

Predicting the yield stress of a 3D printed porous material from its internal structure

Martin Lesueur^{1,*+}, Thomas Poulet^{2,+}, and Manolis Veveakis^{1,+}

¹Duke University, Civil and Environmental Engineering, Durham, 27708, USA

²CSIRO, Mineral Resources, Perth, 6151, Australia

*martin.lesueur@duke.edu

+these authors contributed equally to this work

ABSTRACT

The design of any engineering structure requires the knowledge, and therefore determination, of the yield, i.e. limit of elasticity, for the building material. Whilst destructive experimental testing is currently necessary to do so, our work is part of initiatives which aim at deriving the yield without such laboratory experiments. The seminal work of Gurson (1977) on a simplified pore structure, a single spherical pore, first provided a theoretical relationship between the yield stress and the porosity. Specifically, it showed that the presence of pore space is responsible for the existence of a compression cap in plasticity, in addition to lowering the yield stress. This contribution extends the approach to determine the macroscopic yield of a porous material by taking explicitly into account its internal structure. As the yielding of a porous material is controlled by the geometry of its internal structure, we postulate that it is nearly independent of the constitutive plastic behaviour of the material. Here, we show that the influence of that internal structure on the yield could be retrieved from a finite element computation with just an elasto-plastic ideal material equivalent of the skeleton's. With some basic knowledge about the skeleton's mechanical properties, this process allows the determination of the yield stress without requiring the experimental compression of the material. We showcase the predictive power of the method against experimental testing, initially for a unique spherical void in a 3D-printed cylinder sample following Gurson, before demonstrating its applicability on a complex 3D-printed rock microstructure, reconstructed from segmented micro-Computerised Tomography scans.

Introduction

Many studies are aiming at accounting for the influence of the material's internal structure on its mechanical properties, whether on elastic modulus^{1,2}, strength³⁻⁶, or plastic flow law^{5,6}, to cite only a few. This contribution focuses on the yield value, which is important for two reasons. By definition, it points to the limit of elasticity, which is necessary for the design of structures to prevent them from entering plastic regime and suffer irreversible deformations. On top of this, when the plastic regime is expected, knowledge of the yield surface is necessary to model a material's plastic behaviour⁷.

The only unambiguous determination of mechanical yield point, as a limit of elasticity, is possibly restricted to the simplest case of ideal non-porous linear elastic and ideally plastic materials, like metals for instance. Indeed, experimental compression tests of such materials lead to characteristic stress-strain curves displaying a sharp transition between the linear elasticity and plasticity, where strain increases at constant stress. For more complex materials, however, including real geomaterials like porous rocks, the notion of macroscopic yield stress is more ambiguous and its determination dependent on the method selected. To alleviate this ambiguity we use the following three definitions of yield from the sixth edition of the McGraw-Hill Dictionary of Scientific and Technical Terms⁸:

- **yield** [MECHANICS] *That stress in a material at which plastic deformation occurs.*
- **yield point** [MECHANICS] *The lowest stress at which strain increases without increase in stress.*
- **yield strength** [MECHANICS] *The stress at which a material exhibits a specified deviation from proportionality of stress and strain.*

The first definition, referred to as **initial yield** in this contribution corresponds to the stress when the first region in the material undergoes plasticity. This value is not particularly useful since not easily measurable⁹. The second definition, commonly named **limit load**, points to the state of collapse of the material. The last definition is the **macroscopic yield**, which is the focus in this contribution because it really points to the limit of linear elasticity at the scale of the sample, necessary for structure design. Additionally, it is a necessary parameter for any modelling of plasticity. This yield is classically measured experimentally on stress-strain curves using the classical offset method¹⁰, as the intersection of the curve with a line parallel to

33 the initial linear-elastic part of that curve, shifted by an ad-hoc strain threshold. In this contribution, following Lesueur *et al.*¹¹,
34 the macroscopic yield is measured on stress-strain curves with an energetic method, which provides similar values but with
35 stronger physical meaning. It is necessary to try and reduce the use of destructive methods for its measurement, particularly
36 when material samples are difficult to obtain in sufficient quantity, which is the case in subsurface operations like petroleum
37 engineering for instance. Two alternatives are possible.

38 The first one is to find relationships to link the yield to properties that can be measured with non-destructive tests. Some
39 models already account for the impact on yield from the simplest parameter characterising the internal structure: porosity^{5,6,12}.
40 Through these models, we know that the presence of pore space is responsible for the existence of a compression cap of
41 plasticity, in addition to lowering the yield stress value. We focus specifically in this contribution on Gurson⁵, who developed a
42 criterion to predict the limit load of a material from its porosity. He considered the simplest configuration, which consists of a
43 single spherical pore. To avoid any non-essential complexity and restrict the study to the influence of porosity, he selected a
44 rigid and ideally plastic material, which ensured that the relationships extracted were indeed be attributed solely to the porosity
45 without any interfering influence of any other material properties. The outer geometry of the unit cell was taken as spherical, like
46 the void shape, in order to retain the geometrical isotropy benefits from the symmetry in the analysis. Using the upper-bound
47 limit analysis method, Gurson obtained an approximate upper solution to the limit load of the hollow sphere geometry, which
48 proved to be precise enough to fit experimental data^{13,14}. Extending his work, numerous studies followed which improved on
49 the model, including derivations accounting for other shapes of voids (e.g. elliptical¹⁵), the interaction between voids^{16,17},
50 or the consideration of more complex matrix materials (e.g. viscoplasticity¹⁸). That type of analysis was an important first
51 step in our understanding of the influence of the internal structure. Yet, because of its use of limit analysis, its applicability is
52 restricted to the limit load value and corresponding results have not yet been derived for the more practical definition of yield,
53 the macroscopic yield. In addition, porosity, as a scalar field, only represents one of the characteristics¹⁹ and can therefore not
54 capture all geometrical effects, with more work remaining from a more general perspective. Specifically in this contribution,
55 we define the internal structure as the structure of the porous material's skeleton. In the case of 3D printed porous scaffolds, it
56 is directly the internal geometry of the unit cell. In the case of natural materials bones or rocks, it is delimited in segmented
57 micro-Computerised Tomography (μ CT) scans by the pore space boundary.

58 The second alternative is to numerically obtain the stress-strain curve. With recent computational advances, it is now
59 possible to simulate mechanical deformation of a Representative Element Volume (REV) of the material (e.g.²⁰). At that
60 size, the mechanical behaviour of the volume considered should be representative of the whole sample at the larger scale.
61 Therefore stress-strain curves of the REV can be produced numerically and be comparable to the experimental ones. However,
62 reproducing numerical stress-strain curves of real materials remains difficult. Such simulations require indeed a high mesh
63 resolution to match the REV with accurate grain shapes. The computational cost is amplified by the fact that non-trivial
64 constitutive plastic law are usually implemented to reproduce the behaviour of the material. Indeed, characterising the plastic
65 behaviour of a real material is no easy feat as there exist numerous constitutive models²¹, some of which that require many
66 parameters to be calibrated⁷.

67 We take inspiration from both approaches for the solution we suggest in this contribution. We propose to extend Gurson's
68 type of analysis from limit load to macroscopic yield and from porosity to the explicit geometry of the internal structure. To do
69 so, we simulate mechanical compressions on the digitised internal structure and determine the yield with the energetic method
70 from the stress-strain curves computed. By narrowing our study to the determination of yield, we only need to simulate the
71 initial phase of plasticity. The plastic regime starts at the initial yield since theoretically speaking, the material is, from that
72 point on, undergoing localised plastic deformations. However, for porous materials, it is instinctive that the initial yield does
73 not coincide with the macroscopic yield. Localised heterogeneities in the material will indeed fail before the overall response of
74 the material can visually deviate from linearity⁹. Under ongoing deformation, from the initial yield to the macroscopic yield, an
75 arbitrarily small plastic strain is accumulated (as defined by the offset method). However, we conjecture that plasticity does not
76 noticeably affect the material's response until the macroscopic yield. This hypothesis is tested in this contribution by verifying
77 that the yield of a porous material is equal to the one of a virtual porous material with an equivalent ideal elasto-plastic skeleton,
78 instead of considering its more realistic plastic behaviour (including rheology).

79 The material selected in this contribution is 3D printed polylactic acid (PLA), whose mechanical response from laboratory
80 experiments is plotted in Fig. 1 and modelled in Methods. 3D printing presents great advantages for the experimental validation
81 of our approach. As observed by the superposition of curves in Fig. 1 or Fig. S4, the printed material has a very reproducible
82 behaviour. In addition, the 3D printing technique allows a perfect control of the internal structure of the samples, whose
83 influence we are characterising. 3D printed PLA is particularly well-suited to test our hypothesis because its plastic response
84 is far from ideal plastic (see Fig. 1). This material displays a strong viscoplasticity and we selected a sample size such that
85 deformation pattern is shearbanding, which results in a weakening before reaching the limit load (see Fig. 1), for added
86 complexity. Moreover, the printing process itself influences the plastic properties of the resulting material, as discussed in
87 Methods and Supplementary text, which adds an extra layer of complexity. It is therefore extremely interesting to select this

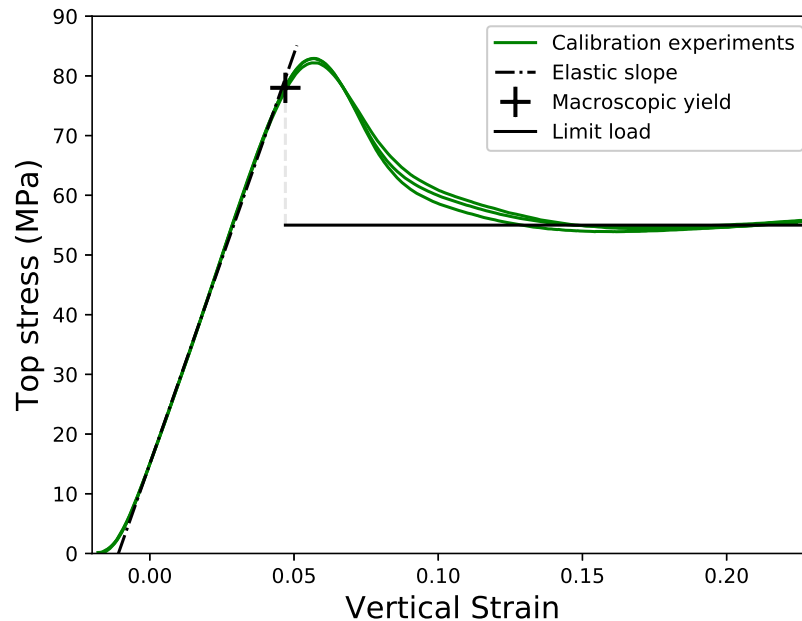


Figure 1. Three stress strain curves of uniaxial compression of 3D printed full cylinders of PLA to observe the plastic response of the material and assess the reproducibility of mechanical tests on 3D printed samples. Our suggested elasto-plastic model is superposed to the curves and determined by three parameters: the slope of the linear elastic part, the macroscopic yield value, and the limit load (displayed on a wide range of strain for visualisation purposes).

88 material to test our approach, which eliminates the need for characterisation of the viscoplasticity and weakening law of the
 89 printed PLA.

90 To validate the approach proposed, we select at first the simplest internal structure which is the unique spherical pore as
 91 Gurson studied. This simple structure allows the 3D printed internal geometry to be very accurate and improve reproducibility of
 92 this benchmark. The second part of this contribution presents an application for a more complex internal structure, reconstructed
 93 from a rock's segmented μ CT scans.

94 Results

95 Prediction of the internal structure influence on 3D-printed PLA yield

96 The objective of this section is to verify if a simplified numerical model can correctly predict the yielding of printed PLA
 97 sample with a given internal structure. The printed samples are cylinders containing a spherical void of varying diameter at
 98 their centre, specifically of 0.6 and 0.7 (normalised to the cylinder diameter). Due to the Fused Deposition Modeling (FDM)
 99 principle of printing, the molten filament is deposited vertically on the sample, which makes it impossible for this technique to
 100 print perfectly any overhanging part with an angle greater than 45° . Unfortunately, this is the case of the spherical void with the
 101 overhang going to 90° at the top of the sphere. To help the printing, FDM usually relies on printing under these overhangs
 102 some support structure that the user can remove after the print is finished. However, our overhang is fully enclosed in the
 103 structure so this technique cannot be used without any support structure and we can only assess the quality of the print visually
 104 by cutting the sample after the experiments, as shown in Fig. 2. We can see that the quality of the print remained acceptable,
 105 even though imperfect. Indeed, during the mechanical compression, this top part of the sphere is the location which experiences
 106 the minimum of stress overall.

107 The samples are subjected to uniaxial compression and the experimental results are plotted in Fig. 3 for the two different
 108 sphere diameters. Note that each test is repeated two times for reliability reasons. The good superposition of all curves shows
 109 that the results of hollow cylinders experiments are as reproducible as the full ones. The resulting curves for the porous cylinders
 110 display sequentially a hardening and a softening phase, before converging to a limit load. Note that an eventual hardening
 111 occurs artificially due to the increase of the surface area of contact, which is disregarded. All in all, the mechanical behaviours
 112 of the porous samples are very similar to the one for a full sample but with increasingly lower and faster transition to plasticity



Figure 2. Visualisation of the printing quality of the top of the spherical void of the hollow cylinder. Only the top half of the hollow cylinder was printed for visualisation purposes and the sample is displayed upside down.

113 as porosity increases.

114 In this section, we are looking at predicting the influence of the internal structure of the hollow cylinders on two characteristic
115 yield points of the printed PLA: the limit load, studied by Gurson, and as an extension of his study, the macroscopic yield. We
116 used the mechanical simulator of the Finite Element platform MOOSE²² for all numerical simulations in this contribution.
117 It solves for the momentum balance of the skeleton of the porous material. In our simulation, the skeleton's material is
118 attributed the elastic parameters measured for the printed PLA and for the plasticity, we use a $J2$ rate-independent model with
119 no hardening or softening, defined by our single parameter of interest, the yield point.

120 The evolutions of the two yields with the internal structure are simulated separately, for each yield. Each simulation is run
121 using the same model with respectively the macroscopic yield and the limit load value of Fig. 1 as the single plastic parameter
122 of the skeleton. The set of simulations is performed twice, for the two geometries of different hole diameters, meshed with
123 second order tetrahedra. The results are displayed in Fig. 3, following the layout introduced in Fig. 1. The stress value of the
124 macroscopic yield of the porous material, measured on the stress-strain curve produced by the first simulation, is reported of
125 Fig. 1 as a cross on the elastic slope measured from the simulation. The stress value of the limit load is retrieved by taking the
126 limit load value of the second simulation and is reported as a solid line in Fig. 1, regardless of the strain value at which it is
127 actually reached.

128 The comparison of the numerical and experimental results of Fig. 3, quantified in Table. 1, shows that the simulation
129 is matching closely the two properties of plasticity considered, the macroscopic yield and the limit load value, obtained
130 experimentally. Interestingly, this perfect fit, with the numerical limit load matching the stress plateau obtained experimentally
131 as shown in Fig. 3, demonstrates that the influence of the microstructure on a real material's yield can be retrieved even with an
132 idealised model without taking into account the intrinsic behaviour of the real material. This verification validates the hypothesis
133 suggested in the introduction that plasticity has little influence on the porous material's behaviour before the macroscopic yield.
134 Particularly, we showed in this section that the hardening and softening behaviour of the 3D printed PLA does not influence
135 the value of the macroscopic yield in this recognised benchmarking example. This conclusion highlights the potential of the
136 numerical approach to extract the impact of the internal structure on the yield despite an idealised modelling of the material.

137 **Application to rock microstructure**

138 The advantage of studying a single pore compared to any arbitrary internal structure is that this geometry was designed as an
139 idealised REV of rock microstructure. Gurson's theory on the hollow sphere proved to be able to fit real data^{13,14}. As a natural
140 extension, we therefore select a rock's microstructure as a more complex geometry in this section for further validation of the
141 suggested approach. The selected rock is a 0.5 mm^3 subsample of the Berea sandstone²³.

142 Using the stack of segmented 2D μCT scan images, the geometry is meshed in 3D following the methodology described
143 by Lesueur *et al.*²⁴. In order to be processed by the Ultimaker 3 machine for printing, the mesh is converted to an STL file
144 format. The sample is printed as a cube of 22 mm^3 size. The quality of the printed sample is quite remarkable in terms of
145 details, capturing very well the overall complexity of the original rock, even though the quality of the print remains imperfect,

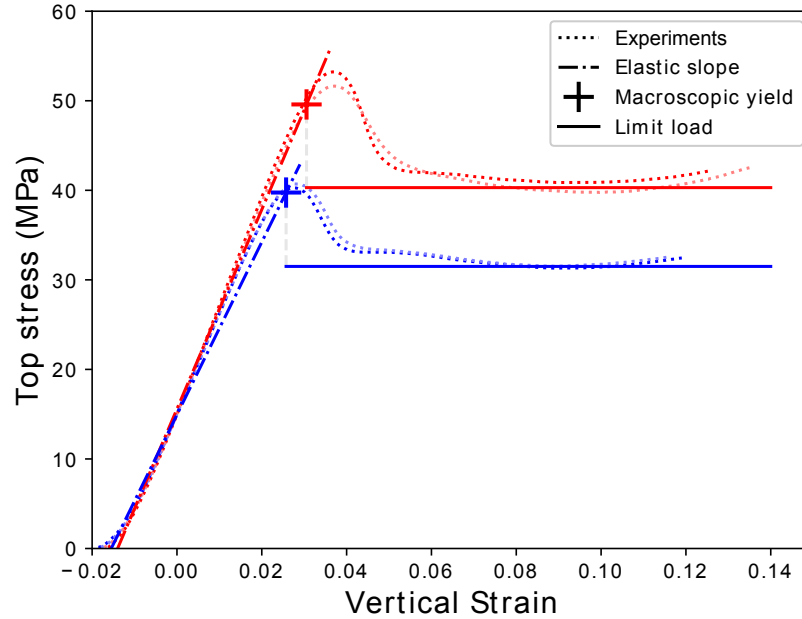


Figure 3. Four stress strain curves of uniaxial compression of 3D printed cylinders of PLA containing a spherical void of different normalised diameters: 0.6 in red and 0.7 in blue. The results of the simulation (elasticity and characteristic plastic points) using the model presented in Methods are superposed to the experimental results.

Table 1. Mechanical properties measured for the experimental and numerical results of the uniaxial compression of hollow cylinders.

Void diameter (normalised)	0.6			0.7		
	exp 1	exp 2	simulation	exp 1	exp 2	simulation
Specimen number						
Young's modulus (MPa)	1215	1181	1113	1123	1122	964
Macroscopic yield (MPa)	49.30	47.34	49.59	36.88	37.58	39.75
Limit load (MPa)	40.85	39.78	40.30	31.3	31.51	31.50

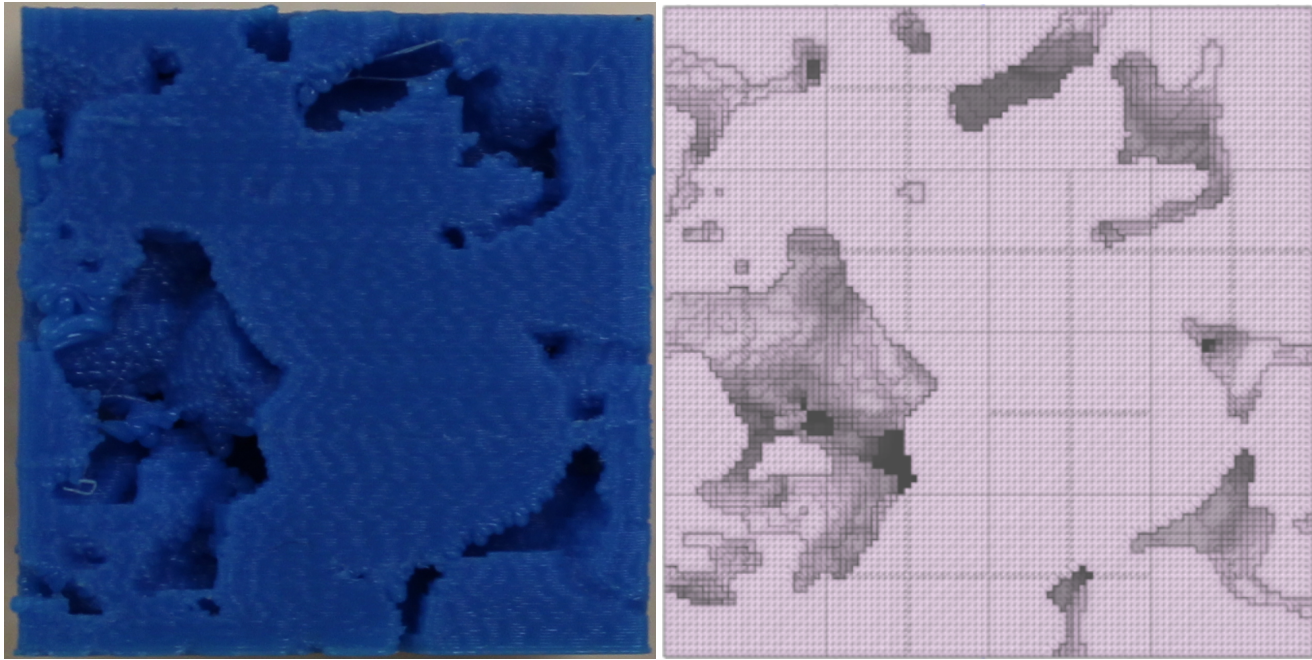


Figure 4. Side face of the printed microstructure (a) compared to the digital rock (b). The whole microstructures can be visualised as 3D figures in Supplementary Material (Fig. S5).

146 as can be seen in Fig. 4, due to the 45° limit of any overhang discussed in Methods. The printing quality can be assessed by
 147 comparing the two 3D figures (see Fig. S5 in Supplementary Material) that visualise the pore space respectively from the
 148 original μ CT scan and the 3D printed version which was μ CT scanned after being printed.

149 Five identically printed samples are then tested in uniaxial compression following the experimental procedure described in
 150 Methods, with a loading speed of 0.08 mm/min. The resulting stress strain curves, plotted in Fig. 5, all have the same general
 151 shape, including the same elastic properties and plastic hardening, but noticeably different values of macroscopic yields. We can
 152 only infer that the lack of reproducibility is due to the insufficient printing resolution and quality because the curves of Fig. 3,
 153 whose samples' printing quality was high, superposed completely. Compared to Fig. 3, the complex internal structure plays a
 154 different role than the idealised single pore: the sample shows no softening nor limit load, but instead hardens continuously. The
 155 complex pore network in the μ CT scan results in a very disperse pore collapse over the whole sample (see plastic deformations
 156 in Fig. 6) that could prevent therefore a homogeneous shearband to form, which would explain the absence of softening.

157 In order to numerically determine the yield of this example, we simulate the same compression on a digital version of that
 158 same microstructure, reconstructed from μ CT scans and meshed following the method of Lesueur *et al.*²⁴. In order to retrieve
 159 exclusively the influence of the internal structure on the yield, we take an ideal J_2 model for the plasticity of the matrix materia.
 160 We calibrate this model to resemble the behaviour of the printed PLA, on the response of 3D printed PLA cubes of 22 mm³ size.
 161 Eventually the elasto-plastic model selected is an isotropic linear elasticity with a Young's modulus of 875 MPa and a Poisson
 162 ratio of 0.45 and a yield point of 70 MPa, calibrated on the yield measured. The resulting stress strain curve is plotted in Fig. 5.

163 We note that for both numerical and experimental approach, we do not reach indeed a limit load like in Fig. 3. This confirms
 164 our interest in extending Gurson's exclusive study of the limit load to the macroscopic yield because the limit load does not
 165 exist for every material.

166 Despite the fact that the match of Fig. 5 is not as impressive as the one from Fig. 3, the numerical and experimental curves
 167 still match qualitatively and display a similar shape. In this more complex example, the porous material appears to be stiffer
 168 and stronger (higher macroscopic yield) with the experimental approach. This could be explained by the reinforcement of the
 169 structure due to the existence of artificial bridges between pores that were created during the imperfect printing process. The
 170 suboptimal printing quality adds to the uncertainty of the experimental results, which brings us more confidence in the value of
 171 elasticity and macroscopic yield determined with the numerical approach.

172 Discussion

173 In this contribution we presented an approach to determine the macroscopic yield of a porous material from finite element
 174 compression of its internal structure, replacing the traditional destructive testing approach. By focusing the study on the

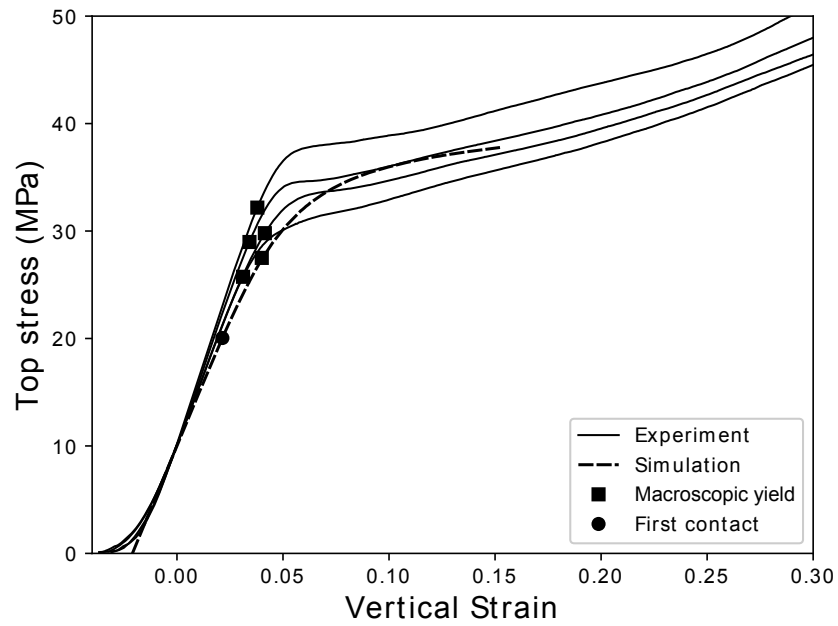


Figure 5. Experimental and numerical stress-strain curves of the uniaxial compression of 3D printed samples of the Berea sandstone²³.

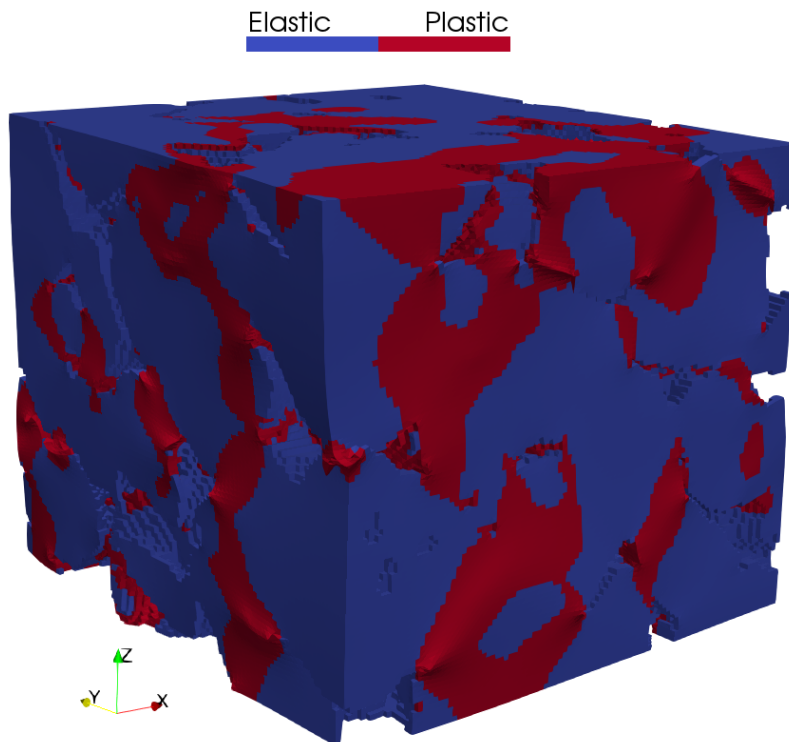


Figure 6. Visualisation of plastic deformations on the numerical uniaxial compression of Fig. 5 at 12% strain.

175 macroscopic yield instead of the full mechanical behaviour, we have shown that the complex skeleton material can be
176 satisfactorily approximated by an equivalent ideal elasto-plastic material before reaching the macroscopic yield. By reducing
177 the complexity of the material implemented, simulations of mechanical compressions become more accessible.

178 The new approach was validated on 3D printed PLA. The homogeneity of this material and the reproducibility of the 3D
179 printing techniques makes it a material very suitable for our approach. Furthermore, we show in the Appendix that the plastic
180 behaviour of the material presents a difficult calibration as it depends on many printing parameters. This justifies the use of our
181 approach which disregards this exact plastic behaviour.

182 To be able to generalise the applicability of the method, we tested the approach against both a simple and complex internal
183 structures. In order to refine the benchmarking of complex internal structure, more suitable 3D printing techniques could
184 be used in order to obtain a quality high enough to obtain perfectly reproducible and therefore trustful results. Indeed, 3D
185 printing μ CT scans with high resolution was achieved for instance by Ishutov *et al.*²⁵.

186 Since the method was validated against an already challenging material, 3D printed PLA, that displays softening and
187 hardening behaviour, we expect the approach to apply also for a wider range of materials such as geomaterials. To improve the
188 accuracy of our approach in this case, contact mechanics could be implemented as we have shown in Fig. 5 that contacts in a
189 real rock microstructure happen early on the stress-strain curve. It is however unclear if this would affect the macroscopic yield.
190 In any case, contact forces have been known to be responsible for the pressure sensitivity of the yield surface at the macro-scale,
191 modelled commonly with a Drucker-Prager which characterises a wide array of geomaterials.

192 In summary, this study aimed at highlighting the predictive power scientific community can develop as 3D printing
193 technology is maturing, at a level of quality where reproducible mechanical compression experiments on 3D printed samples
194 can be performed. We showed that the macroscopic yield can be obtained for a given internal structure from 3D printed
195 reproductions, for high enough resolutions. More importantly, it was shown that it can also be predicted numerically, in a
196 non-destructive manner, using the simplest plasticity model for the actual filler material. This result has striking repercussions
197 for a number of applications, including 3D printed scaffolds, or even more for real materials like bones or geomaterials, whose
198 internal structure can be obtained from μ CT scans. Given that 3D printing and numerical simulations are approaching their
199 originally anticipated goal of providing invaluable insight to the mechanical properties of natural materials, studies like the
200 present one are aiming at opening the door to an enhanced material design era.

201 **Methods**

202 **3D printing and mechanical testing procedure**

203 In recent years, many 3D printing methods have been made available, see review from Dizon *et al.*²⁶. Without loss of generality
204 we choose to work with the standard Fused Deposition Modeling (FDM) on the Ultimaker 3 machines of the Innovation co-Lab
205 of Duke University, with a nozzle of 0.4mm diameter. The machine offers the possibility to print multiple materials (see
206 exhaustive list¹ from the manufacturer), with polylactic acid (PLA) and acrylonitrile butadiene styrene (ABS) two of the most
207 commonly used in mechanical testing of 3D printed materials^{26,27}. Without any preferences, we choose to work with PLA.

208 Many of the printing settings influence the mechanical properties of the printed sample, as can be seen in the extensive
209 review of Popescu *et al.*²⁷ as well as our Supplementary text and the references therein pointing to studies on the influence of
210 slicing parameters, building orientation and temperature conditions. It is therefore important to keep those parameters constant
211 for consistency purposes between all samples preparation. Starting from the default settings of the 3D printer, we keep the infill
212 density at 100% in order to have a non-porous matrix material. For the building orientation, the samples are printed vertically
213 and each layer is printed with a rotation of 90 degrees from the previous one in order to reduce the anisotropy of the printing
214 that you would obtain when stacking directly the filaments on top of each other. For the temperature conditions, we follow the
215 recommendation of the Ultimaker 3 user manual for PLA² for the extruder's temperature at 200°C and the one of the bed table
216 at 60°C. For the slicing parameters, the wall/shell thickness of the sample is taken to be equal to the layer height in order to be
217 printed with a single filament in size. Finally, the printing speed is set to 30 mm/s which produces a sample of good quality.

218 All the compression tests presented in this contribution were performed on the HM3000.3F load frame, manufactured by
219 Humboldt Mfg. Co., with a maximum loading capacity of 50 kN. In order to measure the stress on top of the sample, we use
220 the HM-2300.100 S-Type load cell, which has the same load capacity as the machine and is also manufactured by Humboldt
221 Mfg. Co. The strain is measured directly from the speed of the load plate and it was verified that the deformation of the load
222 cell, which is taken in account with this method, had a negligible effect on the results.

223 **Mechanical model for 3D printed PLA**

224 In this section we propose an elasto-plastic model to fit the stress-strain curves, plotted in Fig. 1, of the PLA cylinders printed
225 as described above. Note that we can appreciate the reproducibility of the experiments as the curves superpose to a precision

¹<http://ultimaker.com/materials>

²<http://ultimaker.com/en/resources/22225-how-to-print-with-ultimaker-pla>

226 level hard to obtain experimentally on natural materials.

227 We can observe on that figure that the material does not behave in a linear elastic manner at first but rather displays a
228 non-linear phase due to strain measurement errors²⁸ (e.g. bedding error). Relatively quickly, however, the material follows a
229 linear elastic response once the top stress value reaches a threshold of approximately 5 MPa, as seen on Fig. 1. In order to
230 remove the inconsistent bedding error and have the elastic part of all curves superposed for assessing the reproducibility of the
231 results, we shift the origin of vertical strain of each stress-strain curve so it corresponds to the stress value of 15 MPa, a safe
232 arbitrary value of above which linear elasticity is fully observed.

233 As can be seen on Fig. 1, the elastic properties are extremely consistent between all samples as the elastic parts of the
234 curves completely superpose. We can then measure a Young's modulus of 1375 MPa on the curves of Fig. 1 since the Young's
235 modulus of the material is directly given by the slope of the elastic part in uniaxial compression. Since Poisson's ratio does
236 not play a role in uniaxial compression, we do not measure it and assume the value reported in the literature of 0.45 for our
237 numerical model, as the material is known to be quite incompressible. As a polymer, PLA naturally remains viscoplastic after
238 the printing process. Since this contribution is not focused on quantifying the rate-dependency of the material, we select an
239 arbitrary loading rate of 0.08mm/min for all experiments in this contribution.

240 The plastic part of the curves in Fig. 1 is decomposed in three different phases. Past the end of elastic behaviour, the material
241 is hardening, which is due to the viscoplasticity of the polymer, until the peak stress is reached. After the peak stress, the
242 material experiences softening, arguably from the shearband forming. Eventually, the material converges to a limit load. Note
243 that the curves on Fig. 1 do not display a limit load per se as the stress increases again at large strains. This is a post-processing
244 artefact stemming from the fact that the sample is compressed above an amount of strain at which the top surface of the sample
245 starts increasing due to the shearband. As a result, even though the stress might have converged to a steady value, the load
246 is increasing because the surface on which the stress is applied enlarges. Due to the technical difficulty to properly take this
247 change of surface area into account in the post-processing step, the stress plotted is calculated by accounting for the initial
248 top surface area only, which can therefore not capture the eventual hardening at large strains. For this reason we decide to
249 disregard the last part of the curves where the load increases and consider for the value of limit load the minimum value of
250 stress achieved after softening.

251 Our contribution focuses on the influence of the microstructure of the material yield. For this reason, we do not necessarily
252 need to capture in the model the parts of the stress strain curves that corresponds to the intrinsic behaviour of the material, i.e.
253 the hardening from viscoplasticity and the softening from shearbanding. We focus instead on the two end points of the plastic
254 behaviour: the macroscopic yield which corresponds to the onset of plasticity, and the limit load. Fig. 1 shows the calibration of
255 the plasticity model for the 3D printed PLA with a value of the macroscopic yield measured at 78 MPa and the value of the
256 limit load at 55 MPa.

257 Note that the model selected here is only suggested for the specific printing settings with the testing procedure detailed
258 above and may not be applicable with other parameters as we have shown – non-exhaustively – that many parameters influence
259 the mechanical properties of the printed PLA.

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