

1        **Distinct roles of cyclones and anticyclones in setting the midwinter**  
2        **minimum of the North Pacific eddy activity. Part II: Eulerian eddy**  
3        **statistics and energetics**

4  
5                    Satoru Okajima,<sup>a</sup> Hisashi Nakamura,<sup>a</sup> Yohai Kaspi,<sup>b</sup>

6                    <sup>a</sup> *Research Center for Advanced Science and Technology, The University of Tokyo, Tokyo, Japan*

7                    <sup>b</sup> *Department of Earth and Planetary Sciences, Weizmann Institute of Science, Rehovot, Israel*

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13       *Corresponding author:* Satoru Okajima, [okajima@atmos.rcast.u-tokyo.ac.jp](mailto:okajima@atmos.rcast.u-tokyo.ac.jp)

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## ABSTRACT

The characteristics and dynamics of midlatitude storm-tracks have been long investigated. Nevertheless, our understandings of the storm-tracks, especially the midwinter minimum of the North Pacific storm-track activity, are still limited, partly because Eulerian eddy statistics are incapable of separating cyclonic and anticyclonic contributions. Here we investigate the detailed seasonal evolution of the contributions from cyclonic and anticyclonic vortices, identified with local wind curvature, to the Eulerian eddy statistics, eddy feedback forcing onto the climatological-mean westerly jet, and the energetics associated with the North Pacific storm-track. We demonstrate that the midwinter minimum of the net energy conversion/generation rate normalized by the eddy total energy over the entire North Pacific is more distinct for anticyclonic vortices than for cyclonic vortices, to which the midwinter minimum of the probability of anticyclonic vortices contributes. The main factors include the reduction of the baroclinic energy conversion into midwinter and the reduction of the net energy outflux into spring. To long-term modulations of the MWM of the NP storm-track activity, both cyclonic and anticyclonic vortices contribute constructively. Together with our companion paper, the present study reveals the primal importance of migratory anticyclones for the midwinter minimum of the North Pacific storm-track activity. We posit a viewpoint of considering roles of both migratory cyclones and anticyclones for the storm-track dynamics, the latter of which has been long disregarded unconsciously.

## 34 **1. Introduction**

35 Midlatitude transient eddies that give rise to day-to-day weather variability are one of the  
36 rudimentary components of the Earth’s climate system. Those migratory cyclones and  
37 anticyclones interact with climatological-mean states and low-frequency, quasi-stationary  
38 anomalies in which they are embedded. Passing over a given midlatitude location recurrently,  
39 they locally account for a major fraction of high-frequency sub-weekly fluctuations of a  
40 given atmospheric variable. Landsberg et al. (1959) found a variability whose period is  
41 between 5-7 days on top of that of 15-25 days from a station temperature data. Based on a 10-  
42 year atmospheric analysis data over the entire extratropical Northern Hemisphere, Blackmon  
43 (1976) decomposed the variance of 500-hPa height into high-pass, band-pass (with periods of  
44 2–6 days), and low-pass filtered components. His result suggested that regions of large band-  
45 pass variance correspond to those of frequent cyclone passage over the North Atlantic (NA)  
46 and North Pacific (NP) basins. Blackmon et al. (1977) further pointed out that those regions  
47 are characterized by prominent poleward eddy heat flux based on the band-pass-filtered  
48 fluctuations in temperature and meridional wind velocity at 850-hPa, referring to them as  
49 “storm-tracks”. A zonally-elongated region of pronounced poleward heat transport associated  
50 with sub-weekly transient eddies, a measure of their baroclinic development, is also found  
51 located over the South Indian Ocean (Nakamura and Shimpo 2004). Those three major storm-  
52 tracks are collocated with the lower-tropospheric eddy-driven westerly jets and associated  
53 lower-tropospheric baroclinic zones (Nakamura et al. 2004), and they are maintained under  
54 the effective restoration by major oceanic frontal zones (Hotta and Nakamura 2011).

55 The Eulerian statistics utilize gridded atmospheric data, including operational analysis  
56 and global reanalysis data, which have been available since the mid-1970s and mid-1990s,  
57 respectively. The Eulerian approach became common in the 1980s. Since then,  
58 climatological-mean seasonality, and interannual and decadal-scale variations have been  
59 investigated through Eulerian statistics (e.g., Chang et al. 2002). For Eulerian eddy statistics,  
60 temporal digital filters are commonly used to extract sub-weekly fluctuations associated with  
61 extratropical high-frequency, synoptic-scale disturbances. Eulerian statistics are compatible  
62 with the quantitative analyses and dynamical diagnostics, including the “Lorenz energy  
63 cycle”, in which exchanges are evaluated between kinetic and available potential energies  
64 and between the zonal-mean averaged and eddy components (Lorenz 1955). For example, the  
65 formation and maintenance mechanisms for the mean state and storm-track activity can be

66 examined through energy budget analysis (e.g., Orlanski and Katzfey 1991; Chang et al.  
67 2002). The (extended) Eliassen-Palm (E-P) flux, which represents the propagation of Rossby  
68 wave packets and associated translation of wave-activity pseudo-momentum, is useful for  
69 assessing the wave-mean flow interaction (Hoskins et al. 1983; Trenberth 1986; Plumb  
70 1986).

71 Expanding Lee and Kim (2003), Nakamura et al. (2004) illustrated the maintenance  
72 mechanisms for the westerly jets and storm-tracks, with particular emphasis on the lower-  
73 boundary condition. They contrasted a regime characterized by a well-defined eddy-driven  
74 subpolar jet separated from a subtropical jet with another regime by a stronger subtropical jet.  
75 In the former regime, a storm-track and an associated eddy-driven westerly jet extend from  
76 the surface to the upper troposphere, anchored along a subarctic oceanic frontal zone, as  
77 typically observed in the summertime Southern Hemisphere, the NP shoulder seasons  
78 (autumn and spring) and the NA. Transient eddies systematically transport westerly  
79 momentum poleward and downward from a subtropical jet to maintain the deep eddy-driven  
80 jet. In the latter regime, as typically observed in the wintertime Southern Hemisphere, the  
81 upper-tropospheric storm-track is trapped into the subtropical jet core, while the low-level  
82 storm-track forms along an oceanic frontal zone. They argued that the midwinter NP is in-  
83 between, characterized by a hybrid jet through a merger of subpolar and subtropical jets.

84 Nakamura (1992) found that the NP storm-track activity, measured by RMS of the band-  
85 pass filtered upper-tropospheric geopotential height fluctuations, exhibits a clear minimum in  
86 midwinter. It is inconsistent with the baroclinic instability theory (Eady 1949), in which the  
87 maximum growth rate of eddies is proportional to the vertical shear of the zonal winds, given  
88 that the westerly jet speed over the NP is clearly maximized in midwinter. It is in sharp  
89 contrast to the midwinter maximum of the NA storm-track activity, which is compatible with  
90 the theory. This counterintuitive phenomenon has been referred to as “midwinter minimum  
91 (MWM)” or “midwinter suppression” of the NP storm-track activity, for which multiple  
92 mechanisms have been proposed. They include barotropic and baroclinic aspects of eddies,  
93 diabatic heating, an upstream influence (or seeding effect), and structures of the jet streams  
94 (see literatures in Okajima et al. 2022a; hereafter referred to as Part I). To date, however, no  
95 single mechanism has been proposed that can explain the existence of the MWM of the NP  
96 storm-track activity.

97 For a quantitative comparison among the proposed mechanisms, energy budget analysis is  
98 a powerful tool. Zhao and Liang (2019) examined the energetics of migratory eddies along  
99 the NP storm-track in relation to its MWM. However, they evaluated only energy  
100 conversion/generation rates, which is inevitably dependent of eddy amplitude. Their findings  
101 are therefore insufficient for full consideration of the mechanisms for the suppressed eddy  
102 activity. Chang (2001) defined “growth rates” as the local eddy energy conversion/generation  
103 rates for the NP storm-track divided by local eddy kinetic energy (EKE). The normalization  
104 by local EKE is, however, debatable from the basin-scale characteristic of the NP storm-track  
105 with obvious downstream translation of eddy energy. Schemm and Rivière (2019) evaluated  
106 “efficiency” of some of the energetic terms, which is independent of eddy amplitude, but  
107 their analysis is rather limited to particular pressure levels and variables. It is therefore  
108 necessary to perform more comprehensive energetics of eddies along the NP storm-track for  
109 clarifying the mechanisms for the MWM. More recently, Okajima et al. (2022; hereafter  
110 ONK22) comprehensively evaluated the detailed seasonal evolution of energy  
111 conversion/generation rates associated with transient eddies along the NP storm-track,  
112 normalized by the eddy total energy. Giving a perspective of the MWM of the NP storm-  
113 track activity that encompasses various mechanisms proposed by previous studies, they found  
114 that the net normalized energy conversion/generation rate is indeed suppressed in midwinter.  
115 Their results demonstrated that conversion of available potential energy (APE) from the  
116 climatological-mean state to eddies plays a substantial role especially in the reduction of the  
117 normalized total energy conversion/generation rate in midwinter from its early-winter peak,  
118 whereas that of the energy influx from the upstream is particularly important for the spring  
119 recovery.

120 Over decades, considerable effort has been devoted to investigating the characteristics  
121 and dynamics of midlatitude storm-tracks. Nevertheless, our understandings of storm-tracks  
122 are still limited, including the MWM of the NO storm-track activity. One of the reasons is  
123 that Eulerian statistics, which facilitate a quantitative diagnosis and analysis, assumes  
124 linearity of eddy-associated fluctuations and therefore cannot handle cyclones and  
125 anticyclones separately. Lagrangian tracking, by contrast, enables to analyze them separately  
126 but is not suited for a quantitative analysis, especially for eddy-mean flow interactions. These  
127 two methods are intrinsically complementary, but they have been applied separately in most  
128 of the previous studies. Some studies conducted analyses of both Eulerian statistics and  
129 Lagrangian tracking, but just in parallel (Chang et al. 2002; Hoskins et al. 2019a, b). In fact, a

130 region of pronounced high-frequency SLP variance in the Eulerian framework does not  
131 necessarily coincide with a region of frequent cyclone tracks in the Lagrangian framework.  
132 Wallace et al. (1988) documented the behavior of high-frequency fluctuations in geopotential  
133 height, to reveal that cyclones tend to move northeastward to the east of Japan, not along the  
134 zonally-elongated “storm-track” identified by Blackmon (1976). Comparison between  
135 Eulerian variance and feature-tracking statistics by Penny et al. (2010) was immensely  
136 simplified, assuming eddies as a series of identical non-overlapping sine-shaped pulses.  
137 Schemm and Rivière (2019) attempted to evaluate the APE conversion only around surface  
138 cyclones by accumulating the local APE conversion within circles of a given threshold radius  
139 centered at a cyclone center, or the outermost closed SLP contour. Nevertheless, surface  
140 cyclones and associated cyclonic vortices are neither circular nor vertical. Additionally, they  
141 did not assess the contribution from anticyclones.

142 Recently, Okajima et al. (2021; hereafter ONK21) proposed a novel method to identify  
143 three-dimensional domains of individual cyclonic and anticyclonic vortices based on local  
144 curvature, and to evaluate separate contributions from cyclonic and anticyclonic vortices to  
145 Eulerian eddy statistics and atmospheric energetics. Local curvature and related quantities  
146 have been utilized for tracking mid- to upper-tropospheric troughs (Lefevre and Nielsen-  
147 Gammon 1995; Schemm et al. 2020), but they have not been used to separate and reconstruct  
148 eddy Eulerian statistics. Combined with the framework for the comprehensive energetics of  
149 storm-tracks independent of the eddy amplitude by ONK22, the “hybrid perspective” method  
150 will effectively illustrate the seasonal evolution of the separate contributions from cyclonic  
151 and anticyclonic vortices to the Eulerian eddy statistics and atmospheric energetics.

152 The present study thus aims to investigate the detailed seasonal evolution of the separate  
153 contributions from cyclonic and anticyclonic vortices to Eulerian eddy statistics, eddy  
154 feedback forcing onto the climatological-mean westerly jet, and energetics associated with  
155 the NP storm-track. In Part I, seasonal evolution of the frequency of surface cyclones and  
156 anticyclones are investigated based on a Lagrangian tracking, focusing especially on the  
157 importance of anticyclones traveling from the Japan Sea into the NP.

158 This paper (Part II) continues in section 2 with a description of data used and analysis  
159 methods. Sections 3 and 4 delineate seasonality of cyclonic and anticyclonic contributions to  
160 eddy Eulerian statistics and eddy feedback forcing on the background westerlies. Section 5

161 investigates separated energetics for the NP storm-track. Section 6 describes the long-term  
162 modulations in the energetics briefly. Section 7 offers a summary and discussions.

## 163 2. Data and analysis methods

### 164 *a. observational data*

165 Same as in Part I, this study utilizes atmospheric variables, including geopotential height,  
166 temperature, wind velocities, and diabatic heating rates in pressure coordinates as well as sea-  
167 level pressure (SLP), from the Japanese 55-year reanalysis (JRA-55; Kobayashi et al. 2015;  
168 Harada et al. 2016) for the period 1958-2017. Those variables are available on a  $1.25^\circ \times 1.25^\circ$   
169 grid.

170 At each grid point, fluctuations of a given variable with synoptic-scale transient eddies  
171 have been extracted from the 6-hourly atmospheric reanalysis as its deviations from their  
172 low-pass-filtered fields through a Lanczos filter with a 121-point window and a cutoff period  
173 of 8 days. Climatological-mean fields are calculated from 31-day running mean fields.

174

### 175 *b. Separation of cyclonic and anticyclonic contributions to eddy Eulerian statistics*

176 Following the methodology developed by ONK21, climatological-mean eddy Eulerian  
177 statistics are calculated separately for cyclonic and anticyclonic contributions. Local  
178 curvature  $\kappa_2$  is calculated at a given vertical level instantaneously from zonal and meridional  
179 winds ( $u, v$ ) as

$$180 \quad \kappa_2 \equiv \frac{1}{R_S} = \frac{1}{V^3} (-uvu_x + u^2v_x - v^2u_y + uvv_y), \quad (1)$$

181 where  $R_S$  denotes the curvature radius,  $V$  scalar wind speed and a subscript zonal or  
182 meridional derivative. This can be derived also from the definition of two-dimensional  
183 curvature with an implicit curve of geostrophic streamfunction (Goldman 2005). Note that  
184 this method does not require any temporal filtering to determine the shape of vortices, and  
185 therefore near-surface curvature corresponds directly to counterclockwise (clockwise)  
186 rotations associated with a cyclone (anticyclone) over the NH.

187 Separate contributions from cyclonic and anticyclonic vortices (or eddies) to Eulerian  
188 statistics is evaluated by accumulating instantaneous contributions only at grid points where  
189 cyclonic or anticyclonic curvature is observed. In the following, the threshold curvature for  
190 cyclonic and anticyclonic rotations are set to be zero so as not to miss any circulation on the  
191 fringe of pressure troughs and ridges, unless otherwise specified. Note that the conclusions



192 are qualitatively similar with the requirement for vortices to accompany corresponding  
 193 surface features (see Appendix).

194

195 *c. Energetics*

196 The framework of the energetics for transient eddies along storm-tracks basically follows  
 197 ONK22. Climatological-mean eddy APE (EAPE) and eddy kinetic energy (EKE) associated  
 198 with sub-weekly disturbances are defined as:

$$199 \quad \text{EAPE} = \frac{R}{pS_p} \left( \frac{(T'^2)_c}{2} \right), \text{EKE} = \frac{(u'^2+v'^2)_c}{2}, \quad (2)$$

200 where primes denote sub-weekly fluctuations, subscripts “c” the climatological-mean fields,  
 201 and  $S_p (\equiv -\bar{T}_c \partial \ln \bar{\theta}_c / \partial p)$  a stability parameter. Overbars denote horizontally-averaged  
 202 quantities over a specific domain. In (2),  $u$ ,  $v$ ,  $T$ , and  $R$  represent zonal and meridional wind  
 203 components, temperature, and the gas constant for dry air, respectively.

204 Energy conversion/generation rates are defined as in (3), assuming quasi-geostrophic  
 205 scaling. Additionally, time tendency of horizontally-averaged climatological-mean static  
 206 stability is omitted. In the following,  $CK$  denotes the barotropic energy conversion (or KE  
 207 conversion),  $CP$  the baroclinic energy conversion (or APE conversion), and  $CQ$  the APE  
 208 generation through diabatic processes:

$$209 \quad CK = \frac{(v'^2-u'^2)_c}{2} \left( \frac{du_c}{dx} - \frac{dv_c}{dy} \right) - (u'v')_c \left( \frac{du_c}{dy} + \frac{dv_c}{dx} \right) \quad (3a)$$

$$210 \quad CP = \frac{R}{pS_p} \left( -(u'T')_c \frac{dT_c}{dx} - (v'T')_c \frac{dT_c}{dy} \right) \quad (3b)$$

$$211 \quad CQ = \frac{R}{pS_p} (Q'T')_c \quad (3c)$$

$$212 \quad ET = -\frac{R}{p} (\omega'T')_c, \quad (3d)$$

213 where  $Q$  signifies temperature tendency due to diabatic processes and  $\omega$  pressure velocity.

214 In this study, the energy conversion/generation terms are integrated vertically from the  
 215 surface to the 100-hPa level and then integrated horizontally over the NP domain  
 216 [130°E–130°W, 20–65°N]. Energy influx or outflux through the lateral boundaries of the

217 domain ( $EF$ ) needs to be evaluated for the energetics within the domain. The flux can be  
 218 expressed as:

$$219 \quad EF = \int_{p_s}^{p_{top}} \{(\Phi' \mathbf{v}')_C + (EAPE + EKE)\mathbf{v}_c\} dp, \quad (4)$$

220 where  $\Phi$  is geopotential,  $p_s$  denotes surface pressure, and  $p_{top}$  is set to be 100 hPa.

221 As in ONK22, the “normalized rate” of each of the energy conversion, generation and  
 222 flux terms is evaluated by dividing the term by the total eddy energy defined as the sum of  
 223 EAPE and EKE:

$$224 \quad \text{Normalized rate} \equiv \frac{\langle \text{Energy conversion/generation term} \rangle}{\langle \text{EAPE+EKE} \rangle}. \quad (5)$$

225 This “normalized rate of conversion (or generation)” represents a replenishing rate in unit of  
 226 [ $\text{day}^{-1}$ ], and its reciprocal thus gives a time scale of how long it would take to replenish the  
 227 total eddy energy within the storm-track domain solely by a given conversion/generation  
 228 process. Hereafter, to improve visibility, we refer to the normalized rate simply as “ $\lambda$ ” and a  
 229 subscript as listed in Table 1 is added in referring to a specific conversion/generation process  
 230 as in ONK22<sup>1</sup>. Because the energy conversion/generation rate, by definition, depends on  
 231 eddy amplitude, it is essential for the energetics to evaluate it as its normalized rate, which is  
 232 independent of eddy amplitude. In our definition negative values of  $\lambda$  mean that eddies are  
 233 giving up energy to the time-mean circulation.

Symbol	Meaning
$\lambda_{CK}$	Rate of CK (barotropic energy conversion) normalized by (EKE+EAPE)
$\lambda_{CP}$	Rate of CP (baroclinic energy conversion) normalized by (EKE+EAPE)
$\lambda_{CQ}$	Rate of CQ (diabatic energy generation) normalized by (EKE+EAPE)
$\lambda_{EF}$	Rate of EF (net energy influx/outflux) normalized by (EKE+EAPE)

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<sup>1</sup> We also calculate barotropic and baroclinic energy conversion rates from low-frequency variability to sub-weekly eddies (“LF” terms in ONK22). It was found to be relatively small and thus omitted in the present study.

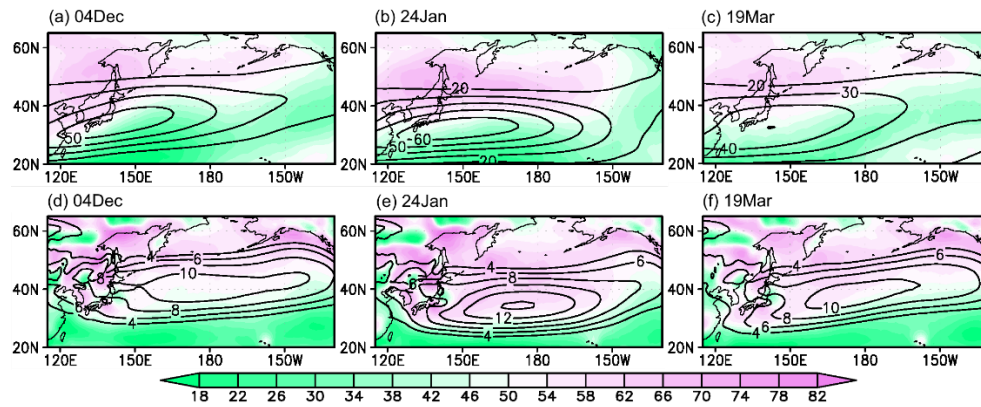
$\lambda_{Tot}$	Rate of total energy conversion/generation rate (sum of the above terms) normalized by (EKE+EAPE)
$\lambda_{ET}$	Rate of ET (energy transfer from EKE to EAPE) normalized by (EKE+EAPE)

234 Table 1. List of the symbols for the normalized rates used in this study. All quantities are  
235 expressed in the unit of [day<sup>-1</sup>].

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### 237 **3. Cyclonic and anticyclonic contributions to Eulerian statistics and their** 238 **seasonality**

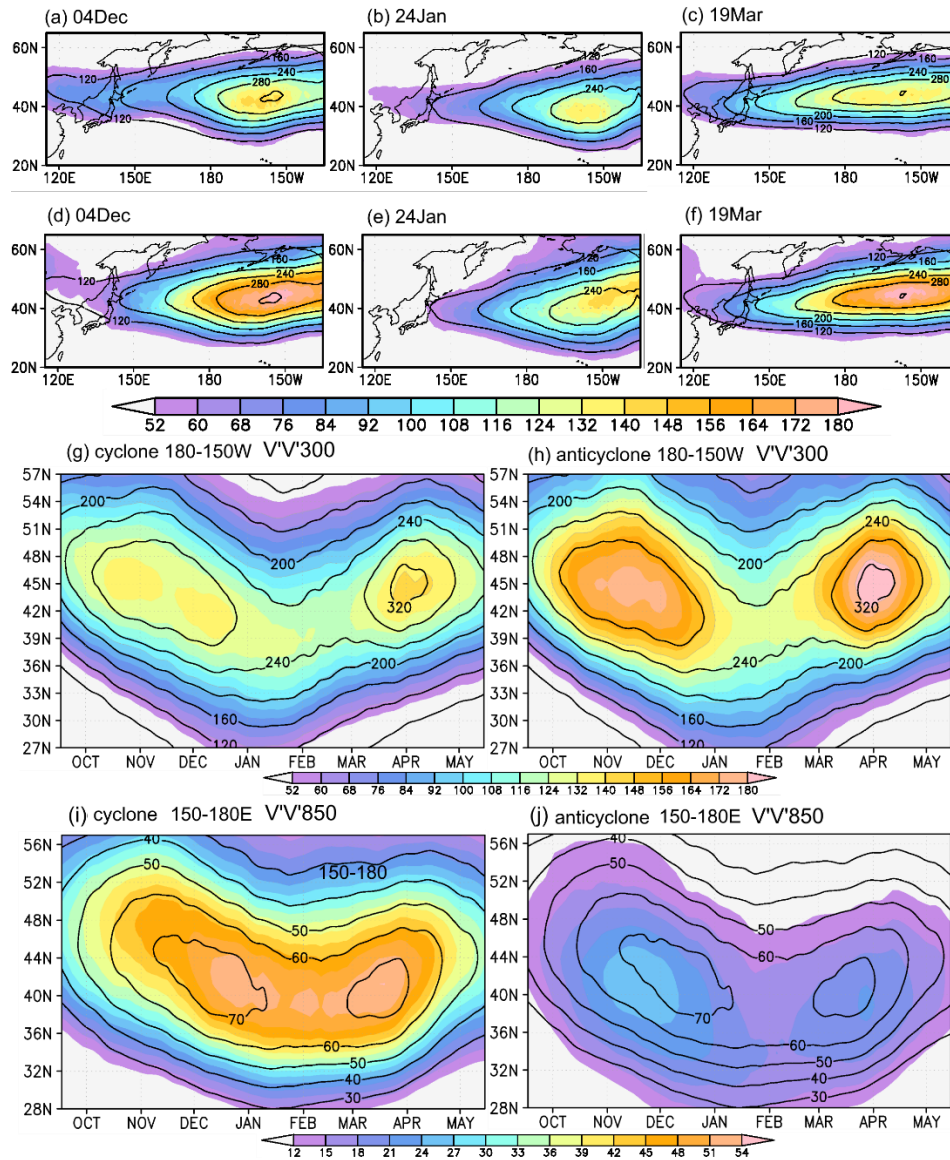
239 First, we examine the seasonal evolution of climatological-mean probability distributions  
240 of cyclonic and anticyclonic vortices over the NP (Fig. 1). The three calendar days indicated  
241 in Fig. 1 correspond to the first peak, MWM and second peak of the NP storm-track activity  
242 (Nakamura 1992; ONK22). As pointed out by ONK21, probability of upper-level cyclonic  
243 vortices is higher poleward of the jet axis in midwinter and vice versa for anticyclonic  
244 vortices (Fig. 1b). This may be attributed to the effect of “trough thinning and broadening” by  
245 meridional shear of the jet (Thorncroft et al. 1993), given that those vortices are identified by  
246 local curvature free from shear vorticity. As shown in Fig. 1e, probability of lower-  
247 tropospheric cyclonic vortices (and anticyclonic counterpart as the local residual) is  
248 consistent with the result by feature tracking at the surface (Part I). In the shoulder seasons,  
249 the upper- and lower-tropospheric probability distributions of vortices basically follow the  
250 poleward displacement of the westerly jet as its seasonal evolution (Fig. 1). Midwinter  
251 equatorward expansion of high probability of cyclonic vortices, equivalently equatorward  
252 retreat of a domain of frequent anticyclonic vortices, also corresponds to those of surface  
253 migratory cyclones and anticyclones, the latter of which is largely regulated by anticyclones  
254 from the Japan Sea (Part I).



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 256 Fig. 1. Climatological probability distributions of cyclonic vortices (colors; with positive  
 257 curvature) at the (a-c) 300-hPa and (d-f) 850-hPa levels, for the indicated calendar days as  
 258 indicated. The probability of anticyclonic vortices can be obtained as the local residual. Black  
 259 contours indicate climatological-mean (a-c)  $U_{300}$  and (d-f)  $U_{850}$  (m/s).

260  
 261 Figure 2 shows climatological-mean distributions of the variance of sub-weekly  
 262 fluctuations of 300-hPa meridional wind ( $V'V'_{300}$ ) evaluated as the separate contributions  
 263 from cyclonic and anticyclonic vortices. For both polarities,  $V'V'_{300}$  maximizes in the eastern  
 264 NP as in the total variance (black contours in Fig. 2). While the maximum value of the  
 265 cyclonic  $V'V'_{300}$  remains largely unchanged from early winter into spring (Figs. 2a-c),  
 266 anticyclonic  $V'V'_{300}$  exhibits its distinct MWM. It is stronger than its cyclonic counterpart in  
 267 the shoulder seasons especially over the central and eastern NP (Figs. 2d-f). These  
 268 characteristics are clearly seen in latitude-season sections (Figs. 2g-h), indicative of the  
 269 primary importance of anticyclonic vortices for the distinct MWM between the dual peaks in  
 270 the climatological-mean  $V'V'_{300}$ .

271 Conversely, cyclonic  $V'V'_{850}$  is much stronger than its anticyclonic counterpart (Figs. 2i-  
 272 j), with their maxima located over the western NP (ONK22). The stronger wind fluctuations  
 273 associated with lower-tropospheric cyclonic vortices are congruous with our perception of  
 274 migratory cyclones as “storms”. The contribution from cyclonic vortices to  $V'V'_{850}$  is nearly  
 275 constant, and the peak latitude largely coincides with the climatological axis of  $V'V'_{850}$ . By  
 276 contrast throughout the cold season, its anticyclonic counterpart is stronger in early winter  
 277 than in spring with a moderate MWM in between, and the peak latitude is displaced to the  
 278 south of the climatological-mean axis of  $V'V'_{850}$ . The contribution of anticyclonic vortices to  
 279 the MWM of the total  $V'V'_{850}$  (black contours in Figs. 2i-j) is evident, though in smaller  
 280 amplitude.



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283 Fig. 2. (a-c) Climatological-mean distributions of  $V'V'_{300}$  (color,  $m^2/s^2$ ) for the calendar  
 284 days as indicated, reconstructed only with cyclonic vortices. Contours denote climatological-  
 285 mean total  $V'V'_{300}$  from all vortices ( $m^2/s^2$ ). (d-f) Same as in (a-c), respectively, but for  
 286 anticyclonic vortices. (g) Climatological seasonality in  $V'V'_{300}$  reconstructed only with  
 287 cyclonic vortices averaged for  $180^\circ-150^\circ W$ . Contours denote the corresponding seasonality  
 288 of total  $V'V'_{300}$  from all vortices. (h) Same as in (g), but for anticyclonic vortices. (i, j) Same  
 289 as in (g, h), respectively, but for  $V'V'_{850}$  averaged over  $150^\circ-180^\circ E$ .

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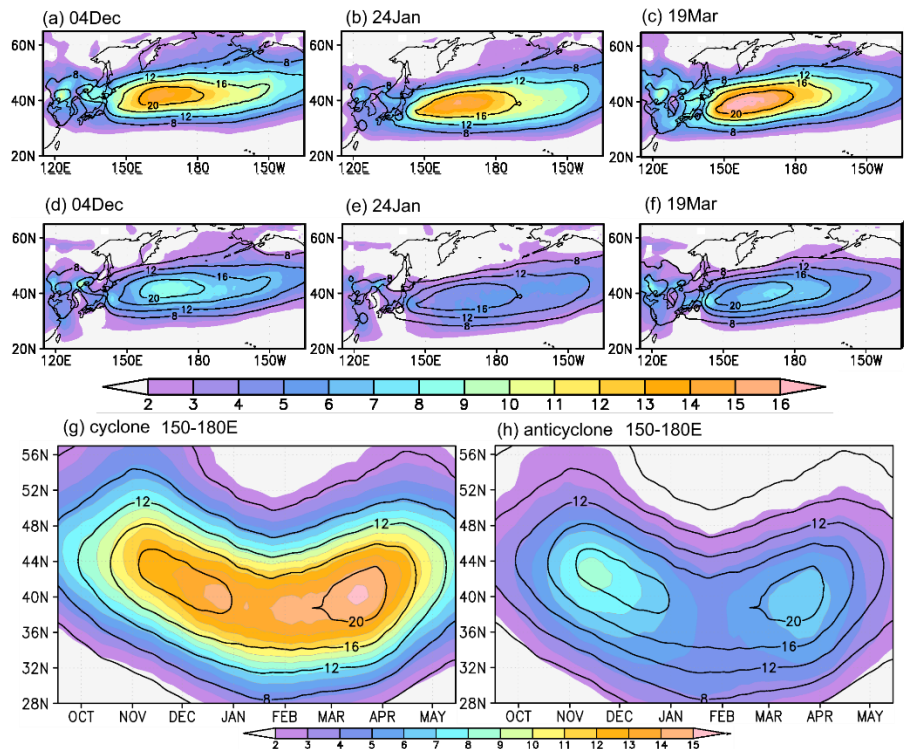
291 Similarly, lower-tropospheric poleward eddy heat flux ( $V'T'_{850}$ ) associated with cyclonic  
 292 vortices is much stronger than its anticyclonic counterpart (Fig. 3). The cyclonic contribution  
 293 is found to maximize in spring with no apparent MWM (Figs. 3a-c). Though weaker in  
 294 general, the anticyclonic contribution, by contrast, exhibits a distinct MWM between its  
 295 primary early-winter peak and its secondary spring peak (Figs. 3d-f). These characteristics  
 are more discernible in the latitude-season sections (Figs. 3g-h), indicative of the

296 indispensable role of anticyclonic vortices in forming the MWM of the climatological-mean  
297  $V'T'_{850}$ .

298 Related to this, seasonality of the upward component of the phase-independent wave-  
299 activity flux defined by Takaya and Nakamura (2001) differs between cyclonic and  
300 anticyclonic vortices, as shown in Fig. 4. This flux component is related to both the down-  
301 gradient eddy heat flux and background static stability. The lower-tropospheric upward wave-  
302 activity flux associated with cyclonic vortices exhibits its distinct midwinter maximum (Fig.  
303 4b), which is contributed to mainly by the reduced static stability (not shown). Meanwhile, its  
304 anticyclonic counterpart maximizes in early winter and gradually weakens into spring (Figs.  
305 4d-f). Just below the westerly jet core ( $\sim 400$  hPa), the upward wave-activity flux exhibits a  
306 MWM for both cyclonic and anticyclonic vortices. In the upper troposphere above the 400-  
307 hPa level, the upward component of the wave-activity flux diminishes, while its eastward  
308 component becomes dominant. Compared with its counterpart associated with cyclonic  
309 vortices, the eastward component associated with anticyclonic vortices is stronger and its  
310 maximum extends farther downstream into the eastern NP. Since the wave-activity flux  
311 plotted in Fig. 4 does not include the term as a product of the phase velocity and wave-  
312 activity pseudo-momentum associated with migratory eddies, our result is indicative of the  
313 greater contribution to the downstream development of transient eddies from anticyclonic  
314 vortices<sup>2</sup>, and the contribution exhibits a MWM.

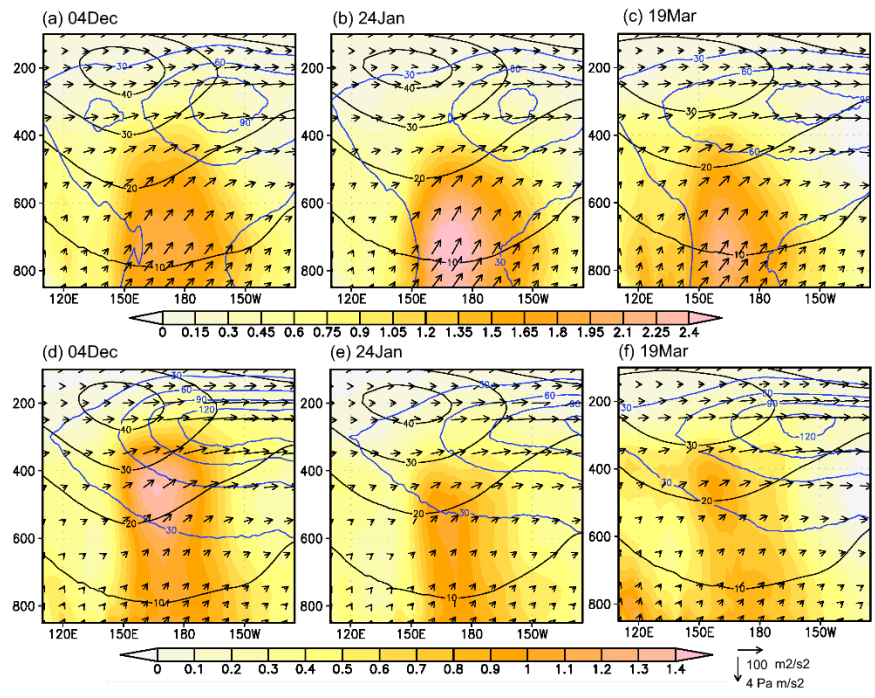
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<sup>2</sup> Horizontal components of the wave-activity flux in this analysis should be interpreted with caution. The flux is evaluated for sub-weekly eddies, whereas the corresponding vortices are identified in unfiltered wind fields, from which the contribution from the NP westerly jet has been removed. Although those eddies and vortices are closely related to each other as seen in Fig. 1, the eastward wave-activity flux may not directly indicate downstream developing vortices.



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Fig. 3. (a-f) Same as in Figs. 2a-f, respectively, but for  $V'T'_{850}$  (K m/s). (g, h) Same as in Figs. 2g-h, respectively, but for  $V'T'_{850}$  (K m/s) averaged over  $150^{\circ}$ – $180^{\circ}$ E.



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324

Fig. 4. (a-c) Zonal sections of climatological-mean wave-activity flux (vectors) formulated by Takaya and Nakamura (2001), which evaluated for sub-weekly disturbances but only from cyclonic vortices for (a) 04Dec, (b) 24Jan, and (c) 19Mar. The flux plotted here does not include the term as a product of the phase velocity and wave-activity pseudo-momentum associated with sub-weekly disturbances. Colors and blue contours indicate its upward ( $\text{Pa m/s}^2$ ) and eastward ( $\text{m}^2/\text{s}^2$ ) components, respectively. Black contours denote

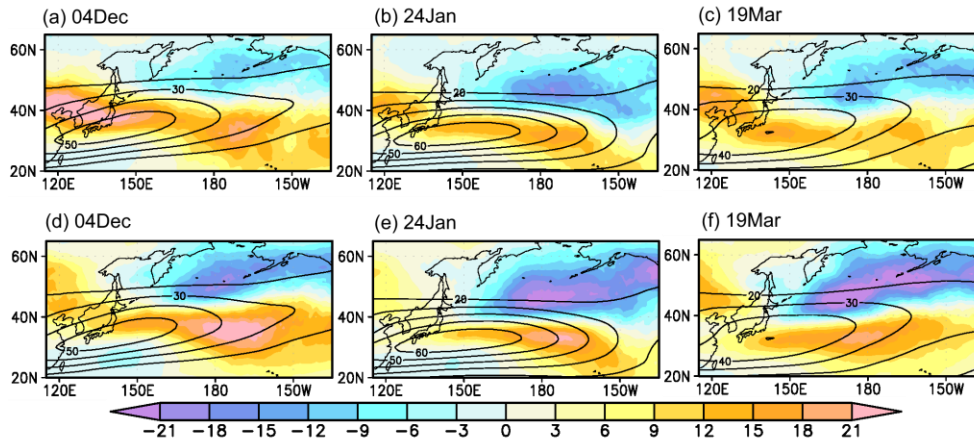
325 climatological-mean westerly wind speed (m/s). Quantities shown are meridionally averaged  
326 in 30°–55°N. (d-f) Same as in (a-c), respectively, but for anticyclonic vortices.

327

328 To investigate westerly momentum transport by cyclonic and anticyclonic vortices, their  
329 contributions to  $U'V'_{300}$  are shown in Fig. 5. The contributions from vortices of both  
330 polarities show a striking meridional contrast characterized by positive (negative)  $U'V'$  to the  
331 south (north) of the climatological-mean NP westerly jet. Zero lines of  $U'V'_{300}$  tend to be  
332 zonally oriented between ~40°N and 45°N, indicating that the westerly acceleration by both  
333 cyclonic and anticyclonic eddies occurs slightly to the north of the NP jet axis. Over East  
334 Asia, positive  $U'V'_{300}$  is clearly shifted poleward of the primary jet axis in the subtropics  
335 without noticeable negative  $U'V'_{300}$ . Obviously, the amplitude of  $U'V'_{300}$ , especially its  
336 negative values north of the jet, is larger for the anticyclonic contribution even in midwinter,  
337 when the magnitude of anticyclonic  $V'V'_{300}$  is comparable to its cyclonic counterpart (Fig. 2).  
338 Consistently, the negative  $U'-V'$  correlation to the north of the jet axis is more prominent for  
339 anticyclonic vortices, while the positive correlation is generally stronger for cyclonic vortices  
340 farther to the south of the jet axis (not shown). For both cyclonic and anticyclonic vortices,  
341 midwinter is the season when the negative correlation north of the jet is most prominent.

342 In the lower troposphere, by contrast, the contribution from cyclonic vortices to the  
343 meridional westerly momentum transport is much larger than that from anticyclonic vortices,  
344 especially on the northern flank of the jet (Fig. 6). This means that cyclonic vortices overall  
345 act to transport westerly momentum effectively into the lower-tropospheric eddy-driven jet  
346 from the Aleutian Low located to the north. Although lower-tropospheric cyclonic vortices  
347 contribute dominantly to the negative  $U'V'$ , the negative  $U'-V'$  correlation to the north of  
348 ~40°N is roughly comparable between cyclonic and anticyclonic vortices (not shown). The  
349 prominent negative lower-tropospheric  $U'V'$  north of the low-level jet axis is a manifestation  
350 of the dominance of cyclonic vortices in sub-weekly wind fluctuations as shown in Figs. 2i-j.



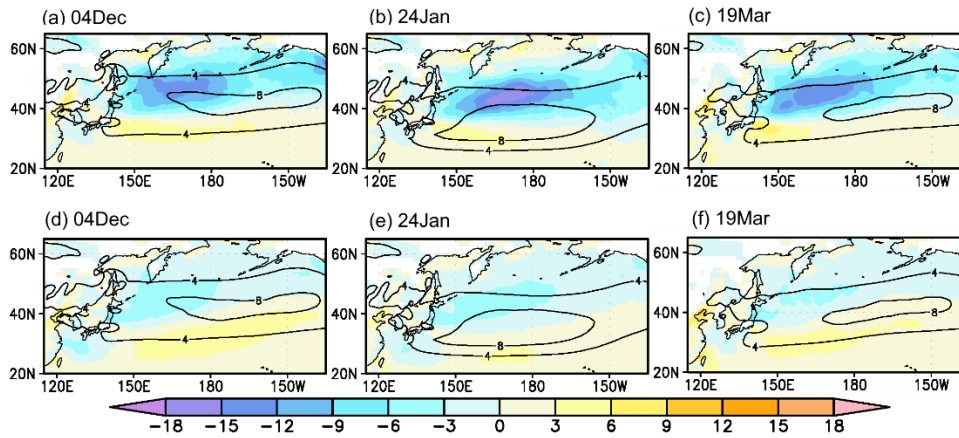


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Fig. 5. Same as in Figs. 2a-f, respectively, but for  $U'V'_{300}$  ( $\text{m}^2/\text{s}^2$ ). Black contours denote climatological-mean  $U_{300}$  ( $\text{m/s}$ ).



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Fig. 6. Same as in Fig. 5, but for  $U'V'_{925}$ .

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#### 4. Cyclonic and anticyclonic contributions to eddy feedback forcing on the westerlies and their seasonality

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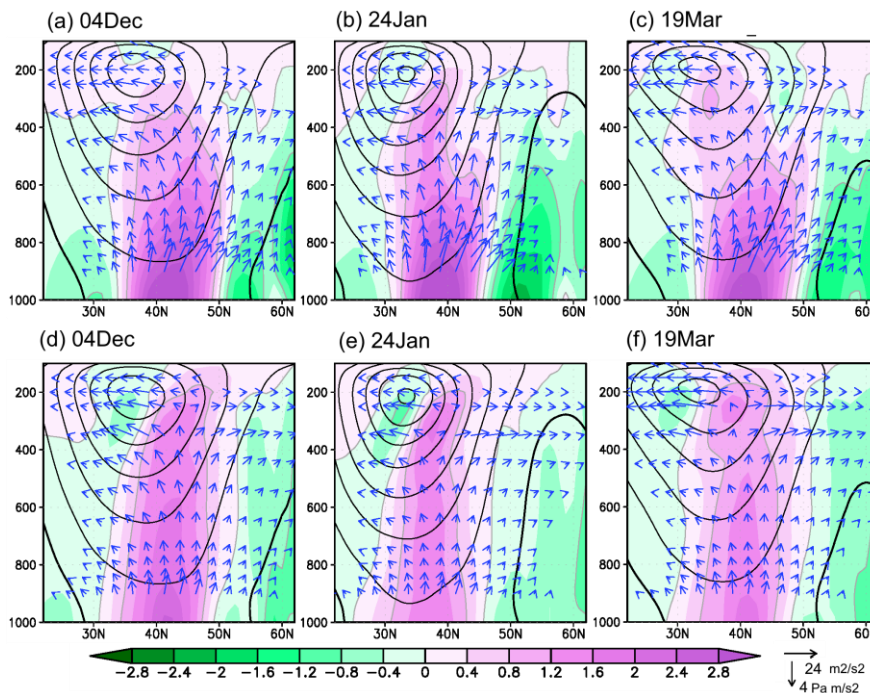
The deformation of cyclonic and anticyclonic vortices implies that they act to yield a converging flux of westerly momentum into the jet as discussed in previous studies (e.g., Blackmon et al. 1977; Lau and Holopainen 1984; ONK21), and their seasonality can exert some influence onto the seasonal march of the NP westerly jet. To quantify the westerly acceleration or deceleration as the separate contributions from cyclonic and anticyclonic vortices, three-dimensional transient eddy feedback forcing onto background westerlies is calculated in following ONK22.

366

367

Figure 7 shows meridional sections of westerly wind acceleration/deceleration by eddy heat and vorticity fluxes over the western NP as separate contributions from cyclonic and

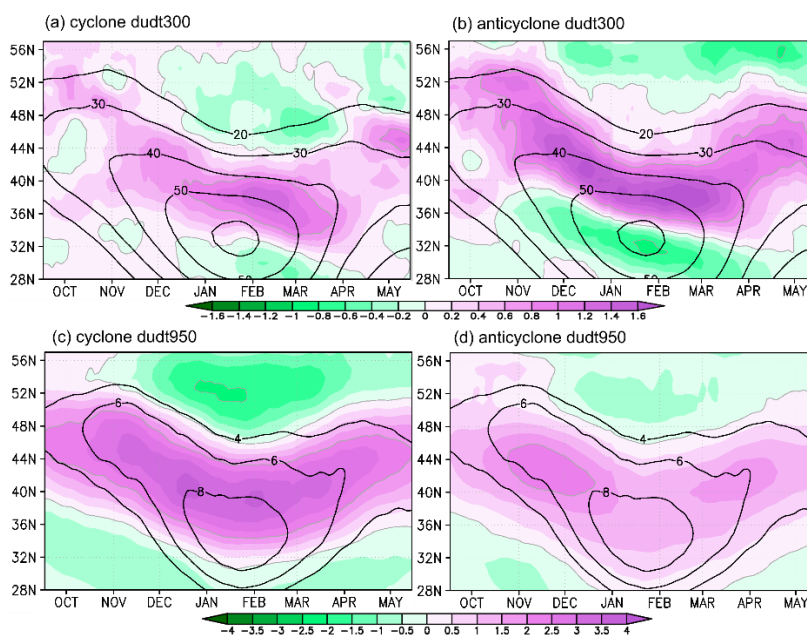
368 anticyclonic vortices. In the cold season, both cyclonic and anticyclonic vortices act to  
 369 reinforce the NP westerly jet slightly to the north of its axis. In the lower troposphere, the  
 370 westerly acceleration centered at  $\sim 40^\circ\text{N}$  by cyclonic vortices (Figs. 7a-c) is much stronger  
 371 than that by anticyclonic vortices (Figs. 7d-f), consistently with the upward extended E-P flux  
 372 (Fig. 7). The westerly deceleration north of  $\sim 50^\circ\text{N}$  also arises mainly from cyclonic vortices,  
 373 in agreement with  $U'V'_{925}$  (Fig. 6). As pointed out by ONK21, the westerly acceleration by  
 374 anticyclonic vortices is more coherent vertically (Figs. 7d-f), and it is comparable with or  
 375 even stronger than that by cyclonic vortices in the upper troposphere. Correspondingly to the  
 376 midwinter maximum of the equatorward transport of westerly momentum by cyclonic  
 377 vortices to the north of the upper and lower-tropospheric jet axes (Figs. 5b and 6b), the  
 378 westerly acceleration/deceleration by cyclonic vortices are strongest in midwinter, especially  
 379 along  $\sim 40^\circ\text{N}$  in the mid- and upper troposphere and  $\sim 50^\circ\text{N}$  in the lower troposphere.  
 380 Conversely, the lower-tropospheric westerly acceleration by anticyclonic vortices is  
 381 somewhat suppressed in midwinter (Fig. 7e), consistent with the suppression of midwinter  
 382  $V'T'_{850}$  (Fig. 3h).



383  
 384 Fig. 7. (a-c) Meridional sections of climatological-mean net westerly wind acceleration or  
 385 deceleration as transient eddy feedback forcing (color, m/s/day) by cyclonic vortices for the  
 386 calendar days as indicated. All the quantities plotted are zonally averaged for  $150^\circ\text{--}180^\circ\text{E}$ .  
 387 Black contours denote climatological-mean westerly wind speed (every 10m/s, thick for 0  
 388 m/s). Vectors indicate the extended E-P flux associated with transient eddies. (d-f) Same as in  
 389 (a-c), respectively, but for anticyclonic vortices.

390

391 The seasonal evolution of westerly wind acceleration/deceleration is highlighted in  
 392 latitude-season sections in Fig. 8. Throughout the cold season, the upper-tropospheric  
 393 westerly acceleration (deceleration) to the north (south) of the NP jet axis by anticyclonic  
 394 vortices is stronger than its cyclonic counterpart. This is indicative of a more important role  
 395 of anticyclonic vortices in transporting westerly momentum poleward from the NP hybrid jet,  
 396 the equatorward portion of which is maintained thermally by the Hadley Cell. The  
 397 predominance of the contribution from anticyclonic vortices to the westerly acceleration is  
 398 slightly reduced in midwinter from the shoulder seasons, reflecting the midwinter maximum  
 399 in the cyclonic contribution to  $U'V'_{300}$  to the north of the jet axis (Fig. 8b). In the lower  
 400 troposphere, westerly acceleration by cyclonic vortices remains rather constant throughout  
 401 the cold season with the midwinter maximum of westerly deceleration to the north (Fig. 8c).  
 402 The westerly acceleration by anticyclonic vortices, in contrast, exhibits its modest MWM  
 403 with a noticeable peak in late autumn through early winter, which may be related to the early-  
 404 winter peak of mid-tropospheric  $V'T'$  (Fig. 4d). Interestingly, southward expansion of the  
 405 near-surface westerly jet ( $\sim 30^\circ\text{N}$ ) in midwinter is attributable only to the acceleration by  
 406 anticyclonic vortices, counteracting the deceleration by cyclonic vortices in the subtropics.

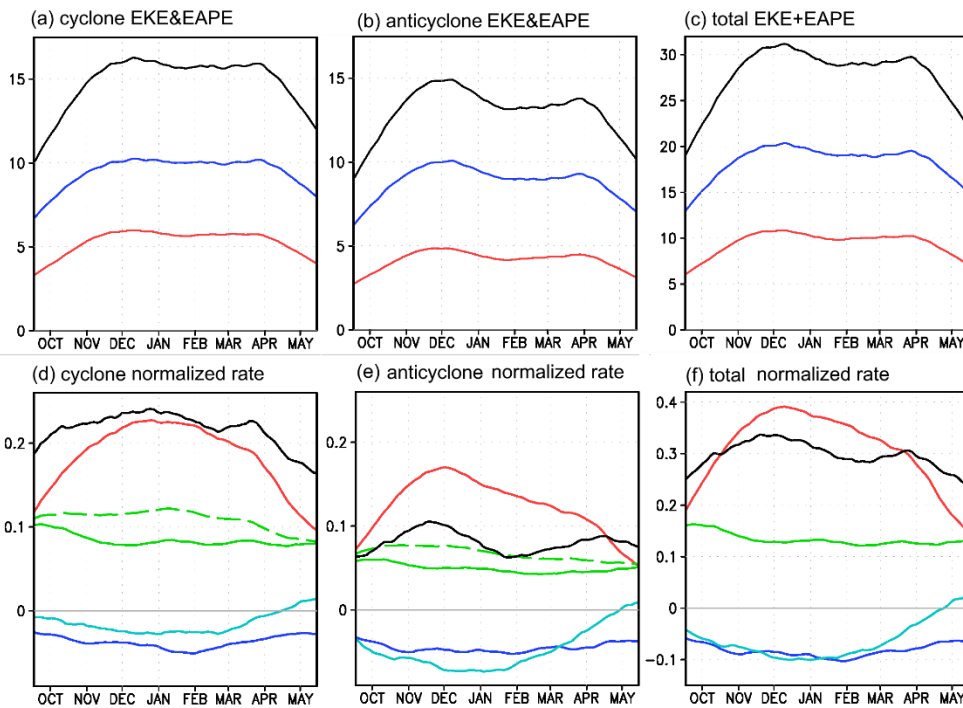


407  
 408 Fig. 8. (a) Climatological seasonality in net westerly wind acceleration or deceleration as  
 409 transient eddy feedback forcing (color, m/s/day) at 300-hPa by cyclonic vortices. Black  
 410 contours denote climatological-mean  $U_{300}$ . All the quantities plotted are averaged for  
 411  $150^\circ\text{--}180^\circ\text{E}$ . (b) Same as in (a), but for anticyclonic vortices. (c, d) Same as in (a, b),  
 412 respectively, but for the results at 950-hPa.

413

## 414 5. Separated energetics for the North Pacific storm-track

415 In this section, we evaluate the energetics for the NP storm-track based on the Eulerian  
 416 statistics that are decomposed into contributions from cyclonic and anticyclonic vortices as in  
 417 ONK21. Figures 9a and 9b show seasonal evolution of three-dimensionally integrated EKE  
 418 and EAPE associated with cyclonic and anticyclonic vortices, respectively. Throughout the  
 419 cold season, cyclonic EKE (blue line) and EAPE (red line) are larger than those associated  
 420 with anticyclonic vortices, presumably because of stronger fluctuations in wind and  
 421 temperature associated with cyclonic vortices in the lower and mid-troposphere. For cyclonic  
 422 vortices, EKE and EAPE are nearly constant throughout the cold season, with a subtle peak  
 423 of total eddy energy (EKE+EAPE; black lines) in early winter (Fig. 9a), which is consistent  
 424 with the nearly constant cyclonic  $V'V'_{300}$  (Fig. 2). Conversely, both the early-winter peak,  
 425 midwinter suppression and a secondary spring peak are more distinct for anticyclonic EKE  
 426 and EAPE (Fig. 9b), reflecting the obvious MWM of  $V'V'_{300}$  (Fig. 2).



427  
 428 Fig. 9. (a) Climatological-mean seasonal evolution of EKE (blue), EAPE (red), and  
 429 EKE+EAPE (black) associated only with cyclonic vortices ( $10^{18}\text{J}$ ). All the quantities plotted  
 430 are integrated three-dimensionally over the North Pacific. A tick mark on the abscissa  
 431 represents the first day of a given calendar month. (b) Same as in (a), but for anticyclonic  
 432 vortices. (c) Same as in (a), but for the sum of the cyclonic and anticyclonic contributions. (d-  
 433 f) Same as in (a-c), respectively, but for “normalized conversion/generation rates” of  $\lambda_{CK}$  (blue),  
 434  $\lambda_{CP}$  (red),  $\lambda_{CQ}$  (solid green),  $\lambda_{CQ}$  associated only with precipitation (dashed green),  
 435 and  $\lambda_{EF}$  (light blue). Black line denotes  $\lambda_{Tot}$  relevant to the budget of EKE+EAPE (viz.  
 436  $CK+CP+CQ+EF$ ). Unit:  $\text{day}^{-1}$ .

437

438 As described in section 2, we calculate the normalized rate (“ $\lambda$ ”) of each of the energy  
439 conversion/generation terms by dividing it by the total energy associated with transient eddies  
440 (without separating cyclonic and anticyclonic contributions), to investigate maintenance  
441 mechanisms for the NP storm-track activity more quantitatively (Figs. 9d-e). Throughout the  
442 cold season, cyclonic vortices tend to be more “energetic” than anticyclonic vortices, with  
443 substantially higher net normalized rate ( $\lambda_{Tot}$ ; black lines) of energy conversion and  
444 generation. Noteworthy,  $\lambda_{Tot}$  for anticyclonic vortices exhibits a well-defined MWM, while  
445 the corresponding midwinter reduction for cyclonic vortices is rather modest. This seems  
446 consistent with the seasonality of their total energy (Figs. 9a-b).

447 For both cyclonic and anticyclonic vortices,  $\lambda_{CP}$  (red lines) is much higher than any other  
448 conversion/generation rates during the cold seasons. In other words, the baroclinic energy  
449 conversion is the most important maintenance mechanism for cyclonic and anticyclonic  
450 vortices over the North Pacific, consistent with their baroclinic nature. Interestingly,  $\lambda_{CP}$  for  
451 cyclonic vortices is higher but not apparently suppressed in midwinter, while its anticyclonic  
452 counterpart peaks in early winter and undergoes more apparent suppression in midwinter.  
453 This contrasting seasonality is consistent with the more obvious midwinter suppression in  
454  $V'T'_{850}$  for anticyclonic vortices (Fig. 3). The early-winter peak in anticyclonic  $\lambda_{CP}$  is  
455 contributed to partly by the early-winter maximum in mid-tropospheric  $V'T'$  associated with  
456 anticyclonic vortices (Fig. 4). Given that anticyclonic vortices tend to be located near the  
457 westerly jet axis compared to cyclonic vortices in the middle and upper troposphere, the more  
458 distinct reduction in anticyclonic  $\lambda_{CP}$  from early winter to midwinter than the cyclonic  
459 counterpart seems consistent with Hadas and Kaspi (2021).

460 Compared to  $\lambda_{CP}$ ,  $\lambda_{CQ}$  (green lines) is less constructive throughout the winter with no  
461 signature of MWM for both cyclonic and anticyclonic vortices. The higher  $\lambda_{CQ}$  associated  
462 only with precipitation (green dashed lines) than the net  $\lambda_{CQ}$  (green solid lines) reflects the  
463 thermal damping due to eddy heat exchange with the underlying ocean.

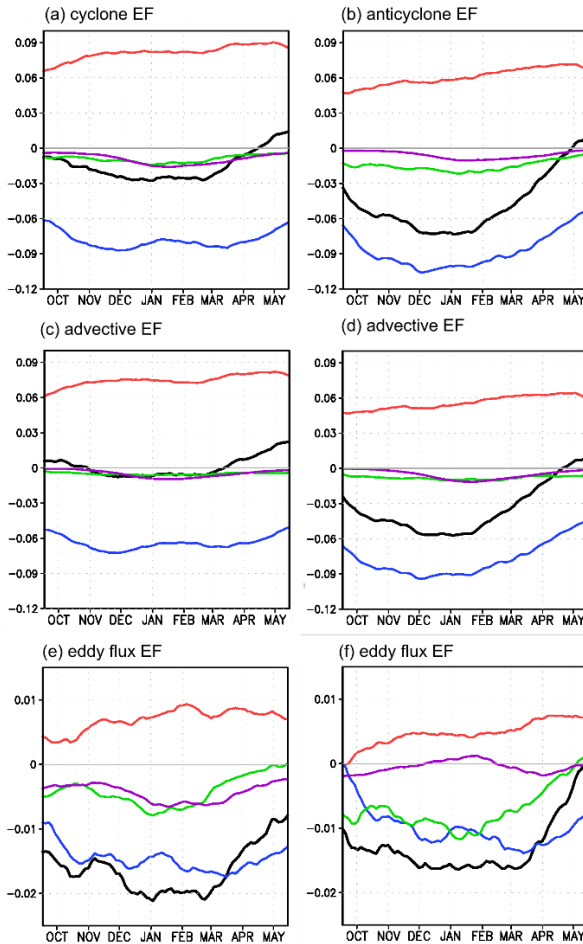
464 For the cyclonic  $\lambda_{Tot}$ , its early-winter peak and subsequent slight decrease into February  
465 arise mainly from  $\lambda_{CP}$ , with a slight contribution from  $\lambda_{CK}$  (blue lines), which becomes  
466 slightly more destructive into midwinter. The spring recovery of the cyclonic  $\lambda_{Tot}$  is, by  
467 contrast, due mainly to  $\lambda_{EF}$  (light blue lines) and additionally to the cyclonic  $\lambda_{CK}$ , which  
468 becomes less destructive into spring. For the anticyclonic  $\lambda_{Tot}$ , its marked decrease into

469 midwinter after its early-winter peak is attributable to both  $\lambda_{CP}$  and  $\lambda_{EF}$ , the latter of which  
470 becomes more destructive into January. The contribution from  $\lambda_{CK}$  is much more modest.  
471 The spring recovery of the anticyclonic  $\lambda_{Tot}$  after its MWM is due primarily to  $\lambda_{EF}$ , which  
472 fully offsets the continuous reduction of  $\lambda_{CP}$  into spring. The contribution from  $\lambda_{CK}$  is again  
473 much more modest.

474 As discussed above,  $\lambda_{EF}$  for anticyclonic vortices is most destructive in January. This  
475 indicates that anticyclonic vortices contribute to the net outward transport of eddy energy out  
476 of the NP domain most effectively in midwinter. As in ONK21,  $\lambda_{EF}$  can be separated into the  
477 contributions from energy fluxes passing through the four lateral boundaries of the  
478 rectangular domain and further into the contributions from the advective and eddy flux  
479 components. For cyclonic vortices (Figs. 10a, c, e), energy fluxes through the western and  
480 eastern boundaries, which are mainly due to the advection by climatological-mean flow,  
481 largely cancel out each other. The small net outflux in winter arises mainly from eddy energy  
482 flux out of the meridional boundaries. The spring increase of the net cyclonic  $\lambda_{EF}$  is  
483 attributable mainly to the advective  $\lambda_{EF}$  (Fig. 10c) with an additional contribution from eddy  
484 flux (Fig. 10e). Toward spring, the advective energy influx slightly increases, while the  
485 advective outflux is reduced. In addition, eddy geopotential fluxes out of the eastern and  
486 meridional boundaries decrease into spring.

487 For anticyclonic vortices (Figs. 10b, d, f), by contrast, energy flux out of the eastern  
488 boundary is substantially larger than the influx through the western boundary, leading to the  
489 net negative  $\lambda_{EF}$  especially in early and mid-winter. The large energy outflux in midwinter  
490 through the eastern boundary is presumably due to the high probability of anticyclonic  
491 vortices downstream of the prominent westerly jet core region (Fig. 1b), which may be  
492 attributable to the poleward deflection of the jet over the eastern NP. The outward flux by the  
493 mean westerlies through the eastern boundary peaks in early winter, when the connection of  
494 storm-track activity between the NP and NA is most discernible (not shown). This advective  
495 flux accounts for the major fraction of the net outward energy flux in winter. The contribution  
496 from eddy geopotential flux to the seasonality of the net anticyclonic  $\lambda_{EF}$  is in the same sense  
497 as the advective contribution but only of secondary importance. The smaller anticyclonic  $\lambda_{EF}$   
498 associated with the eddy geopotential flux through the eastern boundary than its cyclonic  
499 counterpart, despite the stronger eastward wave-activity flux associated with anticyclonic  
500 vortices (Fig. 4), may arise from those terms that are not related to eddy geopotential fluxes

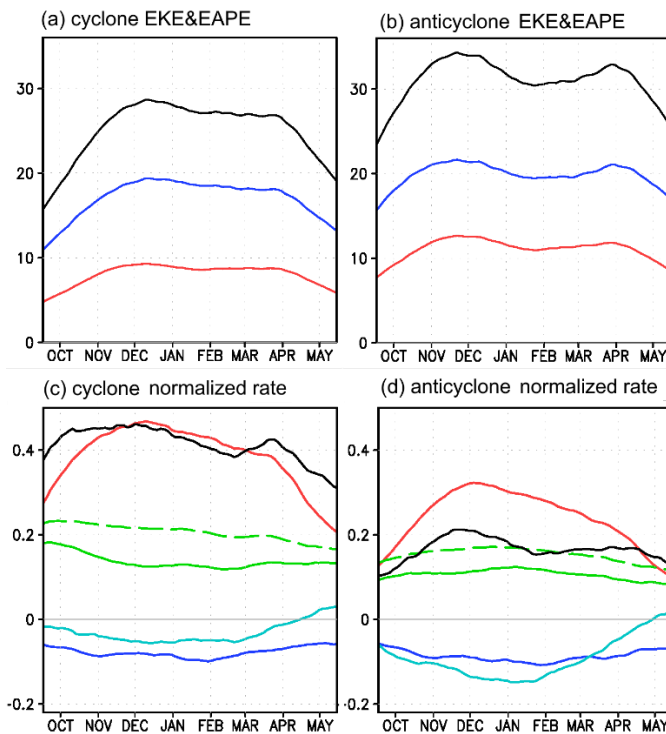
501 (c.f., Takaya and Nakamura 2001). Note that the results regarding  $\lambda_{EF}$  are qualitatively the  
 502 same if the eastern boundary is placed at a different longitude (e.g., 140°W).



503  
 504 Fig. 10. (a, b) Same as in Figs. 9c-d, respectively, but for the energy fluxes through the  
 505 western (red), eastern (blue), southern (purple), and northern (green) boundaries of the NP  
 506 domain. Black line denotes the net total energy flux. Positive (negative) values mean inward  
 507 (outward) fluxes. (c, d) Same as in (a, b), respectively, but for energy flux associated solely  
 508 with energy advection by climatological-mean flow. (e, f) Same as in (a, b), respectively, but  
 509 for energy flux associated solely with eddy geopotential flux.

510  
 511 The above results for transient eddy energetics also include the effect of the seasonality in  
 512 frequencies of cyclonic and anticyclonic vortices. To verify this effect, we reevaluate energy  
 513 and conversion/generation efficiencies by normalizing them with the probability of  
 514 positive/negative curvature averaged within the NP domain [130°–130°W, 20°–65°N], as  
 515 shown in Fig. 11. As shown in Fig. 1 and Part I, frequency of cyclonic (anticyclonic) vortices  
 516 basically increases (decreases) in midwinter: more specifically, probability of positive  
 517 curvature averaged over the reference domain is ~60% in midwinter, while ~50-55% in the  
 518 shoulder seasons. Although the seasonal evolutions of cyclonic and anticyclonic EKE and

519 EAPE thus obtained exhibit certain resemblance to those without the normalization by  
 520 frequency, anticyclonic EKE and EAPE normalized by their frequencies are larger than their  
 521 cyclonic counterpart (Figs. 11a-b), implying that anticyclonic vortices are not necessarily less  
 522 energetic than cyclonic vortices. As for the normalized energy conversion/generation rates,  
 523 the net rate is still higher for cyclonic vortices, owing mainly to the more effective  $\lambda_{CP}$  and  
 524 less destructive  $\lambda_{EF}$ . With the normalization by frequency, the late-winter minimum of the net  
 525 normalized rate for cyclonic vortices becomes more distinct. Conversely, the MWM of  
 526 anticyclonic net normalized rate becomes slightly less discernible. Consequently, the  
 527 seasonality of the net normalized rates of cyclonic and anticyclonic vortices becomes closer  
 528 to each other.



529  
 530 Fig. 11. Same as in Fig. 9, but for the results normalized by frequency of cyclonic and  
 531 anticyclonic vortices averaged over the North Pacific domain.

532

533 Finally, it should be noted that the above results are qualitatively similar with a threshold  
 534 curvature of  $4 \times 10^{-7} \text{ m}^{-1}$ , equivalently radius of 2500km.

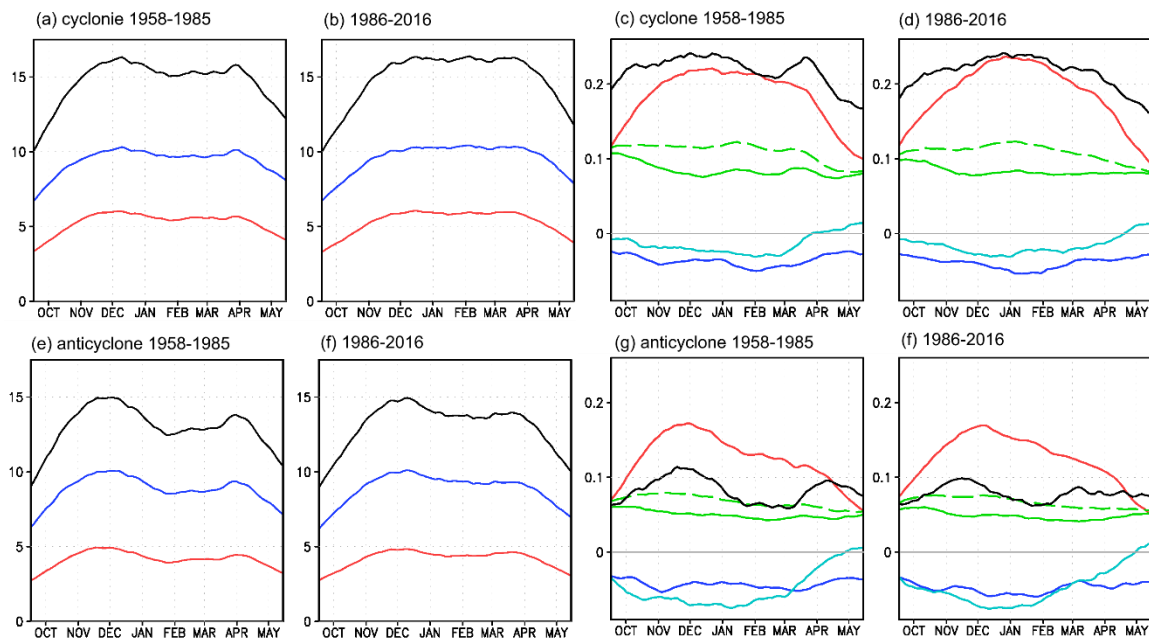
535

## 536 6. Long-term modulations



537 In this section, we briefly describe long-term modulations of the contributions to the  
 538 energetics from cyclone and anticyclone vortices. As in Part I, we divide the full analysis  
 539 period into the earlier (1958/59–1985/86) and later (1986/87–2016/17) periods, following  
 540 Nakamura et al. (2002). The result of the energetics for migratory eddies along the NP storm-  
 541 track separately for the two periods is comparable with ONK22.

542 In the earlier period, when the MWM of the NP storm-track activity was more distinct,  
 543 both cyclonic and anticyclonic EKE and EAPE contributed to the more distinct MWM (Figs.  
 544 12a, e). In the later period, by contrast, their MWM signatures became less distinct or even  
 545 almost vanished (Figs. 12b, f), in association with enhanced midwinter EKE and EAPE. This  
 546 means that both cyclonic and anticyclonic vortices contribute to the long-term modulations of  
 547 the MWM of the NP storm-track activity as evident in the full Eulerian statistics (ONK22).



548  
 549 Fig. 12. (a, b) Same as in Fig. 9a, but for the periods of (a) 1958/59–1985/86 and (b)  
 550 1986/87–2016/17, respectively. Unit:  $10^{18}$ J. (c, d) Same as in Fig. 9c, but for the periods of  
 551 (c) 1958/59–1985/86 and (d) 1986/87–2016/17, respectively. Unit:  $\text{day}^{-1}$ . (e-f) Same as in (a-  
 552 d), respectively, but for anticyclonic vortices.

553  
 554 As in section 4, we apply the energetic framework separately for the two periods. The  
 555 midwinter reduction in  $\lambda_{Tot}$  associated with anticyclonic vortices has also been weakened in  
 556 the later period, though the minimum is still noticeable (Figs. 12g-h). The enhancement of  
 557  $\lambda_{Tot}$  by anticyclonic vortices is contributed to by midwinter increase in  $\lambda_{CP}$ , while  $\lambda_{EF}$   
 558 becomes more destructive in spring and thus responsible for the reduced spring recovery of  
 559  $\lambda_{Tot}$ . The reduced  $\lambda_{Tot}$  in late winter for cyclonic vortices was clearly seen in the earlier

560 period (Fig. 12c), but it has disappeared in the later period (Fig. 12d). This interdecadal  
561 enhancement of the late-winter  $\lambda_{Tot}$  is caused by the enhanced  $\lambda_{CP}$  as well as the reduced  
562 destructive contribution from  $\lambda_{EF}$ .

563 Meanwhile, the destructive midwinter  $\lambda_{CK}$  has become more prominent in the later period  
564 from both cyclonic and anticyclonic vortices, which is mainly due to its “meridional”  
565 component related to  $U'V'$  (not shown). This suggests that the lateral shear of the NP  
566 westerly jet, which is assumed to be essential in the “barotropic governor” mechanism, is not  
567 necessarily relevant to the interdecadal enhancement of the midwinter NP storm-track  
568 activity for both cyclonic and anticyclonic vortices.

569

## 570 **7. Summary and Discussions**

571 As Part II of the paired papers, the present study has investigated the detailed seasonal  
572 evolution of contributions from cyclonic and anticyclonic vortices to Eulerian eddy statistics,  
573 eddy feedback forcing onto the climatological-mean westerly jet, and energetics associated  
574 with the NP storm-track. This has been achieved by combining the framework for the  
575 comprehensive energetics that is independent of eddy amplitude and the method for  
576 evaluating separate contributions from cyclonic and anticyclonic vortices to Eulerian eddy  
577 statistics through identifying three-dimensional domains of cyclonic and anticyclonic vortices  
578 based on local curvature. This “hybrid perspective” effectively illustrates the seasonal  
579 evolution of the contributions from cyclonic and anticyclonic vortices to Eulerian eddy  
580 statistics and atmospheric energetics.

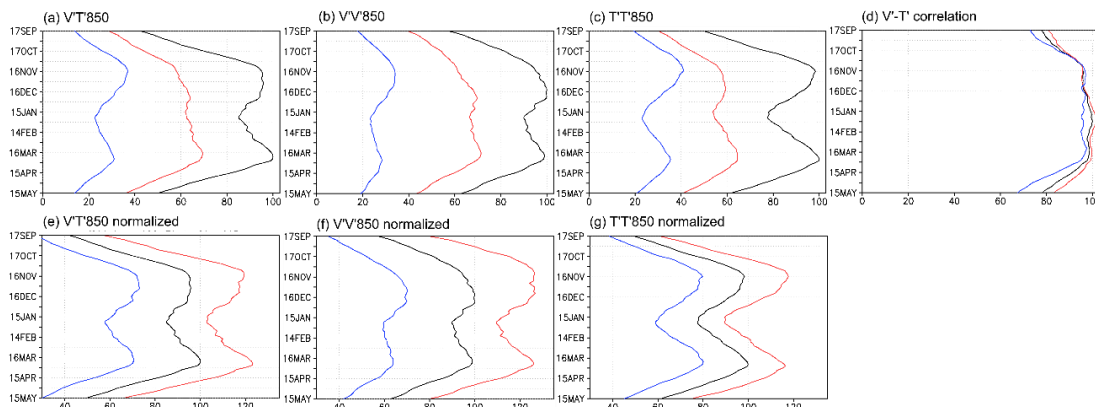
581 The key finding of this study is that the MWM of the net energy conversion/generation  
582 rate normalized by the eddy total energy ( $\lambda_{Tot}$ ) over the entire NP is more distinct for  
583 anticyclonic vortices than for cyclonic vortices. Together with Part I, the present study is the  
584 first to reveal the primal importance of migratory anticyclones for the MWM of the NP  
585 storm-track activity in a quantitative fashion. Given their broader and drier property  
586 compared to cyclones, the conclusion is compatible with the fact that the MWM of the NP  
587 storm-track activity has been reproduced even in coarse-resolution GCMs (Christoph et al.  
588 1997; Zhang and Held 1997). We posit the viewpoint of considering both migratory cyclones  
589 and anticyclones for the storm-track dynamics, where the latter has been long overlooked.

590 As mentioned by Wallace et al. (1988), the term *storm-track* used by Blackmon et al.  
591 (1977) implies that Eulerian eddy statistics are exclusively associated with cyclone tracks and  
592 thus anticyclone tracks are somehow irrelevant. Wallace et al. (1988) pointed out, however,  
593 that it is not necessarily the case, because the Eulerian eddy statistics does not imply any  
594 sense of polarity. The “hybrid perspective” of storm-track dynamics postulated in this study  
595 provides an answer to the long-standing question in a quantitative manner. Specifically,  
596 cyclonic vortices are found to contribute predominantly to  $V'T'_{850}$  as a typical measure of  
597 lower-tropospheric “storm-track” activity, whereas  $V'V'_{300}$  as a typical measure of upper-  
598 tropospheric “storm-track” activity includes a greater contribution from anticyclonic vortices.

599 The present study also provides some insights that complement the results based on the  
600 detailed energetics with the full Eulerian statistics by ONK22. They found that the higher  
601 mid- to upper-tropospheric  $V' - T'$  correlation and associated larger  $V'T'$  along the NP storm-  
602 track in early winter partly contribute to the early-winter peak in  $\lambda_{CP}$ . The present study has  
603 revealed that the high  $V' - T'$  correlation is associated mainly with anticyclonic vortices,  
604 although detailed mechanisms remain unsolved. Additionally, ONK22 pointed out that their  
605 result about the contribution from  $U'T'$  to  $CP$  does not seem consistent with Schemm and  
606 Rivière (2019), who showed the angle between the vectors of background baroclinicity  
607 gradient and eddy heat transport is maximized in midwinter over the western NP, leading to  
608 the MWM of the baroclinic energy conversion ( $CP$ ). However, the sign of zonal eddy heat  
609 flux is opposite between the upper and lower troposphere (ONK22). The present study has  
610 found that a midwinter maximum of positive  $U'T'$  below the jet core and another maximum  
611 of lower-tropospheric negative  $U'T'$  north of  $\sim 40^\circ\text{N}$  are distinct particularly for anticyclonic  
612 and cyclonic vortices, respectively (not shown). The negative lower-tropospheric  $U' - T'$   
613 correlation is more distinct in midwinter for both cyclonic and anticyclonic vortices,  
614 suggesting that they tend to tilt equatorward with height, presumably reflecting upper-  
615 tropospheric eddies trapped into the equatorward-shifted westerly jet, as suggested by some  
616 earlier studies (Nakamura and Sampe 2002; Hadas and Kaspi 2021). Nevertheless, the  
617 cyclonic contribution is dominant for the negative lower-tropospheric  $U' - T'$  correlation,  
618 owing to their higher probability than that of anticyclonic vortices north of  $\sim 40^\circ\text{N}$ . Indeed,  
619 slight positive midwinter  $\lambda_{CP}$  associated with  $U'T'$  pointed out by ONK22 is mostly due to  
620 cyclonic vortices (not shown), which is supportive of the above discussion. In other words,  
621 the tendency of midwinter NP cyclones to be more equatorward-tilted with height acts to

622 attenuate the MWM of the NP storm-track activity, although its contribution is overwhelmed  
 623 by  $CP$  associated with  $V'T'$ .

624 Considering that  $CP$  is the most important process for the maintenance of both cyclonic  
 625 and anticyclonic vortices, it is of interest what brings the seasonality of cyclonic and  
 626 anticyclonic contributions to  $V'T'_{850}$ . Figure 13 shows the decomposed contributions of  
 627  $V'V'_{850}$ ,  $T'T'_{850}$ , and  $V'-T'$  correlation at 850-hPa to  $V'T'_{850}$  ( $= \sqrt{V'V'_{850} \cdot T'T'_{850}} \cdot$   
 628  $\text{Corr}(V'_{850}, T'_{850})$ ). Consistent with Fig. 3, the MWM of  $V'T'_{850}$  is more distinct for  
 629 anticyclonic vortices (Fig. 13a), arising from all the contributions from velocity and  
 630 temperature fluctuations, as well as the  $V'-T'$  correlation. In addition, cyclonic  $T'T'_{850}$  also  
 631 exhibits a clear MWM. The seasonal evolution of the cyclonic and anticyclonic  $V'-T'$   
 632 correlation coefficients implies that cyclonic vortices tend to be slightly more baroclinic in  
 633 the lower troposphere, with highest  $V'-T'$  correlation in midwinter.



634  
 635 Fig. 13. (a) Climatological-mean  $V'T'_{850}$  averaged longitudinally for  $150^{\circ}$ – $180^{\circ}$ E. All the  
 636 quantities plotted are meridionally-averaged within 9 grids whose center is at its maximum in  
 637  $25^{\circ}$ – $55^{\circ}$ N of the total climatological-mean  $V'T'_{850}$  at each longitude. Red, blue, and black  
 638 lines indicate  $V'T'_{850}$  by cyclonic and anticyclonic domains, and the total climatological  
 639 mean, respectively. The values are normalized so that the maximum of the total  
 640 climatological-mean is set to 100%. (b-d) Same as in (a), but for (b)  $V'V'_{850}$ , (c)  $T'T'_{850}$ , and  
 641 (d)  $V'-T'$  correlation at 850-hPa averaged for  $150^{\circ}$ – $180^{\circ}$ E. Again, those quantities are  
 642 averaged around the axis of the total climatological-mean  $V'T'_{850}$  at each longitude. (e-g)  
 643 Same as in (a-c), respectively, but for quantities normalized by probability of cyclonic or  
 644 anticyclonic vortices averaged meridionally around the  $V'T'_{850}$  axis.

645  
 646 Those contributions to  $V'T'_{850}$  are also influenced by the seasonal evolution of lower-  
 647 tropospheric probability of cyclonic and anticyclonic vortices, in which cyclonic vortices are  
 648 more frequent along the NP storm-track axis as seen in Figs. 1d-f. We further examine the  
 649 seasonality of the cyclonic and anticyclonic  $V'T'_{850}$ ,  $V'V'_{850}$ , and  $T'T'_{850}$  by normalizing them  
 650 separately with the probabilities of cyclonic and anticyclonic vortices averaged meridionally

651 around the  $V'T'_{850}$  axis. The seasonality shown in Figs. 13e-g, which is thus independent of  
652 that of the probability, indicates that those eddy statistics related to cyclonic vortices also  
653 substantially decreased in magnitude in midwinter, which is consistent with Hoskins and  
654 Hodges (2019). This is indicative of the primal importance of the different seasonality of  
655 probability between cyclonic and anticyclonic vortices. Additionally, for both cyclonic and  
656 anticyclonic vortices,  $T'T'_{850}$  is more important for the MWM of  $V'T'_{850}$ .

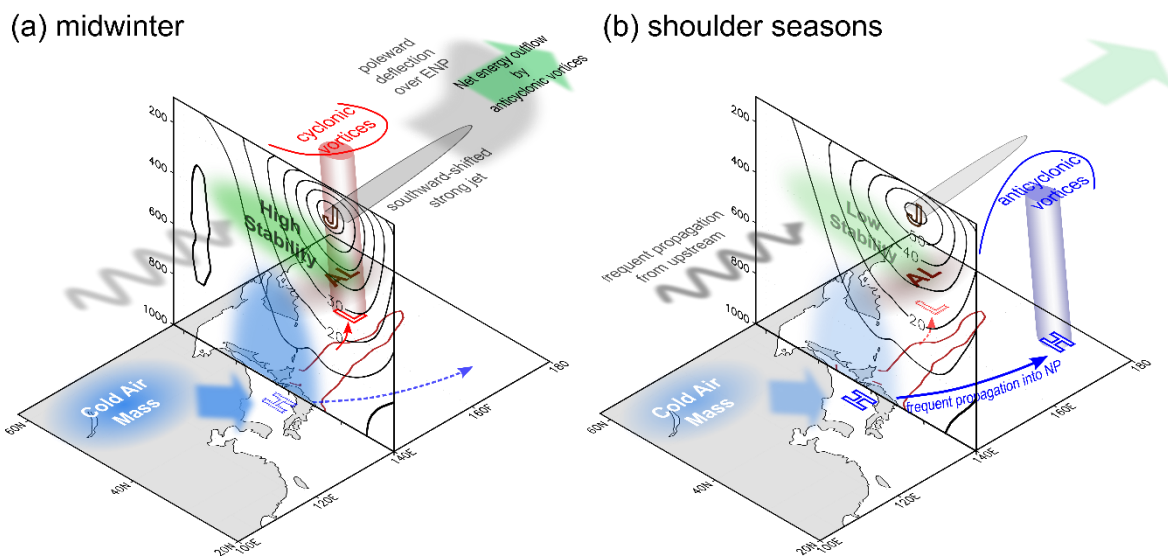
657 On the basis of the results obtained in this paper and Part I, Fig. 14 schematically  
658 illustrates climatological-mean evolution of the NP storm-track activity during the cold  
659 season. In comparison with the shoulder seasons, the midwinter situation is characterized by  
660 the following factors:

- 661 ● Higher mid-tropospheric static stability over the Japan Sea. This is attributable partly  
662 to the stronger Siberian High and the associated cold low-level monsoonal airflow.  
663 The higher static stability retards the vertical connection of upper-tropospheric eddies  
664 with lower-tropospheric disturbances, which leads to less frequent formation of  
665 surface migratory anticyclones over the Japan Sea.
- 666 ● Less frequent upper-tropospheric anticyclonic vortices over the Japan Sea (Part I). In  
667 sharp contrast, cyclonic vortices become more frequent from early- to mid-winter  
668 around  $\sim 40^\circ\text{N}$ .
- 669 ● Enhanced energy outflux downstream from the eastern NP associated mainly with  
670 anticyclonic vortices. This is presumably due to their high probability downstream of  
671 the jet core region. It may be explained by the poleward deflection of the midwinter  
672 westerly jet, on the equatorward flank of which “trough thinning and broadening”  
673 mechanism (Thorncroft et al. 1993) may be favorable for anticyclonic vortices. The  
674 poleward jet deflection may be a manifestation of seasonally intensified planetary  
675 waves (Wang and Ting 1999). Rodwell and Hoskins (2001) stressed the importance  
676 of the interaction between the zonal-mean flow and large-scale mountains for  
677 wintertime subtropical circulation, especially the equivalent-barotropic stationary  
678 anticyclone west of high mountains, while Orlandi (1998) argued that eddy feedback  
679 forcing has a similar effect. Indeed, our analysis indicates that transient eddy  
680 components, especially those associated with anticyclonic vortices, act to shift the jet  
681 axis poleward over the eastern NP. Another possible explanation is that migratory  
682 anticyclones, by nature, tend to propagate farther downstream compared to cyclones,

683 given the poleward-propagating tendency of cyclones away from the jet related to  
 684 diabatic heating (Tamarin and Kaspi 2016). The relationship between the poleward  
 685 deflection of the jet and the high probability of anticyclonic vortices over the eastern  
 686 NP under the eddy-mean flow interaction is to be further examined.

- 687 ● Weaker background baroclinicity for surface cyclones propagating into the quasi-  
 688 stationary AL. This is due to their greater distance from the baroclinic zone owing to  
 689 the equatorward displacement and narrowing of the midwinter westerly jet.

Seasonal evolution of North Pacific storm-track activity and related background conditions



690  
 691 Fig. 14. Schematic diagrams of climatological-mean situations in (a) midwinter and (b)  
 692 shoulder seasons. Black contours on vertical sections indicate zonally-averaged  
 693 climatological-mean westerly wind speed (m/s) for 130°–140°E on (a) 24Jan and (b) 19Mar.  
 694

695 The activity of cyclonic and anticyclonic vortices may be related to cyclonic/anticyclonic  
 696 Rossby wave breaking named LC2 and LC1, respectively (Thorncroft et al. 1993). It is  
 697 suggested that LC1 wave breaking acts to draw westerly momentum from an upper-  
 698 tropospheric subtropical jet into a midlatitude eddy-driven jet. By contrast, LC2 is likely to  
 699 occur under the cyclonic shear of a westerly jet, leading to the development of a cyclonic  
 700 vortex on the poleward flank of the jet, as observed in the midlatitude western NP. The  
 701 dominance of LC1 or LC2 in a numerical model used by Thorncroft et al. (1993) is sensitive  
 702 to a parameter chosen for the structure of an initial westerly distribution (Hartmann and  
 703 Zuercher 1998). Nevertheless, the results by Thorncroft et al. (1993) seem consistent with the  
 704 probability distributions of cyclonic and anticyclonic vortices obtained through our vorticity  
 705 identification based on local curvature (Fig. 1). Moreover, based on a dry idealized model,

706 Garfinkel and Waugh (2014) suggested that LC1 frequency tends to decrease on the  
707 equatorward flank of a westerly jet when shifted equatorward, which is compatible with the  
708 seasonal evolution of the NP storm-track. A linearized barotropic model experimented by  
709 Lorenz (2014) indicated that the stronger a westerly jet, the more poleward-propagating  
710 Rossby waves tend to be reflected around its poleward fringe (i.e., adjacent reflecting  
711 latitude), leading to an anomalous poleward flux of westerly momentum associated with  
712 equatorward-propagating waves after their reflection.

713 The results about the atmospheric energetics obtained in the present study can be  
714 regarded as a decomposition of the normalized energy conversion/generation rates obtained  
715 in ONK22 into cyclonic and anticyclonic contributions. We have confirmed that qualitatively  
716 similar results are obtained on the eddy energetics with a non-zero threshold of curvature.  
717 The present study provides a new perspective for the MWM of the NP storm-track activity by  
718 demonstrating the importance of both anticyclones and cyclones in the investigation of storm-  
719 track dynamics. There remains, however, a question to be answered: *why do the two*  
720 *contrasting seasonal evolutions of the activities of cyclonic and anticyclonic vortices lead to*  
721 *the MWM as their net effect?* It can be rephrased from the viewpoint of decomposed  
722 energetics: is there any necessity for their net contribution to decrease in the presence of an  
723 equatorward-shifted or narrow westerly jet with excessive intensity? Expanding our analysis  
724 by including a viewpoint from Rossby wave breaking will certainly be informative. It will be  
725 also informative to address the fundamental question by, for example, applying our  
726 framework to idealized experiments in which a jet stream can be artificially manipulated. The  
727 incorporation of energy exchange terms between cyclonic and anticyclonic vortices into the  
728 framework proposed in this study is also needed as an important future topic.

729 While some of the factors pointed out in this study are inherent to the geographical  
730 characteristics around the NP, which may be viewed as local effects, the present study can  
731 also fit with Yuval et al. (2018) and Novak et al. (2020), who argued that the MWM of the  
732 storm-track activity can be reproduced in idealized GCM experiments with zonally-  
733 symmetric lower-boundary conditions. Although the excessively strong Pacific jet owes its  
734 existence to such local effects as the climatologically deep East Asian trough, a stronger  
735 westerly jet, even if zonally uniform, may induce a higher probability of anticyclonic wave  
736 breaking on its equatorward flank, as suggested by previous studies.

737 The relationship between the results obtained in the present study and Part I needs to be  
738 further studied. This is because this paper explores contributions from three-dimensional  
739 cyclonic and anticyclonic vortices to Eulerian eddy statistics, whereas Part I focuses on  
740 centers of surface migratory disturbances, in which vortices without any distinct surface  
741 extrema have been excluded. The seasonality of composited structure of surface anticyclones  
742 should be investigated in the future. Additionally, the air-sea interaction associated with  
743 cyclones and anticyclones, especially over the North Pacific oceanic frontal zones, needs to  
744 be examined.

745 Finally, the analyses performed in this study are applicable to storm-tracks over other  
746 ocean basins. Comparison between the NP and NA storm-tracks within the framework of this  
747 study would be beneficial for further deepening of our understanding of storm-track  
748 dynamics and the MWM in storm-track activity. Our preliminary result suggests that, unlike  
749 the NP storm-track, there is no MWM in  $\lambda_{Tot}$  associated with anticyclonic vortices over the  
750 NA. A brief investigation of the long-term modulations of the MWM of the NP storm-track  
751 activity in this study should be complemented by a more detailed analysis in future.  
752 Interannual variability of the storm-track activity is also in our future scope. Furthermore, the  
753 framework of the present study may also be applicable to storm-tracks under the past and  
754 future climate, including the reduced storm-track activity during the last glacial maximum (Li  
755 and Battisti 2008).

756

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770

771 *Data Availability Statement.*

772 The JRA-55 atmospheric reanalysis is available online from the Japan Meteorological  
773 Agency at [https://jra.kishou.go.jp/JRA-55/index\\_en.html](https://jra.kishou.go.jp/JRA-55/index_en.html) as cited in Kobayashi et al. (2015)  
774 and Harada et al. (2016).

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