

Seismic source tracking with six degree-of-freedom ground motion observations

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Key Points:

- We demonstrate concepts for seismic source tracking from combined point measurements of translational and rotational motions.
- We show that S-waves can be used for seismic source tracking due to the wavefield separation in rotational components.
- We discuss the effects of directivity, rupture speed, source-receiver geometry, and present a real data example of tracking traffic sources.

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Abstract

Back azimuth information can be determined from combined measurements of rotations and translations at a single site. Such six degree-of-freedom (6-DoF) measurements are reasonably stable in delivering similar information compared to a small-scale array of three-component seismometers. Here we investigate whether a 6-DoF approach is applicable to tracking seismic sources. While common approaches determining the timing and location of energy sources generating seismic waves rely on the information of P-waves, here we use S-waves. We track back azimuths of directly arriving SH-waves in the 2-D case, P-converted SV-waves, direct SV- and direct SH-waves in the 3-D case. For data analysis, we compare a cross-correlation approach using a grid-search optimization algorithm with a polarization analysis method. We successfully recover the rupture path and rupture velocity with only one station, under the assumption of an approximately known fault location. Using more than one station, rupture imaging in space and time is possible without a priori assumptions. We discuss the effects of rupture directivity, supershear rupture velocity, source-receiver geometry, wavefield interference, and noise. We verify our approach with the analysis of moving traffic noise sources using 6-DoF observations. The collocated classic seismometer and newly-built ring laser gyroscope ROMY near Munich, Germany, allow us to record high-fidelity, broadband 6-DoF (particle velocity and rotational rate) ground motions. We successfully track vehicles and estimate their speed while traveling along a nearby highway using the estimated BAZ as a function of time of a single station observation.

1 Introduction

The path and speed of large earthquakes are crucial factors determining their damage potential. Rupture kinematics can be routinely determined by finite-fault inversion approaches based on close fitting of observations using a large number of free parameters. However, despite recent advances (e.g., Shimizu et al., 2020), kinematic models typically need to pre-define fault geometry, are characterized by inherent non-uniqueness (Mai et al., 2016) and do not ensure mechanical consistency in terms of earthquake dynamics (e.g., Tinti et al., 2005; Causse et al., 2014). Rising computational resources allow the development of observational constrained dynamic rupture models complementing data-driven analyses. Such rupture scenarios provide physically self-consistent descriptions of how complex (yet again, prescribed) faults slip, while their complexity limits the total number of feasible numerical experiments (Uphoff et al., 2017; Wollherr et al., 2019; Ulrich, Gabriel, et al., 2019; Ulrich, Vater, et al., 2019, e.g.,). The rise of large-scale, dense seismic array instrumentation has enabled complementary techniques tracking earthquakes in space and time (e.g., Meng et al., 2012; Kiser & Ishii, 2017; Bao et al., 2019). Such methods image

55 coherent high-frequent energy radiation (not slip) in simple and rapid manners requiring
56 very limited a priori knowledge.

57 Here, we present proofs of concept for seismic source tracking with single-point mea-
58 surements combining rotational and translational components of the seismic wavefield. We
59 study the polarity of directly incoming SH- (in 2-D and 3-D), direct SV-waves and P-
60 converted SV-waves (in 3-D) of synthetic 6-DoF time series. First, we introduce two distinct
61 single-station approaches for estimating the direction of incoming waves, namely a cross-
62 correlation approach using a grid-search optimization algorithm and a polarization analysis
63 method using point measurements. Second, a statistical approach for combining the back
64 azimuth estimates of several stations is presented which shows a high robustness with re-
65 spect to measurement uncertainties. We verify the concept in synthetic 2-D experiments
66 analyzing SH-wave polarity and discuss the applicability and robustness of the developed
67 methodologies. Third, we demonstrate earthquake rupture tracking in 3-D media from the
68 rotation polarization caused by P-converted SV-waves and direct SV- and SH-waves. We
69 analyze the effect of interfering arrivals and non-uniform slip rate distribution. We discuss
70 source-receiver scales and geometry as well as challenges of the method for future global ap-
71 plications. Finally, we show a real data application of the proposed approach, which mimics
72 the rupture imaging process. We track moving traffic seismic sources using a high-resolution
73 6-DoF point measurement.

74 1.1 Seismic source tracking

75 Most common techniques to image earthquake properties using array data can be di-
76 vided into two categories which are both based on analyzing the phase information of P-
77 waves. In contrast to finite slip inversions, no detailed knowledge of Green's functions and
78 source properties is necessary. Methods of the first category are based on conventional ar-
79 ray measurements. Termed back-projection methods, seismic energy radiation is imaged
80 by applying array beam-forming techniques. Back-projection was for the first time suc-
81 cessfully demonstrated for the 2004 Sumatra-Andaman Earthquake (Krüger & Ohrnberger,
82 2005; Ishii et al., 2005). Directivity effects were utilized to characterize faulting mechanisms
83 (Ammon et al., 2005).

84 Methods of the second category track earthquake rupture by estimating the back az-
85 imuth (BAz) of incoming waves with a single-station. In polarization analysis, the three
86 translational components of standard seismometers can be used to estimate the BAz and
87 incidence angle of incoming waves (Flinn, 1965; Montalbetti & Kanasewich, 1970; Vidale,
88 1986; Greenhalgh et al., 2005). Bayer et al. (2012) developed a single-station approach to
89 track moving sources by polarization analysis of local and regional P-wave arrivals. They

90 normalize the BAz variation with respect to a known hypocenter. Frohlich & Pulliam (1999)
 91 pointed out that, compared to travelttime-based methods, classic single-station approaches
 92 suffer from several ambiguities, for example 180° BAz errors. The joint analysis of transla-
 93 tional and rotational motions can help overcome such drawbacks.

94 1.2 6-DoF ground motions

95 The complete wavefield excited by an infinitesimally small deformation can be described
 96 by the three components of translation, three components of rotation, and six components
 97 of strain (Aki & Richards, 2002). However, until recently, seismology is dominated by trans-
 98 lational observations (vertical, N-S, E-W), sometimes combined with strain measurements.
 99 Translational motion is the movement of a particle along an axis. In contrast, rotational
 100 motion describes the particle movement around an axis. Information on rotations has been
 101 widely ignored, mainly, because of measurement difficulties. 6-DoF information is obtained
 102 from measuring in addition to three translational components also three components of ro-
 103 tational motion. This increase in information compared to classical observations has the
 104 potential to improve existing methods and creates new opportunities for research and in-
 105 dustry (e.g., Igel et al., 2015; Schmelzbach et al., 2018). Rotational motions can be derived
 106 from arrays of conventional single or multi-component sensors (e.g., Spudich et al., 1995;
 107 Huang, 2003; Suryanto et al., 2006; Spudich & Fletcher, 2009; Taylor et al., 2020). However,
 108 these methods are limited by array spacing as well as local heterogeneities and site effects.
 109 Classic translational measurements are also sensitive to tilt, i.e., the horizontal components
 110 of the rotation vector (van Driel et al., 2015; Graizer & Kalkan, 2008). Recent advances in
 111 fibre-optic gyroscopes and ring laser-based sensors show that applicable, single-station mea-
 112 surements for translation and rotation are within reach (Schreiber & Wells, 2013; Bernauer
 113 et al., 2012, 2018).

114 The earthquake source process and the interaction of the wavefield with a free surface
 115 or heterogeneities of the Earth can excite rotational ground motions. In isotropic media
 116 the rotational motion $\boldsymbol{\omega} = (\omega_x, \omega_y, \omega_z)^T$ can be described by a linear combination of spatial
 117 derivatives of the translational particle displacement motion $\mathbf{u} = (u_x, u_y, u_z)^T$ (e.g., Cochard
 118 et al., 2006):

$$119 \begin{pmatrix} \omega_x \\ \omega_y \\ \omega_z \end{pmatrix} = \frac{1}{2} \nabla \times \mathbf{u} = \frac{1}{2} \begin{pmatrix} \partial_y u_z - \partial_z u_y \\ \partial_z u_x - \partial_x u_z \\ \partial_x u_y - \partial_y u_x \end{pmatrix}, \quad (1)$$

120 where \times denotes cross product and ∂_k denotes spatial derivatives with respect to x_k . The
 121 same relation is valid for the time derivatives, the rotation rate $\dot{\boldsymbol{\omega}}$ and particle velocity \mathbf{v} . At
 122 the free surface, the stress-free boundary condition is slightly modified (e.g., Schmelzbach
 123 et al., 2018). While inside isotropic media the curl operator separates the S-wave field,

124 in anisotropic media even (quasi-)P-waves can have a rotational component (Pham et al.,
 125 2010). Local phase velocities and the BAz can be estimated from 6-DoF measurements at
 126 a single-station due to the relation between translations and rotations (Pancha et al., 2000;
 127 Igel et al., 2007; Hadziioannou et al., 2012; Edme & Yuan, 2016; Sollberger et al., 2018).

128 2 Methodology

129 2.1 BAz estimation from a single station

130 We first test two different methods to track sources of seismic energy (earthquake
 131 rupture) in simple 2D examples. Both methods are based on a plane wave assumption and
 132 analyze the polarity of directly arriving SH-waves at a single station. Since the region of
 133 energy radiation moves during the rupture across the fault plane, we utilize sliding windows
 134 moving throughout the signal to determine the evolution of the signal source direction. In
 135 each time window, the BAz is estimated and a temporal trend can be derived by comparing
 136 all windows. When this information is combined with a priori knowledge on the fault or
 137 with data from other stations, the rupture propagation and its velocity can be estimated.

138 The CC (cross-correlation) method is a grid-search optimization algorithm that relies
 139 on the interdependence of transverse translational motion and vertical rotation. The CC is
 140 a measure for the similarity between two signals and the CC coefficient provides a measure
 141 for the degree of similarity (see Fig. 1). A CC coefficient of 1 implies perfect similarity, a
 142 value of -1 means anti-correlation.

143 Similar to the approach by Igel et al. (2007), we estimate the BAz by rotating the
 144 horizontal acceleration components in small steps around all possible BAz ($0^\circ - 360^\circ$) and
 145 cross-correlating successively with the vertical rotation rate. A zero-lag normalized CC
 146 coefficient is used. For a noise-free signal, the CC coefficient in the grid-search is a function
 147 without a clear maximum. It is a step-function that jumps from -1 to 1. Therefore, we use
 148 the two zero transitions of the step-function instead of the global maximum. We expect
 149 that the central position between the zero transitions corresponds to the actual BAz.

150 The second method was introduced by Sollberger et al. (2018) and we refer to it here-
 151 inafter as polarization analysis. In comparison to the CC method, it is more flexible and
 152 can be applied to P-, SV-, SH-, Rayleigh- and Love-waves. Instead of a MUSIC likelihood
 153 function (Schmidt, 1986), we use one based on the classical power spectrum, because we
 154 expect a more stable result. We assume that the global maximum of the likelihood function
 155 is related to the actual BAz. It is necessary to define a parameter space for evaluating
 156 the likelihood function in a grid search. While the CC method requires the definition of
 157 the BAz increments, the increments for the S-wave velocity and the incident angle must be

158 additionally defined for the polarization analysis. Both methods are illustrated for a plane
 159 wave in Fig. 1.

160 The difficulty of retrieving BAz (source directivity) from 3-D observations with solely
 161 translational motions is due to two challenges:

162 1) the 180° ambiguity in BAz estimates if only translational motions are recorded
 163 (Langston & Liang, 2008). Considering that rotational motions are essentially the curl of
 164 translational motions, the polarity of rotation will reverse in case of an opposite propagating
 165 direction while translation polarity remains unchanged. Thus the joint analysis of rotation
 166 and translation will help to remove the 180° ambiguity when locating the sources.

167 2) translation records suffer from interfering different types of wavefields at the free
 168 surface, i.e., P- and SV/SW-waves, Rayleigh- and Love-waves are generally intermixed in
 169 recorded horizontal components. However, rotational motions naturally separate P- and
 170 S-waves as P-waves do not generate rotational motions in isotropic media. SV- and SW-
 171 waves (the same as Rayleigh- and Love-waves) are also naturally separated despite unknown
 172 source locations since SV-waves or Rayleigh-waves only generate rotational motions on hor-
 173 izontal components while SW related (Love-wave related) rotational motions being isolated
 174 on vertical components. We can therefore take advantage of the fact that two horizontal
 175 rotational components contain exclusively SV- or Rayleigh-waves.

176 Based on the plane wave assumption, the ratio between the two horizontal rotational
 177 components is directly related to the BAz according to:

$$178 \theta_{BAz} = -\arctan\left(\frac{\dot{\omega}_n}{\dot{\omega}_e}\right), \quad (2)$$

179 where $\dot{\omega}_n$ and $\dot{\omega}_e$ denote the north-south and east-west components of rotational rate, re-
 180 spectively. This simple relationship is specially useful for estimating source directivity and
 181 it is independent of any possible radiation pattern that the source might have. The negative
 182 θ_{BAz} value derived from the inverse tangent function is converted to the value within 0°
 183 and 180° by adding 180° . To remove the 180° ambiguity, we then compare the rotated
 184 transverse component of rotational rate ($\dot{\omega}_t$) based on θ_{BAz} to the vertical component of
 185 acceleration (a_z). If they are positively correlated, 180° should be added to θ_{BAz} .

186 2.2 Combining multiple stations

187 It is possible to track the horizontal propagation of a rupture with only one station,
 188 presupposed the BAz changes correctly determined in a seismogram and the fault position
 189 is known a priori. An infinitesimal thin ray could be constructed for each estimated BAz in
 190 the direction of the directly arriving waves. The intersections of these rays with the fault

191 position would show the temporal evolution of the rupture. In case of an unknown fault,
 192 at least two stations are necessary for the tracking process. But for more than two stations
 193 the rays will not intersect in exactly one point, since measurement errors and inaccuracies
 194 in the methodology can not be excluded completely. We want to take these uncertainties
 195 in the BAZ into account by using wider beams instead of infinitesimal thin rays. We define
 196 the shape function $p(x, y, t)$ of each beam as

$$197 \quad p(x, y, t) = \sum_{i=1}^{N_{\text{stations}}} \frac{1}{\sqrt{2\pi\sigma_i(t)^2}} \exp\left(-\frac{1}{2} \left(\frac{\Phi_i(x, y, t)}{\sigma_i(t)}\right)^2\right), \quad (3)$$

198 where the standard deviation $\sigma_i(t)$ is defined individually for each station i , $\Phi_i(x, y, t) \in$
 199 $[0, 180^\circ]$ denotes the angular distance from an arbitrary point in space to the estimated BAZ
 200 of a specific station i and N_{stations} denotes the number of stations. Note that the shape
 201 function is time-dependent and is defined for each horizontal position (x, y) . The time
 202 framework is defined in such a way that $t = 0$ corresponds to the first arrival at a respective
 203 station. For a specific time-step t_0 the amplitudes of all beams are added up and we assume
 204 that the most likely source position is close to the maximum value of $p(x, y, t_0)$. In Fig. 2
 205 we show examples of shape functions for different BAZ errors and a 2-D representation of
 206 $p(x, y, t_0)$ for two stations using Eq. 3 with $N_{\text{stations}} = 2$.

207 **3 Synthetic case studies in 2-D**

208 In the following, we use data of elastic wave simulations in 2-D to demonstrate possible
 209 applications on a fundamental level and their limitations. 2-D wave propagation simulations
 210 are performed using the spectral element package `se2wave`. The mesh representation and
 211 support for MPI parallelism in `se2wave` is provided via PETSc (Balay et al., 2019, 1997).
 212 We note that while homogeneous 2D media SH rotation rate ought to be proportional to
 213 ground acceleration, the subsequently presented synthetic examples gradually increase in
 214 complexity and benefit from consistent and reproducible numerical calculations.

215 We describe two different test cases. The fault position is known in the first case and we
 216 try to track the spatial and temporal evolution with only a single station. In the second case,
 217 we assume that the fault position is unknown and the individual results of many stations
 218 are combined.

219 **3.1 Rupture tracking with a single 6-DoF station**

220 We model a pure strike-slip earthquake embedded in a 2-D homogeneous medium. The
 221 unilateral rupture has a constant speed of 80% of the shear velocity v_s and it is implemented
 222 as a line of double-couple point sources. We choose the source time function of each point

223 source to be an ordinary Gaussian. Additionally, we slightly randomize the onset time and
 224 the seismic moment to render our synthetic study more realistic.

225 Fig. 3 visualizes the model setup and tracking results of a unilateral rupture that
 226 propagates from north to south. We illustrate the receiver and fault setup in the upper
 227 panel, in which the stations are represented by two blue triangles. The a priori known fault
 228 position is marked by a grey dashed line and the unknown rupture trace by a red line.

229 Each station records the horizontal accelerations a_x, a_y and the vertical rotation rate
 230 $\dot{\omega}_z$ of the directly incoming P- and S-waves (middle panels of Fig. 3). These seismograms
 231 demonstrate that P-waves do not have a rotational component in an isotropic and homoge-
 232 neous medium. Station A only records weak P-wave amplitudes due to the perpendicular
 233 position with respect to the rupture. Due to different BAZ between source and station, the
 234 duration of the SH-arrivals is different for stations A and B. The maximum expected BAZ
 235 variation for station A is about 11.5° and 5° for station B. We estimate the BAZ changes
 236 by moving a sliding window of 1.5 s length through the SH-wave signal of the seismograms.

237 In each window the BAZ is estimated by the polarization analysis and the CC method
 238 and the results are illustrated in the bottom subplots of Fig. 3. Each point in these graphs
 239 represents the central position of a time window. The points are color-coded relative to the
 240 first and last SH-wave arrival. Both methods provide the same linear trend and the results
 241 are nearly perfectly overlapping. We include a graphical representation of the resulting BAZ
 242 estimates in the upper panel of the same figure in the form of color-coded thin rays for each
 243 estimated BAZ, respectively. The rays show a clear trend for both stations from north to
 244 south. The horizontal dimensions of the rupture are tracked correctly.

245 If the rupture or a certain part of the rupture propagates at a constant rupture speed and
 246 the starting and ending point are approximately known, it is possible to estimate the rupture
 247 speed by trigonometric considerations. The rupture speed depends on the S-wave velocity,
 248 the rupture length, the rupture duration measured at the receiver, and the orientation
 249 of rupture direction and receiver. Both bottom plots in Fig. 3 are divided in three sub-
 250 windows. We determine the velocity in each sub-window by fitting a straight line through
 251 the estimated BAZ and express it relative to the known S-wave velocity. In this manner, we
 252 estimate rupture velocities close to the real value of $80\% v_s$ in all sub-windows.

253 **3.2 Direct estimation of rupture velocity**

254 We evaluate the sensitivity of the proposed 6-DoF tracking methods to variations in
 255 earthquake rupture propagation speed across the fault. Reliable, far-field estimation of rup-
 256 ture velocity is important to constrain earthquake dynamics, stress drop, and implications

for seismic hazard but is inherently difficult because of the intermixing of rupture geometry and rise time in controlling the P- and S-wave pulse shapes (e.g., McGuire & Kaneko, 2018).

We test three different rupture scenarios. The model setup is the same as in the previous 2-D tests, but here only the first half of the rupture has a constant speed of 80% v_s . The second half breaks with a constant velocity of 40%, 60% or 150% of v_s . Earthquake ruptures can propagate at sub-Rayleigh or at intersonic speeds (e.g., Archuleta, 1984; Dunham et al., 2003; Gabriel et al., 2012) and a speed of 150% v_s means that the rupture is propagating faster than the radiated SH-waves and close to the P-wave speed. This effect is referred to as supershear rupture speeds.

In Fig. 4 we estimate the BAz changes for each case of velocity variation at station A in the same way as in Fig. 3. Each column represents the result for a specific velocity jump. The BAz results are divided into three sub-windows, in which the rupture velocity is calculated, respectively. The final velocity results, which we express relative to the shear-wave velocity, are illustrated for each sub-window by a red horizontal line in the lower panels. We indicate the true rupture speed by blue dashed lines. In each test scenario there is a significant increase or decrease visible from the starting velocity of 80% v_s in the first sub-window to the final rupture velocity of 40%, 60% and 150% in the last sub-windows. While the speed of the second half of the rupture is determined nearly perfectly, the starting velocity is slightly underestimated.

3.3 Rupture tracking in heterogeneous media

We expect that 6-DoF rupture tracking is more difficult in heterogeneous materials since reflected and scattered energy will contaminate the directly arriving SH-waves. We rerun the simulation of Fig. 3 now perturbing the homogeneous model by adding a normally distributed random material heterogeneity. We add a variation of up to $\pm 5\%$ to density, P-wave and S-wave velocity in the medium. In the numerical simulations quadrilateral elements are employed, each possessing piece-wise constant material properties and edge lengths ~ 100 m. Material properties in each element are perturbed independently of each other and no smoothing of the piece-wise constant properties is applied between neighboring elements.

The seismograms recorded at station A are shown in Fig. 5. Due to reflections in the material, P- and S-waves are no longer perfectly separated. Reflected phases are visible in all components after the dominant SH-arrival. The lower panel shows the tracking result, in which the true starting and ending BAz for station A are marked by blue dashed lines. The spatial dimensions of the rupture are nearly perfectly estimated. Although we expect

291 a nearly straight line for the temporal evolution, there are higher deflections than in the
 292 homogeneous model. However, on average a rupture speed of 77% v_s is determined, which is
 293 very close to the true velocity. The BAz deflections are increasing for stations in the higher
 294 distance and for stations that are placed in a geometrical orientation, in which the SH-wave
 295 amplitudes are less dominant compared to the P-wave amplitudes.

296 **3.4 Rupture tracking for unknown simple and complex rupture paths and** 297 **directivity effects**

298 Rupture tracking with only a single station is possible if the fault or more explicitly
 299 the rupture path is known a priori. In the following, we assume that the fault position is
 300 unknown. Since a single station is not enough to track the rupture in this case, we here
 301 combine the BAz estimates of many stations. The BAz is still calculated in a single-station
 302 approach at each receiver, but the final tracking results of all stations are combined.

303 As described in section 2.2, for a certain time-step we send a virtual beam back from
 304 each station in the direction of the rupture. By using a broad beam instead of a thin ray,
 305 we here take BAz uncertainties into account. We add up the amplitudes of all beams and
 306 the maximum is expected to be the most likely source point. The station coordinates and
 307 the BAz changes at each receiver are the input parameters for Eq. 3.

308 However, there is another issue that is referred to as a consistent time-frame. Due to
 309 different BAz between stations and rupture, the SH-arrivals at each station have a varying
 310 length of time (compare to the seismograms of Fig. 3). The time-shift of the sliding window
 311 has to be corrected at each station for this effect. Otherwise, an offset of the estimated
 312 rupture position from the actual location is expected, even if the BAz is correctly determined.
 313 Previous studies neglect such directivity effects expecting only small deviations.

314 First, we verify that a 6-DoF method provides accurate results for simple and more
 315 complex fault geometry then we discuss the importance of directivity effects. We apply a
 316 time correction to the BAz estimates by assuming that the start and the end of the directly
 317 arriving SH-waves is visible in the signal. This is done in the rotational component of ground
 318 motions due to its high sensitivity to shear motions.

319 In the following, we track a simple unilateral rupture at five stations. The stations are
 320 placed in an asymmetrical pattern around the rupture with different distances to the source.
 321 The stations are situated in such a way that the resolution is about the same for both spatial
 322 dimensions. The medium and rupture parameters are equal to the homogeneous model in
 323 the previous section. The SH-waves at each station are picked manually and the BAz
 324 change is independently determined of the other stations. In Fig. 6 rupture tracking results

325 are shown for three different time-steps (a, b and c). The estimated starting position is
 326 shown in Fig. 6a and the ending position in Fig. 6c. Animation S1 visualizes the continuous
 327 rupture imaging (Movie S1: ms01.avi).

328 In the first subplot of Fig. 6, we show the arrangement of rupture (red line) and stations
 329 (white triangles). The following subplots are zoomed in the source location and the red dots
 330 show the estimated source points. The background color-map represents Eq. 3 as a two
 331 dimensional function. It is more likely that the current rupture position is in the vicinity of
 332 a point with bright colors than of a point with dark colors. The results of all stations are
 333 equally weighted. For each time-step, the current estimated source location is represented
 334 by a black star and previous most likely positions are marked by white crosses. The beams
 335 of all stations intersect nearly perfectly in one position. An unambiguous trend from top to
 336 bottom is visible and the white crosses match the red rupture line.

337 We repeat the same experiment as presented in Fig. 7 for a more complex rupture ge-
 338 ometry. An animation of the continuous rupture imaging for the complex rupture geometry
 339 is provided in the supporting material (Movie S2: ms02.avi). The rupture propagates on
 340 three horizontally displaced segments of different lengths. The four subfigures show the
 341 rupture tracking at different time steps. Even in this more complex situation, the rupture
 342 is correctly tracked and the fault offsets are visible in the final tracking results.

343 The length of time during which body waves arrive directly varies for different station
 344 locations in dependence on the rupture position. Such directivity effects will cause artifacts
 345 in the tracking result if the information of many stations is combined. In both previous
 346 experiments, we correct for directivity effects by picking the start and end time of the SH-
 347 arrivals in the seismograms. We expect that, in real data, it is difficult to determine the last
 348 arrival of the SH-waves, although the rotational observation facilitates the identification of
 349 shear waves.

350 Thus, an important question arises: How is rupture tracking affected if only the first
 351 arrival is visible in the seismograms? Bayer et al. (2012) neglected directivity effects in a
 352 comparable approach with classic 3C data for P-waves and used the same time-shift for
 353 the BAz estimation window. While conventional back-projection does not include a time
 354 correction for directivity effects (e.g., Ishii et al., 2005), P-wave based rupture tracking
 355 utilizing beam-forming has been shown to require correction for the varying locations of the
 356 seismic sources (e.g., Krüger & Ohrnberger, 2005).

357 Fig. 8 illustrates the impact of directivity effects in an additional numerical experiment.
 358 The rupture is tracked by a small array of four stations which has a relatively small opening
 359 angle. The impact of directivity effects is expected to be significantly smaller for an array

360 with a small opening angle. However, if the angle is too small, it is not possible to determine
 361 both spatial dimensions of the rupture in good quality. In Fig. 8 the station positions are
 362 shown in the first map. In the left subplot, we show the final result for BAZ estimation
 363 in which we corrected for directivity effects. The resolution of the x coordinate is not as
 364 good as in the previous results, since a smaller array is used, but the rupture is still tracked
 365 correctly. The right subplot shows the result for the same data, but this time the time-shift
 366 of the sliding window in the BAZ estimation is assumed the same for all stations. The
 367 starting position is still correctly determined but later estimated points show a systematic
 368 deviation from the true rupture path. Even if the rupture area and its linear trend are
 369 roughly tracked, the geometry is not correctly derived. By neglecting directivity effects it
 370 is possible to track the beginning of the rupture, but not its complete spatial evolution.

371 **4 Rupture tracking in 3-D heterogeneous media**

372 We extend the presented 2-D findings by examining the stability and accuracy of 6-DoF
 373 rupture tracking in 3-D heterogeneous media where multi-phases interfere with each other.

374 The opening angle for a given earthquake, i.e., the detected BAZ variation, depends
 375 on epicentral distance and station azimuth. The resolving power at a single station will
 376 potentially decrease with the increasing epicentral distance while increase with increasing
 377 inclination angles. In the following, we first estimate the expected opening angle for 3-D
 378 rupture tracking of earthquakes as a function of rupture length, epicentral distances and
 379 inclination angles, while avoiding the polarization uncertainty dropping below the noise
 380 level. Then we perform 3-D synthetic tests for rupture tracking with 6-DoF measurements.

381 **4.1 3-D single-station opening angles**

382 We define the opening angle of a specific station as the difference between the BAZ for
 383 the starting position of the unilateral rupture and the BAZ for the ending position (Bayer et
 384 al., 2012). A large opening angle is desirable to minimize uncertainties which corresponds
 385 to a short distance between receivers and earthquakes. But the distance has to be large
 386 enough to fulfill the plane wave assumption while keeping in mind that the analysis of the
 387 polarization is less efficient for signals in which many different phases are interfering.

388 The following description is a purely geometrical concept to demonstrate the expected
 389 scaling of opening angles. If we define the rupture path as a straight line on a sphere, we
 390 can describe the geometry between the receiver and rupture by a large triangle. In Fig. 9
 391 we illustrate the opening angle α for fault lengths between 100 and 1000 km (blue lines) at
 392 epicentral distances d of 10° , 30° , 50° and 80° (different plotting windows). The triangle

is not necessarily isosceles, which is expressed by the inclination angle $\delta \in [0, 90^\circ]$. The maximum opening angle occurs for $\delta = 90^\circ$, i.e., the station is situated perpendicular to the center of the rupture and the triangle is isosceles. In general α increases for larger faults as well as for shorter station distances. By applying the spherical law of cosines, the opening angle α can be described by the side lengths s_1 , s_2 and l , where l denotes the rupture length:

$$\cos(\alpha) = \frac{\cos(l) - \cos(s_1) \cos(s_2)}{\sin(s_1) \sin(s_2)}. \quad (4)$$

The side lengths s_1 and s_2 can be expressed by δ , l and d via

$$\cos(s_{1/2}) = \cos(d) \cos(\frac{l}{2}) \pm \sin(d) \sin(\frac{l}{2}) \cos(\delta). \quad (5)$$

In an epicentral distance of 30° the opening angle for a fault of 1000 km is about 18° (see the upper right window of Fig. 9). For example, the mainly unilateral 2004, Great Sumatra-Andaman earthquake ruptured across a length of about 1200 km. If we increase the inclination angle from 0° to 30° at the same distance, the opening angle decreases to about 9° . This may be still sufficient for an estimation of the spatial and temporal evolution of the rupture with a single station. Significantly shorter ruptures or very small inclination angles, however, will lead to an opening angle of only a few degrees challenging to track.

4.2 Synthetic case studies in 3-D

To verify the stability and accuracy of the proposed rupture tracking approach using 6-DoF measurements in 3-D, we calculate synthetic seismograms using `Instaseis`, an efficient tool for generating synthetic global seismograms using Green's function databases generated in 2.5D axisymmetric spectral element simulations using `AxiSEM` (Driel et al., 2015) and utilizing 1-D axisymmetric velocity models. We here assume the 1-D isotropic PREM (Preliminary Reference Earth Model) model with attenuation effects and the highest frequency of up to 0.5 Hz to generate the synthetic data. Although `Instaseis` does not allow the direct analysis of rotational components, we derive them using a densely spaced array, i.e., gradient-based array-derived rotation.

In the following synthetic tests, we place four additional stations surrounding the central station with a spatial interval of 100 m (see the upper left subplot in Fig. 10). The array-derived rotation is calculated based on a finite-difference scheme (Spudich et al., 1995; Langston, 2007), as the rotational motions will be simplified to horizontal spatial gradients of translational motions at the free surface where vertical stress equals zero (Robertsson & Curtis, 2002). `Instaseis` enables us to handle finite ruptures represented by an arbitrary number of point sources. The simulated rupture consists of six subevents and propagates approximately from south-east to north-west (indicated by the black arrow in Fig. 10).

426 All subevents are assigned a uniform faulting mechanism (strike: 336° , rake: 114° , dip:
 427 7°) and are evenly distributed along the fault plane at the same depth (10 km). The total
 428 rupture length is about 236 km. Considering that the rupture speed and radiated energy can
 429 be largely affected by local structural properties and stress conditions, we slightly randomize
 430 the source time functions of each subevent in terms of slip rate and initiation time (see the
 431 source time functions in Fig. 11b, d, f and h. White noise has been added to each synthetic
 432 dataset, which ends up with a signal-to-noise ratio (SNR) of 25. The SNR value is defined
 433 as $SNR = 10\log_{10}(A_{signal}/A_{noise})$ using the logarithmic decibel scale, where A denotes
 434 the root mean square amplitude in a certain time window (150 s after the corresponding
 435 arrival).

436 The estimated BAz at stations ST1–4 with the epicentral distances of 25° , 14° , 50° and
 437 45° , respectively, is expected to continuously increase during the rupture tracking process
 438 (Fig. 10). For stations ST1–3, we use the singular value decomposition (SVD) algorithm
 439 for robust ratio calculations of Eq. 2 (Greenhalgh et al., 2018), with a sliding time win-
 440 drow of 25 s. The recorded rotational motions in the two horizontal rotational components
 441 are mostly resulting from P-converted SV-waves at the Earth surface at ST1–2 and direct
 442 SV-waves at ST3. For station ST4, we apply the CC method to the vertical rotational
 443 component and the two horizontal translational components, in order to focus on direct
 444 SH-waves, with the same sliding window as the one for ST1–3. We select stations ST3–4
 445 with a larger epicentral distance such that the direct SV- and SH-waves can be separated
 446 in time from surface waves. We generate one dataset for each station. The BAz estimate
 447 for each dataset as a function of time is shown in Fig. 11a, c, e and g (dashed black lines
 448 denote the theoretical starting and ending BAz of the rupture at the given station). The
 449 solid dots represent the theoretical BAz of each subevent. The arrows indicate the onset of
 450 direct P- or S-arrivals. The corresponding source time functions are plotted in Fig. 11b, d,
 451 f and h.

452 We show that the estimated BAz during the rupture tracking process is generally ac-
 453 curate and consistent at all four stations (Fig. 11 using either P-converted SV-waves (at
 454 ST1–2), direct SV- (at ST3) or direct SH-waves (at ST3). However, the slope of the esti-
 455 mated BAz is not ideally uniform.

456 This can be attributed to two factors: 1) the uneven onset time and slip rate across the
 457 finite sources. Theoretically, the slope variation of the estimated BAz is supposed to directly
 458 indicate the changes of rupture speed as we have discussed in section 3.1. However, since we
 459 randomize the source time functions of all subevents, the slope of the estimated BAz should
 460 not be strictly invariant. In Fig. 11d, we notice that there is a stronger subevent between
 461 60 s and 80 s, which may lead to the plateau between 260 s and 280 s in Fig. 11c. The actual

462 BAz estimation within this period is a result of the earlier weaker arrival (with a smaller
 463 BAz) and later stronger arrival (with a bigger BAz), which ends up with a bias towards
 464 the bigger BAz when performing SVD analysis. 2) may be the interference of non-direct
 465 arrivals generated by different subevents. The earlier free-surface related multiples and the
 466 reflections from subsurface discontinuities could both interfere with the later direct arrivals.
 467 This is a major issue we may have to deal with in real data analysis. The intensity of
 468 the interference may vary with different factors, such as radiation patterns of earthquakes,
 469 epicentral distances and 3-D velocity structures.

470 The influence of varying levels of noise and of the length of the sliding window is
 471 evaluated for station ST3 (Fig. 10) based on the polarization analysis of direct SV-waves in
 472 the same rupture scenario introduced earlier in this section. As shown in Fig. 12a, the error
 473 of the estimated BAz increases with the decreasing SNR as expected. When comparing the
 474 time-variant BAz and the theoretical BAz of each subevent (solid dots in Fig. 12) we find
 475 that it is quite stable and accurate in cases with a SNR above 20. The overall trend of the
 476 BAz variation is still preserved despite the BAz errors from signals with relatively low SNR
 477 (≤ 10).

478 The length of sliding windows is another factor that may affect the stability and accu-
 479 racy of the proposed BAz estimation. In Fig. 12b, we fix the SNR at 35 and apply three
 480 distinct window lengths (10 s, 20 s, and 30 s) for analysis with the same highest frequency
 481 in the synthetic data (0.5 Hz). A shorter sliding window theoretically provides higher reso-
 482 lution in identifying subevents, but in practice, we tend to apply longer windows to stabilize
 483 BAz estimation and mitigate the effect of interfering arrivals and/or noisy data as is shown
 484 in Fig. 12b.

485 **5 Real case studies - tracking moving traffic noise sources using 6-DoF** 486 **observation**

487 Due to the lack of real 6-DoF observations of large earthquakes at suitable epicentral dis-
 488 tances, we verify the proposed methodology by tracking moving traffic seismic noise sources,
 489 which might be a useful analogy to the real rupture process regardless of the scale differ-
 490 ence. The collocated classic triaxial seismometer (STS-2) at station FUR belonging to the
 491 German Regional Seismic Network and newly-built ring laser gyroscope ROMY (Gebauer
 492 et al., 2020) at the Geophysical Observatory Fürstenfeldbruck near Munich, Germany, allow
 493 us to record high-fidelity, broadband 6-DoF (particle velocity and rotational rate) ground
 494 motions. In Fig. 13a, we show the site map of the 6-DoF station. As the station is not
 495 far from the highway (blue line in the map), traffic noise is expected to be dominant in the
 496 6-DoF observation for certain frequencies. To avoid overlapping issues and better identify

497 each passing vehicle, we specially choose a data example around midnight when there is
 498 no heavy traffic. After converting the particle velocity record to particle acceleration, the
 499 6-DoF data is then detrended, band-pass filtered to 1-20 Hz and downsampled to 50 Hz.

500 In the top three panels of Fig. 13b, we show a nearly one-hour continuous record of
 501 the two horizontal acceleration components and the vertical rotational rate, from which we
 502 can see consistent traffic-induced signals. Considering the distance and the relative position
 503 of the highway and the station, the BAz variations of inbound vehicles (from southeast to
 504 northwest) should roughly decrease from 100° to 0° and then keep decreasing from 360°
 505 to 300° . The outbound ones are just the reverse. As is shown in the bottom panel of Fig. 13b,
 506 we can see two patterns of the estimated BAz as a function of time indicating the inbound
 507 and outbound vehicles along the highway. The colormap represents the CC coefficient for
 508 each sliding window of 10 s at a certain trial BAz and only the estimated BAz (solid dots)
 509 with a maximum CC larger than 0.4 are shown.

510 We only apply the CC method here focusing on the SH/Love waves as the horizontal
 511 rotational components of the ring laser gyroscope are still under improvement. From the
 512 zoom-in plot of one selected time window (Fig. 13c), we can calculate the moving speed of
 513 the vehicle through the ratio of distance and time over a certain BAz variation range. The
 514 estimated speed when BAz changing from 0° to 70° is approximately 90 km/h, which is
 515 empirically a reasonable estimate at this time of night.

516 **6 Discussion**

517 In this study we explore the potential of using 6-DoF observations to track seismic
 518 sources by exploiting 1) the correlation of translational and rotational motion observations
 519 of SH-waves and 2) the polarization filtering effect of pure horizontal rotational motions. We
 520 demonstrate with synthetic and real data that tracking seismic sources is possible provided
 521 that 6-DoF measurements are taken at an appropriate epicentral distance and direction to
 522 the source.

523 In synthetic tests, we show that direct estimates of rupture velocity can be derived
 524 under sub-Rayleigh and supershear speed variations along the fault. We also show that -
 525 as long as the rupture-induced shear waves can be identified, the directivity effect can (and
 526 should) be corrected for.

527 The presented synthetic models are all unidirectional propagating. We do expect con-
 528 siderably more complexity if 1) the rupture is bilateral, 2) the finite source is extremely
 529 complex or 3) in presence of complex 3-D velocity structures. For any of the complex
 530 rupture scenarios listed above, however, more advanced processing techniques need to be

531 developed in future studies. Despite their simplicity, the presented synthetic examples pave
532 the way for successful observational tracking of moving traffic sources using the highly sen-
533 sitive ring laser gyroscope ROMY.

534 The advantage of the proposed method lies in the fast and relatively easily obtained
535 seismic source tracking using point stations and taking S-waves fully into account; potentially
536 complementing P-wave based methods and fully complex data-driven kinematic or dynamic
537 earthquake source inversion (Ide, 2007; Gallovič et al., 2019, and references therein.).

538 Structural discontinuities, e.g. posed by radially stratified media such as included in the
539 3D synthetic examples in section 4.2, as well as 3D heterogeneity, may introduce interfering
540 P-wave and S-wave arrivals from various epicentral distances. In difference to conventional
541 array-based techniques, 6-DoF measurements can mostly mitigate the P-coda interference
542 due to the nature of rotational motions allowing to image seismic source processes using
543 direct S-waves. Severe interference of multiple types of waves may hinder the here pre-
544 sented application of rupture tracking methods using S-waves. However, this issue might be
545 mitigated in case of 6-DoF measurements thanks to the inherent wavefield separation in the
546 rotational components, i.e., only SV-waves or Rayleigh waves are presented in horizontal
547 rotation and only SH-waves or Loves wave are presented in vertical rotation. As is shown in
548 Fig. 11e–h, we are able to capture the rupture process when applying both the polarization
549 and the CC method onto direct SV-waves and direct SH-waves, respectively. This illustrates
550 potential as a useful complement to classical back-projection earthquake rupture imaging
551 which solely relies on P-wave information.

552 Through the traffic-induced seismic noise real data example, we have further verified
553 the feasibility and effectiveness of tracking various seismic sources with 6-DoF point mea-
554 surement. The results show that moving sources can be precisely located. The vehicle
555 speed estimates are also reasonable assuming a known highway location. The real data
556 example resembles the rupture imaging process assuming finite earthquake sources and may
557 be readily applied to future 6-DoF observations of earthquake rupture.

558 Compared to the seismic array method, 6-DoF point measurements are superior for
559 alleviating site effects. We note that the complexity in both earthquake rupture and wave
560 propagation within 3-D Earth may challenge BAZ estimates. With the advent of the first
561 commercial broadband portable rotational seismometer systems (Bernauer et al., 2018; Yuan
562 et al., 2020), direct observations of 6-DoF motions of large earthquakes becomes feasible.
563 The sensitivity of the instrument allows for recording large earthquakes with high signal-
564 to-noise ratios. The presented methods with respect to rotational seismology are applicable
565 to (combined) strain observations. With the increasing accuracy of distributed acoustic

sensing (DAS) type measurements application of this method to DAS observations should be further explored (Lindsey et al., 2017; Jousset et al., 2018; Yu et al., 2019).

7 Conclusions

Six degree-of-freedom (6-DoF) single-station observations allow the extraction of wave-field information comparable to small-scale seismic arrays (e.g., Igel et al., 2015; Schmelzbach et al., 2018; Sollberger et al., 2018). In particular, estimates of phase velocities and sub-receiver physical velocities, incidence and BAZ angles are possible. We show that such 6-DoF observations allow in principle to track the location of sources of seismic energy and discuss sensitivity and challenges to methods based on cross-correlation or polarization analysis, respectively. Investigating the potential of emerging 6-DoF observations in the context of earthquake physics, the here developed approaches can be generalized to arbitrary sources of seismic energy such as environmental sources, volcanic sources and atmospheric sources as well as to distributed acoustic sensing (DAS) type strain measurements. While we demonstrate the potential of the proposed point measurement for fast and easy seismic source tracking, future efforts may focus on accounting for complexities such as bilateral earthquake rupture, geometric source complexity and complex 3-D velocity structures.

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The open-source 2-D wave propagation package `se2wave` is available at <https://bitbucket.org/dmay/se2wave>. The 3-D experiments use the openly accessible software `instaseis` <http://instaseis.net>, which databases are hosted by the IRIS DMC via the Syngine webservice backend Krischer et al. (2017). Synthetic and real datasets used in

600 this study are available from ZENODO: <https://doi.org/10.5281/zenodo.4011759>. The seis-
601 mometer data recorded at station FUR in the observatory is also available by accessing
602 EIDA data archive through <http://eida.gfz-potsdam.de/webdc3/>.

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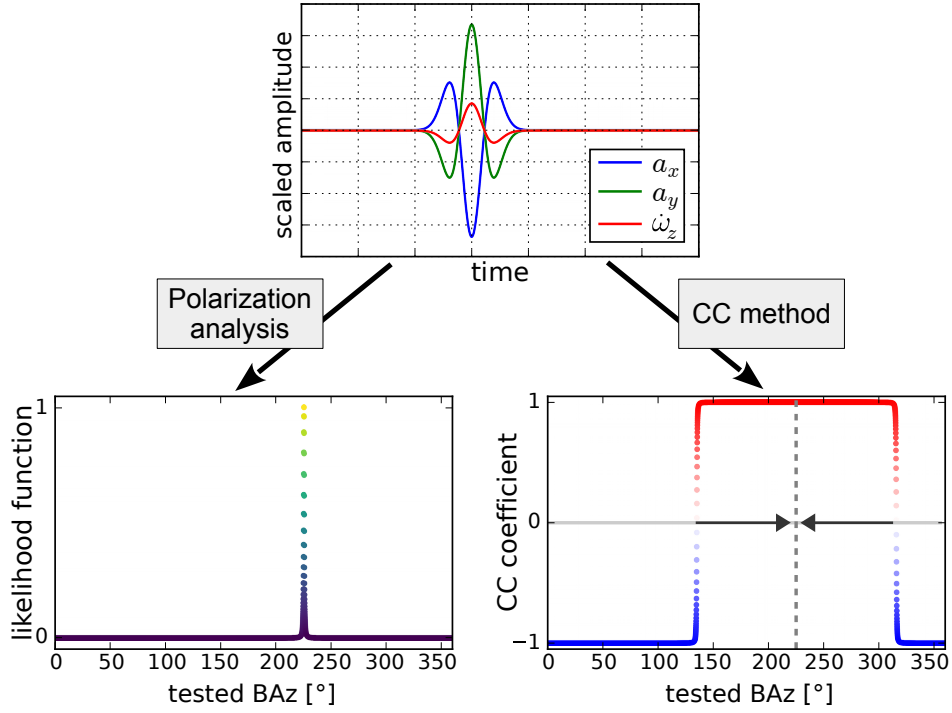


Figure 1. The actual BAz of 225° of a plane SH-wave can be estimated by the CC method and the polarization analysis. **Top:** the seismogram of an ideal SH-wave at the free surface contains three nonzero components (the horizontal accelerations a_x , a_y , and the vertical rotation rate $\dot{\omega}_z$). **Bottom Left:** the grid search result of the polarization analysis. The normalized likelihood function has a maximum at 225° . **Bottom Right:** the grid search result of the CC method. The zero-lag CC coefficient as a function of the tested BAz has the shape of a step function. The BAz is correctly estimated by determining the central point between the zero-transitions (indicated by arrows).

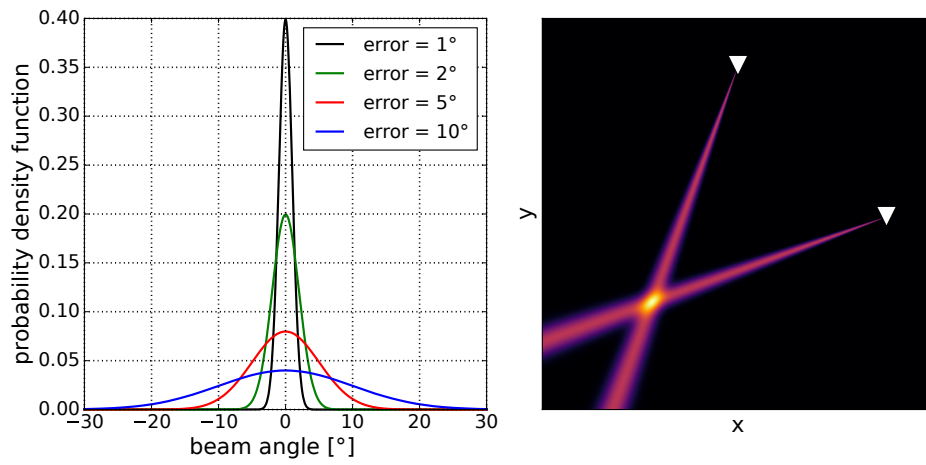


Figure 2. The BAz uncertainties can be described by a shape function. **Left:** each curve shows the shape function for a specific BAz error σ . The beam angle describes the broadness around the estimated BAz which corresponds to 0° . **Right:** Eq. 3 is illustrated in 2-D at a certain time-step for two stations (white triangles). An expected error of $\sigma = 1.5^\circ$ describes the broadness of each beam. Its highest value (yellow area) is expected to be close to the source position.

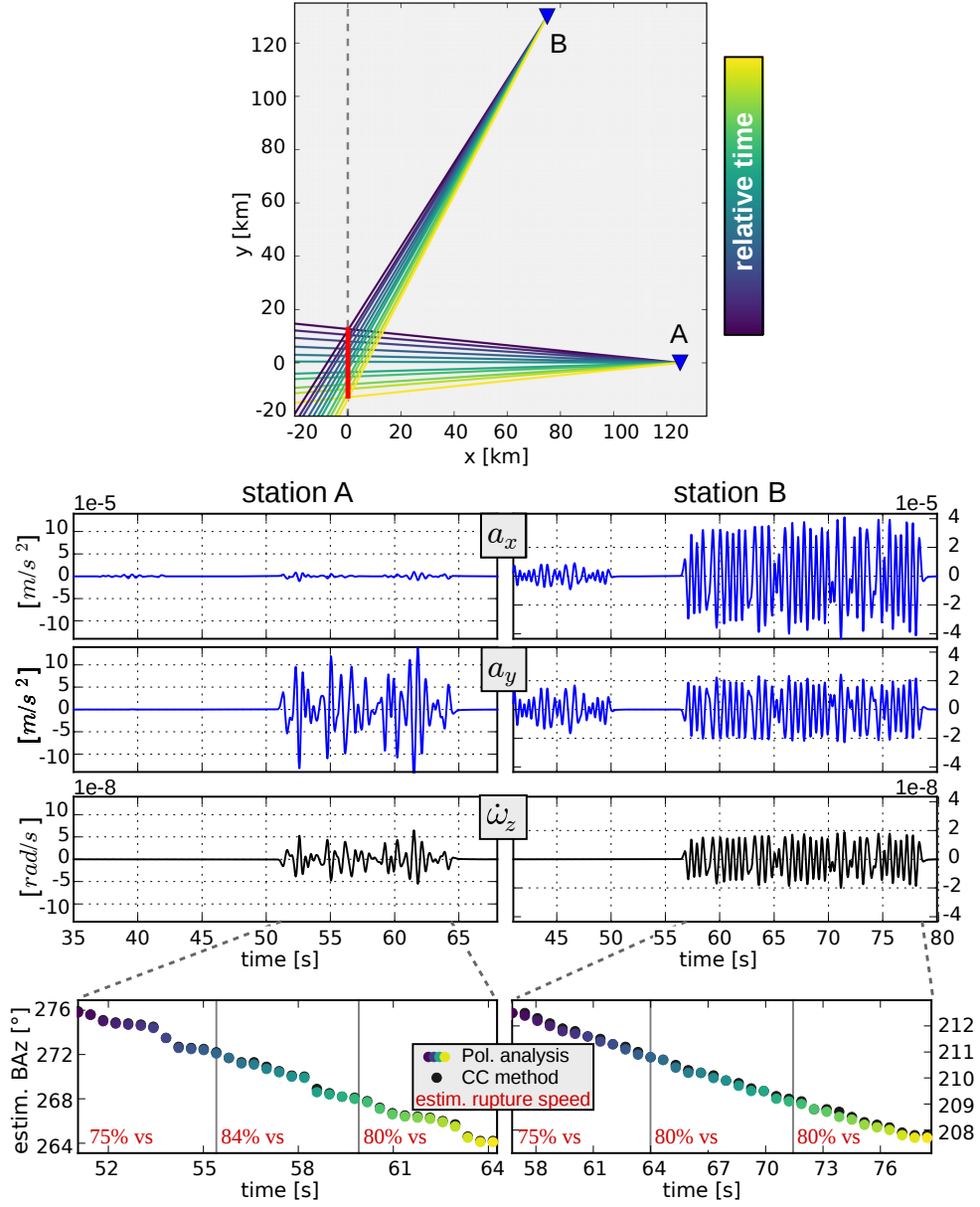


Figure 3. Using SH-waves for rupture tracking in a single station approach. **Top:** rupture (red line) and receiver positions (blue triangles) are shown in horizontal dimensions x, y . The color-coded rays indicate the estimated BAZ variations at each station. The rupture is correctly tracked. **Middle:** recorded seismograms of horizontal accelerations a_x , a_y and vertical rotation rate $\dot{\omega}_z$. **Bottom:** the BAZ is estimated by two different methods from the direct SH-arrivals. The estimates are divided into three sub-windows in each of which rupture speed is determined. The true rupture velocity is $80\% v_s$.

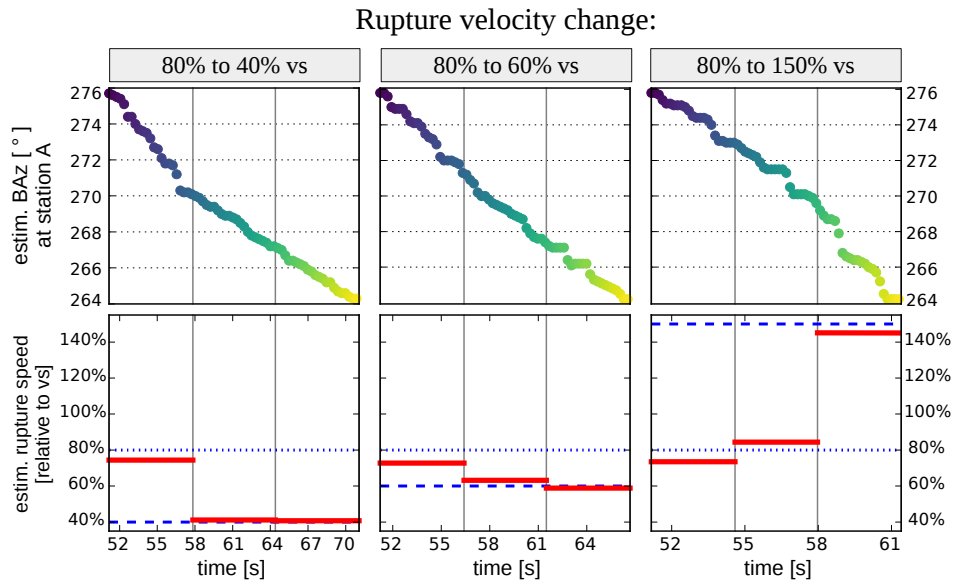


Figure 4. Tracking variations in rupture velocity at station A. The first half of the rupture propagates with 80% of the shear-wave velocity v_s , the second half ruptures with 40%, 60% or 150% v_s . The BAz variation is estimated from the SH-arrivals at station A (see Fig. 3). The results are represented by a red horizontal line in each sub-window. The actual velocities are indicated by blue dashed lines. For consistency reasons, we use the same color-scale from yellow to blue in the upper subplots.

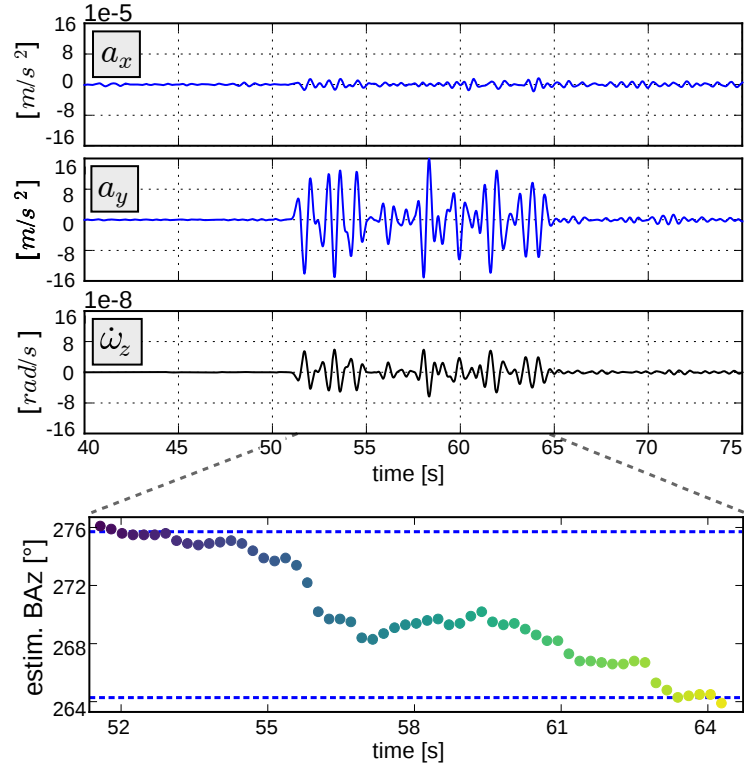


Figure 5. Rupture tracking in a heterogeneous medium at station A. Density, P-wave and S-wave velocity are perturbed independently in each material cell. **Top:** seismograms of horizontal accelerations a_x , a_y and vertical rotation rate $\dot{\omega}_z$. **Bottom:** result of the BAz estimation. The true BAz for start- and end-position of the rupture is marked by blue dashed lines.

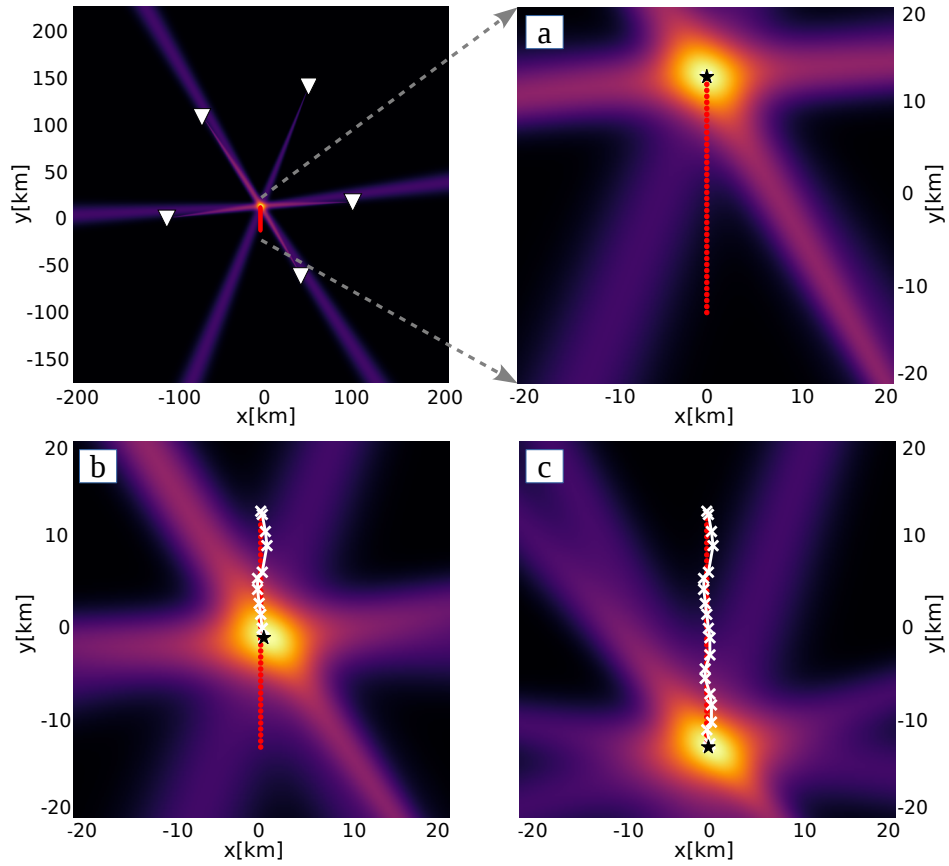


Figure 6. A unilateral rupture (red dots) is observed at five stations (white triangles) in two spatial dimensions x, y . At each station, the BAz change is estimated independently and is corrected for directivity effects. Subplots a, b and c show the result at different time-steps. The most likely source locations are marked by black stars and white crosses indicating current and past time steps, respectively. An animation of the continuous rupture imaging process is provided in the supporting material (Movie S1: ms01.avi).

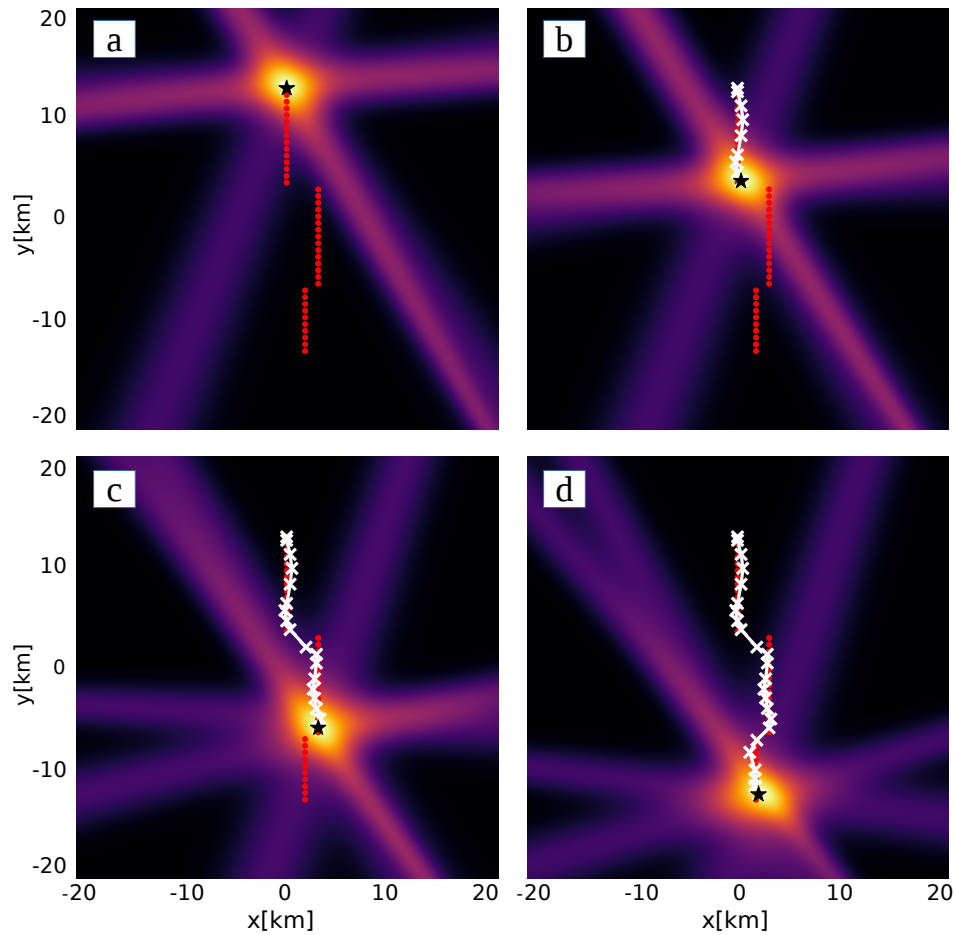


Figure 7. Rupture tracking for a more complex fault geometry. Station locations and model parameters are the same as in Fig. 6. The rupture has two spatial offsets. Subplots a, b, c and d show the tracking at different time steps. An animation of the continuous rupture imaging process is provided in the supporting material (Movie S2: ms02.avi).

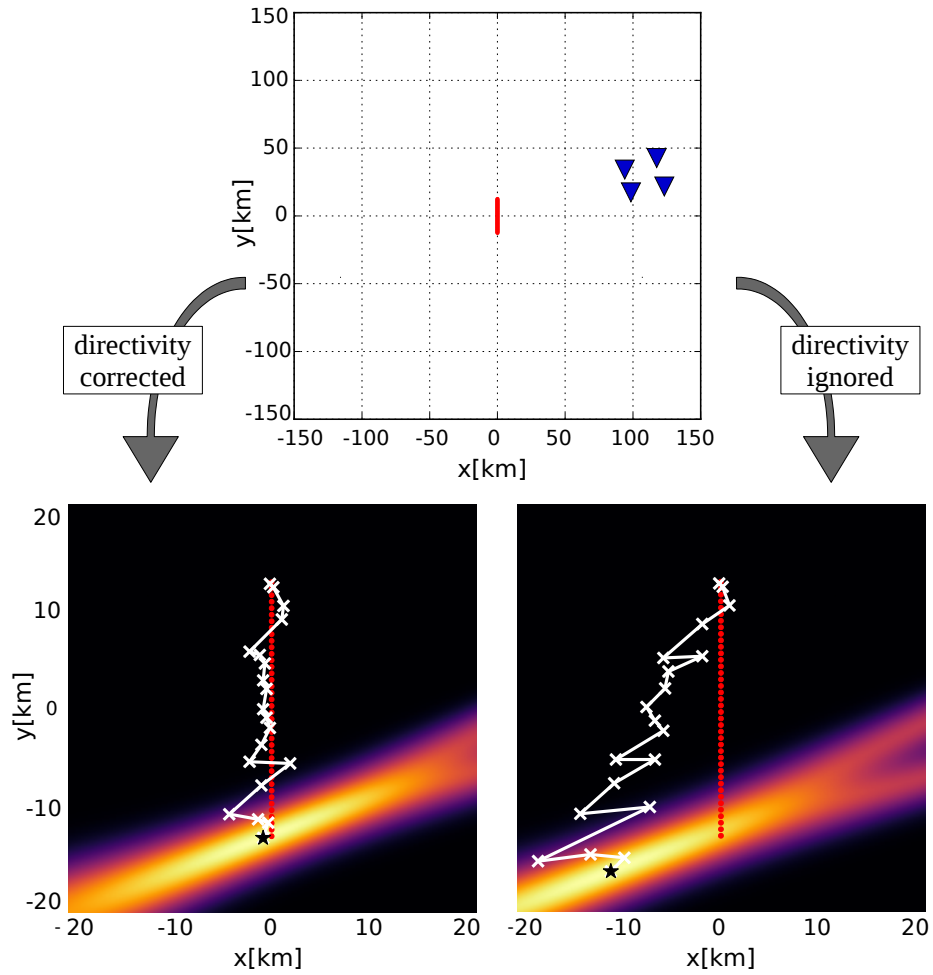


Figure 8. Influences of directivity effects on the combination of tracking results. A rupture is tracked with four stations in a small network (blue triangles). The rupture properties are the same as in Fig. 6. The bottom left shows the tracking result for corrected directivity effects. The bottom right figure shows the results when neglecting directivity effects.

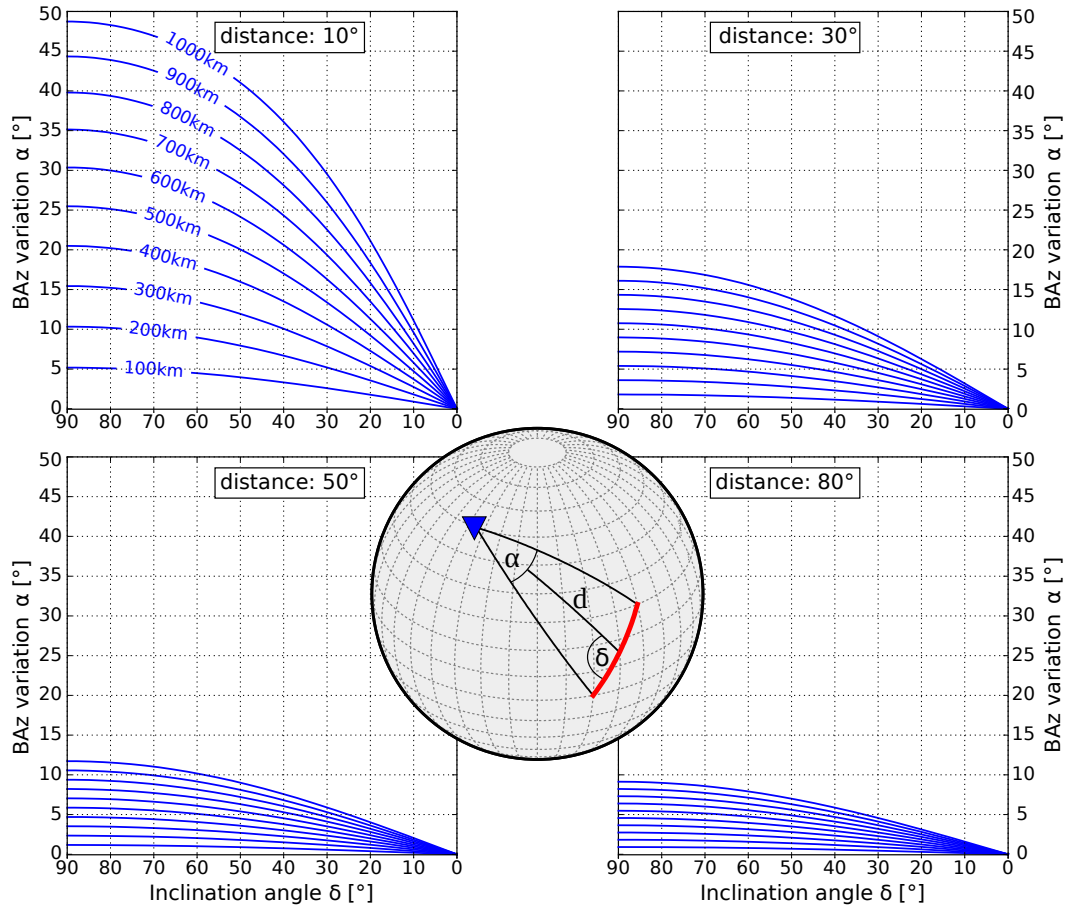


Figure 9. The opening angle α (Eq. 4) for a single station depends on the epicentral distance d , the inclination angle δ and the rupture length (adapted from Bayer et al. (2012)). The rupture path is illustrated as a red line on a sphere. The opening angle is calculated for four different distances and in each subplot for fault lengths between 100 km and 1000 km.

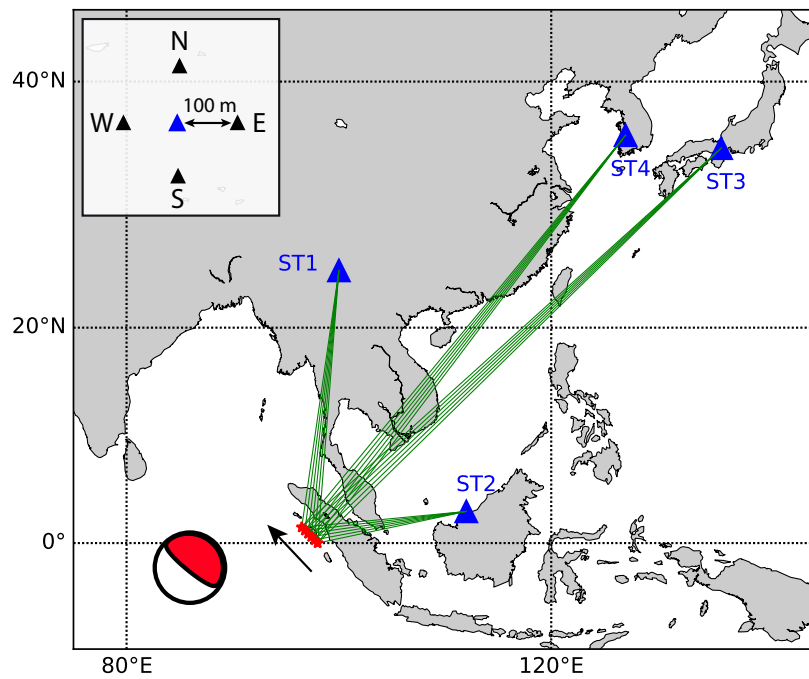


Figure 10. Modeled unilateral earthquake rupture and four seismic stations ST1–4 (blue triangles). Red dots denote subevents of the rupture process. The beach ball in the left-bottom corner denotes the uniform focal mechanism of all subevents. The black arrow indicates the rupture direction. The upper left subplot illustrates the surrounding four stations which are used to derive rotational motions at the central station.

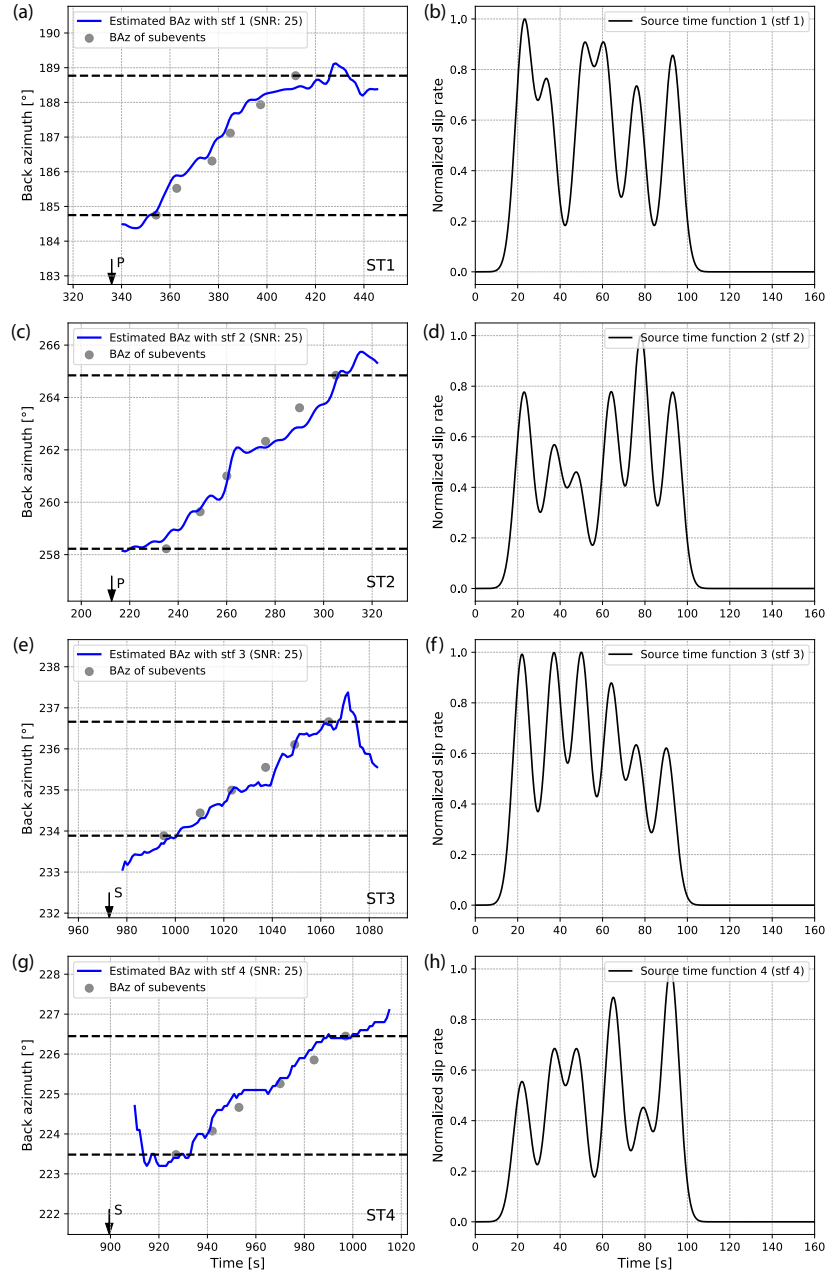


Figure 11. Earthquake rupture tracking in 3-D heterogeneous media. The arrow indicates the onset of P or S waves. Solid dots represent the theoretical BAZ of each subevent. (a–b): Estimated BAZ variation as a function of time for one synthetic dataset at the station ST1 and the corresponding source time function (stf 1). The blue line connects all estimated BAZ of each sliding window (25 s) and the dashed black lines denote the theoretical starting and ending BAZ of the rupture. (c–d): Same as (a–b) but for the station ST2. (e–f): Same as (a–b) but using direct SV-waves recorded at station ST3. (g–h): Same as (a–b) but using direct SH-waves recorded at station ST4.

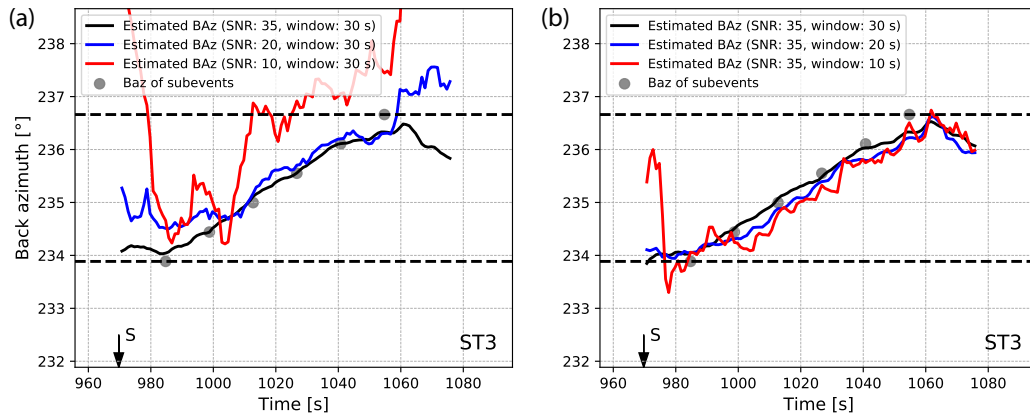


Figure 12. Effect of signal-to-noise ratio (SNR) and sliding window length on BAZ estimation using direct SV-waves at station ST3. (a) The estimated BAZ variation with different white noise levels (SNR: 35, 20 and 10) for a fixed window length (30 s). (b) The estimated BAZ variation with different sliding window lengths (30, 20 and 10 s) for a fixed SNR value (35). Solid dots represent the theoretical BAZ of each subevent.

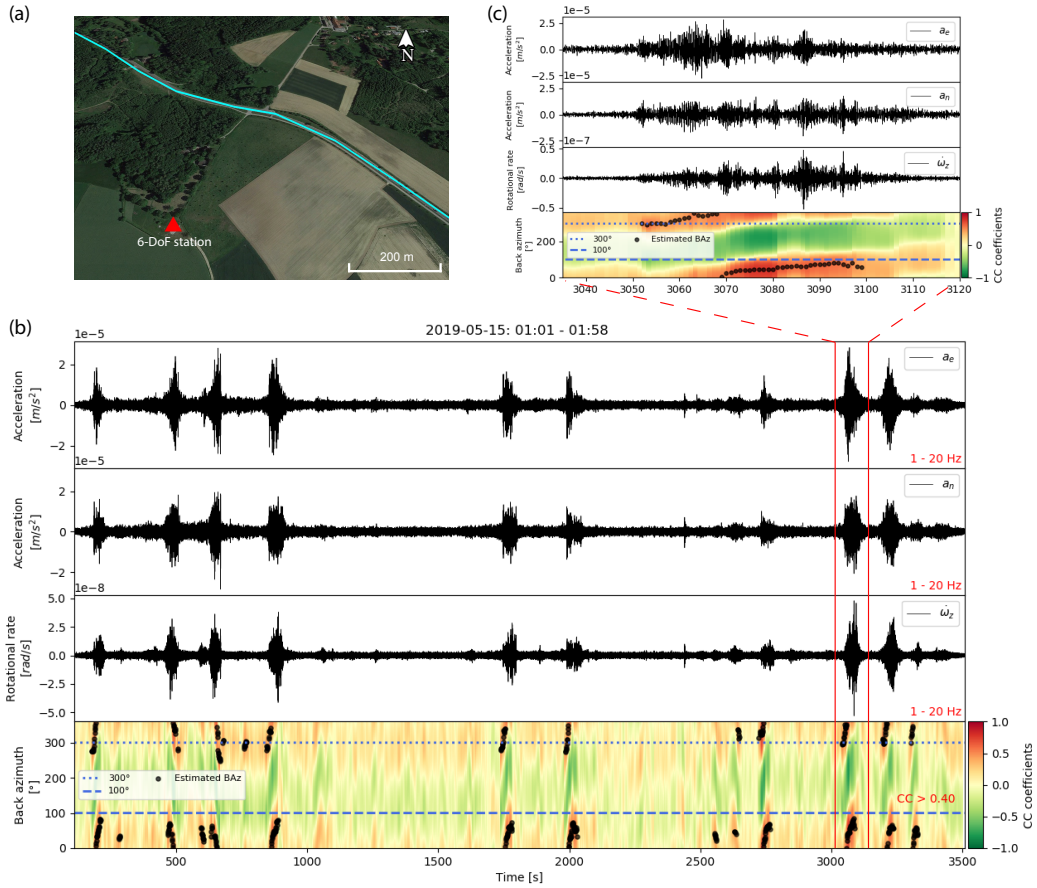


Figure 13. Site map at the Geophysical Observatory Fürstentfeldbruck near Munich, Germany and estimated BAZ of the traffic-induced seismic noise (nearly one-hour continuous data) from 6-DoF point measurement. (a) The red triangle indicates the position where a STS-2 seismometer and a ring laser gyroscope are collocated. The blue curve denotes the highway next to the observatory. (b) From top to bottom: the north-south and east-west components of acceleration, the vertical rotational rate and the estimated BAZ from a_e/a_n and $\dot{\omega}_z$ components. Black dots represent the estimated BAZ for the sliding window, with the CC coefficient being higher than 0.4. The dotted and dashed lines indicate the expected range of the BAZ variation for the inbound and outbound moving cars. (c) The zoom-in plot of (b) within the two solid red lines.