The relationship between lake surface area and maximum depth

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Abstract: Maximum depth varies among lakes from <1 to 1741 meters, but attempts to explain this variation have achieved little predictive power. In this paper, we describe the probability distribution of maximum depths based on recent developments in the theory of fractal Brownian motions. The theoretical distribution is right-tailed and adequately captures variations in maximum depth in a dataset of 8,164 lakes (maximum depths 0.1 to 135 meters) from the northeastern United States. Maximum depth increases with surface area, but with substantial random variation - the 95% prediction interval spans more than an order of magnitude for lakes with any specific surface area. Our results explain the observed variability in lake maximum depths, capture the link between topographic characteristics and lake bathymetry, and provide a means to upscale maximum-depth-dependent processes, which we illustrate by upscaling the diffusive flux of methane from northern lakes to the atmosphere.

Significance Statement: Many foundational ecosystem characteristics in lakes, such as thermal stratification, are constrained by maximum depth. Aquatic scientists have have struggled to both explain this variation and to predict maximum depths for lakes without detailed bathymetric surveys. The morphometry of collections of lakes has previously been shown to follow theoretical predictions that assume Earth's topography approximates a fractal Browning motion. In this study, we demonstrate how these theories relate maximum depth to surface area. Maximum depth increases with surface area, but massive lake-to-lake variation is an inherent characteristic of this pattern. This is because the maximum depth is a single random displacement on a topographic profile. The probability distribution of maximum depths can be used to upscale the patterns and processes for many lakes at the regional or global scale. We demonstrate this for diffusive flux of methane from temperate lakes to the atmosphere - a process that correlates strongly with maximum depth but not surface area and cannot be up-scaled without the results presented in this paper.

¹ Introduction

² Maximum depth varies among lakes between ~0.1 and 1741 meters [27, 13]. This variation engen-³ ders patterns of diverse ecosystem characteristics including mixing and thermal stratification [15], the ⁴ relative sizes of littoral and pelagic habitats [22], and carbon cycling including methane evasion [16]. ⁵ However, there is a paucity of bathymetric data relative to the global abundance of lakes, and em-⁶ pirical relationships that relate maximum depth to lake and landscape characteristics consistently fail ⁷ to develop sufficient predictive power to estimate maximum depth in lakes that have not been depth ⁸ sounded [26]. In particular, there is need to develop scaling relationships that relate maximum depth ⁹ to lake characteristics that are easily measured across broad geographic regions, such as surface area ¹⁰ (e.g., [2, 23]). Such relationships provide the simple rules used to generalize understanding of aquatic ¹¹ ecosystem patterns and processes at regional to global scales [7].

Bathymetric surveys are time consuming, and prohibitively expensive for large numbers of lakes [26, 12 14]. While global perspectives have come to dominate the aquatic sciences over the last twenty years, 13 the proportion of lakes depth sounded has remained low, preventing up-scaling of empirical results 14 from local to global scales [26, 21, 8]. This has spurned a series of empirical studies seeking to 15 predict lake-specific maximum depth based on surface area and other easily mapped characteristic, 16 typically some metric of vertical relief within a buffer zone around each lake (e.g., [21, 14, 27, 19, 18, 17 12). Typically, these studies assume that larger lakes should be deeper than smaller lakes, and that 18 integrative measures of topography (i.e. variance, slope) should relate to lake bathymetry. However, 19 these might not be reliable assumptions. For example, while the deepest lakes all have large surface 20 areas ($> 500 \text{ km}^2$), many lakes with large surface areas are remarkably shallow, often only a few meters 21 deep [13]. Further, correlations between surface area and maximum depth are often weak. For example, 22 the Pearson correlation between the logarithms of surface area and maximum depth for Canadian lakes 23 is r = 0.46 [19]. Based on this correlation, the probability of a larger lake being deeper than a smaller 24 lake is only slightly better (p = 0.65) than a fair coin flip (p = 0.5), if the two lakes are selected at 25 random [9]. Additionally, it is not clear that topography should predict maximum depth. Specifically, 26 many lakes have relatively sudden and catastrophic origins that may not reflect the processes shaping 27 topography [29]. Additionally, maximum depth is essentially a single random perturbation along a 28 combined topographic-bathymetric profile [26]. While integrative measures of topography may be able 29 to predict integrative measures of bathymetry, there is no clear reason to believe that they should 30 accurately predict the value of a single random displacement [26]. The consequence of this is highly 31 uncertain predictions that preclude the application of these equations for upscaling. For example, the 32 error ratio for maximum depth predictions for Canadian lakes is approximately two [3, 19]. This means 33 that a lake predicted to have a maximum depth of 10 meters will have a maximum depth between 5 and 34 20 meters [3]. Error ratios are even higher in other studies (e.g., [12]). This magnitude of prediction 35 error is material across the entire size spectra of lakes. 36

In this paper, we describe the theoretical basis for predicting maximum depth from surface area. This theory is based on the characteristics of fractional Brownian motions, which approximate key features of Earth's topography and are the basis for other lake scaling relationships including mean depth and volume (e.g., [11, 23, 2, 25]). Specifically, we describe how recent developments generalizing the arcsine laws from Brownian motions to fractional Brownian motions relate to the problem of maximum depths [4, 5]. We test goodness-of-fit of the theoretical distribution derived based on these



Figure 1: An example fractal Brownian motion path, with Hurst exponent H = 0.4. This path is taken to be an analog of a transect along a landscape; lakes fill in sections along this transect, which thus sets their maximum depth.

theoretical developments to a bathymetry database with more than 8000 lakes. Finally, demonstrate how this distribution can be applied to advance understanding of global lake characteristics by upscaling diffusive methane flux from temperate lakes to the atmosphere, a process that has previously been shown to correlate with maximum depth but not surface area, and hence a process that is difficult to upscale without the results presented in this paper [16]. Collectively, these results both advance fundamental understanding of patterns of lake morphometry, and provide tools for upscaling depth dependent processes when seeking to place lakes within a global context.

50 Theory

Earth's topography is approximately scale-invariant [30]. This is self-evident for many landforms including lakes. For example, it would be difficult or impossible to determine if a lake is 1 ha or 1,000 ha on a map or image without a scale reference (e.g., a scale bar on a map or a house on an aerial image) because the key characteristics of lake geometry are similar across wide ranges of scales [25, 26, 23]. In particular, Earth's topography is self-affine, which implies different scaling in the vertical and horizontal directions [30]. This is also self-evident; when walking or driving away from a mountain

range, the profile flattens more rapidly than it compresses horizontally [26]. Specifically, the Hurst 57 exponent H captures the relationship between horizontal and vertical scales. For a surface $Z(\vec{X})$, 58 rescaling by any coefficient b conforms, in a statistical sense, to $b^{-H}Z(b\vec{X}) = Z(\vec{X})$. H takes values in 59 the range (0,1). When $H = \frac{1}{2}$, transects across Z have the statistical properties of a one-dimensional 60 Brownian motion; when $H > \frac{1}{2}$ (< $\frac{1}{2}$), surfaces are smoother (rougher) than Brownian trajectories 61 (e.g., [11]). These characteristics engender the derivation of empirically robust and theoretically sound 62 power-law relationships between lake surface area and various aspects of lake morphometry including 63 abundance, volume, mean depth, perimeter, and hydrological connectivity [2, 23, 25, 23]. However, 64 to our knowledge, the relationship between maximum depth and surface area has not been described 65 [26].66

⁶⁷ When topography approximates a Brownian motion $(H = \frac{1}{2})$, the maximum depth over a given interval ⁶⁸ is described by the third arcsine law [4]. Essentially, because the maximum depth is a single random ⁶⁹ displacement on a topographic profile, the arcsine law shows that maximum depth converges as a ⁷⁰ probability distribution based on surface area. Recently, this result has been generalized to cases where ⁷¹ $H \neq \frac{1}{2}$ [4, 5]. This generalization allows application of the arcsine law to predict the distribution of lake ⁷² maximum depth p(z) based on surface area and the Hurst coefficient. Unfortunately, the mathematical ⁷³ form of p(z) is almost comically complex. First, define the normalized maximum depth

$$y = \frac{z}{\sqrt{2}L^H}$$

⁷⁴ and define the probability distribution in terms of y, i.e. p(y) (later we recast this in terms of z as ⁷⁵ well). Doing so,

$$p(z) = \frac{1}{\sqrt{2}L^H}f(y)$$

where $\frac{1}{\sqrt{2}L^H}$ is a normalization constant that ensures the total probability integrates to one, and f(y) π has the form:

$$f(y) = \sqrt{\frac{2}{\pi}} e^{-y/2} y^{(1/H)-2} e^{(H-\frac{1}{2})(4\ln y + \mathcal{G}(y))}$$

⁷⁸ and $\mathcal{G}(y)$ is ithe complex expression:

$$\mathcal{G}(y) = \frac{y^4}{6} \, _2F_2\left(1, 1; \frac{5}{2}, 3; \frac{y^2}{2}\right) - 3y^2 + \pi(1 - y^2) \mathrm{erfi}\left(\frac{y}{\sqrt{2}}\right) + \sqrt{2\pi}e^{y^2/2}y + (y^2 - 2)(\gamma_E + \ln(2y^2))$$

⁷⁹ where F is the hypergeometric function, erfi is the imaginary error function, and γ_E is the Euler-⁸⁰ Mascheroni constant [4, 5].

To date, the arcsine laws for fractal Brownian motions are only proven for the one-dimensional case. 81 Therefore, lake areas must be recast as one-dimensional lengthscales L to empirically test the theo-82 retical depth distribution. The lengthscale L is in a sense the length of a randomly chosen horizontal 83 transect along a lake's bathymetry, which includes the lake's maximum depth (solid black line in Figure 84 1). As surface area $a \, [m^2]$ is the fundamental horizontal metric for lakes, it would be preferable to 85 cast L in terms of a; from unit considerations one necessarily must have $L \propto \sqrt{a}$, but the coefficient 86 relating the two is not easily derived from first principles because lake surfaces often have complex 87 shapes, and their bathymetry can also be complex, including multiple basins and/or having maximum 88 depths far from their centers. We therefore introduce the free parameter ℓ which allows us to relate 89 L to a according to $L = \ell \sqrt{a}$. ℓ should be no larger than what it would be for a circular lake with 90 its maximum depth at the center, i.e. $\ell = \sqrt{2/\pi} \approx 0.8$, but is likely much smaller as lakes take 91 myriad shapes that are often very far from circular, and their maximum depths of course do not have 92 to be in their center [23]. To compare the theoretical distribution to observed depths, it is further 93 useful to used a normalized maximum depth $y = z/\sqrt{2}L^H$, and define the probability distribution in 94 terms of y, i.e. p(y) (which can later be recast in terms of z). Thus in terms of maximum depth and 95 area, normalized maximum depth is: $y = z/\sqrt{2}(\ell\sqrt{a})^H$. This allows an overall test of goodness-of-fit 96 which is needed because of difficulty finding large numbers of lakes with identical surface areas on real 97 landscapes. 98

⁹⁹ Empirical Test

We tested goodness-of-fit of normalized maximum depths to the theoretical distribution using the publicly available LAGOS database, which includes maximum depths for N = 8164 lakes in the northeastern United States [20, 21] (Accessed 28 August 2021). We selected this database because it has been extensively documented and has previously been utilized for developing and testing predictive models for lake maximum depth [21]. Maximum depth spans more than three orders of magnitude among these lakes, from z = 0.1m to z = 135m while surface area spans from a = 4ha to a = 989ha. This database does not distinguish between lakes and reservoirs.

We compare the theoretical and empirical depth distributions using the Kolmogorov-Smirnov statistic 107 $D = \max(P(y) - E(y))$, where P (E) is the theoretical (empirical) cumulative distribution function 108 for y [28]. For a given value of H and ℓ , we calculate y for each lake, then E(y) from the empirical 109 distribution of y values, then p(y) based on that H value, then P(y) by integrating p(y), then finally 110 D by comparing E(y) and P(y). First, we set H = 0.4 as an estimated value for Earth's topography 111 [17, 6], which has also been shown to capture global scale relationships between mean depth and 112 surface area [2]. We then select ℓ by systematically varying ℓ and repeating this procedure until D is 113 minimized. We also test what the best-fit value of H is by systematically varying both H and ℓ until 114 D is minimized. Finally, because small lakes are often not accurately represented in lake databases, 115 we repeat this process but only considering the upper half of the distribution, i.e. using a modified 116 $D_{alt} = \max(P(y > \text{median}(y)) - E(y > \text{median}(y))))$ [26]. This tests the goodness-of-fit of the larger 117 well mapped lakes to the upper tail of the theoretical distribution. 118

119 **Results**

The theoretical distribution p(z) adequately captures the empirical distribution of maximum depth 120 when H = 0.4 (Figure 2; D = 0.033). The distribution is right-skewed and unimodal, with few lakes 121 having a $y \approx 0$, many lakes having a $y \approx 0.5$, and a heavy tail of relatively deep lakes having a 122 $y \in (1,4)$. The best-fitting ℓ value was 0.17, substantially lower than the theoretical limit of $\sqrt{2/\pi}$; 123 this presumably reflects the extent to which these lakes are not circular and their maximum depths are 124 not located in their centers. Furthermore, when H is left as a free parameter rather than externally 125 set by topographic considerations, the best-fitting H = 0.38 is negligibly different to H = 0.4 (Figure 126 A1; D = 0.027; ℓ also changes only slightly, from 0.17 to 0.18). 127

This distribution p(y) can be recast as a distribution for maximum depth z for lakes of a given area a(Figure 3); for a lake with a = 1 ha this corresponds to a median maximum depth of 3.9m and a 95% confidence interval of (0.4m, 10.5m). This wide range of predicted depths arises naturally from random topographic variations, and is consistent with the ranges of variability observed in empirical datasets (Figure 3). In this context of a particular a value, the difference between the theoretical and empirical distributions can be most easily understood by comparing the percentiles of each, which underscores



Figure 2: Theoretical probability distribution versus histogram of normalized maximum depth (y) values for 8164 lakes. The Hurst exponent is fixed at H = 0.4 and the area-lengthscale coefficient ℓ (= 0.17) is left as a free parameter. D = 0.033 is the Kolmogorov-Smirnov statistic.

that the correspondence is indeed quite good. If we treat the empirical and theoretical y distributions as if they were maximum depth distributions for lakes with a = 1ha, the 1st-91st percentiles of the theoretical distribution for maximum depth are all within 8cm of those of the empirical distribution, which is remarkable given that errors in predicting maximum depth are primarily material for shallow and medium depth lakes. Lakes in the 92-99th percentile have larger absolute deviations from their expected values (10-56cm), but as these expected values are larger these correspond to small relative deviations (5-17%).

141 Discussion

Scaling relationships provide simple rules for understanding hydrographic patterns at regional and global scales. Our study contributes to this understanding by describing the relationship between surface area and maximum depth. Prior work has primarily focused on developing empirical relationships with multiple linear regression or similar methods (e.g., [27, 12, 21]), whereas our study provides a rigorous theoretical perspective which is accurate and generalizable among regions. Lake area and the Hurst coefficient are the key factors controlling lake depth, with a large stochastic component. These



Figure 3: Lake surface areas a and maximum depths z, with percentiles of the probability distribution for z (from figure 2, for a given a) superimposed.

factors can be measured without bathymetric surveys, and hence our results can be applied to existing
hydrographic and topographic data sets to estimate characteristics of lake across broad geographic
extents including littoral habitat size, carbon burial, or greenhouse gas evasion [22, 16, 1].

Theoretical lake morphometry, including our study, primarily derives from the assumption that Earth's 151 topography approximates a fractal Brownian surface [11, 2, 25]. This assumption is imperfect because 152 1) it implies that the landscape is a static surface whereas real landscapes are dynamic and evolve over 153 time, 2) it implies scale-invariance whereas real landscapes are shaped by scale dependent processes, 154 and 3) it implies that lakes are formed by flooding preexisting landscape depressions, which is often 155 not the case [11, 29]. Within this context, linear regression analyses based on landscape characteristics 156 essentially are trying to leverage the differences between real and random surfaces to improve maximum 157 depth predictions. However, the low explanatory power of such analyses, and the general consistency 158 between theoretical predictions and empirical patterns, indicates that the imperfect assumptions of 159 the fractal Brownian motion do not materially detract from their application to lake morphometry. 160 Hence, although somewhat stylized, the fractional Brownian motion is an effective starting point for 161 developing predictions about the global characteristics of lakes. 162

In our analysis, the number of medium-depth lakes is slightly overestimated and number of small 163 lakes slightly underestimated by the theoretical distribution, but the overall fit is remarkably good. 164 However, whether H is set by external topographic considerations (H = 0.4) or estimated from the 165 lake y distribution (H = 0.38), there are some discrepancies between the theoretical depth distribution 166 and observations is not perfect. Specifically, at the lower end of the distribution, with the theory 167 slightly overestimating the number of medium-depth lakes $(y \in (1,2))$, underestimating the number of 168 shallow lakes $(y \in (\frac{1}{4}, 1))$, and overestimating the number of very shallow lakes $(y < \frac{1}{4})$. The general 169 ontongeny of lakes is decreased depth overtime due to sedimentation [22]. Empirical deviations from 170 the theoretical depth distribution are consistent with this ontogeny, specifically very shallow lakes 171 are rarer than expected, probably because they have transitioned from lake to wetland or terrestrial 172 ecosystem. Additionally, the accentuated peak for shallow lakes is consistent with patterns expected 173 due sedimentation. There are several other factors that can contribute to discrepancies between theory 174 and empirical patterns, primarily related to the collection of bathymetric data. First, large lake 175 databases typically contain samples strongly biased to certain lake characteristics, and the values in 176 the data we analyzed may not be completely representative of the true maximum depth distribution [26, 177 22]. Second, maximum depths are often only reported to one decimal place, which can lead to significant 178 rounding errors for shallow lakes. This is clearly visible in Figure 3, where there is a regular patterning 179 of shallow maximum depths, but not deep maximum depths (i.e. equal spacing among points along the 180 ordinate). Finally, very shallow lakes may be mis-classified as wetlands and therefore not included in 181 bathymetric databases. Collectively, these factors most strongly impact small and shallow lakes, and 182 this observation is consistent with the excellent fit between theoretical and empirical distributions for 183 deep lakes and somewhat weaker fit across the whole depth distribution. Specifically, when we fit only 184 the upper-half of the y distribution (i.e. using D_{alt} above; Figure A2); the goodness-of-fit was excellent 185 $(D_{alt} = 0.012)$. Advancing global limnology relies on developing robust probability distributions 186 for lake characteristics [7]. The variety of factors causing discrepancies between the theoretical and 187 observed distributions highlight the challenges faced in characterizing these distributions, even in ideal 188 cases like this where theory provides clear guidance on the appropriate distribution. 189

Our approach to understanding the relationship between maximum depth and surface area is useful for predicting characteristics of collections of lakes, but not individual lakes. The only accurate way to measure morphometry for individual lakes remains the detailed bathymetric survey. However, many urgent questions in the aquatic sciences are framed in a global perspective where the characteristics

of large populations are of specific interest [7, 8]. Our approach is well suited for application to these 194 questions. In particular, the characteristics of small numbers of lakes are typically up-scaled based 195 on abundance-area distributions to estimate population level characteristics [7, 24]. This poses a 196 problem for ecosystem characteristics that are closely tied to maximum depth but not surface area. 197 Diffusive methane flux from temperate lakes to the atmosphere is one example of these characteristics 198 [16]. Methane is a potent greenhouse gas and estimating evasion from lakes to the atmosphere is a 199 priority for understanding contributions of lakes to the global carbon cycle. A recent synthesis of 200 diffusive flux measurements for temperate lakes revealed a statistically significant inverse relationship 201 with maximum depth (p < 0.001), but no significant relationship with surface area (p > 0.05) [16]. We 202 applied the empirical relationship for diffusive methane flux from [16] to the full database of lakes 203 from 17 northeastern United States for which surface areas are available (N = 141, 265) by randomly 204 simulating maximum depths for these according to the distribution above (with H = 0.4, and $\ell = 0.17$ 205 estimated from the 8,164 lakes from these same 17 states for which maximum depth measurements were 206 available; Figure 2) and then calculating diffusive methane flux for each lake as a function of surface 207 area and simulated maximum depth. We repeat this process many times to estimate uncertainty due 208 to the randomness of simulated maximum depths. This results in an estimated 0.5 ± 0.04 Tg CH₄/year 209 from these lakes. This is a substantially different estimate than if one just takes an average rate 210 per unit area across all lakes, because the relationship between methane flux and maximum depth is 211 nonlinear [16] and this average will be dominated by lakes with large surfaces, deep maximum depths, 212 and low rates of methane flux per unit area; one may then substantially underestimate overall methane 213 flux as a result. For instance, multiplying the average of 0.9 g $CH_4/m^2/year$ from [16] by the total 214 surface area of the 141,265 lakes above results in an estimate of 0.03 Tg CH_4 /year. Hence, our results 215 represent an important methodological advancement for up-scaling lake characteristics correlated with 216 depth but not surface area. For illustrative purposes (as the parameter ℓ has been fit here to lakes 217 from the northeastern United States rather than lakes globally), applying the relationship from [16] 218 globally to the lake areas from the Hydrolakes database (https://hydrosheds.org) [18] via simulating 219 maximum depths as we describe here results in a global diffusive methane flux of 27 ± 3 Tg CH₄/year. 220 Matlab code demonstrating the application of the maximum depth distribution for up-scaling is in the 221 supplemental materials. 222

Our study highlights the far reaching influence of the Hurst coefficient on global-scale lake characteristics. Specifically, the differences between horizontal and vertical scaling described by the Hurst

coefficient underlie differences in characteristics between large and small lakes - small lakes are typi-225 cally deeper relative to their surface area compared to large lakes which has implications for energy 226 and carbon budgets across the lake size spectra [1]. Additionally, the Hurst coefficient is involved 227 in most other lake scaling relationships, including for abundance, perimeter, volume, and hydrologic 228 connectivity [11, 25, 2, 23]. Despite this, empirical measurements of the Hurst coefficient for Earth's 229 topography and bathymetry are relatively rare and highly variable (H = 0.4-0.7)[11, 10, 31]. Devel-230 oping such measurements should be an important priority for advancing global scale understanding of 231 lakes. These measurements could explain variations in scaling relationships among regions, as well as 232 improve the precision of predictions by reducing uncertainty in parameterization. 233

800 Histogram of Lakes • 700 Theoretical Prediction H = 0.38600 $\ell = 0.18$ D = 0.027500 Counts 400 300 200 100 0 1 1.5 $\mathbf{2}$ 2.53 3.50.50 4 $y \left(= z/\sqrt{2}(\ell\sqrt{a})^H \right)$

²³⁴ Appendix: alternate versions of Figure 2

Figure A1: As in Figure 2, but with H left as a free parameter. The best-fit Hurst exponent is H = 0.38 (0.4 in Figure 2) and the best-fit $\ell = 0.18$ (0.17 in Figure 2); D reduces to 0.027 (0.033 in Figure 2).



Figure A2: As in Figure 2, but with H left as a free parameter and D is only evaluated for y values above the median (i.e. D_{alt} is used). The best-fit Hurst exponent is H = 0.46 (0.4 in Figure 2) and the best-fit $\ell = 0.11$ (0.17 in Figure 2) while the fit to the upper half of the distribution is much improved $(D_{alt} = 0.012 \text{ as opposed to } D = 0.033 - \text{albeit over the full distribution - in Figure 2}).$

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