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1 **Estimating Ocean Surface Currents from Satellite Observable Quantities**  
2 **with Machine Learning**

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## ABSTRACT

10 Global surface currents are usually inferred from directly observed quanti-  
11 ties like sea-surface height, wind stress by applying diagnostic balance rela-  
12 tions (like geostrophy and Ekman flow), which provide a good approximation  
13 of the dynamics of currents at large scales and low Rossby numbers. However,  
14 newer generation satellite altimeters (like the upcoming SWOT mission) will  
15 capture the high wavenumber variability associated with unbalanced compo-  
16 nents, and applying these balances directly may lead to an incorrect estimate  
17 of the surface flow. In this study, we explore Machine Learning (ML) as an al-  
18 ternate route to infer surface currents from satellite observable quantities using  
19 SSH, SST and wind stress from available ocean GCM simulation outputs as  
20 inputs to make predictions of surface currents (u,v), which are then compared  
21 against the true GCM output. We demonstrate that a linear regression model  
22 is ineffective at predicting velocities accurately beyond localized regions. In  
23 comparison, a relatively simple neural network (NN) can predict surface cur-  
24 rents accurately over most of the global ocean, with lower mean errors than  
25 geostrophy+Ekman. Using a local stencil of neighboring grid points as ad-  
26 ditional input features, we can train the deep learning models to effectively  
27 “learn” spatial gradients and the physics of surface currents. By passing the  
28 stenciled variables through convolutional filters we can help the model learn  
29 spatial gradients much faster. Various training strategies are explored using  
30 systematic feature hold out, to understand the effect of each input feature on  
31 the NN’s ability to accurately represent surface flow. Sensitivity analysis of a  
32 reference NN reveals that besides SSH, geographic information is an essential  
33 ingredient required for making accurate predictions of surface currents with  
34 deep learning.

## 35 1. Introduction

36 The most reliable spatially continuous estimates of global surface currents in the ocean come  
37 from geostrophic balance applied to the sea surface height (SSH) field observed by satellite al-  
38 timeters. For the most part, the dynamics of slow, large-scale currents (up to the mesoscale) are  
39 well approximated by geostrophic balance, leading to a direct relationship between gradients of  
40 SSH and near-surface currents. However, current meter observations for the past few decades  
41 and some of the newer generation ultra-high-resolution numerical model simulations indicate the  
42 presence of an energized submesoscale as well as high-frequency waves / tides at smaller spatial  
43 and temporal scales (Rocha et al. 2016). These high-frequency unbalanced motions are likely to  
44 complicate the estimation of surface currents from from SSH in the upcoming Surface Water and  
45 Ocean Topography (SWOT) mission (Morrow et al. 2018). That is, the high-wavenumber SSH  
46 variability may represent a different, ageostrophic regime, where geostrophy might not be the best  
47 route to infer velocities. Motivated by this problem, in this study we explore statistical models  
48 based on machine learning (ML) algorithms for inferring surface currents from satellite observ-  
49 able quantities like SSH, wind and temperature. These algorithms offer a potential alternative to  
50 the traditional physics-based models.

51 The traditional method of calculating surface currents from sea surface height relies on the  
52 following physical principles. Assuming 2D flow and shallow water pressure, the momentum  
53 equation at the ocean surface can be written as:

$$\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} + f \times \mathbf{u} = -g \nabla \eta + \mathbf{F} \quad (1)$$

54 where  $\mathbf{F}$  is the frictional term due to wind stress. For a sufficiently low Rossby number (accel-  
55 eration terms small; a questionable assumption for the submesoscale regime), the leading-order  
56 balances are geostrophy and Ekman flow. The surface flow can be split into a geostrophic and an

57 ageostrophic, Ekman component ( $\mathbf{u} = \mathbf{u}_g + \mathbf{u}_e$ ), and this leading-order force balance can be written  
 58 as

$$f \times \mathbf{u}_g = -g \nabla \eta \quad (2)$$

$$f \times \mathbf{u}_e = F \quad (3)$$

59 Satellite altimetry products typically provide the sea surface height relative to the geoid (SSH,  $\eta$ ),  
 60 with tidally driven SSH signals removed (LeTraon and Morrow 2001). The geostrophic velocities  
 61 associated with the SSH anomalies are given by:

$$fv_g = g \frac{\partial \eta}{\partial x} \quad (4)$$

$$fu_g = -g \frac{\partial \eta}{\partial y} \quad (5)$$

62 Since geostrophic balance does not hold at the equator ( $f \approx 0$ ), typically (Ducet et al. 2000),  
 63 a higher order “equatorial geostrophic” treatment is used to compute velocities near the equa-  
 64 tor (Lagerloef et al. 1999), which is matched to the geostrophic regime away from the equator.  
 65 Usually, the data-assimilative processing algorithms used to map along-track SSH observations to  
 66 gridded maps (e.g. AVISO Ducet et al. 2000) also involve some form of temporal smoothing. The  
 67 process of combining measurements from multiple satellites and filtering can also lead to spurious  
 68 physical signals (Arbic et al. 2012) leading to exaggerated forward-cascades of energy.

69 In addition to the geostrophic velocities, some products like OSCAR (Ocean Surface Current  
 70 Analysis Real Time, Bonjean and Lagerloef 2002), or GEKCO (Geostrophic and Ekman Current  
 71 Observatory, Sudre and Morrow 2008; Sudre et al. 2013) provide an additional ageostrophic com-  
 72 ponent due to Ekman flow. The Ekman velocity is related to friction, which in the upper layer of

73 the ocean is provided by wind stress and can be derived from the following equations:

$$fv_e + \frac{\partial \tau_x}{\partial z} = 0 \quad (6)$$

$$fu_e - \frac{\partial \tau_y}{\partial z} = 0 \quad (7)$$

$$\tau_x = \rho A_z \frac{\partial u}{\partial z} \quad (8)$$

$$\tau_y = \rho A_z \frac{\partial v}{\partial z} \quad (9)$$

74 Since the Coriolis parameter  $f$  changes sign at the equator, the functional relationship between  
 75 velocity and wind stress is different between the two hemispheres. In the Northern Hemisphere  
 76 we derive:

$$u_e = \frac{1}{\rho \sqrt{2A_z |f|}} (\tau_x + \tau_y) \quad (10)$$

$$v_e = \frac{1}{\rho \sqrt{2A_z |f|}} (-\tau_x + \tau_y) \quad (11)$$

77 And in the Southern Hemisphere:

$$u_e = \frac{1}{\rho \sqrt{2A_z |f|}} (\tau_x - \tau_y) \quad (12)$$

$$v_e = \frac{1}{\rho \sqrt{2A_z |f|}} (\tau_x + \tau_y) \quad (13)$$

78 where  $A_z$  is the linear drag coefficient representing vertical eddy viscosity. Alternatively we can  
 79 write these equations in terms of the Ekman layer depth  $h_{Ek}$  which is related to the eddy viscosity  
 80  $A_z$  as:

$$h_{Ek} = \sqrt{\frac{2A_z}{f}} \quad (14)$$

81 Both of these quantities ( $A_z, h_{Ek}$ ) are largely unknown for the global ocean and are estimated based  
 82 on empirical multiple linear regression from Lagrangian surface drifters (Lagerloef et al. 1999;  
 83 Sudre et al. 2013). Typical values of Ekman depth  $h_{Ek}$  in the ocean range from 10 to 40 meters .

84 So geostrophy + Ekman is the essential underlying physical/dynamical “model” currently used  
 85 for calculating surface currents from satellite observations. This procedure, combining observa-

86 tions with physical principles, represents a top-down approach. A more bottom-up approach would  
87 be a data driven regression model that extracts information about empirical relationships from  
88 data. Recently, machine learning (ML) methods have grown in popularity and have been proposed  
89 for a wide range of problems in fluid dynamics: Reynolds-averaged turbulence models (Ling  
90 et al. 2016), detecting eddies from altimetric SSH fields (Lguensat et al. 2017), reconstructing  
91 subsurface flow-fields in the ocean from surface fields (Chapman and Charantonis 2017; Bolton  
92 and Zanna 2018), sub-gridscale modeling of PDEs (Bar-Sinai et al. 2018), predicting the evolu-  
93 tion of large spatio-temporally chaotic dynamical systems (Pathak et al. 2018), parameterizing  
94 unresolved processes, like convective systems in climate models (Gentine et al. 2018), or eddy  
95 momentum fluxes in ocean models (Bolton and Zanna 2018), to name just a few examples.

96 In this study we aim to tackle a simpler problem than those cited above: training a ML model  
97 to “learn” the empirical relationships between the different observable quantities (sea surface  
98 height, wind stress etc.) and surface currents  $(u, v)$ . The hypothesis to be tested is the follow-  
99 ing: Can we use machine learning to provide surface current estimates that resolve small scale  
100 (balanced/unbalanced) turbulent processes better than geostrophy+Ekman? The motivation for  
101 doing this exercise is two-fold:

- 102 1. It will help us understand how machine learning can be applied in the context of traditional  
103 physics-based theories. ML is often criticised as a “black box.” But can we use ML to com-  
104 plement our physical understanding? This present problem serves as a good test-bed since  
105 the corresponding physical model is straightforward and well understood.
- 106 2. It may be of practical value when SWOT mission launches.

107 We see this work as a stepping stone to more complex applications of ML to ocean remote sensing  
108 of ocean surface currents.

109 This paper is organized as follows. In section 2, we introduce the dataset that was used, the  
110 framework of the problem and identify the key variables that are required for training a statistical  
111 model to predict surface currents. In section 3 we describe numerical evaluation procedure for  
112 baseline physics-based model that we are hoping to match/beat. In sections 4 and 5 we discuss the  
113 statistical models that we used. We start with the simplest statistical model - linear regression in  
114 Section 4 before moving on to more advanced methods like neural networks in Section 5. In sec-  
115 tion 6 we summarize some the findings from the present study, discuss some of the shortcomings  
116 of the present approach, propose some solutions as well as outline some of the future goals for this  
117 project.

## 118 **2. Dataset and Input Features**

119 To focus on the physical problem of relating currents to surface quantities, rather than the ob-  
120 servational problems of spatio-temporal sampling and instrument noise, we choose to analyze a  
121 high-resolution global general circulation model (GCM), which provides a fully sampled, noise-  
122 free realization of the ocean state. The dataset used for this present study is the surface fields from  
123 the ocean component of the Community Earth System Model (CESM), called the Parallel Ocean  
124 Program (POP) simulation (Smith et al. 2010) which has a  $\approx 0.1^\circ$  horizontal resolution, with  
125 daily-averaged outputs available for the surface fields. The model employs a B-grid (scalars at  
126 cell centers, vectors at cell corners) for the horizontal discretization and a three-time-level second-  
127 order-accurate modified leap-frog scheme for stepping forward in time. The model solves the  
128 primitive equations of motion, which, for the surface flow, are essentially (1). Further details  
129 about the model physics and simulations can be found in Small et al. (2014); Uchida et al. (2017).



130 We selected this study because of the long time record of available data (approx. 40 years),  
131 although, in retrospect, we found that all our models can be trained completely with just a few  
132 days of output!

133 A key choice in any ML application is the choice of features, or inputs, to the model. In this  
134 paper, we experiment with a range of different feature combinations; seeing which features are  
135 most useful for estimating currents is indeed one of our aims. The features we choose are all  
136 quantities that are observable from satellites: SSH, surface wind stress ( $\tau_x$  and  $\tau_y$ ), sea-surface  
137 temperature (SST,  $\theta$ ) and sea-surface Salinity (SSS). Our choice of features is also motivated by  
138 the traditional physics-based model: the same information that goes into the physics-based model  
139 should also prove useful to the ML model. Just like the physics-based model, all the ML models  
140 we consider are pointwise, local models: the goal is to predict the 2D velocity vector  $u, v$  at each  
141 point, using data from at or around that point.

142 Beyond these observable physical quantities, we also need to provide the models with geo-  
143 graphic information about the location and spacing between the neighboring points. In the physics-  
144 based model, geography enters in two places: 1) in the Coriolis parameter  $f$ , and 2) in the grid  
145 spacing  $dx$  and  $dy$ , which varies over the model domain. Geographic information can be provided  
146 to the statistical models in a few different ways. The first method involves providing the same  
147 kind of spatial information that is provided to the physical models, *i.e.*  $f$  and local grid spacings  
148 -  $dx$  and  $dy$ . We can also encode geographic information (lat, lon) in our input features, using a  
149 coordinate transformation of the form:

$$\begin{bmatrix} X \\ Y \\ Z \end{bmatrix} = \begin{bmatrix} \sin(lat) \\ \sin(lon) \cdot \cos(lat) \\ -\cos(lon) \cdot \cos(lat) \end{bmatrix} \quad (15)$$

150 to transform the spherical polar lat-lon coordinate into a homogeneous three dimensional coordi-  
151 nate (Gregor et al. 2017). This transformation gives the 3D position of each point in Euclidean  
152 space, rather than the geometrically warped lat / lon space (which has a singularity at the poles and  
153 a discontinuity at the dateline). Note that one of the coordinates - $X$ , that comes out of this kind  
154 of coordinate transformation, is functionally the same as the coriolis parameter ( $f$ ) normalized by  
155  $2\Omega$  ( $\Omega$  = Earth’s rotation). Therefore we will use  $X$  as proxy for  $f$  for all the statistical models  
156 throughout this study. We also explored another approach where the only geographic information  
157 provided to the models is  $X$  ( $= \frac{f}{2\Omega}$ ).

158 Since geostrophic balance involves spatial derivatives, it is not sufficient to simply provide SSH  
159 and the local coordinates pointwise. In order to compute derivatives, we also need the SSH of  
160 the surrounding grid points as a local stencil around each grid point. The approach we used  
161 for providing this local stencil is motivated by the horizontal discretization of the POP model.  
162 Horizontal derivatives of scalars (like SSH) on the B grid requires 4 cell centers. At every timestep,  
163 each variable of the The  $1^\circ$  POP model output has  $3600 \times 2400$  data points (minus the land mask).  
164 We can simply rearrange each variable as a  $1800 \times 1200 \times 2 \times 2$  dataset or split it into 4 variables  
165 each with  $1800 \times 1200$  data points, corresponding to the 4 grid cells required for taking spatial  
166 derivatives. The variables that require spatial a spatial stencil for physical models, we will refer to  
167 as the stencil inputs. For the variables for which we do not need spatial derivatives for (like wind  
168 stress), we can simply use every alternate grid point resulting in a dataset of size  $1800 \times 200$ . We  
169 will refer to these variables as point inputs. For the purpose of the statistical models the inputs  
170 need to be flattened and have all the land points removed. This means that each input variable  
171 has a shape of either  $N \times 2 \times 2$  or  $N$  depending on whether or not a spatial stencil is used ( where  
172  $N = 1800 \times 1200$ — the points that fall over land). Alternatively we can think of the stencilled

173 variable as 4 features of length  $N$ . This kind of stencil essentially coarsens the resolution of the  
174 targets, and point variables.

175 Similarly we can also construct a 3 point time stencil, by providing the values at preceding and  
176 succeeding time steps as additional inputs so that each variable that is stencilled in space and time  
177 has a shape of  $N \times 2 \times 2 \times 3$  (or 12 features of length  $N$ ).

178 This data preparation leads to 10 potential features (for some of which we will use a stencil,  
179 which further expands the feature vector space) for predicting  $u, v$  at each point :  $\tau_x, \tau_y, \text{SSH} (\eta),$   
180  $\text{SST} (\theta), \text{SSS} (S),$  the 3 transformed coordinates  $(X, Y, Z)$  and the local grid spacings  $(dx$  and  $dy)$ .

181 For building any statistical / ML model, we need to split the dataset into 2 main parts, *i.e.*  
182 training and testing. For the purpose of training our machine learning models, the first step involves  
183 extracting the above mentioned variables from the GCM output as the input features and the GCM  
184 output surface velocities  $u, v$  as targets for the ML model. The data extracted from the GCM  
185 output for a certain date (or range of dates) is then used to fit the model parameters. This part  
186 of the dataset is called the training dataset. During training, the model minimizes a chosen cost  
187 function (we used mean absolute error for our experiments, but using mean squared error produced  
188 very similar results) and typically involves a few passes through this section of dataset. The trained  
189 models are then used to make predictions of  $u, v$  for a different date (or range of dates) where the  
190 model only receives the input variables. The model predictions are evaluated by comparing with  
191 the true (GCM output) velocity fields for that particular date (date range). This part of the dataset,  
192 which the model has not seen during training, that is used to evaluate model predictions is called  
193 the test dataset.

### 194 3. Baseline Physics-Based Model: Geostrophy + Ekman

195 The two components of the physics-based model used as the baseline for our ML models are  
 196 geostrophy and Ekman flow. In this section we describe how these two components are numeri-  
 197 cally evaluated for our dataset. For the sake of fair comparison, we evaluate the geostrophic and  
 198 Ekman velocities from the same features that are provided to the regression models. With the  
 199 POP model's horizontal discretization, finite-difference horizontal derivatives and averages are  
 200 defined as (Smith et al. 2010) :

$$\psi_x = [\psi(x + \Delta_x/2) - \psi(x - \Delta_x/2)] / \Delta_x \quad (16)$$

$$\overline{\psi^x} = [\psi(x + \Delta_x/2) + \psi(x - \Delta_x/2)] / 2 \quad (17)$$

201 With the data preparation and stencil approach described in the previous section,  $\eta$  now has a  
 202 shape of  $N \times 2 \times 2$  and the  $f, u, v, dx, dy$  are all variables of length  $N$ . Following (4) and (5) the  
 203 geostrophic velocities ( $u_g^j, v_g^j$ ) are calculated on the stencil as as:

$$v_g^j = g/f^j [\eta^j(1, 1) + \eta^j(0, 1) - \eta^j(1, 0) - \eta^j(0, 0)] / 4dx^j \quad (18)$$

$$u_g^j = -g/f^j [\eta^j(1, 1) + \eta^j(1, 0) - \eta^j(0, 1) - \eta^j(0, 0)] / 4dy^j \quad (19)$$

$$(20)$$

204 where  $j \in [1, N]$ . Similarly the Ekman velocity is calculated numerically from the  $\tau_x^j, \tau_y^j$  and  $f^j$  as

$$u_e^j = \begin{cases} \frac{1}{\rho\sqrt{2A_z|f^j|}}(\tau_x^j + \tau_y^j), & \text{if } f^j > 0 \\ \frac{1}{\rho\sqrt{2A_z|f^j|}}(\tau_x^j - \tau_y^j), & \text{if } f^j < 0 \end{cases} \quad v_e^j = \begin{cases} \frac{1}{\rho\sqrt{2A_z|f^j|}}(-\tau_x^j + \tau_y^j), & \text{if } f^j > 0 \\ \frac{1}{\rho\sqrt{2A_z|f^j|}}(\tau_x^j + \tau_y^j), & \text{if } f^j < 0 \end{cases} \quad (21)$$

205 For calculating the Ekman velocity, we used constant values for vertical diffusivity ( $A_z = 8 \times$   
 206  $10^{-3}m^2/s$ ) and density of water at the surface, ( $\rho = 1027kg/m^3$ ). It should be noted that both  
 207 these quantities vary both spatially and temporally in the real ocean. For the vertical diffusivity

208 we came up with this estimate by solving for  $A_z$  that provides the best fit between zonal mean  
 209  $((u, v)_{true} - (u, v)_g)$  and  $(u, v)_e$ . In the CESM high res POP simulations, the parameterized vertical  
 210 diffusivity was capped around  $100 \text{ cm}^2/\text{s}$  (Smith et al. 2010). For plotting spatial maps for both  
 211 the physics based model predictions as well as the statistical model predictions, the velocity fields  
 212 are then reshaped into  $1800 \times 1200$  arrays, after inserting the appropriate land masks.

#### 213 4. Multiple Linear Regression Model

214 The simplest of all statistical prediction models is essentially multiple linear regression, where  
 215 an output or target is represented as some linear combination of the inputs. The input is charac-  
 216 terized by a feature vector  $\mathbf{x}_i^j$  where  $i \in [1, n_f]; j \in [1, N]$ ,  $N$  being the number of samples, and  $n_f$   
 217 being the number of features. We can now write the linear regression problem as  $\mathbf{U}^j = \mathbf{x}_i^j \mathbf{T} \cdot \beta_i + \delta^j$ .  
 218 where  $\beta_i$  are the coefficients or weight vector. For our regression problem, the input features are  
 219 wind stress, sea surface height and the 3 dimensional transformed coordinates. Of those fea-  
 220 tures,  $\eta, X, Y, Z$  are stencil inputs (meaning 4 input columns per feature) and  $\tau_x, \tau_y$  are the point  
 221 inputs, resulting in a total of 18 input features. The aim therefore is to find the coefficients  $\beta_i$   
 222 that minimize the loss (error) represented by  $\delta^j$  for a training set of  $\mathbf{x}_i^j$  and  $\mathbf{U}^j$  ( $\mathbf{x}_{i_{train}}^j, \mathbf{U}_{train}$ ) and  
 223 use these coefficients for a test set of  $\mathbf{x}_i^j$  ( $\mathbf{x}_{i_{test}}^j$ ) to make predictions for  $\mathbf{U}^j$  ( $\mathbf{U}_{pred}^j$ ). For imple-  
 224 menting linear regression model as well as the deep learning models discussed in this study we  
 225 use the Python library Keras (<https://keras.io>) (Chollet et al. 2015), a high-level wrapper around  
 226 TensorFlow (<http://www.tensorflow.org>).

227 Linear regression can be performed in one of 2 different ways

- 228 • The matrix method or Normal equation method (where we solve for the coefficients  $\beta$  that  
 229 minimize the squared error  $\|\delta\|^2 = \|\mathbf{U} - \mathbf{X}^T \cdot \beta\|^2$  and involves computing the pseudo-inverse  
 230 of  $\mathbf{X}^T \cdot \mathbf{X}$ ).

231 • A stochastic gradient descent (SGD) method (which represents a more general procedure that  
 232 can be used for different regression algorithms with different choices for optimizers and is  
 233 more scalable for larger datasets).

234 The normal-equation method is less computationally tractable for large datasets (large number of  
 235 samples) since it requires loading the full dataset into memory for calculating the pseudoinverse  
 236 of  $\mathbf{x}_i^{\mathbf{jT}} \cdot \mathbf{x}_i^{\mathbf{j}}$ , whereas the SGD method works well even for large datasets, but requires tuning of  
 237 the learning rate. Due to the versatility offered by the gradient descent method we used that  
 238 for performing the linear regression although the normal equation method also produced similar  
 239 results. The essential goal for any regression problem is to minimize a predetermined cost / loss  
 240 function (which for our experiments we chose as the mean absolute error) :

$$J = MAE = \overline{(|u_{pred} - u_{true}| + |v_{pred} - v_{true}|)} \quad (22)$$

241 where the overbar denotes the average over all samples. Fig. 2(a) shows a schematic of the linear  
 242 regression model. The number of trainable parameters for our example with 18 inputs and 2  
 243 outputs is 38 ( $18 \times 2$  weights + 2 biases). For this as well as all subsequent models discussed here,  
 244 we used the same optimizer (Adam, (Kingma and Ba 2017)) and loss function (Mean absolute  
 245 error, MAE). All models are trained on 1 day of GCM output data and we use the same date of  
 246 model output as the training data for all models.

247 We start by splitting the global ocean into 3 boxes to zoom into three distinct regions of dynami-  
 248 cal importance in oceanography, namely the Gulf stream, Kuroshio and Southern ocean / Antarctic  
 249 circumpolar current (ACC). The Kuroshio region is chosen to extend south of the equator to in-  
 250 clude the equatorial jets as well as to test whether the models can generalize to large variations  
 251 in  $f$ . The daily averaged GCM output surface speed on a particular reference date, with the three  
 252 regions (marked by three different colored boxes) is shown in Fig. 1. WE then train 3 different

253 linear regression models with training data from these three sub-domains. We also trained a linear  
254 regression model for the whole globe using the same model architecture. During training, the  
255 models are fed a shuffled batch of the training data with 32 samples in each batch and the loss  
256 (MAE) is computed for the batch. For the linear regression model as well as for all the neural net-  
257 works discussed in this study we present here we kept the batch size constant. Changing the batch  
258 size does not significantly alter the loss at the end of training, but smaller batch sizes generally  
259 help the model learn faster. The different models, the number of epochs (an epoch is defined as  
260 one pass through the training dataset) used for each, and losses at the end of training and during  
261 evaluation against a test dataset are summarized in table 1. The evolution of model loss function  
262 during training for the 3 different models are presented in Fig. 3. Linear regression is shown in  
263 the darker colors. The big jumps in the loss function correspond to the end of an epoch. We plot  
264 the models' training progress in the Gulf Stream region for 8 epochs, and for 5 epochs on the  
265 Kuroshio and ACC regions. The trained models are then evaluated for a test dataset (which the  
266 model has not seen, GCM output from a different point in time) and the evaluation loss is plotted  
267 as the horizontal dashed lines. The linear regression model trained on the whole globe is also  
268 evaluated for each subdomain (gulf stream, Kuroshio, ACC) and the global model evaluation loss  
269 is plotted as the dotted line. Comparing the model losses in the 3 different sections, we find that  
270 the linear regression model performs the most poorly for the Kuroshio region (i.e the subdomain  
271 with the most variation in  $f$ ). The model does progressively better for the gulf stream and the  
272 ACC in terms of MAE, where the variations in  $f$  are relatively smaller in comparison. However,  
273 the root mean squared error of predicted velocities is still quite large in all these regions (second  
274 panels of Figs. 4, 5, 6). The linear regression model trained on the global ocean does even worse  
275 during evaluation. Since geostrophy relies on non-linear combination of the Coriolis parameter  
276 ( $f$ ) with the spatial gradients, linear regression is ineffective at predicting velocities beyond local-

277 ized regions with small variation of  $f$  or little mesoscale activity. This shows that a linear model  
278 fails to accurately represent surface currents in any region that includes significant variation in the  
279 Coriolis parameter  $f$ . Even in regions far enough from the equator such that the variation in  $f$   
280 is not significant (like the gulf stream or ACC), the performance of such a linear model does not  
281 improve with more training examples and/or starts overfitting. We also show that a lower MAE  
282 during training does not necessarily guarantee that the model is picking up on the small scale  
283 fluctuations in velocity, as can be seen from the relatively large squared errors especially in and  
284 around high surface current regions (Figs. 4, 5, 6). We suspect that this failure is largely due to  
285 the fact that the linear model is trying to fit the velocities as a linear combination of the different  
286 features, whereas realistic surface current predictions should be based on non linear combinations  
287 of features.

288 These non-linear combinations between the different features can instead be incorporated by  
289 using deep learning or artificial neural networks. In the following section, we demonstrate the  
290 feasibility of using neural networks to extract the nonlinear relationships from data.

## 291 **5. Deep Learning: Artificial Neural Networks**

292 Artificial neural networks (or neural networks for short) are machine learning algorithms that are  
293 loosely modeled after the neuronal structure of a biological brain but on a much smaller scale. A  
294 neural network is composed of layers of connected units or nodes called artificial neurons (LeCun  
295 et al. 2015; Nielsen 2015; Goodfellow et al. 2016) that combine input from the data with a set of  
296 weights and passes the sum through the node’s activation function along with a bias term, to the  
297 subsequent set of nodes, to determine to what extent that signal progresses through the network  
298 and how it affects the ultimate outcome. Neural nets are typically “feed-forward,” meaning that  
299 data moves through them in only one direction. A layer is called densely connected when each



300 node in that layer is connected to every node in the layers immediately above and below it. Deep  
301 learning, or deep neural networks is the name used for “stacked neural networks” - i.e., networks  
302 composed of several layers.

303 In the past few years, there have been several studies applying machine learning tools, and more  
304 specifically deep learning methods to model physical/dynamical processes. For example, deep  
305 neural networks (DNN) have been used to develop Reynolds-averaged turbulence models (Ling  
306 et al. 2016) to show that a neural network can be trained to preserve Galilean invariances. Lguensat  
307 et al. (2017) developed a convolutional neural network (CNN) based architecture for automated  
308 eddy detection and classification from Sea Surface Height (SSH) maps. Chapman and Charan-  
309 tonis (2017) constructed a form of neural network known as a self-organising map to reconstruct  
310 sub-surface velocities in the Southern ocean using satellite altimetry data and Argo floats. Pathak  
311 et al. (2018) used yet another recently developed machine learning algorithm, known as reser-  
312 voir computing, to make predictions for the evolution of a very large spatiotemporally chaotic  
313 dynamical systems. Another recent study (Bar-Sinai et al. 2018) demonstrated the capabilities  
314 of a CNN based method for coarse-graining partial differential equations. This study has strong  
315 potential implications for future data-driven subgrid scale parameterizations in atmospheric and  
316 oceanographic models. In a recent publication, Gentine et al. (2018) used deep neural networks  
317 (DNN), trained with a outputs from a superparameterized climate model, to successfully predict  
318 most of the key features of embedded convection necessary for climate simulation, thereby sug-  
319 gesting a strong future for data-driven convection parameterizations in climate models. On the  
320 oceanographic modeling side, Bolton and Zanna (2018) used CNNs trained on spatio-temporally  
321 degraded data from a high-resolution quasi-geostrophic ocean model to successfully replicate the  
322 spatio-temporal variability of the eddy momentum forcing. Furthermore, the CNN based method  
323 was shown to be generalizable to a range of dynamical behaviours, and could be forced to respect

324 global momentum conservation. One of the common criticisms of deep-learning methods has been  
325 that, they are a “black-box”, i.e., lacking any simple intuitive physical interpretations. Some of  
326 these recent works (Ling et al. 2016; Bolton and Zanna 2018; Gentine et al. 2018) showed that  
327 data-driven approaches, even with limited data, can be used in conjunction with physical models,  
328 to help speed up some of the time intensive/ memory intensive processes in the physical models,  
329 while still respecting physical principles.

330 Our neural network code was written using the Python library Keras (<https://keras.io>) (Chol-  
331 let et al. 2015), a high-level wrapper around TensorFlow (<http://www.tensorflow.org>). The feed-  
332 forward NNs consist of interconnected layers, each of which have a certain number of nodes. The  
333 first layer is the input layer, which in our case is a stacked vector containing the input variables just  
334 like in the linear regression example above. The last layer is the output layer, which is a stacked  
335 vector of the two outputs (U,V). All layers in between are called hidden layers. The activation  
336 function, i.e. the function acting on each node – is a weighted sum of the activations in all nodes  
337 of the previous layer plus a bias term, passed through a non-linear activation function. For our  
338 study, we used the Rectified Linear Unit (ReLU) as an activation function. The output layer is  
339 purely linear without an activation function. Training a NN means optimizing the weight matrices  
340 and bias vectors to minimize a loss function – in our case the MAE - between the NN predictions  
341 and the true values of  $(u, v)$ .

342 The model reduces the loss, by computing the gradient of the loss function with respect to all  
343 weights and biases using a backpropagation algorithm, followed by stepping down the gradient  
344 – using stochastic gradient descent (SGD). In particular we use a version of SGD called Adam  
345 (Kingma and Ba 2014, 2017). Although most neural network strategies involve normalizing the  
346 input variables, we did not use any normalization, since the normalization factors would be largely  
347 dependent on the choice of domain / ocean basin, given that the dynamical parameters (like SSH

348 and wind stress) vary widely across the different ocean basins. Instead we wanted the NN to be  
349 generalizable across the whole ocean.

350 We construct a 3-hidden-layer neural network to replace the linear regression model described  
351 in the previous section. A schematic model architecture for the neural network is presented in  
352 Fig. 2(b). Using the same basic model architecture, we train three NNs on the same three sub-  
353 domains (Gulf Stream, Kuroshio, ACC) along with one which is trained on the global ocean.  
354 Everything including batch size, the training data, the targets, the input features and the number  
355 of epochs the model is trained for in each region is kept exactly the same as what we used for the  
356 linear regression examples. The only thing that we changed is the model, where instead of 1 layer  
357 with no activation we now have three hidden layers with a total of 1812 trainable parameters.

358 Just like we did with the linear regression model, we then tracked the evolution of the models’  
359 loss function as it moved through batches of input data over multiple epochs (Fig. 3, lighter col-  
360 ored lines in all panels). As we can see, in comparison to the linear regression model, the NNs  
361 perform significantly better at reducing the loss in all the ocean regions. What is even more strik-  
362 ing is that the NN trained on the globe (dashed line) consistently outperforms the local models,  
363 predicting surface currents with lower MAE/ MSE than the models trained on the local subdo-  
364 mains. This is especially noticeable for the Kuroshio region (Fig. 3, second panel), where the NN  
365 trained on the globe manages to get the signature of the equatorial currents better than the NN  
366 trained specifically in that region (compare panels 3 and 4 of Fig. 5) and gets the absolute error  
367 down to  $\approx 5\text{cm/s}$ . This shows, that in comparison to the linear model the neural network actually  
368 manages to learn the physics better when it receives a more spatially diverse input data, and is  
369 therefore more generalizable. Even though the linear regression models all manage to get the loss  
370 down to comparable magnitude, looking at the spatial plot of the predicted squared error Figs. 45  
371 and 6 gives us an idea how poorly it does at actually learning the physics of surface currents. In

372 comparison, even a relatively shallow 3-hidden-layer neural network performs remarkably better  
373 with very few localized hotspots of large errors. This is to be expected since the largest order  
374 balance, i.e., geostrophy relies on non-linear combination of the Coriolis parameter ( $f$ ) with the  
375 spatial gradients, and therefore these non-linear combinations are not represented by linear regres-  
376 sion and are better captured by a neural network with dense interconnected layers with non-linear  
377 activation functions.

378 In Fig. 7 we plot the joint histogram of the zonal and meridional velocity predictions against  
379 the true (GCM output) values for the physical model, linear regression model (trained on the local  
380 subdomain) and the locally and globally trained neural networks in the ACC sector. From these  
381 joint histograms, it is obvious that the physical model, the local and global neural networks all  
382 predict velocities that are extremely well correlated with the true velocities in this region. In  
383 addition the root mean squared (rms) errors normalized by the rms velocities are also very well  
384 correlated between the physical model and neural network predictions. This provides us with  
385 reasonable confidence that the model is indeed learning the physics of surface geostrophy and  
386 Ekman flow.

387 We also plotted the squared errors in predicted velocity from the physical model (geostro-  
388 phy+Ekman) and the local Rossby number (expressed as the ratio of the relative vorticity  $\zeta =$   
389  $v_x - u_y$ , to the planetary vorticity  $f$ ) in the three domains (Gulf Stream - Fig. 4; Kuroshio - Fig. 5  
390 and the ACC - Fig. 6). It is interesting to note that the localized regions in large root squared er-  
391 rors in both the neural network and physical models coincide with regions where the local Rossby  
392 number is high. High Rossby numbers indicate unbalanced flow and the specific regions where  
393 we see high Rossby numbers are typically associated with heightened submesoscale activity. We  
394 speculate that the prediction errors in these locations are due to the NN's inability to capture higher

395 order balances (e.g. gradient wind, cyclostrophic balance) that are necessary to fully capture the  
396 small scale variability associated with these motions and close the momentum budget.

397 The NN also generally predicts weaker velocities near the Equator where the true values of the  
398 surface currents are quite large (due to strong Equatorial jets). This can lead to large errors for  
399 the global mean, which get magnified when the differences are squared. However we know that  
400 geostrophic and Ekman balance also doesn't hold near the Equator. A fairer comparison would  
401 therefore involve masking out the near equatorial region ( $5^{\circ}N - 5^{\circ}S$ ) for both the statistical model  
402 (i.e. NN predictions) as well as for the physical model (*geo + ekman*). region sincetfails

#### 403 *a. Neural networks with Convolutional Filters*

404 In Section 2 we explained how we can use the local  $2 \times 2$  stencil to expand the feature vector  
405 space by a factor of 4. We can further expand the feature vector space by passing all the stenciled  
406 input features through  $k$  convolutional filters of shape  $2 \times 2$ . If  $k > 4n_f^s$  where  $n_f^s$  is the number  
407 of input features with a stencil, we end up with more input features that goes into the NN than  
408 before. There is very little functional difference between this kind of training approach and the  
409 one discussed previously, except that we end up with more trainable parameters, which we can  
410 potentially use to extract even more information from the data. We should note that this is techni-  
411 cally not the same as convolutional neural networks, which are typically used for image analysis  
412 and classification, where the convolutional layers serve to *reduce* the feature vector space without  
413 losing information. This is particularly important for problems like image classification where it is  
414 needed to scale down large image datasets without losing feature information. A schematic of this  
415 subcategory of neural network is shown in Fig. 2(c). After applying the convolutional filter and  
416 passing it through a reshape layer in keras the point inputs and filtered stencil inputs are passed  
417 through a Leaky ReLU before being fed into a similar 3-hidden layer NN framework as described

418 before. Using a similar procedure, we can also apply  $k$  3D convolutional filters of shape  $2 \times 2 \times 3$   
419 on the time and space stenciled inputs to effectively end up with  $k$  input features of length  $N$  for  
420 the stencil variables (. 2(d)). The goal with the time stenciled input being to potentially learn time  
421 derivatives and explore how the tendencies can affect the NN projections. In hindsight, this data  
422 set is probably not be the most suited for this kind of approach since the variables we used as input  
423 features are daily averaged and any fast-time scale / tendency effects that we hoped to capture from  
424 multiple snapshots of the same variable are probably filtered out by the time averaging. These two  
425 approaches are virtually identical with slightly different preprocessing of the input data.

#### 426 *b. Model dependence on choice of input features*

427 We then trained these NNs with varying combinations of input features to explore how the  
428 choice of input features can influence the model training rate and loss. Feeding the NN models  
429 varying combination of input features, either as stencilled or as point variables and by selectively  
430 holding out specific features for each training case allowed us to assess the relative importance of  
431 each physical input variable for the neural network’s predictive capability. The different models  
432 with their corresponding input features and the number of trainable parameters for each case are  
433 summarized in table 2. As with all previous examples, we chose mean absolute error as the loss  
434 function for all these experiments. We performed a few training exercises using the mean squared  
435 error instead and did not notice any significant difference. For models numbered 1 - 13, we used  
436 a 2 point space-stencil and for models 1t - 10t, in addition to a stencil in space, we provide a 3  
437 point time stencil with the intention of helping the neural network “learn” time derivatives. The  
438 different experiments listed in table 2, can broadly be categorized into 6 groups based on their  
439 input features. In group 1, is model 1, where the model only sees  $\eta$  (stencil) and wind stress,  
440  $\tau$  (point) as input features. No spatial information is provided. In the second category, we have

441 models that receive  $\eta$  (stencil) and spatial information  $\mathbf{X}$  in some form, but no wind stress. This  
442 includes models 2, 5t and 7t. The third category describes models that receive  $\eta, \theta$  (stencil) and  
443 spatial information  $\mathbf{X}$  and no wind stress and includes models 3, 6t and 8t. The fourth category  
444 describes models that receive SSH ( $\eta$ ), spatial information ( $\mathbf{X}$ ) and wind stress ( $\tau$ ) but no SST  
445 and includes models 4, 6, 7, 10, 1t, 3t. The fifth category of models receive SSH ( $\eta$ ), SST ( $\theta$ ),  
446 spacial information ( $\mathbf{X}$ ) and wind stress and the only input feature these models don't receive in  
447 any form is sea surface salinity ( $S$ ). This includes models 5, 8, 9, 11, 2t, 4t. The sixth and final  
448 category represents models tat receive all the input features ( $\eta, \theta, S, \mathbf{X}, \tau$ ) in some form or another  
449 and includes models 13, 9t and 10t.

450 As mentioned previously, spatial information is provided in one of 3 ways, (a) in the form of 3  
451 dimensional transformed coordinates ( $X, Y, Z$ ), (b) just the coriolis parameter ( $X$  here serves as  
452 a proxy for the coriolis parameter) and (c) with both the Coriolis parameter and local  $dx$  and  $dy$   
453 values. Barring a few examples (models 10, 11) windstress is always provided as a point variable  
454 and apart from models 6, 7, 8, 9, none of the models receive a stencil in the spatial coordinates. We  
455 also trained a few models without SSH as an input feature, but the loss in all these cases was much  
456 larger than those shown here ( $> 50cm/s$ ) and the NNs fail to pick up any functional dependence  
457 on the input features. Those cases are therefore not presented. Each of these models are trained  
458 for 4 Epochs on the same day of data (or 3 consecutive days centered around that date for the time  
459 stencilled cases).

460 In Fig. 9 we summarize the findings from these experiments by plotting the rms error for all  
461 the model predictions along with the rms error for the physical model predictions side by side.  
462 With the exception of 5 models (model 1, 5t, 7t, 6t, and 8t) all our NN model predictions have  
463 lower domain mean squared errors than the physics-based models. In terms of features, model  
464 without spatial information has the largest error, followed by models without wind stress (The

465 absolute largest error is for the model without SSH, which is too big to be considered here).  
466 This signifies that to accurately represent surface currents, apart from SSH, the most important  
467 pieces of information required by the neural networks to successfully learn the physics of surface  
468 currents are spatial information and wind stress. It is striking to see how much the model struggles  
469 without spatial information. This implies that latitude dependence is a critical component for a  
470 NN to be able to predict surface currents accurately. It is only expected since the dynamics of  
471 surface currents do depend very strongly on latitude and therefore it is impossible to construct a  
472 meaningful prediction model based on just snapshots without any knowledge of latitude.

473 The zonal mean rms error for the predictions from some of the representative models from  
474 the 6 categories described above are shown in Fig. 8. The NNs all generally predict weaker  
475 velocities near the Equator where the true values of the surface currents are quite large (due to  
476 strong Equatorial jets). This can lead to large errors for the global mean, which get magnified  
477 when the differences are squared. However we know that geostrophic and Ekman balance also  
478 doesn't hold near the Equator. Therefore to allow for a fair comparison between all the models,  
479 we mask out the rms errors in a  $10^\circ$  latitude band surrounding the equator ( $5^\circ N - 5^\circ S$ ) for both the  
480 physical and statistical models. Out of all the models, model 1 which does not receive any spatial  
481 information ( $\mathbf{X}$ ), has the highest mean squared errors throughout the globe. For the models that  
482 don't see wind stress ( $\tau$ ) as an input feature, the rms errors are comparable if not smaller at most  
483 latitudes when compared to the physics-based model where you only consider geostrophy (dashed  
484 black line). Additionally, all models that receive  $\eta$ ,  $\tau$  and  $\mathbf{X}$  in some form perform consistently  
485 better than geostrophy+Ekman at all latitudes (except for near the equator where the physics-  
486 based models and the NN are all equally inadequate). We noticed that during training, the NN's  
487 minimize the loss function slightly faster when a stencil is provided for the spatial coordinates, but  
488 after a few epochs the differences in training loss between models that receive a spatial stencil and



489 models that dont, diminish very rapidly. During prediction also, the models that receive stencils  
490 in spatial coordinates perform slightly better especially at the high latitudes than the ones where  
491 spatial information is provided pointwise.

492 Therefore among the various strategies tested, for this particular dataset, the models that perform  
493 the best in terms of prediction rms error are the models that receive SSH, wind stress and spatial  
494 information with a space stencil. The three point time stencil does not add anything meaningful  
495 and appears to hurt, rather than help the model overall, which was surprising, even though in  
496 hindsight we speculate that this might be due to the daily averaged nature of the POP model  
497 output. Variables like sea surface temperature and sea surface salinity have very little impact on  
498 the model as well.

499 In terms of choice of features, model 13 stands out as the most practical and physically mean-  
500 ingful training strategy for a few reasons.

- 501 • It is the most complete in terms of features
- 502 • It is the most straightforward to implement, since it does not involve calculating any trans-  
503 formed 3 dimensional coordinates. (All the input variables would be readily available for any  
504 gridded oceanographic dataset.)
- 505 • It is one of the models with the lowest prediction rms errors.

506 For these specific reasons we choose model 13 as the reference for performing a sensitivity anal-  
507 ysis. The purpose of this analysis is to characterize the sensitivity of the model to perturbations  
508 in the different input features during testing/prediction. For the sensitivity tests we simply add a  
509 gaussian noise of varying amplitude to each of the input variables, while keeping the rest of the  
510 input variables fixed. For each of the input variables ( $x_i \in \{\eta, \theta, X, dx, dy, \tau_x, tau_y\}$ ), we chose 3  
511 different zero-mean gaussian noise perturbations with the standard deviations of  $0.5\sigma(x_i)$ ,  $\sigma(x_i)$ ,

512 and  $2\sigma(x_i)$ , where  $\sigma(x_i)$  is the standard deviation of the corresponding input variable  $x_i$ . The  
513 model loss is then evaluated for each of these perturbations and normalized by the amplitude of  
514 perturbations (right panel Fig. 10). This normalization is done to level the playing field for all the  
515 input variables and allow for a one-to-one comparison since the different input variables vary in  
516 orders of magnitude (*e.g.* the amplitude of perturbations in SSH is  $O(100)$ , while the amplitude  
517 of perturbations in wind stress is  $O(1)$  and therefore a perturbation of amplitude  $\sigma(\eta)$  in  $\eta$  would  
518 lead to a much larger model error than a perturbation of  $\sigma(\tau_x)$  in  $\tau_x$  would, as can be seen from  
519 the log scaling of the y-axis in the left panel of Fig. 10).

520 Given what we learned about the importance of spatial coordinates for NN training, it is not  
521 surprising to see that for prediction also, the NN is most sensitive to perturbations in the coriolis  
522 parameter (or X). The input variables that the model is most sensitive to, arranged in descend-  
523 ing order of model sensitivity are coriolis parameter, SSH and wind stress, followed by SST. The  
524 model is not particularly sensitive to perturbations in local grid spacing or salinity. The relative  
525 effect of the input variables, observed in the model sensitivity test closely matches what we saw  
526 in the different model training examples where we selectively held out these features. This again  
527 confirms that in order to train a deep learning model to make physically meaningful and generalis-  
528 able predictions of surface currents it is not sufficient to simply provide it snapshots of dynamical  
529 variables like SSH as images. We also need to provide spatial information like latitude for the  
530 NN's to effectively "learn" the physics of surface currents.

## 531 **6. Summary and Future Directions**

532 The goal of this study was to use machine learning to make predictions of ocean surface currents  
533 from satellite observable quantities like SSH wind stress, SST etc. Our central question was: Can

534 we train deep learning based models to learn physical models of estimating surface currents like  
535 geostrophy, Ekman flow and perhaps do better than the physical models themselves?

536 We used the output from the CESM POP model surface fields as our “truth” data for this study.  
537 As a first order example, we tested a linear regression model for a few of local subdomains ex-  
538 tracted from the global GCM output. Linear regression works well only when the domains are  
539 small and far removed from the equator and gets progressively worse as the domain gets bigger  
540 and the variation in local coriolis parameter gets large. It performs most poorly when  $f$  changes  
541 sign in the domain. reasonably well for small enough regions, far enough from the equator. This  
542 showed that unsurprisingly it is not possible to train a simple linear model to accurately predict  
543 surface currents. In addition, providing more data does not necessarily improve the predictive  
544 ability of a linear model and only made it worse as it starts overfitting. Whereas for the same  
545 kind of domain, a neural network we can minimize the loss (MSE) with fewer data-points and still  
546 remain generalisable, since neural networks can learn functional relationships between regressors  
547 (input features) with only a small amount of data. The model’s ability to make predictions is also  
548 shown to improve with more data. Furthermore, compared to a linear regression model, a NN  
549 even with a relatively small network of densely connected nodes, with a suitable non-linear acti-  
550 vation function (like ReLU), allows us to have a large number of trainable parameters (weights,  
551 biases) that can be optimized to minimize the loss. The activation function is what allows the  
552 different non-linear combinations between the different regressors (input features). Furthermore,  
553 a neural network trained on the entire globe is shown to predict surface currents more accurately  
554 in the sub-domains than neural networks trained in those specific sub domains. In comparison,  
555 a similar approach with a linear regression model produces the opposite result, *i.e.* a globally  
556 trained linear regression model produces higher prediction errors than the one that’s trained on  
557 each specific sub-domain. The fact that spatially diverse data actually makes the neural network

558 perform better is indicative of the fact that a neural network can actually “learn” the functional  
559 relationships needed for calculating surface currents rather than simply memorizing some target  
560 values for different combination of input features. By examining the dependence of the NNs on  
561 the choice of input features and by looking at the sensitivity of a NN model to perturbations in  
562 the input features, we established that apart from SSH, the physical location of the input features  
563 is one of the most crucial elements for the NN to “learn” the physics of surface currents. It is  
564 further demonstrated that with a careful and deliberate choice of input features the neural network  
565 can even beat the physics-based models at predicting surface currents accurately in most regions  
566 of the global ocean. A key ingredient for calculating the Ekman part of the flow using current  
567 physics based models is the vertical diffusivity, which is largely unknown for most of the global  
568 ocean. Most observational ocean current estimates that include the Ekman part of the flow relies  
569 on inferring the vertical diffusivity based on empirical multiple linear regressions with Lagrangian  
570 surface drifter data, The neural network approach, by comparison does not suffer from the same  
571 kind of limitation, since in this framework, we do not need to provide  $A_z$  as an input feature, which  
572 is one more added advantage for this method.

573 In this study, we wanted to see whether we can train a statistical model like a NN with data to  
574 essentially match or perhaps beat the baseline physics based models we currently use to estimate  
575 surface currents. By examining the errors in surface current predictions from our NN predictions  
576 and comparing them with predictions from physically motivated models (like geostrophy and Ek-  
577 man dynamics), we showed that a relatively simple NN captures most of the large scale flow  
578 features equally well if not better than the physical models, with only one day of training data for  
579 the globe.

580 However, some key aspects of the flow, associated with mesoscale and sub-mesoscale turbulence  
581 are not reproduced. We speculate that this is possibly caused by the fact that the neural network

582 framework can not capture the higher order balances (gradient wind) that are likely at play in these  
583 regions since these hotspots of high errors are collocated with regions of High Ro where balance  
584 breaks down (see Figs. 4-6).

585 One of the biggest hurdles associated with these studies is figuring out efficient strategies to  
586 stream large volumes of earth system model data into a NN framework. So before diving headfirst  
587 into the highest resolution global ocean model (currently available), we wanted to test the feasi-  
588 bility of using a regression model based on deep learning as a framework for estimating surface  
589 currents with a lower resolution model data (smaller/more managable dataset), while still being  
590 eddy resolving. Hence we chose the CESM POP model data for this present study. In the future,  
591 we propose to train a NN with data from a higher spatio-temporal resolution global ocean model  
592 like the MITgcm llc4320 model (Rocha et al. 2016). As a further step, we could coarse-grain  
593 such a model to SWOT-like resolutions, or use the SWOT simulator, train NNs on that, and make  
594 predictions for global surface currents.

595 As for the weak surface currents predicted by our NN at the equator, we need to keep in mind  
596 that geostrophic balance (defined by the first order dervatives of SSH) only holds away from the  
597 equator and satellite altimetry datasets (e.g. AVISO, Ducet et al. 2000) typically employ a higher  
598 order balance (Lagerloef et al. 1999) at the equator, to match the flow regime with the geostrophic  
599 regime away from the equator. One possible way to train the NN to learn these higher order  
600 balances would be by increasing the stencil size around each point. Since our primary goal was  
601 for the model to learn geostrophy, we started with a spatial stncil in SSH. We also explored training  
602 approaches where we provided stencils in SST, wind and SSS, with the intention of helping the  
603 model learn about wind-stress curl and thermal wind balance. In practice, however these didn't  
604 payoff as much and these additional stencils did not significantly improve model performance. In

605 future approaches we can try to provide separate stencils of varying size to each of these input  
606 variables, to test whether we can further improve the model accuracy.

607 As another future step, we also aim to incorporate recursive neural networks (RNNs) in con-  
608 junction with convolutional filters of varying kernel sizes, to train the models on cyclostrophic or  
609 gradient wind balance. This recursive neural network approach would be analogous to iteratively  
610 solving the gradient wind equation (Knox and Ohmann 2006), a technique which was originally  
611 developed for numerical weather prediction before advances in computing allowed for integrating  
612 the full non-linear equations.

613 The present work demonstrates that to a large extent, a simple neural network can be trained  
614 to extract functional relationships between SSH, wind stress etc. and surface currents with quite  
615 limited data. The field of deep learning as of now is rapidly evolving. It remains to be seen, if with  
616 some clever choices of training strategies and by using some of the other more recently developed  
617 deep learning techniques, we can improve upon this. In this study, we propose a few approaches  
618 that can be implemented to improve upon our current results and would like to investigate this in  
619 further detail in future studies. In addition, we believe that data driven approaches, like the one  
620 shown in this present study, have strong potential applications for various practical problems in  
621 physical oceanography, and require further exploration. Insights gained from this type of analy-  
622 sis could be of great potential significance, especially for future satellite altimetry missions like  
623 SWOT.

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## 624 **References**

- 625 Arbic, B. K., J. G. Richman, J. F. Shriver, P. G. Timko, E. J. Metzger, and A. J. Wallcraft,  
626 2012: Global modeling of internal tides: Within an eddy ocean general circulation model.  
627 *Oceanography*, **25** (2), 20–29, URL <http://www.jstor.org/stable/24861340>.
- 628 Bar-Sinai, Y., S. Hoyer, J. Hickey, and M. P. Brenner, 2018: Data-driven discretization: a method  
629 for systematic coarse graining of partial differential equations. *ArXiv e-prints*, 1808.04930.
- 630 Bolton, T., and L. Zanna, 2018: Applications of deep learning to ocean data inference and sub-grid  
631 parameterisation. EarthArXiv, URL [eartharxiv.org/t8uhk](http://eartharxiv.org/t8uhk), doi:10.31223/osf.io/t8uhk.
- 632 Bonjean, F., and G. S. E. Lagerloef, 2002: Diagnostic model and analysis of  
633 the surface currents in the tropical pacific ocean. *Journal of Physical Oceanog-*  
634 *raphy*, **32** (10), 2938–2954, doi:10.1175/1520-0485(2002)032<2938:DMAAOT>2.0.CO;2,  
635 URL [https://doi.org/10.1175/1520-0485\(2002\)032<2938:DMAAOT>2.0.CO;2](https://doi.org/10.1175/1520-0485(2002)032<2938:DMAAOT>2.0.CO;2), [https://doi.org/](https://doi.org/10.1175/1520-0485(2002)032<2938:DMAAOT>2.0.CO;2)  
636 [10.1175/1520-0485\(2002\)032<2938:DMAAOT>2.0.CO;2](https://doi.org/10.1175/1520-0485(2002)032<2938:DMAAOT>2.0.CO;2).
- 637 Chapman, C., and A. A. Charantonis, 2017: Reconstruction of subsurface velocities from satel-  
638 lite observations using iterative self-organizing maps. *IEEE Geoscience and Remote Sensing*  
639 *Letters*, **14** (5), 617–620, doi:10.1109/LGRS.2017.2665603.
- 640 Chollet, F., and Coauthors, 2015: Keras. <https://keras.io>.
- 641 Ducet, N., P. Y. Le Traon, and G. Reverdin, 2000: Global high-resolution map-  
642 ping of ocean circulation from topex/poseidon and ers-1 and -2. *Journal of Geo-*

643 *physical Research: Oceans*, **105 (C8)**, 19 477–19 498, doi:10.1029/2000JC900063,  
644 URL <https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2000JC900063>, <https://agupubs.onlinelibrary.wiley.com/doi/pdf/10.1029/2000JC900063>.  
645

646 Gentine, P., M. Pritchard, S. Rasp, G. Reinaudi, and G. Yacalis, 2018: Could ma-  
647 chine learning break the convection parameterization deadlock? *Geophysical Re-*  
648 *search Letters*, **45 (11)**, 5742–5751, doi:10.1029/2018GL078202, URL <https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2018GL078202>, <https://agupubs.onlinelibrary.wiley.com/doi/pdf/10.1029/2018GL078202>.  
649  
650

651 Goodfellow, I., Y. Bengio, A. Courville, and Y. Bengio, 2016: *Deep learning*, Vol. 1. MIT press  
652 Cambridge.

653 Gregor, L., S. Kok, and P. M. S. Monteiro, 2017: Empirical methods for the estimation of southern  
654 ocean CO<sub>2</sub>: support vector and random forest regression. *Biogeosciences*, **14 (23)**, 5551–5569,  
655 doi:10.5194/bg-14-5551-2017, URL <https://www.biogeosciences.net/14/5551/2017/>.

656 Kingma, D. P., and J. Ba, 2014: Adam: A method for stochastic optimization. *CoRR*,  
657 **abs/1412.6980**, URL <http://arxiv.org/abs/1412.6980>, 1412.6980.

658 Kingma, D. P., and J. Ba, 2017: Adam: A method for stochastic optimization. 1412.6980.

659 Knox, J. A., and P. R. Ohmann, 2006: Iterative solutions of the gradient wind equation. *Computers*  
660 *Geosciences*, **32 (5)**, 656–662, doi:<https://doi.org/10.1016/j.cageo.2005.09.009>, URL <https://www.sciencedirect.com/science/article/pii/S0098300405001974>.  
661

662 Lagerloef, G. S. E., G. T. Mitchum, R. B. Lukas, and P. P. Niiler, 1999: Tropical pa-  
663 cific near-surface currents estimated from altimeter, wind, and drifter data. *Journal of*  
664 *Geophysical Research: Oceans*, **104 (C10)**, 23 313–23 326, doi:10.1029/1999JC900197,



665 URL <https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/1999JC900197>, <https://agupubs.onlinelibrary.wiley.com/doi/pdf/10.1029/1999JC900197>.

666

667 LeCun, Y., Y. Bengio, and G. Hinton, 2015: Deep learning. *nature*, **521 (7553)**, 436.

668 LeTraon, P., and R. Morrow, 2001: Ocean currents and eddies. **69**, 171–215.

669 Lguensat, R., M. Sun, R. Fablet, E. Mason, P. Tandeo, and G. Chen, 2017: Eddynet: A deep  
670 neural network for pixel-wise classification of oceanic eddies. *CoRR*, **abs/1711.03954**, URL  
671 <http://arxiv.org/abs/1711.03954>, 1711.03954.

672 Ling, J., A. Kurzawski, and J. Templeton, 2016: Reynolds averaged turbulence modelling using  
673 deep neural networks with embedded invariance. *Journal of Fluid Mechanics*, **807**, 155–166,  
674 doi:10.1017/jfm.2016.615.

675 Morrow, R., D. Blurmstein, and G. Dibarboure, 2018: Fine-scale altimetry and the future swot  
676 mission. *New Frontiers in Operational Oceanography*, 191–226.

677 Nielsen, M. A., 2015: *Neural networks and deep learning*, Vol. 25. Determination press USA.

678 Pathak, J., B. Hunt, M. Girvan, Z. Lu, and E. Ott, 2018: Model-free prediction of large  
679 spatiotemporally chaotic systems from data: A reservoir computing approach. *Phys. Rev.  
680 Lett.*, **120**, 024 102, doi:10.1103/PhysRevLett.120.024102, URL [https://link.aps.org/doi/10.  
681 1103/PhysRevLett.120.024102](https://link.aps.org/doi/10.1103/PhysRevLett.120.024102).

682 Rocha, C. B., T. K. Chereskin, S. T. Gille, and D. Menemenlis, 2016: Mesoscale to submesoscale  
683 wavenumber spectra in drake passage. *Journal of Physical Oceanography*, **46 (2)**, 601–620.

684 Sinha, A., 2019: Temporal variability in ocean mesoscale and submesoscale turbulence. Ph.D.  
685 thesis, Columbia University in the City of New York, doi:<https://doi.org/10.7916/d8-bngk-r215>,  
686 URL <https://academiccommons.columbia.edu/doi/10.7916/d8-bngk-r215>.

687 Small, R. J., and Coauthors, 2014: A new synoptic scale resolving global climate simula-  
688 tion using the community earth system model. *Journal of Advances in Modeling Earth Sys-*  
689 *tems*, **6** (4), 1065–1094, doi:<https://doi.org/10.1002/2014MS000363>, URL [https://agupubs.](https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1002/2014MS000363)  
690 [onlinelibrary.wiley.com/doi/abs/10.1002/2014MS000363](https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1002/2014MS000363), [https://agupubs.onlinelibrary.wiley.](https://agupubs.onlinelibrary.wiley.com/doi/pdf/10.1002/2014MS000363)  
691 [com/doi/pdf/10.1002/2014MS000363](https://agupubs.onlinelibrary.wiley.com/doi/pdf/10.1002/2014MS000363).

692 Smith, R., and Coauthors, 2010: The parallel ocean program (pop) reference manual ocean com-  
693 ponent of the community climate system model (ccsm) and community earth system model  
694 (cesm). *Rep. LAUR-01853*, **141**, 1–140.

695 Sudre, J., C. Maes, and V. Garcon, 2013: On the global estimates of geostrophic  
696 and ekman surface currents. *Limnology and Oceanography: Fluids and Environments*,  
697 **3** (1), 1–20, doi:[10.1215/21573689-2071927](https://doi.org/10.1215/21573689-2071927), URL [https://aslopubs.onlinelibrary.wiley.com/](https://aslopubs.onlinelibrary.wiley.com/doi/abs/10.1215/21573689-2071927)  
698 [doi/abs/10.1215/21573689-2071927](https://aslopubs.onlinelibrary.wiley.com/doi/abs/10.1215/21573689-2071927), [https://aslopubs.onlinelibrary.wiley.com/doi/pdf/10.1215/](https://aslopubs.onlinelibrary.wiley.com/doi/pdf/10.1215/21573689-2071927)  
699 [21573689-2071927](https://aslopubs.onlinelibrary.wiley.com/doi/pdf/10.1215/21573689-2071927).

700 Sudre, J., and R. A. Morrow, 2008: Global surface currents: a high-resolution product for inves-  
701 tigating ocean dynamics. *Ocean Dynamics*, **58** (2), 101, doi:[10.1007/s10236-008-0134-9](https://doi.org/10.1007/s10236-008-0134-9), URL  
702 <https://doi.org/10.1007/s10236-008-0134-9>.

703 Uchida, T., R. Abernathey, and S. Smith, 2017: Seasonality of eddy kinetic energy in  
704 an eddy permitting global climate model. *Ocean Modelling*, **118**, 41 – 58, doi:[https://](https://doi.org/10.1016/j.ocemod.2017.08.006)  
705 [doi.org/10.1016/j.ocemod.2017.08.006](https://doi.org/10.1016/j.ocemod.2017.08.006), URL [http://www.sciencedirect.com/science/article/pii/](http://www.sciencedirect.com/science/article/pii/S1463500317301221)  
706 [S1463500317301221](http://www.sciencedirect.com/science/article/pii/S1463500317301221).

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712 TABLE 1. Table summarizing model errors from the the physics based model (geostrophy + Ekman flow) and  
 713 the two types of regression models - linear regression and neural network (Panel (a) and (b) in Fig. 2).

Model(training region)	Number of trainable parameters	Epochs	MAE (train)[ $cm/s$ ]	MAE (eval)GS [ $cm/s$ ]	MAE (eval)Kuroshio [ $cm/s$ ]	MAE (eval)ACC [ $cm/s$ ]
LR (Gulf Stream)	38	8	10.7	11.4	-	-
NN (GS)	1812	8	2.3	3.7	-	-
LR (Kuroshio)	38	5	12.9	-	13.4	-
NN (Kuroshio)	1812	5	5.8	-	7.0	-
LR (ACC)	38	5	7.5	-	-	-
NN (ACC)	1812	5	1.9	-	-	-
NN (global)	1812	4	3.0	2.4	5.1	-
<i>geo + Ek</i> (global)	-	-	-	6.1	29.2	-

TABLE 2. Table summarizing the different CNNs and the training strategies explored

Model No.	Stencil inspace (2s)	Stencil intime (3t)	Stencil Variables	Point Variables	Number of trainable parameters
1	✓	×	$\eta$	$\tau_x, \tau_y$	4772
2	✓	×	$\eta$	$X (= \frac{f}{2\Omega})$	4732
3	✓	×	$\eta, \theta$	$X$	5052
4	✓	×	$\eta$	$X, \tau_x, \tau_y$	4812
5	✓	×	$\eta, \theta$	$X, \tau_x, \tau_y$	5132
6	✓	×	$\eta, X, Y, Z$	$\tau_x, \tau_y$	5732
7	✓	×	$\eta, X$	$\tau_x, \tau_y$	5092
8	✓	×	$\eta, \theta, X, Y, Z$	$\tau_x, \tau_y$	6052
9	✓	×	$\eta, \theta, X$	$\tau_x, \tau_y$	5412
10	✓	×	$\eta, \tau_x, \tau_y$	$X$	5372
11	✓	×	$\eta, \theta, \tau_x, \tau_y$	$X$	5692
12	✓	×	$\eta, \theta$	$X, dx, dy, \tau_x, \tau_y$	5212
13	✓	×	$\eta, \theta, S$	$X, dx, dy, \tau_x, \tau_y$	5532
1t	✓	✓	$\eta$	$\tau_x, \tau_y, X, dx, dy$	5532
2t	✓	✓	$\eta, \theta$	$\tau_x, \tau_y, X, dx, dy$	6492
3t	✓	✓	$\eta$	$\tau_x, \tau_y, X$	5452
4t	✓	✓	$\eta, \theta$	$\tau_x, \tau_y, X$	6412
5t	✓	✓	$\eta$	$X, dx, dy$	5452
6t	✓	✓	$\eta, \theta$	$X, dx, dy$	6412
7t	✓	✓	$\eta$	$X$	5372
8t	✓	✓	$\eta, \theta$	$X$	6332
9t	✓	✓	$\eta, \theta, S$	$\tau_x, \tau_y, X$	7372
10t	✓	✓	$\eta, \theta, S$	$\tau_x, \tau_y, X, dx, dy$	7452

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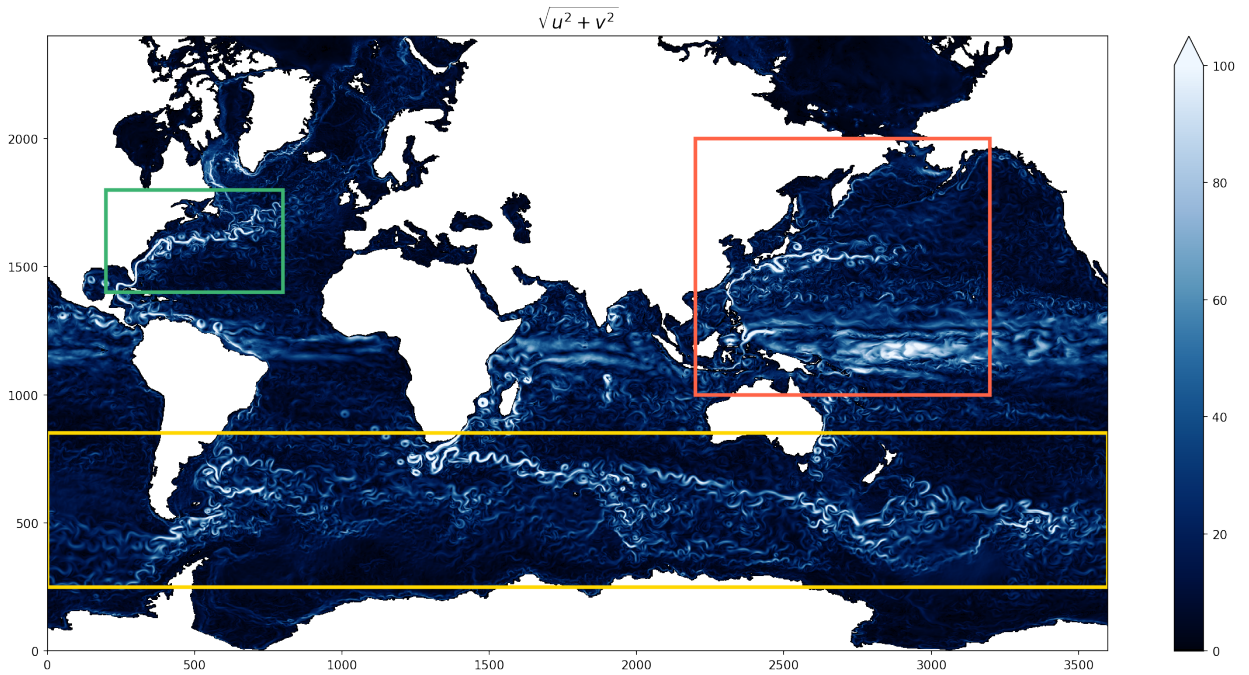
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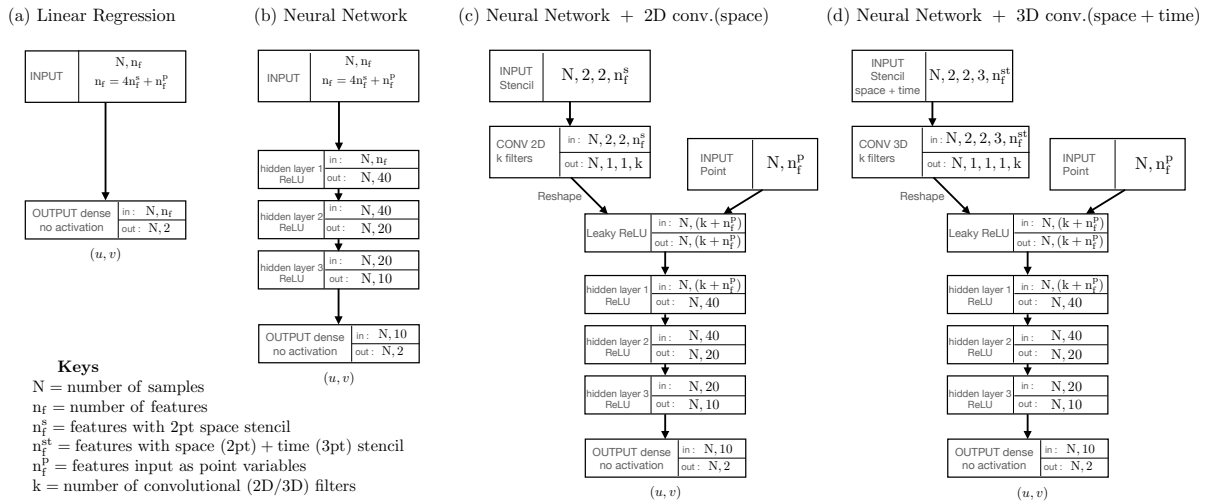
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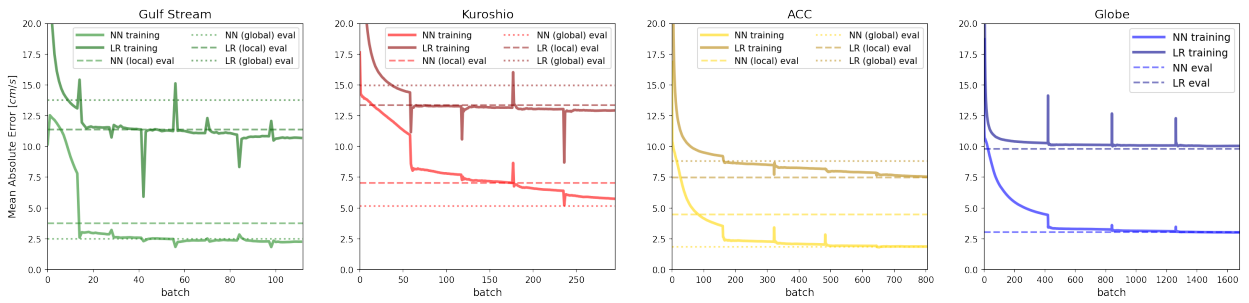


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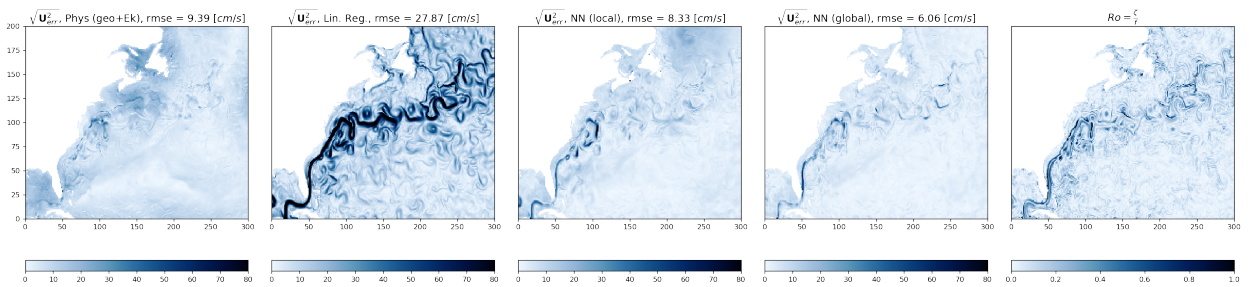




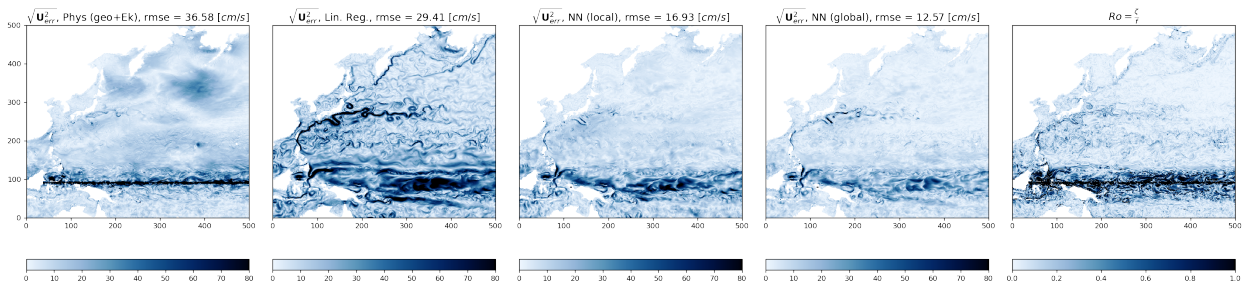
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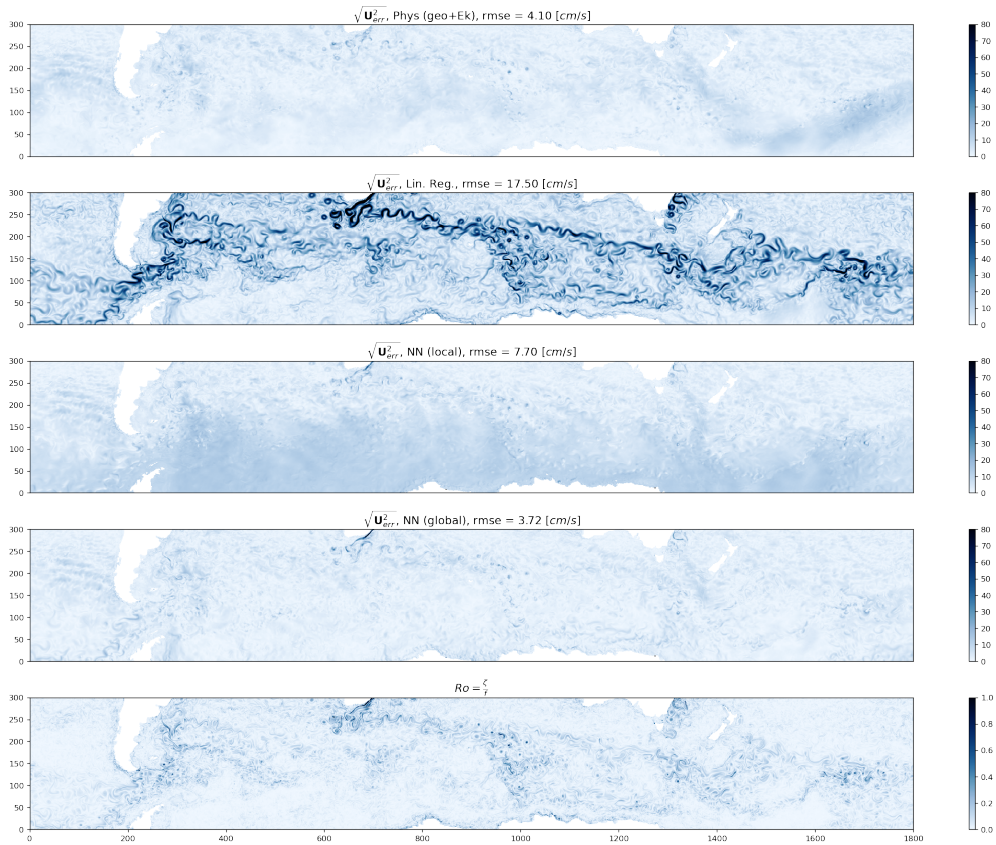
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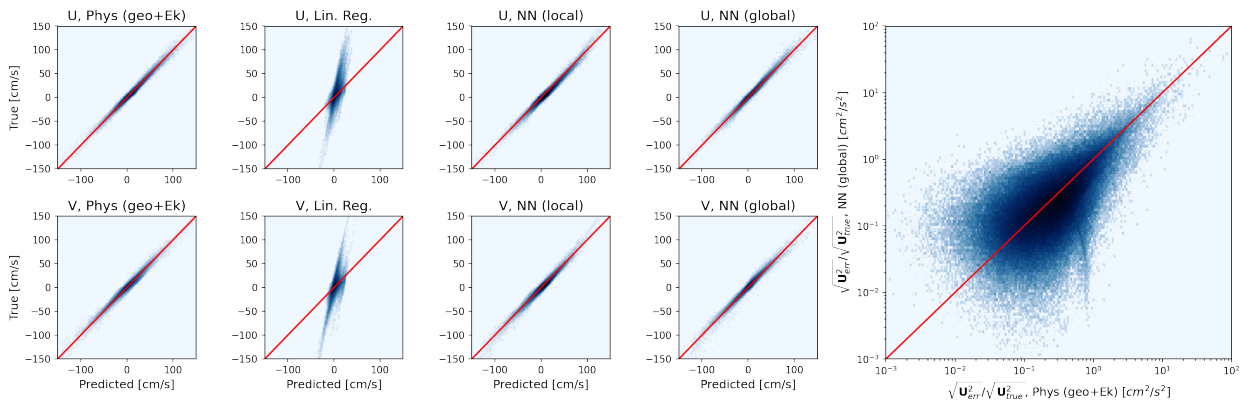
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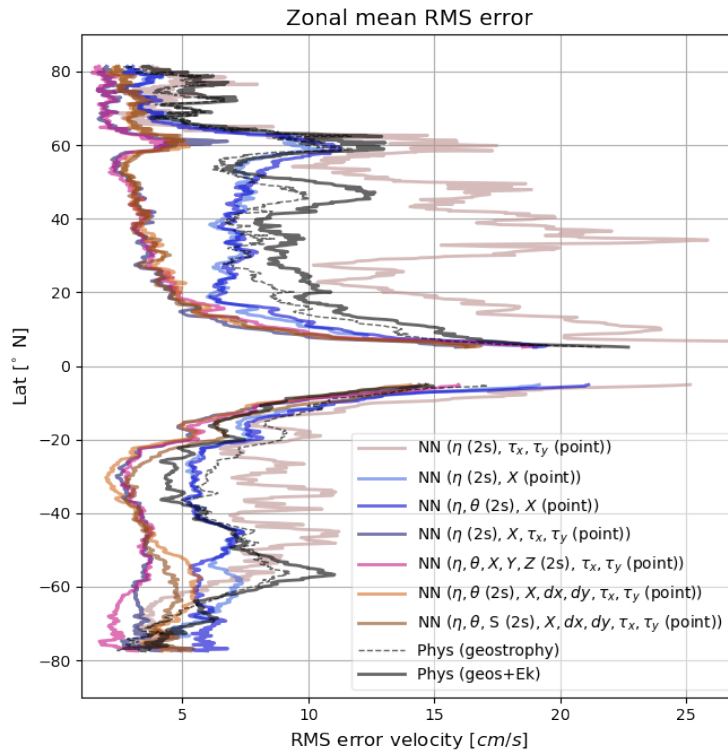
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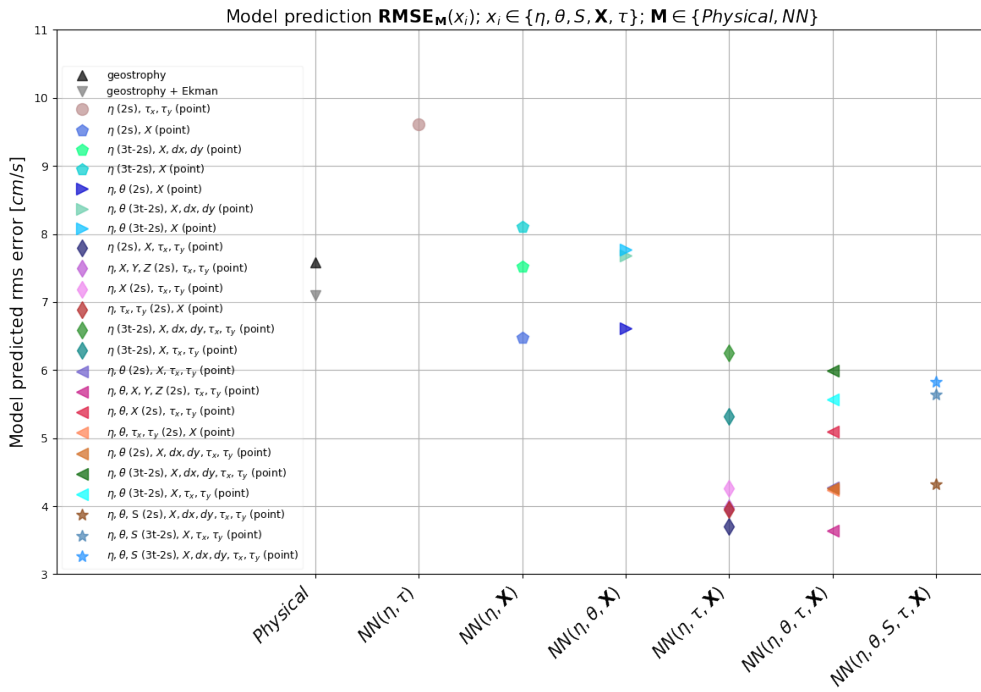
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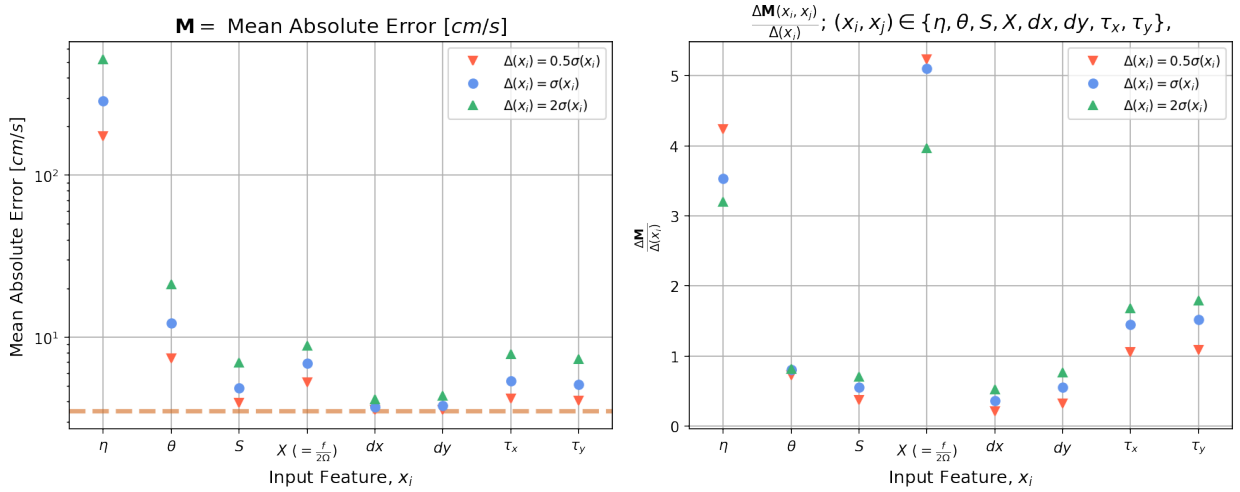


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