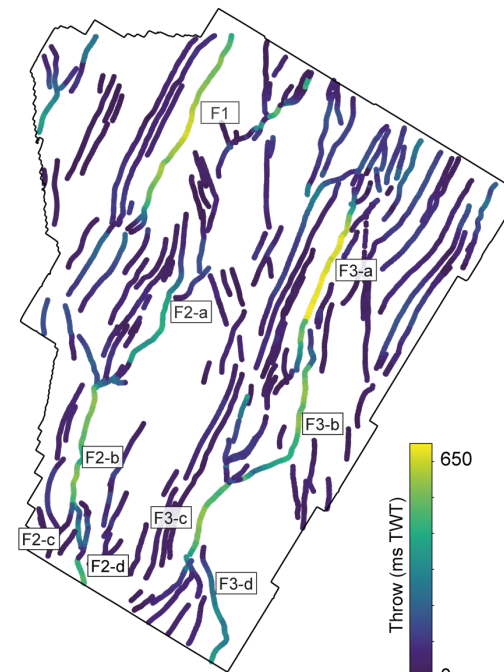
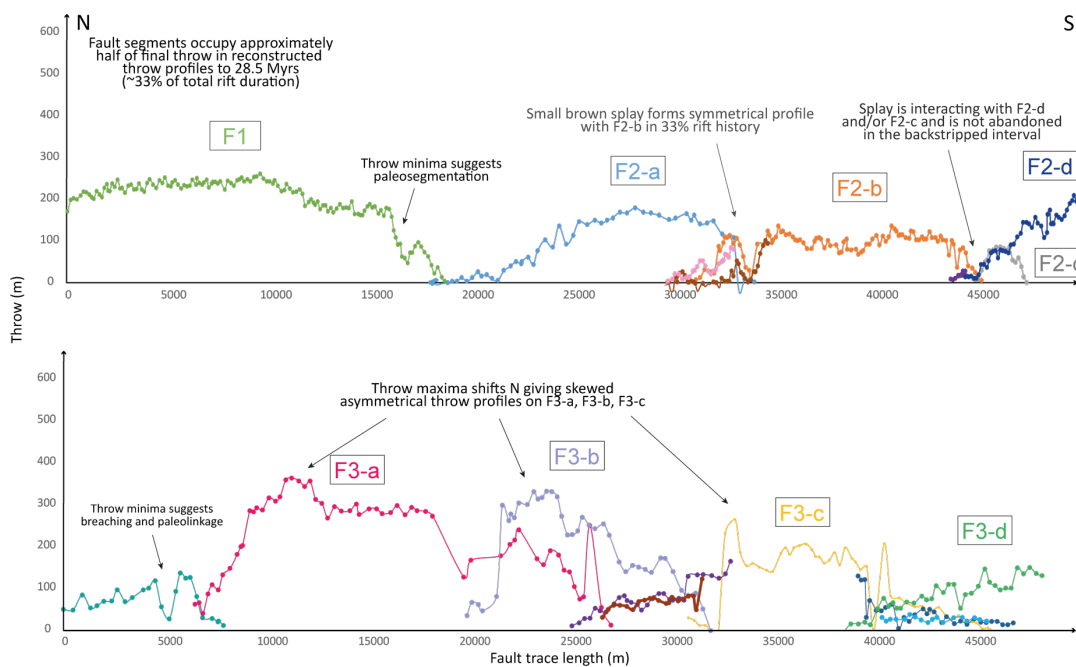
**Figure 7.** Dip direction of the fault array. The rose diagram shows the average strike azimuth and dip direction per fault segment. The rose diagram is split into two halves, whereby the blue coloured faults are east dipping, and the red coloured faults are west dipping.



a) Present day throw distribution; 100% of rift history



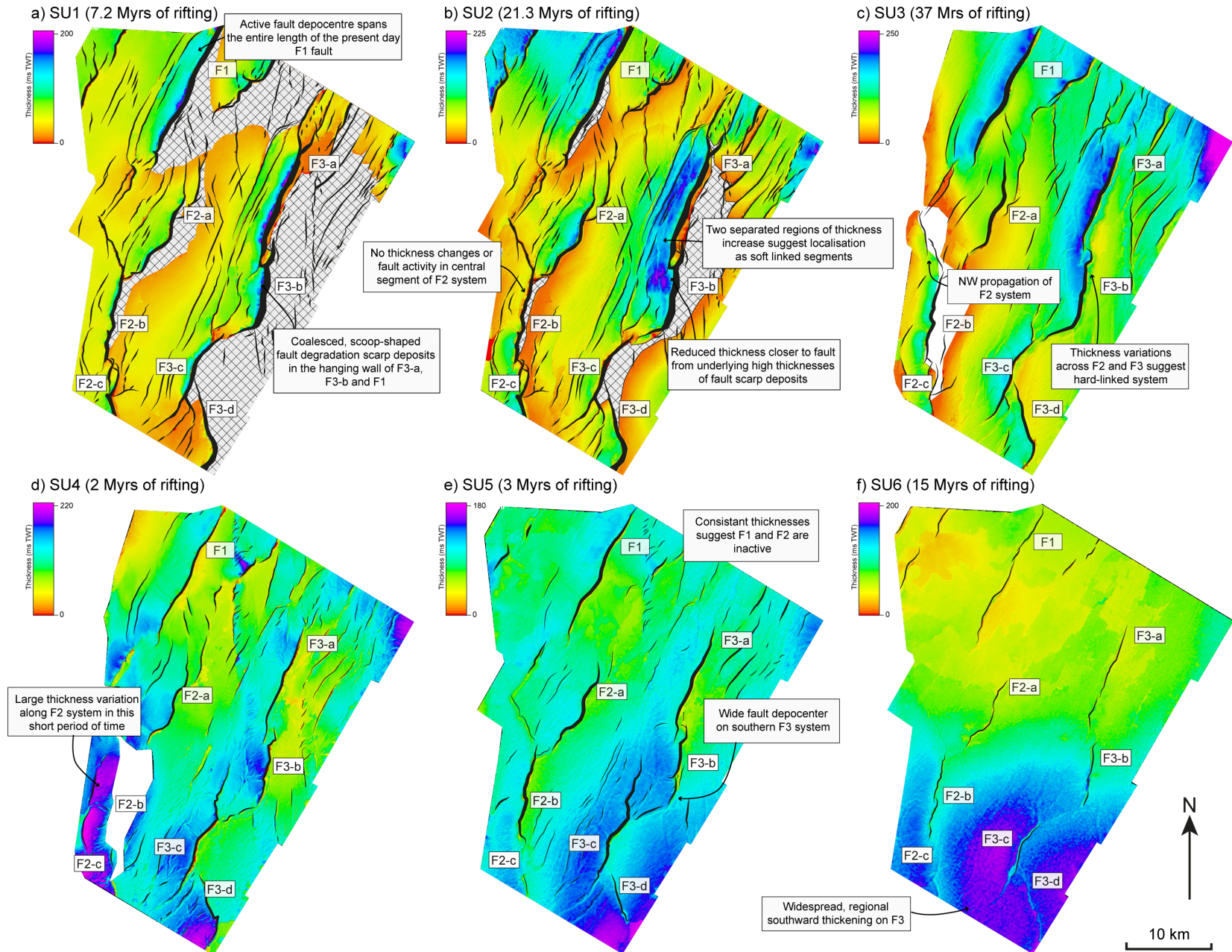
b) Backstripped throw distribution; 33% of rift history

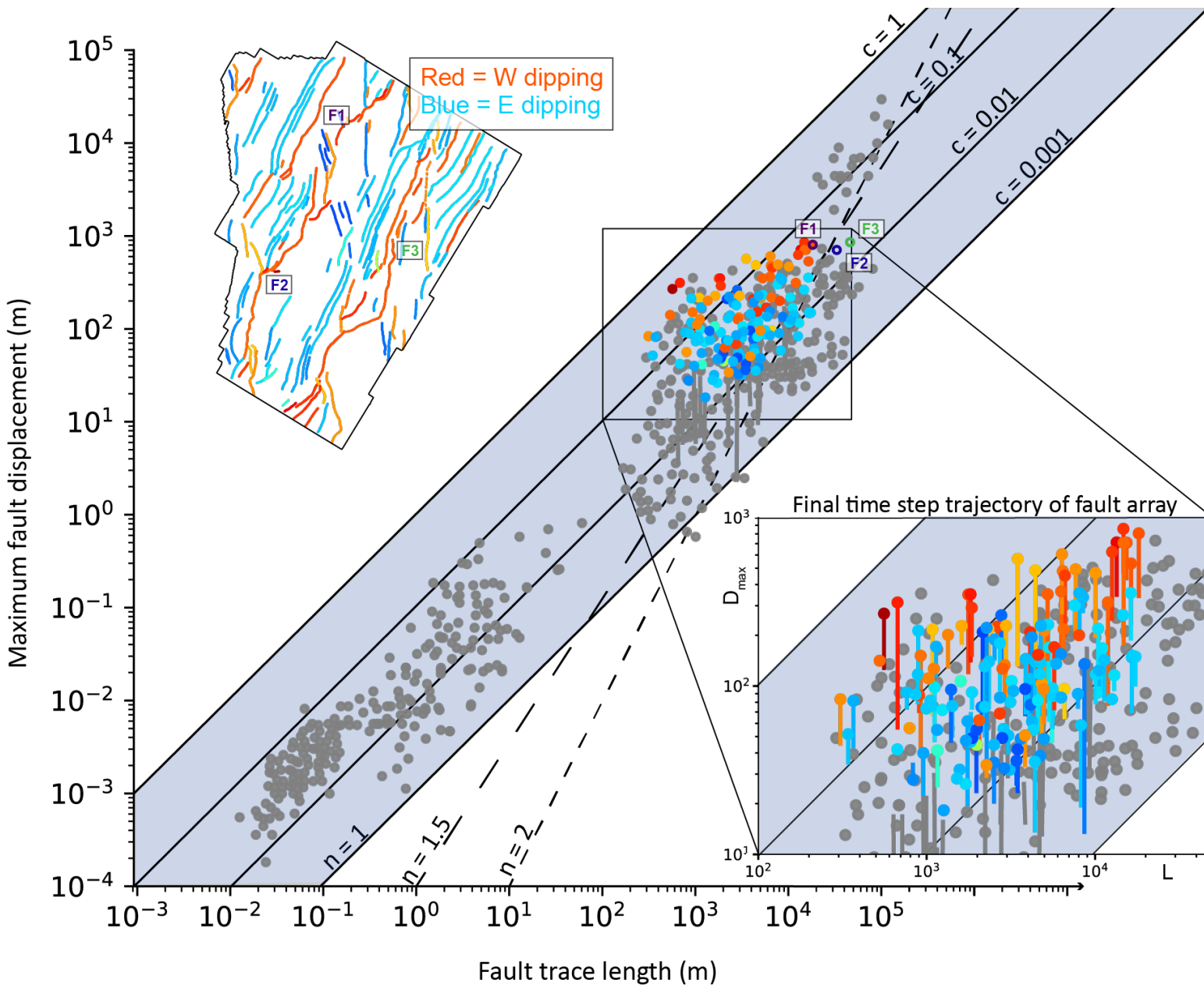


**Figure 8.** Throw length profiles (where fault segments are coloured by Fig 3b) with a map view of throw distribution: (a) The present day fault array, i.e. the final product of rift history; and (b) The reconstructed, backstripped fault array, where faults are backstripped to 28.5 Myrs (H4; Top Murat).

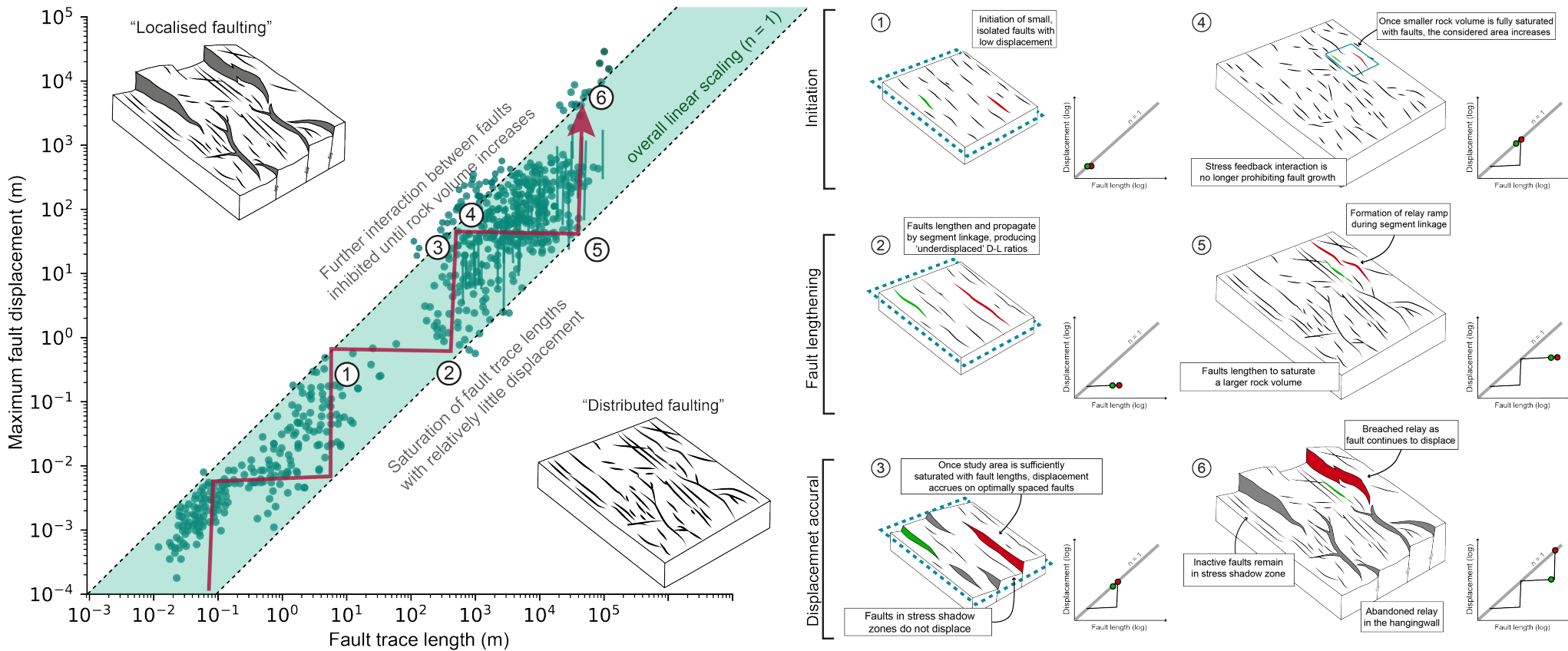


**Figure 9.** Isochron maps showing thickness changes between syn-depositional horizons. The syn-rift unit is divided into six, labelled as SU1-SU6.





**Figure 10.** Our results plotted against global displacement-length data (in grey). Data is coloured by their dip (see Figure 7) whereby blue-coloured faults are east dipping, and red-coloured faults are west dipping. The inset shows the last 67% of rift history, as we connect our present-day values of fault geometry with their respective backstripped values.



**Figure 11.** Our new proposed fault growth model, where the fault trajectory shows growth via alternating phases of fault lengthening (producing a near horizontal trajectory) and fault displacement (producing a near vertical trajectory). Stage 1, 2 and 3 are examples of initiation, linkage and localisation over a smaller region. Faults initiate as small numerous fault segments (1), then lengthen via segment linkage until the rock volume is sufficiently saturated, producing relatively under-displaced D-L values (2), displacement accrues on optimally located and spaced faults (3). Once the smaller region is saturated by fault lengths and displacement, the considered region increases, allowing the fault array to undergo initiation (4), lengthening (5) and localisation (6) again at a larger scale.