

1        **Array-based iterative measurements of SmKS travel**  
2        **times and their constraints on outermost core structure**

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7                                **Key Points:**

- 8                                • We develop an array-based iterative method to measure SmKS-SKKS (m=3-5)  
9                                differential travel times.
- 10                               • 3D mantle structure effects must be considered in studies of SmKS differential travel  
11                               times.
- 12                               • Our measurements support a low  $V_p$  at the top of outer core.

## Abstract

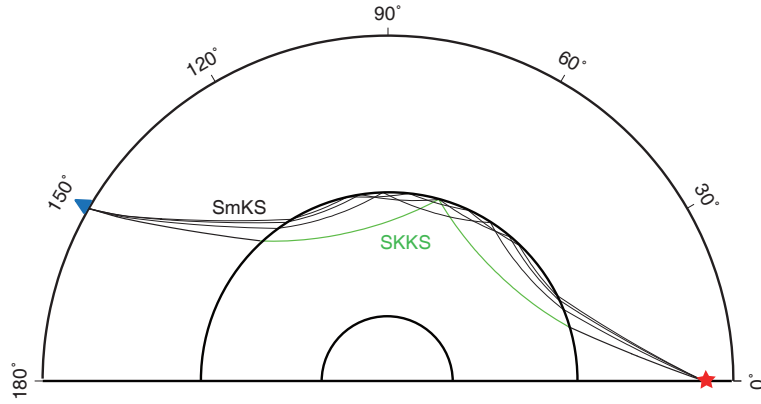
Vigorous convection in Earth’s outer core led to the suggestion that it is chemically homogeneous. However, there is increasing seismic evidence for structural complexities close to the outer core’s upper and lower boundaries. Both body waves and normal mode data have been used to estimate a P-wave velocity,  $V_p$ , at the top of the outer core (the E’ layer), which is lower than that in the Preliminary Reference Earth Model. However, these low  $V_p$  models do not agree on the form of this velocity anomaly. One reason for this is the difficulty in retrieving and measuring SmKS arrival times. To address this issue, we propose a novel approach using data from seismic arrays to iteratively measure SmKS-SKKS differential travel times. This approach extracts individual SmKS signal from mixed waveforms of the SmKS series, allowing us to reliably measure differential travel times. We successfully use this method to measure SmKS time delays from earthquakes in the Fiji-Tonga and Vanuatu subduction zones. SmKS time delays are measured by waveform cross-correlation (CC) between SmKS and SKKS and the CC coefficient allows us to access measurement quality. We also apply this iterative scheme to synthetic SmKS seismograms to investigate the 3D mantle structure’s effects. The mantle structure corrections are not negligible for our data and neglecting them could bias the  $V_p$  estimation of uppermost outer core. After mantle structure corrections, we can still see substantial time delays of S3KS, S4KS and S5KS, supporting a low  $V_p$  at the top of Earth’s outer core.

## 1 Introduction

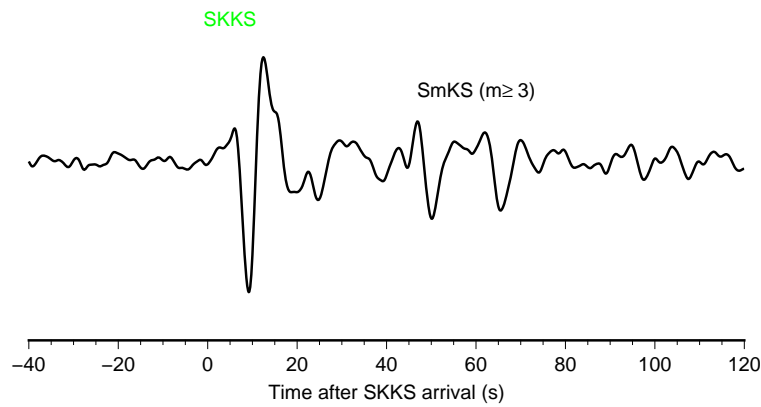
The liquid outer core in the Earth plays a critical role in the geodynamo and in thermochemical interactions between the mantle and core. Seismic studies can provide important constraints on the physical properties of the core and therefore improve our understanding of the composition and state of the core (Hirose et al., 2013). Due to vigorous convection, the bulk of the outer core is believed to be well mixed and therefore chemically homogeneous (Stevenson, 1987). However, there is increasing seismic evidence for structural complexities close to its top and bottom boundaries. A stratified layer with a lower  $V_p$  gradient than the Preliminary Reference Earth Model (PREM; Dziewonski & Anderson, 1981), labeled the F-layer, has been documented using body seismic wave observations (Souriau & Poupinet, 1991b; Song & Helmberger, 1995; Zou et al., 2008; Ohtaki & Kaneshima, 2015). Another stratified layer, the E’ layer, is hypothesized to

45 exist at the top of outer core and its properties may be constrained by geomagnetic sec-  
46 ular variations (Gubbins, 2007; Buffett, 2014), but the seismic evidence, especially SmKS  
47 differential arrival times, for this layer is contradictory and controversial (e.g. Eaton &  
48 Kendall, 2006; Alexandrakis & Eaton, 2010; Helffrich & Kaneshima, 2010; Kaneshima  
49 & Helffrich, 2013).

50 SmKS waves ( $m=1, 2, 3, \dots$ ) travel as S-waves in the mantle, are converted to com-  
51 pressional waves entering the outer core, reflected  $m-1$  times on the underside of the core-  
52 mantle boundary (CMB), and reconvert to S-waves to travel through the mantle (Fig.  
53 1a). SmKS waves are sensitive to the structure of outer core and their arrival times have  
54 been used to investigate  $V_p$  (compressional wave velocity) in the shallow outer core (Choy,  
55 1977). SKS absolute arrival times have a large scatter, especially due to 3D mantle struc-  
56 ture (e.g. Garnero et al., 2016), which results in large uncertainties in their constraints  
57 on outer core structure. SKKS and SKS have similar raypaths near the source, so their  
58 differential arrival times can partially remove the source effects and constrain the  $V_p$  of  
59 shallow outer core better. Hales & Roberts (1971) compiled SKKS-SKS differential ar-  
60 rival times and found a low  $V_p$  in the outermost core. However, the reliability of this  
61 study is reduced by the uncorrected phase shifting between SKS and SKKS (Choy & Richards,  
62 1975; Choy, 1977).



(a)



(b)

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65 **Figure 1:** Ray paths of SmKS waves and one example of SmKS waveforms. (a) Ray paths of  
66 SmKS. The red star is an earthquake and the blue triangle represents a seismic station. The  
67 green line shows the ray path of SKKS traveling in the outer core. The black lines are ray paths  
68 of SmKS ( $m=3, 4$  and  $5$ ) and sections of SKKS ray path traveling in the mantle and crust. (b)  
69 A band-pass filtered (bp 0.05-0.7 Hz) seismogram of SmKS data from station ASSE with an  
70 epicentral distance of  $150.5^\circ$  from the event #110729 (Table S1). Time zero is the SKKS arrival  
71 predicted by PREM. The predicted arrival time of S3KS is 39 s after SKKS. S4KS arrives at 52 s  
72 and S5KS is only 5 s after the S4KS.

74 Although the ray paths of SKS and SKKS are close to each other near the source,  
75 they diverge further in the lower mantle, where lateral heterogeneities could affect their  
76 different travel times (Garnero et al., 1988; Souriau & Poupinet, 1991a). Compared to

77 SKKS and SKS, SmKS and S(m-1)KS with  $m > 2$ , e.g. S3KS-SKKS, have closer raypaths  
 78 (Fig. 1a) and therefore their differential arrival times are less affected by 3D mantle struc-  
 79 tures. With the high quality seismic data accumulated in the last few decades, many more  
 80 observations of SmKS ( $m \geq 2$ ) waves has been reported and their differential travel times  
 81 have been used to investigate the stratification of the top outer core. However, the con-  
 82 clusions of various studies are not consistent. For example, Alexandrakis & Eaton (2007)  
 83 exploited the Empirical Transfer Function (ETF) technique to precisely measure SmKS  
 84 differential travel times and found no evidence for stratification, consistent with some  
 85 other SmKS studies (e.g. Souriau & Poupinet, 1991a; Alexandrakis & Eaton, 2010). In  
 86 contrast, other reports support a layer with lower  $V_p$  than that of PREM in the outer-  
 87 most core (e.g. Garnero et al., 1993; Tanaka, 2004; Eaton & Kendall, 2006; Tanaka, 2007;  
 88 Helffrich & Kaneshima, 2010; Kaneshima & Helffrich, 2013; Tang et al., 2015; Kaneshima  
 89 & Matsuzawa, 2015; Kaneshima, 2018), although the thickness and amplitude of the  $V_p$   
 90 anomaly varies from one study to another.

91 There are at least two reasons for the preceding contradictory results. The first one  
 92 is the difficulty in extracting each individual SmKS phase and precisely measuring the  
 93 differential arrival times. For high orders  $m \geq 3$ , SmKS series constitute a whispering-  
 94 gallery mode and consecutive SmKS phases have very close arrival times (e.g. S4KS and  
 95 S5KS in Fig. 1a), which makes separating consecutive SmKS waveforms difficult. An-  
 96 other problem is contamination from lateral heterogeneities in mantle structure. Although  
 97 ray paths of SmKS and S(m-1)KS ( $m > 2$ ) series are closer to each other than that of SKS  
 98 and SKKS, there are still differences in the mantle, especially the heterogeneous  $D''$  re-  
 99 gion (Garnero & Helmberger, 1995). These mantle heterogeneities could cause large un-  
 100 certainty or bias in the differential arrival time measurements made using individual seis-  
 101 mograms (Garnero et al., 1993) or small-aperture arrays (Eaton & Kendall, 2006). Stack-  
 102 ing of data from large-scale arrays (Helffrich & Kaneshima, 2010; Kaneshima & Helffrich,  
 103 2013; Kaneshima & Matsuzawa, 2015; Kaneshima, 2018) or global networks (Alexandrakis  
 104 & Eaton, 2010) tends to average out perturbations due to mantle heterogeneities and  
 105 therefore mitigate the possible bias. Alternatively, the bias can be evaluated using ray  
 106 theory (e.g. Helffrich & Kaneshima, 2010; Kaneshima & Helffrich, 2013; Kaneshima &  
 107 Matsuzawa, 2015; Kaneshima, 2018) or sophisticated waveform modeling (Tanaka, 2004,  
 108 2007), based on either known 3D mantle tomography model or hypothesized structure.

109 Combining array stacking and accurate 3D mantle corrections would be an optimal so-  
 110 lution to suppress 3D mantle effects, which has not been reported before.

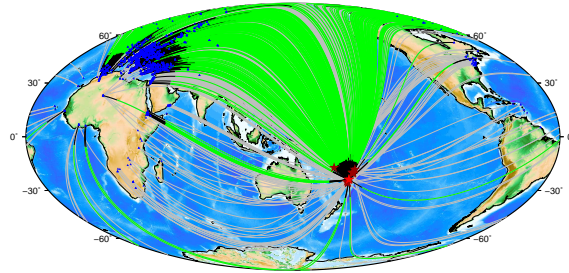
111 To ameliorate these problems, we develop an iterative method to separate individ-  
 112 ual SmKS phases from the SmKS wavetrain in array data and use normalized cross-correlation  
 113 (CC) to measure the differential travel times between SmKS ( $m=3, 4$  and  $5$ ) and SKKS.  
 114 We carefully select good quality data to successfully obtain each SmKS phase. The it-  
 115 erative method provides us with accurate waveform-based measurements of differential  
 116 arrival times and important information to assess the measurement quality. We use two  
 117 methods, ray theory and the Spectral Element Method (SEM), to investigate the effects  
 118 of lateral heterogeneities in the mantle, using the 3D tomography model S40RTS (Rit-  
 119 sema et al., 2011), and also assess the effect of choosing a different mantle model (S362ANI  
 120 Kustowski et al., 2008). The measured differential arrival times, after correction for 3D  
 121 mantle structure effects, are compared to the predictions of body-wave derived model  
 122 KHOMC (Kaneshima & Helffrich, 2013) and normal-mode constrained model EPOC (Irv-  
 123 ing et al., 2018).

## 124 **2 Data**

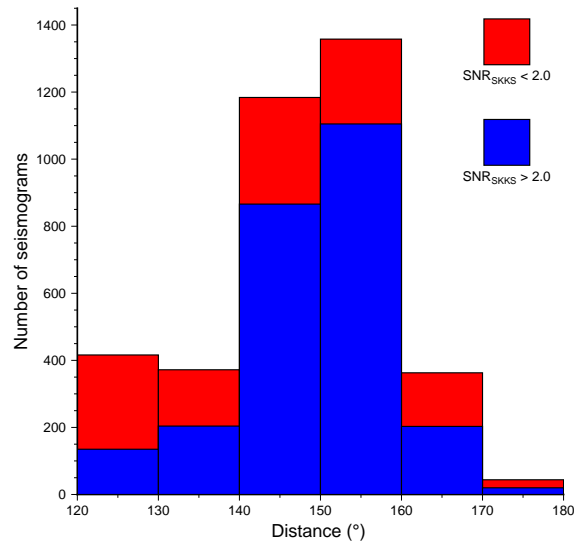
125 We collected more than 320,000 seismograms from global stations from 500 earth-  
 126 quakes in the subduction zones of Fiji-Tonga, Vanuatu, New Britain and Solomon with  
 127 depths  $\geq 150$  km and  $M_w \geq 5.5$  (Global Centroid-Moment Tensor catalog, Ekström et al.,  
 128 2012) in the period 2000-2016 (Supporting Information Fig. S1). We select events with  
 129 depths  $\geq 150$  km to avoid contamination from depth phases sSmKS ( $m \geq 2$ ). The seis-  
 130 mograms have a distance range of  $120-180^\circ$ , where waveforms SmKS ( $m \geq 2$ ) are read-  
 131 ily observed.

132 We remove instrument responses and rotate the two horizontal components to get  
 133 the radial displacement, on which SmKS primarily appears. Then a band-pass filter ( $0.05$   
 134 -  $0.7$  Hz) is applied to the data with Signal-Noise Ratio (SNR) computation. From these  
 135 500 earthquakes, we find 11 events with a large number of good observations of SKKS  
 136 (Fig. 2a). Here, good observation means SNR larger than 2, a large number means 100  
 137 or more seismograms, and we carefully inspect the data to rule out any possible contam-  
 138 ination from small local earthquakes. The SNR is defined as the peak-to-peak amplitude  
 139 ratio of SKKS to noise. We measure SKKS amplitude in a time window between 20 s

140 before and 50 s after the SKKS arrival time predicted by PREM (Fig. 3a). The time win-  
 141 dow of noise is taken between 70 s and 20 s before the SKKS arrival. There are total 3741  
 142 radial components from these 11 events and 2535 of them have  $SNR_{SKKS} > 2.0$  (Fig.  
 143 2b). Limited by the geographic distribution of seismic stations, most of these clear SKKS  
 144 data are from stations in Europe with a distance range of  $140^\circ$ - $160^\circ$  and their ray paths  
 145 sample the northeastern Pacific, Asia and Europe.

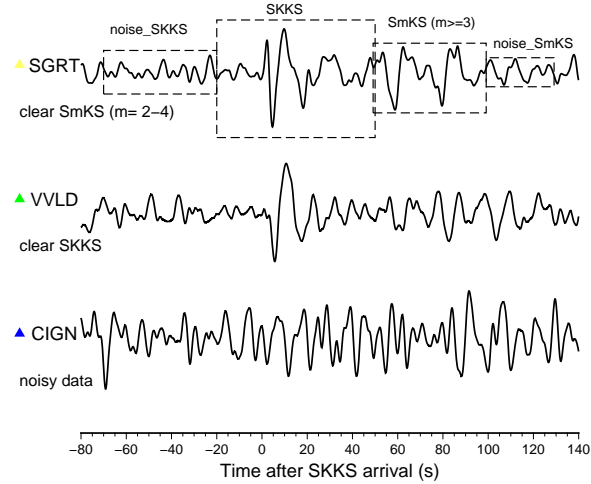


(a)

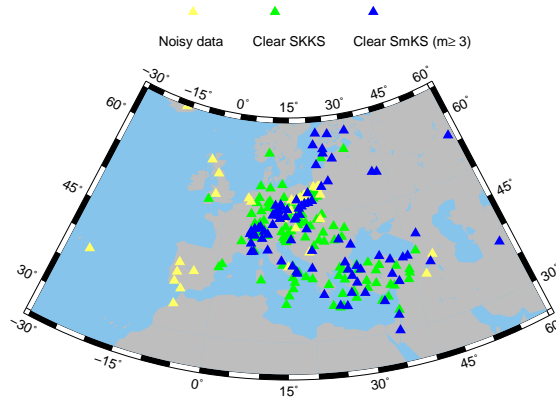


(b)

148 **Figure 2:** Map and histogram of clear SmKS data. (a) Map of good SKKS data ( $SNR \geq 2.0$ )  
 149 from the ten earthquakes. The blue triangles and red stars show the stations and earthquakes, re-  
 150 spectively. The lines connecting stations and earthquakes are ray paths of SKKS. The green lines  
 151 show the ray paths of SKKS traveling in the outer core from the event #110729. (b) Histogram  
 152 of SKKS data in (a). The blue bar portions correspond to SKKS data with high  $SNR \geq 2.0$  and  
 153 the red bar portions show the ones with  $SNR < 2.0$ .



(a)



(b)

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**Figure 3:** SmKS data from stations in Europe from event #071016. (a) An example of waveforms with poor SKKS with  $\text{SNR} < 2.0$  (bottom panel), good SKKS with  $\text{SNR} \geq 2.0$  only (middle panel), and high SNRs for both SmKS ( $m=3$  and higher) and SKKS (top panel). The SNR of SKKS is defined as the peak-to-peak amplitude ratio of SKKS (20 s before to 50 s after the SKKS arrival predicted by PREM) to that of the noise (70 s to 20 s before the SKKS arrival). Similarly, the SNR of SmKS ( $m \geq 3$ ) is obtained by measuring SmKS signals (0 s to 50 s after S3KS arrival) and the associated noise (50 s to 85 s after S3KS arrival). The time zero is the SKKS arrival predicted by PREM. (b) Map of stations in Europe from event #071016. Stations with noisy SKKS, good SKKS only and high SNRs for both SmKS ( $m=3$  and higher) and SKKS are shown as yellow, green and blue triangles respectively.



168 Following previous studies (e.g. Tanaka, 2004; Eaton & Kendall, 2006; Helffrich &  
 169 Kaneshima, 2010; Alexandrakis & Eaton, 2010; Kaneshima, 2018), we use SKKS as a  
 170 reference phase to investigate the arrivals of SmKS ( $m>2$ ), so clear SmKS ( $m>2$ ) sig-  
 171 nals are also important for high quality measurements. We compute SNR of SmKS ( $m>2$ )  
 172 and only use the data with clear SmKS ( $SNR_{SmKS} \geq 2.0$ , see Fig. 3). In contrast to  
 173 the  $SNR_{SKKS}$  computation, we take the noise window starting after the predicted S2KS  
 174 arrival time for  $SNR_{SmKS}$  (by 100 seconds) and some *SmKS* coda waves are included  
 175 in this time window. Thus, the data with strong *SmKS* coda due to significant unwanted  
 176 source and wave propagation complexities would have low  $SNR_{SmKS}$  and therefore be  
 177 discarded. Then, we use the method described in section 3 to measure these data with  
 178 clear SmKS ( $m \geq 2$ ). Most of our clear data are from Europe and our array-based method  
 179 needs a number of records to form an array, so here we focus on stations in Europe and  
 180 north Africa to investigate the SmKS arrivals.

### 181 **3 Array-based iterative method to measure SmKS-SKKS differential** 182 **arrival times**

#### 183 **3.1 Workflow of the array-based iterative method**

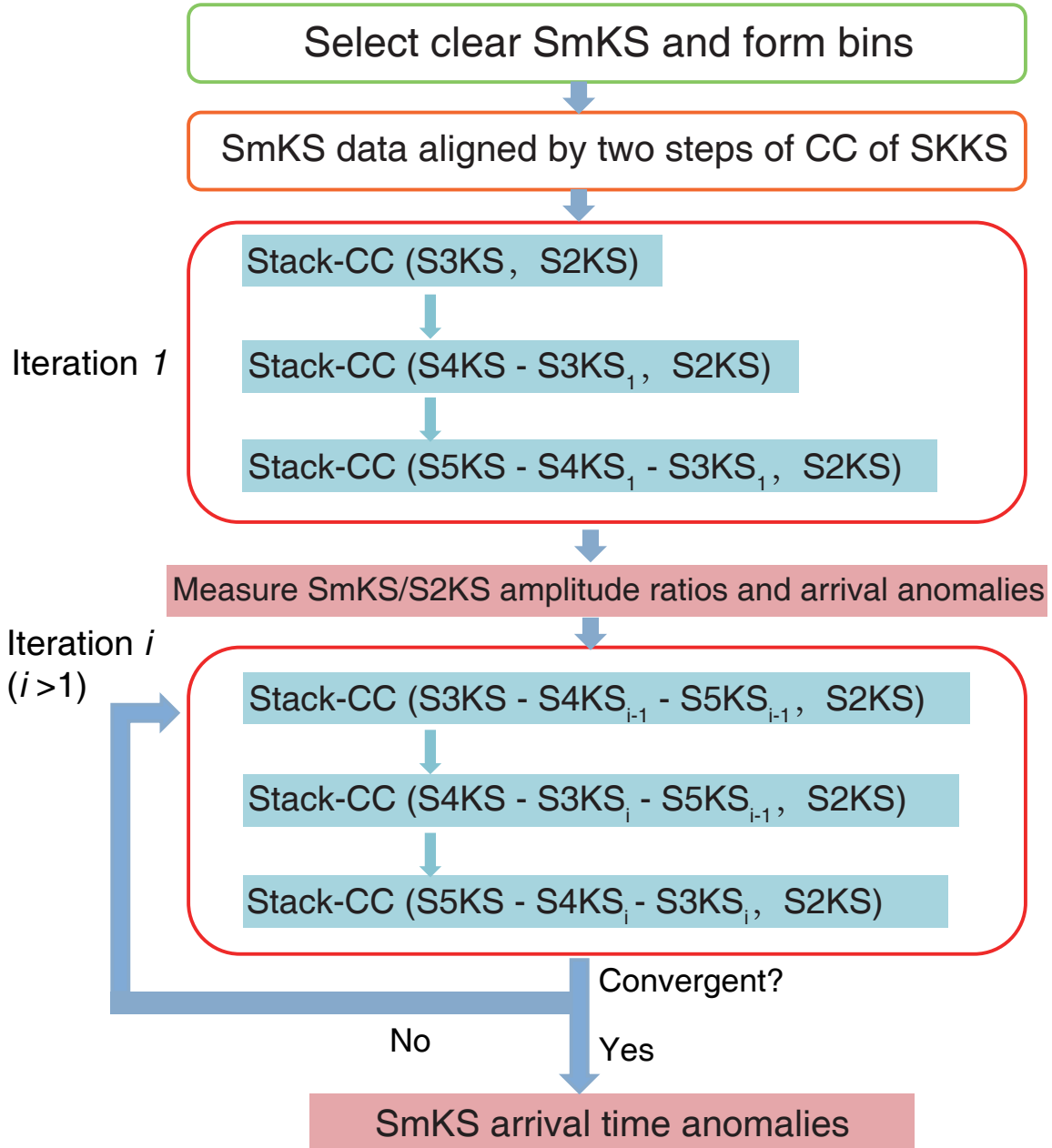
184 SmKS ( $m \geq 2$ ) series travel in the mantle and upper outer core, so their arrivals are  
 185 sensitive to the Vs in the mantle and Vp in the outer core. The ray paths of SKKS and  
 186 SmKS ( $m>2$ ) are close to each other in the mantle and further apart in the outer core  
 187 (Fig. 1a), so taking arrival time differences between SKKS and SmKS ( $m>2$ ),  $t_{SmKS}-$   
 188  $t_{SKKS}$ , instead of absolute travel time, can significantly reduce the effects of 3D Vs struc-  
 189 ture in the mantle and improve the constrains on the Vp in outer core. On the other hand,  
 190 these spatially close ray paths result in small time separations between consecutive SmKS  
 191 signals, which can make identifying individual SmKS phase and measuring its arrival time  
 192 difficult. For example, the arrival time difference between S3KS and S4KS at station ASSE  
 193 from event #110729 is only 13 s (Fig. 1a). The difference between S4KS and S5KS is  
 194 even smaller and their waveforms are mixed with each other. Many previous efforts have  
 195 been made to retrieve individual SmKS phase and accurately measure their arrival times  
 196 (e.g. Eaton & Kendall, 2006; Helffrich & Kaneshima, 2010; Kaneshima & Helffrich, 2013).  
 197 In particular, array stacking techniques have been used to analyze slownesses and arrival  
 198 times of SmKS signals (e.g. Eaton & Kendall, 2006; Helffrich & Kaneshima, 2010; Kaneshima  
 199 & Helffrich, 2013). Here, we take the advantage of the large number of stations with good

200 data to form one or more arrays or bins and develop an iterative method to retrieve in-  
 201 individual SmKS and measure their arrival times. This iterative strategy has been used  
 202 to extract direct S-waves and CMB reflected ScS waves (Z. Yu et al., 2012).

203 Arrivals in the SmKS series share many factors, such as source time function, 3D  
 204 wave propagation effects, site responses etc., due to their similar ray paths in the crust  
 205 and mantle. Although their ray paths diverge further in the outer core, the outer core  
 206 is believed to be highly laterally homogeneous. Thus, SKKS and SmKS ( $m > 2$ ) usually  
 207 have very coherent waveforms (after a  $\pi(m-2)/2$  phase shift is applied to SmKS with  
 208  $m=3, 4$  and  $5$ ). This property helps us significantly simplify the problem and separate  
 209 individual SmKS waveform. In our iterative method, the reference phase SKKS is as-  
 210 sumed to be perfectly coherent with each SmKS ( $m > 2$ ) waveform after the phase-shift  
 211 is applied and only two unknown parameters, SmKS arrival time anomalies and SmKS/SKKS  
 212 amplitude ratios, are measured in each iteration. We note that another alternative mea-  
 213 suring strategy would be attempting to measure SmKS-S( $m-1$ )KS (i.e. S3KS-SKKS, S4KS-  
 214 S3KS and S5KS-S4KS), which have even closer raypaths than those of SmKS-SKKS ( $m=3,$   
 215  $4$  and  $5$ ). However, this strategy suffers from the problem of weak and noisy reference  
 216 phases S3KS and S4KS, which would affect the performance of our method. Thus, we  
 217 choose the clearer SKKS waveforms as the reference phase.

218 This workflow of our iterative method is composed of data preparation and then  
 219 iterative measuring (Fig. 4). As described in subsection 2, we set an SNR threshold of  
 220 2 for both SKKS and SmKS ( $m > 2$ ) to obtain good quality data. Following (Helffrich &  
 221 Kaneshima, 2010; Kaneshima & Helffrich, 2013), we divide the clear SmKS data from  
 222 the same event into several bins and stack traces in each bin to further improve the SNR  
 223 (an example of one bin is shown in Fig. 5). Before stacking the traces, two steps of CC  
 224 are carried out on SKKS waveforms to align the data. In the first step of CC, we choose  
 225 one typical trace (i.e. station with the median distance of the bin) as a template (e.g.  
 226 black line in Fig. 5b) and compute CC of SKKS between this template and other traces  
 227 in this bin with shifting times. Then these traces are aligned on the time with the max-  
 228 imum CC values. In the next step, we stack the aligned SKKS with normalized ampli-  
 229 tudes to form a new template (e.g. red line in Fig. 5b) and then repeat the CC process-  
 230 ing to align the SmKS data again (Fig. 5c). The time window of SKKS used in CC is  
 231 5 s before and 30 s after the arrival time of SKKS and the maximum allowed time shift  
 232 is 5 s. Data with maximum CC coefficients lower than 0.8 are not used in the following

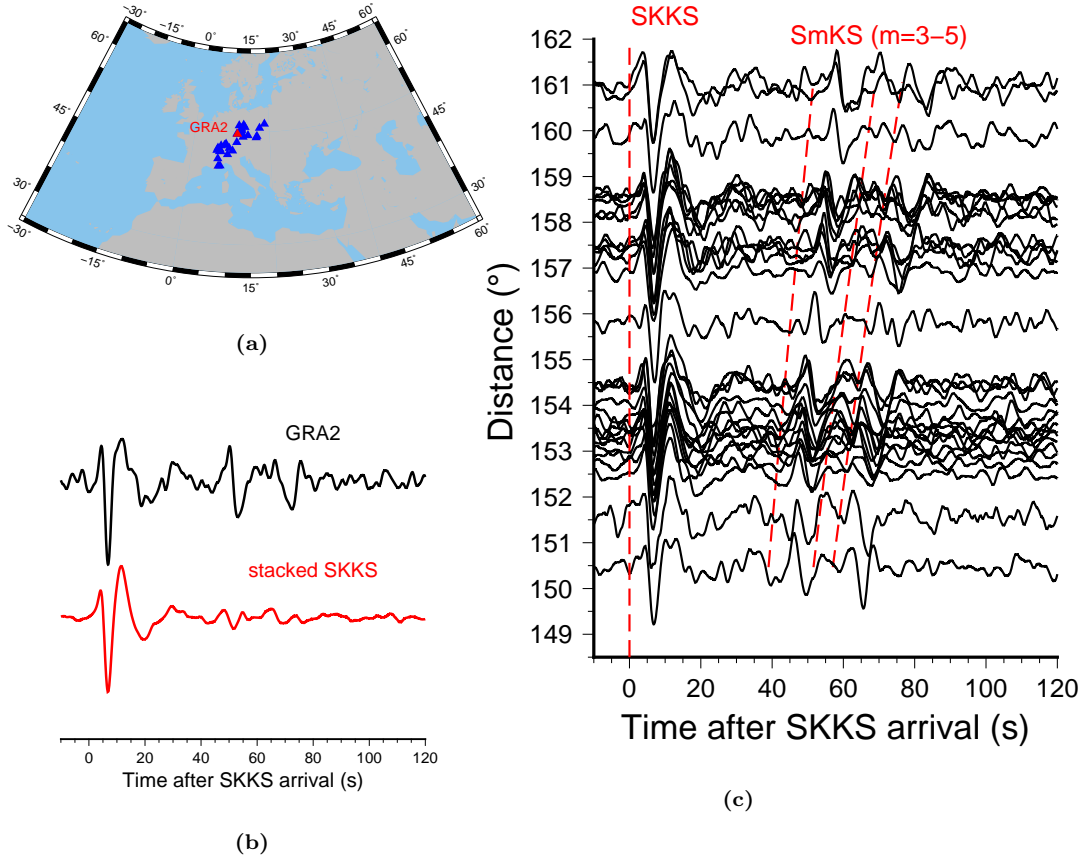
233 iterative measuring, because their low waveform similarities, due to complex site struc-  
 234 ture or/and instrumental issues, could decrease the quality of stacking and affect the mea-  
 235 surements. In these two steps of alignment, the shifted times are primarily due to 3D  
 236 structures near the stations, source mislocation and/or clock time errors and these fac-  
 237 tors are shared by SKKS and SmKS ( $m > 2$ ). Thus, shifting the traces are not expected  
 238 to significantly affect the measurements of differential arrival times.



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**Figure 4:** Workflow of the array-based iterative method.



242 **Figure 5:** An example of aligning SKKS by two steps of cross-correlation. (a) A map of a bin  
 243 of stations with clear SmKS from the event #071016. The other stations with clear SmKS are  
 244 shown in Fig. 2b. (b) The SKKS waveforms from the reference station GRA2 (upper trace, epi-  
 245 central distance of  $154.4^\circ$ ) and stacked SKKS after alignment by CC with GRA2 (lower panel).  
 246 (c) Distance profile of SmKS data (0.05-0.7 Hz) aligned on SKKS by two steps of CC. The cor-  
 247 responding stations are shown in (a). The time zero is the SKKS arrival. The other red dashed  
 248 lines are the SmKS ( $m=3-5$ ) arrivals predicted by PREM.

250 Next, we use these aligned SmKS data to iteratively retrieve individual SmKS phases,  
 251 measure differential travel time anomalies and assess quality of each measurement. In  
 252 the first iteration (“*iteration1*” in Fig. 4), we stack the data in a bin and use three CC  
 253 processes to measure S3KS, S4KS and S5KS one by one. For S3KS measurement, we stack  
 254 S3KS using  $t_{S3KS} - t_{SKKS}$  predicted by PREM (e.g. see the second red dashed line Fig.  
 255 5c), apply the Hilbert transform on them to correct the  $90^\circ$  phase shift and then com-

256 pute CC between stacked SKKS and S3KS to get S3KS/SKKS amplitude ratio and time  
 257 delay of S3KS. Then we cut out the SKKS waveform at each station, scale them using  
 258 the previously measured S3KS/SKKS amplitude ratio and apply the phase shift to get  
 259 S3KS waveform estimation. This estimated S3KS is subtracted from the data to retrieve  
 260 a 'clean' S4KS and then a similar stack-CC processing is applied on the retrieved S4KS  
 261 for measurement. Once S3KS and S4KS have been measured, we can estimate both S3KS  
 262 and S4KS, remove them in the data and then measure S5KS. After *iteration1*, we ob-  
 263 tain initial estimations of SmKS/S2KS (m=3, 4 and 5) amplitude ratios and their time  
 264 delays. In the next iteration, these information are used to retrieve the target SmKS and  
 265 more accurately measure them. This iteration is repeated until the measurements are  
 266 convergent.

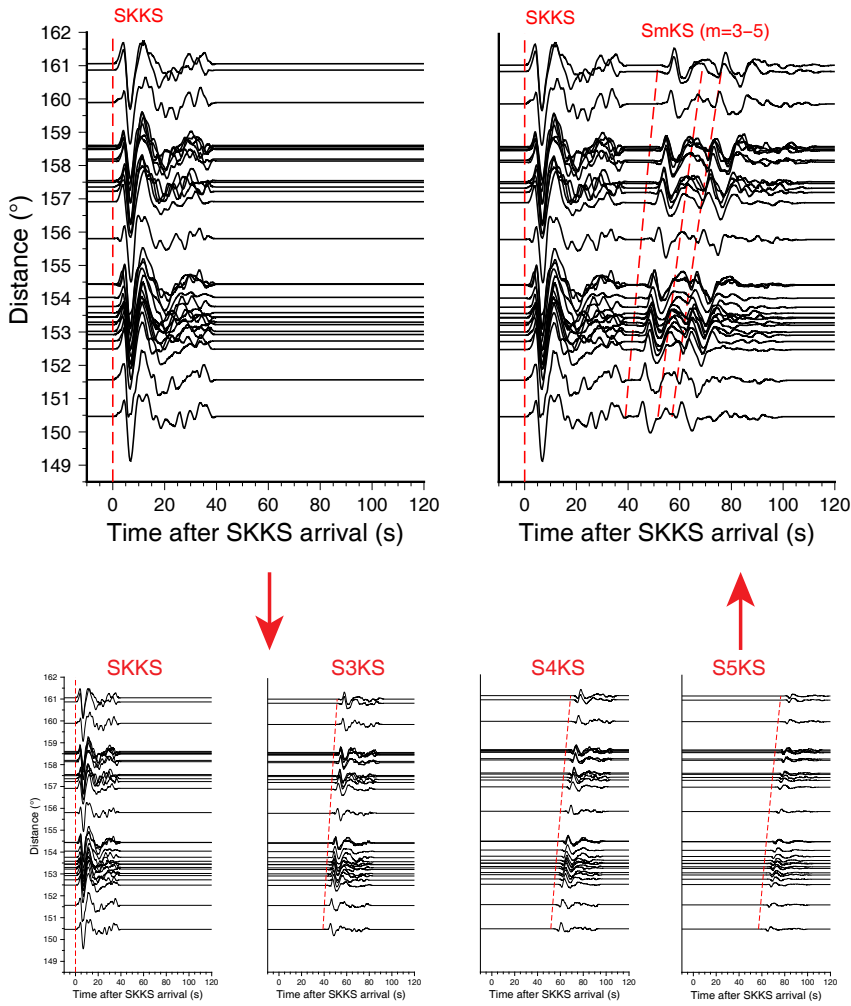
267 This array-based iterative method uses good quality data and has the advantages  
 268 of enhancing SNR by stacking and retrieving target SmKS signals well by removing other  
 269 SmKS interfering signals. Note that we use theoretical slowness derived from PREM to  
 270 stack array data, because Vp anomaly in the uppermost outer core only causes small slow-  
 271 ness deviation and slowness measurements could have large uncertainties. A large slow-  
 272 ness anomaly would result in less coherent stacking, which would be reflected in the CC  
 273 coefficient. In the first step of data preparation, we set strict criteria to rule out the data  
 274 with potential issues that might affect the validity of our method. For example, the re-  
 275 quirement of  $SNR_{SKKS} \geq 2.0$  allows us discard the data with high noise before SKKS.  
 276 In addition to that, the other two thresholds of  $SNR_{SmKS} \geq 2.0$  and  $CC \geq 0.8$  rule out  
 277 more bad quality data (e.g. complex SmKS waveforms and/or strong SmKS coda waves  
 278 due to 3D heterogeneity or source or station structures). Stacking the data with high  
 279 CC value further increases SNR and extracting individual SmKS phase from mixed sig-  
 280 nals allows us reduce uncertainties in measurements. More importantly, this method pro-  
 281 vides us two critical parameters to assess qualities of measurements. The most impor-  
 282 tant parameter is the CC values between S2KS and target SmKS (S3KS, S4KS and S5KS).  
 283 A low CC value means a bad quality measurement and we should either discard it, or  
 284 be careful when using it. Low CC values could be due to a failure of the assumptions  
 285 we made, weak target signals (e.g. near the nodal plane of radiation pattern of earth-  
 286 quake), insufficient number of traces in a bin etc. In addition to CC values, the ampli-  
 287 tude information is also helpful to assess measurement quality. More details are discussed  
 288 in section 5.

289           Uncertainty of differential arrival time for each bin is estimated by bootstrapping  
 290 (Efron & Tibshirani, 1991), which reflects the variance in the bin. For each bin, we ran-  
 291 domly select  $N$  seismograms, with replacement, from the original  $N$  seismograms and  
 292 measure the differential arrival times. This process is repeated 300 times and we com-  
 293 pute the standard deviation of these 300 measurements as an estimation of variance in  
 294 that bin.

295           In next section, we demonstrate the validation of our method by testing synthetic  
 296 seismograms and then apply it to data.

### 297           **3.2 Synthetic tests**

298           In this subsection, we cut real SKKS waveforms from data, use them to make SmKS  
 299 ( $m=2, 3, 4$  and  $5$ ) synthetics and then validate our iterative method. Fig. 5c shows SmKS  
 300 data of a bin from event #071016. We cut and taper the SKKS waveforms from 0 s to  
 301 40 s as input to generate S3KS, S4KS and S5KS (Fig. 6). S3KS is formed by scaling the  
 302 input signals with a prescribed S3KS/SKKS amplitude ratio of 0.42, applying a  $90^\circ$  phase  
 303 shift and a prescribed time shift, which is 1.13 s greater than to the PREM S3KS-SKKS  
 304 differential arrival time. Similarly, S4KS and S5KS are made with different amplitude  
 305 ratios and time delays. Then, complete SmKS synthetic seismograms are generated by  
 306 summing SKKS, S3KS, S4KS and S5KS.



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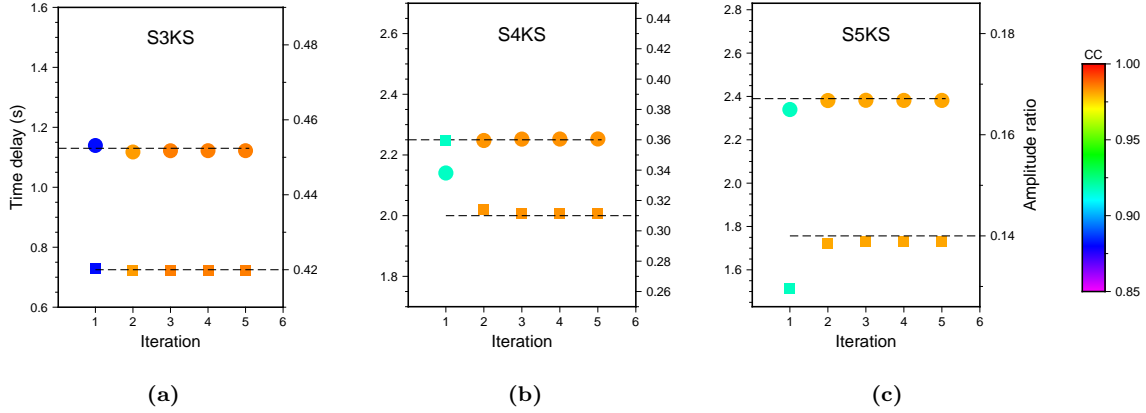
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**Figure 6:** Making SmKS ( $m=2-5$ ) synthetics using the SKKS data from the event #071016. The upper left figure shows the cut-out and tapered SKKS waveforms (complete data shown in Fig. 5c). Each sub-figure in the lower panel corresponds to the synthetics of individual SmKS phase. We take the tapered SKKS data (upper left figure), scale them using given amplitude ratios and apply the corresponding phase shift and time shift to form each SmKS phase. The S3KS/SKKS amplitude ratio is given as 0.42 and its time delay is 1.13 s. For S4KS, the amplitude ratio is 0.31 and the time delay is 2.25 s. For S5KS, the amplitude ratio is 0.14 and the time delay is 2.39 s. These SmKS phases are added together to form the complete synthetics of SmKS series (upper right figure). The dashed red line at the time zero in each figure is the SKKS arrival, the same as Fig. 5c. The other red dashed lines are the SmKS ( $m=3-5$ ) arrivals predicted by PREM.

320 Then we apply our iterative method to these synthetic seismograms and check its  
321 validity. In the step of searching maximum CC values, we take a time window of 0-30  
322 s after the target SmKS arrival time and the maximum allowed time shift is 5 s. In pre-  
323 vious studies, the time delays of SmKS ( $m=3, 4$  or  $5$ ) are less than 5 s and most of them  
324 are less than 3 s (e.g. Eaton & Kendall, 2006; Helffrich & Kaneshima, 2010; Kaneshima  
325 & Helffrich, 2013).

326 Fig. 7 shows the measurements from the first five iterations. We can see that both  
327 amplitude ratios and time delays are successfully retrieved and the CC values for S3KS,  
328 S4KS and S5KS are higher than 0.95 after the second iteration. In the first iteration,  
329 there are some differences between the measured results and true values. For example,  
330 the measured time delay of S4KS is  $\sim 2.15$  s, which is  $\sim 0.1$  s smaller than the input 2.25  
331 s. The CC value for S3KS measurement,  $CC_{3,2}$ , is 0.88, lower than  $CC_{4,2}=0.92$  for S4KS  
332 and  $CC_{4,2}=0.92$  for S5KS, because the S3KS measurement is affected by the presence  
333 of S4KS and S5KS signals. In the second iteration, the CC values are significantly in-  
334 creased and the measurements are close to the true values. The measurements become  
335 almost constant in the next three iterations, showing they reach convergence. After the  
336 first two iterations, waveforms are successfully retrieved and the time delays are accu-  
337 rately measured (e.g. see the waveform cross-correlations between SKKS and S5KS in  
338 Fig. S2). These results demonstrate the validation of our method. Of course, real data  
339 may be more complex than the synthetic SmKS here, e.g. different noise signals may be  
340 present in data, and therefore measurement quality might be not as good as in these syn-  
341 thetic tests. However, CC values indicate this complexity, demonstrating their impor-  
342 tance.





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**Figure 7:** A synthetic test to validate the array-based iterative method. The colored circles indicate the measured time delay in each iteration. The colored squares are the measured amplitude ratios. The upper black dashed line in each figure is the prescribed time delay and the lower black dashed line corresponds to the given amplitude ratio. Note that the time delays are relative to SmKS-SKKS differential arrival times predicted by PREM. The color represents the CC values between the single SmKS phase and transformed SKKS (e.g. results of S5KS from the first four iterations shown in Fig. S2). Both time delays and amplitude ratios of S3KS, S4KS and S5KS converge to the input values after five iterations and this is also reflected in the high CC coefficients.

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### 3.3 Correcting 3D mantle structure effects

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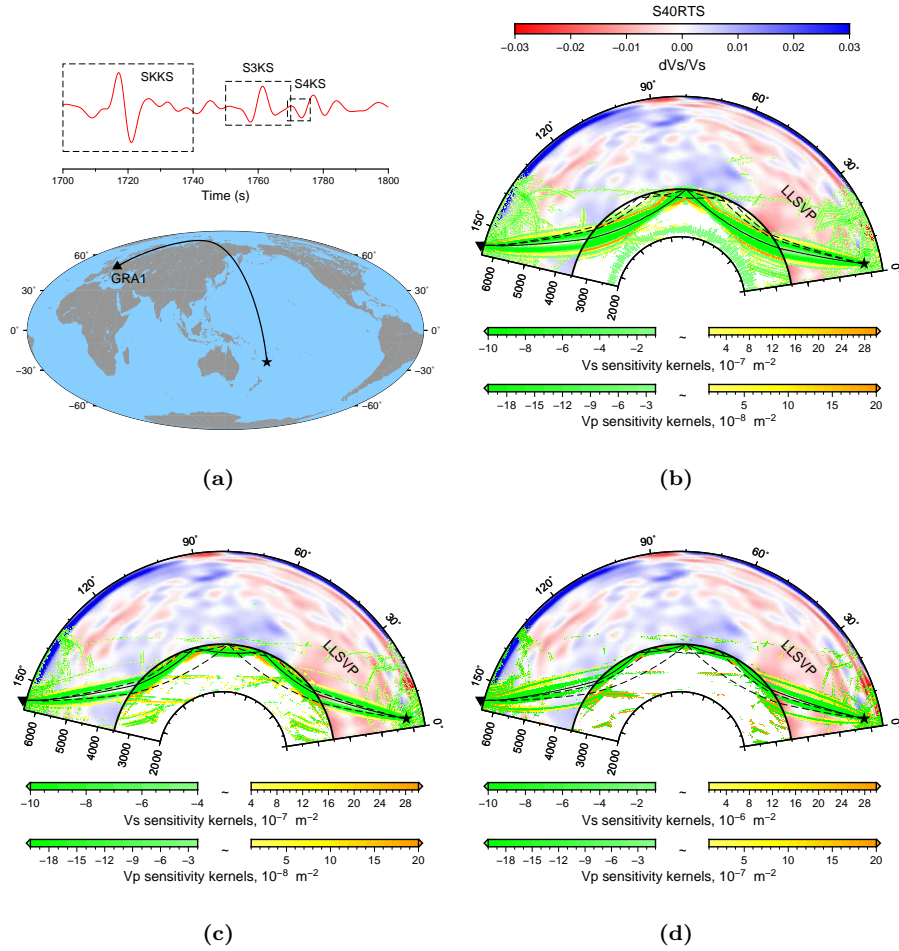
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Because the ray paths of SmKS ( $m=2-5$ ) are close to each other in the mantle, many previous studies assume that the effects of 3D mantle structures are the same for SKKS and SmKS ( $m>2$ ). Thus, the measured time delays of SmKS ( $m>2$ ) are only due to the  $V_p$  anomalies in the top outer core. However, we know that the ray paths between SKKS and SmKS ( $m>2$ ) are not exactly the same and the 3D mantle structures must affect the arrival time difference between SKKS and SmKS ( $m>2$ ). Kaneshima & Matsuzawa (2015) used ray theory to investigate these mantle effects at receiver-side and source-side. At the receiver side, they found that the mantle effects on  $dt_{3,2}$  are much less than 0.4 s. However, the presence of a Large Low Shear Velocity Province (LLSVP) beneath the Pacific could cause some time delays of SmKS ( $m\geq 2$ ) and affect the measurements.

365 To investigate 3D mantle effects, we use two different methods, ray theory and SEM,  
366 to compute the travel time delays of SmKS and compare their differences. We use SPECFEM3D\_globe  
367 to compute synthetic seismograms and evaluate the 3D mantle effects present in the to-  
368 mography models S40RTS (Ritsema et al., 2011) and S362ANI (Kustowski et al., 2008).  
369 As a spectral element method, the SPECFEM3D\_globe package solves the weak form  
370 of the seismic wave propagation equation and has the advantages of high accuracy, fast  
371 computation speed, handling discontinuity topography etc. (Komatitsch & Tromp, 1999,  
372 2002; Tromp et al., 2008). The adjoint source technique is part of SPECFEM3D\_globe,  
373 allowing the efficient computation of global scale sensitivity kernels of seismic signals in  
374 a given time window and frequency band (Tromp et al., 2008; Luo et al., 2013). We set  
375 the mesh parameters *NEX\_XI* and *NEX\_ETA* to 896 and the minimum resolved period  
376 is about 4.9 s. We use source parameters from GCMT (Global Centroid Moment Ten-  
377 sor, Ekström et al., 2012) and SPECFEM3D\_globe to compute the synthetic seismogram  
378 and the SmKS travel time sensitivity kernels. We note that GCMT solutions do not con-  
379 tain detailed inversions for source duration, so we reestimate the source duration using  
380 teleseismic P-waves from global stations. For comparison, we also compute the 3D man-  
381 tle structure corrections based on ray theory using the S40RTS model. To simplify the  
382 problem, we use PREM to get the ray path and compute the arrival time perturbations  
383 along that ray path. In other words, we assume that the ray path is not dramatically  
384 distorted by 3D structures.



387 **Figure 8:** Map and travel time sensitivity kernels of SmKS ( $m=2-4$ ) at station GRA1 from  
 388 event #110729. (a) SEM synthetic seismogram (upper panel) and map (lower panel). The red  
 389 line in the upper figure is the radial component synthetic seismogram of station GRA1 at 0.05-  
 390 0.2 Hz. The triangle in the map shows the location of station GRA1 and the star is the centroid  
 391 location of the event #110729. The black line in the map shows the great circle path of SmKS.  
 392 The arrival times predicted by PREM are 1717.0 s for SKKS, 1758.5 s for S3KS and 1772.1 s for  
 393 S4KS. The centroid time is 2.5 s, half of our re-estimated duration (Table S1), after the origin  
 394 time for this event. (b) Travel time sensitivity kernels of SKKS. Sensitivity to Vs is shown in the  
 395 mantle and to Vp in the core. The red-blue colors illustrate the depth cross-section of  $dV_s/V_s$   
 396 ( $V_s$  perturbation) of the 3D model S40RTS. The green-yellow colors show the travel time sen-  
 397 sitivity kernels of SKKS and its ray path is plotted with the black line. The dashed black lines  
 398 are the ray paths of S3KS and S4KS. (c) Travel time sensitivity kernels of S3KS. (d) Travel time  
 399 sensitivity kernels of S4KS.

401 Fig. 8 shows a depth cross-section of fractional velocity anomaly,  $dV_s/V_s$ , from the  
 402 3D model S40RTS (Ritsema et al., 2011) along the great circle connecting station GRA1  
 403 and event #110729. The ray paths of SmKS ( $m=2, 3$  and  $4$ ) sample the LLSVP at the  
 404 source side, where  $dV_s$  is lower than  $-1\%$ . Using this 3D  $V_s$  mantle model, we can use  
 405 ray theory to compute the arrival time anomalies of SmKS ( $m=2, 3, 4$  and  $5$ ) along their  
 406 ray paths. However, ray theory only works at infinite frequency. Indeed, seismic waves  
 407 at a finite frequency are sensitive to a Fresnel zone, a region centered at its ray path. To  
 408 demonstrate the Fresnel zones of SmKS, we use the SPECFEM3D\_globe package (Ko-  
 409 matitsch & Tromp, 1999) to compute sensitivity kernels of SmKS in the mantle and outer  
 410 core.

411 Fig. 8a shows the SPECFEM3D\_globe synthetic seismogram at GRA1 from the  
 412 event #110729 and three time windows used to compute the sensitivity kernels of SKKS,  
 413 S3KS and S4KS. We use the GCMT solution (Ekström et al., 2012) as the input of source  
 414 parameters, but reestimate its source duration (Fig. S4). The S40RTS model is used to  
 415 describe mantle heterogeneity and attenuation simulation is disabled to speed up the com-  
 416 putation. We use 1536 CPU cores to run the SEM simulation, taking about 14 hours for  
 417 forward modeling and 27 hours for each adjoint simulation. At frequency 0.05-0.2 Hz,  
 418 the first Fresnel zone of SKKS (the green band centered at SKKS ray path) has a width  
 419 of  $\sim 18$  deg ( $\sim 900$  km) on the CMB and its upper boundary approaches the ray paths  
 420 of S3KS and S4KS (Fig. 8). The sensitivity kernels of S3KS and S4KS have similar di-  
 421 mensions (i.e. the width of the first Fresnel zone), but more complex patterns than SKKS.  
 422 Compared to SKKS, S3KS and S4KS are more sensitive to the shallower outer core, re-  
 423 flected in the distribution of sensitivity kernels. The wide dimensions and complex pat-  
 424 terns of SmKS sensitivity kernels in Fig. 8 indicate that the 3D mantle structure cor-  
 425 rections based on ray theory may cause systematic biases and uncertainties. We will dis-  
 426 cuss the detailed 3D mantle structure correction of each bin and the comparison of ray  
 427 theory and SEM results in subsection 4.2.

## 428 4 Results

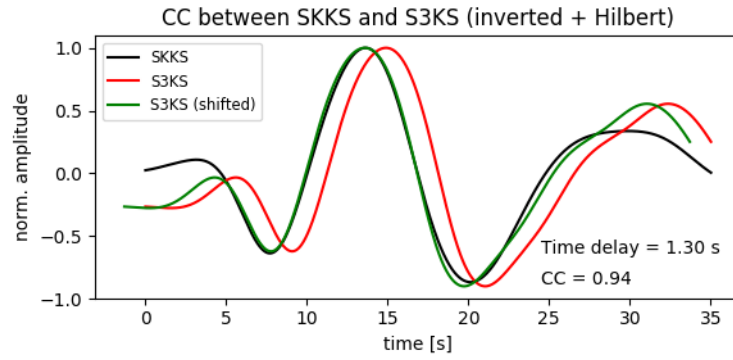
### 429 4.1 Measuring SmKS-S2KS differential arrival times

430 We apply the iterative method to data at three frequency bands (0.05-0.2 Hz, 0.05-  
 431 0.7 Hz and 0.1-0.7 Hz) and investigate the time delays of S3KS, S4KS and S5KS. For

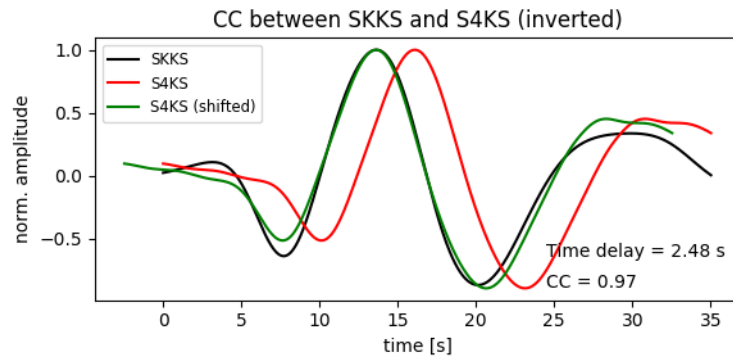
432 each frequency band, we compute the SNRs ( $SNR_{SKKS}$  and  $SNR_{SmKS}$ ), take clear SmKS  
 433 data of each event to form bins (one example of event #141101 shown in Fig. S3) and  
 434 apply the iterative method to each bin.

435 We only use data at epicentral distances greater than  $140^\circ$ . At shorter distances,  
 436 S3KS arrival times are close to SKKS (i.e. arrival time difference smaller than 27 s) and  
 437 therefore might affect quality of cut SKKS waveforms. Based on the number of clear SmKS  
 438 traces and the station distribution, we divide the data from each event into several ge-  
 439 ographical bins. For example, the event #141101 provides more than 100 clear SmKS  
 440 traces (0.05-0.2 Hz) and we divide them into four bins (see Table S1 and Fig. S3). For  
 441 some bins (e.g. bin 2 from event #010526 in Table S1), the number of clear SmKS traces  
 442 is too few (i.e.  $<10$ ) to provide reliable measurements, so we do not use the results of  
 443 these bins.

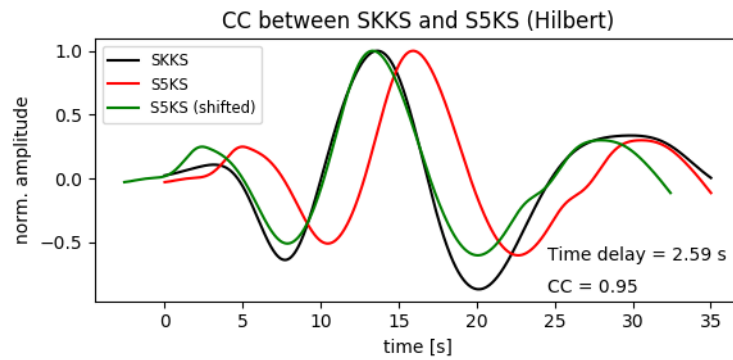
444 At frequencies 0.05-0.2 Hz, we eventually have twenty five effective bins from the  
 445 eleven events (Table S1). We use the same parameters (i.e. a time window of 40 s to cut  
 446 SKKS and 30 s for CC computation) as in synthetic testing and apply the iterative method  
 447 to each bin. In the synthetic testing, the measured results are almost constant after the  
 448 second iteration. Thus, here we conduct six iterations and take the results from the fifth  
 449 iteration (detailed measurements listed in Table S1). For each bin, we check the results  
 450 and make sure that there is no substantial difference between the fourth and fifth and  
 451 six iterations. In the twenty five bins, the measured S3KS time delay,  $dt_{3,2}$ , ranges from  
 452 -0.03 s to 2.83 s and the S3KS/SKKS amplitude ratios,  $A_{3,2}$ , are between 0.35-0.71. Nine-  
 453 teen bins have  $CC_{3,2} \geq 0.90$  and most of the time delays are positive values, except bin  
 454 1 from event #010428. S4KS and S5KS are more difficult to retrieve and measure. This  
 455 is reflected in the generally lower CC values and larger measurement scatter than S3KS.  
 456 Fig. 9 shows an example of bin 4 from the event #141101. All the three CC values are  
 457 higher than 0.94, indicating good quality measurements. For this bin, the measured time  
 458 delays are 1.30 s for S3KS, 2.48 s for S4KS and 2.59 s for S5KS. The median epicentral  
 459 distance of this bin is  $145.29^\circ$  and those time delays would indicate a slower  $V_p$  than in  
 460 PREM in the topmost outer core, consistent with previous studies (e.g. Eaton & Kendall,  
 461 2006; Helffrich & Kaneshima, 2010; Kaneshima & Helffrich, 2013).



(a)



(b)



(c)

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464

465 **Figure 9:** Time delays of S3KS, S4KS and S5KS measured on bin 4 data after five iterations  
466 from event #141101 (0.05-0.2 Hz). (a) CC between SKKS and S3KS (after Hilbert transform and  
467 polarity inverted). The black line is the stacked SKKS and the red line represents the stacked  
468 S3KS. The green line shows the shifted S3KS with the maximum CC value. The time shift be-  
469 tween the red line and green line is 1.30 s and the corresponding CC value is 0.94. Note that  
470 the time delay is relative to S3KS-SKKS differential arrival time predicted by PREM. (b) CC  
471 between SKKS and S4KS (polarity inverted). (c) CC between SKKS and S5KS (after Hilbert  
472 transform).

474 The measurement qualities are primarily indicated by their CC coefficients. In ad-  
 475 dition to CC coefficients, amplitude information is also useful to assess the measurement  
 476 quality. If other factors, such as source radiation pattern, are the same, the amplitude  
 477 of the SmKS phase decreases with its order  $m$ , due to the energy loss at each reflection  
 478 on the underside of the CMB. All the measurements with good quality at 0.05-0.2 Hz  
 479 follow this trend of  $A_{3,2} > A_{4,2} > A_{5,2}$  (amplitude information in Table S1 and the  
 480 good quality measurements are listed in Tables S2,S3 and S4).

## 481 4.2 3D mantle structure corrections

482 We run SPECSEM3D\_globe to obtain the synthetic seismograms corresponding to  
 483 the data with good quality measurements. Here, good quality means that more than ten  
 484 traces are used in a bin and  $CC_{3,2} > 0.90$  (Table S1). Most source parameters used in  
 485 the SEM simulations are from GCMT, but the source durations are replaced with our  
 486 estimated values. Then we apply our iterative method to these synthetic seismograms  
 487 to obtain the time delays, amplitude ratios and corresponding CC values. For most bins,  
 488 we successfully retrieve signals of S3KS, S4KS and S5KS and get high CC coefficients  
 489 (Table S2). For example, Fig. S5 shows the measurements using synthetic seismograms  
 490 corresponding to the bin 4 from the event #141101. The CC coefficients are 0.95 for S3KS,  
 491 0.96 for S4KS and 0.94 for S5KS, indicating good measurement quality. The S3KS time  
 492 delay,  $3dM_{3,2}^{SEM}$ , is as large as 0.60 s and the S4KS time delay,  $dt_{4,2}^{SEM}$ , is even larger,  
 493 1.13 s. The time delays measured on the data are 1.30 s for S3KS and 2.48 s for S4KS  
 494 (Table S1). Thus, 3D mantle structure corrections are large, up to nearly half the size  
 495 of the observations, and can not be ignored for this bin. The S3KS measurements on syn-  
 496 thetic seismograms of other bins are listed in Table S2 and almost all the bins have  $CC_{3,2}^{SEM}$   
 497 higher than 0.95, except the bin 1 from #010428 and bin 1 from #140721. The S3KS/SKKS  
 498 amplitude ratios range from 0.30 to 0.51 and the corrections to S3KS time delays are be-  
 499 tween -0.95 s and 0.04 s. Most of the corrections have negative values, indicating that  
 500 S3KS are delayed more than SKKS by the 3D mantle structure. The results for S4KS  
 501 and S5KS are listed in Tables S3 and S4.

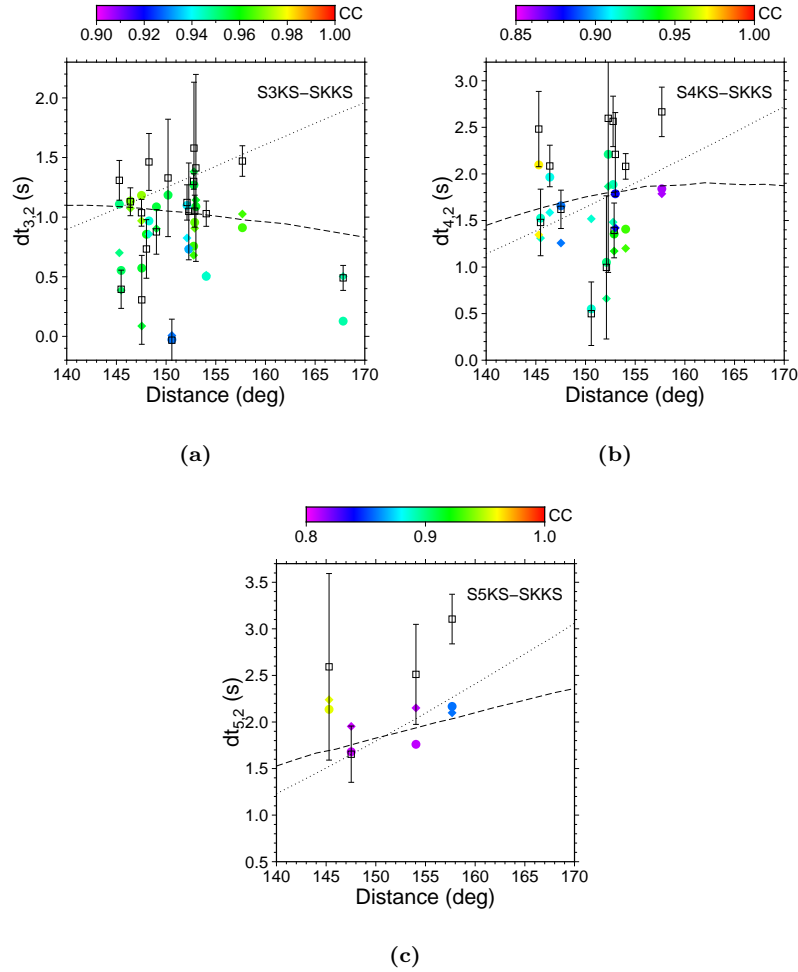
502 We also use ray theory to compute 3D mantle structure corrections for data at in-  
 503 dividual stations in each bin and take the average value to represent the correction for  
 504 that bin. These corrections are close to that measured on SEM synthetic seismograms  
 505 (Tables S2, S3, S4 and Fig. S6). However, large discrepancies are present for some bins.

506 For example, the correction to S3KS time delay based on ray theory is 0.27 s for the bin  
 507 5 from #141101, but it is -0.22 s using SEM synthetic seismograms.

508 Fig. 10 shows the SmKS (m=3, 4 and 5) time delays measured on the data with  
 509 high CC coefficients and the results after 3D mantle structure corrections. Here, we re-  
 510 quire  $CC_{3,2} \geq 0.90$  for a good quality of S3KS measurement. For S4KS, we only take  
 511 the bins with  $CC_{4,2} \geq 0.85$  and  $CC_{3,2} \geq 0.90$ , because a good quality of S4KS measure-  
 512 ment relies on a well-retrieved S3KS. Similarly, we require  $CC_{3,2} \geq 0.90$ ,  $CC_{4,2} \geq 0.85$   
 513 and  $CC_{5,2} \geq 0.80$  for good quality of S5KS measurements.

514 Most of the bins with good qualities of  $dt_{3,2}$  measurements have uncertainties smaller  
 515 than 0.4 s (Table S2). It is not surprising that  $dt_{4,2}$  and  $dt_{5,2}$  generally show larger un-  
 516 certainties than  $dt_{3,2}$ , due to their smaller *SNR* and/or incomplete separation of SmKS  
 517 (m=2, 3, 4 and 5) waveforms of our method. In spite of this, the uncertainties are still  
 518 much smaller than the anomalies (Tables S4 and S5), because the bins with large errors,  
 519 resulted from poor phase stripping and/or low SNRs, are discarded by the CC require-  
 520 ments. We note that bootstrapping results only help us infer variance in the dataset, but  
 521 not able to estimate systematic bias. The systematic bias could be due to strong man-  
 522 tle heterogeneities and source complexities etc, which can be assessed by investigating  
 523 global data from earthquakes at various places.





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525

526 **Figure 10:** SmKS time delays measured at 0.05-0.2 Hz. The empty squares represent the  
 527 SmKS ( $m=3$  in a, 4 in b and 5 in c) time delays measured on the data. Note that error bars  
 528 are symmetric and in a few cases extend beyond the limits of the figure. The solid circles are  
 529 SmKS time delays after the 3D mantle structure corrections based on ray theory and using  
 530 S40RTS model. The time delays are relative to SmKS-SKKS differential arrival times predicted  
 531 by PREM. The solid diamonds are SmKS time delays after the corrections measured on the  
 532 SEM synthetic seismograms made using S40RTS. The color shows the corresponding CC values  
 533 measured on the data. More detailed information is displayed in Tables S3-5. The black dashed  
 534 line in each figure is the corresponding SmKS time delay predicted by KHOMC (Kaneshima &  
 535 Helffrich, 2013). The black dotted lines show the EPOC predictions. The source depth used in  
 536 the KHOMC and EPOC predictions is 150 km. (a) S3KS time delays. (b) S4KS time delays. (c)  
 538 S5KS time delays.

### 4.3 Comparison between observations and predictions of two 1D models, EPOC and KHOMC

From Fig. 10, we can see that S3KS, S4KS and S5KS time delays predicted by EPOC (Irving et al., 2018) and KHOMC (Kaneshima & Helffrich, 2013) are close to each other at distance  $140^{\circ}$ – $155^{\circ}$ , where most of our data are located. The measured S3KS time delays are generally positive and consistent with the KHOMC and EPOC predictions supporting a slower  $V_p$  in the top outer core. The 3D mantle structure corrections, using either ray theory or SEM synthetic seismograms, are primarily negative and therefore reduce the measured S3KS time delays. After the corrections, travel time anomalies are less than the EPOC predictions and they seem to fit the KHOMC predictions better than that of EPOC. The difference between these two models can be better resolved using data at distances  $> 160^{\circ}$ , where their difference is larger than 0.5 s. Unfortunately, we have only one such datum, at a distance of  $167.8^{\circ}$ , so we can not clearly distinguish between EPOC and KHOMC. For S4KS and S5KS, the 3D mantle structure corrections are also primarily negative and they make measurements closer to the KHOMC and EPOC predictions. However, there are a few measurements dramatically departing from the EPOC and KHOMC predictions. For example,  $dt_{4,2}^{SEM}$  of bin 1 from #140721 is -0.22 s while the EPOC prediction is 1.22 s. For this bin, the two types of 3D mantle structure corrections have a large difference, 0.12 s from ray theory computation and -0.95 s from SEM synthetic seismograms. This large difference could be due to the limitation of ray theory, uncertainty in the S40RTS model, or poor performance of our method on the synthetic seismograms of this bin. The  $CC_{3,2}^{SEM}$  is only 0.91, much lower than that of other bins, which indicates a poor measurement quality. However,  $CC_{3,2}$  from data is a high value of 0.96 and its  $dt_{3,2}^{ray}$  is close to the EPOC and KHOMC predictions. This big difference is most likely due to a large uncertainty in the 3D mantle corrections using SEM synthetic seismograms. Some other bins, including S3KS time delays of bin 1 from #010428 and S5KS time delays of bin 1 from #010516, have similar issues. Note that the S5KS time delay of bin 1 from #010516 is beyond the y-axis range and not plotted in Fig. 10c.

We also apply two other filters, 0.05-0.7 Hz and 0.1-0.7 Hz, to the data and repeat the measurements. Because running SPEC-FEM3D\_globe to resolve a frequency of 0.7 Hz is very computationally expensive, we only compute the 3D mantle structure corrections using ray theory. Similar to the results at 0.05-0.2 Hz, the S3KS, S4KS and S5KS

572 measurements at 0.05-0.7 Hz are close to the EPOC and KHMOC predictions after the  
 573 3D mantle structure corrections (see Fig. S7).

574 Note that the bins shown in Fig. S7 are not the same as 0.05-0.7 Hz, because the  
 575 SNRs of data may change with frequency band and the measurement qualities could also  
 576 be different. Comparing to 0.05-0.2 Hz and 0.05-0.7 Hz, the number of bins with good  
 577 measurement qualities is lower at 0.1-0.7 Hz, indicating lower SNRs of data and/or re-  
 578 duced performance of our iterative method for this high frequency band for the data used  
 579 here. Relatively long period SmKS waves have been stacked to investigate outermost core  
 580 structure (e.g. 0.02-0.1 Hz in Tanaka, 2007). Shorter period waves have the potential  
 581 to resolve finer seismic structure. However, source rupture processes and propagation  
 582 effects due to lateral heterogeneities could give rise to more waveform complexities at  
 583 shorter period waves reducing the waveform coherencies of SmKS phases and therefore  
 584 affecting measurement qualities. This might explain the lower number of good quality  
 585 measurements (Figs. S7d-f) at 0.1-0.7 Hz than that at 0.05-0.2 Hz (Fig. 10) and 0.05-  
 586 0.7 Hz (Figs. S7a-c).

## 587 **5 Discussion**

588 SmKS differential arrival times are sensitive to the outer core structure, but accu-  
 589 rate measurements of differential arrival times are hampered by their mixed waveforms  
 590 as a whispering-gallery mode. To extract each individual SmKS phase, Eaton & Kendall  
 591 (2006) use SKKS as reference waveform and apply deconvolution to SmKS series to con-  
 592 vert their waveforms into simple pulses. However, the deconvolution method either re-  
 593 quires very high SNR and or has reduced resolution. Here, we develop an iterative method  
 594 to isolate individual SmKS waveforms with a high resolution. Our method keeps wave-  
 595 form features of each SmKS and therefore allows us to measure SmKS time delays by  
 596 CC.

597 We use two different methods, ray theory and SEM synthetic seismograms, to com-  
 598 pute effects of mantle heterogeneities and make these corrections to the measurements.  
 599 The corrections are between -0.5 s and 0.5 s for most bins, but some bins have 3D man-  
 600 tle perturbations even greater than 1.0 s. Furthermore, we see big differences between  
 601 the two types of corrections for some data (e.g. > 1 s for bin 1 from event #140721),  
 602 although they generally have positive correlation (Fig. S6). We also used another 3D man-

603 tle model S362ANI (Kustowski et al., 2008) to compute the SEM synthetic seismograms  
604 and measure the 3D mantle structure corrections. Compared to our results using S40RTS,  
605 the corrections to S3KS-SKKS, S4KS-SKKS and S5KS-SKKS differential arrival times  
606 using S362ANI are generally stronger. Consequently, the corrected SmKS-SKKS ( $m=3$ ,  
607 4 and 5) time delays become even smaller (Fig. S8,9). Helffrich & Kaneshima (2013) used  
608 earthquakes in Fiji and Argentina to investigate SmKS-SKKS time delays. Their mea-  
609 sured S3KS-SKKS time delays from Fiji are generally larger than that from Argentina.  
610 The earthquakes in our study are geographically close to Fiji and the 3D mantle correc-  
611 tions to S3KS-SKKS time delays tend to reduce the S3KS-SKKS time delays (Fig. 10a).  
612 Thus, the higher S3KS-SKKS time delays from events in Fiji by (Helffrich & Kaneshima,  
613 2013) can be largely explained with the 3D mantle structure. 3D mantle structure cor-  
614 rections should be routinely considered to reduce bias in the  $V_p$  estimation of uppermost  
615 outer core.

616 After correcting for 3D mantle structure, there are still significant SmKS-SKKS time  
617 delays at all of the three frequency bands (Figs. 10 and S7), indicating a lower  $V_p$  than  
618 PREM model in the shallow outer core. Strong locally concentrated heterogeneities, such  
619 as the previously detected Ultra Low Velocity Zones (ULVZs) at the source side of our  
620 study region (see, for example, the compilations by S. Yu & Garnero, 2018), are not ac-  
621 curately represented in the smooth global tomography model of S40RTS and could af-  
622 fect the measurements. However, these ULVZ effects have been investigated by Tanaka  
623 (2007) and they are expected to be smaller than our measured time delays. In addition,  
624 such strong heterogeneities would decrease the coherencies between SKKS and SmKS  
625 ( $m=3$ , 4 and 5) and only the results with high CC values are selected in our method. Fur-  
626 ther quantitative investigations will rely on better constraints on the properties and ge-  
627 ographical distributions of ULVZ and more detailed numerical waveform modeling. Thus,  
628 we do not believe that the SmKS-SKKS travel time delays are solely due to ULVZs, but  
629 do indeed indicate a seismically slow uppermost outer core.

630 Although scatter and uncertainty are present, our measurements are generally con-  
631 sistent with the predictions by the KHOMC and EPOC models. Assuming the outer core  
632 is homogeneous, Irving et al. (2018) use a physically consistent equation-of-state (EoS)  
633 to parameterize the elastic properties of outer and carry out inversions for seismic nor-  
634 mal mode data. This normal mode derived EPOC model shows lower  $V_p$  and higher den-  
635 sity than PREM at the top of outer core. Although EPOC does not use body-wave data,

636 its fit to SmKS data is better than PREM (see Fig. 3 in Irving et al. (2018) and Figs.  
 637 10 and S7 in this study). KHOMC is derived from SmKS body-wave travel time anoma-  
 638 lies and has higher depth resolution than EPOC. We note that both EPOC and KHOMC  
 639 models have a low  $V_p$  at the top of outer core, but they have different depth gradients  
 640 of  $V_p$ . KHOMC seems to fit our results better than EPOC. For example, EPOC over-  
 641 predicts most S3KS-SKKS time delays after 3D mantle corrections. However, given the  
 642 scatter present in our measurements, either EPOC or KHOMC fits the data well. The  
 643 contrast between stratified or homogeneous structure has important implications for un-  
 644 derstanding the thermochemical status of core and the associated geodynamo. A strat-  
 645 ified outer core would change the flow in the outer core and therefore affect the secular  
 646 variation of geomagnetic field (e.g. Braginsky, 1993; Buffett, 2014; Buffett et al., 2016).  
 647 However, the detailed effects of such stratification on the geodynamo and the compat-  
 648 ibility between seismic and geomagnetic observations (e.g. the thickness of stratified layer)  
 649 are still inconclusive (Gubbins, 2007; Buffett, 2014; Chulliat & Maus, 2014; Lesur et al.,  
 650 2015). Additionally, the mechanism for the formation of stratification is also under de-  
 651 bate. For example, high concentrations of light elements, including S, O, Si, C and H,  
 652 at the top of outer core could cause a stratification (e.g. Fearn & Loper, 1981; Buffett  
 653 & Seagle, 2010; Gubbins & Davies, 2013; Nakagawa, 2018; Helffrich & Kaneshima, 2013),  
 654 but how these light elements change  $V_p$  is still under debate (Helffrich, 2012; Brodholt  
 655 & Badro, 2017). In this study, we cannot easily distinguish between the EPOC and KHOMC  
 656 models, but these two models do give different predictions of SmKS-SKKS differential  
 657 arrival times. Thus, both gathering more observations, e.g. S3KS-SKKS differential times  
 658 at a distance  $> 160^\circ$ , and considering other geophysical probes of the outer core, for ex-  
 659 ample normal mode observations, will be critical to better resolve the uppermost outer  
 660 core's density and  $V_p$ , providing vital data to constrain the thermochemical status of  
 661 the outer core.

## 622 **6 Conclusions**

663 We introduce an array-based iterative method to measure SmKS-SKKS ( $m=3, 4$   
 664 and 5) differential arrival times and use them to investigate the  $V_p$  in Earth's uppermost  
 665 outer core. We validate this method by testing synthetic seismograms and apply this method  
 666 to data at stations in Europe from eleven earthquakes in Fiji-Tonga, Vanuatu, New Britain  
 667 and Solomon Islands. Using the SKKS signal as a reference, S3KS, S4KS and S5KS wave-

668 forms are successfully extracted and S3KS-SKKS, S4KS-SKKS and S5KS-SKKS differ-  
 669 ential arrival times are measured by waveform cross-correlation. This iterative method  
 670 not only gives us the measurements of differential arrival times, but also allow us to as-  
 671 sess measurement qualities based the CC coefficients and amplitude information. SmKS-  
 672 SKKS differential arrival times are sensitive to  $V_p$  at the top of the outer core, but 3D  
 673 mantle structures could also affect the arrival times. We use the 3D mantle model S40RTS  
 674 and two different methods, ray theory and SEM synthetic seismograms, to estimate these  
 675 anomalies for the frequency of 0.05-0.2 Hz. The results show that the arrival time anoma-  
 676 lies due to 3D mantle structure effects are large (e.g.  $> 0.5$  s) for some data and some-  
 677 times there are big differences between the corrections calculated using ray theory and  
 678 SEM synthetics. After corrections for 3D mantle structure, we still see large positive S3KS-  
 679 SKKS, S4KS-SKKS and S5KS-SKKS differential arrival times, indicating a lower  $V_p$  than  
 680 in PREM at the top of outer core. Our measurements are consistent with the predic-  
 681 tions of KHOMC and EPOC models. EPOC has a homogeneous outer core while KHOMC  
 682 contains a stratified layer at the top of outer core. Based on the data in this study, we  
 683 cannot clearly distinguish the KHOMC and EPOC models, so more data, e.g. S3KS-SKKS  
 684 differential time at a distances  $> 160^\circ$ , will be necessary to help us distinguish between  
 685 them.

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 692 the References) have been collected using the Python toolbox ObsPy (Beyreuther et al.,  
 693 2010).

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832 *The waveform data used in this study is from the following networks:*

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 834 *BN, BS (doi:10.7914/SN/BS), BW (doi:10.7914/SN/BW), C, C1*  
 835 *(doi:10.7914/SN/C1), CA (doi:10.7914/SN/CA), CB (doi:10.7914/SN/CB),*  
 836 *CH (doi:10.12686/sed/networks/ch), CM, CN (doi:10.7914/SN/CN),*  
 837 *CR, CU (doi:10.7914/SN/CU), CX (doi:10.14470/PK615318), CZ*  
 838 *(doi:10.7914/SN/CZ), DK, DR (doi:10.7914/SN/DR), DZ, EB,*  
 839 *EE, EI (doi:10.7914/SN/EI), FN, FR (doi:10.15778/RESIF.FR), G*  
 840 *(doi:10.18715/GEOSCOPE.G), GB, GE (doi:10.14470/TR560404), GR, GS*  
 841 *(doi:10.7914/SN/GS), GT (doi:10.7914/SN/GT), GU (doi:10.7914/SN/GU),*  
 842 *HE (doi:10.14470/UR044600), HL (doi:10.7914/SN/HL), HT*  
 843 *(doi:10.7914/SN/HT), HU (doi:10.14470/UH028726), IB (doi:10.7914/SN/IB),*  
 844 *II(doi:10.7914/SN/II), IM, IP, IS, IU(doi:10.7914/SN/IU), IV*  
 845 *(doi:10.13127/SD/X0FXnH7QfY), KC (doi:10.7914/SN/KC), KN, KO*  
 846 *(doi:10.7914/SN/KO), KP (doi:10.7914/SN/KP), KR (doi:10.7914/SN/KR),*  
 847 *KW, KZ (doi:10.7914/SN/KZ), LD, LI (doi:10.7914/SN/LI), LX,*  
 848 *MC, MD (doi:10.7914/SN/MD), MN (doi:10.13127/SD/fBBBtDtd6q),*  
 849 *MX (doi:10.21766/SSNMX/SN/MX), N4 (doi:10.7914/SN/N4), NA*  
 850 *(doi:10.21944/dffa7a3f-7e3a-3b33-a436-516a01b6af3f), NE (doi:10.7914/SN/NE),*  
 851 *NI (doi:10.7914/SN/NI), NJ (doi:10.7914/SN/NJ), NL (doi:10.21944/e970fd34-*  
 852 *23b9-3411-b366-e4f72877d2c5), NM, NO, NR (doi:10.7914/SN/NR), NU*  
 853 *(doi:10.7914/SN/NU), OE (doi:10.7914/SN/OE), OV, OX (doi:10.7914/SN/OX),*  
 854 *PE (doi:10.7914/SN/PE), PL, PM, PR (doi:10.7914/SN/PR), PZ*  
 855 *(doi:10.7914/SN/PZ), RD (doi:10.15778/RESIF.RD), RO (doi:10.7914/SN/RO),*  
 856 *SI, SJ, SK (doi:10.14470/FX099882), SL (doi:10.7914/SN/SL),*  
 857 *SP (doi:10.7914/SN/SP), SS, SV, SX (doi:10.7914/SN/SX), TA*  
 858 *(doi:10.7914/SN/TA), TH (doi:10.7914/SN/TH), TR, TT, TU, UK,*  
 859 *UP (doi:10.18159/SNSN), US (doi:10.7914/SN/US), VE, VI, WC, WI*  
 860 *(doi: doi:10.18715/antilles.WI), WM (doi:10.14470/JZ581150), X5,*  
 861 *X6 (doi:10.7914/SN/X6\_2007), X7 (doi:10.15778/RESIF.X72010),*  
 862 *XB (doi:10.7914/SN/XB\_2009), XE (doi:10.7914/SN/XE\_2009), XI*  
 863 *(doi:10.7914/SN/XI\_2011), XJ (doi:10.12686/sed/networks/xh), XK*

864 (doi:10.7914/SN/XK\_2012), XN (doi:10.7914/SN/XN\_2008), XO  
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 871 YG, YH (doi:10.7914/SN/YH\_2012), YI (doi:10.7914/SN/YI\_2003 and  
 872 doi:10.15778/RESIF.YI2008), YJ, YK, YO (doi:10.7914/SN/YO\_2014), YP,  
 873 YQ (doi:10.7914/SN/YQ\_2013), YR (doi:10.15778/RESIF.YR1999), YS  
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 875 Z4 (doi:10.7914/SN/Z4\_2009), Z9 (doi:10.7914/SN/Z9\_2010), ZA,  
 876 ZC (doi:10.7914/SN/ZC\_2013), ZD (doi:10.7914/SN/ZD\_2010), ZE  
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 880 ZU, ZV, ZX, ZZ (doi:10.14470/MM7557265463), 1E, 4F, 6D, 6E, 7A  
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