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Secular change of true polar wander over the past billion years

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Abstract:

The rate of movement of Earth's solid shell relative to its spin axis, or true polar wander, depends on variations in mantle convection and viscosity. We report paleomagnetic and geochronologic data from South China that constrain the rate of rapid true polar wander ($>5^\circ$ Myr⁻¹) from 832–821 million years ago. Analysis of the paleomagnetic database demonstrates secular change of true polar wander related to mantle cooling and thermal structure across supercontinent cycles. True polar wander rates are relatively muted with a partially insulated mantle during supercontinent assembly and accelerate as mantle thermal mixing re-establishes with supercontinent breakup. Decreasing true polar wander rate through the Neoproterozoic was succeeded by smaller variations in the Phanerozoic. We propose that extensive Neoproterozoic plate-tectonic activities enhanced mantle cooling, giving rise to a reduction in mantle convective forcing, an increase in mantle viscosity, and a decrease in true polar wander rates into the Phanerozoic.

One Sentence Summary:

Mantle cooling and thermal structure control secular change of true polar wander rates.

1 Introduction

2 True polar wander (TPW) is the rotation of a planet's or moon's entire solid exterior relative to its
3 spin axis in response to changes in its moment of inertia associated with mass redistribution (1).
4 Two first-order controls have been proposed to dictate Earth's TPW rate in geological time: the
5 magnitude of internal inertia perturbations, particularly convective loading that scales with the
6 vigor of mantle convection, and the viscosity of the lower mantle, which is temperature-dependent
7 (2–5). These two properties coevolve over Earth's history and are modulated by the secular cooling
8 of the mantle, through which decreasing mantle temperature leads to less vigorous convection and
9 higher viscosity, which should limit rapid TPW (2–5).

10 Paleomagnetic data constrain TPW to $\leq 3^\circ$ million year⁻¹ (Myr⁻¹) during the Phanerozoic (ca. 539
11 million years ago [Ma] to present) (6–9). More rapid TPW ($>4^\circ$ Myr⁻¹) has been revealed mostly
12 from Neoproterozoic rock records (1,000–539 Ma) (10–15). For example, the hypothesized ca.
13 810–795 Ma Bitter Springs TPW based on paleomagnetic results from the Akademikerbreen
14 Group in Svalbard implied $\sim 83^\circ$ of cumulative continental motion in ~ 15 Myr (10). However,
15 without direct dates on paleomagnetic poles, uncertainty has lingered in the inferred TPW rates
16 (10–12). Similarly, the significance of ca. 850–800 Ma pole motions of Baltica ($\sim 90^\circ$) remains
17 unresolved without more precise age constraints (16, 17). Moreover, discerning rapid TPW events
18 depends on the geocentric axial dipole (GAD) hypothesis, and large Ediacaran pole shifts have
19 been alternatively speculated to reflect the component of non-GAD fields (18, 19). Lastly,
20 determining the pure TPW signal involves removing simultaneous tectonic movements from the
21 composite pole shift that comprises the two, which has been a longstanding issue of concern to
22 circumvent (20).

23 To estimate TPW rates from the pole shifts complicated by tectonic movements, which remain
24 largely uncertain in the Neoproterozoic, new paleomagnetic poles with records of uniformitarian
25 geomagnetic fields and tight geochronologic controls are required. With precise age
26 determinations of a pole pair, tectonic movements during the time interval can be modeled and
27 subtracted to quantify the pure TPW signal. We have acquired high-precision U-Pb chemical
28 abrasion thermal-ionization mass-spectrometry (CA-TIMS) zircon dates coupled with
29 paleomagnetic poles from Neoproterozoic mafic sills, which intruded in the lower Fanjingshan
30 Group (Gp), Guizhou province, South China. Integrated with existing global data, our new results
31 discover an exceptionally rapid TPW event ($>5^\circ$ Myr⁻¹ based on the best estimation) between 832
32 and 821 Ma, leading to further analysis of the global paleomagnetic database to explore the
33 mechanisms for the changing TPW rate in Earth's history. Our analysis suggests a geodynamic
34 linkage that relates the observed overall decrease and periodic fluctuations in TPW rates with
35 mantle cooling and alternating thermal structures across supercontinent cycles. The resolved
36 geodynamic coupling provides new arrays of predictions for Earth's thermal history and
37 implications for the co-evolution of Earth's interior systems and rotational stability.

38 Results

39 The Fanjingshan mafic sills experienced sub-greenschist-grade metamorphism and were folded in
40 the Jiangnan Orogeny (ca. 830–815 Ma) (21, 22), permitting regional fold tests for the
41 paleomagnetic results. We sampled oriented paleomagnetic cores from a total of twelve diabase
42 sills in five sections, with intruded host-rock samples from the Jinzhanping section for a baked
43 contact test (fig. S1). The intrusive contacts of the sampled sills are parallel to bedding (fig. S1,

44 fig. S2). Five zircon crystals were selected from a geochronologic sample, TS01C, collected from
45 a gabbro-diorite sill for CA-TIMS analysis (fig. S1).

46 **Geochronology**

47 Zircon grains selected for CA-TIMS analysis were euhedral, elongated, similar-sized (~200–300
48 μm) and with sharp terminations, interpreted to represent crystallization from the same magma.
49 Isotope and age results are reported in Table S1. $^{206}\text{Pb}/^{238}\text{U}$ versus $^{207}\text{Pb}/^{235}\text{U}$ ratios show
50 concordant and coherent results (fig. S3A), indicating successful removal of potential lead loss in
51 zircon grains after annealing and chemical abrasion procedures. The weighted mean $^{206}\text{Pb}/^{238}\text{U}$
52 date calculated from the zircon analyses gives 831.51 ± 0.32 Ma (2σ uncertainty) (fig. S3, Table
53 S1), representing the crystallization age of the mafic sill.

54 **Paleomagnetism**

55 Seven sills from which 125 specimens were sampled yielded stable and consistent high-
56 temperature components. Two low-temperature components, LT1 and LT2, were first removed
57 during thermal demagnetization. LT1, separated mostly below ~ 450 °C, is present in most samples
58 and identified as a viscous remanent magnetization (VRM) of the recent geomagnetic field (fig.
59 S5, fig. S6). LT2 (300–350 °C), a southeast-up direction in geographic coordinates, is only found
60 from the five sills in the Huguosi and Jingzhanping sections (fig. S5, fig. S6). LT2 fails a fold test
61 and is also secondary (fig. S6), but its age and origin are relatively unconstrained.

62 A stable high-temperature component, HT1, carried by magnetite, is identified from all seven sills
63 (~ 500 – 580 °C) and their baked contacts (~ 480 – 550 °C) (fig. S5). HT1 passes the McFadden fold
64 test (23) at 95% confidence level and a progressive unfolding test (24) that shows the tightest
65 distribution obtained at 98.9% unfolding (fig. S7), suggesting a magnetization acquired before the
66 folding of the Fanjingshan Gp by ca. 815 Ma (12). Paleomagnetic results from the intruded host-
67 rocks in the Jinzhanping section demonstrate a positive baked contact test: baked siltstones record
68 a similar HT1 from that of the sills, while unbaked siltstones preserve a distinct high-temperature
69 component (HT2) from HT1 and other directions determined in the region (fig. S9). The positive
70 baked contact test and fold tests indicate that HT1 is a primary magnetization acquired during the
71 cooling of the sills at 831.51 ± 0.32 Ma.

72 The corresponding paleomagnetic pole of HT1 (the Fanjingshan pole, 34.7°S , 118.2°E , $A_{95} = 8.6$)
73 does not resemble any younger pole reported from South China, further supporting a primary
74 origin (Fig. 1A, fig. S11). Multiple dispersion parameters of virtual geomagnetic poles (VGPs) fall
75 within expected values of paleosecular variation models (25, 26) (fig. S11), supporting that the
76 Fanjingshan pole has sufficiently recorded the time-averaged position of the geomagnetic pole at
77 ca. 832 Ma, during which the fidelity of the GAD field is also independently substantiated by
78 magnetostratigraphic studies of coeval sedimentary successions (17) and comprehensive analyses
79 of the Precambrian geomagnetic field (e.g., 27).

80 **Discussion**

81 **Discovery of rapid TPW between 832 and 821 Ma**

82 The new Fanjingshan pole places the sampling location at a paleolatitude of $\sim 27^\circ$ at ca. 832 Ma
83 (Fig. 1A). The Fanjingshan pole and the 821.64 ± 0.2 Ma Xiaofeng dykes pole (28, 29) quantify
84 the paleolatitudinal displacement of South China to $\sim 41^\circ$ between 832 and 821 Ma, at an average
85 rate of $\geq 46.5 \pm 15.7$ cm yr^{-1} (Fig. 1A). Observational and theoretical speed limits of modern plates

86 are $\sim 19 \text{ cm yr}^{-1}$ and $\sim 20 \text{ cm yr}^{-1}$, respectively (30, 31). The mean latitudinal plate velocity implied
87 by the Fanjingshan and Xiaofeng poles is at least twice the upper limits of the known plate-tectonic
88 speeds. It is thus interpreted to record a combined signal of tectonic movements and TPW. The
89 total pole shift between 832 and 821 Ma is quantified to be $54.7 \pm 13.9^\circ$ with a mean rate of $5.5 \pm$
90 $1.4^\circ \text{ Myr}^{-1}$.

91 South China was either disconnected entirely from Rodinia (12, 32, 33) or located at its
92 periphery during the Tonian (1,000–720 Ma) (34). The direction and speed of the tectonic
93 motions of South China between 832 and 821 Ma are currently unknown, due to a paucity of
94 exactly coeval data from other continents for comparison of relative tectonic motion.
95 Considering that signals of TPW and tectonic motion might enhance or counteract each other
96 (20), implementing a Monte Carlo method, we simulate tectonic motions of South China
97 between 832 and 821 Ma, by which the total pole shifts are corrected for to constrain
98 quantitatively the pure TPW component (Fig. 2, fig. S15) (Materials and Methods). The
99 disentangled 832–821 Ma TPW rate has an estimated 95% confidence interval of $3.5\text{--}7.8^\circ \text{ Myr}^{-1}$
100 (Fig. 2C). The rate is comparable to the hypothesized Bitter Springs TPW constrained to be
101 $\geq 3.6\text{--}4.9^\circ \text{ Myr}^{-1}$ for the entire Bitter Springs Stage (from ca. 810 to $<795 \text{ Ma}$) (10, 11) and $\geq 5.4^\circ$
102 Myr^{-1} for the first portion of the TPW oscillation (ca. 810–795 Ma) (12).

103 Large paleolatitudinal continental motions during the Tonian were also revealed from the
104 paleomagnetic poles of Baltica (16, 17, 35, 36). These poles, dated roughly between 850 and 800
105 Ma, demonstrate a total rotation of $\sim 90^\circ$. The pole comparison is illustrated in Fig. 1 C–D, which
106 shows the lining up of the apparent polar wander paths (APWPs) of South China and Baltica
107 between ca. 850 and 810 Ma, reconstructed in relative configurations consistent with those
108 proposed for Rodinia (Euler rotation parameters and references see Table S7). This ca. 850–810
109 Ma Rodinia APWP is further supported by the ca. 850 Ma Manso dykes pole from West Africa
110 (37) and the ca. 820–810 Ma lower Grusdievbreen Formation pole from East Svalbard (10),
111 which both fall closely along the APWP defined by the South China and Baltica poles (Fig. 1 C–
112 D). Both East Svalbard and Baltica are considered part of assembled Rodinia in the Tonian (10,
113 36). The discovery of comparable prominent paleogeographic shifts of congruent Rodinia and
114 presumably separate South China supports the hypothesis that these movements record an
115 extended period of rapid TPW from 850 to 795 Ma: 850–820 Ma equator-ward migration of
116 Rodinia and its associated geoid high ('type I' TPW) (Fig. 1 C–D) followed by 810–795 Ma
117 oscillatory motions about the center of Rodinia (inertial interchange [II] TPW) once the positive
118 geoid docked and stabilized in the equatorial zone (3, 5, 20, 38, 39). Similar to Rodinia, Pangaea
119 was proposed to have witnessed a similar fashion of TPW succession from northward migration
120 (400–250 Ma, I TPW) to equator-centered oscillations (250–100 Ma, II TPW) while at lower
121 average rates (3, 39) (Fig. 3A).

122 **Secular change of TPW over the past billion years**

123 The findings of 850–795 Ma rapid TPW events enable contextualizing the Tonian observations
124 in the evolution of TPW over time. The inferred high Tonian TPW rates exceed the speed limit
125 of observed TPW through the Phanerozoic ($\sim 3^\circ \text{ Myr}^{-1}$) (6–9) and theoretical constraints ($\sim 2.4^\circ$
126 Myr^{-1}) (2) with present-day lower mantle viscosity estimates. Numerical simulations suggest that
127 rapid TPW at $>4^\circ \text{ Myr}^{-1}$ with a max rate of $\sim 6^\circ \text{ Myr}^{-1}$ can be achieved with substantially
128 increased convective flow (4) and/or reduced lower mantle viscosity to $\sim 3\text{--}10 \cdot 10^{21} \text{ Pa s}$, 10–30%
129 of the present value (3) consistent with thermal history modeling for the Tonian (40). These
130 comparisons support a more vigorously convecting and less viscous mantle in the Tonian that

131 allowed more rapid reorganization between principal moments of inertia. We synthesize and
132 analyze the extant probable estimates for the TPW rate over the past billion years (Fig. 3A, Table
133 S8). The time series manifests a general decreasing nature of the rate since the Tonian, a
134 significant drop toward the Cambrian (from 832 to 505 Ma), and a relatively more stable state
135 through the Phanerozoic (505–0 Ma) (Fig. 3A). Two TPW peaks occur at around 832–795 Ma
136 and 170–80 Ma, coinciding with the early breakup of Rodinia and Pangaea, respectively (41, 42).
137 The Ediacaran rapid polar wander rates (conservative estimates) appear to be compatible with an
138 implied TPW reduction from the Tonian to the early Phanerozoic if TPW rather than non-GAD
139 fields was the driving mechanism (Fig. 3A) (13, 15).

140 We propose a geodynamic link to explain both the systematic decrease in the TPW rate and
141 episodic variations imposed over the mean trend since the late Mesoproterozoic, in which mantle
142 cooling coupled with supercontinent cycles. The observed overall secular TPW decline is
143 consistent with TPW rate controlled primarily by increasing lower mantle viscosity and declining
144 vigor of mantle convection as Earth's interior cooled (3–5) (Fig. 3A, fig. S14). The two high
145 TPW stages, between 832–795 Ma and 170–80 Ma (Fig. 3A), coincide with transitions in the
146 Rodinia and Pangaea supercontinent cycles: preceded by protracted supercontinent thermal
147 insulation and succeeded shortly by supercontinent dispersal (42–44). During the 1800–850 Ma
148 Nuna-Rodinia lifetime (44, 45) and the 320–180 Ma Pangaea lifetime (7), supercontinents
149 rimmed by peripheral subduction zones may have partially insulated the underlying mantle and
150 altered mantle thermal structures: warming the subcontinental mantle and cooling the suboceanic
151 mantle (temperature variation on the order of 100 °C or greater) (46–48). This insulating regime
152 resulted in inhibited convective mixing between the subcontinental and suboceanic domains,
153 localized mantle flow patterns relative to the supercontinent, raised viscosity of voluminous
154 suboceanic mantle, and associated decreased velocity of oceanic plates (46, 47), setting the stage
155 for slow TPW (Fig. 3A–B). The 1,110–1,080 Ma TPW rates appear to have been restrained more
156 considerably – a presumed consequence of longer-lived insulation by Nuna and Rodinia
157 compared to Pangaea, creating greater lateral mantle thermal and viscosity variations (46, 47).

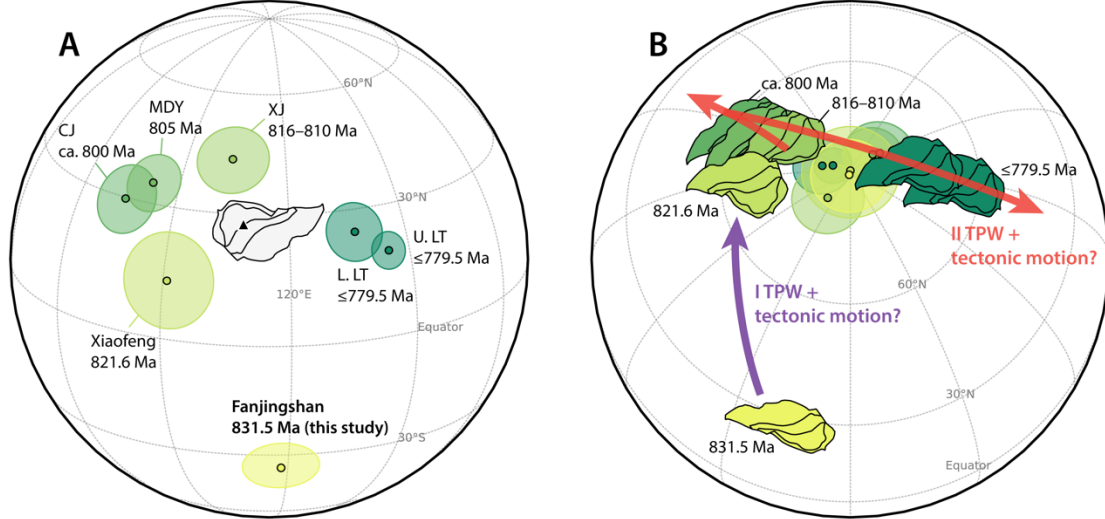
158 As heat accumulated beneath the subcontinental region and the lateral hydrostatic pressure
159 gradient built up in the mantle, the peripheral subduction girdle was destabilized and a
160 supercontinent approached the wake of dispersal (39, 46–48). Breakdown of a subduction girdle
161 would re-establish large-scale advective exchange between the subcontinental and suboceanic
162 mantle, relaxing mantle thermal and viscosity anomalies (46). This mantle thermal mixing
163 involves unleashing a pulse of increased convective vigor driven by the preexisting lateral
164 temperature gradient, accelerated overturn of oceanic plates, and enhanced global convective
165 motions (45–48), invigorating mass redistributions in Earth's interior and on the surface that
166 speed up TPW. Meanwhile, as mantle thermal re-equilibrium proceeds, warming the once cooled
167 and more viscous suboceanic mantle might reduce the maximum viscosity of the lower mantle
168 (rate-limiting TPW), which potentially relaxes the restriction on TPW rates (5). Such a
169 changeover of mantle thermal state could take over ~100 Myr to accomplish (46–48) and may
170 explain the accelerated TPW rates prompted at 832–795 Ma, 170–145 Ma, and ca. 80 Ma,
171 succeeding the breakup of Rodinia and Pangaea (Fig. 3). Following supercontinent breakup,
172 elevated mantle flow, surface heat flux, and lithospheric mobility would have amplified mantle
173 cooling rates initially, followed by gradual restoration to baseline cooling trend as convective
174 forcing weakens toward the completion of mantle thermal mixing (46–48). The relaxed
175 convective vigor contributes to the subsequent decline and stabilization of TPW rates until the
176 formation of the next supercontinent cycle (Fig. 3A). This model, if valid, underscores the

177 increasing importance of accounting for large-scale and varying lateral thermal heterogeneities
178 associated with mantle thermal structures in future numerical modeling of TPW.

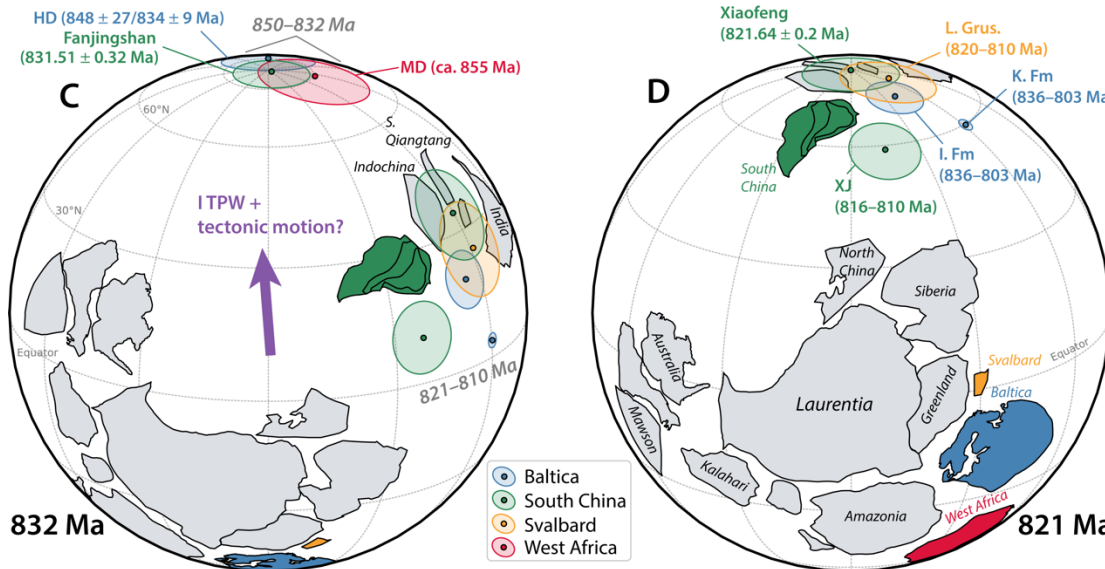
179 Notably, the striking reduction in TPW rate during the Neoproterozoic might suggest an
180 irreversible transition, after which the mantle cooled enough to render a slow-TPW regime in the
181 Phanerozoic ($<3^\circ \text{ Myr}^{-1}$), also characterized by an overall smaller amplitude of variation in TPW
182 rates (Fig. 3A). This implied mantle thermal decline in the Neoproterozoic resonates with coeval
183 lithospheric evidence showing a rising dominance of low thermobaric ratios (T/P) of rock
184 records since the Tonian that subsequently persisted through the Phanerozoic (44, 49). The
185 promoted and continual prevalence of low T/P metamorphic records, i.e., the blueschist-facies
186 and ultrahigh-pressure metamorphism, is inferred to be directly linked to cooling of the upper
187 mantle, leading to colder, denser, and strengthened oceanic lithosphere and deeper slab breakoff
188 in the collision zones (49, 50). Numerical models suggest such a transition in the style of
189 orogenesis from hotter to colder may have taken place as the mantle temperature fell to ~ 80 –
190 100° C higher than present-day values (50), corresponding to an inferred Neoproterozoic timing
191 constrained by igneous rock records and thermal modeling (51). Such mantle heat loss may have
192 been facilitated by plate-tectonic activity since the Tonian associated with the breakup of
193 Rodinia (peaking at ca. 800–600 Ma) and the assembly of Gondwana (ca. 650–500 Ma) (52). In
194 particular, cooling may have been facilitated by extensive circum-supercontinent subduction
195 zones and internal rifting (52). Consequently, by the early Cambrian, progressively reduced
196 mantle heat, which would have led to suppressed mantle convection and raised viscosity, in turn,
197 deactivated mantle physical conditions favorable for prompting rapid TPW in succeeding periods
198 (Fig. 3A).

199 Characterizing the secular change of TPW raises further opportunities and considerations in
200 paleogeographic reconstructions. The established framework provides testable predictions for
201 facilitating subsequent explorations of TPW to improve and revise our extant understanding.
202 Despite an expanding global TPW database, records of well-defined TPW excursions have
203 remained generally elusive. Nevertheless, the importance of high-resolution sampling in
204 resolving fast while smaller-amplitude TPW ($\sim 10^\circ$) has been demonstrated recently (9), which
205 suggests more potential missing TPW signals to search for, especially during the dynamic TPW
206 stages as our model predicts (Fig. 3). With better characterization of TPW from new
207 observations, full accounting of TPW motions, including periods throughout supercontinent
208 transitions, may become available to aid in the solution of absolute paleolongitude of ancient
209 plates (53). The secular change of TPW also gives insights into tracing the changing rate of plate
210 tectonics. As TPW manifests with coherent motion of global plates, the intervals of large-scale
211 TPW would reinforce the sensitivity in quantifying plate-tectonic rates concurrently with TPW,
212 as already explored by analyzing differences of contemporaneous pole motions (e.g., 12, 54).
213 Most existing estimates of Proterozoic plate velocities are not corrected for possible TPW (e.g.,
214 55), particularly for the Neoproterozoic when TPW was likely faster and more variable.
215 Discerning rates between TPW and tectonic movement in Earth's deep time would enable
216 probing possible coupling of the two in the long term, both functions of evolving mantle thermal
217 history and cyclic transformations in continental configurations and mantle structures (1, 3–5,
218 31, 46).

Apparent polar wander path for South China

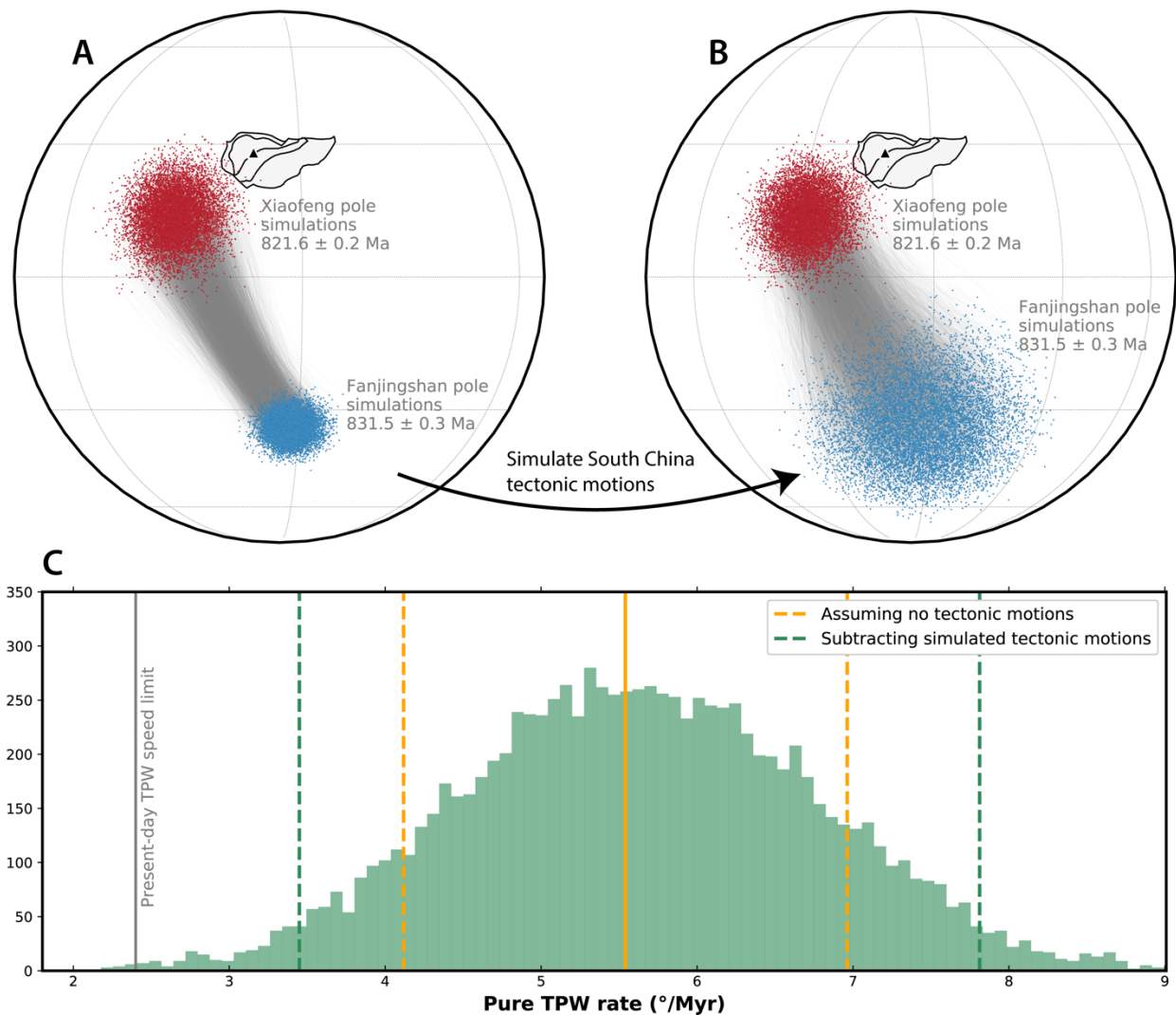


Rodinia paleogeography



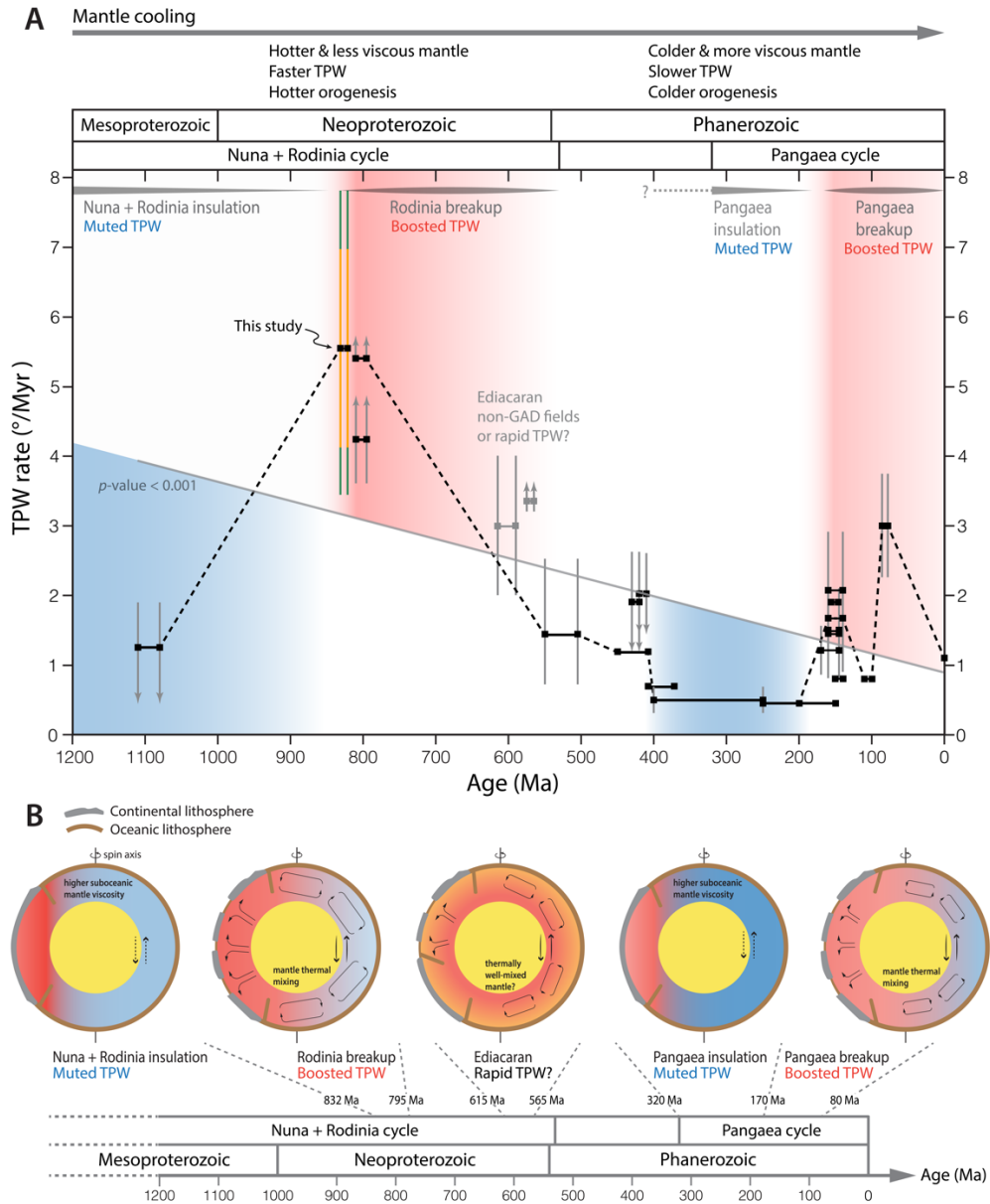
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220 **Figure 1. Mid-Tonian paleogeographic models for South China and Rodinia.** (A) Tonian apparent
 221 polar wander path (APWP) of South China. Ellipses show A_{95} , the 95% confidence cones of the poles.
 222 Triangle signifies the mean sampling locality of the Fanjingshan sills (27.92°N, 108.69°E) for paleolatitude
 223 conversion. See Table S3 for details of data. (B) The proposed paleogeographic model for South China
 224 reconstructed with the APWP in (A). Paleogeographic changes between 816 and 780 Ma follow the
 225 preferred ‘minimize difference’ model from (12). Hybrid TPW + tectonic motions are interpreted for 832–
 226 821 Ma and 816–795 Ma. (C and D) Rodinia paleogeographic changes in the context of 832–821 Ma ‘type
 227 I’ TPW. Poles and Euler rotation parameters are provided in Table S3 & S4 & S7. Pole abbreviations:
 228 CJ—Chengjiang Formation; MDY—Madiyi Formation; XJ—Xiajiang Group; L. LT—lower Liantuo
 229 Formation; U. LT—upper Liantuo Formation; HD—Hunnedalen dykes; MD—Manso dykes; L. Grus.—
 230 lower Grusdievbreen Formation; K. Fm—Katav Formation; I. Fm—Inzer Formation.



231

232 **Figure 2. 832–821 Ma TPW rate estimated using Monte Carlo analysis.** (A) Resampled Fanjingshan
 233 and Xiaofeng poles from Fisher distribution ($n = 10,000$ for each). Pole pairs are randomly drawn between
 234 the two pole groups, connected with fine grey lines representing angular distance. Triangle indicates the
 235 mean sampling locality of the Fanjingshan sills (27.92°N, 108.69°E). (B) Simulation of South China
 236 tectonic motions between 832 and 821 Ma assuming a speed limit of implied fastest Neoproterozoic plate-
 237 tectonic velocities ($\sim 31 \text{ cm yr}^{-1}$) (56) with unconstrained directions and application of these displacements
 238 to the resampled Fanjingshan poles from (A). The resultant new pole pairs reflect the pure TPW component
 239 after subtracting simulated tectonic motions. Dividing the pure TPW angular distances by the age
 240 differences between the pole pairs yields pure TPW rate estimates. (C) Histogram of pure TPW rate
 241 estimates from (B). The green dashed lines bound the 95% confidence intervals of the TPW rate after
 242 accounting for simulated tectonic motions ($3.5\text{--}7.8^\circ \text{ Myr}^{-1}$). The orange solid and dashed lines show the
 243 mean rate and 2σ uncertainties of the total pole shift between the Fanjingshan and Xiaofeng poles using
 244 standard error propagation ($5.5 \pm 1.4^\circ \text{ Myr}^{-1}$) (Materials and Method), which could reflect TPW estimates
 245 only when assuming no tectonic motions of South China. The grey solid line shows the proposed TPW
 246 speed limit ($\sim 2.4^\circ \text{ Myr}^{-1}$) set by present-day mantle viscosity estimates (2). See Materials and Methods for
 247 details.



248

249 **Figure 3. Secular change of the TPW rate and our proposed geodynamic coupling.** (A) Probable TPW
 250 rate estimates since the late Mesoproterozoic. TPW rates and uncertainties are summarized in
 251 Supplementary Text and tabulated in Table S8. The grey solid line indicates a linear regression fit of the
 252 data (p -value < 0.001) that characterizes the first-order decreasing trend (the fit excludes the Ediacaran data
 253 of debated origins, shown as grey squares). See fig. S14 for sensitivity tests that also consider variously
 254 interpreted Ediacaran TPW rate estimates. The black dashed line traces the second-order variation in TPW
 255 connected by representative rate estimates of individual time intervals. Grey vertical bars show the
 256 uncertainties of rate estimates. Up/down arrows indicate that the rate could be above/below the shown range
 257 defined by the error bars. For the TPW rate uncertainty of this study, the orange bar represents considering
 258 no tectonic motions and the green bar reflects accounting for simulated tectonic motions (corresponding to
 259 those shown in Fig. 2C). (B) Schematic illustration of secular change of TPW related to mantle cooling and
 260 alternating mantle thermal structures linked to assembly and breakup of continents (supercontinent
 261 aggregation–mantle thermal insulation–muted TPW; supercontinent breakup–mantle thermal mixing–
 262 boosted TPW).

263 **Materials and Methods**

264 **Field sampling**

265 Paleomagnetic cores were collected using a portable gasoline-powered drill with non-
266 magnetic, diamond-rimmed bites. The intrusive contacts of the sampled mafic sills are parallel to
267 the bedding of the country rocks (sandstone and siltstone of the lower Fanjingshan Group). Sills
268 with clearly observable contacts were prioritized for sampling. The seven sills that yielded stable
269 and consistent high-temperature components used to calculate the pole all have direct bedding
270 controls. Multiple bedding measurements were made for the country rocks closely adjacent to the
271 contacts to calculate the mean bedding for each sill (Table S2). The thickness of the sampled sills
272 is typically over ten meters. Each sill was drilled evenly across its thickness to average the
273 characteristic remanent magnetization. Where measured close to the sampling sites, no significant
274 bedding difference was observed for the upper and lower contacts of each sill, indicating a
275 consistent emplacement of each sill along bedding at each of the sampling localities. Therefore,
276 we adopt a common mean bedding for samples collected from the same sill. We measured the
277 orientation of the cores in the field using a combination of the magnetic compass and sun compass
278 when the sun was available. No significant difference is observed between the two orientation
279 methods.

280 **U-Pb geochronology**

281 Zircon separation from block samples was performed at Harvard University using standard
282 techniques. Zircon U-Pb geochronology was conducted in the Isotope Laboratory at NERC Isotope
283 Geoscience Laboratories (NIGL) at British Geological Survey. Separated zircon crystals were
284 annealed in a muffle furnace at 900 °C for 60 hours in quartz beakers. Selected zircon grains were
285 transferred to 3 ml Hex Savillex beakers. They were processed through ultrasonic bathing and
286 rinsing with 30% HNO₃ and then moved to 300 µl Teflon PFA micro-capsules to be leached in an
287 about 5:1 mix of 29 M HF and 30% HNO₃ for 12 hours at 180 °C. After removing the leachate,
288 the residues were rinsed again with 30% HNO₃ and 6 M HCl and spiked with the EARTHTIME
289 ²³⁵U–²³³U–²⁰⁵Pb tracer (57). Each zircon grain was then fully dissolved in about 120 µl of 29M HF
290 with a trace amount of 30% HNO₃ at 220 °C for 48 hours. The solutions were then dried, and
291 fluorites in the samples were converted to chlorides in 3 M HCl at about 180 °C overnight.
292 Standard HCl-based anion-exchange chromatographic procedures were performed to separate U
293 and Pb from each sample, then loaded on a single Re filament in a silica-gel–phosphoric acid
294 mixture. Isotope ratios of U and Pb were analyzed using the Thermo-Electron Triton TIMS system
295 in NIGL. U isotopes were measured in static Faraday mode or on a single SEM detector, depending
296 on the uranium contents. Pb isotopes were measured using the peak-hopping method with a single
297 SEM detector. Data reductions were processed in Tripoli (58). U-Pb concordia diagrams were
298 plotted using IsoplotR (59).

299 **Paleomagnetic and rock magnetic experiments**

300 Paleomagnetic and rock magnetic experiments were performed in Paleomagnetism and
301 Environmental Magnetism Laboratory at the China University of Geosciences, Beijing. Oriented
302 paleomagnetic cores were first sliced into cylindrical specimens of ~1–2.2 cm thickness using non-
303 magnetic, diamond-rimmed saw blades. Remanent magnetizations of all specimens were
304 measured with a 2G 755-4 K three-axis cryogenic magnetometer in a magnetic shield room with
305 a <300 nT residual field. Specimens were processed through stepwise thermal demagnetization in
306 either an ASC TD-48 or MMTDSC furnace that holds an internal residual field of <10 nT. Most

specimens were completely demagnetized in ~20–35 steps by up to ~580 °C. Temperature increment between each step varies between 20–40 °C in the low-temperature range (<500 °C) and 3–10 °C in the high-temperature range (500–580 °C). Some specimens were subject to small demagnetizing increments of 5–10 °C between 300–350 °C to characterize better the removal of low-temperature components in this interval. To identify the carriers of magnetic remanences, temperature-dependent low field susceptibility of representative samples (ground to fine powders) was measured in an argon atmosphere with an AGICO KLY-4S Kappabridge with CS-3 high-temperature device. Heating and cooling curves were acquired between room temperature and 700 °C.

316 **Paleomagnetic analyses and calculations**

Remanence components were fit using linear principal component analysis (60). Unless specified, we force the best-fit lines to pass through the origin when fitting the high-temperature components. Only specimens with ≥ 4 stable points that decay to the origin with MAD (maximum angle of deviation) smaller than 20° in the high-temperature interval are regarded as credible and used to calculate specimen directions. Specimen directions were then averaged using Fisher statistics (61) to calculate the site mean direction of each sill/site of sedimentary rocks. Specimen analysis was done using the Paleomag OS X program (62). Paleomagnetic analyses and plotting were completed in Python facilitated by the PmagPy package (63).

To quantify the rate of pole shift between a pair of poles, we incorporate full error propagation in the calculation considering uncertainties in both the ages and positions of the poles. Given the great-circle distance [$\Delta\theta$ (°)] and age difference [Δt (Myr)] between two poles, a and b , the mean rate of pole shift (° Myr⁻¹) with 2σ error is calculated as follows:

$$329 \quad \text{rate} = \frac{\Delta\theta}{\Delta t} \left(1 \pm \sqrt{\frac{\sigma_{(a)}^2 + \sigma_{(b)}^2}{(\Delta t)^2} + \frac{A_{95(a)}^2 + A_{95(b)}^2}{(\Delta\theta)^2}} \right) \quad (1)$$

330 Where A_{95} is the radius of the 95% confidence cone of the pole; σ is the 2σ uncertainty on the pole
331 age.

332 We take a Monte Carlo approach to further constrain the range of pure TPW rates between 832
333 and 821 Ma inferred from the Fanjingshan and Xiaofeng poles. From Fisher distributions, a large
334 number (10,000) of random virtual geomagnetic poles (VGPs) are simulated for each
335 paleomagnetic pole according to their pole positions and associated uncertainties (Table S3),
336 paired with random draws from Gaussian distributions for the pole ages (831.51 ± 0.32 Ma and
337 821.64 ± 0.2 Ma). Simulated VGPs from each pole group are paired randomly, connected by grey
338 arcs shown in Fig. 2A.

339 To account for the potential pole shift produced by tectonic movement and disentangle the pure
340 TPW component, we simulate tectonic movements of South China between 832–821 Ma and
341 subtract them from the total pole motion. The direction and velocity of tectonic movements of
342 South China during this period are currently unclear. As such, for simulating the tectonic speed,
343 we adopt a triangular distribution (fig. S15) assuming:

- 344 1. The lower speed limit at 0 cm yr⁻¹.
- 345 2. The mode at 10 cm yr⁻¹ that corresponds to the fastest plate velocity of South China
346 extracted from the Neoproterozoic plate kinematic models by (55).

347 3. The upper speed limit at 31 cm yr^{-1} that represents the possible fastest tectonic movement
348 implied since the late Mesoproterozoic (56).

349 Unconstrained directions of South China tectonic movements are assumed to accommodate the
350 largest effect of tectonic motions. We take 10,000 random draws from the tectonic speed
351 distribution described above (fig. 15), and directions of the tectonic motions are randomly sampled
352 from $0\text{--}360^\circ$ from uniform distributions. These simulated tectonic speeds paired with random
353 directions are multiplied by the simulated age differences to obtain the distances of tectonic
354 movements between 832 and 821 Ma. The series of tectonic movements are then randomly applied
355 to the sampled Fanjingshan poles in Fig. 2A to convert the effect of plate motions to its
356 paleomagnetic expression. Fig. 2B shows the displaced Fanjingshan poles after correcting for
357 simulated tectonic motions, leaving the new angular difference between each pole pair an estimate
358 of the pure TPW component. Finally, the pure TPW angular distance is divided by the age
359 difference between the pole pair to estimate the pure TPW rate (Fig. 2C). The 2.5 and 97.5
360 percentiles are taken from the rate distribution, interpreted as the 95% confidence intervals (Fig.
361 2C). Accounting for tectonic motions yields a larger spectrum of TPW rate estimations compared
362 with assuming zero plate displacement (Fig. 2C). The expanded upper tail corresponds to scenarios
363 where the total pole motion is produced by TPW counteracted by tectonic motion, resembling the
364 case interpreted for South China between 815 and 800 Ma (12) and North China in the late Jurassic
365 (54). Whereas the reduced lower tail corresponds to TPW and tectonic motion moving along
366 similar directions and acting additively, similar to the interpretation for Laurentia between ca.
367 1,110 and 1,080 Ma (56).

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698 **Supplementary Materials:**

699 Supplementary Text
700 Figs. S1 to S15
701 Tables S1 to S8