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1	Enhanced iceberg discharge in the western North Atlantic during all Heinrich
2	events of the last glaciation
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17	Abstract
18	A series of catastrophic iceberg discharges to the North Atlantic, termed Heinrich events,
19	punctuated the last ice age. During Heinrich events, coarse terrigenous debris released from the
20	drifting icebergs was preserved in deep-sea sediments, serving as an indicator of iceberg passage.
21	Quantifying the vertical flux of ice-rafted debris (IRD) in open-ocean settings can resolve

22 questions regarding the timing and spatial variation in ice sheet calving intensity. In this study, ²³⁰Th_{xs}-based IRD flux throughout the last glacial period was measured in a deep-sea sediment 23 24 core from the western North Atlantic, and complemented by data spanning 0-32 ka from a sediment 25 core in the Labrador Sea. The cores were recovered from sites downstream from Hudson Strait, a likely conduit for icebergs calving from the Laurentide ice sheet (LIS). We compare our results 26 with equivalent existing data from the eastern North Atlantic and show that the two cores in our 27 study have higher IRD fluxes during all Heinrich events, notably including events H3 (~31 ka) and 28 29 H6 (~60 ka). This study demonstrates that the LIS played a role in all Heinrich events, and raises 30 the likelihood that a single mechanism can account for the genesis of these events.

31 Keywords: Heinrich event; western North Atlantic; IRD flux; thorium normalization

32 1. Introduction

The last glaciation was characterized by two types of millennial-scale climate change. The classic 33 Greenland ice core records document Dansgaard-Oeschger (D-O) events, rapid shifts in air 34 35 temperature over the ice cap of ~8 °C within a few decades (Dansgaard et al., 1993; Grootes et al., 36 1993; Johnsen et al., 1992; NGRIP members, 2004). Throughout the same interval, subpolar North 37 Atlantic deep-sea sediments preserve evidence of century-scale episodes of catastrophic iceberg 38 discharge known as Heinrich events (HEs). The passage of icebergs is marked by ice-rafted debris 39 (IRD), which forms layers of detrital material in glacial sediment sequences (Bond et al., 1992; 40 Broecker et al., 1992; Heinrich, 1988; Hemming, 2004).

41 The icebergs discharged during HEs drifted and melted across the subpolar North Atlantic Ocean, 42 an important region of deep-water formation that is sensitive to disruption by surface water 43 freshening (Manabe and Stouffer, 1997; Rahmstorf, 1995). Indeed, the freshwater flux associated with these events is thought to have dramatically and repeatedly weakened the Atlantic Meridional 44 45 Overturning Circulation (AMOC) (Henry et al., 2016; McManus et al., 2004). As a result, the 46 Northern Hemisphere cooled (Bard et al., 2000; Bond et al., 1993) while the Southern Hemisphere 47 warmed (Barker et al., 2009; Buizert et al., 2018). The oceans, cryosphere, and atmosphere all respond in connected ways during HEs, making these events ideal targets for examining how 48 49 abrupt climate change unfolds.

An unsolved mystery of HEs is the apparently anomalous behavior associated with H3 (~31 ka) and H6 (~60 ka). These two events are often regarded as atypical because of their relatively low IRD concentrations and the absence of detrital carbonates in the eastern subpolar North Atlantic (Broecker et al., 1992). In this paper, we will refer to the other HEs, H1, H2, H4, and H5, as the "typical" HEs and H3 and H6 as "atypical". The observed differences between the typical and the 55 atypical events have been explained by the variable behavior of the ocean or the cryosphere. The 56 robustness of these divergent observations is relevant to understanding the mechanism(s) accounting for these dramatic events. Similar behavior of the events would fit more readily into a 57 58 single mechanism, while contrasting behavior suggests different, even drastically divergent, 59 mechanisms or causes at different times. From a glaciological perspective, the provenance of IRD 60 in H3 and H6 suggests a European ice sheet origin that led to smaller iceberg fluxes in the western 61 North Atlantic (Grousset et al., 2000, 1993; Peck et al., 2007; Snoeckx et al., 1999), although this 62 claim has been disputed (Jullien et al., 2006). A complementary hypothesis suggests that H3 and 63 H6 occurred at the onset of Marine Isotope Stage (MIS) 2 and MIS4 when the LIS was just starting to grow, resulting in a smaller magnitude of calving (Gwiazda et al., 1996). From an oceanographic 64 perspective, assuming that all the events originated from the LIS, warmer sea-surface temperatures 65 might have melted the icebergs closer to their western source during the two atypical HEs (Bond 66 67 et al., 1992). Consistent with that view, Pb isotope data suggest that, in the eastern North Atlantic, 68 the relatively IRD-rich layers are actually the results of foraminifera-dissolution events (Gwiazda 69 et al., 1996). The lack of foraminifera in these layers would cause an apparent increase in the proportion of IRD, the other major coarse constituent in North Atlantic sediments. 70

Reconstructing the IRD deposition rate provides a straightforward test of these competing hypotheses. If the increased IRD concentrations in H3 and H6 are solely the result of foraminifera dissolution, the IRD fluxes (deposition rate per unit area per unit time) during these periods should not increase. If the icebergs originated and mostly melted in the west during the atypical HEs, the IRD fluxes would have been higher in the west than the east. The atypical HE IRD fluxes might also have been even higher than during the typical HEs at the western sites, assuming the total magnitude of calving remained the same. If the Laurentide ice sheet calving increased by a limited amount during the atypical HEs, the IRD fluxes would have shown a visible but smaller increase
than the typical HEs throughout the basin. Lastly, if the two events did originate from the European
ice sheet, the IRD fluxes in the east should have been higher than the west.

Resolving IRD fluxes during HEs provides a way to test the aforementioned hypotheses regarding 81 82 the two atypical events. This hypothesis-testing can be achieved with ²³⁰Th normalization. The rapid scavenging of ²³⁰Th results in a near balance between its burial and local production, allowing 83 84 it to be used to reconstruct vertical mass flux (Bacon, 1984; Bacon and Anderson, 1982; Bacon and Rosholt, 1982; Costa et al., 2020; Francois et al., 2004). IRD flux can therefore be 85 reconstructed from the ²³⁰Th-normalized burial flux of sediment that is uniquely identified as IRD. 86 Previously, century-scale measurements of IRD have only been made in the eastern North Atlantic 87 88 (McManus et al., 1998). Here we present a 135-thousand-year (kyr) ice-rafting record off the coast 89 of Newfoundland, 2200 km downstream from the Hudson Strait. We complement it with a record from Orphan Knoll that is ~1700 km downstream from the Hudson Strait and extends through H3. 90 91 By comparing our ice-rafting records with the directly equivalent existing data in the eastern North 92 Atlantic (McManus et al., 1998), we demonstrate that H3 and H6 are mechanistically consistent 93 with the other HEs of the last glaciation.

94 **2.** Methods

EW9303-37JPC (43.68°N, 46.28°W, 3981m, IGSN: DSR000507, JPC37 hereafter) is a 13.315m-long jumbo piston core retrieved at the foot of a continental slope off the coast of Newfoundland,
Canada (Figure 1). DY081-GVY001 (50°09'36''N, 45°30'36''W, 3721m, GVY001 hereafter) is
a gravity core retrieved near Orphan Knoll (Figure 1, Hendry et al. 2019). Both sites are presently
under the influence of the surficial Labrador Current, which transports icebergs southward in the

Labrador Sea. At depth, the deep western boundary current (DWBC) also flows southward pastthe sites (McCartney, 1992).

Samples of 8-10 g were taken at 2 cm intervals from JPC37, and at 5 cm intervals from GVY001, 102 103 then freeze dried, weighed, and washed through 63 µm sieves to separate coarse and fine fractions. The proportion of coarse sediment (% coarse) was calculated by dividing the dried >63 µm fraction 104 weight by the total dry weight. Subsequently, the coarse fraction was dry-sieved at $>150 \mu m$ and 105 split so that 300-400 foraminifera shells and ~100 IRD grains could be identified and counted 106 under a microscope. Foraminifera counts were used to calculate the relative abundance of the polar 107 108 planktic foraminifera, Neogloboquadrina pachyderma (hereafter N. pachy.), calculated as the 109 number of *N. pachy*. specimens in a sample divided by the total number of planktic foraminifera shells. The relative abundance of IRD (hereafter % IRD) was calculated as the number of IRD 110 111 grains divided by the sum of IRD grains plus planktic foraminifera shells.

112 During microscope counting work, eight to ten specimens of *N. pachy.* were picked for δ^{18} O and 113 δ^{13} C analysis in the Lamont-Doherty Earth Observatory of Columbia University (LDEO) stable 114 isotope laboratory, using a Thermo Delta V Plus equipped with a Kiel IV individual acid-bath 115 sample preparation device. Measurements made on standard carbonate NBS19 yield a standard 116 deviation of 0.06% for δ^{18} O and 0.03% for δ^{13} C.

Elemental intensities from JPC37 were measured using an X-ray fluorescence (XRF) core scanner (ITRAX, Cox Ltd., Sweden) at LDEO at 2mm resolution, using an integration time of 10s and a molybdenum x-ray source set to 30 kV and 50 mA. The intensities were calibrated with flux fusion concentration measurements from JPC37, following the procedure of Murray et al. (2000). Sediment samples spanning the range of XRF intensities were randomized, interspersed with standard reference materials (e.g., JLS-1, JDO-1, SCO-1, AGV-2, DTS-2b, W-2a),
ashed, digested in HF and HNO₃ overnight, and diluted for analysis on an Agilent 720 Inductively
Coupled Plasma Optical Emission Spectrometer (ICP-OES) at LDEO. The ICP-OES data provide
a robust calibration (R²=0.80 for Ca, R²=0.90 for Sr) for the high-resolution XRF records.
Elemental intensities from GVY001 were measured from a similar ITRAX XRF core scanner at
the British Ocean Sediment Core Research Facility (BOSCORF; Hendry et al. 2019). The
elemental intensities from GVY001 are uncalibrated and presented as ratios.

Bulk wet density and magnetic susceptibility from JPC37 were obtained at 1cm resolution from a Geotek multi-sensor core logger at the Lamont-Doherty Core Repository. Local anomalously lowdensity peaks were identified as cracks formed due to drying and omitted in data processing. Bulk wet density and magnetic susceptibility from GVY001 were obtained at 1cm resolution from a similar Geotek instrument at BOSCORF.

Each bulk sediment sample of ~100 mg was spiked, digested (Fleisher and Anderson, 1991),
purified (Lao et al., 1993), and analyzed for uranium and thorium isotope activities. The last step
was done on an Element Plus inductively coupled plasma mass spectrometer (ICP-MS) at LDEO.
The conversion from raw counting data to activities and associated error propagation has been
packaged into a Python script named ThxsPy accessible at https://github.com/yz3062/ThxsPy.

Sediment fluxes were calculated using ²³⁰Th_{xs}, which is the ²³⁰Th derived from the decay of ²³⁴U
in seawater and subsequently scavenged by adsorption onto settling particles. To calculate ²³⁰Thxs,
other sources of non-excess ²³⁰Th need to be quantified and removed (Bacon, 1984; Costa and
McManus, 2017; Francois et al., 2004, 1990).

First, detrital ²³⁰Th is the ²³⁰Th produced from the radioactive decay of ²³⁸U in mineral lattices.
Detrital ²³⁸U can be estimated by assuming a constant detrital ²³⁸U/²³²Th and that all measured

145 232 Th is detrital. A range of $(^{238}\text{U}/^{232}\text{Th})_{\text{detrital}}$ between 0.47 and 0.7 has been used by previous 146 studies in this region (Table 1). We conducted a leaching experiment in JPC37 to isolate the detrital 147 uranium and thorium and determined $(^{238}\text{U}/^{232}\text{Th})_{\text{detrital}}$ to be 0.48. We further determined from the 148 leaching experiment the disequilibrium in $(^{230}\text{Th}/^{238}\text{U})_{\text{detrital}}$ caused by α recoil to be 0.81 (details 149 in Discussion).

- Next, ²³⁰Th can also be produced from the radioactive decay of ²³⁸U that precipitated from the
 soluble form U(VI) to its insoluble form U(IV) in anoxic, reducing sediments (Barnes and Cochran,
 1990; Klinkhammer and Palmer, 1991). Assuming the non-detrital portion of ²³⁸U is authigenic,
 the derived ²³⁰Th (abbreviated to authigenic ²³⁰Th) can be calculated using a seawater ²³⁴U/²³⁸U of
 1.1468 (Andersen et al., 2010) and applying the classic radio-decay equations.
- 155 In summary, the calculation of 230 Th_{xs} is thus
- $156 \qquad {}^{230}Th_{xs}={}^{230}Th_{measured}-{}^{230}Th_{detrital}-{}^{230}Th_{authigenic}$
- $157 = {}^{230}Th_{measured} ({}^{238}U/{}^{232}Th)_{detrital} * ({}^{230}Th/{}^{238}U)_{detrital} * {}^{232}Th_{measured} ({}^{238}U_{measured} {}^{238}U_{measured} + {}^{230}Th/{}^{238}U_{measured} + {}^{230}Th/{}^{23}U_{measured} + {}^{230}Th/{}^{23}U_{measured} + {}^{230}Th/{}^{23}U_{measured} + {}^{230}Th/{}^{23}U_{measured} + {}^{230}Th/{}^{23}U_{measured} + {}^{23$
- $158 \quad (^{238}U/^{232}Th)_{detrital} * (^{230}Th/^{238}U)_{detrital} * ^{232}Th_{measured}) * [(1-e^{-\lambda 230^{*}t}) + \lambda_{230}/(\lambda_{230}-\lambda_{234})(e^{-\lambda 234^{*}t} e^{-\lambda_{234}})(e^{-\lambda_{2$
- 159 λ^{230*t})(($^{234}U/^{238}U$)seawater-1)]
- 160 Where (²³⁸U/²³²Th)_{detrital} is 0.48, (²³⁰Th/²³⁸U)_{detrital} is 0.81, λ is the isotope decay constant,
 161 (²³⁴U/²³⁸U)_{seawater} is 1.1468, and t is the time of decay since deposition.
- The values of (²³⁸U/²³²Th)_{detrital} and (²³⁰Th/²³⁸U)_{detrital} used by this study are different from most previous studies and are based on experimental results. We conducted leaching experiments throughout core JPC37 to obtain representative lattice-bound isotopic ratios. Bulk samples were leached with 5mL of 1N or 3N HCl and sonicated for 20 minutes. After a 5-minute centrifuge, the supernatant was decanted. In some cases, the supernatant was filtered using 0.42 μm filters. The

167 rest of the procedure is the same as described above for typical sediment U-Th analysis. The strength of acid and time of sonication used were shown previously to remove authigenic uranium 168 effectively without leaching lattice-bound uranium and thorium (Robinson et al., 2008). 169

From ²³⁰Th_{xs}, we can calculate the vertical mass flux

 $F = \beta * Z /^{230} T h_{xs,0}$ 171

170

Where F is the vertical mass flux, β is the production rate of ²³⁰Th, Z is the water depth, and 172 ²³⁰Th_{xs},0 is ²³⁰Th_{xs} corrected for decay since deposition using the independent age model. 173 Everywhere else in this paper, we use 230 Th_{xs} as a shorthand of 230 Th_{xs,0}. We can further calculate 174 the IRD flux 175

IRD flux = F * #IRD / M176

where #IRD is the total number of IRD grains, and M is the dry bulk mass. 177

Cross correlation was performed on the time series of δ^{18} O of *N. pachy.* and ²³⁰Th_{xs} to identify any 178 179 lead or lag between the two. The 65-10 ka portions of both proxies were first extrapolated to a 180 common time step of 500 years. The pairwise correlation was then computed for the two time 181 series, first on their original chronology, then by offsetting the two by both positive and negative time steps. In each direction, up to 15 steps of offset were tested and the correlation coefficient 182 183 calculated in each. The cross correlation was then bootstrapped 1000 times, allowing sampling 184 interval and time step to vary.

3. Chronology 185

The chronology of JPC37 from modern to ~46 ka is based on radiocarbon dating. A monospecific 186 sample of 350-400 specimens of N. pachy., the most abundant foraminifera species in this core, 187 were picked from the >150µm fraction. These foraminifera shells were sonicated in water and 188

189 ethanol to remove any fine-grained sediment, including detrital carbonate that can potentially bias 190 the results. Radiocarbon analyses were performed at NOSAMS-WHOI facility (see Appendix A 191 for a link to data). The calibration to calendar ages uses Marine 13 and the CALIB program 192 (Stuiver et al., 2019). We used 400 years as the marine reservoir correction (Reimer et al., 2013). 193 The chronology beyond the range of radiocarbon is determined by aligning our % N. pachy. with 194 an alkenone-based sea-surface temperatures (SST) record from MD01-2444 (Martrat et al., 2007). 195 The polar foraminifera N. pachy. lives in the coldest environment among planktic species and its 196 abundance in JPC37 signals low SST (Ericson, 1959). MD01-2444 (37°33.68'N, 10°08.53'W, 2637 m) is on the Iberian margin and its latitude is comparable to JPC37. Its 197 198 chronology is based on temperature alignment with the North Greenland Ice Sheet Project (NGRIP) 199 ice core during the last glacial (Anderson et al., 2006; Johnsen et al., 2001; NGRIP members, 2004; 200 Rasmussen et al., 2006) and before that, Antarctica Dome C (EPICA Community Members, 2004; 201 Parrenin et al., 2004). Even though the two sites sit on the opposite side of the North Atlantic basin, 202 the prevailing westerlies and the resulting current put MD01-2444 downstream from JPC37. As a 203 result, the temperatures at the two sites should be closely correlated during regional-scale climate changes. The advantage of aligning with this SST record instead of directly with NGRIP is twofold. 204 205 First, both records reconstruct SST so the correlation does not need to consider meridional or air-206 sea signal propagation. Second, NGRIP stops at MIS 5e, whereas the SST records extend beyond 207 that, allowing us to align the late MIS 6 segment of our core as well.

The core chronology is corroborated with tephrochronology. We found several concentrated zones of plates of glassy, bubble-wall shards, consistent with the description of Ruddiman and Glover (1972). The shards are mostly non-existent outside of the zones. The youngest of these zones, at 70 cm, is identified Ash Zone 1. The commonly assumed age of Ash Zone 1 (12.2 cal ka BP) (Andrews and Voelker, 2018) is only 700 years apart from the radiocarbon-based age model. Ash
Zone 2 is identified at 628.5 cm and was previously dated to 55.4 ka in the NGRIP record based
on the GICC05 age model (Svensson et al., 2008) and 54.5 ka by Ar-Ar dating of the volcanic ash
(Southon, 2004), whereas in our age model it is at 57.4 ka. Outside of these two zones, glass shard
counts are also high at H3 and H6. We cannot rule out that the shards at Heinrich layer 3 were
delivered by a gravity flow.

The chronology of GVY001 is based on radiocarbon and tephrochronology. Radiocarbon analysis (n=7) in this core follows the same procedure as in JPC37, except for the youngest sample at the depth of 2 cm. At that depth we used *G. bulloides* instead of *N. pachy*. since not enough *N. pachy*. were found. Although age offsets have been identified in co-occurring planktic species (Ausín et al., 2019; Brocker et al., 1988; Costa et al., 2017), our age model is not likely to be sensitive to offsets of centuries or even millennia at this one horizon. Ash Zone 1 is identified at 58 cm, and we use 12.2 ka BP as its age (Andrews and Voelker, 2018).

4. Results

226 The age models indicate moderately high sedimentation rates in both cores: ~10 cm/kyr in JPC37 227 and ~12 cm/kyr in GVY001 (Figure 2 and Figure S1). Core JPC37 captured the whole last glacialinterglacial cycle (MIS 1-5) and the very end of MIS 6, whereas core GVY001 extends to just 228 229 beyond 30 ka. Crossed beddings of foraminifera-rich sands are found in JPC37 around depths 288 230 cm – near the depth where H3 is found – (Figure S2) and 1312 cm. These sands are likely caused by turbidites and could influence the interpretation of results from this particular interval. Data 231 232 from H3 in JPC37 is therefore considered tentatively except for when it is supported by 233 independent evidence from GVY001, which is from a different location not influenced by any 234 turbidite deposits. GVY001 has a similarly high sedimentation rate as JPC37.

235 In JPC37, % IRD varies between 0% and 100% (Figure 3e). At five depths, % IRD reached 100%. At three other depths, % IRD reached 50%. The majority of the % IRD peaks are found in the 236 237 upper (glacial) part of the core. Heinrich layers are identified by peaks in % IRD, as well as peaks 238 of Ca/Sr that are indicative of detrital carbonate (Hodell et al., 2008), magnetic susceptibility, density, and % coarse (Figure 3). Layers with prominent increases in all of these proxies are present 239 240 in MIS 2-4 and late MIS 6. H4 and H5 show the strongest signals in all of these proxies, consistent 241 with previous findings (Hemming et al., 2004). Several episodes of smaller magnitude can be 242 found during early MIS 5. The strongest peak of Ca/Sr appears at 1220 cm depth and is matched with a ²³⁰Th_{xs} low. However, no signals in magnetic susceptibility, % coarse, % IRD, % *N. pachy.*, 243 or δ^{18} O are present at this depth (density data does not extend to this depth). This discrepancy 244 245 suggests the delivery of fine detrital carbonates at a high rate without freshwater flux or lower SST. 246 A possible explanation is a gravity flow that originates from a region with high concentration of 247 detrital carbonates. In GVY001, % IRD ranges between 0% and 100% (Figure 3e). Compared to 248 JPC37, % IRD in GVY001 is higher on average. Ca/Sr, magnetic susceptibility, density, and % 249 coarse help identify HEs. Possibly due to GVY001's more poleward location and setting within 250 the cold Labrador current, % IRD and % N. pachy. are saturated to 100% during most of the last 251 glacial period.

In JPC37, % *N. pachy.* has a range between 1% and 99% (Figure 3f). Peaks of % *N. pachy.* can be found at the same or slightly above the depths of % IRD peaks. δ^{18} O of *N. pachy.* ranges between 1.2‰ and 4.1‰ and displays a hybrid signal that combines the typical sawtooth glacial-interglacial pattern and episodes of depletion corresponding to HEs. In each of the identified HEs, % *N. pachy.* increased while the δ^{18} O of *N. pachy.* decreased. During HEs, % *N. pachy.* continues to vary even when % IRD stabilizes at a high value. Our leaching experiment on JPC37 sediments suggests that the recoil-related losses of 234 U and ²³⁰Th are about 10% each on average in Heinrich layer detrital sediments and higher in between (Figure S3). The variations in 234 U/ 238 U are smaller than those of 230 Th/ 234 U. We use 0.9 for ²³⁰Th/ 234 U and 234 U/ 238 U and 0.81 for 230 Th/ 238 U for both cores. The leaching experiment also suggests a potentially broad range in the detrital 238 U/ 232 Th ratio (Figure S4), but the five leaching experiments made in Heinrich layer 4 have an average of 0.48, which we use in this study.

The ²³⁰Th_{xs} profile from JPC37 contains values that vary between 0-5 dpm/g. Low ²³⁰Th_{xs} values (<0.5 dpm/g) are observed at 16.3, 24.0, 38.1, 45.8, 60.6, 67.8 ka during the last glacial period. The low ²³⁰Th_{xs} episodes have ages one thousand years within the previously determined HE ages (Hemming, 2004), which is broadly within the uncertainty of radiocarbon dating. One sample in H4 and another in late MIS 6 have ²³⁰Th_{xs} so low that their 95% uncertainty range barely reach above 0. GVY001 resembles JPC37 in ²³⁰Th_{xs} for the most part, except the periods before and after H2.

271 In JPC37, mass flux has a range between 2 g/cm²kyr and 133 g/cm²kyr, although higher mass 272 fluxes are associated with higher uncertainties. There are nine peaks of mass flux. Outside of the 273 mass flux peaks, mass flux fluctuates around 4 g/cm²kyr. IRD flux varies between 1 grain/cm²kyr 274 and 560,000 grains/cm²kyr. At eight depth ranges, the values of IRD flux reach above 50,000 275 grains/cm²kyr. These depths coincide with HEs. The IRD fluxes during H5 and H11 are the highest 276 among HEs, with values at or above 500,000 grains/cm²kyr. The IRD fluxes during H2 and H4 277 are around 250,000 grains/cm²kyr. The IRD fluxes at H1, H6, and the event prior to H6 are lower at 100,000 grains/cm²kyr, 41,000 grains/cm²kyr, and 42,000 grains/cm²kyr, respectively. The IRD 278 fluxes during each HE and the event prior to H6 are statistically distinct from the ambient IRD 279 flux, even at the smallest event (H6), where the p-value of the "Student's" t-Test is 4×10^{-6} . In 280

281 GVY001, mass flux ranges from 3 g/cm²kyr to 40 g/cm²kyr. The peaks of mass flux at GVY001 are lower than the peaks of JPC37, while the mass flux between the peak values clusters around 7 282 283 g/cm²kyr. During H1-3, the mass fluxes are around 20-30 g/cm²kyr. However, unlike JPC37, 284 GVY001 shows a clearer signal of mass flux increase during Younger Dryas. The IRD flux from 285 GVY001 ranges between 0 grains/cm²kyr and 190,000 grains/cm²kyr. It is the highest during H1 286 and H2 at or above 160,000 grains/cm²kyr, followed by the Last Glacial Maximum at 120,000 grains/cm²kyr, and H3 at 90,000 grains/cm²kyr. A "Student's" t-Test shows that the H3 IRD flux 287 is distinctly different from the ambient sediment (p-value= 1×10^{-10}). 288

289

5. Discussion

290 5.1. Interpretation of IRD flux

HE sedimentary layers have been viewed as the possible result of intervals of decreased 291 foraminifera productivity, foraminifera dissolution, increased IRD deposition, or some 292 combination of those influences. Previously, the only ²³⁰Th_{xs}-based quantification of the 293 294 depositional flux of grains that are uniquely identified as IRD during the HEs is from V28-82 in 295 the eastern subpolar North Atlantic (McManus et al., 1998). The data from V28-82 showed that 296 H1, H2, H4, and H5 were at least in part increased ice-rafting events, but left open the question for the two atypical events. Our new results from JPC37 and GVY001 reaffirm McManus et al.'s 297 298 conclusion, showing increased IRD flux for the four most typical HEs (Figure 4 c). Additionally, 299 here we show for the first time that, at least in the western North Atlantic, the sediment layers associated with H3 and H6 were also the result of increased ice-rafted deposition, rather than solely 300 301 the result of reduced productivity near the sea surface or enhanced foraminifera dissolution on the seafloor. 302

303 IRD fluxes in V28-82 are approximately half of those in the IRD flux of the western cores during H1 and H2, and about half of the IRD flux of JPC37 during H4 and H5. Even larger differences 304 305 between IRD fluxes in the east and west are observed during H3, H6, and H11. The increases in 306 IRD fluxes in V28-82 are muted during these events, whereas the IRD flux from JPC37 during H3 is thirty times higher than V28-82, during H6 nine times higher, and during H11 twenty times 307 308 higher. Although H3 in JPC37 is near what may be a turbidite deposit, an increase in IRD flux 309 during this interval also occurs in core GVY001, which has no evidence of turbidite deposition. 310 At GVY001, the IRD flux during H3 is eight times higher than in V28-82. A comparison of IRD 311 concentration and mass flux, two variables used to calculate IRD flux, reveals that much of the 312 difference in IRD flux between V28-82 and the two western cores comes from the difference in 313 mass flux. This comparison confirms that while IRD was an important sedimentary component at 314 all three core locations, much more of it was deposited in the west than the east (Figure 4 a and b). 315 Since the eastern core displays much lower IRD flux during H3 and H6, a stronger zonal flux 316 gradient may have existed during these two periods – more icebergs melted in the western basin, with fewer icebergs reaching and depositing IRD in the eastern region. The distinct melting 317 318 patterns could be explained by any one, or a combination of, the following three factors: The 319 calving flux from the Laurentide may have been smaller in magnitude during H3 and H6 so the 320 majority of the drifting ice melted in the western NA without making it to the east; some different 321 ice sheet(s) may have contributed to or dominated calving during this interval; or the surface ocean 322 current patterns were different, causing changed iceberg trajectories.

323 If the calving of icebergs from the Laurentide ice sheet increased during H3 and H6, but to a lesser 324 extent than during the other events, the IRD fluxes associated with these two events would display 325 a visible but smaller increase than in the typical HEs across the subpolar Atlantic and a possible 326 gradient from west to east. At face value, our data are consistent with this hypothesis. H3 and H6 327 do display increases in IRD flux that are smaller in magnitude than the other events, and there is a 328 clear depositional gradient. Several hypotheses have been proposed to explain the smaller 329 magnitude of calving during these events. Gwiazda et al. (1996) speculated that the ice sheet 330 volume during H3 and H6 could have been smaller, leading to the smaller magnitude of calving. 331 The logic of Gwiazda et al. is based on the observation that H3 and H6 occurred at the onset of the 332 MIS 2 and 4, when the Laurentide ice sheet was just starting to regrow after periods of interglacial 333 warmth. A similar yet distinct hypothesis states that H4, which in our western cores appears to be 334 a particularly large event, "gutted" the Hudson Strait, removing it of all ice (Kirby and Andrews, 1999). As a result, H3 occurred when the Laurentide was still in a growth phase. The lack of a 335 336 signal of H3 and H6 in V28-82 to the east might also reflect the greater distance from the source 337 of icebergs, or the possibility that H3 and H6 icebergs were less dirty (Andrews, 2000; Kirby and 338 Andrews, 1999). However, we are not aware of any evidence to suggest that H3 and H6 icebergs 339 were particularly clean compared to icebergs in other HEs.

It has also been suggested that H3 and H6 may have had a European ice sheet origin (Grousset et al., 2000, 1993; Snoeckx et al., 1999). Our data do not support an entirely European origin for H3 and H6, since the deposition of IRD during each was demonstrably greater in the western North Atlantic. An iceberg trajectory model supports the unlikeliness of this scenario and suggests that the icebergs originating from the European ice sheet are mostly confined to the Norwegian Sea (Death et al., 2006). However, we cannot rule out that associated or precursor events of a European origin took place.

Our data do not directly support the hypothesis that warmer SST or a reorganization of thecirculation pattern led to more iceberg melting in the west during H3 and H6. Although IRD fluxes

increased during H3 and H6 at our western sites, the fluxes are not higher than during the other
HEs. Combined with the essential lack of increased IRD deposition at the eastern site, the total
IRD flux during H3 and H6 appears to be lower compared to the other events.

352 IRD flux reconstruction provides a potentially important constraint on modeled iceberg discharge 353 during HEs as a freshwater delivery mechanism (Death et al., 2006). While most HE modeling 354 studies have focused on meltwater (Ganopolski and Rahmstorf, 2001; Prange et al., 2004; Roberts 355 et al., 2014) and iceberg calving (Alvarez-Solas et al., 2013, 2010; Bassis et al., 2017; Marshall 356 and Koutnik, 2006), the absolute fluxes of meltwater and iceberg calving are challenging to 357 reconstruct with paleo proxies. Meltwater proxies, including δ^{18} O (Cortijo et al., 1997; Roche et al., 2004), %C37:4 (Naafs et al., 2011; Rodrigues et al., 2017; Stein et al., 2009), and ¹⁰Be/⁹Be 358 359 (Valletta et al., 2018), while valuable, are currently only qualitative. IRD grain concentration can 360 only provide a relative measure of the magnitude of iceberg calving (e.g., Bond et al., 1992). In 361 contrast, IRD flux is a proxy that is both quantitative and accessible, with the potential for 362 comparable model output. A better understanding of the IRD entrainment and delivery 363 mechanisms could be developed by incorporating IRD fluxes into ice sheet calving simulations. 364 The direct data-model comparison of IRD flux can potentially help assess ice sheet calving 365 simulations and improve the understanding of calving behaviors.

366 5.2.Sea surface proxies

The relative abundance of polar foraminifera (% *N. pachy.*) at GVY001 is saturated at 100% for most of the last glacial period, possibly due to the site's more poleward location (Figure 3). The % *N. pachy.* data from JPC37 exhibit more variability, with increases in every HE layer, indicating repeated sea-surface cooling. During H1, H2, and H4, % *N. pachy.* seems to show a double peak structure. The double peak might suggest that the SST is lowest at the beginning and end of these HEs but briefly returned to warmer temperatures midway through the events. We do not see a
similar sequence of delayed onset in the deposition of IRD after the initial increase in *N. pachy*.
abundance, as reported previously (Barker et al., 2015).

375 The δ^{18} O of *N. pachy*. from JPC37 displays both the typical sawtooth glacial-interglacial cycle and episodes of depletion associated with HEs (Figure 3h). Given the % N. pachy. increase and 376 377 therefore implied SST decrease during HEs, the depletion of δ^{18} O cannot be the result of 378 temperature changes. The likely explanation is that the melting icebergs released a large amount 379 of δ^{18} O-depleted freshwater during each HE. A depletion of planktic δ^{18} O is not obvious in H6 380 and HQ (discussed in the next section), but neither did it become enriched during these periods. In light of the concurrent % N. pachy. increases, we suggest that just as during the other HEs, an 381 382 influx of freshwater reached this site during H6 and HQ. The depletion of N. pachy. δ^{18} O is 383 consistent with previous studies on North Atlantic cores (Bond et al., 1992; Cortijo et al., 1997; Hillaire-Marcel et al., 1994; Labeyrie et al., 1999). Cortijo et al. (1997) mapped out changes in N. 384 pachy. δ^{18} O across the North Atlantic during H4. The magnitude of changes observed in JPC37 385 386 during H4 (~ 1 ‰) is comparable to a nearby core from that study (1.1 ‰ at SU90-11).

A cross-correlation of *N. pachy*. δ^{18} O and ²³⁰Th_{xs} from JPC37 during the period between H1-6 shows little lag between the two variables (Figure S5). This simultaneity is robust on a range of interpolation intervals and does not change when the starting and ending times are varied by a few thousand years. The deposition of debris and local freshening of the sea surface were thus contemporaneous events on the time scale resolvable by our record. That is consistent with both phenomena being the direct consequence of melting icebergs, although we cannot rule out the accompanying presence of additional meltwater.

Although H1-6 are the most commonly known and studied HEs of the last glacial period (e.g., Bond et al., 1992; Broecker et al., 1992; Grousset et al., 2000; Marcott et al., 2011), Heinrich (1988) postulated the existence of additional events in the early glacial and late interglacial period. A recent ice sheet modeling study (Bassis et al., 2017) predicted a previously unidentified HE that they named HQ at ~65 ka. Our IRD flux from JPC37 provides the strongest evidence yet of the existence of HQ.

401 The basis of our age model is partially the alignment of millennial-scale cooling events (C16-C24) 402 identified throughout the North Atlantic region (McManus et al., 1994), including on the Iberian 403 margin (Martrat et al., 2007; Greenland Stadial (GS) 17-25). Bassis et al. predicted the occurrence 404 of HQ during GS19, or C18 in the framework of the cooling events. There are two distinct IRD-405 flux events within C16 and C18. If the younger event is H6, consistent with an age of ~60 ka 406 (Bassis et al., 2017; Bond et al., 1999; Hemming 2004; Martrat et al., 2007), it thus seems likely 407 that the older event of the two is the hypothesized event HQ. Given that Bassis et al. identified 408 H7b prior to HQ within C19, which we identified with our % N. pachy. record, it is unlikely that 409 we mistake HQ for an earlier HE. These two IRD flux events also cannot be H5a (Rashid et al., 410 2003) since we found Ash Zone 2 at a shallower depth in the core. Ash Zone 2 is between H5a and 411 H6, which gives us confidence with the designation of H6. The fact that HQ is potentially found 412 in JPC37 but was absent in previous studies suggests that the influence of the event may be 413 regionally limited, but it also raises the possibility that HQ and H6 may be mistaken for one another 414 in previous studies. Therefore, we recommend that ocean sediment studies that seek to identify H6 or HQ should have an accompanying SST proxy to resolve cooling events C16 and C18. 415

416 During the late last interglacial period at 70, 78, 87, 105, 109, and 117 ka, we found six additional 417 IRD flux increases in JPC37 (Figure S6). These IRD flux increases were two orders of magnitude 418 smaller than HEs, and unlike the typical HEs, changes in Ca/Sr, magnetic susceptibility, density, % 419 coarse, % IRD, and % N. pachy., do not accompany the IRD flux increases or have a temporal 420 offset with the IRD flux increases. Another difference, relative to the HEs, is that most of these 421 interglacial IRD flux increases were preceded by discernable mass flux increases (Figure S7). This 422 repeated sequence of events may give us a clue as to their origins. According to the turbidite-IRD 423 sequence previously proposed to explain IRD layers in the Labrador Sea (Rashid et al., 2012), the 424 early high mass flux could be caused by the initial meltwater discharge and the ensuing turbidity 425 and nepheloid flow. This mechanism would increase mass flux without bringing in IRD. Following 426 the turbidite facies, IRD would have been deposited as icebergs were discharged. The number of 427 these locally high IRD flux events we identified corresponds to the number of cooling events 428 during the last interglacial period proposed by McManus et al. (1994), which we tentatively marked in Figure S6 and S7. 429

430 The lack of associated signals in other typical HE proxies (Figure S6 and S7) during most of these cooling events raises questions about their nature. Among other explanations, the magnitude of 431 432 the events could have been too small to have detectable changes in typical HE proxies. 433 Alternatively, the source region of the delivered materials may have changed, which could have 434 led to the muted responses in Ca/Sr and magnetic susceptibility. These events could also have been 435 the result of meltwater outbursts, similar to the 8.2 ka event (Alley et al., 1997; Ellison et al., 2006; 436 Keigwin et al., 2005). A fourth potential explanation is that they were triggered by deep turbidity 437 currents, as suggested by Hillaire-Marcel et al. (1994). Given that most of the events were

438	associated with increases in % N. pachy., it is unlikely that they were caused by the deep turbidity
439	currents alone, although a combination of the above mechanisms is still possible.
440	6. Conclusions
441	(1) The IRD flux in the western North Atlantic cores JPC37 and GVY001 increased during
442	each HE during the last glacial cycle. A single mechanism may therefore account for all
443	HE during this period.
444	(2) Compared to the only other available 230 Th _{xs} -based IRD flux record, which is in the eastern
445	North Atlantic, the western sites experienced much higher IRD flux during all HEs, notably
446	including H3 and H6. We suggest that these two events, in the western North Atlantic at
447	least, were the result of increased ice calving, rather than solely the result of other
448	mechanisms such as increased foraminifera dissolution or reduced productivity.
449	(3) IRD fluxes during H3 and H6 in the western North Atlantic are smaller than the other
450	typical HEs. This result is most consistent with the hypothesis that the calving of icebergs
451	from the LIS increased during H3 and H6 but to a lesser extent than during the other events.
452	(4) All HEs were accompanied by surface cooling and freshening in the western subpolar
453	North Atlantic.
454	(5) A series of previously identified cooling events during the MIS 5 interglacial were found
455	in JPC37, accompanied by evidence for increased ice rafting that was two orders of
456	magnitude smaller than HEs. These findings confirm the regional extent and limited impact

457 of these stadial events.

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465 Appendix A. Supplementary material

All data are archived in PANGAEA and will be made public upon publication. Reviewers can
temporarily access them with this link:
<u>https://www.pangaea.de/tok/f2915574eee9270de2552d634d19422d0680d4da</u>



Figure 1. North Atlantic map with core locations. Stars depict the locations of cores used by this study:

471 JPC37 (43°58'N, 46°25'W, 3981 m); GVY001 (50°09'36''N, 45°30'36''W, 3721 m). The white dot is

472 the core used for comparison of IRD flux: V28-82 ($49^{\circ}27$ 'N, $22^{\circ}16$ 'W, 3935 m) (McManus et al., 1998).

473 The frosted area represents ice sheet extent during the Last Glacial Maximum (Ehlers et al., 2011). Red

and blue arrows are the warm and cold surface circulation, respectively, after (Hemming et al., 2002).
Aqua arrow leaving the Hudson Strait represents the calving of icebergs from the LIS. Contours delineate

475 Aqua arrow leaving the Hudson Strait represents the calving of icebergs from the LIS. Contours delineate476 the Ruddiman IRD belt (Ruddiman, 1977). Basemap from NASA Blue Marble June image (Stockli et al.,

477 2005).



Figure 2. Chronology of JPC37. Variations in *N. pachy*. relative abundance are correlated with an

480 alkenone unsaturation SST record (Martrat et al., 2007; Iberian margin stadials marked in black numbers)

(a). That SST record was previously tied to the North Greenland Ice Core Project (NGRIP) chronology.

Blue numbers denote cooling events (McManus et al., 2002, 1994, see Figure S8 for details). Tie points to

483 our core are marked by thin gray lines. The lower panel contains the compilation of all age control points,

484 including radiocarbon dating, tephrochronology, and tuning with alkenone record (b).



Figure 3. GVY001 (cyan) and JPC37 (other colors) Ca/Sr with JPC37 age control points marked by red (ash zone), brown (radiocarbon), and blue (SST tie points) triangles (a), magnetic susceptibility with GVY001 age control point marked by red (ash zone) and brown (radiocarbon) triangles (b), bulk wet sediment density (c), coarse (>63µm) fraction (d), IRD abundance (e), *N. pachy.* abundance, of which the early last interglacial (95-125 ka) data are partially from McManus et al. (2002) (f), ²³⁰Th_{xs} with shading marking 2 σ uncertainty (g), and *N. pachy.* δ^{18} O (h). Gray bars are Younger Dryas (YD), H1-6, H11, as well as HQ, as predicted by Bassis et al. (2017).



495 concentration (b), and IRD flux (c). Triangles in (a) are mass flux data points too high to quantify. The bar plot inset in (c) compares maximum IRD flux during each HE relatively. Although the Y-axis for the inset is omitted for simplicity, the bar heights are drawn proportionally.

Study	(²³⁸ U/ ²³² Th) _{detrital} used
Veiga-Pires and Hillaire-Marcel, 1999	0.58
Thomson et al., 1995	0.67
Thomson et al., 1999	0.7
McManus et al., 2004	0.57
Henderson and Anderson, 2003	0.6
Böhm et al., 2015	0.47
Bourne et al., 2012	0.55
Gherardi et al., 2009	0.6
Lippold et al., 2009	0.5
Lippold et al., 2011	0.5

	1
Lippold et al., 2016	0.6
Guihou et al., 2010	0.5
Guihou et al., 2011	0.5
Roberts et al., 2014	0.6

500 *Table 1.* Detrital U/Th used by previous studies. The range of the values is 0.47 - 0.7.

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786	Supporting Information for
787	Enhanced iceberg discharge in the western North Atlantic during all Heinrich
788	events of the last glaciation
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806 Introduction

807	The text presented here discusses the U-Th systematics, and the figures further characterize the
808	physical and chemical properties of the two cores used by this study.
809	In Text S1, we discuss the choices made in calculating 230 Th _{xs} , including the values of
810	$(^{230}\text{Th}/^{238}\text{U})_{\text{detrital}}$ and $(^{238}\text{U}/^{232}\text{Th})_{\text{detrital}}$.
811	In Figure S1, we show the age model of GVY001.
812	In Figure S2, we show the core image of the turbidite sequence found at 280-290 cm depth
813	in JPC37.
814	In Figure S3, we show the leaching experiment results from JPC37 for determining the
815	detrital ²³⁰ Th/ ²³⁸ U ratio.
816	In Figure S4, we show the leaching experiment results from JPC37 for determining the
817	detrital U/Th ratio.
818	In Figure S5, we show the phasing lag correlations of planktic $\delta^{18}O$ with $^{230}Th_{xs}$ from
819	JPC37.
820	In Figure S6, the same as Figure 3, except we zoom in on the cooling events and the panel
821	that displayed ²³⁰ Th _{xs} shows IRD flux instead.
822	In Figure S7, the same as Figure 4, except we zoom in on the cooling events.
823	In Figure S8, we show the last interglacial % N. pachy. from JPC37 and V29-191 side by
824	side (McManus et al., 1994).
825	
826	

Text S1. The calculation of 230 Th_{xs} commonly assumes that the detrital decay chain of 238 U is in 827 secular equilibrium ($^{238}U_{det} = ^{230}Th_{det}$). This assumption is valid when an isotopic system is closed 828 829 for a sufficient length of time. However, α recoil disrupts the closed system by potentially ejecting the decay product (²³⁴U or ²³⁰Th in this case) or leaving it within a damaged crystal lattice site. 830 Recently, it has been suggested that 4% of the decayed 234 U is lost due to α recoil, and the decay 831 of ²³⁴U further ejects 4% of ²³⁰Th (Bourne et al., 2012). These authors reasoned that the average 832 value between 0.92 (1×0.96^2) and 1 (secular equilibrium), 0.96, should be used as the value of 833 (²³⁰Th/²³⁸U)_{detrital}. Deep-sea sediments have also been observed within this range and somewhat 834 lower (DePaolo et al., 2006). Our leaching experiment on sediments suggests that the effect of α 835 recoil may be stronger than previously thought (Figure S3). The losses of ²³⁴U and ²³⁰Th are 10% 836 each on average, although the uncertainty on ²³⁰Th loss is greater than on ²³⁴U. These data imply 837 that, when calculating ²³⁰Th_{xs}, the detrital correction is smaller, leading to more ²³⁰Th counted 838 839 towards the scavenged portion. The result is a higher 230 Th_{xs}. This correction, compared to the equilibrium assumption, increases 230 Th_{xs} during non-HE period by ~5%. Because the higher burial 840 841 fluxes during HEs often result in very low ²³⁰Th_{xs}, the proportional change associated with this 842 correction is even larger (up to 70%), although the absolute change is small.

A wide range of $(^{238}U/^{232}Th)_{detrital}$ has been used in the $^{230}Th_{xs}$ calculation in the North Atlantic, ranging from 0.47 to 0.7 (Table 1). These studies use either the $^{238}U/^{232}Th$ minimum measured, which is thought to reflect the minimal influence of authigenic ^{238}U , or a vaguely defined basinwide value. More recently, it has been suggested that $(^{238}U/^{232}Th)_{detrital}$ may vary through time (Missiaen et al., 2018). Consistent with that study's conclusion, our leaching experiment from JPC37 (Figure S4) shows large variations in the detrital ratio. Since HE mass flux is the focus of this study, we choose to use the $(^{238}U/^{232}Th)_{detrital}$ that produces a conservative yet realistic estimate of mass flux throughout. The five leaching experiments made in Heinrich layer 4 provide a mean ratio of 0.48, which we use in this study. While this is toward the low end of the range of values previously applied for North Atlantic sediments, it is indeed within that range. Using higher ratios would yield negative 230 Th_{xs} in some cases, implying net sedimentary loss of 230 Th from settling particles to the water column, which we consider unlikely. The higher detrital ratios in other Heinrich layers could be due to the lower sampling resolution, which may not capture the lowest detrital ratios.

The relatively conservative but nonetheless very low ²³⁰Th_{xs} results we obtained are useful to 857 858 inform us that the subpolar western North Atlantic had episodically high fluxes of ice rafting. However, we are hindered in using 230 Th_{xs} to normalize the burial of 231 Pa_{xs}. In other words, we 859 860 cannot use Pa/Th as a tracer for circulation strength at this site. The low concentrations are 861 overwhelmed by the magnitude of the combined, propagated uncertainties for Pa and Th, rendering 862 our results so ambiguous that they are uninterpretable. Future studies aiming to use this approach to reconstruct rates of deep ocean circulation associated with iceberg discharges from the 863 Laurentide should focus on sites further east or south to avoid being similarly overwhelmed by the 864 increased IRD flux. 865









Figure S2. Turbidite sequence found at 280-290 cm depth in JPC37, around the depth of Heinrich layer 3. The distance between the two core depth tags is 10 cm.



Age (ka BP)
 Figure S3. Leaching experiment results from JPC37 for determining the detrital ²³⁰Th/²³⁸U ratio. Dashed lines are the ratios used by this study (0.9 for ²³⁰Th/²³⁴U and ²³⁴U/²³⁸U and 0.81 for ²³⁰Th/²³⁸U).



877 878

Figure S4. Leaching experiment results from JPC37 for determining the detrital U/Th ratio. The dashed line is the average ratio (0.48) of the high-resolution measurements within H4, which we use in 230 Th_{xs} normalization calculations in this study.





884 **Figure S5**. Phasing lag correlations of planktic δ^{18} O with ²³⁰Th_{xs} from JPC37, bootstrapped 1000 times allowing sampling start time, end time, and time step to vary. The correlation coefficient in the positive

(negative) direction is calculated when planktic δ^{18} O leads (lags) 230 Th_{xs}.



- Figure S6. Same as Figure 3 except (g) displays IRD flux, limited to the late last interglacial period for
- the locally-high IRD flux events from JPC37, and assigned cooling events numbering according to
- McManus et al. (1994) and McManus et al. (2002). Shadings are HQ and cooling events.



893

Figure S7. Same as Figure 4 but for the locally-high IRD flux events of the late last interglacial period from JPC37 and assigned cooling events tentatively according to McManus et al. (1994). Shadings are HQ and cooling events.



- Figure S8. Correlation of last interglacial % *N. pachy.* between JPC37 and V29-191 (McManus et al.,
 1994).
- 900
- 901 **References**
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