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Global trends in novel stable isotopes in basalts: theory and observations

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Abstract

The geochemistry of global mantle melts suggests that both mid-ocean ridge basalts (MORB) and ocean island basalts (OIB) sample lithological and temperature heterogeneities originating in both the upper and lower mantle. Recently, non-traditional stable isotopes have been suggested as a new tool to complement existing tracers of mantle heterogeneity (e.g., major and trace elements, radiogenic isotopes), because mineral- and redox-specific equilibrium stable isotope fractionation effects can link the stable isotope ratios of melts to their source mineralogy and melting degree. Here, we investigate five stable isotope systems (Mg-Ca-Fe-V-Cr) that have shown promise in models or natural samples as tracers of mantle temperature and/or lithological heterogeneity. We use a quantitative model, combining thermodynamically self-consistent mantle melting and equilibrium isotope fractionation models, to explore the behaviour of the isotope ratios of these elements during melting of three mantle lithologies (peridotite, and silica-excess and silica-deficient pyroxenites), responding to changes in mantle mineralogy, oxygen fugacity, temperature and pressure. We find that, given current analytical precision, the stable isotope systems examined here are not predicted to be sensitive to mantle potential temperature variations through equilibrium isotope fractionation processes. By contrast, source lithological heterogeneity is predicted to be detectable in some cases in the stable isotope ratios of erupted basalts, although generally only at proportions of > 10% MORB-like pyroxenite in the mantle source, given current analytical precision. Magnesium and Ca stable isotopes show most sensitivity to a garnet-bearing source lithology, and Fe and Cr stable isotopes are potentially sensitive

28 to the presence of MORB-like pyroxenite in the mantle source, although the behaviour of
29 Cr isotopes is comparatively under-constrained and requires further work to be applied with
30 confidence to mantle melts. When comparing the magnitude and direction of predicted equi-
31 librium isotopic fractionation of peridotite and pyroxenite melts to natural MORB and OIB
32 data, we find that aspects of the natural data (including the mean Mg-Ca-Fe-V isotopic com-
33 position of MORB, the range of Mg-Ca isotopic compositions seen in MORB data, the mean
34 Mg-Ca-Cr isotopic composition of OIB, and the range of Mg-V-Cr isotopic compositions in
35 OIB data) can be matched by equilibrium isotope fractionation during partial melting of
36 peridotite and pyroxenite sources – with pyroxenite required even for some MORB data.
37 However, even when considering analytical uncertainty on natural sample measurements,
38 the range in stable isotope compositions seen across the global MORB and OIB datasets
39 suggests that kinetic isotope fractionation, or processes modifying the isotopic composition
40 of recycled crustal material such that it is distinct from MORB, may be required to ex-
41 plain all the natural data. We conclude that the five stable isotope systems considered here
42 have potential to be powerful complementary tracers to other geochemical tracers of the
43 source lithology of erupted basalts. However, continued improvements in analytical preci-
44 sion in conjunction with experimental and theoretical predictions of isotopic fractionation
45 between mantle minerals and melts are required before these novel stable isotopes can be
46 unambiguously used to understand source heterogeneity in erupted basalts.

47 1 Introduction

48 A key window into Earth’s mantle is through the petrology and geochemistry of its melts, which
49 include both mid-ocean ridge basalts (MORB) and ocean island basalts (OIB). MORB may be
50 able to tell us about the composition and temperature of the upper mantle across multiple length
51 scales (e.g., Allègre et al., 1984; Zindler and Hart, 1986; Mahoney et al., 1994; Agranier et al.,
52 2005; Janney et al., 2005; Herzberg et al., 2007; Gale et al., 2013; Shorttle, 2015). Key ques-
53 tions about MORB-source mantle concern the various roles of temperature, crustal thickness,
54 lithological heterogeneity, melt mixing and melt-rock reaction and the degree to which they can
55 explain the trace, major element, and radiogenic isotope arrays of global MORB. Decades of in-
56 vestigation have produced a broad consensus that in the MORB-source mantle thermally-driven
57 variations in melting degree create global signals in major and/or trace elements (Klein and
58 Langmuir, 1987; Dalton et al., 2014; Gale et al., 2014). On a local scale MORB also show con-

59 siderable isotopic and trace element variability, linked to small-scale lithological heterogeneity in
60 the upper mantle (e.g., Zindler et al., 1984; Langmuir et al., 1986; Fornari et al., 1988; Hekinian
61 et al., 1989; Schiano et al., 1997; Castillo et al., 2000; Waters et al., 2011; Gill et al., 2016; Liu
62 and Liang, 2017; Jiang et al., 2021; Zhong et al., 2021). OIB also sample a heterogeneous mantle
63 over short and long lengthscales including significant lithological heterogeneity (e.g., Cohen and
64 O’Nions, 1982; Weaver, 1991; Chauvel et al., 1992; Kogiso et al., 2003; Sobolev et al., 2007; Jack-
65 son and Dasgupta, 2008; Dasgupta et al., 2010; Day and Hilton, 2011; Shorttle and MacLennan,
66 2011), elemental and isotopic heterogeneity (e.g., Zindler and Hart, 1986; Dupuy et al., 1988;
67 Hauri and Hart, 1993; Hofmann, 1997, 2003; Stracke et al., 2005; Willbold and Stracke, 2006;
68 Prytulak and Elliott, 2007; Jackson and Dasgupta, 2008; MacLennan, 2008; Dasgupta et al.,
69 2010; Jackson et al., 2012; Mundl et al., 2017), and record large temperature variations relative
70 to MORB (e.g., Putirka, 2005; Herzberg et al., 2007; Putirka, 2008a,b; Herzberg and Gazel,
71 2009; Herzberg and Asimow, 2015; Matthews et al., 2021).

72 In this study, we investigate the use of novel stable isotopes of major and minor elements in
73 basalts (magnesium [$\delta^{26/24}\text{Mg}$], calcium [$\delta^{44/40}\text{Ca}$], iron [$\delta^{57/54}\text{Fe}$], vanadium [$\delta^{51/50}\text{V}$], chromium
74 [$\delta^{53/52}\text{Cr}$]) as new tools for studying temperature and lithological heterogeneity in the mantle,
75 complementary to existing work on MORB and OIB mantle sources using major and trace el-
76 ements (e.g., Klein and Langmuir, 1987; Langmuir et al., 1992; McKenzie et al., 2004; Putirka,
77 2005; Herzberg et al., 2007; Jackson and Dasgupta, 2008; Niu and O’Hara, 2008; Dasgupta et al.,
78 2010; Le Roux et al., 2011; Jackson et al., 2012; Lambart et al., 2013; Gale et al., 2014; Yang
79 et al., 2019; Mallik et al., 2021). The chosen isotope systems may trace thermal and mineralogi-
80 cal heterogeneity in the source region of melts through their temperature and mineral-dependent
81 equilibrium fractionation factors (e.g., Schauble, 2004; Young et al., 2015), and we include detail
82 on the choice of these isotope systems in Section 1.3.

83 1.1 Temperature and lithological heterogeneity in the mantle

84 Mantle-derived melts, particularly MORB, show global geochemical arrays that can be explained
85 by the conditions of melt generation in the mantle (e.g., temperature, lithology; Klein and Lang-
86 muir, 1987; Langmuir et al., 1992; Niu and O’Hara, 2008; Arevalo Jr and McDonough, 2010;
87 Gale et al., 2014) and/or subsequent melt transport and storage processes (e.g., melt mixing,

88 melt-rock reaction, fractional crystallisation; Devey et al., 1994; Arevalo Jr and McDonough,
89 2010; Liang et al., 2011; Kimura and Sano, 2012; Till et al., 2012; Shorttle, 2015; Bo et al.,
90 2018; Mallik et al., 2021; Stracke, 2021). Once corrected for crystal fractionation, moderately
91 incompatible element concentrations (e.g., Na, Al) in global MORB correlate positively with
92 ridge depth, commonly interpreted to equate to a control by mantle potential temperature
93 variations of around 220 °C (Klein and Langmuir, 1987; Langmuir et al., 1992; Brandl et al.,
94 2013; Gale et al., 2014). Correlations between fractionation-corrected major element composi-
95 tions in MORB (e.g., FeO-Na₂O, CaO-Al₂O₃) are consistent with the same mantle temperature
96 variations (Langmuir et al., 1992; Gale et al., 2014). Although mantle composition has also
97 been invoked instead of temperature variation to explain these trends (by applying a different
98 fractionation correction; Niu and O’Hara, 2008), major element compositions measured in OIB
99 and their constituent olivine, as well as olivine crystallisation thermometry and geophysical ob-
100 servations, argue in favour of there being regions of mantle up to 250 °C hotter than typical
101 mid-ocean ridge mantle (e.g., Putirka, 2005; Herzberg et al., 2007; Putirka, 2008a; Dalton et al.,
102 2014; Herzberg and Asimow, 2015; Spice et al., 2016; Ball et al., 2021; Matthews et al., 2021).

103 Although temperature may be a dominant control on MORB geochemical variability, man-
104 tle compositional heterogeneity is also clear in the elemental and isotopic compositions of
105 both MORB and OIB. For example, in MORB highly incompatible element enrichment (e.g.,
106 fractionation-corrected K concentrations) correlates roughly with radiogenic isotopes, which
107 must relate to source heterogeneity, such as through the addition or removal of low-degree
108 melts (e.g., McKenzie and O’Nions, 1995; Donnelly et al., 2004; Gale et al., 2011, 2013, 2014).
109 Radiogenic and stable isotopes and trace elements have also been used to argue for the pres-
110 ence of incompatible trace element-enriched, recycled components in the MORB source (e.g.,
111 Hirschmann and Stolper, 1996; Schiano et al., 1997; Salters and Dick, 2002; Bezard et al., 2016).
112 In OIB, there is extensive evidence for mantle heterogeneity in the form of radiogenic and stable
113 isotope compositions and the concentrations of major elements in basalt and trace elements in
114 olivine (e.g., Zindler and Hart, 1986; Weaver, 1991; Hauri, 1996; Hofmann, 1997; Sobolev et al.,
115 2005; Stracke et al., 2005; Sobolev et al., 2007; Jackson and Dasgupta, 2008; Day et al., 2009;
116 Herzberg, 2011; Shorttle and MacLennan, 2011; Konter et al., 2016; Mundl et al., 2017; Neave
117 et al., 2018; Nebel et al., 2019; Gleeson et al., 2020). This mantle heterogeneity is thought to
118 relate largely to the recycling of crust, and the reaction between recycled lithologies (and their

119 melts) and ambient mantle peridotite (e.g., Cohen and O’Nions, 1982; Hofmann and White,
120 1982; Allègre and Turcotte, 1986; Schiano et al., 1997; Sobolev et al., 2005, 2007; Herzberg,
121 2011; Mallik and Dasgupta, 2012; Stracke, 2012; Rosenthal et al., 2014; Lambart, 2017).

122 Many tools are available to study the contributions of temperature and lithological heterogeneity
123 in these two types of basalts, for example: rare earth element inversion modelling (INVMEL:
124 McKenzie and O’Nions, 1991, 1995), major element calculations (PRIMELT: Herzberg and Asi-
125 mow, 2008, 2015; the thermodynamic model of Jennings and Holland, 2015; Jennings et al.,
126 2016), trace element and radiogenic isotope composition forward modelling (REEBOX PRO:
127 Brown and Lesher, 2016), and models combining trace elements, crystallisation temperatures
128 and magma productivity (Shorttle et al., 2014; Matthews et al., 2016, 2021). However, achiev-
129 ing success in linking the elemental and radiogenic isotope variability in MORB and OIB to
130 temperature and/or lithological heterogeneity is complicated by uncertainty in the nature of
131 enriched lithologies, metasomatism by small volumes of melt (which are usually highly enriched
132 in incompatible elements, so can overwhelm evidence of the original source lithology), magma
133 recharge and mixing, diffusional re-equilibration and fractional crystallisation (e.g., Niu and
134 O’Hara, 2003; Workman et al., 2004; Niu and O’Hara, 2008; Lambart et al., 2013; Matzen et al.,
135 2017; Gleeson and Gibson, 2019). Stable isotopes of major and minor elements complement
136 radiogenic isotopes and major element calculations and could provide new constraints on under-
137 standing global basalt chemistry: stable isotope systems are time-independent, and since their
138 equilibrium partitioning is a function of mineral chemistry, they are the natural extension of
139 major element calculations based on thermodynamic data. In addition, some of these stable
140 isotopes are major elements whose budgets in the mantle source and subsequent erupted melts
141 reflect contributions from both enriched and depleted lithologies, and are not dominated by
142 small-degree melt metasomatism.

143 **1.2 Stable isotopes as a probe of mantle temperature and compositional** 144 **heterogeneity**

145 Stable isotopes may trace thermal and mineralogical heterogeneity in the mantle through their
146 temperature and mineral-dependent equilibrium fractionation factors (e.g., Schauble, 2004; Young
147 et al., 2015). Our understanding of the equilibrium high temperature fractionation effects ex-

148 perience by novel stable isotope systems is improving, and many studies (both theoretical and
149 from natural samples) suggest controls from mantle temperature and lithology on the isotopic
150 composition of the melts produced. Previous studies have shown considerable variability in
151 stable isotope ratios in MORB and OIB even where the geochemical signals of fractional crys-
152 tallisation are considered to be negligible, or have been removed (Fig. 1), but understanding
153 the origin of the variability remains challenging. For example, Fig. 1 shows that natural basalt
154 stable isotope data for the isotope systems considered here have different systematics relative
155 to the bulk silicate Earth (BSE). Magnesium and V stable isotope compositions in MORB and
156 OIB scatter both sides of the BSE value, whereas Fe stable isotope compositions of basalts are
157 heavier than the BSE, and for Cr and Ca basalts are (mostly) isotopically lighter than the BSE.
158 The Mg and Fe stable isotope compositions of OIB show greater variability than MORB for the
159 same isotope system, whereas the opposite is true for V stable isotopes (although this may be
160 a sampling bias, given the greater amount of V isotope data for MORB than OIB).

161 Modelling of expected stable isotopic behaviour has led to questions about the origin of isotopic
162 signatures observed in natural samples. For example, some studies have modelled Ca stable
163 isotope fractionation using composition-dependent inter-mineral fractionation factors and con-
164 cluded that partial melting of eclogite cannot produce the variability measured in OIB (Chen
165 et al., 2020a), whereas others have predicted and measured that garnet has a higher $\delta^{44}\text{Ca}$ than
166 coexisting clinopyroxene (Antonelli et al., 2019; Huang et al., 2019; Kang et al., 2019; Wang
167 et al., 2019; Chen et al., 2020a; Dai et al., 2020; Smart et al., 2021; Tappe et al., 2021), and have
168 used this observation to predict that melts from recycled oceanic crust (garnet-bearing eclogite)
169 will have low $\delta^{44}\text{Ca}$, thus explaining the low $\delta^{44}\text{Ca}$ measured in some OIB (Kang et al., 2019;
170 Dai et al., 2020). Lithological heterogeneity (specifically, recycled pyroxenite) has also been
171 linked to Mg and Fe isotope variability in MORB and OIB (e.g., Williams and Bizimis, 2014;
172 Konter et al., 2016; Zhong et al., 2017; Stracke et al., 2018; Nebel et al., 2019; Gleeson et al.,
173 2020; Sun et al., 2020; Zhong et al., 2021). However, some recent models of Fe and Mg isotope
174 fractionation have led to uncertainty in whether equilibrium fractionation associated with the
175 presence of garnet can unambiguously identify a garnet-bearing pyroxenite source lithology in
176 natural OIB samples (Stracke et al., 2018; Soderman et al., 2021), and whether the heaviest
177 Fe isotope data in the global dataset can be matched by models of mantle melting (Sun et al.,
178 2020; Soderman et al., 2021).

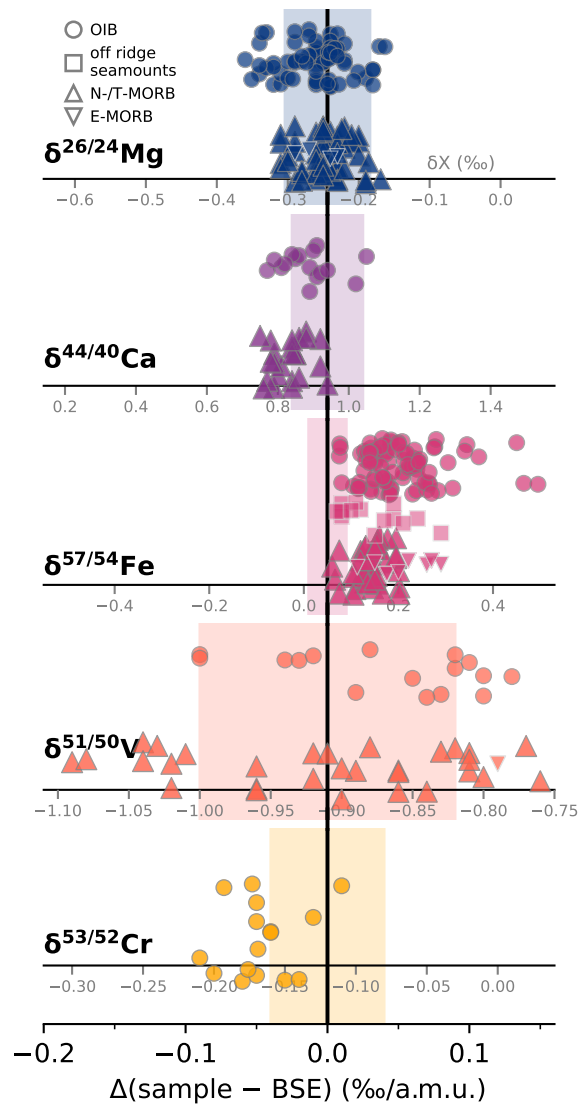


Figure 1: Literature MORB and OIB data for the isotope systems (for sources, see Table S1). Samples are filtered for $7.2 < \text{MgO wt}\% < 16$, but included if MgO is not known, to limit the effects of fractional crystallisation and olivine accumulation. Where MORB type is known, samples are classified as N-/T- or E-type (as classified by Teng et al., 2013); if classification is not known, samples are plotted as N-/T-type. Ca and Mg isotope data is filtered to exclude samples with $^{87}\text{Sr}/^{86}\text{Sr} > 0.7037$, which excludes samples with $> 1\%$ carbonate component in the source (Huang et al., 2011b). The main axis shows the variability in natural data from the BSE value for each isotope system, scaled by the difference in the atomic mass units (a.m.u.) of the isotopes being measured (e.g., for $\delta^{26/24}\text{Mg}$, this a.m.u. factor is 2). The coloured bars show 2 S.D. on the BSE estimates. BSE values used are $\delta^{26}\text{Mg} = -0.24 \pm 0.06\text{‰}$ (Stracke et al., 2018), $\delta^{44}\text{Ca} = 0.94 \pm 0.10\text{‰}$ (Kang et al., 2017; Chen et al., 2019), $\delta^{57}\text{Fe} = 0.05 \pm 0.04\text{‰}$ (Sossi et al., 2016), $\delta^{51}\text{V} = -0.91 \pm 0.09\text{‰}$ (Qi et al., 2019), $\delta^{53}\text{Cr} = -0.12 \pm 0.04\text{‰}$ (Jerram et al., 2020). The pale secondary axis for each isotope shows the raw isotope ratios. The isotope ratios and reference materials used in the delta notation throughout are: $\delta^{26/24}\text{Mg}$ [DSM3]; $\delta^{44/40}\text{Ca}$ [SRM915a]; $\delta^{57/54}\text{Fe}$ [IRMM-014]; $\delta^{51/50}\text{V}$ [AA]; $\delta^{53/52}\text{Cr}$ [NIST979].

179 Given the framework of MORB and OIB data that we now have (Fig. 1), growing theoretical
 180 information on bond strength and predicted inter-mineral fractionations for multiple stable iso-
 181 tope systems, and a variety of published isotope fractionation models that both can and cannot
 182 explain all the natural data, it is timely to develop a self-consistent model for the behaviour of

183 multiple stable isotopes during mantle melting. Here we use a thermodynamically self-consistent
184 model for mantle melting of three lithologies (peridotite and two different pyroxenites), combined
185 with a model for equilibrium Mg-Ca-Fe-V-Cr stable isotope fractionation responding to changes
186 in mantle mineralogy, intrinsic variations in oxygen fugacity (i.e., at constant O content), tem-
187 perature and pressure. These models allow us to investigate the potential for the stable isotope
188 composition of basalts to be a tracer of mantle temperature and lithological heterogeneity.

189 **1.3 Approach**

190 This contribution explores the behaviour of Mg-Ca-Fe-V-Cr stable isotopes during mantle melt-
191 ing, expanding on work on Fe stable isotope behaviour presented in Soderman et al. (2021).
192 These isotope systems were chosen as they have been documented to show resolvable mineral-
193 specific fractionation effects that may make them sensitive to partial melting or source lithology
194 effects (e.g., Konter et al., 2016; Xia et al., 2017; Stracke et al., 2018; Wu et al., 2018; Kang
195 et al., 2019; Nebel et al., 2019; Dai et al., 2020; Gleeson et al., 2020; Shen et al., 2020; Novella
196 et al., 2020), and their bonding environments and/or expected isotopic fractionation in mantle
197 minerals and melt are sufficiently well-studied to provide reasonable inputs for an isotopic frac-
198 tionation model. The isotope systems represent both major and trace elements in basalts, and
199 monovalent and heterovalent elements. The major elements (Mg, Ca, Fe) are abundant in the
200 mantle with comparable or lower concentrations in pelitic sediments and crustal material (Plank
201 and Langmuir, 1998; Rudnick and Gao, 2003; Workman and Hart, 2005), meaning the isotopic
202 signatures of mantle components are not easily affected by metasomatism, and their isotopic
203 compositions in basalts should track the bulk mass contributions of both fertile and depleted
204 mantle lithologies to the melt. This is an important property of major element stable isotope
205 systems that contrasts with radiogenic isotope systems of incompatible trace elements, which
206 cannot probe the proportions of different lithologies directly, as incompatible element concen-
207 tration vary widely between enriched and depleted mantle lithologies. We note that carbonates,
208 whose presence have been invoked in the source regions of mantle melts (e.g., Huang et al.,
209 2011b; Liu et al., 2017a; Wang et al., 2018), have Ca concentrations ~ 10 times higher than
210 the mantle (e.g., Huang et al., 2011b), but we do not discuss carbonates here; instead we filter
211 natural data used in this study by $^{87}\text{Sr}/^{86}\text{Sr}$ to exclude significant contributions from carbonate

212 components.

213 Section 2 contains a description of the combined thermodynamic melting and equilibrium isotope
214 fractionation model. The model was used to calculate the equilibrium isotopic composition of
215 mantle melts from different lithologies over P-T space, and a summary of the results is presented
216 in this section.

217 Sections 3 and 4 address the degree to which different stable isotope measurements (individual
218 isotope systems, or in combination) can identify the relative importance of mantle temperature
219 variability and mantle lithological heterogeneity in generating the observed variability in basalts.
220 We also discuss where stable isotope systems have the potential to be useful in identifying these
221 processes, if measurement uncertainties can be reduced.

222 Finally, Section 5 combines MORB and OIB data for the isotope systems studied with the
223 modelled equilibrium melts, to assess to what extent our present understanding of the isotope
224 behaviour can explain global basalt variability. We do not include arc basalts in our discussion
225 for a number of reasons, including but not limited to the complicating effects of H₂O, slab fluids
226 and redox variability in an arc setting, which have been linked to stable Fe and Mg isotope
227 variability recorded in arc basalts (e.g., Dauphas et al., 2009; Nebel et al., 2013, 2015; Sossi
228 et al., 2016; Teng et al., 2016; Li et al., 2017; Brewer et al., 2018; Hu et al., 2020) and which
229 are beyond the applicable scope of the modelling presented here.

230 **2 Modelling equilibrium isotopic composition of mantle melts**

231 We calculated the equilibrium melting isotopic fractionation over P-T space for $\delta^{26/24}\text{Mg}$,
232 $\delta^{44/40}\text{Ca}$, $\delta^{57/54}\text{Fe}$, $\delta^{51/50}\text{V}$ and $\delta^{53/52}\text{Cr}$ for three representative mantle lithologies, following
233 the model outlined in Soderman et al. (2021). We used the calculated modal mineralogies over
234 P-T space of KLB1 peridotite (a commonly used experimental composition used as an analogue
235 for the upper mantle; Davis et al., 2009) and G2 silica-excess pyroxenite (an important MORB-
236 like bulk composition in melting experiments and models from Pertermann and Hirschmann,
237 2003*a,b* and Lambart et al., 2016) presented in Soderman et al. (2021). Following the methods
238 used for KLB1 and G2, we also calculated the P-T-dependent modal mineralogy of MIX1G, a

239 silica-deficient pyroxenite which plots close to the average global pyroxenite composition and
240 can be considered as a mixture between KLB1 and MORB (Lambart et al., 2016), i.e., recycled
241 crust mixed with ambient mantle. Calculations were performed using the dataset of Holland and
242 Powell (2011) and activity-composition (a-X) models of Holland et al. (2018), implemented in
243 THERMOCALC (Powell et al., 1998); see Section 2 in Appendix for details. This set of lithologies
244 spans a range of fertile compositions, both ambient mantle and those derived from recycled crust,
245 thought to be present in the mantle (Hirschmann and Stolper, 1996; Shorttle and MacLennan,
246 2011; Lambart et al., 2016; Mallik et al., 2021).

247 The THERMOCALC output includes the proportion of each phase present and information on the
248 composition of each phase at any P-T point, including the distribution of elements on available
249 cation sites. These results allow the equilibrium isotopic composition of the phases present
250 (including melts) to be calculated for each isotope system, based on calculated equilibrium
251 isotope fractionation factors, α^{A-B} , between the phases A and B that are present, where

$$\alpha^{A-B} = \frac{\left(\frac{X_2}{X_1}\right)_A}{\left(\frac{X_2}{X_1}\right)_B}, \quad (1)$$

252 and X represents the amount of each isotope 1 or 2 of an element. To facilitate comparisons
253 between different lithologies and isotope systems, results are all expressed relative to a bulk
254 system of 0‰, for convenience (although we recognise that different mantle components may
255 have non-zero initial bulk stable isotope compositions for the stable isotope systems considered
256 here). Iron isotope compositions of minerals and melts were calculated for MIX1G, following
257 Soderman et al. (2021); and Ca, Mg, Cr and V isotope compositions were newly calculated for
258 all three lithologies (full set of results for isotopes and lithologies in Section 3 of the Appendix,
259 and Electronic Annex). The approach of the model for each isotope is described briefly below,
260 with full details in the Appendix.

261 We cannot consider H₂O in our phase-equilibrium calculations, as the THERMOCALC peridotite
262 dataset used does not account for H₂O accommodation in the mantle phases (Holland et al.,
263 2018). However, for the comparison of our models to MORB and OIB, we consider this dry
264 system to be suitable: non-arc mantle is generally considered to have small amounts of water
265 stored in nominally anhydrous minerals (Bell and Rossman, 1992; Hirschmann et al., 2005).

266 While H₂O, if present in the mantle source, would result in a deeper onset of melting than in
267 dry melting, it does not result in significant changes to the maximum extent of melting and
268 most water would be extracted in the first degrees of melting (e.g., Hirth and Kohlstedt, 1996;
269 Hirschmann et al., 1999; Katz et al., 2003), so we consider that the H₂O concentrations will be
270 very low throughout most of the melting region. Therefore, the conclusions reached here will
271 not be significantly changed by the inclusion of water in the thermodynamic system.

272 2.1 Ca, Mg and Fe

273 The partitioning of Ca, Mg and Fe (Fe²⁺ and Fe³⁺) between the phases present at any P-T
274 point was taken from the THERMOCALC results. To calculate the isotope fractionation factors
275 (α) for Ca and Mg isotopes between each phase, published temperature-dependent β -values (a
276 reduced partition function ratio; Young et al., 2015) are used, where

$$\alpha^{A-B} = \frac{\beta(T)_A}{\beta(T)_B}. \quad (2)$$

277 Bulk β -values for each mineral (i.e., not using crystal site occupancy information) are calculated
278 based on vibrational frequencies from first-principle calculations, and are taken from Antonelli
279 et al. (2019) for Ca isotopes and Huang et al. (2013) for Mg isotopes. Where suitable β -values
280 are not published, as for Fe, α factors can be calculated using cation-oxygen bond force constants
281 in each crystal site for the minerals,

$$\ln(\alpha^{A-B}) = C \frac{[K_A - K_B]}{T^2}, \quad (3)$$

282 where C is a constant dependent on the isotopes and element being considered (e.g., Sossi and
283 O'Neill, 2017), T is the temperature in Kelvin, and K_A and K_B refer to the cation-oxygen force
284 constants in minerals A and B. These force constants are based on an ionic bonding model for
285 Fe (Sossi and O'Neill, 2017), as employed in Macris et al. (2015); Young et al. (2015); Soderman
286 et al. (2021); Williams et al. (2021).

287 **2.2 V and Cr**

288 The THERMOCALC a-X models of Holland et al. (2018) do not model V partitioning, nor het-
289 erovalent Cr (Cr is assumed to have a valence of 3+ in the Holland et al., 2018 models). There-
290 fore, the proportion of $V^{3+/4+/5+}$ was calculated following Toplis and Corgne (2002), using the
291 FeO/Fe₂O₃ ratio of the melt (calculated in THERMOCALC) at each point in P-T space. The
292 proportion of $Cr^{2+/3+}$ was calculated using the parameterisation of Berry et al. (2021), with
293 fO_2 at each point calculated using the ‘ fO_2 melt’ software (Holland et al., 2018), hence Cr spe-
294 ciation also depends on the Fe redox equilibria. Here, we note that the requirement to handle
295 heterovalent Cr and V outside of the THERMOCALC calculations means that the calculations
296 of the redox state (hence, ultimately isotopic fractionation) are not fully self-consistent for the
297 redox-sensitive elements. The $Cr^{3+}/Cr_T, V^{4+}/V_T$ and V^{5+}/V_T ratios change across P-T space
298 without taking oxygen from anywhere else (i.e., the Fe^{3+}/Fe_T ratio calculated at each P-T point
299 does not vary when Cr and V redox equilibria are considered). Instead, the calculations of the
300 valence states of Cr and V at each P-T point are dictated by the calculated Fe redox equilibria
301 at that point. Nonetheless, tying the redox equilibria of Cr and V to that of Fe is a reasonable
302 approximation of the natural system, given that the abundance of Fe in the mantle is signifi-
303 cantly greater than that of the other redox-variable elements (Anenburg and O’Neill, 2019; see
304 also Appendix Section 3.5 for the oxygen distribution in our models). As a consequence of our
305 model’s implementation of Cr/V redox, the changes in Cr/V speciation in the model will represent
306 upper limits: i.e., in a more complex system coupled Fe-Cr-V redox equilibria will dampen the
307 variability of $Fe^{3+}/Fe_T, Cr^{3+}/Cr_T$ and $V^{4+,5+}/V_T$ compared to our models.

308 The total V and Cr in the system were partitioned, dependent on the relative proportions of
309 each valence state, between the phases present based on valence-specific partition coefficients
310 from Mallmann and O’Neill (2009). The α factors for V were calculated following the approach
311 used for Mg and Ca above, using β^V values for V in solution systems from Wu et al. (2015),
312 extrapolated to magmatic systems at high temperature following Wu et al. (2018). For Cr, α
313 factors were calculated from Cr-O force constants following an ionic model (Shen et al., 2018,
314 2020), as for Fe isotopes (Soderman et al., 2021).

315 **2.3 Results**

316 The results of the THERMOCALC calculations and subsequent melt stable isotope composition
 317 calculations are given in the Electronic Annex and Section 3 of the Appendix. A summary of
 318 the modelled equilibrium isotopic fractionation of peridotite melts (relative to the source), is
 319 shown in Fig. 2. Melt isotope ratios are shown along isentropic decompression melting paths
 320 calculated from THERMOCALC output – the isentropic approach (McKenzie, 1984; McKenzie
 321 and Bickle, 1988) assumes the adiabatic upwelling of the mantle is reversible, which is generally
 322 considered to be a good approximation of decompression melting in the mantle (e.g., McKenzie,
 323 1984; McKenzie and Bickle, 1988; Asimow, 2002).

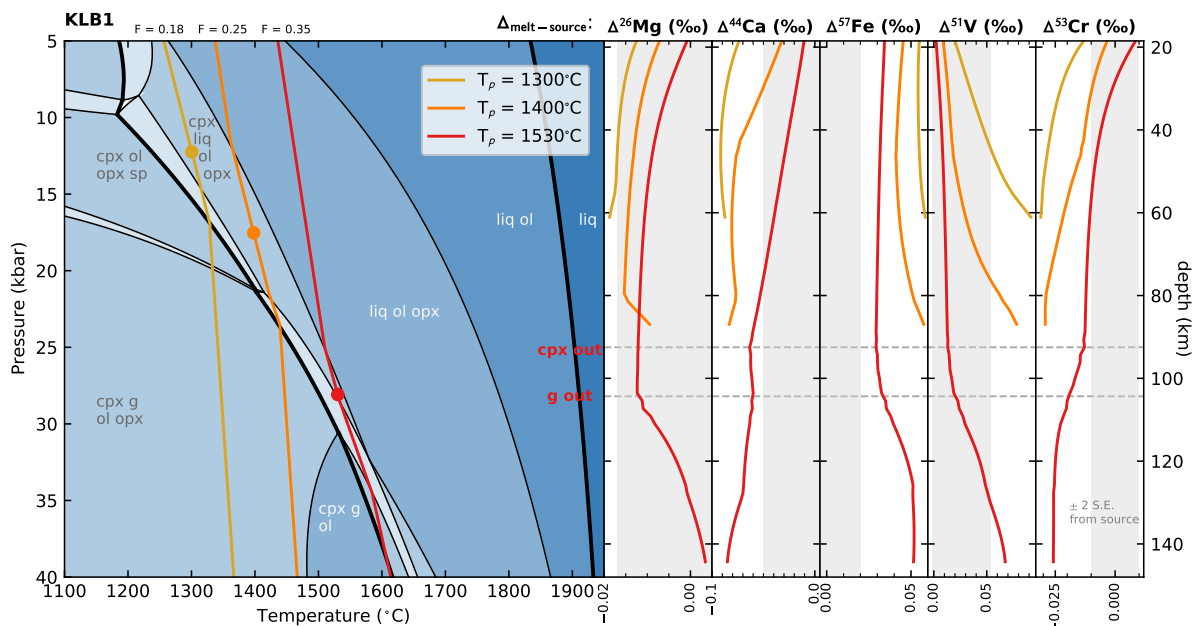


Figure 2: The left hand panel shows the pseudosection for KLB1 peridotite calculated in THERMOCALC, with the solidus and liquidus marked in bold, and three isentropic melting paths also calculated using THERMOCALC. The average pressure of melting for each isentrope (see Appendix Section 6) is shown by the small circles, and maximum melt fraction (F) labelled at the top of each isentrope. The right hand panels show the modelled instantaneous equilibrium melt isotopic fractionation relative to the bulk isotopic composition along both isentropes. Key mineralogical controls on the isotopic composition of the melt are highlighted (colour-coded by the isentropic path they refer to; here both for $T_p = 1530^\circ\text{C}$). Vertical grey bars show the source composition (0‰) with typical sample 2 S.E. analytical precision for each isotope system. The analytical precisions used are based on representative errors from recent studies: for Mg from Wang et al. (2021), for Ca from Zhu et al. (2018, 2020a); Chen et al. (2020a), for Fe from Soderman et al. (2021), for V from Wu et al. (2018); Novella et al. (2020), for Cr from Sossi et al. (2018). Note that long-term analytical 2 S.D. uncertainty is usually greater than sample 2 S.E. for each isotope system; see Table 1.

324 Table 1 summarises how the modelled stable isotope compositions of pyroxenite melts compare
 325 to the peridotite melts from Fig. 2. Due to the wider range of melting degrees reached by
 326 pyroxenites compared to peridotite for a given T_p , generally the pyroxenite melts show a wider

327 range of isotope ratios than peridotite melts for the same T_p .

Table 1: Maximum magnitude $\Delta_{\text{melt-source}}$ for modelled isentropic decompression melting of KLB1 (peridotite), MIX1G (silica-deficient pyroxenite) and G2 (silica-excess pyroxenite), all following KLB1 isentropes, when $T_p = 1300^\circ\text{C}$ ('cold'), $T_p = 1400^\circ\text{C}$ ('mid') and $T_p = 1530^\circ\text{C}$ ('hot'). Typical analytical uncertainties used throughout this study are also given.

Lithology	Temperature	Maximum modelled melt-source fractionation				
		$\Delta^{26}\text{Mg}$	$\Delta^{44}\text{Ca}$	$\Delta^{57}\text{Fe}$	$\Delta^{51}\text{V}$	$\Delta^{53}\text{Cr}$
KLB1	cold	-0.019	-0.092	0.058	0.092	-0.031
	mid	-0.015	-0.083	0.058	0.078	-0.029
	hot	-0.012	-0.085	0.052	0.067	-0.026
MIX1G	cold	-0.081	-0.081	0.050	0.093	-0.058
	mid	-0.024	-0.137	0.045	0.090	-0.049
	hot	0.099	-0.087	0.026	0.022	-0.036
G2	cold	0.088	-0.125	0.048	0.074	-0.095
	mid	0.088	-0.113	0.036	0.041	-0.054
	hot	0.00	0.00	0.00	0.00	0.00
1 S.D. long-term analytical precision		0.03	0.07	0.025	0.04	0.02
2 S.E. typical sample error		0.02	0.05	0.02	0.05	0.01

328 3 Stable isotope sensitivity to mantle temperature variations

329 To explore the sensitivity of stable isotopes in basalts to variations in mantle potential tem-
 330 perature, we use the results of the equilibrium melt fractionation model for KLB1, since for
 331 MORB, the dominant source lithology is peridotite (Hirschmann and Stolper, 1996; Sobolev
 332 et al., 2007). Generally, all equilibrium source-melt isotopic fractionations are expected to de-
 333 crease in magnitude with increased temperature (e.g., Bigeleisen and Mayer, 1947), although
 334 this effect will be small at high temperatures (typically scaling with $1/T^2$). However, the com-
 335 plexity of mineral reactions during melting and the mineral-specific fractionation effects shown
 336 by the stable isotope systems means some isotopes may show larger temperature (and extent
 337 of melting) sensitivity than others as modal mineralogy and mineral chemistry changes across
 338 P-T space. When using stable isotopes as a tool to understand mantle melting processes and
 339 source compositions, we are interested in two aspects of equilibrium stable isotope fractionation:
 340 1) the mineralogy-dependent fractionation on melting, and 2) the degree-of-melting-dependent
 341 fractionation.

342 MORB generated from different mantle potential temperatures will vary in 1) the modal miner-
 343 alogy of their mantle source, 2) source oxygen fugacity ($f\text{O}_2$), as this changes with pressure and
 344 temperature (e.g., Stolper et al., 2020), and 3) the degree of partial melting of the mantle source

(Fig. 2; also Fig. S9). For MORB, the mean pressure and composition of melts produced in the mantle, hence the aggregate MORB composition, can be calculated using equations for 2D melting regimes (e.g., Klein and Langmuir, 1987; McKenzie and Bickle, 1988; Langmuir et al., 1992; Asimow et al., 2001), and show that the mean melting pressure varies from ≈ 12 kbar for mantle potential temperatures around 1300°C , to > 25 kbar for potential temperatures over 1500°C . The THERMOCALC results show that increasing mantle temperature, with a corresponding deepening of the average pressure of melting reflecting the ‘aggregate’ melt, leads to progressively increasing melt fraction and progressively decreasing $\text{Fe}^{3+}/\text{Fe}_T$ in the melt (due to changing mineral chemistry; Stolper et al., 2020), with minimal changes in the proportions of solid phases in the stable assemblage at the mean melting pressure for each T_p (Figs. 2, 3: see Section 6 in Appendix for details of the average melt pressure calculation). Therefore, the stable isotope compositions of melts could be dependent on mantle temperature through an $f\text{O}_2$ dependence, which will arise due to the existence of multiple valence states, and the different melt partitioning behaviour of each valence state of the cation (e.g., Fe, Cr and V). Isotopes that are most sensitive to melting degree, such as those with large magnitude melt-source isotope fractionations at low melting degrees, may also show temperature-related melt isotope composition variations.

Fig. 4 shows how the estimated isotopic composition of aggregate melts relative to the source, taken as the isotope ratio of the melt at the mean pressure of melting for each temperature, varies with mantle potential temperature. Using the mean melting pressure approach, which is a simplification of polybaric melting occurring in the Earth (Langmuir et al., 1992), gives similar results to polybaric decompression melting (see Appendix Section 6 for a comparison of the approaches). The approach used here does not reflect accumulated fractional melting – the mean pressure of melting is in the spinel stability field, and therefore the isotope ratios of the final aggregate melts do not show a clear garnet signature. The isotope fractionation models show that relatively large magnitude temperature-related effects are seen in the heterovalent Fe and V isotope systems, in agreement with predictions above that $f\text{O}_2$ sensitivity will be the driver of T_p -related isotopic fractionation variation. The Cr isotope ratio of the melt does not show large variability with potential temperature despite being a heterovalent system – this is because the $\text{Cr}^{2+}/\text{Cr}_T$ ratio in the system is dependent on pressure and temperature, as well as $f\text{O}_2$ (Berry et al., 2021). Therefore, the $f\text{O}_2$ effect is not dominant when changing T_p , in contrast

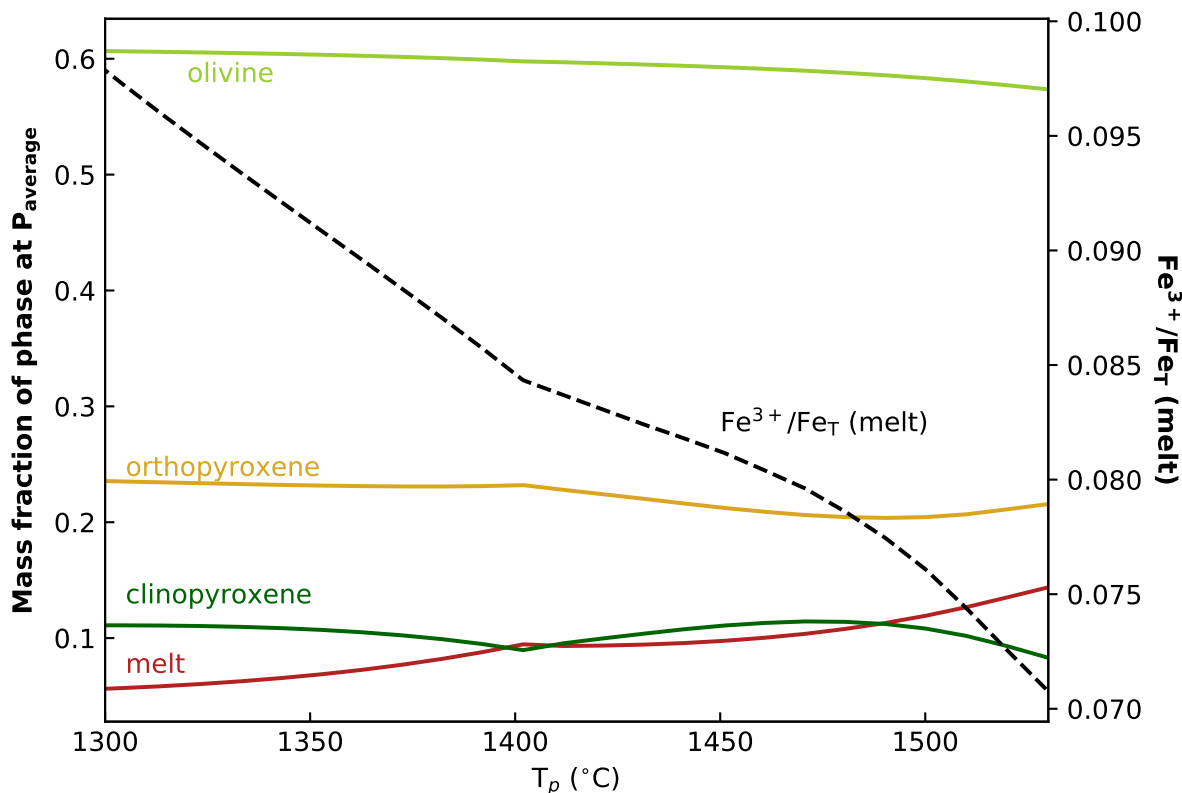


Figure 3: The mass fraction of phases present in the equilibrium KLB1 mineral assemblage and melt Fe^{3+}/Fe_T , at the mean pressure of melt generation calculated for increasing mantle potential temperature. Mantle potential temperature, T_p , is calculated by assuming a linear interpolation between the isentropes calculated in THERMOCALC for $T_p = 1300, 1400$ and 1530 °C. The solid phase proportions are fractions of the whole assemblage (including melt). The average pressure of melting was calculated for $T_p = 1300, 1400$ and 1530 °C (see details in Appendix Section 6), and interpolated between.

376 to V stable isotopes where the V valence state parameterisation is only linked to the Fe^{3+}/Fe_T
 377 at any point. Calcium isotope compositions of accumulated mantle melts also show dependence
 378 on temperature. This variation is linked to the increase in average melt fraction with increasing
 379 temperature; this effect is large for Ca compared to Mg isotopes, the other monovalent system,
 380 because of the much larger magnitude of melt-source Ca isotope fractionation (nearly 0.1 ‰ at
 381 low melt fraction), which results in an increased sensitivity to melting degree.

382 The relative enrichment of the heavy stable isotopes of Fe and V in the melt relative to the
 383 source (because oxidised, isotopically heavy cations behave more incompatibly during mantle
 384 melting than the reduced, isotopically lighter cations of the same element; Canil, 1999; Toplis
 385 and Corgne, 2002; Williams et al., 2004; Dauphas et al., 2009; Mallmann and O'Neill, 2009,
 386 2013; Davis et al., 2013) decreases with increasing T_p (Fig. 4), and is qualitatively consistent
 387 with the positive $\delta^{51}V-Na_{8.0}$ correlation in global MORB (Wu et al., 2018; Qi et al., 2019;

388 Novella et al., 2020; $\text{Na}_{8,0}$ is the Na_2O content corrected for crystallisation to $\text{MgO} = 8 \text{ wt } \%$, and
389 therefore removes the variable chemical effects of low-pressure processes allowing for more direct
390 comparison between samples). However, the predicted magnitude of V isotope variability with
391 T_p is smaller than the range of MORB V isotope data presented by Wu et al. (2018). Due to the
392 enrichment of melts with heavy isotopes, MORB should have a higher $\delta^{51}\text{V}$ than fertile mantle.
393 This is in agreement with observed data: Wu et al. (2018) measure an average global MORB
394 value of $-0.84 \pm 0.10 \text{ ‰}$ (2 S.D.), Qi et al. (2019) propose BSE is $-0.91 \pm 0.09 \text{ ‰}$ (2 S.D.) based
395 on the isotopic composition of peridotite xenoliths and komatiites. Similar behaviour is observed
396 for Fe isotopes: global MORB, after correction for fractional crystallisation, has $\delta^{57}\text{Fe} = 0.10 \pm$
397 0.01 ‰ (2 S.E.); BSE is estimated at $0.05 \pm 0.01 \text{ ‰}$ (2 S.E.) from primitive lherzolites (Sossi
398 et al., 2016). The magnitude of the modelled $\Delta^{51}\text{V}_{\text{melt-source}}$ and $\Delta^{57}\text{Fe}_{\text{melt-source}}$ of aggregate
399 KLB1 melts at $T_p \approx 1300 \text{ °C}$ ($\sim 0.055 \text{ ‰}$ for both isotope systems) is similar to the 0.07 ‰
400 and 0.05 ‰ difference observed between BSE and average MORB for V and Fe, respectively.
401 However, we note that the large uncertainties on the $\delta^{51}\text{V}$ values estimated for the BSE and
402 MORB mean that the $\Delta^{51}\text{V}_{\text{melt-source}}$ predicted by the model would match the BSE and MORB
403 data for all T_p .

404 Fig. 4b shows how the magnitude of isotopic variability across the T_p range compares to
405 typical long-term analytical precision for each isotope system (see Table 1 for the errors used),
406 highlighting that the magnitude of variability in the isotope ratio of melts produced between
407 $T_p = 1300 \text{ °C}$ and 1530 °C (a reasonable T_p range for melt generation at mid-ocean ridges and
408 mantle plumes) is currently analytically unresolvable. Even considering analytical uncertainty
409 in individual measurement sessions (i.e., the typical sample 2 S.E. rather than the long-term
410 uncertainty; Table 1), the temperature-related variations associated with varying melt $\text{Fe}^{3+}/\text{Fe}_T$
411 and modal mineralogy are not resolvable in the modelled average melt isotope ratios. The
412 analytical precision therefore currently precludes mantle temperature variations in the global
413 MORB dataset being resolved; this is consistent with Sossi et al. (2016) finding a homogenous
414 global MORB $\delta^{57}\text{Fe}$, despite global MORB major element chemistry indicating 200 °C variability
415 (Gale et al., 2014). We note that the Mg and Cr stable isotope compositions of basalts are
416 unlikely to be sensitive to temperature effects even with improved analytical precision; significant
417 improvements in V and Ca analytical precision (both currently systems with relatively large
418 errors), down to 2 S.D. of $< 0.04 \text{ ‰}$, offer the most potential to resolve the predicted temperature-

419 related variations.

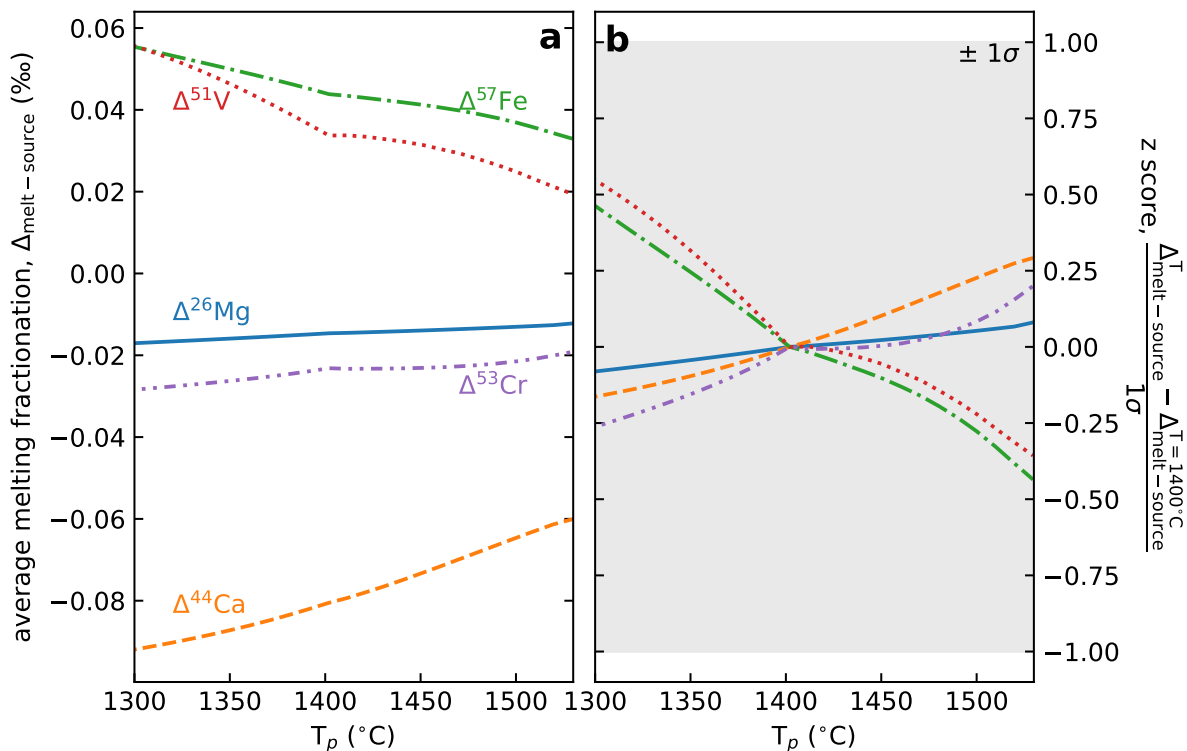


Figure 4: a) The estimated average isotopic fractionation of accumulated model peridotite melts relative to the source, across reasonable mid-ocean ridge and mantle plume temperatures. Mantle potential temperature, T_p , is calculated as in Fig. 3, and the average melt fractionation is taken as the fractionation at the mean melting pressure at each temperature (see details in the text, Fig. 3, Appendix Section 6). b) The z score is calculated as $(\Delta_{\text{melt-source}}(T) - \Delta_{\text{melt-source}}(T = 1400^\circ\text{C})) / 1\sigma$ typical analytical precision, for each isotope system. This calculation reflects how the magnitude of expected isotopic variability compares to reasonable measurable isotopic differences: a z score of 1 represents predicted isotopic variability resolvable at typical 1σ uncertainty. The errors used are the representative long-term analytical uncertainty for each isotope system; see Table S2.

420 4 Lithological heterogeneity

421 We next investigate the use of stable isotopes as tracers of lithological heterogeneity in the
 422 mantle, using two pyroxenite lithologies, MIX1G (silica deficient, SD pyroxenite) and G2 (silica
 423 excess, SE pyroxenite). Silica-deficient pyroxenites may be the most representative type of
 424 pyroxenites found in the convecting mantle (the majority of pyroxenite xenoliths sampled in
 425 OIB are SD; Lambart et al., 2016 – natural SD pyroxenites are considered to be formed by
 426 metasomatism of peridotite by eclogite melts, as shown in experiments; Yaxley and Green, 1998;
 427 Yaxley, 2000; Mallik and Dasgupta, 2012). Natural SE pyroxenites include eclogites (Lambart
 428 et al., 2016).

4.1 Mineralogic control on melt isotope composition

Fig. 5 shows the pseudosections and melt isotope ratios along peridotite isentropic decompression melting paths (see Fig. 2) for the two pyroxenite lithologies considered in this study. We have used peridotite isentropes since the thermal budget of multi-lithologic mantle will be buffered by melting of the dominant lherzolitic peridotite lithology, not pyroxenite melting (Phipps Morgan, 2001; Katz and Rudge, 2011). For example, peridotite (lherzolite + harzburgite) is estimated to make up $\geq 95\%$ of the multi-lithologic mantle in the MORB mantle source (e.g., Hirschmann and Stolper, 1996; Matthews et al., 2016), and $> 90\%$ of the Icelandic mantle source (e.g., Shorttle et al., 2014). While a multi-lithologic mantle will not follow pure peridotite isentropic decompression paths, we consider this approach the most realistic reflection of mantle melting processes. Isentropic decompression melting paths for pure MIX1G and G2 are shown in Figs. S10 and S11 for reference. The melt isotope ratios shown in Fig. 5 are all calculated relative to a bulk system isotopic composition of 0‰, for convenience. In this respect Fig. 5 is strictly looking at the direct lithological (i.e., mineralogic) control on the stable isotope composition of melts independent of any additional source compositional effects that may have been imprinted by the processes forming the source in the first place. For example, MORBs (and hence MORB-derived mantle components) are isotopically distinct from the BSE, due to their formation by partial melting (Fig. 2), crystal fractionation, and low temperature alteration by the time of their subduction (e.g., as for $\delta^{57}\text{Fe}$; Teng et al., 2013; Sossi et al., 2016, $\delta^{44}\text{Ca}$; Zhu et al., 2018; Chen et al., 2020b; Antonelli and Simon, 2020 and $\delta^{51}\text{V}$; Wu et al., 2018). If MORB, or MORB-derived components, then form part of the mantle source during a subsequent melting episode, the new melts will be isotopically fractionated compared to lherzolite melts both because of their distinct mineralogy and their inherited initial isotope compositions. This subject is considered further in Sections 4.2 and 5.

We first consider how the isotopic composition of pyroxenite melts compares to peridotite melts (Table 1) in the context of the two questions posed in Section 3: the magnitude of melt-source isotopic fractionation, and how melt-source isotopic fractionation varies with progressive melting. Generally, the pyroxenite melts show a wider range of melt isotope ratios than peridotite melts for the same T_p (compare Figs. 2 and 5; see also Fig. S13), because the pyroxenite melts represent a wider range of melting degrees than the peridotite and a greater proportion of

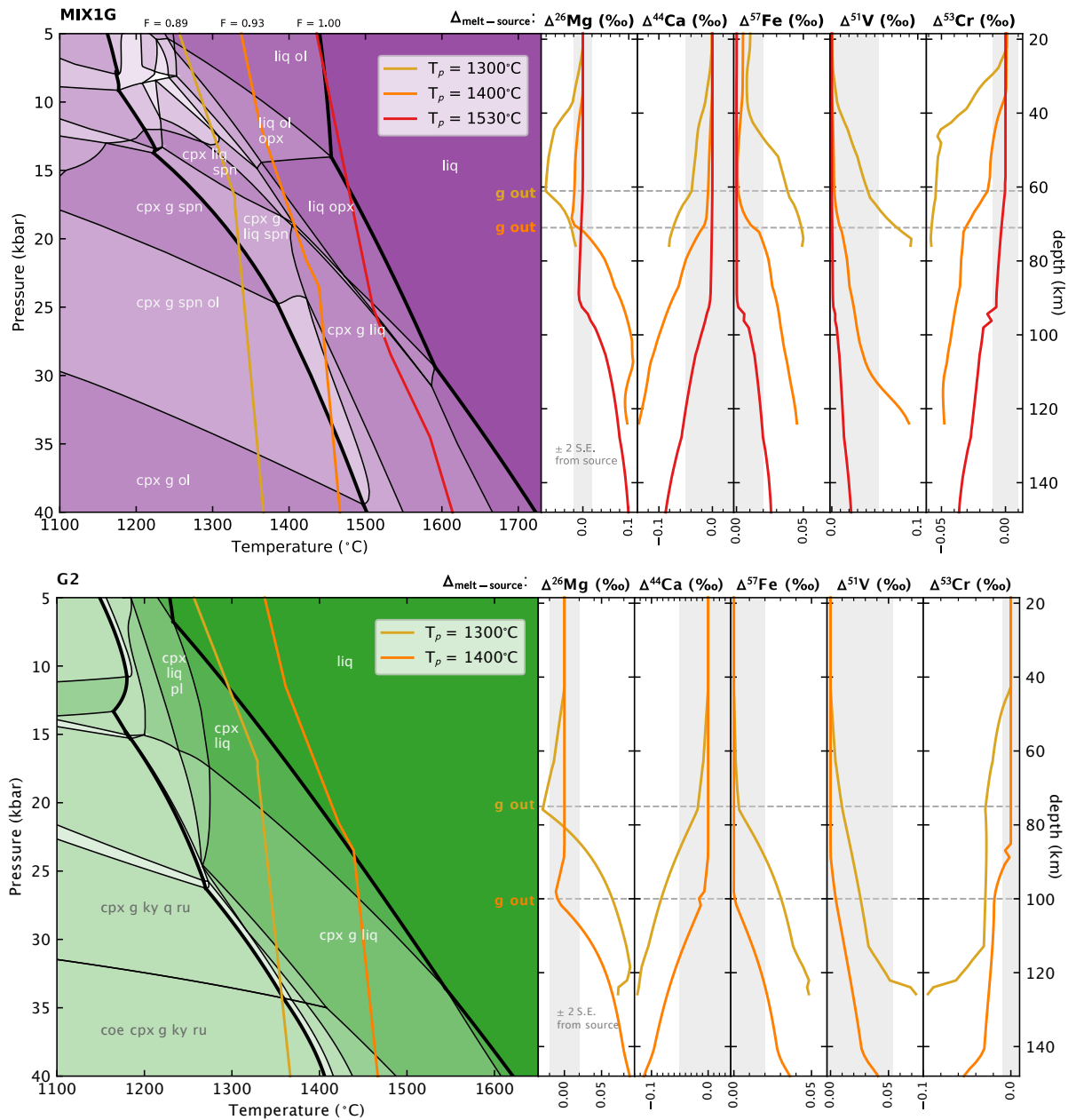


Figure 5: As in Fig. 2, the left hand panels show the pseudosections calculated in THERMOCALC for MIX1G (top) and G2 (bottom) pyroxenites, with the solidus and liquidus marked in bold. Isentropic melting paths for peridotite calculated using THERMOCALC for $T_p = 1300^{\circ}\text{C}$, 1400°C and 1530°C are superimposed. For MIX1G, maximum melt fractions (F) are marked at the top of each isentrope (for G2 $F = 1$). The right hand panels show the modelled instantaneous equilibrium melt isotopic fractionation relative to the bulk isotopic composition ($\Delta_{\text{melt-source}}$) along the isentropes. The $T_p = 1530^{\circ}\text{C}$ isentrope has not been shown on G2 as the isentrope sits hotter than the liquidus. Key mineralogical controls on the isotopic composition of the melt are highlighted (colour-coded by the isentropic path they refer to). Vertical grey bars show the source composition (0 ‰ for convenience) with typical 2 S.E. sample error for each isotope system – see Table S2 for details.

459 these melts are generated in the garnet stability field (garnet generally imparts a large isotopic
 460 fractionation compared to olivine and pyroxene; see full results in Appendix). While we find
 461 that V stable isotopes are predicted to have no direct sensitivity to different source lithologies

462 (because the modelled $\Delta^{51}\text{V}_{\text{melt-source}}$ for peridotite and pyroxenite lithologies are both similar
463 to each other and well within current analytical precision), our results show that Fe, Cr, Ca and
464 Mg could show some direct sensitivity to mantle lithology (Fig. S13). Progressive melting of a
465 garnet-bearing source (such as G2, or MIX1G at depth) produces the opposite sense of melt-
466 source isotope fractionation in Mg and Ca isotopes (both in direction, i.e., positive or negative
467 $\Delta_{\text{melt-source}}$, and trend, i.e., $\Delta_{\text{melt-source}}$ becoming more negative or more positive; Figs. 2, 5)
468 due to the presence of isotopically light and heavy garnet respectively ($\Delta^{26}\text{Mg}_{\text{garnet-source}} \sim$
469 -0.3‰ ; $\Delta^{44}\text{Ca}_{\text{garnet-source}} \sim 0.25\text{‰}$). Therefore, if lithological heterogeneity (specifically, the
470 proportion of garnet in a melt source) is the principal control on the chemical heterogeneity of
471 a suite of basalts, the two isotope systems should show a negative correlation. This relationship
472 is shown in Fig. 6 for MIX1G melts. Fig. 6 also shows how the isotopic compositions of the
473 melts might be predicted to relate to a more traditional geochemical tracer of garnet in the
474 source region, Sm/Yb ratio (e.g., Pertermann et al., 2004; details of the calculation in Appendix
475 Section 7), and similar patterns would be predicted for other garnet tracers such as Dy/Yb
476 (Fig. S10). The correlation between Ca isotopes and garnet-sensitive trace element ratios is in
477 agreement with measurements and predictions by Wang et al. (2019); Antonelli et al. (2021).
478 Magnesium-Ca isotope systematics of basalts could therefore have future potential as a tracer
479 of garnet (hence, depth of melting and/or lithological heterogeneity) in the mantle source, as
480 predicted for both stable isotope systems individually (e.g., Stracke et al., 2018; Kang et al.,
481 2019; Wang et al., 2019; Dai et al., 2020; Antonelli et al., 2021).

482 However, further work on the effect of garnet on $\Delta^{44}\text{Ca}_{\text{melt-source}}$ by taking garnet composi-
483 tion into account, following measurements by Chen et al. (2020a) and predicted composition-
484 dependent β^{Ca} from Antonelli et al. (2019), is required. We note that Antonelli et al. (2019) pre-
485 dict that composition-dependent isotope fractionation effects ('concentration effects') for other
486 mineral solid solutions (olivine, pyroxene) are smaller using their PBE (Perdew, Burke, Ernzerhof)
487 functionals to calculate β relative to previous studies that have employed different ap-
488 proaches (LDA functionals; e.g., Feng et al., 2014; Wang et al., 2017), which supports only small
489 concentration effects in garnet. Currently, typical analytical precision on Ca isotope measure-
490 ments (long-term external 2 S.D. of $\delta^{44}\text{Ca} = 0.14\text{‰}$ on the SRM915b standard, although typical
491 2 S.E. on measurements of individual samples is smaller, average 0.05‰, Chen et al., 2019; Dai
492 et al., 2020) precludes some of the predicted equilibrium variability in $\Delta^{44}\text{Ca}_{\text{melt-source}}$ due to

493 the presence of garnet from being resolved (Fig. 6). However, as analytical precision improves,
 494 the garnet signature in Mg-Ca isotopes could be further tested.

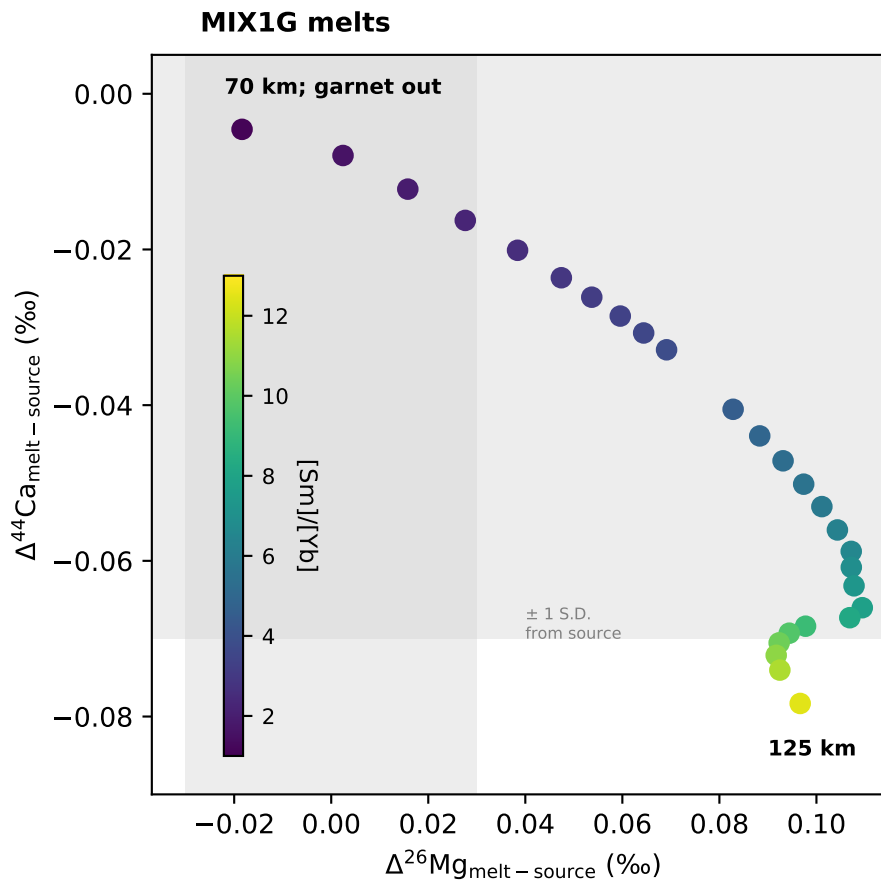


Figure 6: Instantaneous equilibrium melt-source Mg and Ca stable isotope fractionation for MIX1G melts produced at depths greater than 70 km (i.e., where garnet is present in the source) along a peridotitic isentropic decompression melting path ($T_p = 1400^\circ\text{C}$). Grey shading shows the source (0‰) \pm typical 1 S.D. long-term analytical error (Table S2). Details of trace element modelling in Section 7 of the Appendix.

495 However, whether melts from pyroxenite lithologies will ultimately be resolvable in an erupted
 496 melt will depend on two factors in addition to the direct mineralogic control: the bulk isotope
 497 composition of a pyroxenite lithology relative to peridotite, and the amount of each element
 498 that each of the melting lithologies contributes to the final melt.

499 4.2 Resolving pyroxenite source components in erupted melts

500 Fig. 7 shows the results of mass balance calculations for the aggregate bulk melt-source isotope
 501 fractionation of a multi-lithologic mantle relative to a pure pyroxenite source, with varying
 502 proportions of pyroxenite. The peridotite lithology is taken to have a bulk isotope composition

503 equal to BSE (see Fig. 5 caption for data sources; isotopic reference materials used are DSM3,
504 SRM915a, IRMM-014, AA, NIST979 for Mg-Ca-Fe-V-Cr respectively); pyroxenite lithologies
505 are taken to have a bulk isotope composition of average MORB for each isotope, except for
506 Cr where no known MORB data are published and the modelled aggregate KLB1 melt isotope
507 composition for $T_p = 1300^\circ\text{C}$ is used. Full details of the mass balance calculation are found in
508 Section 9 of the Appendix. The right hand panels show the same data referenced to the pure
509 peridotite melt, allowing a simple assessment of the required analytical precision if pyroxenite
510 is to be detected by each system. Fig. S14 shows the same data as Fig. 7, instead scaled to
511 typical current 1 S.D. analytical precision. Each isotope system is shown with a range of isotope
512 compositions, reflecting two end-member conditions of thermal behaviour of the pyroxenite in
513 multi-lithologic mantle. The bold line reflects a pyroxenite decompressing along a peridotite
514 decompression path, as used in Fig. 5. This approximation will be most valid for small (
515 10%) pyroxenite fractions. The other end-member is pyroxenite following its own decompression
516 melting path (the paths shown in Figs. S10 and S11). We acknowledge that these results are still
517 an approximation, as peridotite is assumed to be melting along its own geotherm regardless of
518 pyroxenite fraction. From Figs. 7 and S14, we find that none of the considered isotope systems
519 will be able to resolve the presence of minor (less than $\sim 5\%$) pyroxenite components in mantle
520 source regions, given levels of current analytical precision as given in Table S2.

521 4.2.1 Isotope systems with minimal sensitivity to pyroxenite: V and Ca

522 Our results show that the presence of even moderate amounts of pyroxenite in the mantle
523 source cannot be resolved using V stable isotopes, assuming a MORB-like bulk pyroxenite
524 isotope composition. This result is consistent with Novella et al. (2020) who show that $\delta^{51}\text{V}$
525 cannot detect pyroxenite in the source of Reykjanes Ridge basalts. Given the similar predicted
526 V isotope compositions of pyroxenite and peridotite melts, with pyroxenite fractions up to 40%
527 producing bulk melt isotope compositions that only differ from a peridotite source by $< 0.03\text{‰}$
528 (Fig. 7), improvements in V stable isotope analytical precision are unlikely to make V isotopes
529 sensitive to the pyroxenite lithologies modelled here.

530 We also predict that partial melting of peridotite and pyroxenite produces unresolvable Ca iso-
531 topic fractionation relative to the source (for peridotite, $\Delta^{44}\text{Ca}_{\text{melt-source}}$ smaller than -0.09‰ ;

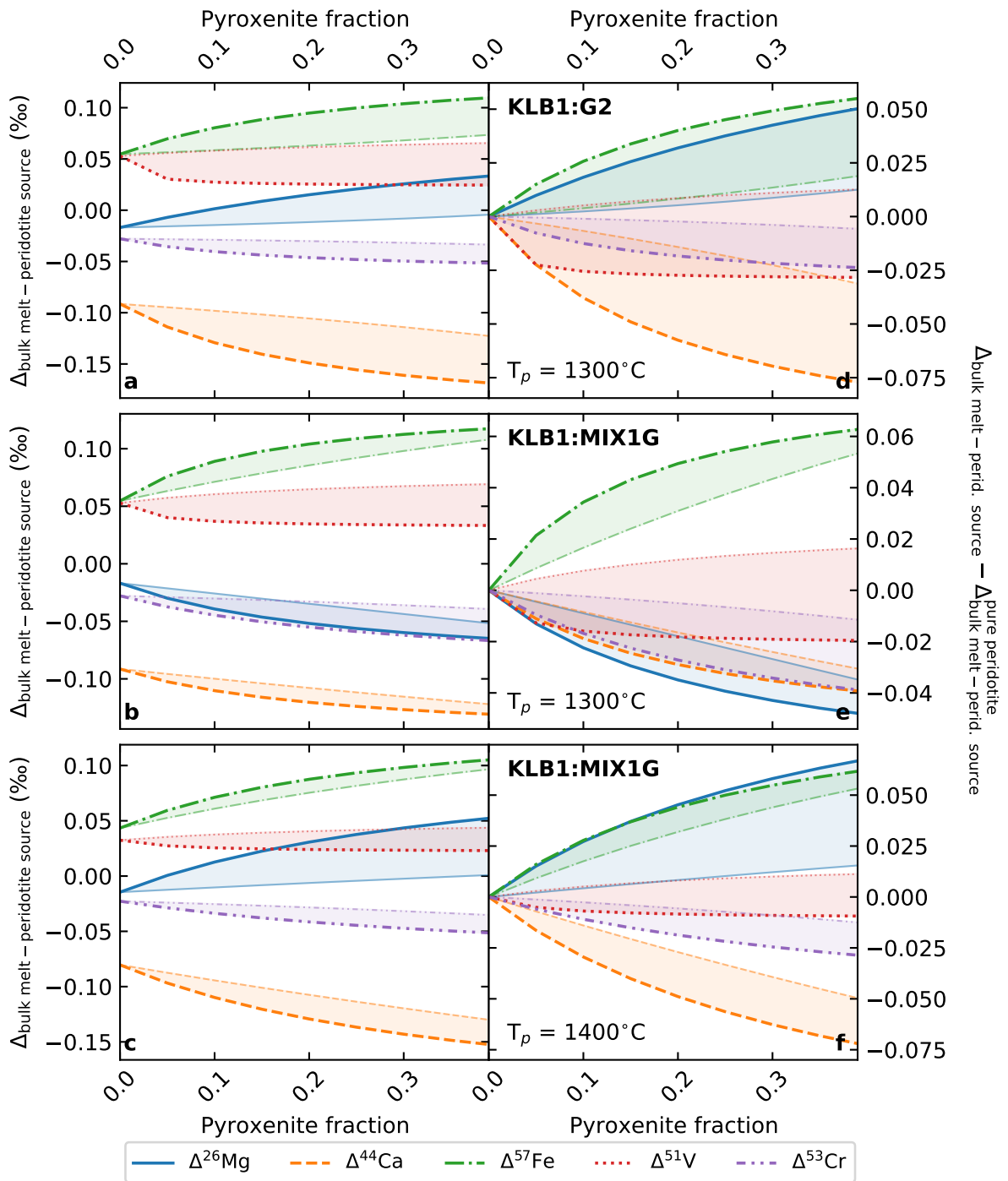


Figure 7: Aggregate bulk melt isotope fractionation (relative to a pure peridotite source, left hand panels; relative to pure peridotite melt; right hand panels) for multi-lithologic mantle. At each T_p , the aggregate melt composition (including isotope composition, and element proportion) at the average pressure and temperature of melting for each lithology is used to calculate the resulting bulk melt isotope composition by mass balance. A range of bulk melt-source isotope fractionations are given for each isotope system, reflecting two endmembers for the thermal behaviour of pyroxenite upwelling in a peridotite host. The bolder line represents the case for pyroxenite following peridotite geotherms (as used throughout this study); the paler line reflects pyroxenite following its own geotherm, as in Figs. S10 and S11.

532 Fig. 2, in agreement with Kang et al., 2017; Chen et al., 2019 but smaller than $\Delta^{44}\text{Ca}_{\text{melt-residue}}$
533 $= -0.09 - -0.2\text{‰}$ suggested by Zhu et al., 2018; Ionov et al., 2019; Chen et al., 2020b). Al-
534 though bulk MORB-like pyroxenite may have an isotopically lighter composition than peridotite
535 (see Fig. 8), which is further amplified by negative melt-source isotopic fractionation to produce
536 pyroxenite melts that are isotopically lighter than peridotite melts, Fig. 7 suggests that Ca
537 isotope ratios measured with current typical long-term uncertainty will only be able to detect
538 MORB-like pyroxenite in a mantle source if the proportion of pyroxenite is greater than 30 %
539 (which may be unrealistic of multi-lithologic mantle, e.g., Sobolev et al., 2005; Pietruszka et al.,
540 2013; Shorttle et al., 2014; Matthews et al., 2016; Brown et al., 2020). When considering a cur-
541 rent best 2 S.E., this detection limit only improves to $\sim 20\%$ of pyroxenite in the source. The
542 small variability in melt isotope ratios with increasing pyroxenite fraction compared to resolv-
543 able variability is consistent with DMM, EM1 and HIMU basalts recording no resolvable $\delta^{44}\text{Ca}$
544 variability (Valdes et al., 2014), and measured mantle pyroxenites and associated peridotites
545 recording indistinguishable $\delta^{44}\text{Ca}$ (Dai et al., 2020). However, even a small improvement in
546 analytical precision could enable greater sensitivity to pyroxenite fraction: for example Figs. 7a
547 and S14 show that $\sim 15\%$ of G2 pyroxenite in a mantle source produces an isotopic offset in the
548 bulk melt of -0.05‰ relative to a pure peridotite source. Given that the typical 2 S.E. for $\delta^{44}\text{Ca}$
549 is around 0.05‰ (Table S2), there is potential for these small pyroxenite fraction signatures to
550 be theoretically resolved with only minor improvements in analytical precision. We note that
551 currently resolvable variability in the measured $\delta^{44}\text{Ca}$ of igneous rocks is found (specifically in
552 ultramafic rocks including dunites and komatiites), but that pyroxenite source contributions are
553 not proposed to be the origin of this $\delta^{44}\text{Ca}$ heterogeneity (Antonelli and Simon, 2020).

554 4.2.2 Distinguishing lithologies: Fe, Mg, Cr

555 The Fe isotope variability of MORB and OIB has been previously linked to source heterogeneity
556 (e.g., Williams and Bizimis, 2014; Konter et al., 2016; Nebel et al., 2019; Sun et al., 2020; Zhong
557 et al., 2021), through the inherent (initial; presumably derived from its formation or protolith)
558 Fe isotope compositions of the different source components, rather than directly through any
559 large mineralogy-dependent partial melting Fe isotope fractionation. Consistent with these
560 conclusions, we find that the calculated partial melting $\Delta^{57}\text{Fe}$ fractionations are small ($\lesssim 0.06\text{‰}$)

561 for peridotite and pyroxenite lithologies (in agreement with previous studies, e.g., Dauphas et al.,
562 2009; Williams and Bizimis, 2014; Sossi et al., 2016; Gleeson et al., 2020; Soderman et al., 2021).
563 We also find that the main driver of melt Fe isotope ratio sensitivity to pyroxenite fraction
564 (Fig. 7) is the expectation that MORB-like pyroxenite will have a heavier bulk Fe isotope
565 ratio than peridotite, which in combination with a positive melt-source isotopic fractionation
566 drives pyroxenite melts to heavier isotope compositions than peridotite melts. The contribution
567 of high $\delta^{57}\text{Fe}$ pyroxenite melts (from both SD and SE pyroxenites) compared to melting of a
568 pure peridotite source is resolvable at greater than 8–9% pyroxenite with current typical 1 S.D.
569 (Figs. 7, S14), making Fe isotopes one of the most sensitive systems to source considered here,
570 although it should be noted that equilibrium isotopic fractionation associated with pyroxenite
571 sources cannot generate the full range of natural stable isotope OIB data from a single stage
572 melting model (Soderman et al., 2021, see Section 5). For G2 pyroxenite however, the choice of
573 thermal behaviour of the pyroxenite (i.e., whether the upwelling pyroxenite component follows
574 a peridotite or pyroxenite decompression path) strongly affects the sensitivity to pyroxenite
575 fraction. At high proportions of G2 pyroxenite, when a pyroxenite decompression path becomes
576 more appropriate, the multi-lithologic mantle melt may no longer be resolvable from a pure
577 peridotite source (Fig. 7d).

578 In agreement with most existing studies (e.g., Teng et al., 2007, 2010; Huang et al., 2011a;
579 Stracke et al., 2018), but in contrast to Zhong et al. (2017), we find that partial melting of
580 peridotite produces unresolvable Mg isotopic fractionation ($\Delta^{26}\text{Mg} < 0.02\text{‰}$; Fig. 2). Melting
581 of MIX1G and G2 pyroxenite lithologies produces larger magnitude melt-source fractionations
582 than peridotite (Fig. 5) due to the presence of isotopically light garnet and isotopically heavy
583 clinopyroxene in the source, in varying proportions over P-T space (Fig. S1). By compari-
584 son, olivine and orthopyroxene, which dominate peridotite melting, have less extreme $\delta^{26}\text{Mg}$
585 compositions. The presence of garnet producing more extreme isotopic fractionations than a
586 non-garnet-bearing source is consistent with melting models from Stracke et al. (2018); however,
587 a key difference is that the modelled melting of garnet pyroxenite by Stracke et al. (2018) pro-
588 duces isotopically light melts (i.e., a negative $\Delta_{\text{melt-source}}$), compared to the mostly isotopically
589 heavy ($\Delta_{\text{melt-source}} > 0\text{‰}$) melts at depths where garnet is present in the source in this model
590 (Fig. 5). This difference in behaviour is due to the choice of initial $\alpha_{\text{melt-source}} = 1$ in Stracke
591 et al. (2018). We do not use this initial condition, and in the melting assemblage considered

592 here, garnet is largely a residual phase (i.e., not contributing to melting), and therefore initial
593 melts are isotopically heavy relative to the source. Both models agree that with progressive
594 melting of isotopically light garnet in the source, melts will become isotopically lighter. Some
595 isotopically light (i.e., $\Delta^{26}\text{Mg}_{\text{melt-source}} < 0\text{‰}$) melts are produced by melting G2 and MIX1G
596 in this model, but these are at depths where garnet is nearly or completely absent from the
597 stable assemblage (Fig. 5). It is this garnet effect that is responsible for the opposite behaviours
598 of $\Delta^{26}\text{Mg}_{\text{melt-source}}$ in Fig. 7d, f compared to Fig. 7e: the aggregate pyroxenite melt sits in
599 the spinel stability field for MIX1G at $T_p = 1300\text{°C}$ (Fig. 7b), but in the garnet stability
600 field for MIX1G at $T_p = 1400\text{°C}$ (Fig. 7c) and G2 at $T_p = 1300\text{°C}$ (Fig. 7d). Melting of
601 multi-lithologic mantle (with MIX1G pyroxenite) is resolvable from pure peridotite at current 1
602 S.D. at $\sim 10\text{--}15\%$ pyroxenite fraction (Fig. S14); our results show that Mg isotopes are more
603 sensitive to MIX1G than G2 pyroxenite. With continued improvements in analytical precision,
604 such as the 0.02‰ 2 S.E. reported by Wang et al. (2021), pyroxenite fractions $< 10\%$ could be
605 resolved in the Mg isotope composition of the bulk melt, as a 10% pyroxenite fraction produces
606 a bulk melt $> 0.02\text{‰}$ different from a pure peridotite melt (Figs. 7e, f).

607 Chromium stable isotope fractionation during partial melting of different lithologies is relatively
608 understudied compared to Mg and Fe, however Fig. 7 shows that melt Cr isotopic composi-
609 tion could show resolvable sensitivity to the fraction of pyroxenite in the source. This is due
610 both to variable modelled melting fractionations dependent on lithology, and comparatively
611 small analytical uncertainty. Equilibrium Cr stable isotope fractionation is controlled by both
612 bonding environment and redox conditions (e.g., Shen et al., 2018, 2020; Bonnand et al., 2020).
613 Chromium stable isotope fractionation during partial melting of peridotite is probably small
614 (magnitude of $\Delta^{53}\text{Cr} < 0.07\text{‰}$; Schoenberg et al., 2016; Xia et al., 2017; Shen et al., 2018; Bon-
615 nand et al., 2020), in agreement with our results. Isotopically heavy spinel, where present, has a
616 large influence on the isotopic fractionation between melt and residue due to the compatibility
617 of Cr in spinel, and this has been used to explain observations from peridotite xenoliths that
618 partial melting appears to generate isotopically light melts ($\Delta^{53}\text{Cr}_{\text{melt-source}} < 0$) and heavy
619 residues (Xia et al., 2017; Shen et al., 2020). However, due to the small area of the KLB1
620 melting regime in which spinel is stable in our model (Figs. 2, S4), spinel does not control the
621 magnitude of stable isotope fractionation associated with melting in this case. We find a slightly
622 larger magnitude of melting-induced Cr stable isotope fractionation for MIX1G pyroxenite than

623 for peridotite (as large as -0.06‰). Like MIX1G, G2 pyroxenite melts all have a negative
624 $\Delta^{53}\text{Cr}$ with respect to the source, with melting fractionations as large as $\Delta^{53}\text{Cr} = -0.1\text{‰}$ (Fig.
625 5). This is because the main Cr bearing phase in much of the melting region is isotopically
626 heavy (relative to olivine) clinopyroxene, and this produces consistently isotopically light melts,
627 consistent with observations by Xia et al. (2017); Shen et al. (2020). Therefore, combined with
628 a predicted bulk pyroxenite isotope composition that is also isotopically lighter than peridotite,
629 bulk melts are predicted to become progressively Cr-isotopically lighter with increasing pyrox-
630 enite fraction in the mantle source. This effect is generally small ($< 0.04\text{‰}$ offset from pure
631 peridotite melts even at 30% pyroxenite; Fig. 7) and at current typical 2 S.D. is only resolvable
632 at more than 20% MIX1G in the source at comparatively low temperatures ($T_p = 1300\text{°C}$).
633 However, as with Mg, recent improvements in analytical precision for Cr isotopes (Sossi et al.,
634 2018; Wagner et al., 2021) may permit greater sensitivity to pyroxenite components: for a 2
635 S.E. of 0.01‰ , multi-lithologic mantle melts become resolvable from pure peridotite melts at
636 around 8% pyroxenite upwards, assuming upwelling pyroxenite follows the peridotite geotherm
637 (Figs. 7, S14).

638 5 Comparisons to natural data

639 Fig. 8 shows how literature MORB and OIB data compare to modelled melts. The modelled
640 melts are filtered for those produced at cooler temperatures than the $T_p = 1530\text{°C}$ peridotite
641 isentrope for a given pressure, to limit the contribution of unrealistic high melt fraction and high
642 pressure melts. The same potential temperature isentrope has been chosen for all lithologies
643 for consistency, although because of its fusibility this filter includes some 100% melts of G2 at
644 all pressures. Natural samples are filtered for those with $7.2 < \text{MgO (wt\%)} < 16$ to minimise
645 the effects of fractional crystallisation and olivine accumulation, which can result in measurable
646 isotopic fractionations in several isotope systems (e.g., Sossi et al., 2016; McCoy-West et al.,
647 2018; Nebel et al., 2019; Shen et al., 2020). The KLB1 melts are shown as the calculated
648 melting fractionation applied to an accessible mantle isotopic composition (presumptive bulk
649 silicate earth, BSE); G2 and MIX1G melts are shown as the melting fractionation applied to
650 the average fresh MORB isotopic composition, since G2 is a MORB-like eclogite in composition.
651 For Cr, where we can find no published MORB data, we use the isotope ratio of BSE + the

652 aggregate KLB1 melt isotope ratio for $T_p = 1300^\circ\text{C}$ (Fig. 4) to predict the isotopic composition
653 of fresh MORB, and reference G2 and MIX1G melts to this value.

654 **5.1 Can peridotite melting explain MORB Mg-Ca-Fe-V isotope ratios?**

655 The mean MORB Mg isotope ratio agrees well with the average peridotite melt produced (BSE
656 plus the melt-source isotope fractionation) when $T_p = 1300^\circ\text{C}$. The mean MORB Ca and V
657 isotope ratios are also well within analytical uncertainty of the average peridotite melt produced
658 (BSE plus the melt-source isotope fractionation) when $T_p = 1300^\circ\text{C}$ (see Figs. 4, 8), since both
659 Ca and V have relatively large analytical uncertainties. For Fe, the isotope ratios of mean
660 MORB and average peridotite melt at $T_p = 1300^\circ\text{C}$ agree just within analytical uncertainty,
661 requiring almost the maximum long-term error on natural sample measurements to match. Con-
662 sidering typical 2 S.E. analytical precision, mean MORB $\delta^{57}\text{Fe}$ are slightly higher than predicted
663 by the model even when considering only extremely small melt fractions (hence maximum par-
664 tial melting isotopic fractionation). This discrepancy could be due to the effect of fractional
665 crystallisation on the Fe isotope ratio of melts, since even high MgO MORB are almost never
666 primary liquids (e.g., Langmuir et al., 1992), and fractional crystallisation produces isotopically
667 heavy residual melts compared to the primary liquid (Teng et al., 2008; Chen et al., 2019b).
668 Sossi et al. (2016) use a fractional crystallisation correction on the MORB dataset and calculate
669 an average primitive MORB liquid with $\delta^{57}\text{Fe} = 0.10\text{‰}$, compared to 0.15‰ in the uncorrected
670 (and non-MgO filtered) dataset (maximum correction on individual samples is $\sim 0.07\text{‰}$). This
671 result matches well with our calculated average peridotite melt at $T_p = 1300^\circ\text{C}$. However, the
672 fractional crystallisation correction applied by Sossi et al. (2016) (e.g., $\Delta^{57}\text{Fe}_{\text{ol-melt}} \sim -0.18\text{‰}$
673 at $T = 1200^\circ\text{C}$) is likely to represent the maximum limit of isotopic fractionation during olivine
674 differentiation (Soderman et al., 2021).

675 In addition to comparing our results with mean MORB isotope compositions, we also examine
676 how the ranges of measured MORB stable isotope compositions (for the systems Mg-Ca-Fe-
677 V) compares to the spread of modelled melts. MORB samples have a wider range of isotopic
678 compositions than the equilibrium KLB1 model melts (Fig. 8). In the Ca isotope system,
679 some of this variability may be attributed to analytical uncertainty on measurements: here the
680 width of the measured distributions of MORB data are within long-term analytical precision of

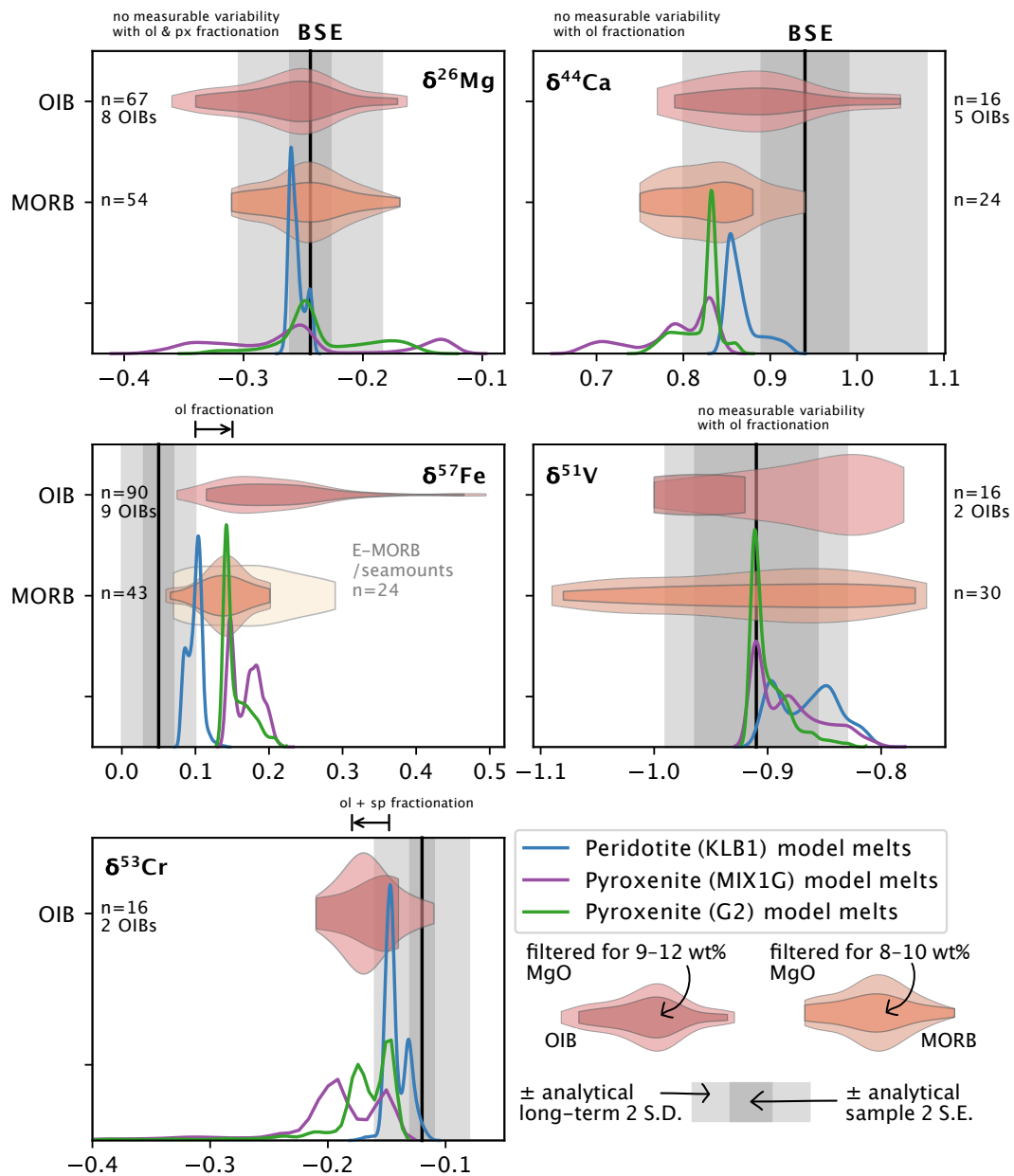


Figure 8: Literature MORB (N- and T- MORB, if classification is known) and OIB data for the isotope systems (data sources in Table S1), compared to the model instantaneous melts (all data shown as ‰ variation from the reference materials outlined in Fig. 1). Samples are filtered for $7.2 < \text{MgO wt \%} < 16$, but included if MgO is not known. ‘n’ refers to the total number of samples included, and the number of OIBs given shows how many unique ocean island settings are sampled. The central darker violin shows a more limited subset of samples: 9–12 wt % MgO in OIB, 8–10 wt % MgO in MORB, to highlight isotope compositions of the most primitive lavas. Ca and Mg isotope data is filtered to exclude samples with $^{87}\text{Sr}/^{86}\text{Sr} > 0.7037$, which excludes samples with $> 1\%$ carbonate component in the source (Huang et al., 2011b). Modelled melts are filtered for those cooler than the $T_p = 1530^\circ\text{C}$ peridotite isentrope at any given pressure. KLB1 melts are shown assuming a bulk source isotope composition of BSE; G2 and MIX1G melts are shown assuming a bulk source isotope composition of average MORB (for Cr, the isotope ratio of BSE + the aggregate KLB1 melt isotope ratio for $T_p = 1300^\circ\text{C}$ is used due to the lack of published MORB data). BSE values are $\delta^{26}\text{Mg} = -0.24\text{‰}$ (Stracke et al., 2018), $\delta^{44}\text{Ca} = 0.94\text{‰}$ (Kang et al., 2017; Chen et al., 2019), $\delta^{57}\text{Fe} = 0.05\text{‰}$ (Sossi et al., 2016), $\delta^{51}\text{V} = -0.91\text{‰}$ (Qi et al., 2019), $\delta^{53}\text{Cr} = -0.12\text{‰}$ (Jerram et al., 2020). The pale outer shaded grey region shows typical long term 2 S.D. analytical uncertainty in recent studies, the darker inner shaded grey region shows typical sample 2 S.E. (S.E. = S.D./ \sqrt{n}) in recent studies (see Table 1). The fractionation arrows for Fe and Cr isotopes represent the direction and approximate magnitude of isotopic fractionation associated with differentiation to form samples with minimum 7 wt % MgO. Fractionation information from Sossi et al. (2016) for Fe, Teng et al. (2007, 2010); Wang et al. (2021) for Mg, Zhang et al. (2018) for Ca, Shen et al. (2020) for Cr, Ding et al. (2020) for V.

681 the modelled KLB1 melt compositions (although, the errors on individual sample Ca isotope
682 measurements are smaller and hence cannot account for all the difference between MORB and
683 peridotite melts). In the Fe, Mg and V systems, the variability of measured MORB isotope com-
684 positions relative to the single-stage equilibrium peridotite melts is greater than can be explained
685 by analytical uncertainty alone (Sossi et al., 2016) – this is particularly evident in $\delta^{51}\text{V}$ (Fig. 8).
686 Stable isotope fractionation during differentiation may play a role, as discussed for Fe isotopes
687 above, but V isotopes are not thought to be significantly fractionated by olivine differentiation
688 (Prytulak et al., 2013; Ding et al., 2020). Fractional crystallisation of clinopyroxene and Fe-Ti
689 oxides, although unlikely for the MgO-filtered data here, is shown to drive melts to higher $\delta^{51}\text{V}$
690 (Wu et al., 2018; Ding et al., 2020), which may explain some of the samples with high $\delta^{51}\text{V}$.
691 However, the isotopically light (relative to BSE) MORB $\delta^{51}\text{V}$ values remain unexplained by our
692 calculations of peridotite melting. For Fe, even with a large Fe isotope fractionation correction
693 applied (up to 0.07 ‰ for some MORB samples; Sossi et al., 2016) some MORB samples with
694 extreme high $\delta^{57}\text{Fe}$ values still remain isotopically heavy relative to modelled KLB1 melts that
695 we consider most realistic of mantle melting (Soderman et al., 2021; see also Chen et al., 2019b).
696 We find that neither long-term analytical uncertainty nor equilibrium peridotite partial melting
697 from realistic regions of P-T space (see Fig. 6 in Soderman et al., 2021) can reproduce the
698 heaviest MORB Fe isotope ratios.

699 The Mg-Ca-Fe-V-Cr stable isotope compositions of mantle-derived melts are not expected to
700 show large heterogeneity as a consequence of the mantle potential temperature variability (with
701 associated changes in mineralogy and melt redox state; Fig. 4) of around 220 °C known from
702 major and trace element studies in MORB (Klein and Langmuir, 1987; Langmuir et al., 1992;
703 Gale et al., 2014). Therefore, we do not predict that the differences in either mean MORB
704 stable isotope composition or the variance in MORB isotope ratios from predicted peridotite
705 melts generated at $T_p = 1300^\circ\text{C}$ can be explained by potential temperature variability. Instead,
706 additional processes beyond single-stage equilibrium peridotite partial melting are required to
707 explain the full range of Mg-Ca-Fe-V stable isotope ratios of MORB. In some cases, the ad-
708 dition of single-stage pyroxenite melts (derived from a source with a MORB-like bulk isotope
709 composition) may explain some of the MORB range, and small amounts of pyroxenite in the
710 MORB source is consistent with conclusions using other tracers of lithological heterogeneity
711 (e.g., Zindler et al., 1984; Langmuir et al., 1986; Castillo et al., 2000; Waters et al., 2011). For

712 example, 10% pyroxenite in the MORB source (at $T_p = 1300^\circ\text{C}$) could produce melts with
713 $\delta^{26}\text{Mg} = 0.02\text{‰}$ and $\delta^{44}\text{Ca} = 0.04\text{‰}$ (using G2 pyroxenite) lower than, and $\delta^{57}\text{Fe} = 0.03\text{‰}$
714 higher than (using MIX1G pyroxenite), pure peridotite melts (Fig. 7), explaining some but not
715 all of the MORB data range for these three isotope systems. However, some of the MORB Fe
716 isotope data (including the seamounts and E-MORB data), would require the source of pyroxenite
717 melts to be isotopically fractionated relative to a protolith of pristine MORB ($\sim 0.05\text{--}0.15\text{‰}$
718 heavier than mean MORB to explain the full range of data, depending on proportion of pyroxenite
719 contributing to the final melt and P-T conditions of melting). A pyroxenite component in
720 the MORB source that is isotopically fractionated from pristine recycled MORB has previously
721 been proposed to explain MORB Fe and Mg isotope data (Sun et al., 2020; Zhong et al., 2021).
722 Multi-stage melting processes and reaction of pyroxenite-derived melts with ambient peridotite
723 have also been proposed to fractionate the bulk Fe isotope composition of mantle pyroxenite
724 away from that of the subducted material (e.g., Konter et al., 2016; Nebel et al., 2019; Sun et al.,
725 2020), although achieving the suggested large isotopic fractionations (e.g., $\Delta^{57}\text{Fe}_{\text{melt-source}} \geq$
726 0.1‰ ; Konter et al., 2016) during repeated partial melting events may be unrealistic (Soderman
727 et al., 2021).

728 We note that the large range in MORB V isotope ratios cannot be explained by pyroxenite in
729 the source, and that the lowest $\delta^{44}\text{Ca}$ pyroxenite melts (which appear able to explain the full
730 range of Ca isotope MORB data not matched by peridotite melting, Fig. 8) are only generated
731 at pressures > 35 kbar. For $T_p = 1300^\circ\text{C}$, pyroxenite fractions in the MORB source would
732 need to be $\geq 40\%$ to match all the Ca isotope MORB data (Fig. 7); alternatively, a pyroxenite
733 component with a bulk isotope composition $\sim 0.05\text{‰}$ lighter than average pristine MORB could
734 reduce the pyroxenite fraction required in the source to explain the Ca isotope data to lower
735 levels. While the $^{87}\text{Sr}/^{86}\text{Sr}$ filter we use where Sr isotope data is available attempts to remove
736 effects from recycled carbonates, we also note that since marine carbonates can have very light
737 $\delta^{44}\text{Ca}$ (up to 2‰ lighter than BSE; e.g., DePaolo, 2004; Fantle and DePaolo, 2005; Antonelli and
738 Simon, 2020), a contribution from recycled carbonates could also be involved in the generation
739 of light Ca isotopic compositions of basalts, as reported by e.g., Huang et al. (2011b); Liu et al.
740 (2017a).

741 The range of Mg-Ca-Fe-V isotopic compositions of MORB with respect to modelled equilibrium

742 melts leaves open the possibility of disequilibrium processes playing a role in the erupted stable
743 isotope composition of MORB. Disequilibrium associated with melt transport can arise as low
744 melt fractions are extracted from the mantle and rapidly migrate (Faul, 2001; Connolly et al.,
745 2009), as has been recorded using U-series disequilibria in MORB (e.g., McKenzie, 1985; Lund-
746 strom, 2003; Van Orman et al., 2006), and could be responsible for some major element trends
747 in the global MORB dataset (Oliveira et al., 2020). As disequilibrium during melt extraction
748 and transport is related to the timescales of chemical diffusion of major and trace elements in
749 the melt, disequilibrium could generate kinetic (rather than equilibrium) stable isotope frac-
750 tionation as seen in chemical diffusion studies (e.g., Richter et al., 2003, 2009; Watkins et al.,
751 2017). Disequilibrium processes could therefore influence both the isotopic composition of melt
752 source regions via metasomatism, and the isotopic composition of melts as they migrate and
753 erupt. In samples recording kinetic isotope fractionation signatures, the direction of isotopic
754 fractionation from the source isotope ratio will be the same for all isotope systems which behave
755 similarly during mantle melting, as the light isotopes of a given element will diffuse faster than
756 the heavier isotopes of the same element (e.g., Dauphas, 2007; Richter et al., 2008, 2009; Teng
757 et al., 2011), and the recorded isotopic fractionation should scale with the mass of the isotopes
758 considered. Kinetic processes in melts can produce much larger isotopic fractionations than
759 equilibrium processes (see Watkins et al., 2017) even for small concentration gradients in the
760 diffusing element (and therefore can be seen even in the diffusion of major elements in silicate
761 melts; Richter et al., 2003; Xiao et al., 2013; Watkins et al., 2014; Kang et al., 2020). We would
762 expect to see positive correlations between all isotope ratios measured for the same samples if
763 their isotope ratio was controlled by kinetic fractionation and chemical diffusion during disequi-
764 librium (for pairs of elements diffusing in the same direction, such as solid to melt), as reported
765 in e.g., Zhao et al. (2017a,b); Kang et al. (2020). Key pairs of isotopes to investigate to identify
766 a kinetic vs equilibrium process recorded in suites of natural samples would be those that be-
767 have similarly (from an elemental perspective) during partial melting but would show opposite
768 directions of isotopic fractionation from the source under equilibrium conditions (e.g., Fe and
769 Ca, with positive and negative melt-source isotopic fractionation respectively). Supporting this
770 prediction, kinetic effects during metasomatism have been proposed to generate positively cor-
771 related $\delta^{44}\text{Ca}$ - $\delta^{57}\text{Fe}$ values measured in ultramafic rocks (with a low $\delta^{44}\text{Ca}$, $\delta^{57}\text{Fe}$ metasomatic
772 endmember; Zhao et al., 2017a).

773 Low temperature alteration (relative to mantle temperatures) could also be responsible for caus-
774 ing significant stable isotope variability in measured MORB. Due to the $1/T^2$ dependence of
775 equilibrium stable isotope fractionations (Bigeleisen and Mayer, 1947), low temperature pro-
776 cesses generate larger magnitude isotope fractionation effects than magmatic processes. For
777 example, Fe leaching during hydrothermal fluid circulation is generally shown to preferentially
778 remove light Fe isotopes from the host rock, recorded in low $\delta^{57}\text{Fe}$ hydrothermal fluids (e.g.,
779 Rouxel et al., 2003; Severmann et al., 2004). Correspondingly high $\delta^{57}\text{Fe}$ is measured in heavily
780 leached basalts (up to 2.05 ‰ with 80 ‰ Fe loss; Rouxel et al., 2003), although Rouxel et al.
781 (2003) also show that on a scale larger than a few metres, bulk hydrothermally altered oceanic
782 crust is indistinguishable in $\delta^{57}\text{Fe}$ from MORB. Without significant Fe loss, altered crust may
783 not show such extreme Fe isotope variations (Williams et al., 2009), although alteration can
784 generate measurable Fe isotope fractionation in komatiites (Dauphas et al., 2010). By contrast,
785 Ca isotope data from mid-ocean ridge hydrothermal fluids suggest that the fluids and MORB
786 have similar $\delta^{44}\text{Ca}$ (Schmitt et al., 2003; Amini et al., 2008; Scheuermann et al., 2018). A Mg
787 isotope study of altered oceanic crust from the East Pacific Rise also shows that limited bulk
788 rock Mg isotope fractionation occurs during alteration, irrespective of alteration temperature
789 (Huang et al., 2015), but we note that the measured $\delta^{26}\text{Mg}$ range of altered crustal samples does
790 exceed the range of MORB data. Altered oceanic crust also records MORB-like $\delta^{51}\text{V}$ (Wu et al.,
791 2018), and Prytulak et al. (2013) show that $\delta^{51}\text{V}$ is generally insensitive to the range of common
792 alteration processes (serpentinisation, seafloor weathering – although an effect of increased $\delta^{51}\text{V}$
793 by 0.2–0.3 ‰ may be seen in extreme cases – and hydrothermal alteration) in the oceanic crust.
794 By contrast, alteration and serpentinisation in near-surface settings generate large variations
795 in $\delta^{53}\text{Cr}$ in ultramafic rocks (Farkaš et al., 2013; Wang et al., 2016), but altered oceanic crust
796 does not show the same isotopic variability (Wang et al., 2016). Given these studies, for the iso-
797 tope systems considered here mildly altered MORB is not expected to show significant isotopic
798 fractionation from erupted values, but given the potential for large low temperature isotopic
799 fractionations, alteration (both of the basaltic rocks analysed and oceanic crustal material that
800 serves as a protolith for mantle pyroxenite components) remains a consideration.

801 **5.2 Understanding the diversity of OIB Mg-Ca-Fe-V-Cr stable isotope com-** 802 **positions**

803 For Mg, Ca and Fe, measured OIB stable isotope ratios are more variable relative to those
804 measured in MORB (Fig. 8). This variability in OIB compared to MORB has been interpreted
805 as a potential signature of lithological heterogeneity in the OIB source (Mg: Zhong et al., 2017;
806 Stracke et al., 2018; Wang et al., 2018, Ca: Kang et al., 2019; Wang et al., 2019, Fe: Williams
807 and Bizimis, 2014; Konter et al., 2016; Nebel et al., 2019; Gleeson et al., 2020). Carbonates
808 and/or their derived melts have also been invoked for Mg and Ca isotopic heterogeneity in OIB
809 (e.g., Huang et al., 2011b; Liu et al., 2017a; Wang et al., 2018) but this topic is not discussed
810 here. An important caveat of interpreting data from Fig. 8 is that the ocean islands sampled
811 for each isotope system are different, and therefore the same localities are not present in each
812 OIB dataset – Mg, Ca and Fe have data for a wider range of OIBs than V and Cr, and so it
813 is possible that more variability would be seen in both V and Cr if more OIB are explored.
814 However, despite the wide variability, the mean OIB is analytically indistinguishable (using
815 long-term 2 S.D.) from mean MORB for Mg, Ca and V isotopes.

816 The mean measured OIB values can be explained both by modelled peridotite and pyroxenite
817 melts for the Mg, Ca, V and Cr stable isotope systems. For Fe, modelled pyroxenite melts
818 match mean OIB isotopic compositions better than peridotite melts, due to the use of mean
819 MORB as the bulk isotopic composition for the pyroxenites, offsetting the resulting pyroxenite
820 melts to heavier isotope ratios than peridotite melts. This potential Fe isotope signature of
821 lithological heterogeneity in some OIB data due to a bulk isotopic offset of the source relative
822 to ambient mantle has been suggested for several mantle plume systems (e.g., Hawaii, Williams
823 and Bizimis, 2014; Samoa, Konter et al., 2016; Pitcairn, Nebel et al., 2019; Galapagos, Gleeson
824 et al., 2020).

825 The greater diversity of OIB Mg-Ca-Fe-V-Cr stable isotope compositions displayed in OIB com-
826 pared to MORB and the modelled melts raises further questions, and we discuss only a few
827 here. Firstly, for Cr, the contribution of isotopically light pyroxenite melts to multi-lithologic
828 mantle melting means that a peridotite + pyroxenite mantle source is better able to explain the
829 measured OIB data range than pure peridotite melts, especially considering the good analytical

830 precision (2 S.D. $< 0.04\text{‰}$) compared to sample variability achieved in the Cr isotope studies
831 shown in Fig. 8. However, it is possible that the apparently better fit of a pyroxenite model seen
832 here is instead a result of fractionation from a peridotite-derived melt, as cotectic crystallisation
833 of olivine and spinel in basalts has been proposed to drive residual melts to lighter Cr isotope
834 ratios (Shen et al., 2020). Given the MgO filter used for OIB data in Fig. 8 however, spinel crys-
835 tallisation is unlikely to have affected all the OIB samples. The Cr stable isotope data therefore
836 provide tantalising evidence for the presence of lithological heterogeneity in OIB mantle source
837 regions, particularly in the case of the lightest measured Cr isotope compositions. Nonetheless,
838 without equivalent MORB Cr isotope data for comparison, and with a very limited published
839 sample set, it is difficult to confirm how well natural Cr isotope data matches equilibrium iso-
840 tope fractionation models. Fig. 7f suggests that even 20% SD pyroxenite in a mantle source
841 melting at $T_p = 1400\text{°C}$ will produce a bulk melt only 0.02‰ lighter than a pure peridotite
842 melt, which would not be able to explain the full range of Cr isotope OIB data. An expanded
843 study of MORB and OIB Cr isotope ratios would be an important step in both learning more
844 about Cr isotope behaviour and exploring stable isotopes as a new tool for studying mantle
845 melting conditions, given the potential for Cr isotope sensitivity to lithology discussed in this
846 study.

847 For Mg stable isotopes, the extra range in isotope composition displayed by OIB relative to
848 MORB is small ($< \pm 0.05\text{‰}$) relative to analytical precision, and the extreme MORB and OIB
849 isotope ratios are only resolvable from each other using a typical sample 2 S.E., rather than values
850 quoted for long-term reproducibility (see Table S2). The pyroxenite melts show a much wider
851 range of melt stable Mg isotope ratios than peridotite melts, and therefore could be contributing
852 to the small extra variability in OIB compared to MORB: for example, Fig. 7 shows that 10–
853 20% pyroxenite could generate melt isotopic variability of $\pm 0.04\text{–}0.05\text{‰}$ compared to a pure
854 peridotite source. Noticeably in the Mg isotope data, and particularly in comparison to Fe, the
855 natural data is less variable than the full range predicted for pyroxenite melts by our models (Fig.
856 8). This is despite the sampled OIB including St Helena, Society and the Cook-Austral Islands,
857 thought to sample various recycled crustal components (which will form pyroxenites in the
858 mantle) based on multiple studies including radiogenic isotopes, trace elements and other stable
859 isotopes (e.g., Chauvel et al., 1992; Stracke et al., 2005; Kawabata et al., 2011; Teng et al., 2013;
860 Hanyu et al., 2014). The limited Mg isotope variability in OIB compared to possible pyroxenite

861 melts could relate to the location of the melting region, and the dilution of pyroxenite melts with
862 peridotite, as highlighted in Fig. 7. For example, the most negative $\Delta^{26}\text{Mg}_{\text{melt-source}}$ for MIX1G
863 is generated at low melt fraction close to the solidus, at temperatures of around 1200–1250 °C
864 and pressures of 8–15 kbar, and the most positive $\Delta^{26}\text{Mg}_{\text{melt-source}}$ at depths of over 100 km
865 (Fig. S1). The isotopically lightest melts are not sampled by pyroxenite decompressing along
866 a peridotite isentrope even at $T_p = 1300$ °C (Fig. 5), and it is unlikely that deep, isotopically
867 heavy melts will be extracted without mixing with shallower, isotopically lighter melts. This
868 is in agreement with Fig. 7, which shows that typical aggregate melts from melting of multi-
869 lithologic mantle at 1300–1400 °C will have a much narrower range in isotopic composition than
870 the full range of pyroxenite melts shown in Fig. 8. Therefore, the process of melt extraction
871 and mixing from a multi-lithologic mantle will act to minimise the isotopic signature from any
872 deep pyroxenite melts, and result in relatively homogeneous OIB Mg isotope ratios seen in the
873 OIB erupted at the surface.

874 In contrast to the Mg isotope system, OIB isotope ratios for Fe, Ca and V extend outside the
875 range of any of the modelled peridotite or pyroxenite melts. For Ca and V isotopes, some
876 OIB data fall to the opposite side of BSE than expected for mantle melting. In both cases,
877 measurement uncertainties (particularly long-term 2 S.D.) are large and could account for all
878 the difference between modelled melts and natural data, and it will be important to continue
879 to improve analytical precision to better understand the natural data and be able to make
880 comparisons between OIB sample sets, especially given the minimal sensitivity to lithology
881 predicted for these isotope systems (Section 4.2.1). In the case of V, fractional crystallisation
882 of Fe-Ti oxides has been shown to drive melts up to 0.3 ‰ higher than their initial isotope
883 composition, therefore able to match the spread of V isotope data to values higher than BSE,
884 although this effect has only been recorded in basalts with < 5 wt % MgO (Wu et al., 2018; Ding
885 et al., 2020). Therefore it is unlikely that the basalts with V isotope ratios heavier than BSE
886 shown in Fig. 8 reflect Fe-Ti oxide crystallisation. For Ca, non-pyroxenitic origins of isotopic
887 variability have been proposed for komatiite data, such as the generation of $\delta^{44}\text{Ca}$ as low as
888 0.64 ‰ from serpentinised oceanic crust in the mantle source (John et al., 2012; Amsellem et al.,
889 2019). If improved analytical precision confirms the Ca isotope data variability shown in Fig. 8
890 is real, similar non-pyroxenitic origins may be able to explain some or all of the spread of OIB
891 data.

892 For Fe, measurement uncertainties are small compared to the variability in the OIB Fe isotope
893 data, which extends to much higher $\delta^{57}\text{Fe}$ values than matched by the melts presented here
894 (Fig. 7; see also Soderman et al., 2021). Heavier Fe stable isotope compositions than can be
895 explained by single-stage melting of garnet-bearing lithologies are also seen in seamounts from
896 close to the EPR (Sun et al., 2020) – the heaviest isotope ratios are found in samples with <
897 7 wt % MgO, and therefore not shown in Fig. 8, but Sun et al. (2020) find that these are still
898 isotopically heavier than can be explained by single-stage melting once they correct their samples
899 for differentiation. One solution to explaining the heavy Fe isotope data could be to invoke a
900 mantle source with unusually high $\delta^{57}\text{Fe}$ (i.e., Fe-isotopically heavier than average MORB as
901 assumed in Figs. 7 and 8), generated for example as a frozen melt that has been through multiple
902 melting and remobilisation stages, and now stalled in the lithosphere. While the limitations
903 of multiple melting events with regards to mass balance and plausible melting degrees are
904 discussed in Soderman et al. (2021), multiple melting stages of garnet-bearing lithologies and
905 reaction with ambient mantle have been proposed to generate the required isotopically heavy
906 mantle sources ($\delta^{57}\text{Fe} > 0.3\text{‰}$) to explain some of the highest $\delta^{57}\text{Fe}$ in the global dataset
907 (Nebel et al., 2019; Sun et al., 2020), although based on our modelled small melt-source isotope
908 fractionations the isotopically heaviest Fe isotope OIB data would require a mantle source with
909 an even heavier overall bulk composition of $\delta^{57}\text{Fe} > 0.4\text{‰}$. Alternatively, large stable isotope
910 effects may be attributed to disequilibrium and/or low temperature alteration processes, as
911 discussed for MORB above. A further possibility to explain heavy Fe isotope ratios in OIB
912 could be the mixing of isotopically heavy components derived from the lower mantle (such as
913 bridgmanite cumulates with $\delta^{57}\text{Fe} \sim 0.3\text{‰}$, Williams et al., 2021) into upwelling plumes. High
914 $\delta^{44}\text{Ca}$ lower mantle components from Ca-perovskite crystallisation related to early Earth magma
915 ocean processes have also been proposed as the origin of isotopically heavy Ca isotope komatiite
916 data (up to 1.54 ‰; Amsellem et al., 2019).

917 **6 Summary**

918 We have combined a thermodynamically self-consistent model for mantle melting of peridotite
919 and pyroxenite with a model for equilibrium Mg-Ca-Fe-V-Cr stable isotope fractionation to
920 quantitatively predict melt-source stable isotope fractionation over P-T conditions relevant to

921 mantle melting at mid-ocean ridges and within mantle plumes. The results allow us to assess
922 the potential of these stable isotope systems for investigating mantle temperature variations,
923 source heterogeneity, and their record in the erupted basalts (MORB, OIB).

924 We used our results for the melting of KLB1 peridotite to explore the sensitivity of melt sta-
925 ble Mg-Ca-Fe-V-Cr isotope compositions in each isotope system to variations in mantle po-
926 tential temperature. Redox-sensitive isotopes (Fe, V) show most variability in accumulated
927 melt isotopic composition with potential temperature, due to the variation in $\text{Fe}^{3+}/\text{Fe}_T$ (and
928 $\text{V}^{4+,5+}/\text{V}_T$) as the average depth of melting increases with T_p . The predicted variability in the
929 V isotopic composition of melts is qualitatively consistent with $\delta^{51}\text{V}-\text{Na}_{8.0}$ trends in natural
930 data. However, analytical precision improvements upon current levels would be necessary for
931 variability to be clearly resolved in the compositions of erupted basalts. At present, the equi-
932 librium fractionation of the stable isotopes considered here during peridotitic mantle melting is
933 not predicted to be a useful tool in identifying temperature variations in MORB or OIB melting
934 regimes.

935 We also obtained model results for Mg-Ca-Fe-V-Cr fractionation during melting of MIX1G
936 silica-deficient pyroxenite and G2 silica-excess pyroxenite. Our results show that Mg and Ca
937 isotopes have the most sensitivity to the presence of garnet in a melt source, with melts pre-
938 dicted to record a negative Mg-Ca isotope correlation if garnet is present in the mantle source,
939 although improvements in analytical precision would be needed to clearly resolve this trend.
940 Our results also demonstrate that Fe, Mg and Cr stable isotopes are sensitive to the proportion
941 of pyroxenite (with a MORB-like bulk isotope composition) in a mantle source, although with
942 current analytical precision only pyroxenite fractions greater than around 10% are resolvable.
943 Moderate improvements in analytical precision could result in Ca stable isotopes being added to
944 this list, and potentially allow pyroxenite fractions $< 10\%$ to be identified in these isotope sys-
945 tems. However, Cr stable isotope partitioning in the mantle is generally understudied (both in
946 measurements of natural samples, and studies of its bonding and fractionation in relevant man-
947 tle phases): while this system has the potential to be sensitive to lithological heterogeneity, we
948 require improvements in the certainty of model inputs and/or more natural data to confidently
949 use this isotopic tool in the mantle.

950 Comparison of the global datasets of stable isotope data for MORB and OIB with modelled

951 melt isotope ratios generated by equilibrium fractionation highlight both agreements and dis-
952 agreements between equilibrium melting models and measured data. Generally, mean MORB
953 Mg-Ca-Fe-V-Cr stable isotope compositions, within typical long-term analytical precision, can
954 be matched to the mean peridotite melt isotope ratio produced for a mantle potential temper-
955 ature of 1300 °C. However, the range of MORB data is wide compared to modelled peridotite
956 melts – even after analytical uncertainty on data and fractional crystallisation are considered,
957 some MORB record stable isotope ratios that appear unable to be matched by equilibrium peri-
958 dotite melting (notably in Mg, Fe and V, and Ca may show a similar result as measurement
959 precision improves). Lithological heterogeneity (pyroxenite) in the MORB source can contribute
960 to some of this variability but not all, and disequilibrium and/or low temperature processes may
961 also be required to match the observed data. For OIB, the variability in natural data, even for
962 isotope systems such as Fe and Cr where analytical precision is good compared to the data
963 variability, could suggest that the involvement of pyroxenite may be recorded in the existing
964 Mg, Fe and Cr stable isotope datasets. However, differences between the range of natural data
965 and modelled pyroxenite melts suggest that additional processes such as the offset of the isotope
966 ratio of recycled components from pristine MORB values and kinetic fractionations may also
967 be important in understanding the current Mg-Ca-Fe-V-Cr stable isotope datasets available for
968 primitive global oceanic basalts.

969 We conclude that stable isotopes in basalts have variable potential as a tracer of heterogeneity
970 in the mantle, complementing existing tools well – however there are knowledge and data gaps
971 that need to be filled before we can use these systems to their full potential.

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APPENDIX

Supplementary Material to ‘Global trends in novel stable isotopes in basalts: theory and observations’ by Soderman et al.

1 Data sources

The data sources used in the natural data compilation in Figs. 1 and 8 in the main text are given in Table S1.

Table S1: Literature data sources for stable isotope data compilation.

Isotope	Type	Literature sources
$\delta^{26}\text{Mg}$	MORB	Wiechert and Halliday (2007); Bourdon et al. (2010); Teng et al. (2010); Zhong et al. (2021b)
	OIB	Bourdon et al. (2010); Teng et al. (2010); Zhong et al. (2017); Wang et al. (2021)
$\delta^{44}\text{Ca}$	MORB	Zhu et al. (2018); Chen et al. (2020b); Zhu et al. (2020a)
	OIB	Huang et al. (2011b); Valdes et al. (2014, 2019); Feng et al. (2017); Chen et al. (2020b)
$\delta^{57}\text{Fe}$	MORB	Teng et al. (2013); Nebel et al. (2013); Zhong et al. (2021) (seamounts: Sun et al., 2020)
	OIB	Teng et al. (2008); Schuessler et al. (2009); Konter et al. (2016); Nebel et al. (2019) Peters et al. (2019); Gleeson et al. (2020); Soderman et al. (2021)
$\delta^{51}\text{V}$	MORB	Prytulak et al. (2013); Wu et al. (2018); Novella et al. (2020)
	OIB	Prytulak et al. (2013); Ding et al. (2020)
$\delta^{53}\text{Cr}$	OIB	Bonnand et al. (2020); Shen et al. (2020)

Table S2 gives the references for the typical sample 2 S.E. and long-term 2 S.D. used throughout the manuscript.

Table S2: Literature data sources for stable isotope data errors.

Isotope	Type	Value (‰)	Literature sources
$\delta^{26}\text{Mg}$	sample 2 S.E.	0.02	Wang et al. (2021)
	long-term 2 S.D.	0.06	Wang et al. (2021)
$\delta^{44}\text{Ca}$	sample 2 S.E.	0.05	Zhu et al. (2018, 2020a)
	long-term 2 S.D.	0.14	Chen et al. (2019)
$\delta^{57}\text{Fe}$	sample 2 S.E.	0.02	Soderman et al. (2021)
	long-term 2 S.D.	0.05	Soderman et al. (2021)
$\delta^{51}\text{V}$	sample 2 S.E.	0.05	Wu et al. (2018); Novella et al. (2020)
	long-term 2 S.D.	0.08	Qi et al. (2019)
$\delta^{53}\text{Cr}$	sample 2 S.E.	0.01	Sossi et al. (2018)
	long-term 2 S.D.	0.04	Xia et al. (2017); Sossi et al. (2018)

2 Phase equilibria calculations

In order to explore stable isotope fractionation taking place during the melting of pyroxenite and peridotite lithologies within the mantle, we first constructed pseudosections and calculated mineral chemistries for MIX1G (KLB1 and G2 calculations are presented in Soderman et al., 2021). The calculations were performed in the KNCFMASSTOcr system using THERMOCALC v3.47 (Powell et al., 1998), and the latest version of the thermodynamic dataset of Holland and Powell (2011) with the a-X models of Holland et al. (2018).

The composition of the MIX1G lithology (Hirschmann et al., 2003; Lambart et al., 2016) used in the phase-equilibria calculations is given in Table S3. The ferric iron content is taken as between that of KLB1 peridotite and MORB, following the approach used for KG1 by Jennings et al. (2016), where $\text{Fe}^{3+}/\text{Fe}_T = 0.1$.

Table S3: MIX1G composition, in KNCFMASSTOcr system (mol % oxides). Composition from Lambart et al. (2016), renormalised without MnO.

SiO ₂	Al ₂ O ₃	CaO	MgO	mole % oxide					
				FeO _t	K ₂ O	Na ₂ O	TiO ₂	Cr ₂ O ₃	O
38.5	1.78	2.82	50.58	5.52	0.01	0.25	0.07	0.11	0.09

After calculation of the phase boundaries, THERMOCALC was run over a grid of P-T space, from 5–40 kbar and 1100–1730 °C.

3 Details and results of the main stable isotope fractionation model

The equilibrium isotopic composition of the phases present at any point in P-T space (including melts) was calculated for each isotope system, based on calculated equilibrium isotope fractionation factors, α^{A-B} , between the phases that are present. α^{A-B} was calculated either from reduced partition function ratios (β factors), or from calculated cation-oxygen force constants, as outlined for Mg, Ca, V and Cr isotopes below. The Fe isotope fractionation input parameters are presented in the supplementary information of Soderman et al. (2021). The mineral abbreviations used in the figures in the following subsections are: olivine (ol), clinopyroxene (cpx), orthopyroxene (opx), spinel (spn), garnet (g), plagioclase (pl).

3.1 Mg

The β^{Mg} factors used for Mg-bearing mantle minerals (spinel, clinopyroxene, orthopyroxene, olivine, garnet) are given in Table S4. These β^{Mg} factors are taken from Huang et al. (2013), with the exception of spinel, which is taken from Schauble (2011) and scaled relative to a reference β^{Mg} of olivine which is calculated in both studies. Pigeonite (a stable phase in MIX1G) is modelled with the same β^{Mg} as clinopyroxene; liquid (melt) is taken as having no fractionation from olivine in the main model (Stracke et al., 2018, see Appendix Section 4 for discussion), so $\beta_{\text{melt}}^{\text{Mg}} = \beta_{\text{olivine}}^{\text{Mg}}$.

The full results of the Mg model, showing the isotopic composition of each phase calculated in

Table S4: Parameters for β^{Mg} factors from Huang et al. (2013). $1000\ln\beta = a/T^6 + b/T^4 + c/T^2$, where T = temperature (Kelvin).

Mineral	a	b	c
spinel	-2.74×10^{16}	2.08×10^{11}	2.93×10^6
orthopyroxene	-8.10×10^{15}	6.65×10^{10}	2.25×10^6
clinopyroxene	-1.17×10^{16}	1.03×10^{11}	2.24×10^6
olivine	-1.19×10^{16}	1.05×10^{11}	2.07×10^6
garnet	-5.50×10^{15}	4.92×10^{10}	1.30×10^6
melt	-1.19×10^{16}	1.05×10^{11}	2.07×10^6

40 KLB1, MIX1G and G2, are in Fig. S1.

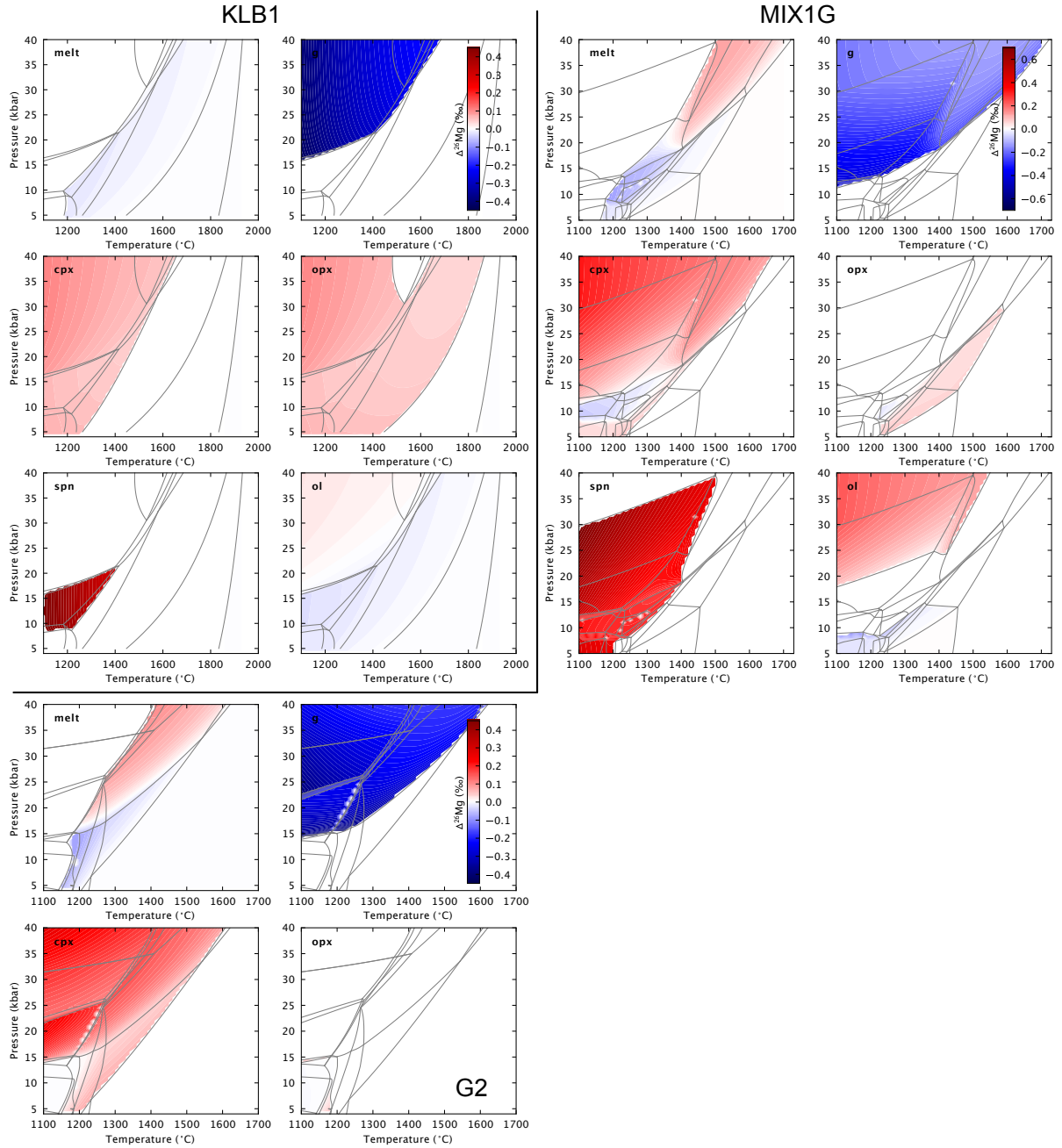


Figure S1: Full results of the Mg isotopic fractionation model, $\Delta^{26}\text{Mg}_{\text{phase-source}}$, for the three lithologies.

Table S5: β^{Ca} factors (given as $1000\ln\beta$, at a temperature of 1000K) from Antonelli et al. (2019).

Mineral	Structure used	$1000\ln\beta$
plagioclase	anorthite	1.06
orthopyroxene	Ca/Mg = 1/32	1.84
clinopyroxene	diopside	1.32
olivine	forsterite Ca/Mg = 1/64	2.11
garnet	pyrope Ca/Mg = 1/24	2.00
liquid		1.19

41 3.2 Ca

42 The β^{Ca} factors used for Ca-bearing mantle minerals (plagioclase, clinopyroxene, orthopyroxene,
 43 olivine, garnet) are given in Table S5, and are taken from Antonelli et al. (2019). Where options
 44 exist for minerals with different structures and/or compositions, we use those with most similar
 45 Ca/Mg to the typical phase compositions calculated by THERMOCALC for KLB1. Pigeonite (a
 46 stable phase in MIX1G) is modelled with the same β^{Ca} as clinopyroxene; liquid is taken as having
 47 β^{Ca} intermediate between clinopyroxene and plagioclase (anorthite), following observations from
 48 natural samples (Zhang et al., 2018; Antonelli et al., 2019b) and consistent with the approach
 49 used by Antonelli et al. (2021). β^{Ca} factors are made temperature-dependent by

$$\ln\beta_{(T)} = \ln\beta_{(1000\text{K})} \times \frac{10^6}{T^2} \quad (1)$$

50 following Antonelli et al. (2021).

51 The full results of the Ca model, showing the isotopic composition of each phase calculated in
 52 KLB1, MIX1G and G2, are in Fig. S2.

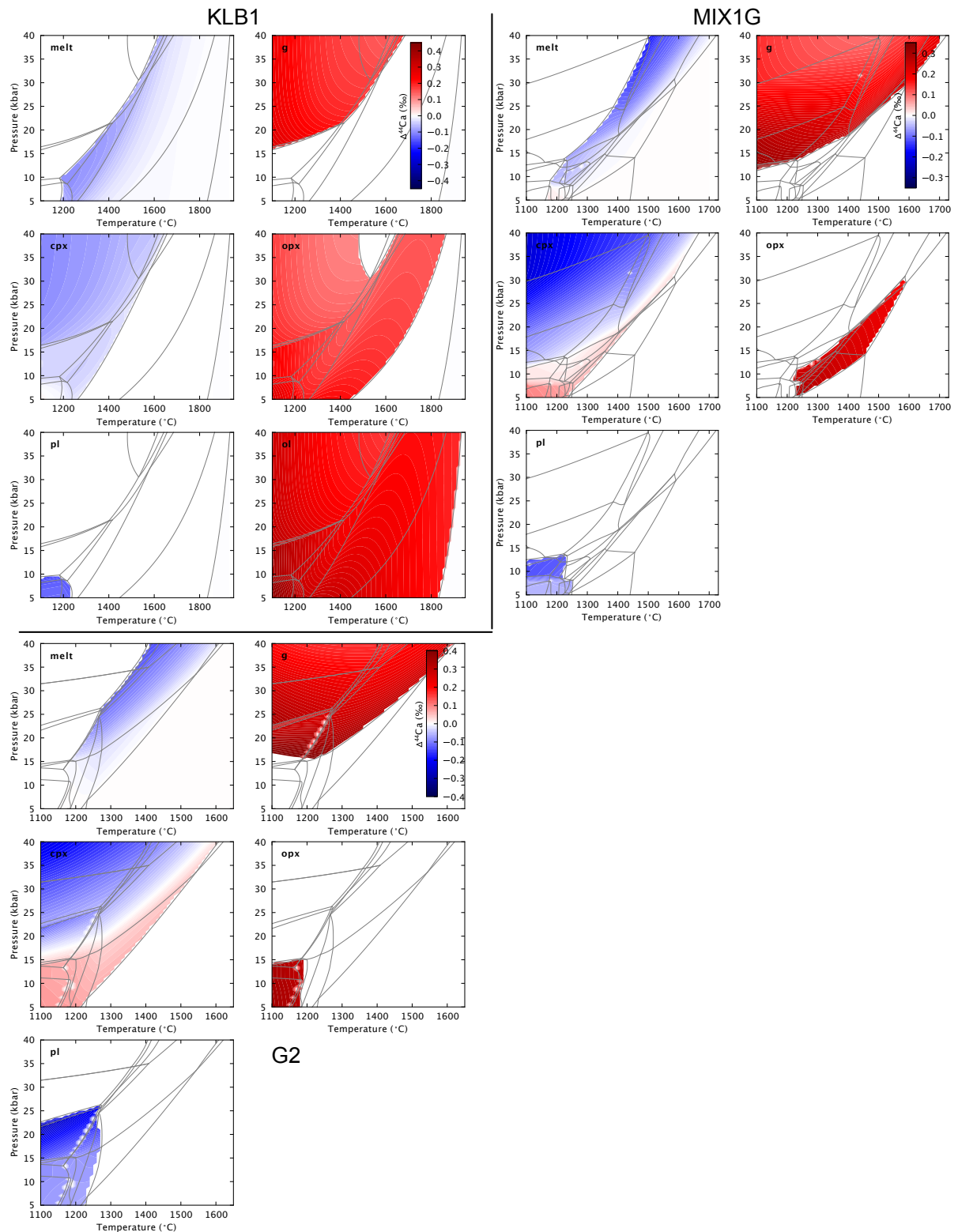


Figure S2: Full results of the Ca isotopic fractionation model, $\Delta^{44}\text{Ca}_{\text{phase-source}}$, for the three lithologies.

53 **3.3 V**

54 The proportion of $V^{3+/4+/5+}$ was calculated following Toplis and Corgne (2002), using the
55 FeO/Fe_2O_3 ratio of the melt (calculated by THERMOCALC) at each point in P-T space with

$$\log(XV^{3+}/XV^{4+}) = \log(XFe^{2+}/XFe^{3+}) - 2.4 \quad (2)$$

56 and

$$\log(XV^{4+}/XV^{5+}) = \log(XFe^{2+}/XFe^{3+}) - 0.3 \quad (3)$$

57 where X is the mole fraction of cation.

58 The partition coefficients and β^V factors of each valence state of V in the V-bearing mantle
59 minerals (olivine, clinopyroxene, orthopyroxene, garnet, spinel) are given in Table S6. Partition
60 coefficients are from Mallmann and O'Neill (2009), with garnet assumed to have similar V
61 partitioning to clinopyroxene (Mallmann and O'Neill, 2007, 2009). The β^V factors are taken from
62 Wu et al. (2015) for V in different coordination environments in solution systems; alternative
63 choices of β^V due to uncertainties in V coordination number are discussed in Section 4 below.
64 Pigeonite (a stable phase in MIX1G) is modelled as clinopyroxene.

Table S6: Parameters for V partitioning and isotopic fractionation. Partition coefficients (D) from Mallmann and O'Neill (2009). β^V factors are from Wu et al. (2015). $1000\ln\beta = a/T^6 + b/T^4 + c/T^2$, where T = temperature (Kelvin).

Mineral	Valence	$D_{\min-liq}$	a	b	c
olivine	V^{3+}	0.18	1.0×10^{14}	-7.70×10^9	6.28×10^5
	V^{4+}	0.0227	4.0×10^{14}	-1.80×10^{10}	9.43×10^5
	V^{5+}	0.00255	7.0×10^{14}	-3.06×10^{10}	1.26×10^6
orthopyroxene	V^{3+}	1	1.0×10^{14}	-7.70×10^9	6.28×10^5
	V^{4+}	0.674	4.0×10^{14}	-1.80×10^{10}	9.43×10^5
	V^{5+}	0.00338	7.0×10^{14}	-3.06×10^{10}	1.26×10^6
clinopyroxene	V^{3+}	3.88	1.0×10^{14}	-7.70×10^9	6.28×10^5
	V^{4+}	1.13	4.0×10^{14}	-1.80×10^{10}	9.43×10^5
	V^{5+}	0.0142	7.0×10^{14}	-3.06×10^{10}	1.26×10^6
spinel	V^{3+}	15.3	1.0×10^{14}	-7.70×10^9	6.28×10^5
	V^{4+}	1.44	4.0×10^{14}	-1.80×10^{10}	9.43×10^5
	V^{5+}	0.00303	7.0×10^{14}	-3.06×10^{10}	1.26×10^6
garnet	V^{3+}	3.88	1.0×10^{14}	-7.70×10^9	6.28×10^5
	V^{4+}	1.13	4.0×10^{14}	-1.80×10^{10}	9.43×10^5
	V^{5+}	0.0142	7.0×10^{14}	-3.06×10^{10}	1.26×10^6
melt	V^{3+}	1	1.0×10^{14}	-7.70×10^9	6.28×10^5
	V^{4+}	1	4.0×10^{14}	-1.80×10^{10}	9.43×10^5
	V^{5+}	1	7.0×10^{14}	-3.06×10^{10}	1.26×10^6

65 The full results of the V isotope model, showing the isotopic composition of each phase calculated
66 in KLB1, MIX1G and G2, are in Fig. S3. Note that stable isotope fractionation can only be
67 calculated where there is liquid present, i.e., where the proportion of each V valence state can
68 be calculated.

69 **3.4 Cr**

70 The proportion of $Cr^{2+/3+}$ was calculated using the parameterisation given in Berry et al. (2021)

$$\frac{Cr^{2+}}{Cr_T} = \frac{1}{1 + 10^{(0.25 \log fO_2 + \log K')}} \quad (4)$$

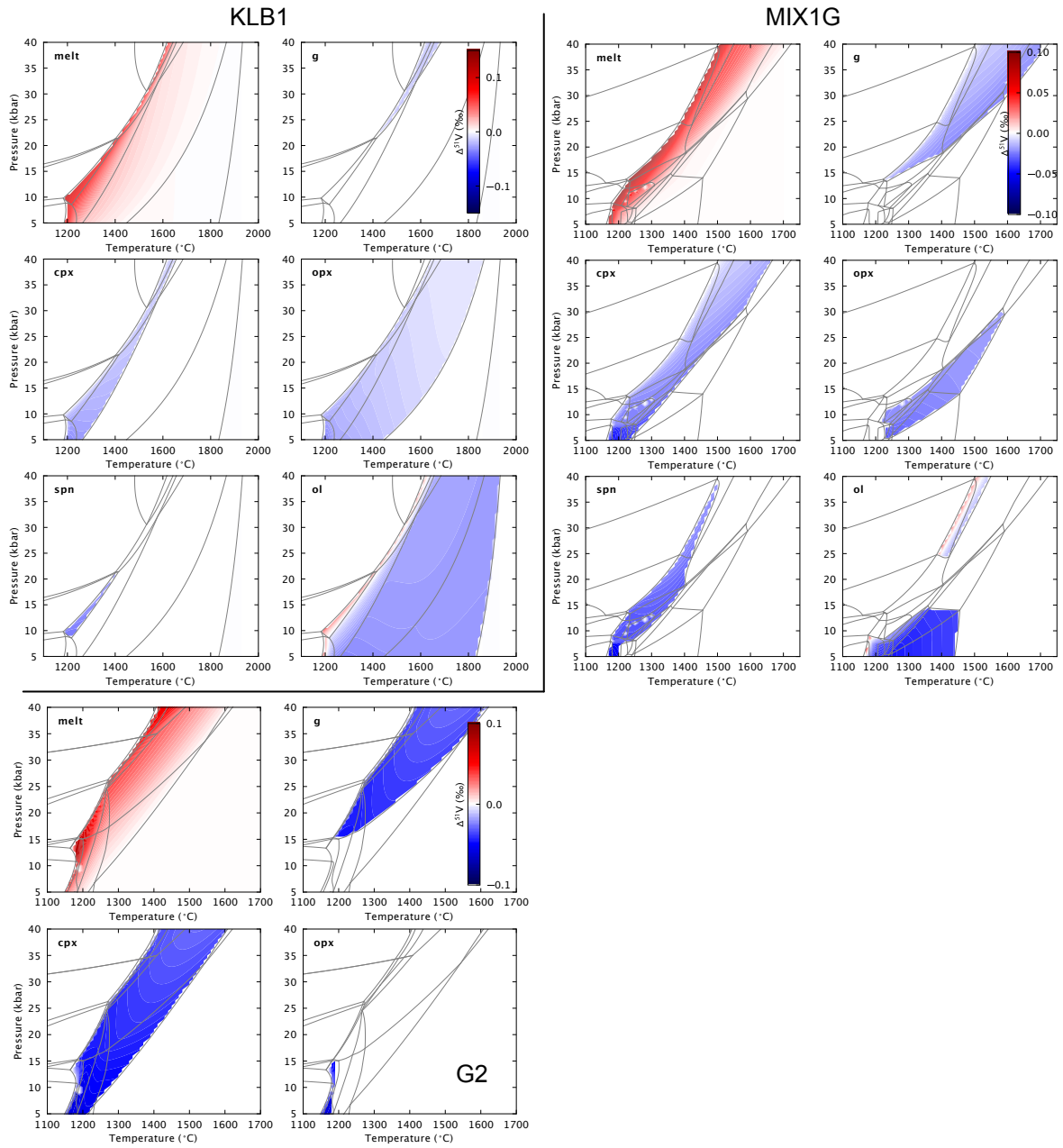


Figure S3: Full results of the V isotopic fractionation model, $\Delta^{51}\text{V}_{\text{phase-source}}$, for the three lithologies.

Table S7: Optical basicity values for each oxide, from Mills (1993).

optical basicity, Λ									
SiO ₂	TiO ₂	Al ₂ O ₃	Fe ₂ O ₃	Cr ₂ O ₃	FeO	CaO	MgO	Na ₂ O	K ₂ O
0.48	0.61	0.60	0.75	0.58	1	1	0.78	1.15	1.4

71 with fO_2 calculated at each P-T point using the fO_2 melt software (Holland et al., 2018), and
 72 the modified equilibrium constant calculated by

$$\log K' = \frac{9770}{T} - 7.69 + 6.22\Lambda + \frac{(900P - 172P^2)}{T} \quad (5)$$

73 where T is in Kelvin, P is in GPa. Λ is the optical basicity, a composition-dependent term (Duffy,
 74 1993) which reflects equilibrium between bridging oxygens, non-bridging oxygens and free oxide
 75 anions in a melt (Humphreys et al., 2015). Since optical basicity values can be assigned to
 76 individual oxides (Duffy, 1993), an ideal optical basicity for a melt can be calculated, following
 77 Mills (1993), as

$$\Lambda = \frac{\sum X_i n_i \Lambda_i}{\sum X_i n_i} \quad (6)$$

78 where X_i is the mole fraction of oxide i, n is the number of O associated with oxide i, and Λ is
 79 the theoretical Λ of oxide i (see Table S7). At each P-T point, the calculated Λ was used with
 80 equations 3 and 4 to calculate Cr redox state.

81 The partition coefficients of each valence state of Cr in the Cr-bearing mantle minerals (olivine,
 82 clinopyroxene, orthopyroxene, garnet, spinel) are given in Table S8. Partition coefficients for
 83 olivine, clinopyroxene, orthopyroxene are from (Mallmann and O'Neill, 2009); for spinel from
 84 the supplement of Shen et al. (2018), for garnet from Sutton et al. (2008). The ionic model
 85 inputs (Table S8) are based on the supplement of Shen et al. (2018); for garnet, where no ionic
 86 inputs are presented in Shen et al. (2018), we take coordination information from Deer et al.
 87 (2013), and use the relevant ionic radii as used by Shen et al. (2018). The force constants for
 88 Cr-O bonds were then calculated based on the equations presented in Sossi and O'Neill (2017)
 89 for Fe isotopes.

Table S8: Parameters for Cr partitioning and isotopic fractionation model. Partition coefficients (D) from Mallmann and O'Neill (2009); Shen et al. (2018); Sutton et al. (2008). Mineral site and bonding information is based on Shen et al. (2018). cpx = clinopyroxene, opx = orthopyroxene.

Mineral	Cr species	$D_{\text{min-melt}}$	Site	Cr coordination	Bond length (\AA)	O coordination	No. sites
spinel	2+	0	n/a	n/a	n/a	n/a	n/a
spinel	3+	220	M	6	1.995	4	2
garnet	2+	0	n/a	n/a	n/a	n/a	n/a
garnet	3+	12	M2	6	1.995	4	2
olivine	2+	0.85	M2	6	2.18	4	1
olivine	3+	0.85	M1	6	1.995	4	1
opx	2+	0.843	M2	7	2.305	3.57	1
opx	3+	3.52	M1	6	1.988	3.67	1
cpx	2+	0.587	M2	7	2.305	3.57	1
cpx	3+	12.6	M1	6	1.988	3.67	1
liquid	2+	1		4	1.913	4	
liquid	3+	1		6	1.995	4	

90 The full results of the Cr model, showing the isotopic composition of each phase calculated in
 91 KLB1, MIX1G and G2, are in Fig. S4. Note that stable isotope fractionation can only be
 92 calculated where there is liquid present, i.e., where the proportion of each Cr valence state can
 93 be calculated.

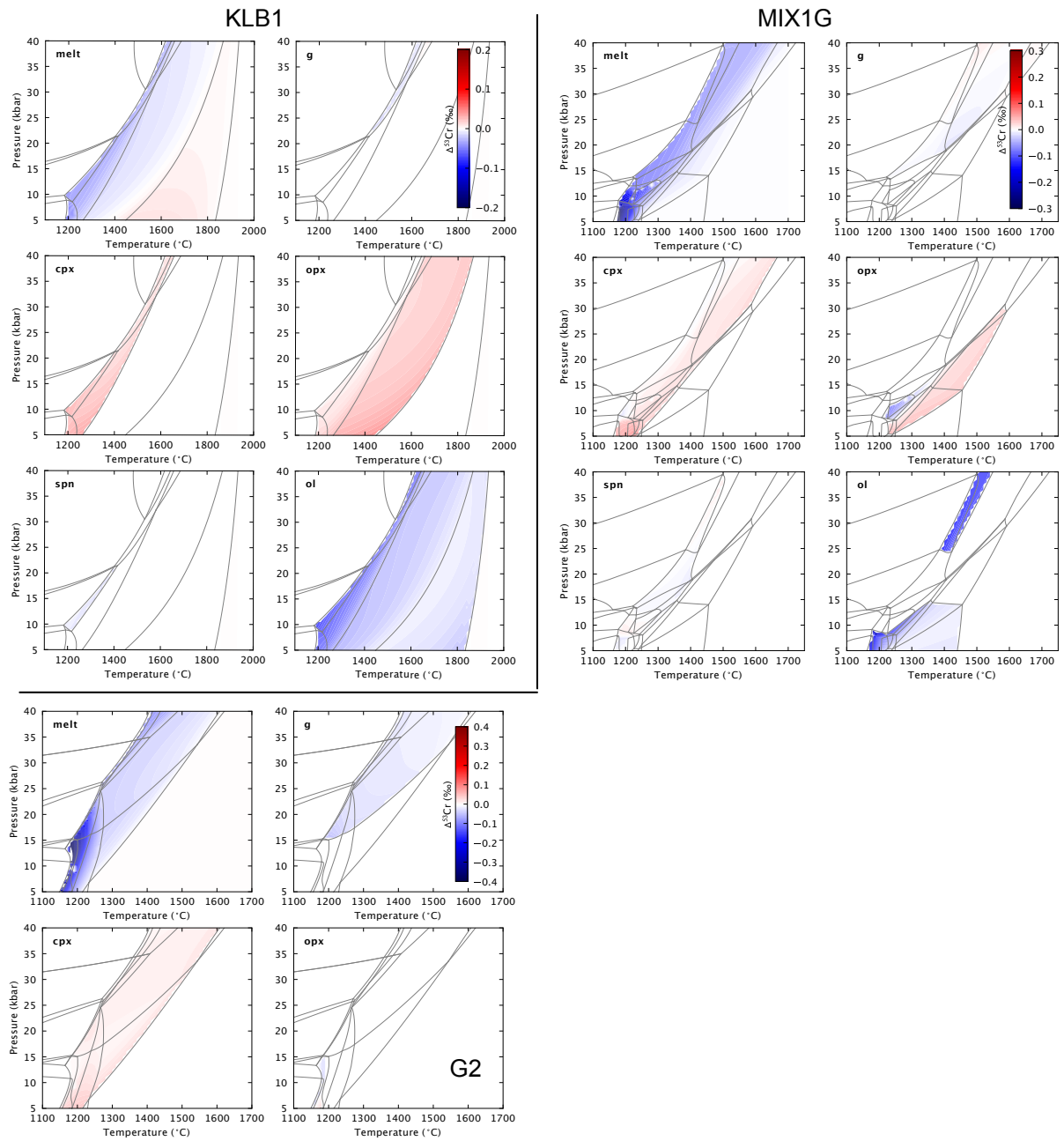


Figure S4: Full results of the Cr isotopic fractionation model, $\Delta^{53}\text{Cr}_{\text{phase-source}}$, for the three lithologies.

94 3.5 Consistency of valence state calculations

95 As discussed in the main text, the calculated $\text{Fe}^{3+}/\text{Fe}_T$, $\text{Cr}^{3+}/\text{Cr}_T$ and $\text{V}^{4+,5+}/\text{V}_T$ are not fully
 96 self-consistent, as the speciation of both Cr and V are calculated based on the $\text{Fe}^{3+}/\text{Fe}_T$ at
 97 any point. This means that O (oxygen) is being partitioned into the oxidised Cr and V oxides
 98 without affecting Fe redox state. We consider tying the Cr and V valence states to the Fe redox
 99 equilibria to be reasonable, as Fe is the dominant redox-sensitive element in the mantle with an
 100 abundance considerably greater than that of Cr or V. To quantify this, we have calculated the
 101 fraction of oxygen associated with oxidised valence states of the each redox-sensitive element
 102 out of the total oxygen associated with the oxidised valence states, e.g.,

$$\frac{\text{O}_{\text{Fe}^{3+}}}{\text{O}_{\text{Fe}^{3+}} + \text{O}_{\text{Cr}^{3+}} + \text{O}_{\text{V}^{4+}} + \text{O}_{\text{V}^{5+}}} \quad (7)$$

103 where O_X represents the number of moles of oxygen associated with cation X at any given P-T
 104 point. At $T = 1300^\circ\text{C}$ and $P = 12\text{ kbar}$, we calculate that 73% of the oxygen associated with
 105 oxidised cations is associated with Fe^{3+} , 27% with Cr^{3+} , and $< 0.01\%$ with $\text{V}^{4+,5+}$, for total
 106 Fe, Cr and V contents in peridotite (Davis et al., 2009; Salters and Stracke, 2004).

107 However, although our use of the Fe redox equilibria to calculate Cr and V speciation is a
 108 reasonable approximation, it is important to note that considering the combined Fe-Cr-V redox
 109 equilibria as in real systems would reduce the variability in the redox state of each element across
 110 P-T space compared to our model. This is because the combined effect of considering three redox
 111 equilibria able to accommodate oxygen via their oxidised cations together is that each system
 112 will have to shift its redox equilibria less compared to a scenario where only one redox-sensitive
 113 element is considered. To approximately quantify the effect of this redox feedback, we took an
 114 extreme case of reducing all the Cr_2O_3 in the melt at a given P-T point, accommodating the
 115 released oxygen by increasing the amount of Fe_2O_3 and reducing the amount of FeO accordingly.
 116 We then calculated the $f\text{O}_2$ of this new composition through the $f\text{O}_2\text{melt}$ software (Holland
 117 et al., 2018), and recalculated the expected $\text{Cr}^{3+}/\text{Cr}_T$ given the new $f\text{O}_2$. We find that, for this
 118 calculation performed using at $T = 1350^\circ\text{C}$ and $P = 12\text{ kbar}$, the $f\text{O}_2$ change produces a 1%
 119 change in $\text{Cr}^{3+}/\text{Cr}_T$ and negligible impact on the calculated $\delta^{53}\text{Cr}$.

120 4 Alternative input parameters and results

121 The inputs into the isotope fractionation models are better constrained for some of the stable
 122 isotope systems considered relative to others. These uncertainties limit our ability to predict the
 123 behaviour of some isotope systems in high temperature settings, and therefore limit the utility
 124 of stable isotopes as a tool for exploring global basalt systematics. Here we discuss the results
 125 of using alternative inputs.

126 4.1 Mg: $\alpha_{\text{orthopyroxene-melt}} = 1$

127 While the bonding of Mg in crystal lattices can be used to calculate a self-consistent set of β^{Mg}
 128 factors (Schauble, 2011; Huang et al., 2013) to input into the melting model, β^{Mg} factors do not
 129 exist for silicate liquid. However, Stracke et al. (2018) show that $\alpha_{\text{olivine, orthopyroxene-melt}} \sim 1$
 130 based on equilibrated peridotite xenoliths. Mg is also similarly coordinated ([5–6] fold) in melt
 131 and olivine/orthopyroxene (George and Stebbins, 1998; Shimoda et al., 2007; Stracke et al.,

132 2018), providing further evidence for minimal olivine/orthopyroxene – melt Mg stable isotope
133 fractionation at magmatic temperatures.

134 We have chosen to use $\alpha_{\text{olivine-melt}} = 1$; i.e., using the same β^{Mg} values for liquid as for olivine.
135 The alternative is to use $\alpha_{\text{orthopyroxene-melt}} = 1$. Fig. S5 compares the $\Delta^{26}\text{Mg}_{\text{melt-source}}$ for an
136 input of $\alpha_{\text{olivine-melt}} = 1$ (left hand panels; these are the results used in the main text) with an
137 input of $\alpha_{\text{orthopyroxene-melt}} = 1$ (right hand panels). The new input parameters replace the melt
138 in Table S4 with Table S9.

Table S9: Alternate parameters for Mg isotope β^{Mg} factors from Huang et al. (2013). $1000\ln\beta = a/T^6 + b/T^4 + c/T^2$, where T = temperature (Kelvin).

Mineral	a	b	c
melt	-8.10×10^{15}	6.65×10^{10}	2.25×10^6

139 For KLB1 peridotite, the magnitude of the difference in melt isotope composition between the
140 two models is $\approx 0.05\text{‰}$ (small compared to the average analytical uncertainty of 0.1‰ suggested
141 in Stracke et al., 2018, and similar to the long-term 2 S.D. of 0.06‰ given in Wang et al., 2021),
142 however the direction of partial melting fractionation is different. In the case of $\alpha_{\text{olivine-melt}} = 1$,
143 KLB1 melts are isotopically lighter than the bulk, and get heavier with increasing degree of
144 melting (the latter part of this behaviour is consistent with fractional melting modelled by
145 Stracke et al., 2018). For $\alpha_{\text{orthopyroxene-melt}} = 1$, partial melts are isotopically heavier than the
146 bulk, and get progressively lighter with increasing degree of melting (the former part of this
147 behaviour is consistent with Stracke et al., 2018).

148 For KLB1, the similar differences in modelled melting isotopic fractionation relative to the ana-
149 lytical uncertainty means that the choice of $\alpha_{\text{olivine-melt}} = 1$ or $\alpha_{\text{orthopyroxene-melt}} = 1$ for liquid
150 does not greatly alter any conclusions drawn from the peridotite melting. Minimal difference
151 between the models is also seen for MIX1G. However, low pressure melts from G2 (< 15 kbar)
152 show effectively no stable isotope fractionation when $\alpha_{\text{orthopyroxene-melt}} = 1$ is used, compared
153 to fractionations as large as $\Delta^{26}\text{Mg} = -0.1\text{‰}$ with $\alpha_{\text{olivine-melt}} = 1$. This is because, for these
154 shallow depths, the only stable Mg-bearing phases in G2 are clinopyroxene and liquid (garnet
155 is not stable). Orthopyroxene and clinopyroxene have similar β^{Mg} factors and therefore, when
156 liquid is given the same β^{Mg} factor as orthopyroxene, the coexisting liquid and clinopyroxene
157 are only minimally isotopically fractionated from each other. These results highlight how impor-
158 tant accurate experimental $\alpha_{\text{mineral-melt}}$ are for understanding the behaviour of stable isotope
159 systems, and the important role analytical precision plays in being able to identify these effects
160 in nature. Magnesium stable isotopes would benefit from further work in both of these areas.

161 4.2 Ca

162 The β^{Ca} value approach reduces uncertainties associated with an ionic model, as the bond lengths
163 and Ca, O coordination numbers in the relevant mineral structures required for the ionic model
164 are variable in the literature (e.g., Feng et al., 2014; Huang et al., 2019). Huang et al. (2019)
165 present an alternative set of β^{Ca} values for the Ca-bearing minerals we are concerned with,
166 although Antonelli et al. (2019) argue that the PBE functionals used in their own study better
167 represent Ca isotope fractionation between species than the less accurate LDA functionals used
168 by Huang et al. (2019). Additionally, Antonelli et al. (2019) calculate β^{Ca} for pyrope garnet,
169 which is more representative of mantle garnet than the grossular endmember used in Huang
170 et al. (2019). Our model uses $\beta_{\text{plagioclase}}^{\text{Ca}} < \beta_{\text{melt}}^{\text{Ca}} < \beta_{\text{clinopyroxene}}^{\text{Ca}}$ (see explanation in Section

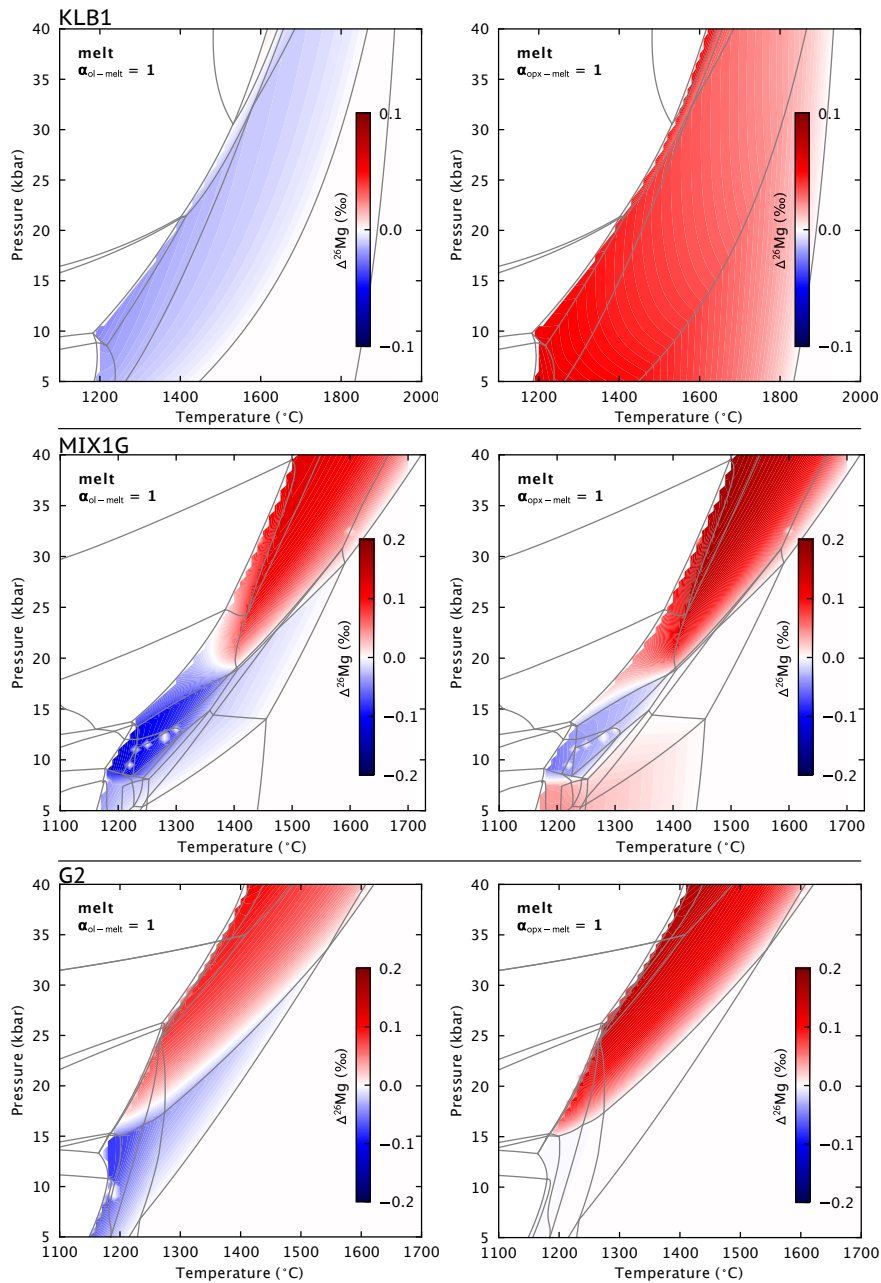


Figure S5: $\Delta^{26}\text{Mg}_{\text{melt-source}}$ compared for a model where $\alpha_{\text{olivine-melt}} = 1$ (left hand plots) and $\alpha_{\text{orthopyroxene-melt}} = 1$ (right hand plots) for all three lithologies considered.

171 3.2), but some studies do suggest that $\beta_{\text{melt}}^{\text{Ca}} = \beta_{\text{clinopyroxene}}^{\text{Ca}}$ (e.g., Chen et al., 2019) based on
 172 observations from natural equilibrated peridotite and pyroxenite samples that there is minimal
 173 Ca isotope fractionation between clinopyroxene and basaltic melt. Independent estimates of the
 174 force constants in silicate liquids and further experimental constraints on the behaviour of Ca
 175 isotopes in silicate melts are therefore needed to improve the Ca isotope fractionation model.

176 Fig. S6 shows the results of using the β^{Ca} values from Huang et al. (2019) with $\beta_{\text{melt}}^{\text{Ca}} =$
 177 $\beta_{\text{clinopyroxene}}^{\text{Ca}}$ (equivalent to $\alpha_{\text{clinopyroxene-melt}} = 1$). The use of isotopically lighter grossular
 178 garnet and isotopically heavier melt (relative to the Antonelli et al. (2019) inputs) result in
 179 overall less negative melt-source isotope fractionation than in the main model, making all Ca

180 isotope fractionations harder to resolve with current analytical precision. However, none of the
 181 conclusions presented in the main text are altered by these alternative parameter choices.

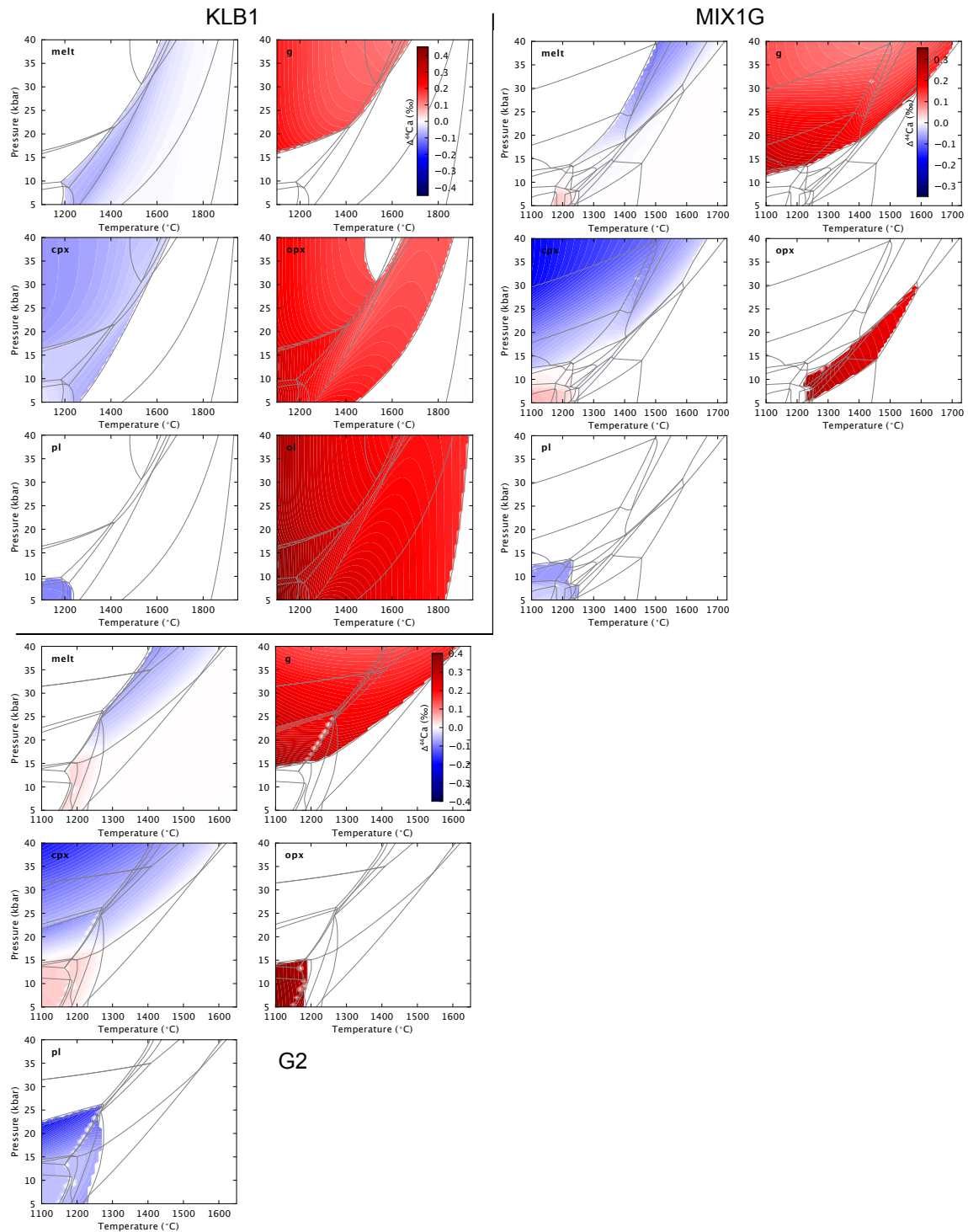


Figure S6: Model results of $\Delta^{44}\text{Ca}_{\text{phase-source}}$ for β^{Ca} inputs from Huang et al. (2019), and using $\beta_{\text{melt}}^{\text{Ca}} = \beta_{\text{clinopyroxene}}^{\text{Ca}}$. Generally, the calculated $\Delta^{44}\text{Ca}_{\text{phase-source}}$ is less negative than using the β^{Ca} inputs from Antonelli et al. (2019) as shown in Fig. S2.

182 The composition of the phases involved further complicates models of Ca isotope fractiona-
 183 tion. Feng et al. (2014) and Wang et al. (2017) show that Ca-O bond strength, hence isotopic
 184 fractionation, in orthopyroxene is dependent on the Ca and Fe content of the pyroxene, al-

185 though the concentrations of these elements in natural orthopyroxenes may produce negligible
186 concentration effects on isotopic fractionation (Wang et al., 2017). The isotopic fractionation
187 associated with garnet, in particular, is sensitive to Ca content for typical mantle compositions,
188 with a garnet-composition-sensitive fractionation factor produces smaller melting fractionations
189 in eclogite than using a fixed composition model (Chen et al., 2020a). Our model has not
190 implemented composition-specific fractionation factors, although does use the pyrope garnet
191 endmember which is most representative of mantle garnet. We also note that Antonelli et al.
192 (2019) calculate compositional effects for olivine and pyroxene solid solutions that are smaller
193 than those predicted using LDA functionals (e.g., Feng et al., 2014; Wang et al., 2017).

194 4.3 Fe

195 The model inputs for Fe isotope fractionation are relatively well constrained. The bonding
196 environment of Fe in mineral phases is well studied, making an ionic model approach possible
197 (e.g., Sossi and O'Neill, 2017). The model based on ionic inputs produces isotopic fractionation
198 and force constant predictions consistent with predictions by other methods and observations
199 (e.g., Dauphas et al., 2014; Macris et al., 2015). In addition, the Fe-O force constants in
200 basaltic glass as a function of Fe³⁺ content have been experimentally determined (Dauphas
201 et al., 2014), allowing the melt isotopic fractionation to be independently parameterised into
202 the model (although we note the potential complication induced by quenching effects, and by
203 melt generation in the mantle at non-atmospheric pressure). As discussed by Sossi and O'Neill
204 (2017), some terms in the ionic model approach (such as the ionicity of the bonds) are less well
205 constrained than the bonding environment of Fe in mineral structures, however the resultant
206 fractionation factors are minimally affected by these uncertainties.

207 4.4 V: V_{melt}⁵⁺ = [4] coordinated

208 It has been suggested that vanadium isotopes may track magmatic redox (Prytulak et al., 2013,
209 2017; Sossi et al., 2018b; Wu et al., 2018) due to the existence of three valence states (3+, 4+,
210 5+) in terrestrial magmatic systems (Canil, 1999), though a better understanding of V isotope
211 behaviour during magmatic processes is required to exploit the system's full potential as a redox
212 tracer (Prytulak et al., 2017). The successful modelling of V isotope fractionation therefore relies
213 on an understanding of the behaviour of each valence state of V.

214 Although slightly variable V partition coefficients for each valence state into mineral phases
215 can be found from experimental studies (Mallmann and O'Neill, 2009), these variations make
216 no significant difference to the calculated melt isotopic compositions. The main uncertainty
217 in the V fractionation model inputs is the coordination of V⁴⁺ and V⁵⁺ in silicate melts (and
218 associated β^V factors). V⁴⁺ in silicate liquid is [5–6] coordinated (Sutton et al., 2005), and
219 V⁴⁺ coordination of both [5.33] and [6] have been used in existing studies discussing V isotope
220 behaviour during melting (Qi et al., 2019; Wu et al., 2018 respectively). V⁵⁺ coordination in
221 silicate melts is [4–6] (Sutton et al., 2005; Righter et al., 2006), with Qi et al. (2019) quoting
222 an average of [5] coordination. As one of the main causes of V isotope fractionation during
223 melting is predicted to be the relative compatibility of V⁵⁺ (and to a lesser extent V⁴⁺) in
224 melts relative to minerals (e.g., Canil, 2002; Mallmann and O'Neill, 2009; Wu et al., 2018), the
225 bonding environment and force constants of V species in the melt is a particularly important
226 input. The choice of coordination number for the V species controls the $\ln \beta$ value chosen for
227 that species: β^V values are taken from those calculated for different bonding environments in

228 aqueous solutions (Wu et al., 2015), assumed to be a good approximation to the case of silicate
 229 melts following Wu et al. (2018). [5] coordination of V^{5+} is not modelled in Wu et al. (2015),
 230 and therefore only the cases of V^{5+} in [4] and [6] coordination can be used, for which β^V values
 231 have been calculated. No calculations are available for V^{4+} in [5] coordination, and therefore
 232 the model only considers V^{4+} in [6] coordination.

233 Fig. S7 compares the $\Delta^{51}V_{\text{melt-source}}$ for an input of $V_{\text{melt}}^{5+} = [6]$ -fold coordination (left hand
 234 panels; these are the results used in the main text) with an input of $V_{\text{melt}}^{5+} = [4]$ -fold coordination
 235 (right hand panels). There are two aqueous complexes with 4-fold V^{5+} in Wu et al. (2015). Here
 236 we have chosen the β^V factors for the $[V^{5+}O_2(OH)_2]^-$ (Table S10) complex since this complex
 237 has the greatest $\ln\beta^V$ difference from the [6]-fold coordinated case, therefore covering the greatest
 238 range of input uncertainty.

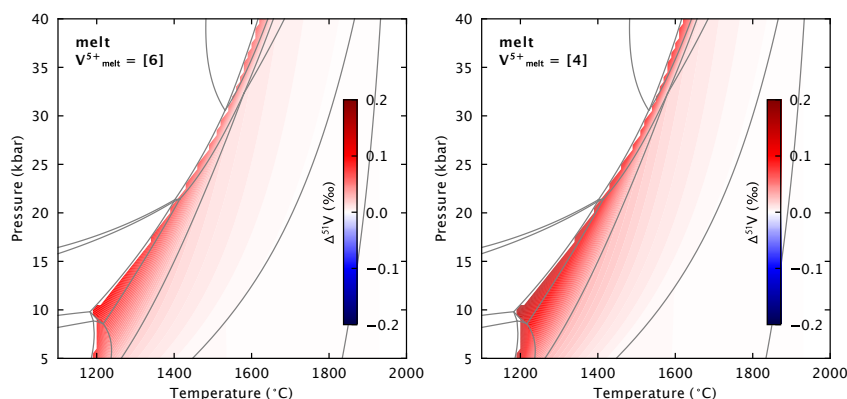


Figure S7: $\Delta^{51}V_{\text{melt-source}}$ compared for a model where $V_{\text{melt}}^{5+} = [6]$ coordinated (left hand plots) and $V_{\text{melt}}^{5+} = [4]$ coordinated (right hand plots) for KLB1.

Table S10: Alternate parameters for V isotope β^V factors from Wu et al. (2015). $1000\ln\beta = a/T^6 + b/T^4 + c/T^2$, where T = temperature (Kelvin).

Mineral	Valence	a	b	c
melt	V^{5+}	7.00×10^{14}	-3.25×10^{10}	1.42×10^6

239 The model discussed in the main text is the result of using the upper limit of [6] coordinated
 240 V^{5+} : this will provide the minimum melting fractionation, since a higher coordination number
 241 results in weaker V-O bonds, and less extreme isotopic fractionation. Using an input of [4]
 242 coordinated V^{5+} results in isotopically heavier melts, with the maximum KLB1 melting stable
 243 isotope fractionation increasing from ≈ 0.1 ‰ in the [6] coordinated case to ≈ 0.15 ‰ close to
 244 the solidus around 10 kbar (Fig. S6 for full results) in the [4] coordinated case. However, the
 245 difference between the two models decreases at higher temperature, and for the temperature
 246 range of interest here (potential temperatures from 1300 – 1530 °C), at 5 kbar the difference
 247 in melt isotopic composition between the two models is ≈ 0.01 ‰, and therefore negligible.
 248 However, to improve understanding of high temperature V behaviour and link natural data to
 249 specific processes, experimental and/or theoretical work on equilibrium isotopic fractionation
 250 factors for silicate melt and minerals for different V species is needed.

251 4.5 Cr: $\text{Cr}_{\text{melt}}^{2+} = [6]$ coordinated

252 As chromium is a first-group transition element, the properties controlling isotopic fractionation
 253 behaviour (e.g., valence, ionic radius hence bonding environment) of Cr^{2+} and Cr^{3+} ions may
 254 be expected to be similar to Fe^{2+} and Fe^{3+} ions (Xia et al., 2017). The main Cr-bearing
 255 phases in the mantle are spinel, garnet and pyroxenes, although some Cr will partition into
 256 olivine (Mallmann and O'Neill, 2009), and as with vanadium, partition coefficients depend on
 257 the valence state of Cr (Mallmann and O'Neill, 2009).

258 As a set of β^{Cr} factors for mantle minerals does not exist, our models of Cr stable isotope
 259 fractionation use force constants in each mineral calculated from an ionic bonding model (Shen
 260 et al., 2018), following the approach used for Fe isotopes. An ionic model has been shown to
 261 produce inter-mineral fractionation in agreement with the magnitude of isotope fractionation
 262 inferred from studies of lherzolite mineral pairs (Shen et al., 2018). As with V isotopes, the
 263 main uncertainty relating to species coordination is in liquid. Cr^{2+} in silicate melt is reported
 264 to be in [4] coordination (Miletich et al., 1999; O'Neill and Berry, 2006; Shen et al., 2020), but
 265 may also be in [6] coordination in a square planar arrangement (O'Neill and Berry, 2006). The
 266 isotope fractionation model presented here calculates force constants using effective anionic and
 267 cationic radii based on Shannon (1976), following the approach outlined in Shen et al. (2018),
 268 and these radii depend negatively on the coordination number chosen (Gibbs et al., 2014).

269 Fig. S8 compares the $\Delta^{53}\text{Cr}_{\text{melt-source}}$ for an input of $\text{Cr}_{\text{melt}}^{2+} = [4]$ coordinated (left hand panels;
 270 these are the results used in the main text) with an input of $\text{Cr}_{\text{melt}}^{2+} = [6]$ coordinated (right
 271 hand panels). The input parameters for the [6] coordinated case are in Table S8.

Table S11: Alternate parameters for Cr isotopic fractionation model. Mineral site and bonding information is based on Shen et al. (2018).

Mineral	Cr species	Site	Cr coordination	Bond length (Å)	O coordination
melt	2+	n/a	6	2.18	4

272 The model presented in the main text uses an input of Cr^{2+} in [4] coordination in the melt.
 273 Inputting [6] coordinated Cr^{2+} in the melt results in isotopically lighter peridotite melts across
 274 all of P-T space (Fig. S8) than the [4] coordinated case as a consequence of weaker Cr-O bonds
 275 in a [6] coordinated environment than [4] coordinated. A similar effect is seen in MIX1G and
 276 G2 melts. The melt-source isotopic fractionation is always negative, consistent with modelling
 277 by Shen et al. (2018), and melts get progressively isotopically heavier (towards 0‰ fractiona-
 278 tion from the source) with increasing degree of melting, consistent with observations from OIB
 279 (Bonnand et al., 2020). The cationic radii are likely to be overestimates, as the radii for [6]
 280 coordinated Cr used in Shen et al. (2018) assume octahedral, not square planar, coordination.
 281 Square planar bonds would be expected to be longer than octahedral, and this would result
 282 in even isotopically lighter melts. Therefore, although the model uses the most appropriate
 283 available inputs for the more commonly quoted [4] coordination of Cr^{2+} in melts, our under-
 284 standing of the behaviour of Cr isotopes in silicate melts would benefit from further work on
 285 the force constants of Cr^{2+} in glasses, or the inputs required to calculate these such as bonding
 286 environment or effective Cr-O bond lengths.

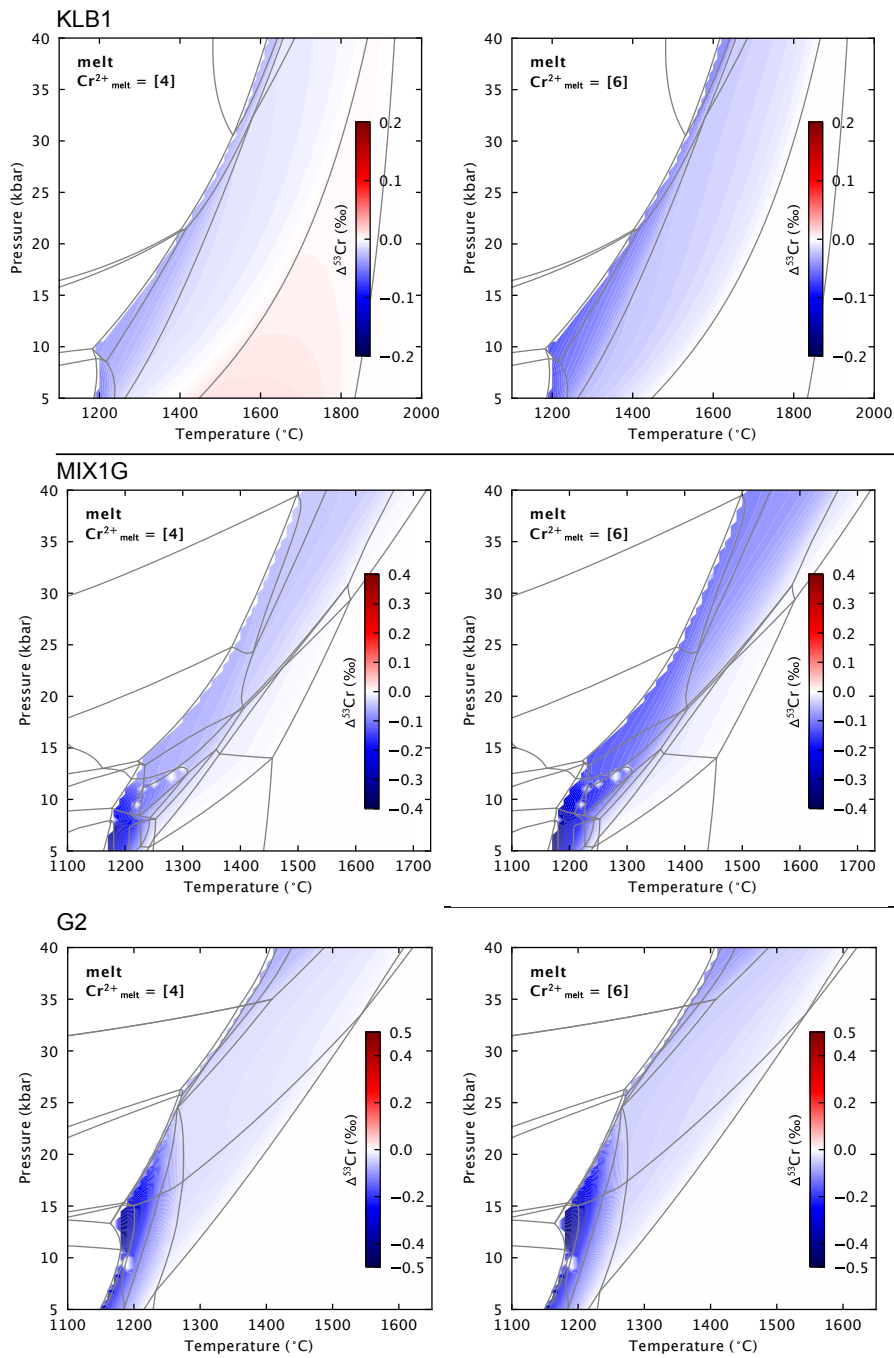


Figure S8: $\Delta^{53}\text{Cr}_{\text{melt-source}}$ compared for a model where $\text{Cr}^{2+}_{\text{melt}} = [4]$ coordinated (left hand plots) and $\text{Cr}^{2+}_{\text{melt}} = [6]$ coordinated (right hand plots) for all three lithologies.

287 5 Other potential isotopes

288 Other stable isotope systems not considered here might also have potential for investigating
 289 mantle heterogeneity (in either lithology or temperature). However, we require a better under-
 290 standing of the bonding environment in melts or relevant minerals before these isotope systems
 291 can be implemented into the stable isotope fractionation models used here.

292 Like other stable isotope systems, the equilibrium $\delta^{30}\text{Si}$ value of a phase will depend on the

293 strength of the Si-O bonds. A more polymerised (or, more Si-rich) phase will have more Si-O
 294 bonds, therefore a stronger average Si-O bond force constant, hence high $\delta^{30}\text{Si}$ (Grant, 1954).
 295 However, density functional theory (DFT) calculations show this simple prediction is compli-
 296 cated by the role of other network-modifying cations in silicates (Méheut et al., 2009; Méheut
 297 and Schauble, 2014). Silicon isotope fractionation factors have been calculated using DFT for
 298 appropriate mantle minerals (Huang et al., 2014), however there is a lack of work on the crucial
 299 mineral-melt fractionation information required for the equilibrium fractionation model here.
 300 Furthermore, it has been suggested that Si isotopic fractionation in liquids may be correlated
 301 with chemical composition (Qin et al., 2016), an effect that is too complicated (and under-
 302 constrained) to model here. This knowledge gap could provide scope for future work, however
 303 the mantle is currently thought to be homogeneous in Si isotopes (within current limits of an-
 304 alytical precision), and there seems to be only a small partial melting effect (Savage et al.,
 305 2014).

306 Titanium stable isotopes have been suggested as tracer of magmatic processes (Millet et al.,
 307 2016), since Ti can exist in multiple coordination states in silicate melts and Ti-oxides (Farges
 308 and Brown Jr, 1997). However, Ti stable isotope fractionation in magmatic systems is dominated
 309 by the behaviour of Fe-Ti oxides, with no partial melting fractionation from mantle lithologies
 310 expected unless under conditions of Ti-oxide saturation. DFT calculations predict no significant
 311 Ti isotope fractionation between pyroxenes, olivine or garnet (Wang et al., 2020), and both
 312 eclogites and MORB have the same measured $\delta^{49}\text{Ti}$ value (Millet et al., 2016). Therefore, Ti
 313 isotopes are not discussed further here, as there is no existing data to support for their use in
 314 tracing temperature or lithological heterogeneity in the mantle associated with MORB or OIB.

315 6 Average depth of melting

316 The mean melt pressure, \bar{P} , for isentropic melting at any T_p is calculated by the following
 317 integral, based on Asimow et al. (2001) with an additional geometric weighting factor to account
 318 for the triangular shape of the melting region

$$\bar{P} = \frac{\int_{P_0}^{P_f} P \Delta F g dP}{\int_{P_0}^{P_f} \Delta F g dP} \quad (8)$$

319 where P_0 and P_f are the initial and final pressures of melting respectively, ΔF is the amount
 320 of melt produced at each pressure step (dP) along the isentropic melting path, and g is the
 321 geometric weighting factor. P_0 is taken as the intersection of the isentrope with the solidus,
 322 P_f is used as 5 kbar in this model as melt compositions have not been calculated for shallower
 323 pressures - this is an overestimate (i.e., final melting will be shallower) for $T_p = 1300^\circ\text{C}$ and an
 324 underestimate for $T_p = 1530^\circ\text{C}$ (Asimow et al., 2001), but will provide the correct behaviour.
 325 Pressure steps used are the 0.5 kbar intervals used in the P-T grid. ΔF is calculated for each
 326 step from the difference in melt fractions calculated by THERMOCALC for consecutive dP steps.
 327 The geometric weighting factor is a linear scaling running from 1 at the base of the melting
 328 region to 0 at the top, to account for the wider base of the triangular melting region beneath
 329 ridges. We calculate the average melting pressure for the three T_p s for which we have calculated
 330 isentropes in KLB1, and lineally interpolate between them (this will not be a linear relationship
 331 in reality, e.g., Asimow et al., 2001, but our approach will provide a good approximation) to
 332 get average melt pressures for temperatures between $T_p = 1300^\circ\text{C}$ and 1530°C . For KLB1, we
 333 calculate an average melting pressure of approximately 12 kbar at $T_p = 1300^\circ\text{C}$, and 28 kbar
 334 at $T_p = 1530^\circ\text{C}$. The average $\Delta_{\text{melt-source}}$ for each T_p is then taken as the $\Delta_{\text{melt-source}}$ at the

335 calculated average melting pressure.

336 The approach above, however, uses an equilibrium melt isotope composition at a single pressure
 337 to approximate the fractional, multibaric melting thought to occur in the mantle (McKenzie,
 338 1984; von Bagen and Waff, 1986; Langmuir et al., 1992; Asimow et al., 2001). To assess the
 339 suitability of the equilibrium melting at single pressure approach, we have also calculated the
 340 average melt isotopic composition by considering the stable isotope composition and elemental
 341 concentration for each new step of melt generated during isentropic decompression at $T_p =$
 342 1300°C and 1530°C by

$$\bar{\Delta}_{\text{melt-source}} = \frac{\int_{P_0}^{P_f} \Delta_x c \Delta F g dP}{\int_{P_0}^{P_f} c \Delta F g dP} \quad (9)$$

343 where Δ_x is the $\Delta_{\text{melt-source}}$ of the melt produced in the pressure step, and c is the wt % of
 344 the element being considered in the melt, and other terms as above. The approach of equation
 345 7, while considering polybaric melting, is still not a fractional melting model, but fractional
 346 melting is beyond the scope of this work. Nonetheless, the chemical composition of pooled
 347 fractional melts from a ridge melting regime has been shown to closely resemble that of a simple
 348 equilibrium batch melt (Langmuir et al., 1992). Equation 7 therefore allows us to broadly assess
 349 whether the approach taken in equation 6 produces estimates of mean melt isotope ratios very
 350 different from a more realistic polybaric approach, although it should be noted that the polybaric
 351 method of equation 7 only allows us to calculate the average melt isotopic composition for the
 352 two T_p s for which we have calculated isentropes.

353 Fig. S9 shows the comparison of the two calculations of average $\Delta_{\text{melt-source}}$. The dots show the
 354 average melt isotope composition when calculated as a sum along the isentrope; the solid lines
 355 show the approximation using the melt composition at the average pressure of melting. The
 356 calculation using the average melt pressure is a good approximation of the calculated average
 357 melt isotope compositions. Therefore we use this pressure approach in the main text: the
 358 limitation of the calculated average melt isotope composition is that we only have information for
 359 the T_p s for which we have calculated isentropes, hence use a linear interpolation between them,
 360 whereas the average pressure approach allows us to use more information from the calculated
 361 P-T map of melt compositions, reflecting the changing of stable mineral assemblage, for example.

362 7 Trace element calculations

363 Table S12 gives the mineral-melt partition coefficients and mantle concentrations for Sm and
 364 Yb used in this model.

Table S12: Input parameters (partition coefficients relative to melt, mantle concentrations) for calculating the Sm, Dy and Yb concentration of trace elements in the model melts. $D_{\text{gt,cpx,ol,opx}}$ and element concentrations from Gibson and Geist (2010); D_{spinel} from Liu et al. (2019) (average values from high P experiments); D_{plag} from Aigner-Torres et al. (2007) (from ‘run 31’); D_{rutile} from Klemme et al. (2005) (Sm), Foley et al. (2000) (Dy, Yb); D_{ilmenite} from Zack and Brumm (1998). Conc.DM = concentration in depleted mantle; mineral abbreviations as previous with rut. = rutile, ilm. = ilmenite.

Element	D_{gt}	D_{cpx}	D_{ol}	D_{opx}	D_{spn}	D_{plag}	D_{rut}	D_{ilm}	Conc.DM
Sm	0.25	0.293	0.0011	0.02	0.198	0.081	0.00135	0.00059	0.299
Dy	2.2	0.4	0.0027	0.011	0.069	0.049	0.00076	0.01	0.525
Yb	6.6	0.4	2	0.08	0.812	0.025	0.0093	0.17	0.347

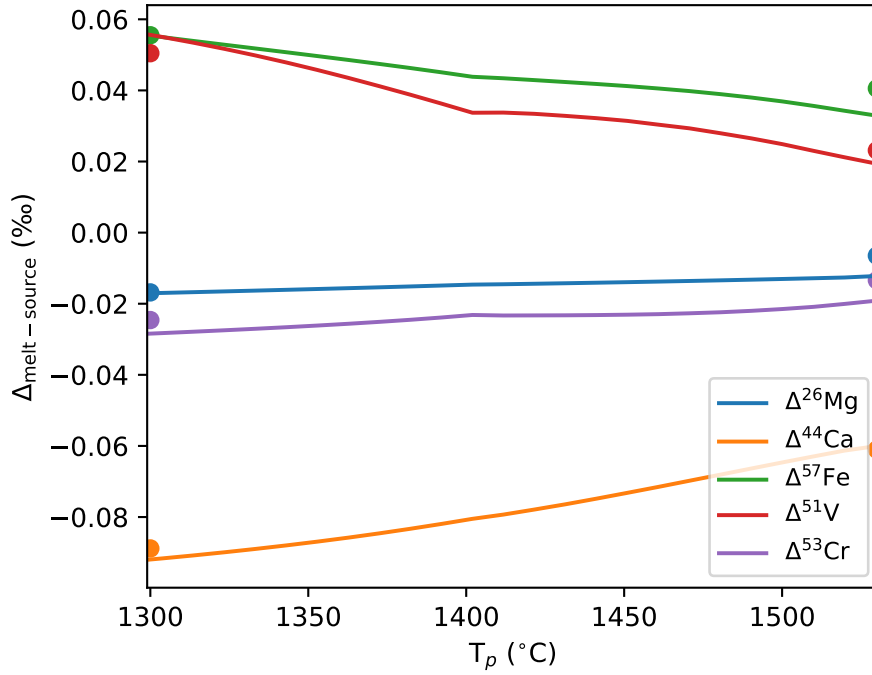


Figure S9: $\Delta_{\text{melt-source}}$ for KLB1 calculated by taking the isotopic fractionation at the average melt pressure (solid lines) compared to a calculation of the average melt composition by considering polybaric melts (equation 7), shown as coloured dots.

365 The concentration of each element in the melt was calculated at each P-T point by:

$$C_{\text{liquid}} = \frac{C_{\text{bulk}}}{\sum D_{\text{mineral}} F_{\text{mineral}} + F_{\text{liquid}}} \quad (10)$$

366 where C = concentration, D = partition coefficient, F = fraction of phase.

367 8 Additional reference figures

368 Fig. S10 provides a reference for how oxygen fugacity, melt fraction, and selected trace element
 369 ratios vary along the isentropes shown in Figs. 1 and 4 of the main text. The trace element
 370 concentrations are calculated using depleted mantle (DM) inputs for KLB1, and plume inputs
 371 for the enriched MIX1G and G2 lithologies (see Table S12).

372 We choose to use peridotite isentropic decompression paths for the two pyroxenite lithologies, as
 373 discussed in the main text. Figs. S11, S12 show results for the other endmember case for the two
 374 pyroxenite lithologies, i.e., where the lithologies decompress along isentropic paths calculated
 375 for each pure pyroxenite lithology.

376 Fig. S13 compares the melt-source isotope fractionation with progressive melting (along peri-
 377 dotite isentropes) for pairs of isotopes in each lithology, allowing assessment of the ‘direct’
 378 sensitivity to lithology. By direct lithological sensitivity, we refer to sensitivity related to miner-
 379 alogy – these isotope systems may still trace lithology via each lithology’s different bulk isotope
 380 composition. On balance, given that melts from any pyroxenite lithology are likely to be diluted
 381 with ambient peridotite melts when extracted from the mantle (e.g., Hirschmann and Stolper,

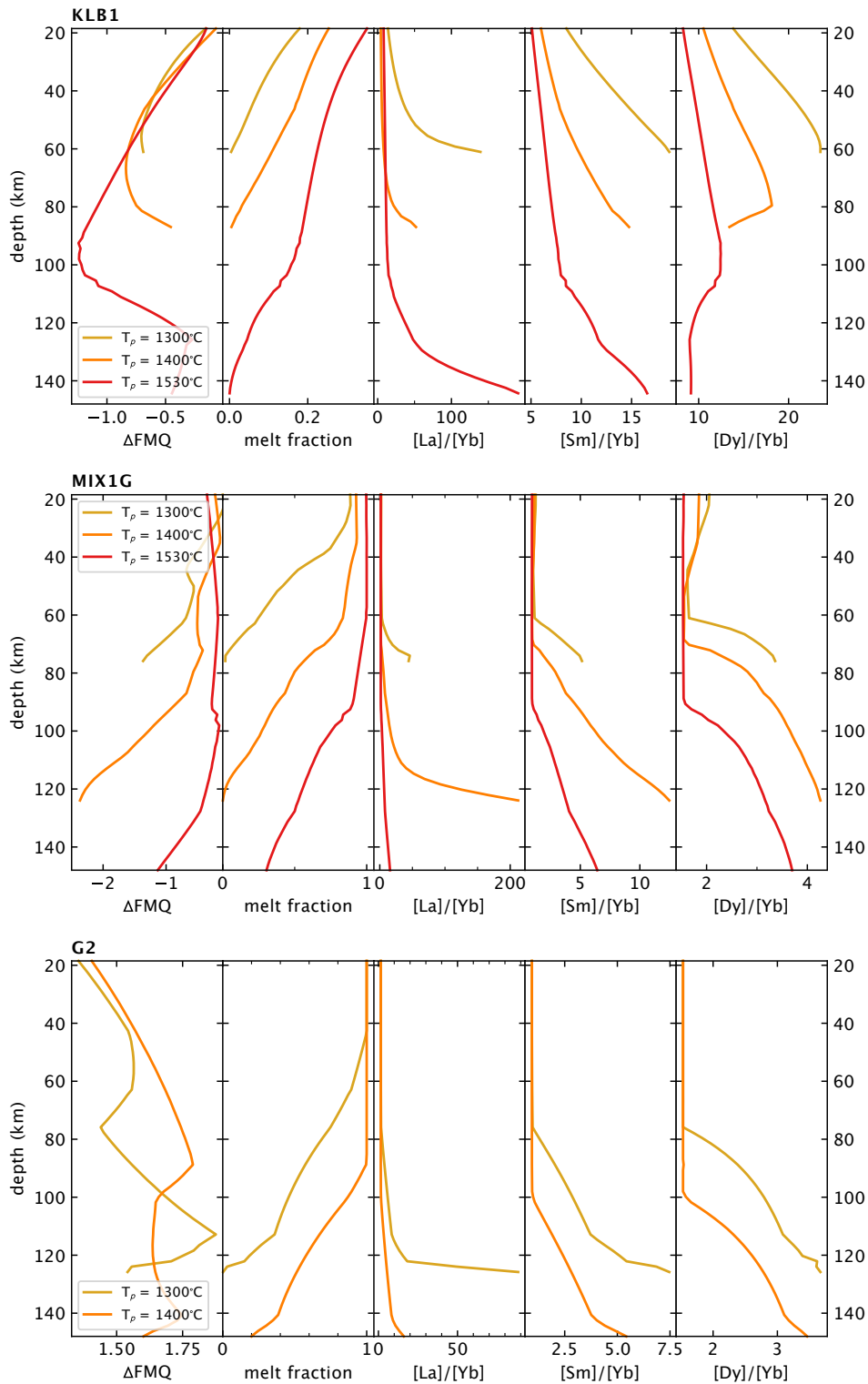


Figure S10: Parameters along the peridotite isentropic decompression paths for all three lithologies: oxygen fugacity (from the $f\text{O}_2$ melt software; Holland et al., 2018), melt fraction (from THERMOCALC output), and selected trace element ratios in the melts (calculated using the model described in the section above). The isentropes match those in Figs. 2 and 5 in the main text.

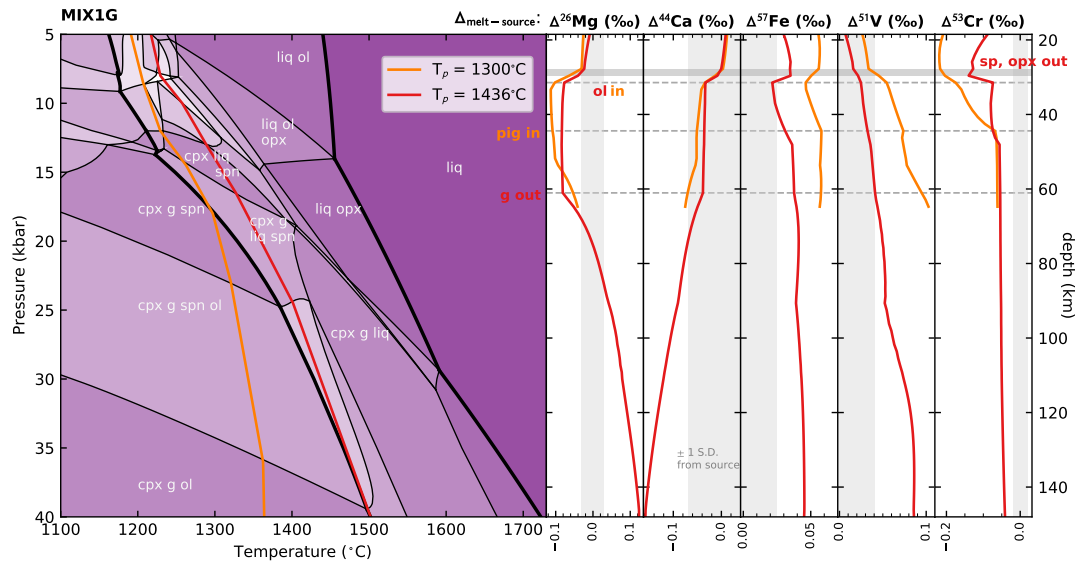


Figure S11: Isentropic melting paths for pure MIX1G, and $\Delta_{\text{melt-source}}$ along these isentropes. Errors are shown as 1 S.D. long-term uncertainty.

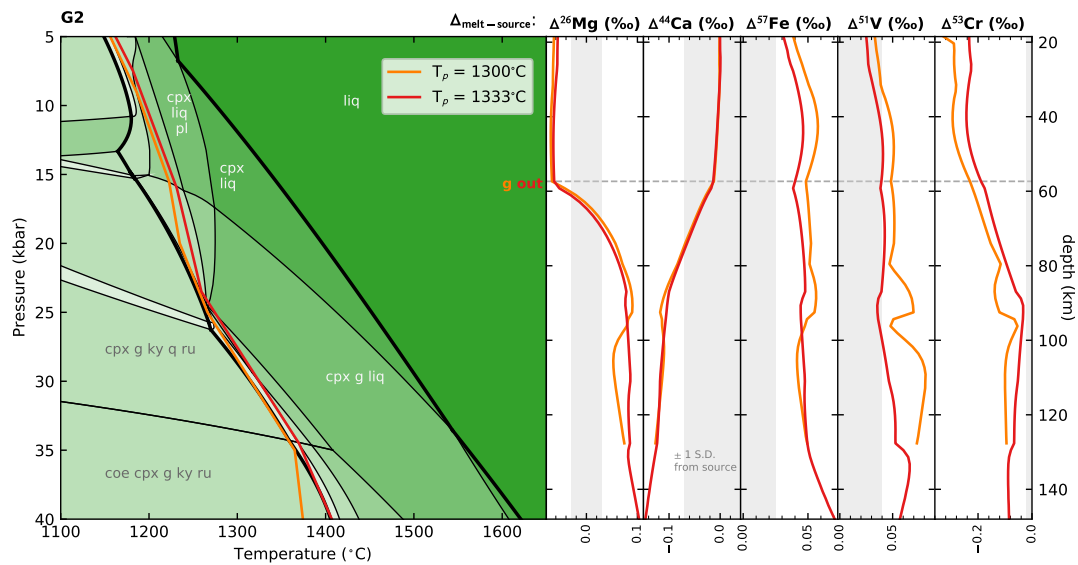


Figure S12: Isentropic melting paths for pure G2, and $\Delta_{\text{melt-source}}$ along these isentropes. Errors are shown as 1 S.D. long-term uncertainty.

1996; Sobolev et al., 2005; Shorttle and MacLennan, 2011; Soderman et al., 2021), our results in Fig. S13 are probably representative of maximum isotope fractionation signatures we might see in natural systems, and Fig. 7 in the main text (which takes into account both this dilution effect and non-identical bulk isotope compositions of the lithologies) may better represent the predicted resolvability of each isotope system to lithology. Generally, Fig. S13 suggests that the stable isotopes considered here may have poor direct sensitivity to lithological heterogeneity in the source regions of basalts, assuming equilibrium fractionation behaviour and in the absence of any pre-imposed source heterogeneity (i.e., where all lithologies have the same bulk isotope ratio).

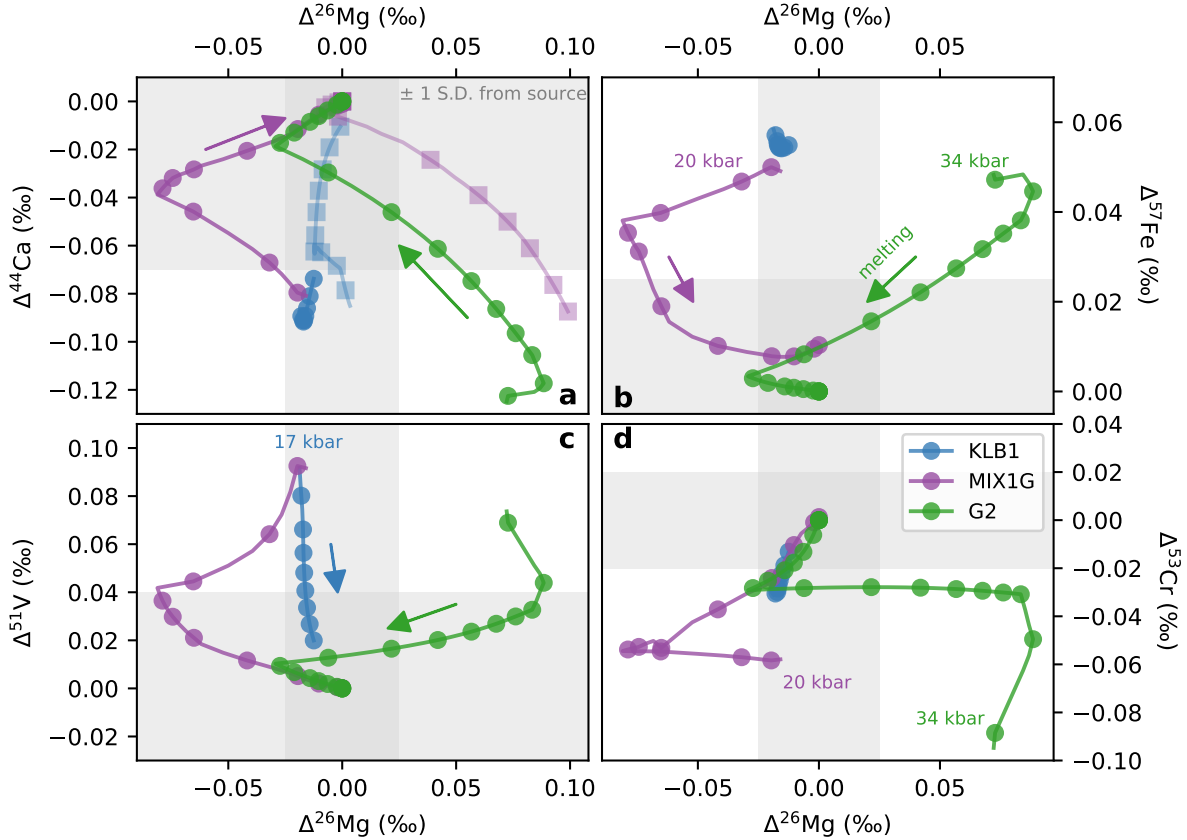


Figure S13: Melting fractionation ($\Delta_{\text{melt-source}}$) plotted along the $T_p = 1300^\circ\text{C}$ isentrope; squares in Mg–Ca space show $T_p = 1530^\circ\text{C}$ isentrope for reference. High or low pressure end of isentrope is marked, and arrows show the direction of progressive melting (decompression). Markers are plotted every 3 pressure steps (each pressure step is 0.5 kbar). Grey shaded regions show the source (0‰) \pm typical long-term analytical 1 S.D. (see Table S2 for details).

391 9 Element budgeting between lithologies

392 To produce Fig. 7 in the main text, we use a mass balance calculation for each isotopic system,
393 X, where

$$\delta X_{\text{bulk}} = (F_{\text{pyroxenite}}^X \times \delta X_{\text{pyroxenite}}) + (F_{\text{peridotite}}^X \times \delta X_{\text{peridotite}}) \quad (11)$$

394 and F^X is the fraction of the total element in the final melt coming from each lithology, $\delta X_{\text{lithology}}$
395 is the isotopic composition of each lithology (bulk isotope composition of lithology + melt-source
396 isotope fractionation). The F term considers both the amount of the element in the melt from

Table S13: Inputs used to calculate bulk melt-source isotope fractionation in Figs. 7 and S14. * indicates inputs using the pure pyroxenite isentropes at that T_p , rather than peridotite isentropes. The MIX1G 1400* inputs are actually for an isentrope of $T_p = 1436^\circ\text{C}$. The P-T given are the closest P-T pair in the model output files to the average P-T of melting along each isentrope, as calculated in Section 6. For V and Cr, where element partitioning is not calculated by THERMOCALC, the melt element concentrations are calculated as described in Appendix Section 3, using the following bulk compositions for each lithology: KLB1 $\text{Cr}_2\text{O}_3 = 0.32\text{ wt } \%$, G2 $\text{Cr}_2\text{O}_3 = 0.08\text{ wt } \%$, MIX1G $\text{Cr}_2\text{O}_3 = 0.11\text{ wt } \%$ (Jennings et al., 2016); KLB1 V = 85 ppm (typical primitive mantle value from Lee et al., 2003; Prytulak et al., 2013), MIX1G and G2 V = 350 ppm (typical MORB value; Prytulak et al., 2013).

T_p ($^\circ\text{C}$):	KLB1		MIX1G				G2	
	1300	1400	1300	1300*	1400	1400*	1300	1300*
Pressure (kbar)	12	17.5	15	13.5	24.5	22	27	25.5
Temperature ($^\circ\text{C}$)	1300	1400	1317	1248	1440	1380	1345	1269
Melt fraction	0.061	0.099	0.323	0.093	0.369	0.155	0.458	0.048
Melt element conc.								
Mg (mol % oxide)	16.50	20.43	15.44	11.01	19.42	16.59	5.92	4.48
Ca (mol % oxide)	12.72	13.62	12.66	10.63	13.09	11.83	9.47	7.53
Fe (mol % oxide)	6.36	7.64	8.93	7.46	9.27	10.21	4.22	3.07
V (wt % element)	0.002	0.003	0.014	0.008	0.015	0.009	0.021	0.011
Cr (wt % element)	0.005	0.009	0.004	0.001	0.007	0.002	0.004	0.001
Melt-source isotopic fractionation (‰)								
$\Delta^{26}\text{Mg}$	-0.02	-0.02	-0.08	-0.10	0.08	0.02	0.06	0.09
$\Delta^{44}\text{Ca}$	-0.09	-0.08	-0.04	-0.05	-0.08	-0.08	-0.08	-0.12
$\Delta^{57}\text{Fe}$	0.06	0.04	0.03	0.06	0.03	0.04	0.03	0.06
$\Delta^{51}\text{V}$	0.05	0.03	0.03	0.08	0.02	0.05	0.03	0.07
$\Delta^{53}\text{Cr}$	-0.03	-0.02	-0.05	-0.06	-0.04	-0.05	-0.03	-0.12
Bulk isotopic composition of lithology (‰)								
$\delta^{26}\text{Mg}$	-0.24		-0.25				-0.25	
$\delta^{44}\text{Ca}$	0.94		0.83				0.83	
$\delta^{57}\text{Fe}$	0.05		0.14				0.14	
$\delta^{51}\text{V}$	-0.91		-0.92				-0.92	
$\delta^{53}\text{Cr}$	-0.12		-0.15				-0.15	

397 each lithology, and the melt fraction of each lithology. Table S13 gives the inputs used: these
398 inputs are taken from the full model results at the average P-T melting condition (as in Section
399 6) for each T_p . This calculation cannot be done for a T_p of 1530°C (MIX1G & G2) or 1400°C
400 (G2) because the isentropic melting paths intersect the solidus at pressures below the base of
401 our model, so an accurate average pressure of melting cannot be calculated. The bulk isotope
402 compositions of each lithology are the BSE values for peridotite (see Fig. 1 caption in main text
403 for sources), and average MORB composition for pyroxenite (except for Cr where no known
404 MORB data are published and the modelled aggregate KLB1 melt isotope composition for T_p
405 = 1300°C is used).

406 Fig. S14 is an alternative to Fig. 7 in the main text, with the right hand panels scaled to typical
407 current 1 S.D. analytical precision.

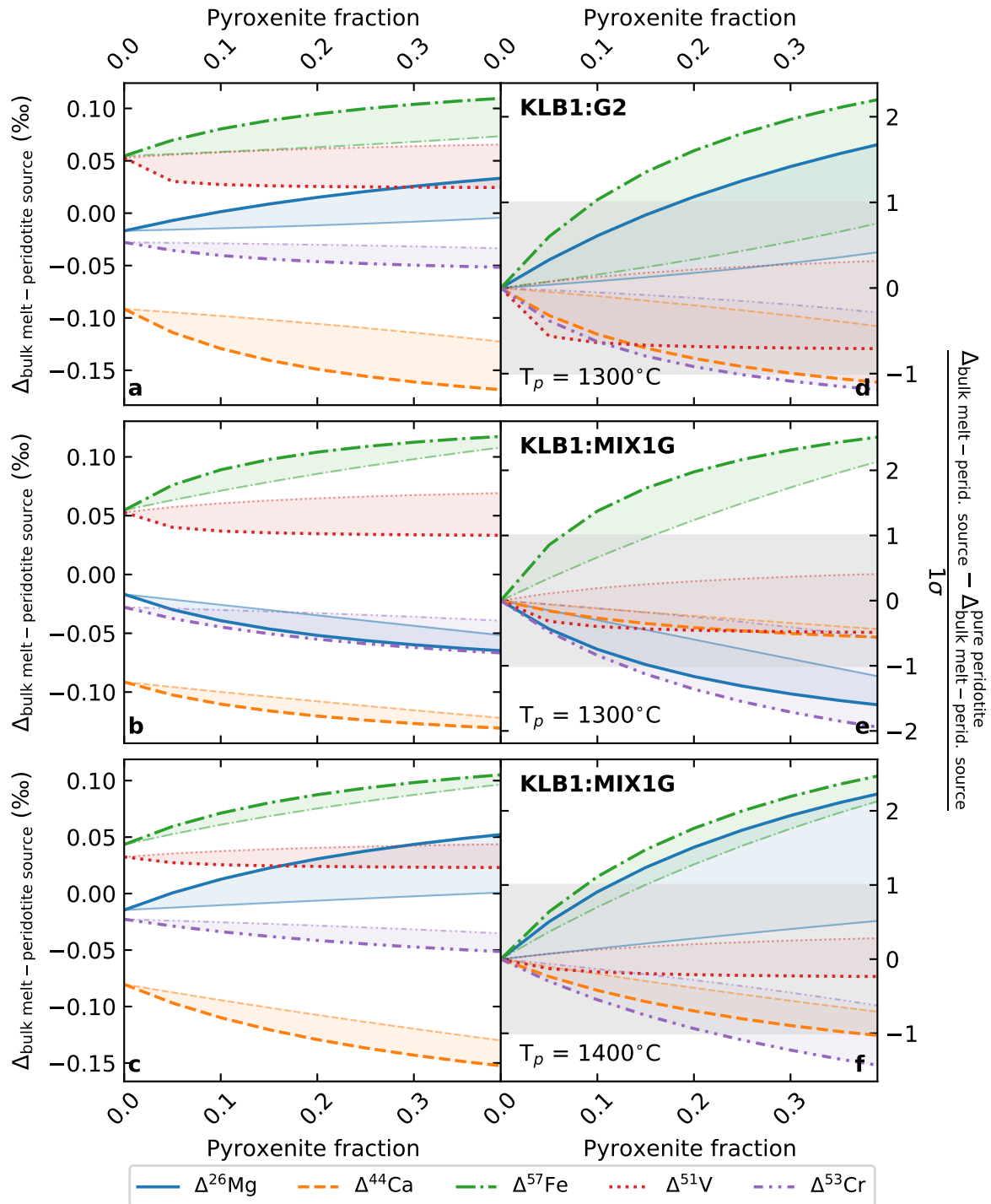


Figure S14: Aggregate bulk melt isotope fractionation (relative to a pure peridotite source, left hand panels; relative to a pure peridotite melt and scaled to current typical long-term analytical 1 S.D., right hand panels) for multi-lithologic mantle. At each T_p , the aggregate melt composition (including isotope composition, and element proportion) at the average pressure and temperature of melting for each lithology is used to calculate the resulting bulk melt isotope composition by mass balance. A range of bulk melt-source isotope fractionations are given for each isotope system, reflecting two endmembers for the thermal behaviour of pyroxenite upwelling in a peridotite host. The bolder line represents the case for pyroxenite following peridotite geotherms (as used throughout this study); the paler line reflects pyroxenite following its own geotherm, as in Figs. S10 and S11. Grey horizontal bars in the right hand panels highlight 1 S.D.

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APPENDIX

Supplementary Material to ‘Global trends in novel stable isotopes in basalts: theory and observations’ by Soderman et al.

1 Data sources

The data sources used in the natural data compilation in Figs. 1 and 8 in the main text are given in Table S1.

Table S1: Literature data sources for stable isotope data compilation.

Isotope	Type	Literature sources
$\delta^{26}\text{Mg}$	MORB	Wiechert and Halliday (2007); Bourdon et al. (2010); Teng et al. (2010); Zhong et al. (2021b)
	OIB	Bourdon et al. (2010); Teng et al. (2010); Zhong et al. (2017); Wang et al. (2021)
$\delta^{44}\text{Ca}$	MORB	Zhu et al. (2018); Chen et al. (2020b); Zhu et al. (2020a)
	OIB	Huang et al. (2011b); Valdes et al. (2014, 2019); Feng et al. (2017); Chen et al. (2020b)
$\delta^{57}\text{Fe}$	MORB	Teng et al. (2013); Nebel et al. (2013); Zhong et al. (2021) (seamounts: Sun et al., 2020)
	OIB	Teng et al. (2008); Schuessler et al. (2009); Konter et al. (2016); Nebel et al. (2019) Peters et al. (2019); Gleeson et al. (2020); Soderman et al. (2021)
$\delta^{51}\text{V}$	MORB	Prytulak et al. (2013); Wu et al. (2018); Novella et al. (2020)
	OIB	Prytulak et al. (2013); Ding et al. (2020)
$\delta^{53}\text{Cr}$	OIB	Bonnand et al. (2020); Shen et al. (2020)

Table S2 gives the references for the typical sample 2 S.E. and long-term 2 S.D. used throughout the manuscript.

Table S2: Literature data sources for stable isotope data errors.

Isotope	Type	Value (‰)	Literature sources
$\delta^{26}\text{Mg}$	sample 2 S.E.	0.02	Wang et al. (2021)
	long-term 2 S.D.	0.06	Wang et al. (2021)
$\delta^{44}\text{Ca}$	sample 2 S.E.	0.05	Zhu et al. (2018, 2020a)
	long-term 2 S.D.	0.14	Chen et al. (2019)
$\delta^{57}\text{Fe}$	sample 2 S.E.	0.02	Soderman et al. (2021)
	long-term 2 S.D.	0.05	Soderman et al. (2021)
$\delta^{51}\text{V}$	sample 2 S.E.	0.05	Wu et al. (2018); Novella et al. (2020)
	long-term 2 S.D.	0.08	Qi et al. (2019)
$\delta^{53}\text{Cr}$	sample 2 S.E.	0.01	Sossi et al. (2018)
	long-term 2 S.D.	0.04	Xia et al. (2017); Sossi et al. (2018)

8 2 Phase equilibria calculations

9 In order to explore stable isotope fractionation taking place during the melting of pyroxenite
10 and peridotite lithologies within the mantle, we first constructed pseudosections and calculated
11 mineral chemistries for MIX1G (KLB1 and G2 calculations are presented in Soderman et al.,
12 2021). The calculations were performed in the KNCFMASSTOcr system using THERMOCALC
13 v3.47 (Powell et al., 1998), and the latest version of the thermodynamic dataset of Holland and
14 Powell (2011) with the a-X models of Holland et al. (2018).

15 The composition of the MIX1G lithology (Hirschmann et al., 2003; Lambart et al., 2016) used
16 in the phase-equilibria calculations is given in Table S3. The ferric iron content is taken as
17 between that of KLB1 peridotite and MORB, following the approach used for KG1 by Jennings
18 et al. (2016), where $\text{Fe}^{3+}/\text{Fe}_T = 0.1$.

Table S3: MIX1G composition, in KNCFMASSTOcr system (mol % oxides). Composition from Lambart et al. (2016), renormalised without MnO.

SiO ₂	Al ₂ O ₃	CaO	MgO	mole % oxide					
				FeO _t	K ₂ O	Na ₂ O	TiO ₂	Cr ₂ O ₃	O
38.5	1.78	2.82	50.58	5.52	0.01	0.25	0.07	0.11	0.09

19 After calculation of the phase boundaries, THERMOCALC was run over a grid of P-T space, from
20 5–40 kbar and 1100–1730 °C.

21 3 Details and results of the main stable isotope fractionation 22 model

23 The equilibrium isotopic composition of the phases present at any point in P-T space (including
24 melts) was calculated for each isotope system, based on calculated equilibrium isotope fraction-
25 ation factors, α^{A-B} , between the phases that are present. α^{A-B} was calculated either from
26 reduced partition function ratios (β factors), or from calculated cation-oxygen force constants,
27 as outlined for Mg, Ca, V and Cr isotopes below. The Fe isotope fractionation input parameters
28 are presented in the supplementary information of Soderman et al. (2021). The mineral abbrev-
29 viations used in the figures in the following subsections are: olivine (ol), clinopyroxene (cpx),
30 orthopyroxene (opx), spinel (spn), garnet (g), plagioclase (pl).

31 3.1 Mg

32 The β^{Mg} factors used for Mg-bearing mantle minerals (spinel, clinopyroxene, orthopyroxene,
33 olivine, garnet) are given in Table S4. These β^{Mg} factors are taken from Huang et al. (2013), with
34 the exception of spinel, which is taken from Schauble (2011) and scaled relative to a reference
35 β^{Mg} of olivine which is calculated in both studies. Pigeonite (a stable phase in MIX1G) is
36 modelled with the same β^{Mg} as clinopyroxene; liquid (melt) is taken as having no fractionation
37 from olivine in the main model (Stracke et al., 2018, see Appendix Section 4 for discussion), so
38 $\beta_{\text{melt}}^{\text{Mg}} = \beta_{\text{olivine}}^{\text{Mg}}$.

39 The full results of the Mg model, showing the isotopic composition of each phase calculated in

Table S4: Parameters for β^{Mg} factors from Huang et al. (2013). $1000\ln\beta = a/T^6 + b/T^4 + c/T^2$, where T = temperature (Kelvin).

Mineral	a	b	c
spinel	-2.74×10^{16}	2.08×10^{11}	2.93×10^6
orthopyroxene	-8.10×10^{15}	6.65×10^{10}	2.25×10^6
clinopyroxene	-1.17×10^{16}	1.03×10^{11}	2.24×10^6
olivine	-1.19×10^{16}	1.05×10^{11}	2.07×10^6
garnet	-5.50×10^{15}	4.92×10^{10}	1.30×10^6
melt	-1.19×10^{16}	1.05×10^{11}	2.07×10^6

40 KLB1, MIX1G and G2, are in Fig. S1.

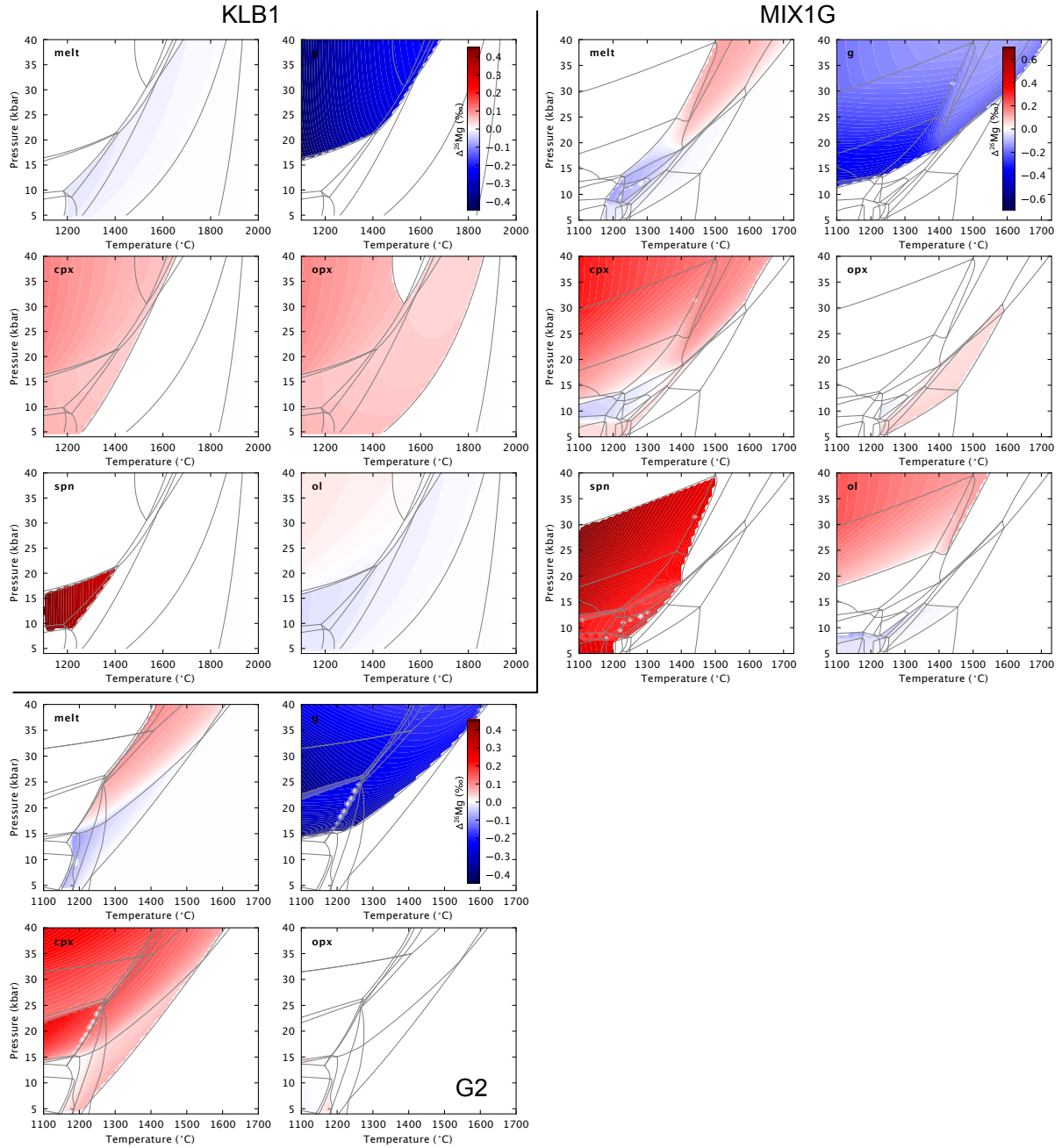


Figure S1: Full results of the Mg isotopic fractionation model, $\Delta^{26}\text{Mg}_{\text{phase-source}}$, for the three lithologies.

Table S5: β^{Ca} factors (given as $1000\ln\beta$, at a temperature of 1000K) from Antonelli et al. (2019).

Mineral	Structure used	$1000\ln\beta$
plagioclase	anorthite	1.06
orthopyroxene	Ca/Mg = 1/32	1.84
clinopyroxene	diopside	1.32
olivine	forsterite Ca/Mg = 1/64	2.11
garnet	pyrope Ca/Mg = 1/24	2.00
liquid		1.19

41 3.2 Ca

42 The β^{Ca} factors used for Ca-bearing mantle minerals (plagioclase, clinopyroxene, orthopyroxene,
43 olivine, garnet) are given in Table S5, and are taken from Antonelli et al. (2019). Where options
44 exist for minerals with different structures and/or compositions, we use those with most similar
45 Ca/Mg to the typical phase compositions calculated by THERMOCALC for KLB1. Pigeonite (a
46 stable phase in MIX1G) is modelled with the same β^{Ca} as clinopyroxene; liquid is taken as having
47 β^{Ca} intermediate between clinopyroxene and plagioclase (anorthite), following observations from
48 natural samples (Zhang et al., 2018; Antonelli et al., 2019b) and consistent with the approach
49 used by Antonelli et al. (2021). β^{Ca} factors are made temperature-dependent by

$$\ln\beta_{(T)} = \ln\beta_{(1000\text{K})} \times \frac{10^6}{T^2} \quad (1)$$

50 following Antonelli et al. (2021).

51 The full results of the Ca model, showing the isotopic composition of each phase calculated in
52 KLB1, MIX1G and G2, are in Fig. S2.

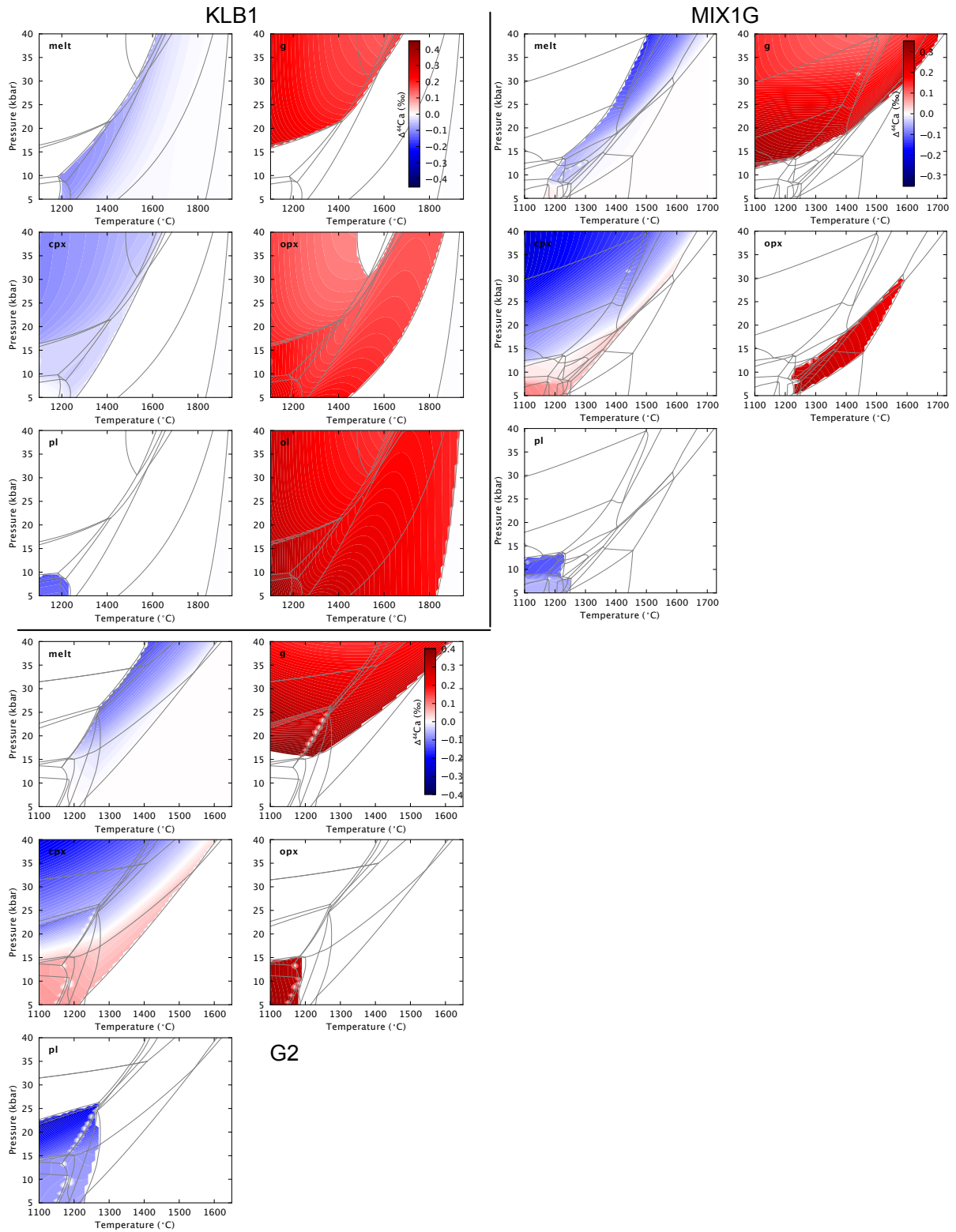


Figure S2: Full results of the Ca isotopic fractionation model, $\Delta^{44}\text{Ca}_{\text{phase-source}}$, for the three lithologies.

53 3.3 V

54 The proportion of $V^{3+/4+/5+}$ was calculated following Toplis and Corgne (2002), using the
 55 FeO/Fe_2O_3 ratio of the melt (calculated by THERMOCALC) at each point in P-T space with

$$\log(XV^{3+}/XV^{4+}) = \log(XFe^{2+}/XFe^{3+}) - 2.4 \quad (2)$$

56 and

$$\log(XV^{4+}/XV^{5+}) = \log(XFe^{2+}/XFe^{3+}) - 0.3 \quad (3)$$

57 where X is the mole fraction of cation.

58 The partition coefficients and β^V factors of each valence state of V in the V-bearing mantle
 59 minerals (olivine, clinopyroxene, orthopyroxene, garnet, spinel) are given in Table S6. Partition
 60 coefficients are from Mallmann and O'Neill (2009), with garnet assumed to have similar V
 61 partitioning to clinopyroxene (Mallmann and O'Neill, 2007, 2009). The β^V factors are taken from
 62 Wu et al. (2015) for V in different coordination environments in solution systems; alternative
 63 choices of β^V due to uncertainties in V coordination number are discussed in Section 4 below.
 64 Pigeonite (a stable phase in MIX1G) is modelled as clinopyroxene.

Table S6: Parameters for V partitioning and isotopic fractionation. Partition coefficients (D) from Mallmann and O'Neill (2009). β^V factors are from Wu et al. (2015). $1000\ln\beta = a/T^6 + b/T^4 + c/T^2$, where T = temperature (Kelvin).

Mineral	Valence	$D_{\min-liq}$	a	b	c
olivine	V^{3+}	0.18	1.0×10^{14}	-7.70×10^9	6.28×10^5
	V^{4+}	0.0227	4.0×10^{14}	-1.80×10^{10}	9.43×10^5
	V^{5+}	0.00255	7.0×10^{14}	-3.06×10^{10}	1.26×10^6
orthopyroxene	V^{3+}	1	1.0×10^{14}	-7.70×10^9	6.28×10^5
	V^{4+}	0.674	4.0×10^{14}	-1.80×10^{10}	9.43×10^5
	V^{5+}	0.00338	7.0×10^{14}	-3.06×10^{10}	1.26×10^6
clinopyroxene	V^{3+}	3.88	1.0×10^{14}	-7.70×10^9	6.28×10^5
	V^{4+}	1.13	4.0×10^{14}	-1.80×10^{10}	9.43×10^5
	V^{5+}	0.0142	7.0×10^{14}	-3.06×10^{10}	1.26×10^6
spinel	V^{3+}	15.3	1.0×10^{14}	-7.70×10^9	6.28×10^5
	V^{4+}	1.44	4.0×10^{14}	-1.80×10^{10}	9.43×10^5
	V^{5+}	0.00303	7.0×10^{14}	-3.06×10^{10}	1.26×10^6
garnet	V^{3+}	3.88	1.0×10^{14}	-7.70×10^9	6.28×10^5
	V^{4+}	1.13	4.0×10^{14}	-1.80×10^{10}	9.43×10^5
	V^{5+}	0.0142	7.0×10^{14}	-3.06×10^{10}	1.26×10^6
melt	V^{3+}	1	1.0×10^{14}	-7.70×10^9	6.28×10^5
	V^{4+}	1	4.0×10^{14}	-1.80×10^{10}	9.43×10^5
	V^{5+}	1	7.0×10^{14}	-3.06×10^{10}	1.26×10^6

65 The full results of the V isotope model, showing the isotopic composition of each phase calculated
 66 in KLB1, MIX1G and G2, are in Fig. S3. Note that stable isotope fractionation can only be
 67 calculated where there is liquid present, i.e., where the proportion of each V valence state can
 68 be calculated.

69 3.4 Cr

70 The proportion of $Cr^{2+/3+}$ was calculated using the parameterisation given in Berry et al. (2021)

$$\frac{Cr^{2+}}{Cr_T} = \frac{1}{1 + 10^{(0.25 \log fO_2 + \log K')}} \quad (4)$$

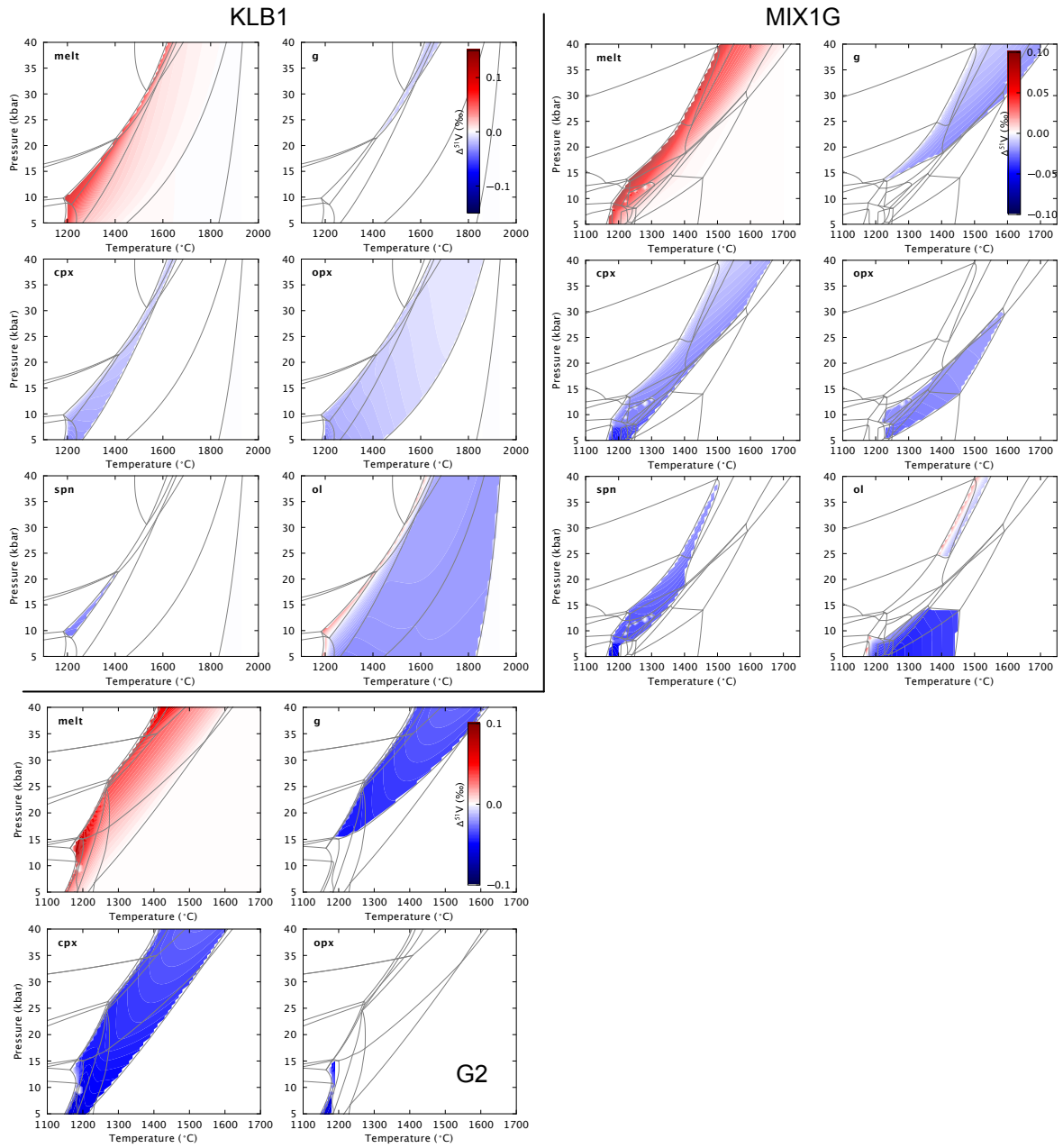


Figure S3: Full results of the V isotopic fractionation model, $\Delta^{51}\text{V}_{\text{phase-source}}$, for the three lithologies.

Table S7: Optical basicity values for each oxide, from Mills (1993).

optical basicity, Λ									
SiO ₂	TiO ₂	Al ₂ O ₃	Fe ₂ O ₃	Cr ₂ O ₃	FeO	CaO	MgO	Na ₂ O	K ₂ O
0.48	0.61	0.60	0.75	0.58	1	1	0.78	1.15	1.4

71 with fO_2 calculated at each P-T point using the fO_2 melt software (Holland et al., 2018), and
 72 the modified equilibrium constant calculated by

$$\log K' = \frac{9770}{T} - 7.69 + 6.22\Lambda + \frac{(900P - 172P^2)}{T} \quad (5)$$

73 where T is in Kelvin, P is in GPa. Λ is the optical basicity, a composition-dependent term (Duffy,
 74 1993) which reflects equilibrium between bridging oxygens, non-bridging oxygens and free oxide
 75 anions in a melt (Humphreys et al., 2015). Since optical basicity values can be assigned to
 76 individual oxides (Duffy, 1993), an ideal optical basicity for a melt can be calculated, following
 77 Mills (1993), as

$$\Lambda = \frac{\sum X_i n_i \Lambda_i}{\sum X_i n_i} \quad (6)$$

78 where X_i is the mole fraction of oxide i, n is the number of O associated with oxide i, and Λ is
 79 the theoretical Λ of oxide i (see Table S7). At each P-T point, the calculated Λ was used with
 80 equations 3 and 4 to calculate Cr redox state.

81 The partition coefficients of each valence state of Cr in the Cr-bearing mantle minerals (olivine,
 82 clinopyroxene, orthopyroxene, garnet, spinel) are given in Table S8. Partition coefficients for
 83 olivine, clinopyroxene, orthopyroxene are from (Mallmann and O'Neill, 2009); for spinel from
 84 the supplement of Shen et al. (2018), for garnet from Sutton et al. (2008). The ionic model
 85 inputs (Table S8) are based on the supplement of Shen et al. (2018); for garnet, where no ionic
 86 inputs are presented in Shen et al. (2018), we take coordination information from Deer et al.
 87 (2013), and use the relevant ionic radii as used by Shen et al. (2018). The force constants for
 88 Cr-O bonds were then calculated based on the equations presented in Sossi and O'Neill (2017)
 89 for Fe isotopes.

Table S8: Parameters for Cr partitioning and isotopic fractionation model. Partition coefficients (D) from Mallmann and O'Neill (2009); Shen et al. (2018); Sutton et al. (2008). Mineral site and bonding information is based on Shen et al. (2018). cpx = clinopyroxene, opx = orthopyroxene.

Mineral	Cr species	$D_{\text{min-melt}}$	Site	Cr coordination	Bond length (\AA)	O coordination	No. sites
spinel	2+	0	n/a	n/a	n/a	n/a	n/a
spinel	3+	220	M	6	1.995	4	2
garnet	2+	0	n/a	n/a	n/a	n/a	n/a
garnet	3+	12	M2	6	1.995	4	2
olivine	2+	0.85	M2	6	2.18	4	1
olivine	3+	0.85	M1	6	1.995	4	1
opx	2+	0.843	M2	7	2.305	3.57	1
opx	3+	3.52	M1	6	1.988	3.67	1
cpx	2+	0.587	M2	7	2.305	3.57	1
cpx	3+	12.6	M1	6	1.988	3.67	1
liquid	2+	1		4	1.913	4	
liquid	3+	1		6	1.995	4	

90 The full results of the Cr model, showing the isotopic composition of each phase calculated in
 91 KLB1, MIX1G and G2, are in Fig. S4. Note that stable isotope fractionation can only be
 92 calculated where there is liquid present, i.e., where the proportion of each Cr valence state can
 93 be calculated.

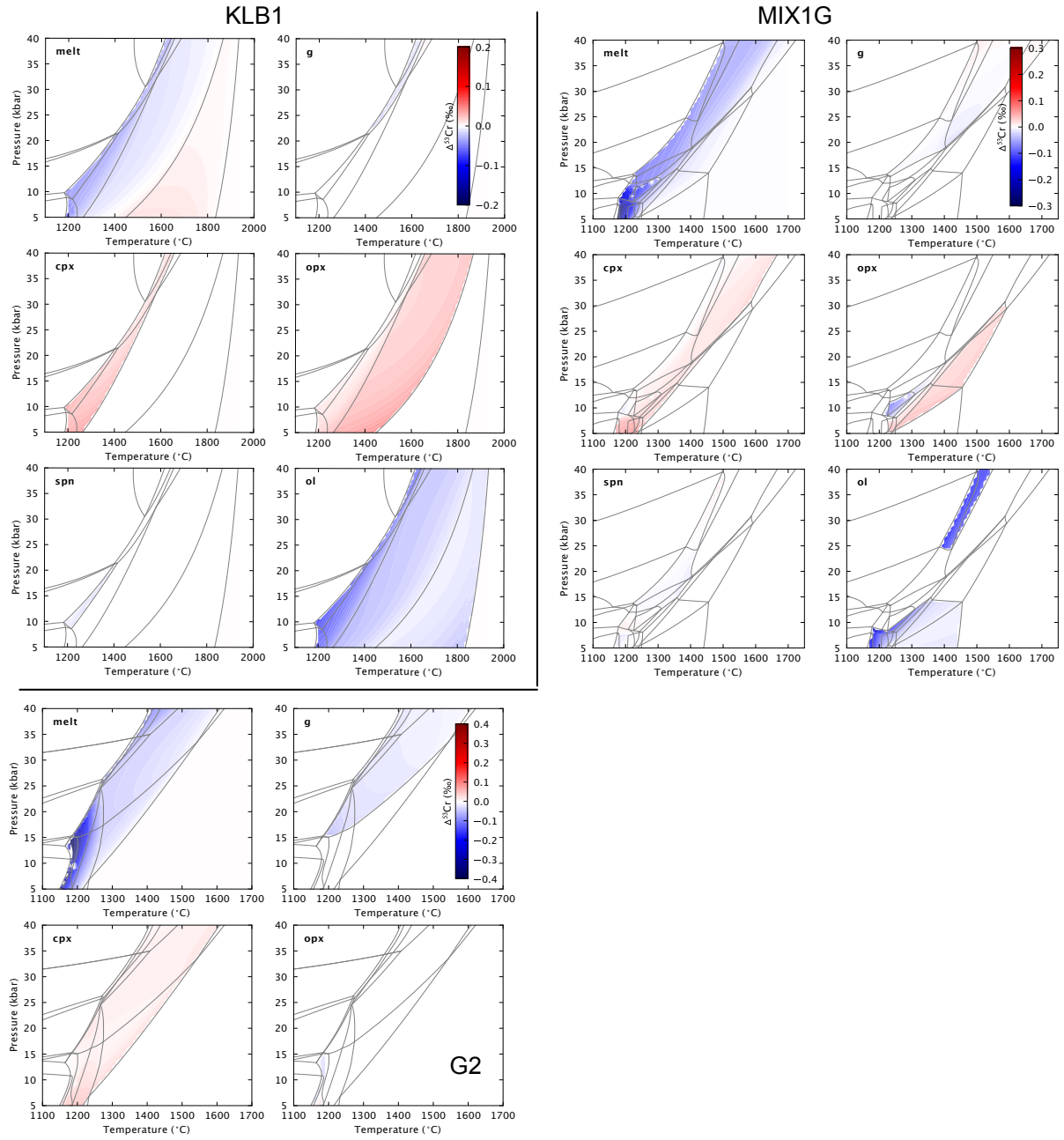


Figure S4: Full results of the Cr isotopic fractionation model, $\Delta^{53}\text{Cr}_{\text{phase-source}}$, for the three lithologies.

94 3.5 Consistency of valence state calculations

95 As discussed in the main text, the calculated $\text{Fe}^{3+}/\text{Fe}_T$, $\text{Cr}^{3+}/\text{Cr}_T$ and $\text{V}^{4+,5+}/\text{V}_T$ are not fully
96 self-consistent, as the speciation of both Cr and V are calculated based on the $\text{Fe}^{3+}/\text{Fe}_T$ at
97 any point. This means that O (oxygen) is being partitioned into the oxidised Cr and V oxides
98 without affecting Fe redox state. We consider tying the Cr and V valence states to the Fe redox
99 equilibria to be reasonable, as Fe is the dominant redox-sensitive element in the mantle with an
100 abundance considerably greater than that of Cr or V. To quantify this, we have calculated the
101 fraction of oxygen associated with oxidised valence states of the each redox-sensitive element
102 out of the total oxygen associated with the oxidised valence states, e.g.,

$$\frac{\text{O}_{\text{Fe}^{3+}}}{\text{O}_{\text{Fe}^{3+}} + \text{O}_{\text{Cr}^{3+}} + \text{O}_{\text{V}^{4+}} + \text{O}_{\text{V}^{5+}}} \quad (7)$$

103 where O_X represents the number of moles of oxygen associated with cation X at any given P-T
104 point. At $T = 1300^\circ\text{C}$ and $P = 12\text{ kbar}$, we calculate that 73% of the oxygen associated with
105 oxidised cations is associated with Fe^{3+} , 27% with Cr^{3+} , and $< 0.01\%$ with $\text{V}^{4+,5+}$, for total
106 Fe, Cr and V contents in peridotite (Davis et al., 2009; Salters and Stracke, 2004).

107 However, although our use of the Fe redox equilibria to calculate Cr and V speciation is a
108 reasonable approximation, it is important to note that considering the combined Fe-Cr-V redox
109 equilibria as in real systems would reduce the variability in the redox state of each element across
110 P-T space compared to our model. This is because the combined effect of considering three redox
111 equilibria able to accommodate oxygen via their oxidised cations together is that each system
112 will have to shift its redox equilibria less compared to a scenario where only one redox-sensitive
113 element is considered. To approximately quantify the effect of this redox feedback, we took an
114 extreme case of reducing all the Cr_2O_3 in the melt at a given P-T point, accommodating the
115 released oxygen by increasing the amount of Fe_2O_3 and reducing the amount of FeO accordingly.
116 We then calculated the $f\text{O}_2$ of this new composition through the $f\text{O}_2\text{melt}$ software (Holland
117 et al., 2018), and recalculated the expected $\text{Cr}^{3+}/\text{Cr}_T$ given the new $f\text{O}_2$. We find that, for this
118 calculation performed using at $T = 1350^\circ\text{C}$ and $P = 12\text{ kbar}$, the $f\text{O}_2$ change produces a 1%
119 change in $\text{Cr}^{3+}/\text{Cr}_T$ and negligible impact on the calculated $\delta^{53}\text{Cr}$.

120 4 Alternative input parameters and results

121 The inputs into the isotope fractionation models are better constrained for some of the stable
122 isotope systems considered relative to others. These uncertainties limit our ability to predict the
123 behaviour of some isotope systems in high temperature settings, and therefore limit the utility
124 of stable isotopes as a tool for exploring global basalt systematics. Here we discuss the results
125 of using alternative inputs.

126 4.1 Mg: $\alpha_{\text{orthopyroxene-melt}} = 1$

127 While the bonding of Mg in crystal lattices can be used to calculate a self-consistent set of β^{Mg}
128 factors (Schauble, 2011; Huang et al., 2013) to input into the melting model, β^{Mg} factors do not
129 exist for silicate liquid. However, Stracke et al. (2018) show that $\alpha_{\text{olivine, orthopyroxene-melt}} \sim 1$
130 based on equilibrated peridotite xenoliths. Mg is also similarly coordinated ([5–6] fold) in melt
131 and olivine/orthopyroxene (George and Stebbins, 1998; Shimoda et al., 2007; Stracke et al.,

132 2018), providing further evidence for minimal olivine/orthopyroxene – melt Mg stable isotope
 133 fractionation at magmatic temperatures.

134 We have chosen to use $\alpha_{\text{olivine-melt}} = 1$; i.e., using the same β^{Mg} values for liquid as for olivine.
 135 The alternative is to use $\alpha_{\text{orthopyroxene-melt}} = 1$. Fig. S5 compares the $\Delta^{26}\text{Mg}_{\text{melt-source}}$ for an
 136 input of $\alpha_{\text{olivine-melt}} = 1$ (left hand panels; these are the results used in the main text) with an
 137 input of $\alpha_{\text{orthopyroxene-melt}} = 1$ (right hand panels). The new input parameters replace the melt
 138 in Table S4 with Table S9.

Table S9: Alternate parameters for Mg isotope β^{Mg} factors from Huang et al. (2013). $1000\ln\beta = a/T^6 + b/T^4 + c/T^2$, where T = temperature (Kelvin).

Mineral	a	b	c
melt	-8.10×10^{15}	6.65×10^{10}	2.25×10^6

139 For KLB1 peridotite, the magnitude of the difference in melt isotope composition between the
 140 two models is $\approx 0.05\text{‰}$ (small compared to the average analytical uncertainty of 0.1‰ suggested
 141 in Stracke et al., 2018, and similar to the long-term 2 S.D. of 0.06‰ given in Wang et al., 2021),
 142 however the direction of partial melting fractionation is different. In the case of $\alpha_{\text{olivine-melt}} = 1$,
 143 KLB1 melts are isotopically lighter than the bulk, and get heavier with increasing degree of
 144 melting (the latter part of this behaviour is consistent with fractional melting modelled by
 145 Stracke et al., 2018). For $\alpha_{\text{orthopyroxene-melt}} = 1$, partial melts are isotopically heavier than the
 146 bulk, and get progressively lighter with increasing degree of melting (the former part of this
 147 behaviour is consistent with Stracke et al., 2018).

148 For KLB1, the similar differences in modelled melting isotopic fractionation relative to the ana-
 149 lytical uncertainty means that the choice of $\alpha_{\text{olivine-melt}} = 1$ or $\alpha_{\text{orthopyroxene-melt}} = 1$ for liquid
 150 does not greatly alter any conclusions drawn from the peridotite melting. Minimal difference
 151 between the models is also seen for MIX1G. However, low pressure melts from G2 ($< 15\text{ kbar}$)
 152 show effectively no stable isotope fractionation when $\alpha_{\text{orthopyroxene-melt}} = 1$ is used, compared
 153 to fractionations as large as $\Delta^{26}\text{Mg} = -0.1\text{‰}$ with $\alpha_{\text{olivine-melt}} = 1$. This is because, for these
 154 shallow depths, the only stable Mg-bearing phases in G2 are clinopyroxene and liquid (garnet
 155 is not stable). Orthopyroxene and clinopyroxene have similar β^{Mg} factors and therefore, when
 156 liquid is given the same β^{Mg} factor as orthopyroxene, the coexisting liquid and clinopyroxene
 157 are only minimally isotopically fractionated from each other. These results highlight how impor-
 158 tant accurate experimental $\alpha_{\text{mineral-melt}}$ are for understanding the behaviour of stable isotope
 159 systems, and the important role analytical precision plays in being able to identify these effects
 160 in nature. Magnesium stable isotopes would benefit from further work in both of these areas.

161 4.2 Ca

162 The β^{Ca} value approach reduces uncertainties associated with an ionic model, as the bond lengths
 163 and Ca, O coordination numbers in the relevant mineral structures required for the ionic model
 164 are variable in the literature (e.g., Feng et al., 2014; Huang et al., 2019). Huang et al. (2019)
 165 present an alternative set of β^{Ca} values for the Ca-bearing minerals we are concerned with,
 166 although Antonelli et al. (2019) argue that the PBE functionals used in their own study better
 167 represent Ca isotope fractionation between species than the less accurate LDA functionals used
 168 by Huang et al. (2019). Additionally, Antonelli et al. (2019) calculate β^{Ca} for pyrope garnet,
 169 which is more representative of mantle garnet than the grossular endmember used in Huang
 170 et al. (2019). Our model uses $\beta_{\text{plagioclase}}^{\text{Ca}} < \beta_{\text{melt}}^{\text{Ca}} < \beta_{\text{clinopyroxene}}^{\text{Ca}}$ (see explanation in Section

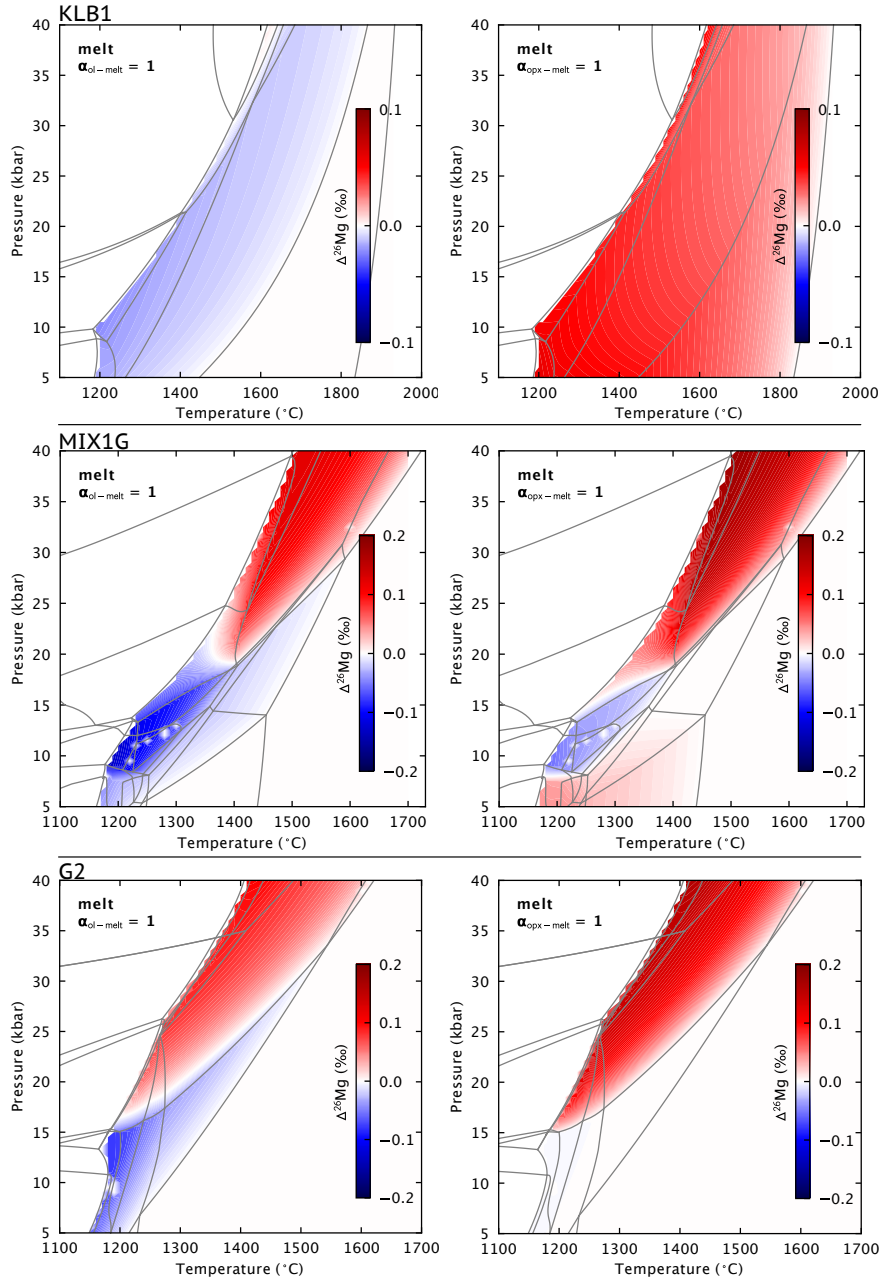


Figure S5: $\Delta^{26}\text{Mg}_{\text{melt-source}}$ compared for a model where $\alpha_{\text{olivine-melt}} = 1$ (left hand plots) and $\alpha_{\text{orthopyroxene-melt}} = 1$ (right hand plots) for all three lithologies considered.

171 3.2), but some studies do suggest that $\beta_{\text{melt}}^{\text{Ca}} = \beta_{\text{clinopyroxene}}^{\text{Ca}}$ (e.g., Chen et al., 2019) based on
 172 observations from natural equilibrated peridotite and pyroxenite samples that there is minimal
 173 Ca isotope fractionation between clinopyroxene and basaltic melt. Independent estimates of the
 174 force constants in silicate liquids and further experimental constraints on the behaviour of Ca
 175 isotopes in silicate melts are therefore needed to improve the Ca isotope fractionation model.

176 Fig. S6 shows the results of using the β^{Ca} values from Huang et al. (2019) with $\beta_{\text{melt}}^{\text{Ca}} =$
 177 $\beta_{\text{clinopyroxene}}^{\text{Ca}}$ (equivalent to $\alpha_{\text{clinopyroxene-melt}} = 1$). The use of isotopically lighter grossular
 178 garnet and isotopically heavier melt (relative to the Antonelli et al. (2019) inputs) result in
 179 overall less negative melt-source isotope fractionation than in the main model, making all Ca

180 isotope fractionations harder to resolve with current analytical precision. However, none of the
 181 conclusions presented in the main text are altered by these alternative parameter choices.

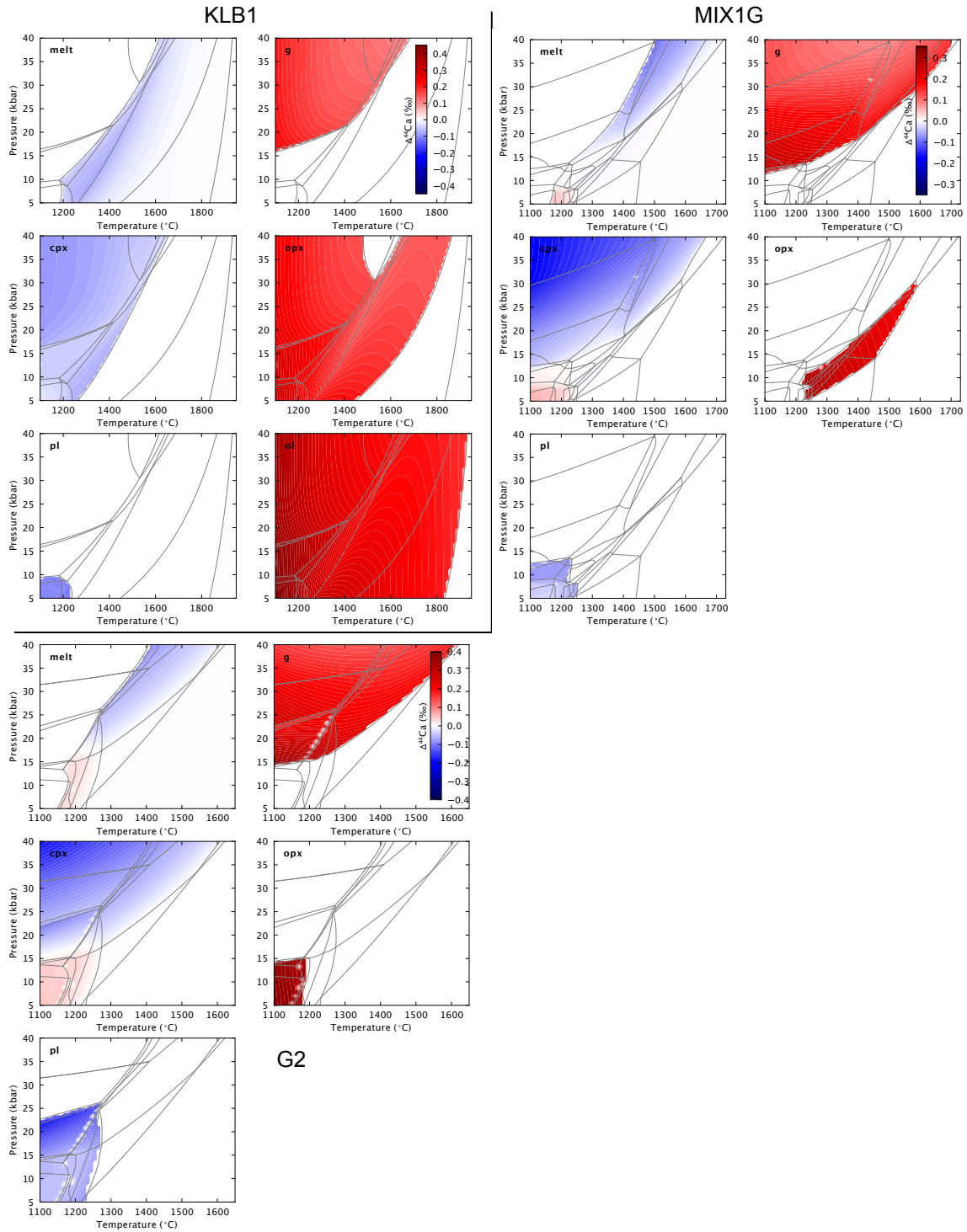


Figure S6: Model results of $\Delta^{44}\text{Ca}_{\text{phase-source}}$ for β^{Ca} inputs from Huang et al. (2019), and using $\beta_{\text{melt}}^{\text{Ca}} = \beta_{\text{clinopyroxene}}^{\text{Ca}}$. Generally, the calculated $\Delta^{44}\text{Ca}_{\text{phase-source}}$ is less negative than using the β^{Ca} inputs from Antonelli et al. (2019) as shown in Fig. S2.

182 The composition of the phases involved further complicates models of Ca isotope fractiona-
 183 tion. Feng et al. (2014) and Wang et al. (2017) show that Ca-O bond strength, hence isotopic
 184 fractionation, in orthopyroxene is dependent on the Ca and Fe content of the pyroxene, al-

185 though the concentrations of these elements in natural orthopyroxenes may produce negligible
186 concentration effects on isotopic fractionation (Wang et al., 2017). The isotopic fractionation
187 associated with garnet, in particular, is sensitive to Ca content for typical mantle compositions,
188 with a garnet-composition-sensitive fractionation factor produces smaller melting fractionations
189 in eclogite than using a fixed composition model (Chen et al., 2020a). Our model has not
190 implemented composition-specific fractionation factors, although does use the pyrope garnet
191 endmember which is most representative of mantle garnet. We also note that Antonelli et al.
192 (2019) calculate compositional effects for olivine and pyroxene solid solutions that are smaller
193 than those predicted using LDA functionals (e.g., Feng et al., 2014; Wang et al., 2017).

194 4.3 Fe

195 The model inputs for Fe isotope fractionation are relatively well constrained. The bonding
196 environment of Fe in mineral phases is well studied, making an ionic model approach possible
197 (e.g., Sossi and O'Neill, 2017). The model based on ionic inputs produces isotopic fractionation
198 and force constant predictions consistent with predictions by other methods and observations
199 (e.g., Dauphas et al., 2014; Macris et al., 2015). In addition, the Fe-O force constants in
200 basaltic glass as a function of Fe³⁺ content have been experimentally determined (Dauphas
201 et al., 2014), allowing the melt isotopic fractionation to be independently parameterised into
202 the model (although we note the potential complication induced by quenching effects, and by
203 melt generation in the mantle at non-atmospheric pressure). As discussed by Sossi and O'Neill
204 (2017), some terms in the ionic model approach (such as the ionicity of the bonds) are less well
205 constrained than the bonding environment of Fe in mineral structures, however the resultant
206 fractionation factors are minimally affected by these uncertainties.

207 4.4 V: V_{melt}⁵⁺ = [4] coordinated

208 It has been suggested that vanadium isotopes may track magmatic redox (Prytulak et al., 2013,
209 2017; Sossi et al., 2018b; Wu et al., 2018) due to the existence of three valence states (3+, 4+,
210 5+) in terrestrial magmatic systems (Canil, 1999), though a better understanding of V isotope
211 behaviour during magmatic processes is required to exploit the system's full potential as a redox
212 tracer (Prytulak et al., 2017). The successful modelling of V isotope fractionation therefore relies
213 on an understanding of the behaviour of each valence state of V.

214 Although slightly variable V partition coefficients for each valence state into mineral phases
215 can be found from experimental studies (Mallmann and O'Neill, 2009), these variations make
216 no significant difference to the calculated melt isotopic compositions. The main uncertainty
217 in the V fractionation model inputs is the coordination of V⁴⁺ and V⁵⁺ in silicate melts (and
218 associated β^V factors). V⁴⁺ in silicate liquid is [5–6] coordinated (Sutton et al., 2005), and
219 V⁴⁺ coordination of both [5.33] and [6] have been used in existing studies discussing V isotope
220 behaviour during melting (Qi et al., 2019; Wu et al., 2018 respectively). V⁵⁺ coordination in
221 silicate melts is [4–6] (Sutton et al., 2005; Righter et al., 2006), with Qi et al. (2019) quoting
222 an average of [5] coordination. As one of the main causes of V isotope fractionation during
223 melting is predicted to be the relative compatibility of V⁵⁺ (and to a lesser extent V⁴⁺) in
224 melts relative to minerals (e.g., Canil, 2002; Mallmann and O'Neill, 2009; Wu et al., 2018), the
225 bonding environment and force constants of V species in the melt is a particularly important
226 input. The choice of coordination number for the V species controls the $\ln \beta$ value chosen for
227 that species: β^V values are taken from those calculated for different bonding environments in

228 aqueous solutions (Wu et al., 2015), assumed to be a good approximation to the case of silicate
 229 melts following Wu et al. (2018). [5] coordination of V^{5+} is not modelled in Wu et al. (2015),
 230 and therefore only the cases of V^{5+} in [4] and [6] coordination can be used, for which β^V values
 231 have been calculated. No calculations are available for V^{4+} in [5] coordination, and therefore
 232 the model only considers V^{4+} in [6] coordination.

233 Fig. S7 compares the $\Delta^{51}V_{\text{melt-source}}$ for an input of $V_{\text{melt}}^{5+} = [6]$ -fold coordination (left hand
 234 panels; these are the results used in the main text) with an input of $V_{\text{melt}}^{5+} = [4]$ -fold coordination
 235 (right hand panels). There are two aqueous complexes with 4-fold V^{5+} in Wu et al. (2015). Here
 236 we have chosen the β^V factors for the $[V^{5+}O_2(OH)_2]^-$ (Table S10) complex since this complex
 237 has the greatest $\ln\beta^V$ difference from the [6]-fold coordinated case, therefore covering the greatest
 238 range of input uncertainty.

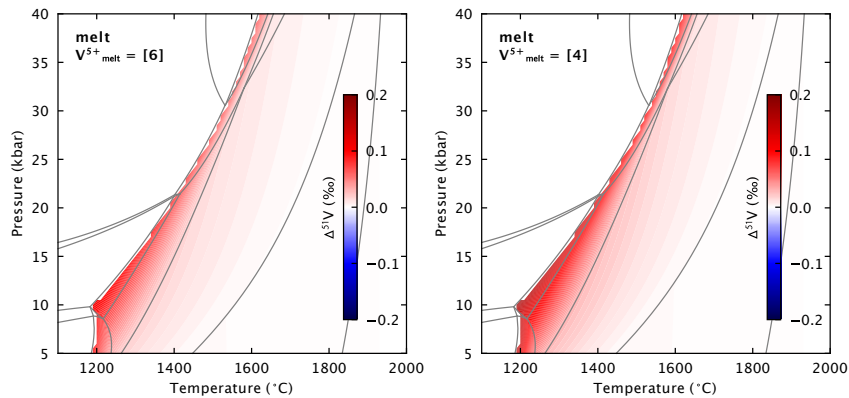


Figure S7: $\Delta^{51}V_{\text{melt-source}}$ compared for a model where $V_{\text{melt}}^{5+} = [6]$ coordinated (left hand plots) and $V_{\text{melt}}^{5+} = [4]$ coordinated (right hand plots) for KLB1.

Table S10: Alternate parameters for V isotope β^V factors from Wu et al. (2015). $1000\ln\beta = a/T^6 + b/T^4 + c/T^2$, where T = temperature (Kelvin).

Mineral	Valence	a	b	c
melt	V^{5+}	7.00×10^{14}	-3.25×10^{10}	1.42×10^6

239 The model discussed in the main text is the result of using the upper limit of [6] coordinated
 240 V^{5+} : this will provide the minimum melting fractionation, since a higher coordination number
 241 results in weaker V-O bonds, and less extreme isotopic fractionation. Using an input of [4]
 242 coordinated V^{5+} results in isotopically heavier melts, with the maximum KLB1 melting stable
 243 isotope fractionation increasing from ≈ 0.1 ‰ in the [6] coordinated case to ≈ 0.15 ‰ close to
 244 the solidus around 10 kbar (Fig. S6 for full results) in the [4] coordinated case. However, the
 245 difference between the two models decreases at higher temperature, and for the temperature
 246 range of interest here (potential temperatures from 1300 – 1530 °C), at 5 kbar the difference
 247 in melt isotopic composition between the two models is ≈ 0.01 ‰, and therefore negligible.
 248 However, to improve understanding of high temperature V behaviour and link natural data to
 249 specific processes, experimental and/or theoretical work on equilibrium isotopic fractionation
 250 factors for silicate melt and minerals for different V species is needed.

251 4.5 Cr: $\text{Cr}_{\text{melt}}^{2+} = [6]$ coordinated

252 As chromium is a first-group transition element, the properties controlling isotopic fractionation
 253 behaviour (e.g., valence, ionic radius hence bonding environment) of Cr^{2+} and Cr^{3+} ions may
 254 be expected to be similar to Fe^{2+} and Fe^{3+} ions (Xia et al., 2017). The main Cr-bearing
 255 phases in the mantle are spinel, garnet and pyroxenes, although some Cr will partition into
 256 olivine (Mallmann and O’Neill, 2009), and as with vanadium, partition coefficients depend on
 257 the valence state of Cr (Mallmann and O’Neill, 2009).

258 As a set of β^{Cr} factors for mantle minerals does not exist, our models of Cr stable isotope
 259 fractionation use force constants in each mineral calculated from an ionic bonding model (Shen
 260 et al., 2018), following the approach used for Fe isotopes. An ionic model has been shown to
 261 produce inter-mineral fractionation in agreement with the magnitude of isotope fractionation
 262 inferred from studies of lherzolite mineral pairs (Shen et al., 2018). As with V isotopes, the
 263 main uncertainty relating to species coordination is in liquid. Cr^{2+} in silicate melt is reported
 264 to be in [4] coordination (Miletich et al., 1999; O’Neill and Berry, 2006; Shen et al., 2020), but
 265 may also be in [6] coordination in a square planar arrangement (O’Neill and Berry, 2006). The
 266 isotope fractionation model presented here calculates force constants using effective anionic and
 267 cationic radii based on Shannon (1976), following the approach outlined in Shen et al. (2018),
 268 and these radii depend negatively on the coordination number chosen (Gibbs et al., 2014).

269 Fig. S8 compares the $\Delta^{53}\text{Cr}_{\text{melt-source}}$ for an input of $\text{Cr}_{\text{melt}}^{2+} = [4]$ coordinated (left hand panels;
 270 these are the results used in the main text) with an input of $\text{Cr}_{\text{melt}}^{2+} = [6]$ coordinated (right
 271 hand panels). The input parameters for the [6] coordinated case are in Table S8.

Table S11: Alternate parameters for Cr isotopic fractionation model. Mineral site and bonding information is based on Shen et al. (2018).

Mineral	Cr species	Site	Cr coordination	Bond length (Å)	O coordination
melt	2+	n/a	6	2.18	4

272 The model presented in the main text uses an input of Cr^{2+} in [4] coordination in the melt.
 273 Inputting [6] coordinated Cr^{2+} in the melt results in isotopically lighter peridotite melts across
 274 all of P-T space (Fig. S8) than the [4] coordinated case as a consequence of weaker Cr-O bonds
 275 in a [6] coordinated environment than [4] coordinated. A similar effect is seen in MIX1G and
 276 G2 melts. The melt-source isotopic fractionation is always negative, consistent with modelling
 277 by Shen et al. (2018), and melts get progressively isotopically heavier (towards 0‰ fractiona-
 278 tion from the source) with increasing degree of melting, consistent with observations from OIB
 279 (Bonnand et al., 2020). The cationic radii are likely to be overestimates, as the radii for [6]
 280 coordinated Cr used in Shen et al. (2018) assume octahedral, not square planar, coordination.
 281 Square planar bonds would be expected to be longer than octahedral, and this would result
 282 in even isotopically lighter melts. Therefore, although the model uses the most appropriate
 283 available inputs for the more commonly quoted [4] coordination of Cr^{2+} in melts, our under-
 284 standing of the behaviour of Cr isotopes in silicate melts would benefit from further work on
 285 the force constants of Cr^{2+} in glasses, or the inputs required to calculate these such as bonding
 286 environment or effective Cr-O bond lengths.

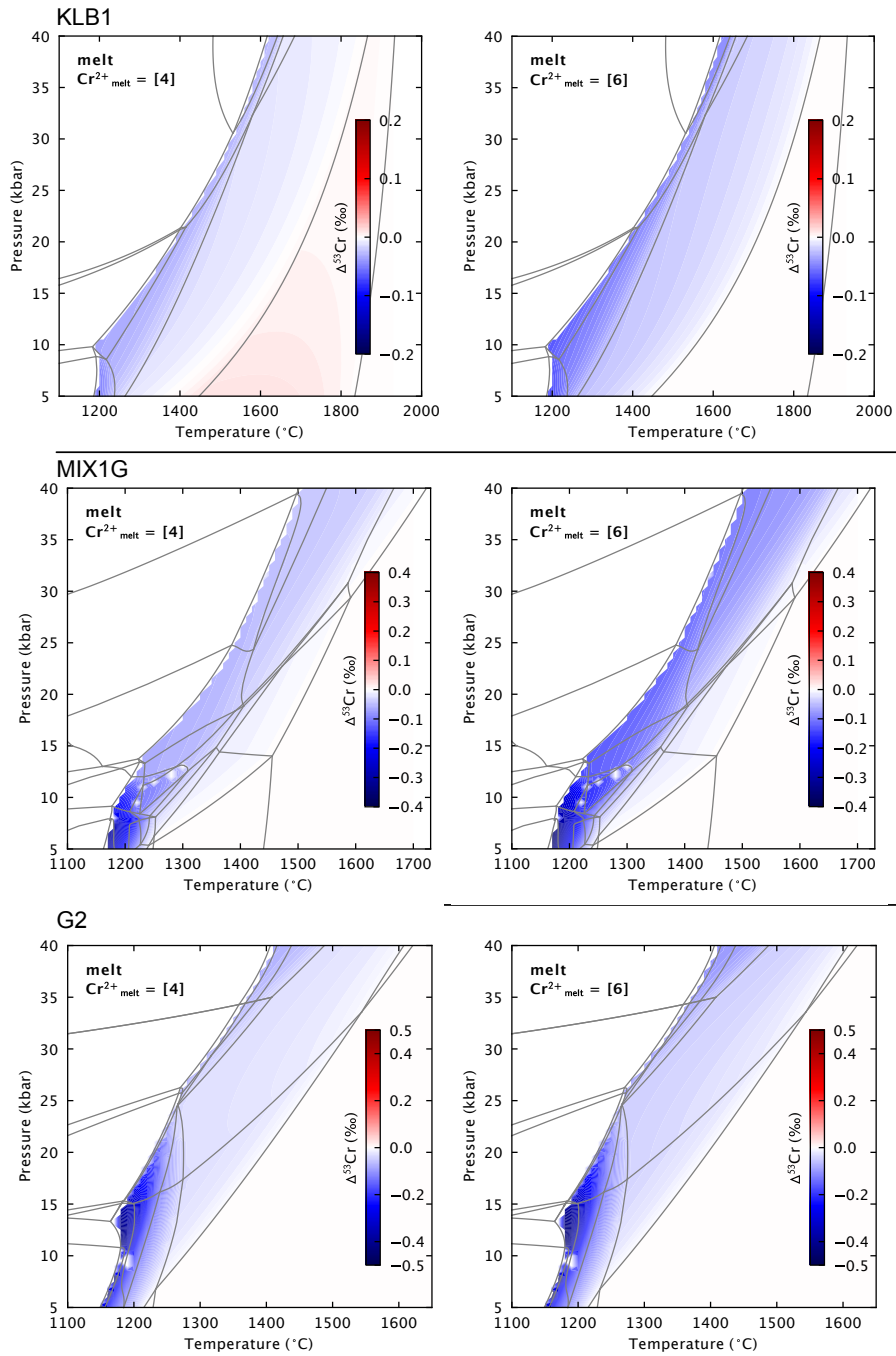


Figure S8: $\Delta^{53}\text{Cr}_{\text{melt-source}}$ compared for a model where $\text{Cr}^{2+}_{\text{melt}} = [4]$ coordinated (left hand plots) and $\text{Cr}^{2+}_{\text{melt}} = [6]$ coordinated (right hand plots) for all three lithologies.

287 5 Other potential isotopes

288 Other stable isotope systems not considered here might also have potential for investigating
 289 mantle heterogeneity (in either lithology or temperature). However, we require a better under-
 290 standing of the bonding environment in melts or relevant minerals before these isotope systems
 291 can be implemented into the stable isotope fractionation models used here.

292 Like other stable isotope systems, the equilibrium $\delta^{30}\text{Si}$ value of a phase will depend on the

293 strength of the Si-O bonds. A more polymerised (or, more Si-rich) phase will have more Si-O
 294 bonds, therefore a stronger average Si-O bond force constant, hence high $\delta^{30}\text{Si}$ (Grant, 1954).
 295 However, density functional theory (DFT) calculations show this simple prediction is compli-
 296 cated by the role of other network-modifying cations in silicates (Méheut et al., 2009; Méheut
 297 and Schauble, 2014). Silicon isotope fractionation factors have been calculated using DFT for
 298 appropriate mantle minerals (Huang et al., 2014), however there is a lack of work on the crucial
 299 mineral-melt fractionation information required for the equilibrium fractionation model here.
 300 Furthermore, it has been suggested that Si isotopic fractionation in liquids may be correlated
 301 with chemical composition (Qin et al., 2016), an effect that is too complicated (and under-
 302 constrained) to model here. This knowledge gap could provide scope for future work, however
 303 the mantle is currently thought to be homogeneous in Si isotopes (within current limits of an-
 304 alytical precision), and there seems to be only a small partial melting effect (Savage et al.,
 305 2014).

306 Titanium stable isotopes have been suggested as tracer of magmatic processes (Millet et al.,
 307 2016), since Ti can exist in multiple coordination states in silicate melts and Ti-oxides (Farges
 308 and Brown Jr, 1997). However, Ti stable isotope fractionation in magmatic systems is dominated
 309 by the behaviour of Fe-Ti oxides, with no partial melting fractionation from mantle lithologies
 310 expected unless under conditions of Ti-oxide saturation. DFT calculations predict no significant
 311 Ti isotope fractionation between pyroxenes, olivine or garnet (Wang et al., 2020), and both
 312 eclogites and MORB have the same measured $\delta^{49}\text{Ti}$ value (Millet et al., 2016). Therefore, Ti
 313 isotopes are not discussed further here, as there is no existing data to support for their use in
 314 tracing temperature or lithological heterogeneity in the mantle associated with MORB or OIB.

315 6 Average depth of melting

316 The mean melt pressure, \bar{P} , for isentropic melting at any T_p is calculated by the following
 317 integral, based on Asimow et al. (2001) with an additional geometric weighting factor to account
 318 for the triangular shape of the melting region

$$\bar{P} = \frac{\int_{P_0}^{P_f} P \Delta F g \, dP}{\int_{P_0}^{P_f} \Delta F g \, dP} \quad (8)$$

319 where P_0 and P_f are the initial and final pressures of melting respectively, ΔF is the amount
 320 of melt produced at each pressure step (dP) along the isentropic melting path, and g is the
 321 geometric weighting factor. P_0 is taken as the intersection of the isentrope with the solidus,
 322 P_f is used as 5 kbar in this model as melt compositions have not been calculated for shallower
 323 pressures - this is an overestimate (i.e., final melting will be shallower) for $T_p = 1300^\circ\text{C}$ and an
 324 underestimate for $T_p = 1530^\circ\text{C}$ (Asimow et al., 2001), but will provide the correct behaviour.
 325 Pressure steps used are the 0.5 kbar intervals used in the P-T grid. ΔF is calculated for each
 326 step from the difference in melt fractions calculated by THERMOCALC for consecutive dP steps.
 327 The geometric weighting factor is a linear scaling running from 1 at the base of the melting
 328 region to 0 at the top, to account for the wider base of the triangular melting region beneath
 329 ridges. We calculate the average melting pressure for the three T_p s for which we have calculated
 330 isentropes in KLB1, and lineally interpolate between them (this will not be a linear relationship
 331 in reality, e.g., Asimow et al., 2001, but our approach will provide a good approximation) to
 332 get average melt pressures for temperatures between $T_p = 1300^\circ\text{C}$ and 1530°C . For KLB1, we
 333 calculate an average melting pressure of approximately 12 kbar at $T_p = 1300^\circ\text{C}$, and 28 kbar
 334 at $T_p = 1530^\circ\text{C}$. The average $\Delta_{\text{melt-source}}$ for each T_p is then taken as the $\Delta_{\text{melt-source}}$ at the

335 calculated average melting pressure.

336 The approach above, however, uses an equilibrium melt isotope composition at a single pressure
 337 to approximate the fractional, multibaric melting thought to occur in the mantle (McKenzie,
 338 1984; von Bagen and Waff, 1986; Langmuir et al., 1992; Asimow et al., 2001). To assess the
 339 suitability of the equilibrium melting at single pressure approach, we have also calculated the
 340 average melt isotopic composition by considering the stable isotope composition and elemental
 341 concentration for each new step of melt generated during isentropic decompression at $T_p =$
 342 1300°C and 1530°C by

$$\bar{\Delta}_{\text{melt-source}} = \frac{\int_{P_0}^{P_f} \Delta_x c \Delta F g dP}{\int_{P_0}^{P_f} c \Delta F g dP} \quad (9)$$

343 where Δ_x is the $\Delta_{\text{melt-source}}$ of the melt produced in the pressure step, and c is the wt % of
 344 the element being considered in the melt, and other terms as above. The approach of equation
 345 7, while considering polybaric melting, is still not a fractional melting model, but fractional
 346 melting is beyond the scope of this work. Nonetheless, the chemical composition of pooled
 347 fractional melts from a ridge melting regime has been shown to closely resemble that of a simple
 348 equilibrium batch melt (Langmuir et al., 1992). Equation 7 therefore allows us to broadly assess
 349 whether the approach taken in equation 6 produces estimates of mean melt isotope ratios very
 350 different from a more realistic polybaric approach, although it should be noted that the polybaric
 351 method of equation 7 only allows us to calculate the average melt isotopic composition for the
 352 two T_p s for which we have calculated isentropes.

353 Fig. S9 shows the comparison of the two calculations of average $\Delta_{\text{melt-source}}$. The dots show the
 354 average melt isotope composition when calculated as a sum along the isentrope; the solid lines
 355 show the approximation using the melt composition at the average pressure of melting. The
 356 calculation using the average melt pressure is a good approximation of the calculated average
 357 melt isotope compositions. Therefore we use this pressure approach in the main text: the
 358 limitation of the calculated average melt isotope composition is that we only have information for
 359 the T_p s for which we have calculated isentropes, hence use a linear interpolation between them,
 360 whereas the average pressure approach allows us to use more information from the calculated
 361 P-T map of melt compositions, reflecting the changing of stable mineral assemblage, for example.

362 7 Trace element calculations

363 Table S12 gives the mineral-melt partition coefficients and mantle concentrations for Sm and
 364 Yb used in this model.

Table S12: Input parameters (partition coefficients relative to melt, mantle concentrations) for calculating the Sm, Dy and Yb concentration of trace elements in the model melts. $D_{\text{gt,cpx,ol,opx}}$ and element concentrations from Gibson and Geist (2010); D_{spinel} from Liu et al. (2019) (average values from high P experiments); D_{plag} from Aigner-Torres et al. (2007) (from ‘run 31’); D_{rutile} from Klemme et al. (2005) (Sm), Foley et al. (2000) (Dy, Yb); D_{ilmenite} from Zack and Brumm (1998). Conc.DM = concentration in depleted mantle; mineral abbreviations as previous with rut. = rutile, ilm. = ilmenite.

Element	D_{gt}	D_{cpx}	D_{ol}	D_{opx}	D_{spn}	D_{plag}	D_{rut}	D_{ilm}	Conc.DM
Sm	0.25	0.293	0.0011	0.02	0.198	0.081	0.00135	0.00059	0.299
Dy	2.2	0.4	0.0027	0.011	0.069	0.049	0.00076	0.01	0.525
Yb	6.6	0.4	2	0.08	0.812	0.025	0.0093	0.17	0.347

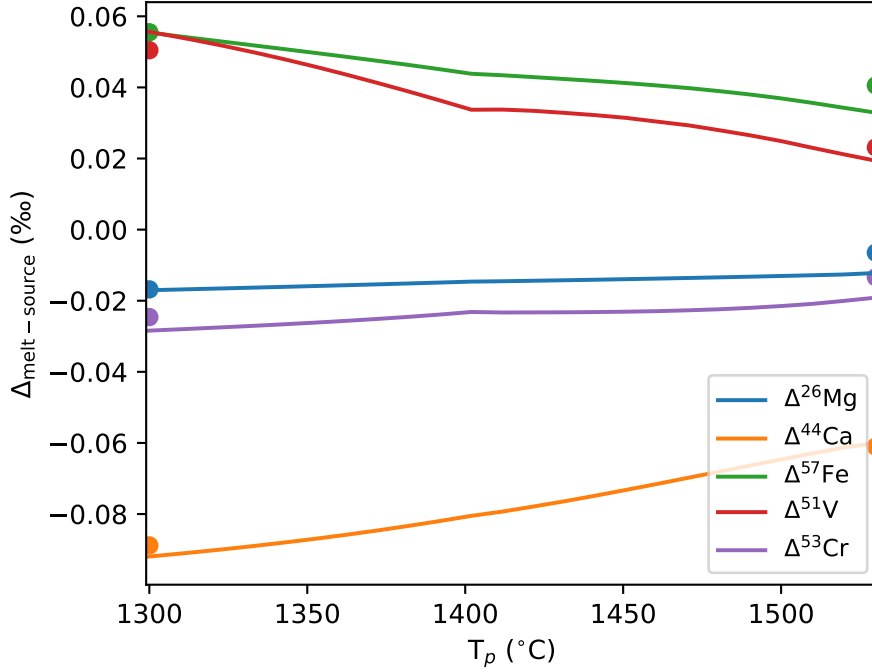


Figure S9: $\Delta_{\text{melt-source}}$ for KLB1 calculated by taking the isotopic fractionation at the average melt pressure (solid lines) compared to a calculation of the average melt composition by considering polybaric melts (equation 7), shown as coloured dots.

365 The concentration of each element in the melt was calculated at each P-T point by:

$$C_{\text{liquid}} = \frac{C_{\text{bulk}}}{\sum D_{\text{mineral}} F_{\text{mineral}} + F_{\text{liquid}}} \quad (10)$$

366 where C = concentration, D = partition coefficient, F = fraction of phase.

367 8 Additional reference figures

368 Fig. S10 provides a reference for how oxygen fugacity, melt fraction, and selected trace element
 369 ratios vary along the isentropes shown in Figs. 1 and 4 of the main text. The trace element
 370 concentrations are calculated using depleted mantle (DM) inputs for KLB1, and plume inputs
 371 for the enriched MIX1G and G2 lithologies (see Table S12).

372 We choose to use peridotite isentropic decompression paths for the two pyroxenite lithologies, as
 373 discussed in the main text. Figs. S11, S12 show results for the other endmember case for the two
 374 pyroxenite lithologies, i.e., where the lithologies decompress along isentropic paths calculated
 375 for each pure pyroxenite lithology.

376 Fig. S13 compares the melt-source isotope fractionation with progressive melting (along peri-
 377 dotite isentropes) for pairs of isotopes in each lithology, allowing assessment of the ‘direct’
 378 sensitivity to lithology. By direct lithological sensitivity, we refer to sensitivity related to miner-
 379 alogy – these isotope systems may still trace lithology via each lithology’s different bulk isotope
 380 composition. On balance, given that melts from any pyroxenite lithology are likely to be diluted
 381 with ambient peridotite melts when extracted from the mantle (e.g., Hirschmann and Stolper,

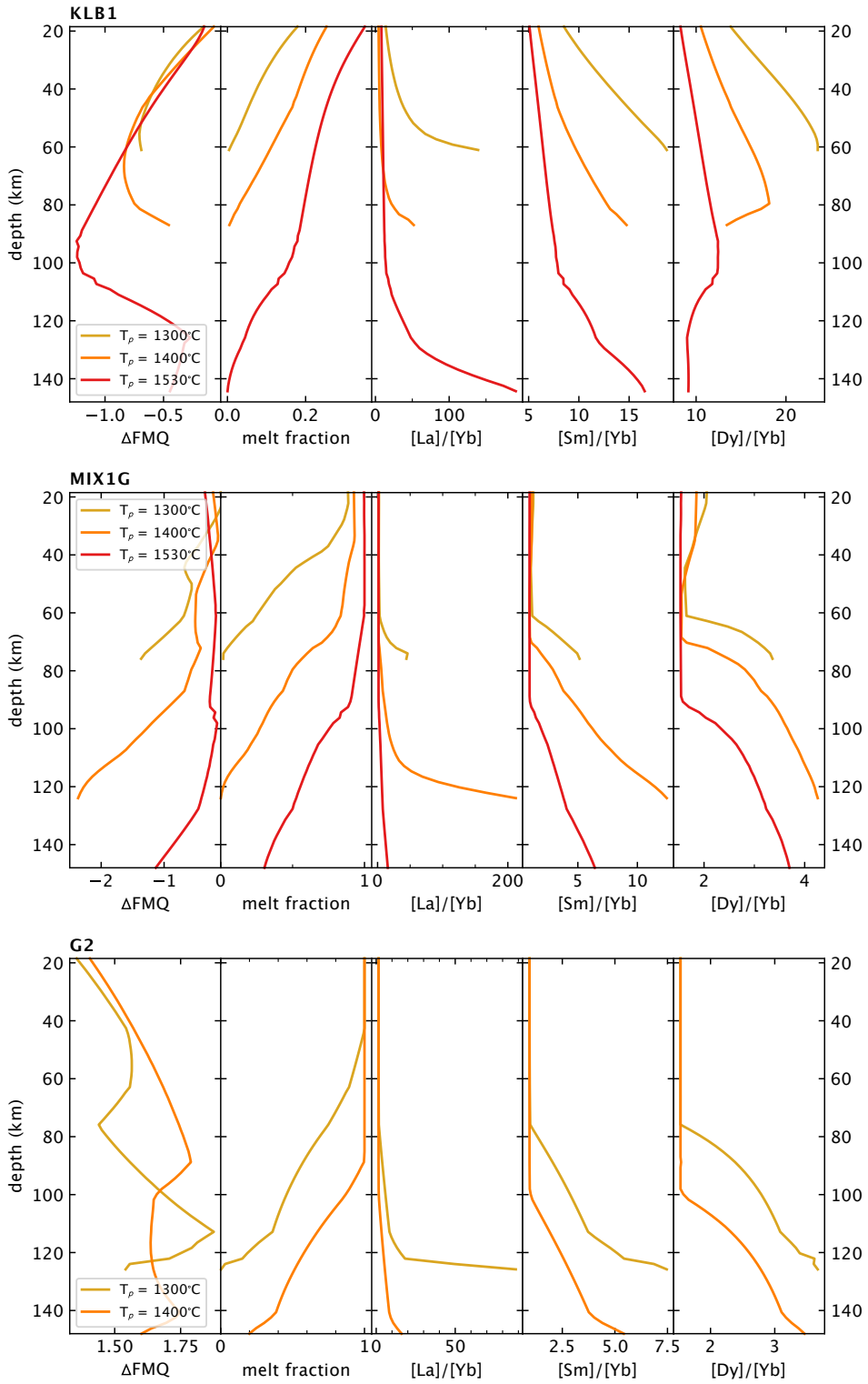


Figure S10: Parameters along the peridotite isentropic decompression paths for all three lithologies: oxygen fugacity (from the $f\text{O}_2$ melt software; Holland et al., 2018), melt fraction (from THERMOCALC output), and selected trace element ratios in the melts (calculated using the model described in the section above). The isentropes match those in Figs. 2 and 5 in the main text.

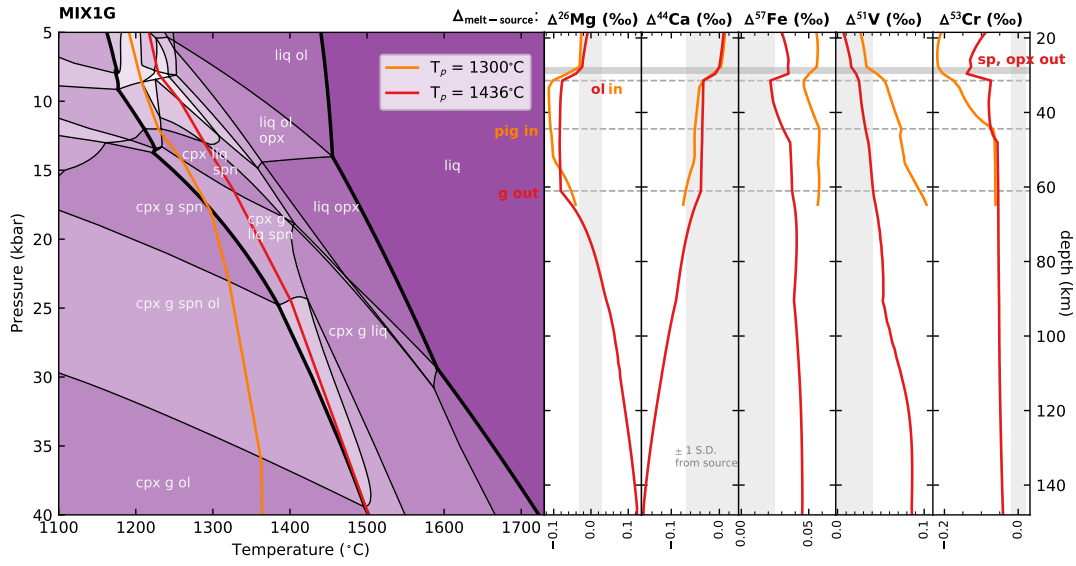


Figure S11: Isentropic melting paths for pure MIX1G, and $\Delta_{\text{melt-source}}$ along these isentropes. Errors are shown as 1 S.D. long-term uncertainty.

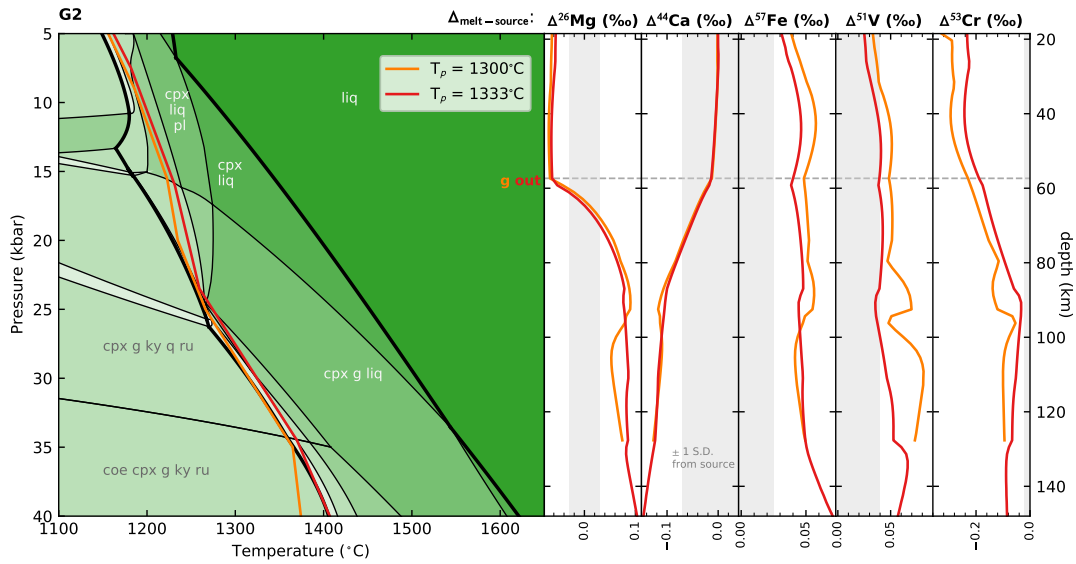


Figure S12: Isentropic melting paths for pure G2, and $\Delta_{\text{melt-source}}$ along these isentropes. Errors are shown as 1 S.D. long-term uncertainty.

1996; Sobolev et al., 2005; Shorttle and MacLennan, 2011; Soderman et al., 2021), our results in Fig. S13 are probably representative of maximum isotope fractionation signatures we might see in natural systems, and Fig. 7 in the main text (which takes into account both this dilution effect and non-identical bulk isotope compositions of the lithologies) may better represent the predicted resolvability of each isotope system to lithology. Generally, Fig. S13 suggests that the stable isotopes considered here may have poor direct sensitivity to lithological heterogeneity in the source regions of basalts, assuming equilibrium fractionation behaviour and in the absence of any pre-imposed source heterogeneity (i.e., where all lithologies have the same bulk isotope ratio).

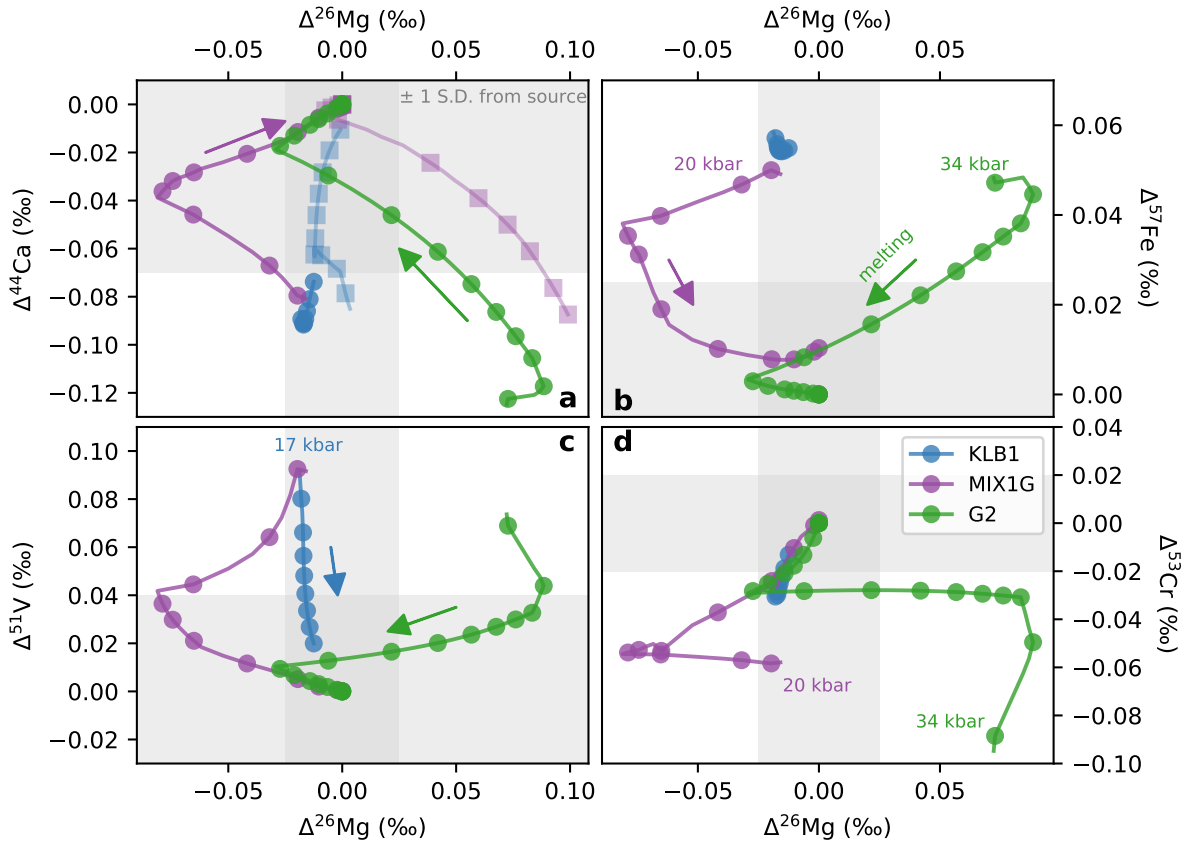


Figure S13: Melting fractionation ($\Delta_{\text{melt-source}}$) plotted along the $T_p = 1300^\circ\text{C}$ isentrope; squares in Mg–Ca space show $T_p = 1530^\circ\text{C}$ isentrope for reference. High or low pressure end of isentrope is marked, and arrows show the direction of progressive melting (decompression). Markers are plotted every 3 pressure steps (each pressure step is 0.5 kbar). Grey shaded regions show the source (0‰) \pm typical long-term analytical 1 S.D. (see Table S2 for details).

391 9 Element budgeting between lithologies

392 To produce Fig. 7 in the main text, we use a mass balance calculation for each isotopic system,
 393 X, where

$$\delta X_{\text{bulk}} = (F_{\text{pyroxenite}}^X \times \delta X_{\text{pyroxenite}}) + (F_{\text{peridotite}}^X \times \delta X_{\text{peridotite}}) \quad (11)$$

394 and F^X is the fraction of the total element in the final melt coming from each lithology, $\delta X_{\text{lithology}}$
 395 is the isotopic composition of each lithology (bulk isotope composition of lithology + melt-source
 396 isotope fractionation). The F term considers both the amount of the element in the melt from

Table S13: Inputs used to calculate bulk melt-source isotope fractionation in Figs. 7 and S14. * indicates inputs using the pure pyroxenite isentropes at that T_p , rather than peridotite isentropes. The MIX1G 1400* inputs are actually for an isentrope of $T_p = 1436^\circ\text{C}$. The P-T given are the closest P-T pair in the model output files to the average P-T of melting along each isentrope, as calculated in Section 6. For V and Cr, where element partitioning is not calculated by THERMOCALC, the melt element concentrations are calculated as described in Appendix Section 3, using the following bulk compositions for each lithology: KLB1 $\text{Cr}_2\text{O}_3 = 0.32\text{ wt } \%$, G2 $\text{Cr}_2\text{O}_3 = 0.08\text{ wt } \%$, MIX1G $\text{Cr}_2\text{O}_3 = 0.11\text{ wt } \%$ (Jennings et al., 2016); KLB1 V = 85 ppm (typical primitive mantle value from Lee et al., 2003; Prytulak et al., 2013), MIX1G and G2 V = 350 ppm (typical MORB value; Prytulak et al., 2013).

T_p ($^\circ\text{C}$):	KLB1	MIX1G				G2		
	1300	1400	1300	1300*	1400	1400*	1300	1300*
Pressure (kbar)	12	17.5	15	13.5	24.5	22	27	25.5
Temperature ($^\circ\text{C}$)	1300	1400	1317	1248	1440	1380	1345	1269
Melt fraction	0.061	0.099	0.323	0.093	0.369	0.155	0.458	0.048
Melt element conc.								
Mg (mol % oxide)	16.50	20.43	15.44	11.01	19.42	16.59	5.92	4.48
Ca (mol % oxide)	12.72	13.62	12.66	10.63	13.09	11.83	9.47	7.53
Fe (mol % oxide)	6.36	7.64	8.93	7.46	9.27	10.21	4.22	3.07
V (wt % element)	0.002	0.003	0.014	0.008	0.015	0.009	0.021	0.011
Cr (wt % element)	0.005	0.009	0.004	0.001	0.007	0.002	0.004	0.001
Melt-source isotopic fractionation (‰)								
$\Delta^{26}\text{Mg}$	-0.02	-0.02	-0.08	-0.10	0.08	0.02	0.06	0.09
$\Delta^{44}\text{Ca}$	-0.09	-0.08	-0.04	-0.05	-0.08	-0.08	-0.08	-0.12
$\Delta^{57}\text{Fe}$	0.06	0.04	0.03	0.06	0.03	0.04	0.03	0.06
$\Delta^{51}\text{V}$	0.05	0.03	0.03	0.08	0.02	0.05	0.03	0.07
$\Delta^{53}\text{Cr}$	-0.03	-0.02	-0.05	-0.06	-0.04	-0.05	-0.03	-0.12
Bulk isotopic composition of lithology (‰)								
$\delta^{26}\text{Mg}$	-0.24		-0.25				-0.25	
$\delta^{44}\text{Ca}$	0.94		0.83				0.83	
$\delta^{57}\text{Fe}$	0.05		0.14				0.14	
$\delta^{51}\text{V}$	-0.91		-0.92				-0.92	
$\delta^{53}\text{Cr}$	-0.12		-0.15				-0.15	

397 each lithology, and the melt fraction of each lithology. Table S13 gives the inputs used: these
398 inputs are taken from the full model results at the average P-T melting condition (as in Section
399 6) for each T_p . This calculation cannot be done for a T_p of 1530°C (MIX1G & G2) or 1400°C
400 (G2) because the isentropic melting paths intersect the solidus at pressures below the base of
401 our model, so an accurate average pressure of melting cannot be calculated. The bulk isotope
402 compositions of each lithology are the BSE values for peridotite (see Fig. 1 caption in main text
403 for sources), and average MORB composition for pyroxenite (except for Cr where no known
404 MORB data are published and the modelled aggregate KLB1 melt isotope composition for T_p
405 = 1300°C is used).

406 Fig. S14 is an alternative to Fig. 7 in the main text, with the right hand panels scaled to typical
407 current 1 S.D. analytical precision.

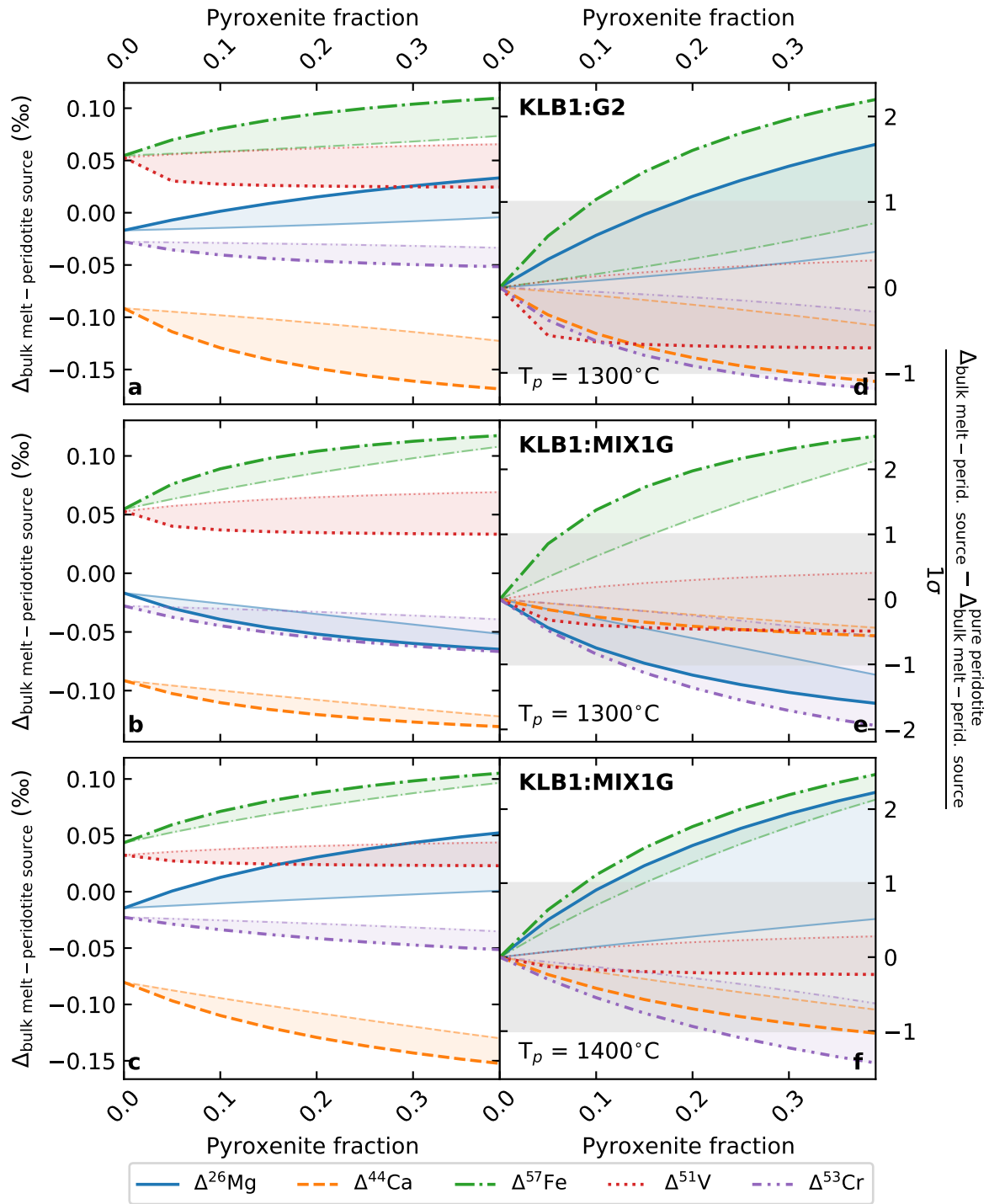


Figure S14: Aggregate bulk melt isotope fractionation (relative to a pure peridotite source, left hand panels; relative to a pure peridotite melt and scaled to current typical long-term analytical 1 S.D., right hand panels) for multi-lithologic mantle. At each T_p , the aggregate melt composition (including isotope composition, and element proportion) at the average pressure and temperature of melting for each lithology is used to calculate the resulting bulk melt isotope composition by mass balance. A range of bulk melt-source isotope fractionations are given for each isotope system, reflecting two endmembers for the thermal behaviour of pyroxenite upwelling in a peridotite host. The bolder line represents the case for pyroxenite following peridotite geotherms (as used throughout this study); the paler line reflects pyroxenite following its own geotherm, as in Figs. S10 and S11. Grey horizontal bars in the right hand panels highlight 1 S.D.

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