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# Uncertainty quantification of geological model parameters in <sup>1</sup> 3D gravity inversion by Hessian informed Markov chain Monte Carlo <sup>2</sup> Zhouji Liang\*, Florian Wellmann\* and Omar Ghattas<sup>†</sup> <sup>3</sup> \*Computational Geoscience and Reservoir Engineering(CGRE),

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# ABSTRACT

Geological modeling has been widely adopted to investigate underground geometries. How-12 ever, modeling processes inevitably have uncertainties due to scarcity of data, measurement 13 errors and simplification of modeling method. Recent developments in geomodeling methods 14 have introduced a Bayesian framework to constrain the model uncertainties by consider-15 ing additional geophysical data into the modeling procedure. Markov chain Monte Carlo 16 (MCMC) methods are normally used as tools to solve the Bayesian inference problem. To 17 achieve a more efficient posterior exploration, advances in MCMC methods utilize deriva-18 tive information. Hence, we introduce an approach to efficiently evaluate second-order 19

derivatives in geological modeling and introduce a Hessian-informed MCMC method, the generalized preconditioned Crank-Nicolson (gpCN), as a tool to solve the 3D model-based gravity Bayesian inversion problem. The result is compared with two other widely applied MCMC methods, random walk Metropolis-Hasting and Hamiltonian Monte Carlo, on a synthetic three-layer geological model. Our experiment demonstrates that superior performance is achieved by the gpCN, which has the potential to be generalized to more complex models.

# INTRODUCTION

In many geoscience applications, inversion methods are used to estimate subsurface prop-27 erties (e.g., structure, density and porosity) from observed geophysical data. Conventional 28 geophysical inversion aims to find the best-fit parameter sets that minimize the error be-29 tween observed geophysical data and simulation results. However, in practical cases, ob-30 servational data have the general difficulties of sparsity and noise. The uncertainties in the 31 data are combined with the geometry errors from the geological model constructions, lead-32 ing to numerous possible solutions within the uncertainty range, and hence ill-posedness. 33 Geoscientists are therefore interested in not only the best-fit model parameters but also a 34 quantification of uncertainties associated with these parameters (e.g., Sen and Stoffa, 1996; 35 Wellmann et al., 2018; Witter et al., 2019). 36

In this study, we focus on the uncertainty quantification in model-based inversion prob-37 lems. Developments in geomodeling have provided us with tools to generate 3D structural 38 geological models based on interface and orientation information from observed or inferred 39 data (Caumon et al., 2009; Caumon, 2010; Wellmann and Caumon, 2018). In contrast to 40 directly inverting the property field (e.g. density or thermal conductivity field), parametric 41 geological models have the advantage of easy interpretability and low parameter dimension-42 ality. Geophysical data can be used as an additional constrained to the developed model? 43 The gravity method has been widely applied in geophysics to investigate subsurface geome-44 tries and properties (Nabighian et al., 2005). While conventional gravity inversion often 45 suffers from ill-posedness due to insufficient data and intrinsic ambiguity (Boulanger and 46 Chouteau, 2001), it is suitable as additional data to constrain a geological model. 47

48

Recent developments in structural geomodeling methods allow geoscientists to quantify

the uncertainties in geological models based on prior geological knowledge and additional 49 geophysical data in a Bayesian inference approach (de la Varga and Wellmann, 2016). The 50 Bayesian approach has long been used as a tool to quantitatively combine uncertainties from 51 various sources in a probabilistic workflow (Tarantola and Valette, 1982; Mosegaard and 52 Tarantola, 1995; Sambridge and Mosegaard, 2002; Sambridge et al., 2013). In a Bayesian 53 inference problem, the parameters of interest (prior information) and the observational data 54 (likelihood function) are treated probabilistically. The inference provides us the updated 55 uncertainty range of the parameters of interest given the observational data. This output 56 uncertainty range is known as the posterior, which is also a distribution. However, it is 57 often not possible to calculate the posterior distribution analytically. The Markov chain 58 Monte Carlo (MCMC) is used to sample from the posterior distribution by generating 59 Markov chains of samples e.g. according to the Metroplis-Hastings algorithm (Metropolis 60 et al., 1953; Hastings, 1970). The posterior distribution can then be approximated by 61 the density distribution of the generated samples. The samples can also provide useful 62 statistical information about the posterior probability density function (PDF) (e.g., mean 63 and standard deviation). 64

The random walk Metropolis-Hasting (RMH) algorithm is likely the most popular 65 MCMC method. It is the simplest algorithm and has been widely adopted from the pi-66 oneering applications in traditional geophysical inversion (Mosegaard and Tarantola, 1995; 67 Malinverno and Leaney, 2000; Malinverno, 2002) through the recent applications of geo-68 logical modeling inversion (de la Varga and Wellmann, 2016; Wellmann et al., 2018; de la 69 Varga et al., 2019). While the classical RMH algorithm is simple in its intuition and imple-70 mentation, it is often inefficient for complex posterior distributions. Complexity may result 71 from high dimensionality or strong nonlinearities of the model. This leads to slow conver-72

<sup>73</sup> gence of the MCMC chain (Ruggeri et al., 2015). The inefficiency results primarily from <sup>74</sup> the mismatch between the proposal distribution used in the random walk method and the <sup>75</sup> actual posterior distribution. The recent work of Scalzo et al. (2019) used a preconditioned <sup>76</sup> Crank–Nicolson (pCN) method to address the high dimensionality problem in the context <sup>77</sup> of the joint geophysical inversion problems.

In cases where the gradient of the negative logarithm of the posterior with respect to 78 the parameters can be obtained easily and efficiently, gradient information can accelerate 79 an MCMC method by biasing samples toward higher probability regions. The Metropolis 80 adjusted Langevin algorithm (MALA) (Rossky et al., 1978; Roberts et al., 1996), which 81 involves Langevin diffusions, utilizes gradient information to confine the induced pseudo-82 dynamics (Betancourt, 2019). MALA can significantly increase the efficiency of MCMC. 83 However, its performance deteriorates for strongly anisotropic distributions since the gradi-84 ent may no longer point in a globally useful direction. A recent example of the application 85 of this algorithm to the seismic inversion problem can be found in the work of Mosser 86 et al. (2020). Another popular gradient-informed algorithm, the Hamiltonian Monte Carlo 87 (HMC) method, also known as hybrid Monte Carlo, was first introduced by Duane et al. 88 (1987). HMC analogizes the Hamiltonian flow of a trajectory in a potential field, which helps 89 to draw more independent samples, at the cost of a leapfrog integrator for the Hamiltonian 90 dynamics at each iteration (Betancourt, 2019). In addition, the parameters for the leapfrog 91 iterations require careful tuning, and the computational cost of leapfrog steps can be sub-92 stantial (Girolami and Calderhead, 2011). Recent studies have shown the applications of 93 HMC in geophysical inversion (Muir and Tkalcic, 2015; Fichtner et al., 2019). 94

The use of higher-order derivative information in MCMC, in particular the Hessian of the negative log-posterior has been shown to lead to more rapid exploration of the posterior

by adapting to the posterior curvature and anisotropy (Geweke and Tanizaki, 1999; Qi and 97 Minka, 2002). Explicitly constructing the Hessian for high dimensional inverse problems 98 is prohibitive. Instead, for many ill-posed inverse problems, the Hessian of the negative 99 log-likelihood admits a low rank approximation, and this can be efficiently computed using 100 a randomized eigensolver along with higher order adjoints (Martin et al., 2012; Petra et al., 101 2014; Isaac et al., 2015). A number of advanced Hessian-based MCMC methods for large-102 scale inverse problems have appeared in the past decade, for example, the stochastic Newton 103 MCMC method (Martin et al., 2012), dimension-independent, likelihood-informed MCMC 104 (Cui et al., 2016) and adaptive Gaussian process emulated geometric Monte Carlo (Lan 105 et al., 2016). These Hessian-based methods overcome the slow convergence of gradient-only 106 MCMC for ill-posed problems and are applicable in high-dimensional parameter space. 107

Developing first and higher order adjoints for complex simulations can be challenging. 108 As an alternative, automatic differentiation (AD) can be employed. The earlier development 109 of geological modeling methods by de la Varga et al. (2019) utilized automatic differenti-110 ation (AD) to make gradient evaluations tractable. Recent work of Güdük et al. (2021) 111 presented the successful application of this technique to model-based probabilistic inversion 112 using HMC. In this work, we introduce the Hessian-based MCMC method as a tool to con-113 duct efficient uncertainty quantification in model-based geophysical simulation. We follow 114 the implicit geological modeling method and forward geophysical simulation introduced by 115 De La Varga et al. (2015); de la Varga et al. (2019) and adapted the Hessian-based MCMC 116 algorithm developed by Villa et al. (2018). Here, we first find the maximum a posteri-117 ori (MAP) point by using Adam, a adaptive gradient-based optimization method (Kingma 118 and Ba, 2014). Then we construct the Laplace approximation of the posterior PDF by 119 equating the posterior covariance to the inverse of the Hessian of the negative-log poste-120

rior at the MAP point. A generalized preconditioned Crank-Nicolson proposal (Rudolf and 121 Sprungk, 2018) is then applied to conduct the MCMC sampling. Most importantly, to effi-122 ciently calculate the second-order derivatives, we further extend the geomodeling framework 123 GemPy (de la Varga et al., 2019) to an implementation in TensorFlow (Abadi et al., 2016), a 124 machine-learning framework that allows fast derivative evaluation through AD. The second-125 order derivative is calculated by a forward-over-backward propagation to achieve time and 126 memory efficiency. We establish a conceptual 3D fold model and simulate synthetic gravity 127 as the observational data. We compare the efficiency of gpCN with the RMH and HMC 128 methods for this inverse problem. The results show the potential of Hessian-based MCMC 129 methods in real-world geophysical inversion applications. 130

# PRELIMINARIES

In this section, we briefly introduce the forward model, including the construction of the structural geological model and the gravity forward simulations, and formulate the Bayesian inverse problem. We review two commonly used MCMC algorithms namely the RMH algorithm and HMC algorithm, which were implemented in our study as a comparison.

# 135 Implicit modeling

To generate the 3D geological models, we applied a implicit surface representation method, specifically the universal co-kriging method, originally introduced by Lajaunie et al. (1997), which has been implemented in the recently developed open-source software Gempy (de la Varga et al., 2019). A geological model can be characterized by a series of interface points and orientations. Compared with explicitly defining the geometry of the model in an explicit method, the implicit method handles the unknown area and complex geometries automatically using an interpolation function, thereby simplifying the modeling procedure in some
cases (Wellmann and Caumon, 2018). Here, we will briefly describe the universal co-kriging
method.

A 3D geological model can usually be described by several geological units bounded 145 by interfaces or faults. These interfaces and faults are typically the result of changing 146 sedimentary environments or tectonic events during geological history. The general idea 147 of the modeling method is to interpolate a scalar field  $Z(\mathbf{x})$  over the entire model in the 148 continuous 3D space  $(x, y, z) \in \mathbb{R}^3$  based on the spatial relationship of the interfaces and 149 orientations, where  $\mathbf{x} = \{x, y, z\}$  in Euclidean coordinates. Two parameters of the interfaces 150 are considered to characterize the scalar field namely the position of the layer interface points 151 and their orientations. We assume that points set  $J_K$  at the same interface K has the same 152 scalar value  $Z(x_j) \forall j \in J_K$ , the gradients of the scalar field  $\nabla Z^K$  are perpendicular to 153 the isocurve of the scalar field and parallel to the orientations of interface point G. These 154 constraints can be formulated as follows (Lajaunie et al., 1997): 155

$$\begin{cases} \nabla Z^{K}\left(\mathbf{x}_{i}\right) = G_{i} & \forall i \in I \\ \left\langle \nabla Z^{K}\left(\mathbf{x}_{i'}\right), \tau_{i'} \right\rangle = 0 & \forall i' \in I' \\ Z^{K}\left(\mathbf{x}_{j}\right) - Z^{K}\left(\mathbf{x}_{j}\right) = 0 & \forall j, j' \in J_{k} \end{cases}$$

Solving the co-kriging system based on the interface variance and covariance information and the above constraints provides a continuous scalar field. A 2D scalar field example is given in Figure 1. Explicit representation of geological surfaces can be extracted as isosurfaces in the scalar field. The full implementation is described in de la Varga et al. (2019).



Figure 1: Interpolated scalar filed using co-kriging of interface points and orientation values. Isocurves represents the scalar values based on the interpolation of the surface points and orientations.

#### <sup>161</sup> Gravity simulation

Gravity method, to date still one of the most widely used geophysical techniques in ex-162 ploration (Nabighian et al., 2005) can be used to integrate geophysical information into 163 geomdeling procedures. One of the most popular methods to include geophysical data is 164 through inversion (Tarantola, 2005). By including additional geophysical information, we 165 can validate the existing model or reduce its uncertainty range. A forward simulation of 166 the gravity field is required to conduct the inversion. To calculate the numerical solution 167 of the gravity field, we first discretize the model into cells and then assign densities to 168 each of these cells. Subsequently, we adapt the formulation presented by Nagy (1966) and 169 implementation by de la Varga et al. (2019) to simulate the gravity attraction from these 170 rectangular prisms. 171

$$F_z = \rho \cdot t_z$$
$$t_z = |||x \ln(y+r) + y \ln(x+r) - z \arctan\left(\frac{xy}{zr}\right) |_{x_1}^{x_2} |_{y_1}^{y_2} |_{z_1}^{z_2}$$

where x, y, and z are the Cartesian components of the targeting prism,  $\rho$  is density and r stands for the Euclidian distance from the center of the cell to the receiver points. In this work we use virtual receivers on the ground surface, which is usually the top of the model. By summing up the attractions of all the cells, the gravity at that receiver point can be simulated.

The resolution of the model, which reflects the size of the cells, has an impact on the simulation. Low-resolution cells introduce more uncertainties into the model. Additionally, receivers close to the boundary of the model will have boundary effects. Therefore, it is necessary to further extend the model to minimize the boundary effects. A model with a small cell size as well as a large model extent would certainly be ideal; however, our computational resources are limited. The attraction force decays exponentially with respect to distance r from the gravitation equation:

$$\Delta F = G\rho \frac{\Delta v}{r^2}$$

Hence, instead of using a regular grid (which means all cells have the same size), we apply a irregular grid (termed as centered grid in this paper) (see also (Güdük et al., 2021,p.5)). To simplify the calculation, we make this grid isotropic around a center position where the gravity response is calculated (illustrated in Figure 2). Both the height and widths of the prism cells grow exponentially with distance from the center. The scalar field is queried in these extended regions. For each receiver at the surface, a corresponding centered grid is created.

# <sup>184</sup> Bayesian Inference and Markov-chain Monte Carlo

The geophysical observations  $d_{obs} \in D$  in any forward problem can be described by  $d_{obs} =$ 185  $f(m) + e_d$ , where  $m \in M$  is the model parameters, f is the parameter-to-observation map 186 that represents the modeling or simulation procedure, and  $e_d$  is due to both the error in 187 modeling and the noise in the data. Geophysical inversion problems are often ill-posed 188 (Tikhonov and Arsenin, 1977) because of sparse and erroneous data, imperfect forward 189 modeling and loss of information in the parameter-to-observation map. Therefore, in many 190 cases we are interested not only in the best-fit parameters set but also other plausible solu-191 tions that fit the observations within an uncertainty range. Solving such inversion problems 192 then amounts to exploring the distribution of plausible model parameters. Bayesian infer-193



Figure 2: Illustration of centered grid with respect to one receiver. The grids are centered around the receiver, and extended further to alleviate the boundary effect.

ence, provides a systematic framework to quantify uncertainties of the inferred parameters,
given uncertainties in the forward model, the observational data, and any prior information
on the model parameters.

Bayesian inference starts from Bayes' theorem,

$$p(\boldsymbol{m} \mid \boldsymbol{d}_{obs}) = \frac{p(\boldsymbol{d}_{obs} \mid \boldsymbol{m}) p(\boldsymbol{m})}{p(\boldsymbol{d}_{obs})}$$

where  $p(\mathbf{d}_{obs} | \mathbf{m})$  is known as the likelihood of the data  $\mathbf{d}_{obs}$  given model parameters  $\mathbf{m}$ , and  $p(\mathbf{m})$  represents the prior probability of the model parameters  $\mathbf{m}$ . The posterior is normalized by the denominator  $p(\mathbf{d}_{obs})$ , which is called the marginal likelihood or evidence, so that the integral of the right-hand side is one. However, the denominator, marginal likelihood, is well known to be intractable in most cases because of the infinite possibility of

observations. Sampling methods, such as MCMC, are therefore often used to approximate 202 the integrals by sampling the posterior space. In the random walk Metropolis-Hasting 203 algorithm, which is the most well-known and intuitively simple MCMC algorithm, a chain 204 of samples is drawn based on a Gaussian proposal distribution, and by an accept/reject 205 step, the sampling chain can move from the previous state to a new state based on an 206 acceptance probability a. The accept/reject step restores the balance of the reversibility 207 condition; therefore, the target distribution can be represented by the chain of samples. 208 The steps of the Metropolis-Hasting algorithm can be expressed as follows: 209

	Al	gorithm 1: Metropolis-Hasting algorithm
	1 W	while $k < number of samples do$
210	2	Set $k = 0$ and initiate with $\boldsymbol{m}^{(0)}$
	3	Propose $\boldsymbol{m}^{cand} = \boldsymbol{m}^k + \beta \boldsymbol{\xi}^k,  \boldsymbol{\xi} \sim \mathcal{N}(0, C)$
	4	Set $\boldsymbol{m}^{k+1} = \boldsymbol{m}^{cand}$ with acceptance probability $a$
	5	Set $\boldsymbol{m}^{k+1} = \boldsymbol{m}^k$ otherwise
	6 e	nd

where acceptance probability  $a(\boldsymbol{m}^{cand}, \boldsymbol{m}^k) = \min\left(1, \frac{p(\boldsymbol{m}^k | \boldsymbol{d}_{obs})}{p(\boldsymbol{m}^{cand} | \boldsymbol{d}_{obs})}\right), \beta$  is the step size, and *C* is the covariance matrix of the proposal distribution.

The random-walk method has limitations in cases of high dimensionality and high correlations. Gradient information is therefore employed to assist in the posterior exploration. The most popular gradient-based method, the HMC algorithm, is designed to draw independent samples, and therefore efficiently explore the state space (Duane et al., 1987; Neal, 1993; Chen et al., 2014; Betancourt, 2017). HMC introduces auxiliary momentum variables  $\boldsymbol{r}$ , and therefore we can write the Hamiltonian of a particle as:

$$H(\boldsymbol{m},\boldsymbol{r}) = K(\boldsymbol{r}) + V(\boldsymbol{m})$$

where K and V are kinetic and potential energies respectively and are defined as:

$$K(\mathbf{r}) = \frac{1}{2} \sum_{i,j=1}^{n} r_i \mathbf{M}_{ij}^{-1} r_j$$
$$V(\mathbf{m}) = -\log p(\mathbf{m} \mid \mathbf{d})$$

Here, the mass matrix M is often set as the identity matrix (Chen et al., 2014)

Hamiltonian dynamics can be simulated by following Hamilton's equations with an artificially introduced time variable t:

$$\frac{\mathrm{d}\boldsymbol{m}}{\mathrm{d}t} = \frac{\partial H}{\partial \boldsymbol{r}}$$

$$\frac{\mathrm{d}\boldsymbol{r}}{\mathrm{d}t} = -\frac{\partial H}{\partial \boldsymbol{m}}$$
(1)

The HMC algorithm requires that Hamilton's equations are integrated using a symplectic integrator. In practice, the "leapfrog" integrator is often used. The algorithm can be <sup>218</sup> summarized as follows (Chen et al., 2014):

Algorithm 2: HAMILTONIAN MONTE CARLO ALGORITHM 1 while k < number of samples doSet k = 0 and initiate with  $\boldsymbol{m}^0$  and step size  $\epsilon$  $\mathbf{2}$ Sample momentum  $\boldsymbol{r}^k \sim \mathcal{N}(0, \boldsymbol{M})$ 3 Simulate discretization of Hamiltonian dynamics in Eq. 1:  $\mathbf{4}$  $\boldsymbol{r}_0 = \boldsymbol{r}_0 - \frac{\epsilon}{2} \nabla V\left(\boldsymbol{m}_0\right)$  $\mathbf{5}$ for i=1, i < leapfrog steps n do 6  $egin{aligned} m{m}_i &= m{m}_{(i-1)} + \epsilon m{M}^{-1} m{r}_{i-1} \ & m{r}_i &= m{r}_{i-1} - \epsilon 
abla V\left(m{m}_i
ight) \end{aligned}$  $\mathbf{7}$ 219 8 end 9 Set  $\boldsymbol{r}_n = \boldsymbol{r}_n - \frac{\epsilon}{2} \nabla V(\boldsymbol{m}_n)$ 10 Propose  $(\boldsymbol{m}^{cand}, \boldsymbol{r}^{cand}) = (\boldsymbol{m}_n, \boldsymbol{r}_n)$ 11 Set  $\boldsymbol{m}^{k+1} = \boldsymbol{x}^{cand}$  with acceptance probability a $\mathbf{12}$ Set  $\boldsymbol{m}^{k+1} = \boldsymbol{m}^k$  otherwise 13 14 end

220

where the acceptance probability 
$$a(\boldsymbol{m}^{cand}, \boldsymbol{m}^{k}) = \min\left(1, e^{-H(\boldsymbol{m}^{cand}|r^{cand})+H(\boldsymbol{m}^{k}|r^{k})}\right)$$

RMH is simple to implement and gradient-free, while HMC uses gradient information to obtain low autocorrelated samples; however, both methods are popular and widely applied. In this study, we will compare the efficiency of the generalized preconditioned Crank-Nicolson (gpCN) with both these methods and discuss their applicability.

# **METHODS**

In this work, we conducted an end-to-end procedure from the model construction to the 225 geophysical simulation and finally to Bayesian inference. We first generated a geological 226 model by using the implicit modeling method. A synthetic gravity field was then simulated 227 based on the model. Finally, we apply three different MCMC algorithms, namely RMH, 228 HMC, and gpCN to solve the Bayesian inference on the same synthetic model. In this sec-229 tion, we describe the Automatic differentiation method which links the geological modeling 230 and Bayesian inference framework and makes the high-order derivative evaluation tractable, 231 and present, how the high-order derivative information can be used to accelerate MCMC. 232

#### 233 gpCN

Although it is intuitively simple, the MH algorithm suffers from the problem of dimensional-234 ity (Cotter et al., 2013; Betancourt, 2017). As the dimension increases, the naive symmetric 235 proposal used in the MH algorithm will result in most of the samples being rejected, and 236 thereby resulting in an inefficient exploration of the posterior. A decent volumetric expla-237 nation was established by (Betancourt, 2017). Thus, researchers have been searching for an 238 alternative to the MH algorithm. Cotter et al. (2013) introduced the preconditioned Crank-239 Nicolson (pCN) proposal, which is a slight modification to the original RMH algorithm but 240 provides its scalability to high-dimensional problems. 241

In recent years, the idea of employing the geometry of the posterior to accelerate the exploration has attracted many researchers (Girolami and Calderhead, 2011; Martin et al., 2012; Law, 2014; Cui et al., 2016), where the first-order or second-order derivative information or either both (Martin et al., 2012) are utilized. Among those studies, Rudolf and Sprungk (2018) introduced the gpCN as an extension to the pCN algorithm. The key idea of this algorithm is to employ the covariance of the posterior  $C_v$  to construct a proposal distribution that adapts the posterior geometry Figure 3. The algorithm can be summarized as follows:

Algorithm 3: Generalized preconditioned Crank–Nicolson algorithm

1 Set 
$$k = 0$$
 and initiate with  $\boldsymbol{m}^{(0)}$ 

<sup>250</sup> 2 Propose 
$$m^{cand} = m_{\nu} + \sqrt{(1-\beta^2)}(m^k - m_{\nu}) + \beta \xi^k, \, \xi \sim \mathcal{N}(0, C_v)$$

**3** Set  $\boldsymbol{m}^{k+1} = \boldsymbol{m}^{cand}$  with acceptance probability a

4 Set  $\boldsymbol{m}^{k+1} = \boldsymbol{m}^k$  otherwise

where the acceptance probability 
$$a(\boldsymbol{m}^{cand}, \boldsymbol{m}^{k}) = \min \{1, \exp \left(\Delta \left(\boldsymbol{m}^{k}\right) - \Delta \left(\boldsymbol{m}^{cand}\right)\right)\},$$

and 
$$\Delta(\boldsymbol{m}) = \Phi(\boldsymbol{m}, \boldsymbol{d}_{obs}) + \frac{1}{2} \|\boldsymbol{m} - \boldsymbol{m}_{prior}\|_{C_{prior}^{-1}}^2 - \frac{1}{2} \|\boldsymbol{m} - \boldsymbol{m}_{\nu}\|_{C_{\nu}^{-1}}^2$$
,

 $\Phi(\boldsymbol{m}, \boldsymbol{d}_{obs})$  is the negative log likelihood function:

$$\Phi(\boldsymbol{m}, \boldsymbol{d}_{obs}) = rac{1}{2} \left\| \boldsymbol{f}(\boldsymbol{m}) - \boldsymbol{d}_{obs} 
ight\|_{\Gamma_{noise}^{-1}}^2$$

Subsequently, Villa et al. (2018) introduced the method to evaluate the posterior covariance by using the Laplacian approximation with the Hessian of the log likelihood  $\mathcal{H}_{misfit}$ at the MAP point  $m_{\nu}$  and covariance of the prior  $C_{prior}$ 

$$C_{\nu} = \left(\mathcal{H}_{misfit}\left(\boldsymbol{m}_{\nu}\right) + C_{prior}^{-1}\right)^{-1}$$

$$oldsymbol{m}_{
u} := rgmin \mathcal{J}(oldsymbol{m}) := \left(rac{1}{2} \left\|oldsymbol{f}(oldsymbol{m}) - \mathbf{d}_{obs}
ight\|_{\Gamma_{noise}^{-1}}^2 + rac{1}{2} \left\|oldsymbol{m} - oldsymbol{m}^{prior}
ight\|_{C_{prior}^{-1}}^2
ight)$$



Figure 3: Different proposal distribution illustrated on the 2D Gaussian posterior distribution  $d|m \sim \mathcal{N}(\mu, I), \mu := x_1 - 0.7x_2$ , where  $x_i \stackrel{iid}{\sim} \mathcal{N}(0, I)$ . Left shows the Gaussian proposal without preconditioning. Right shows the proposal distribution with Laplacian approximation at MAP.

Here, we applied the recently developed first-order optimization algorithm Nadam (Nesterovaccelerated adaptive moment estimation)(Dozat, 2016) to find the MAP. The update rule of Nadam is given as follows:

$$oldsymbol{m}_{t+1} = oldsymbol{m}_t - rac{lpha}{\sqrt{\hat{v}_t} + \epsilon} \hat{b}_t$$

where  $\alpha$  is the step size.  $\hat{b}_t$  and  $\hat{v}_t$  are the bias-corrected first and second momentum estimates of the gradients, respectively. The momentum term averages the past gradient and thus accelerates the convergence rate compared with the standard gradient descent algorithm, while the second momentum which averages the past squared gradients adapts its learning rate.

#### <sup>265</sup> Automatic Differentiation and Evaluation of Hessian

Implementing derivative-informed MCMC methods such as HMC and gpCN in geomodeling 266 require the derivative information in several steps such as finding the MAP point and 267 constructing the posterior covariance approximation. This derivative refers to the derivative 268 of the posterior (or negative log posterior) with respect to the model input parameters 269 (e.g., surface points and orientation points). By tracing the derivative from the Bayesian 270 inference through the geophysical simulation, geological modeling of the input parameters 271 is nontrivial. An analytical solution of the derivatives is difficult or even impossible to find. 272 Using methods such as finite difference (FD) is not only computationally costly but also 273 can suffer from numerical inaccuracy. Therefore, we adopted the automatic differentiation 274 (AD) technique, which is widely applied in the field of artificial intelligence and is critical 275 to the success of training neural networks. Here we briefly introduce how AD works and 276

<sup>277</sup> how gradient and higher-order derivatives can be evaluated efficiently using AD.

Consider an arbitrary function  $\mathbf{y} = F(\mathbf{x})$ , where F is the mapping function maps the input parameter  $\mathbf{x} = x_1, x_2, ..., x_i, (\mathbf{x} \in \mathbb{R}^{D_0})$  to the cost function  $\mathbf{y} = y_1, y_2, ..., y_j, (\mathbf{y} \in \mathbb{R}^{D_1})$ . The first-order partial derivative, also known as the Jacobian matrix, is given as follows:

$$(\mathcal{J}_F)^i_j(x) = \frac{\partial (F)^i}{\partial x^j}(x)$$

by constructing the composite function  $F = F_N \circ F_{N-1} \circ \cdots \circ F_2 \circ F_1$  and ignoring *i*, the Jacobian matrix can be represented by iteratively applying the chain rule:

$$\mathcal{J}_F(x) = \frac{\partial(F)}{\partial x^j}(x) = \frac{\partial(F_N)}{\partial(F_{N-1})} \dots \frac{\partial(F_2)}{\partial(F_1)} \frac{\partial(F_1)}{\partial x}(x)$$

AD records the derivative of each fundamental operator in a program and builds a computational graph to record the dependency of each operator. The derivative of the output with respect to the input parameter of interest is evaluated by a forward or a backward propagation using the chain rule based on the computational graph. Forward mode and backward mode both have their advantages and disadvantages, by combining these two methods, we get an efficient method for second order derivative evaluation. The approach is summarized in the following.

To evaluate the gradient in the forward mode AD, a tangent vector v is selected at the evaluation point x,

$$\mathcal{J}_F(x) \cdot v = \left(\frac{\partial (F_N)}{\partial (F_{N-1})} \dots \left(\frac{\partial (F_2)}{\partial (F_1)} \left(\frac{\partial (F_1)}{\partial x} (x) \cdot v\right)\right)\right)$$

For example,  $\frac{\partial(F)^i}{\partial(x_1)}$  can be evaluated by a tangent vector  $v = (1, 0, \dots, 0), v \in \mathbb{R}^i$ . Thus, the forward mode AD provides the directional derivative.

In contrast, backward mode AD evaluation is based on a fixed dependent variable, and the quantity of interest is the adjoint

$$\mathcal{J}_{F}^{i}(x) = \left( \left( \left( \frac{\partial(F_{N})}{\partial(F_{N-1})} \right) \dots \frac{\partial(F_{2})}{\partial(F_{1})} \right) \frac{\partial(F_{1})}{\partial x}(x) \right)$$

Evidently, for problem  $\mathbb{R}^{D_I} \to \mathbb{R}^{D_O}$ , where  $D_I$  denotes the dimension of input and  $D_o$ 294 denotes the dimension of output, if  $D_I \gg D_O$ , backward propagation is more efficient, 295 which matches the cases of many machine learning problems and the Bayesian inference 296 in our study, where the output is a single cost function, so  $D_O = 1$ . In contrast, forward 297 propagation is more efficient when  $D_I \ll D_O$ , requiring less memory. An efficient second-298 order derivate calculation for multivariate input and single output problem is given by the 299 combination of a forward-over-backward propagation. Applying backward propagation gives 300 us the Jacobian matrix  $\mathcal{J}_F(x) = \left[\frac{\partial F}{x_1}, \frac{\partial F}{x_2}, \dots, \frac{\partial F}{x_j}\right]$ . The second derivative then evaluates the 301 partial derivative over the Jacobian, which is an  $\mathbb{R}^{D_I} \to \mathbb{R}^{D_I}$  problem. Considering the 302 memory efficiency, we can evaluate the Hessian vector product efficiently by an additional 303 forward sweep over the backward propagation. Each forward-over-backward iteration will 304 return a column of the Hessian matrix; therefore, this method has an  $\mathcal{O}(n)$  complexity, 305 where n is the dimension of the parameter of interest. It is significantly more efficient 306 than evaluating the Hessian by the FD method, which has a complexity of  $\mathcal{O}(n^2)$ . A 307 comprehensive review of AD and its implementation can be found in Betancourt (2018) 308 and Margossian (2019). 309

Programming the geological model by adopting AD, a second-order derivative can be efficiently evaluated. The numerical comparison of the computational efficiency is discussed <sup>312</sup> in the following sections.

#### RESULTS

When the available data are limited, uncertainties are inevitable in 3D geological models. 313 We can constrain the uncertainties by additional geophysical data. To configure such a 314 problem in a Bayesian inference framework as described above, both the prior information 315 and the likelihood must be expressed in terms of a probability distribution. As an example 316 we constructed a 3D geological model using the implicit method described above. The 317 numerical implementation is based on the previously developed modeling methods of the 318 GemPy software (https://www.gempy.org/) and specific extension in the differentiable 319 programming framework TensorFlow (https://www.tensorflow.org/). 320

The example model has a simple dome geometry, replicating a conceptual setting which 321 is often observed in real-world geology. It consists of three layers, representing three distinct 322 lithologies. The model has an extent from 0 m to 1,000 m for all three axes. The positions 323 (x-y coordinates) of the interface points are fixed, and the depth z values of the interface 324 points are the uncertain parameters in this study. Both the layout of the interface points 325 and pseudo-gravity receivers is shown in Figure 4. The ground truth z value of the interface 326 point follows a sinusoidal function. The ground truth 3D model and interface point layout 327 are shown in Figure 5. To simplify the problem, we assume a constant distance between 328 the two surfaces at configuration interfaces positions and only vary the upper interface 329 points. In other words, the lower surface moves parallel in dependence of the upper surface; 330 therefore, a constant thickness of the middle layer is expected. Interpolation could introduce 331 some variation to the thickness at locations where no interface points are given. Sixteen 332 pseudo-gravity receivers are evenly located at the ground surface, which is the top surface 333



Figure 4: Top view of model parameter layout and receivers position. The interfaces point for lower rock is overlap by the top layer. Interface points 2 and 3 are set close to each other to simulate correlations.

of the model. The grid size for gravity calculation was selected to be  $10 \times 10 \times 30$  to have a better resolution in the vertical axis because the parameters of interest vary along the z axis. The simulated ground truth gravity data is shown in Figure 6. Our goal is to estimate the uncertainties of the depth position of the interface points of the upper surface considering the additional gravity information.

We assume a simple Gaussian prior to all eight interface points with the mean at the ground truth location  $m_0$  and a standard deviation (std) of 25 m. Therefore, the prior distribution can be expressed as a multivariate normal distribution centered at  $m_0$  and a diagonal covariance matrix assuming no correlations between the parameters in our prior



Figure 5: 3D plot of the base-case geological model. Gravity receivers placed at the top.



Figure 6: Simulated base-case at top surface.

information. Pseudo-noisy gravity observation data are described by another multivariate normal distribution as the likelihood function. Pseudo-noisy gravity observation data was generated by seven forward gravity simulations with an independent variation of 7 m.

To solve the Bayesian inference by MCMC, specifically by using the gpCN, we first 346 search for the Maximum a Posteriori (MAP) point  $m_{MAP}$  in the posterior space.  $m_{MAP}$ 347 can be sufficiently found by Nadam with the assistance of the gradient of the target negative 348 log posterior with respect to the variable (depth value). The gradient is evaluated through 349 AD. The time cost for each gradient evaluation is within 1 s depending on the model size. 350 To avoid the optimization to get stuck at local minima in the parameter space, an initial 351 status is generated by randomly sampling the prior distribution. The optimization can 352 converge normally within 200 iterations, depending on the model complexity, step size, and 353 the initial status. The full Hessian matrix of the target negative log posterior with respect to 354 the variable is evaluated at the MAP point through the forward-over-backward propagation 355 described above. 356

To compare the performance of gpCN, we implemented the two other most commonly 357 used methods, namely RMH and HMC, to solve the same inference problem. The simula-358 tions were run on a single Tesla P100 GPU. The resulting sampling chain, posterior, and 359 efficiency are compared below. Figure 7 shows the trace plot of the same two interface points 360 for each algorithm. We observe that better-mixed chains are obtained by both HMC and 361 gpCN, but the chain of RMH is poorly mixed, and far from convergent. To quantitatively 362 evaluate the chain, we plotted the autocorrelation of the results from all three methods in 363 Figure 8. A similar low autocorrelation is obtained by HMC and gpCN, and RMH samples 364 are highly autocorrelated with the lowest acceptance rate. One can always tune the step 365 size of MCMC to achieve fewer autocorrelated samples, but this will also lead to a lower 366

367 acceptance rate.

Although equivalent or even better sampling chain is obtained by HMC, gpCN is more 368 efficient in terms of computational cost. In HMC, the leapfrog integrator is usually chosen 369 to solve the numerical integration in the Hamiltonian equation, which is a computationally 370 demanding process. The choice of step size and the number of steps used for the leapfrog 371 integration are essential for the performance and efficiency of HMC. Too few steps will result 372 in a bad integration, and therefore, a lower acceptance rate, while increased number of steps 373 will increase the computational cost. Because HMC is intended to find an uncorrelated point 374 by using several evaluations, the theoretical acceptance rate of HMC is 100%, but slightly 375 lower due to the imperfection of numerical integration. Thus, in addition to the acceptance 376 rate and the autocorrelation, we also compared the effective sample size  $n_{\rm eff}$  (Liu, 2008,p. 377 125) per unit time. The effective sample size  $n_{\rm eff}$  is defined as follows: 378

$$n_{\rm eff} = \frac{n_{\rm max}}{1 + 2\sum_{k=1}^{\infty} \rho_k}$$

where  $\rho_k$  is the autocorrelation coefficient at lag k. A higher number of  $n_{\text{eff}}$  means a better mixing Markov chain. Our numerical experiments indicate that gpCN outperformed HMC by two orders of magnitude of the  $AvgN_{eff}$  (Table 1), while preserving the computational efficiency, even though we only used three leapfrog steps in the experiments for HMC. Although we demonstrate only a simple example, we can expect a better performance of the gpCN in a higher dimension and more correlated case.

The main reason for the superior performance of gpCN is the highly correlated posterior space. If we use the sample of HMC as a benchmark and plot the adopted proposal distribution used in the gpCN (Figure 9), we can see that the proposal distribution captured the



Posterior by RMH, epochs:10000, accept rate:64.85%

Figure 7: Trace plot of RMH, HMC and gpCN. RMH has a poor chain mixing and lowest acceptance ratio. Without considering computational efficiency, HMC has the best chain mixing and acceptance ratio.

epochs



Figure 8: Autocorrelation plot of the MCMC chain of all three methods. The solid lines represent the mean autocorrelation among different variables, and the shadow represents the 95% confidence interval of the autocorrelation on different dimensions.

Methods	Compution time for 10000 samples	$n_{\rm eff}$	$AvgN_{eff}$ per 100s
RMH	191s	4.0	2.1
HMC	1851s	60.8	3.3
gpCN	123s	385.5	312.5

#### Table 1: Computational efficiency of different methods

#### 388 posterior geometry.

Finally, we represent the uncertainties following the information entropy method intro-389 duced by Wellmann and Regenauer-Lieb (2012). In Figure 10, we compared the uncer-390 tainties of the Prior and Posterior. We generated 1000 samples from the prior distribution 391 and build the geological models based on the sampled parameters. The samples drawn by 392 the gpCN MCMC are used to generate geological models for posterior representation. By 393 converting the lithology changes into information entropy, we can see that the uncertainty 394 range of the position of the upper interface is significantly reduced by additional gravity 395 data. 396

To consolidate our conclusion, an additional example with less ideal prior information is conducted and compared. In this case, the prior is chosen to be a uniform distribution with the same variation but a uniform mean at  $z = 780 \ m$ . Analogue to the previous example, pseudo noisy data is generated by tuning the base-case model, and the inference results are compared with three different MCMC methods. Even with relatively poorer prior information, the posterior distribution of the surface can be sufficiently explored by



Figure 9: Posterior sampling of posint 0 and posint 2 by HMC and the proposal distribution used in gpCN in red contour. The red square represent the found MAP.



Figure 10: Information entropy plot of the upper boundary of Prior and Posterior. High entropy represents high uncertainties.



Figure 11: Inference results with uniform Prior distribution. Information entropy plot of the upper boundary of Prior and Posterior. High entropy represents high uncertainties.

the gpCN (Figure 11). Similar inference efficiency can be achieved in this case, with slightly more cost in the MAP searching process. This is because our initial guess drawn from the Prior distribution could deviate more from the target set. Similar posterior results are found in this example because of the low variance in the observation data, which dominates the prior information in the weights of posterior probability.

#### DISCUSSION

The gpCN method we adopted in this study is a MAP-based method. These MAP-based 408 methods utilize a global Hessian evaluated at the MAP point. The primary advantage of 409 such methods compared with other derivative methods such as HMC is that the Hessian 410 matrix is only required to be evaluated once. It works the best in cases where the posterior 411 is Gaussian-like. In addition to the adoption of pCN and the adapted proposal, gpCN out-412 performs RMH in high-dimension problems and a highly correlated posterior. Additionally, 413 finding the MAP and starting the MCMC from the MAP saves the burn-in computation 414 normally required in MCMC, because the MAP is already in the target distribution. 415

However, this global Hessian method is also limited when the model is highly nonlinear 416 and the posterior becomes non-Gaussian, such as a banana distribution (e.g. Lan et al., 417 2016). Hessian evaluated at the MAP point are not as representative, and therefore, such 418 bad adaptation could lead to a less efficient Markov chain. Such cases can be extended 419 into the state-dependent local approximation with additional computational costs (Petra 420 et al., 2014; Rudolf and Sprungk, 2018). Another limitations is the ability to explore the 421 multi-modes posterior distribution. In these cases, the posterior could be non-convex, and 422 the global minimum is normally not easily found by simple gradient-based optimization 423 methods such as Nadam used in this study. This is a common issue in optimization in many 424 applications (Nocedal and Wright, 2006). Second-derivative optimization methods such as 425 Newton's methods, could be potentially be used to assess the MAP exploration(Villa et al., 426 2018). 427

The high computational cost of the high-order derivative is the main obstacle for the application of advanced MCMC methods involving high-order derivatives. The efficient

computation of high-order derivatives remains an open research problem (Laue et al., 2018; 430 Nilsen et al., 2019; Margossian, 2019). In our work, the computation time required for a 431 single Hessian calculation with the current implementation depends on the complexity of 432 the computational graph used in AD and the number of parameters of interest. Although 433 the evaluation of the full Hessian matrix is still more expensive than the gradient evaluation 434 using a single backward propagation in AD, it is significantly faster than computing Hessian 435 by the FD method, regardless of the precision. The numerical analysis is provided in 436 Figure 12. As described previously, AD has an  $\mathcal{O}(n)$  complexity, while FD has a complexity 437 of  $\mathcal{O}(n^2)$ . A full Hessian matrix requires  $\frac{n^2+n}{2} \times 4 \times t$  by using FD, where n represents the 438 dimension of the parameter of interest, and t is the time cost for a single forward simulation. 439

In addition to the computational efficiency, the computational graph employed in AD also allows any variable in the graph to be traced with a derivative with minor modifications to the computer program.

This study is the first step towards the application of Hessian-informed MCMC in geo-443 logical modeling. With a more efficient implementation of the Hessian evaluation, we can 444 expect a good performance of these methods in higher-dimensional problems. This work 445 has constructed the Hessian explicitly, which is appropriate for the size of problems consid-446 ered. However this is not necessary, as shown for large-scale problems in Villa et al. (2021); 447 Isaac et al. (2015); Bui-Thanh et al. (2013). The data misfit Hessian  $(H_{misfit})$  admits a low 448 rank approximation (due to ill-posedness), and this approximation can be obtained with 449  $\mathcal{O}(r)$  Hessian-vector products using the randomized SVD algorithm, where r is the effective 450 rank of  $H_{misfit}$ . Once this low rank approximation is constructed, the Sherman-Morrison 451 formula can be used to efficiently draw samples from the Laplace approximation, as the 452 gpCN algorithm requires. Thus only  $\mathcal{O}(r)$  Hessian-vector products are required, each of 453



Figure 12: Comparison of computation time of a full Hessian matrix by Finite Difference(FD) method and Auto-Diff(AD). FD computation time is approximated by extrapolating the computation time for a single forward simulation. AD computation time is simulated with the same configuration used in the previous example.

which requires the solution of a pair of forward/adjoint models. This is in contrast with explicitly building the Hessian, which requires  $\mathcal{O}(n)$  forward model solutions. The rank rdepends on how informative the data are, i.e. how ill-posed the inverse problem is. For many problems,  $r \ll n$ , and r is independent of the parameter dimension. Therefore one can benefit greatly from Hessian information, even for high-dimensional parameters.

# CONCLUSION

In summary, this study extended the previous development of stochastic geological modeling 459 methods. We used the automatic differential technique implemented in TensorFlow to 460 allow second-order derivative information to be efficiently evaluated in geological models. 461 We applied the recently developed Hessian-informed MCMC, the generalized preconditioned 462 Crank-Nicolson (gpCN), to solve the Bayesian inference problem on a synthetic three-layer 463 geological model. We compared gpCN with the other two commonly used MCMC methods, 464 including RMH and gradient-based method HMC. The results demonstrate with a single 465 Hessian evaluation, gpCN outperforms RMH while preserving the computational efficiency, 466 with no additional computational cost at each step, which has the potential to be generalized 467 to more complex models. 468

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