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We present an original manuscript entitled “Lithospheric structure and melting processes in southeast Australia: new constraints from joint probabilistic inversions of 3D magnetotelluric and seismic data” by M.C. Manassero¹, S. Özaydin^{1,2}, J. C. Afonso^{3,4,1}, J. Shea^{1,5}, A. Kirkby^{6,7}, I. Ezad¹, S. Thiel^{8,9}, I. Fomin¹ and K. Czarnota⁷.

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Yours Sincerely,

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1 **Lithospheric structure and melting processes in**
2 **southeast Australia: new constraints from joint**
3 **probabilistic inversions of 3D magnetotelluric and**
4 **seismic data**

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18 **Key Points:**

- 19 • We apply a novel approach for joint probabilistic inversions of 3D magnetotelluric
20 and seismic data.
- 21 • We use the new method to image the lithosphere-asthenosphere system beneath
22 southeastern of Australia.
- 23 • The imaged lithosphere correlates with the location of volcanic centers and pro-
24 vides insights on the melt production in the region.

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25 Abstract

26 The thermochemical structure of the lithosphere exerts control on melting mech-
 27 anisms in the mantle as well as the location of volcanism and ore deposits. Imaging the
 28 complex interactions between the lithosphere and asthenospheric mantle requires the joint
 29 inversion of multiple data sets and their uncertainties. In particular, the combination
 30 of seismic velocity and electrical conductivity with data proxies for bulk composition and
 31 elusive minor phases is a crucial step towards fully understanding large-scale lithospheric
 32 structure and melting. We apply a novel probabilistic approach for joint inversions of
 33 3D magnetotelluric and seismic data to image the lithosphere beneath southeast Aus-
 34 tralia. Results show a highly heterogeneous lithospheric structure with deep conductiv-
 35 ity anomalies that correlate with the location of Cenozoic volcanism. In regions where
 36 the conductivities have been at odds with sub-lithospheric temperatures and seismic ve-
 37 locities, we observe that the joint inversion provides conductivity values consistent with
 38 other observations. The results reveal a strong relationship between metasomatized re-
 39 gions in the mantle and i) the limits of geological provinces in the crust, which elucidates
 40 the subduction-accretion process in the region; ii) distribution of leucitite and basaltic
 41 magmatism; iii) independent geochemical data, and iv) a series of lithospheric steps which
 42 constitute areas prone to generating small-scale instabilities in the asthenosphere. This
 43 scenario suggests that shear-driven upwelling and edge-driven convection are the dom-
 44 inant melting mechanisms in eastern Australia rather than mantle plume activity, as con-
 45 ventionally conceived. Our study offers an integrated lithospheric model for southeast-
 46 ern Australia and provides insights into the feedback mechanism driving surface processes.

47 Plain Language Summary

48 The lithosphere is the outermost rigid layer of the Earth and the focus of impor-
 49 tant geological processes such as earthquakes (seismic activity), volcanism, and miner-
 50 alization. The location of these processes often coincide with deep discontinuities in the
 51 lithospheric structure. Imaging the structure of the lithosphere using geophysical tech-
 52 niques is then crucial to fully understand the nature of these processes. Obtaining the
 53 most reliable images of the lithospheric structure requires the joint analysis of two or more
 54 geophysical data sets. In particular, the combination of magnetotellurics (an electromag-
 55 netic technique) and seismic data holds great potential due to their complementary sen-
 56 sitivity to the Earth's properties. Combining a joint analysis with a probabilistic approach
 57 help us understand the variability of the lithospheric structure better since they provide
 58 a large number of models that can explain the data. Given the good data coverage in
 59 southeast Australia, we use a new probabilistic approach for the joint analysis of mag-
 60 netotelluric and seismic data to image the lithosphere structure beneath this region. Our
 61 results show a complex lithospheric structure in line with the location of volcanism and
 62 tectonic history of the region. Lithospheric composition derived from the models pro-
 63 vides significant insights into melt production in the area.

64 1 Introduction

65 The magnetotelluric method (MT) has great potential for investigating metasoma-
 66 tism and tectono-magmatic processes in the lithosphere (e.g., Wannamaker et al., 2008;
 67 Comeau et al., 2015; Aivazpourporgou et al., 2015; Wannamaker et al., 2014; Bedrosian,
 68 2016; Kirkby et al., 2020; Blatter et al., 2022; Özaydın & Selway, 2022; Cordell et al.,
 69 2022). Due to its sensitivity to fluid and/or melt content, MT is particularly useful for
 70 probing the connection between deep melt/fluid pathways and their surface expressions,
 71 such as the location of ore deposits (e.g., Griffin et al., 2013; Heinson et al., 2018; Kirkby
 72 et al., 2022) and volcanic centers (Wei et al., 2001; Comeau et al., 2015). However, MT
 73 is not free of limitations. For instance, MT struggles to delineate deep conductivity struc-
 74 tures, especially when they are below shallow conductive features. This is due to the dif-

75 fusive behaviour of electromagnetic waves and the high sensitivity of MT to conductors
76 (Jones, 1999). The MT method is also ambiguous in discerning the different factors that
77 affect electrical conductivity, such as temperature, water/melt content and composition.
78 Unlocking the full potential of the MT method requires the development of methodolo-
79 gies that can assign meaningful physical interpretations to conductivity anomalies and
80 discriminate between their causes (Selway, 2014).

81 A widely adopted approach to reduce feature ambiguity is the combination of MT
82 with other geophysical data sets via joint inversions (e.g. Khan et al., 2006; Gallardo &
83 Meju, 2007; Jegen et al., 2009; Moorkamp et al., 2010; Bennington et al., 2015; Afonso,
84 Rawlinson, et al., 2016; Jones et al., 2017; Blatter et al., 2019; Manassero et al., 2021;
85 Liao et al., 2022; Wu et al., 2022). By exploiting the complementary sensitivities of dif-
86 ferent data sets to the properties of interest, joint inversions minimize the range of ac-
87 ceptable models consistent with the available data and can increase model resolution (e.g.
88 Moorkamp et al., 2007; Afonso et al., 2013a; Afonso, Rawlinson, et al., 2016; Afonso, Moorkamp,
89 & Fulla, 2016). For example, in the case of MT and seismic data, both data sets are
90 sensitive (to different degrees) to the background thermal and compositional structure
91 of the lithosphere. However, only MT is strongly sensitive to minor conductive phases
92 (e.g., hydrous minerals and graphite), hydrogen content or small-scale melt/fluid path-
93 ways (Karato, 1990, 2006; Evans, 2012; Yoshino, 2010; Khan, 2016; Selway, 2014; Man-
94 assero et al., 2021). In this way, joint MT+seismic inversions hold great potential for im-
95 proving the resolution of conductivity structures (e.g., Moorkamp et al., 2007, 2010; Gal-
96 lardo & Meju, 2007), detecting regions of partial melting and fluid pathways in the litho-
97 sphere (cf., Selway & O'Donnell, 2019; Evans et al., 2019; García-Yeguas et al., 2017; Ben-
98 nington et al., 2015), as well as understanding their relationship with the location of ore
99 deposits (e.g., Takam Takougang et al., 2015) and metasomatized lithologies (e.g., Sny-
100 der et al., 2014).

101 In addition to the benefits of joint inversions, valuable information about model
102 uncertainties can be obtained via simulation-based probabilistic approaches (Tarantola,
103 2005; Rosas-Carbalajal et al., 2013; Afonso, Rawlinson, et al., 2016; Manassero et al., 2021).
104 Rather than outputting a single best-fitting model, probabilistic approaches provide a
105 distribution of models and their associated probabilities according to their performance
106 in explaining the observations. Thus, probabilistic inversions naturally address the non-
107 uniqueness problem in geophysics (particularly in MT) and quantify model ambiguity
108 (Tarantola, 2005; Gregory, 2005). However, such probabilistic methods require the eval-
109 uation of millions of possible models, each in turn requiring the computation of a for-
110 ward solution. Consequently, simulation-based probabilistic approaches are limited to
111 problems where fast forward operators are available. In the case of 3D MT inversions,
112 fully probabilistic methods have been infeasible due to the large CPU time required by
113 the associated forward problem (Miensopust et al., 2013). In order to address this lim-
114 itation, Manassero et al. (2020) developed a novel strategy based on reduced order mod-
115 elling (referred to as RB+MCMC) that allows obtaining fast and accurate approxima-
116 tions of the forward solution and performing joint probabilistic inversions of 3D MT data
117 with other data sets (Manassero et al., 2021). Potential applications and the efficiency
118 of the method to solve the joint inverse problem of MT and seismic data were demon-
119 strated with whole-lithosphere synthetic examples in our previous paper (Manassero et
120 al., 2021).

121 In this work, we apply the method of Manassero et al. (2020) and Manassero et
122 al. (2021) to a dense array of collocated MT and seismic data in southeastern Australia.
123 This region is known to have experienced multiple orogenic events that resulted in a com-
124 plex crustal architecture, but its deep lithospheric roots remain poorly characterized (Rawl-
125 inson et al., 2016). It also hosts one of the most voluminous intraplate volcanic provinces
126 on Earth, the Eastern Australian Volcanic Province (Johnson et al., 1989). While half
127 of this volcanism can be linked to a hot mantle plume (e.g., Sutherland et al., 2012; Davies

128 et al., 2015), the melting mechanism responsible for the other half is far less clear (Well-
 129 man & McDougall, 1974; Shea et al., 2022). Here we focus on the mapping of whole-lithosphere
 130 3D structures and sub-lithospheric temperature anomalies in order to investigate i) the
 131 origin of the intraplate magmatism with no clear plume signatures, ii) the connection
 132 between deep melt/fluid pathways and the location of volcanic centers, and iii) how litho-
 133 spheric structure may have influenced melting generation and transport. Southeastern
 134 Australia is also a region with abundant xenolith-derived datasets (cf. Shea et al., 2022,
 135 and references therein) that can be used to validate the results from our joint inversion.

136 2 Geological background

137 The Tasmanides in southeastern Australia are a complex orogenic system that de-
 138 veloped from west to east through repetitive cycles of subduction, accretion and litho-
 139 spheric deformation along the eastern margin of Gondwana (Glen, 2005, 2013; Cham-
 140 pion et al., 2016; Rosenbaum, 2018; Moresi et al., 2014). This region is broadly divided
 141 into the Delamerian Orogen in the west (early-Palaeozoic) and the younger (mid-Paleozoic)
 142 Lachlan Orogen in the southeast (Figure 1.a). Much of the geological complexity in the
 143 area can be explained by a geodynamic model of a micro-continent collision and later
 144 development of an orocline, referred to as the Lachlan Orocline model (Cayley, 2011; Cay-
 145 ley & Musgrave, 2015; Moresi et al., 2014; Musgrave, 2015). The major structures de-
 146 scribed in this model are curved crustal geometries with an eastward rotation (Musgrave,
 147 2015), which have been imaged by gravity, magnetic and potential field data (e.g., Mus-
 148 grave & Rawlinson, 2010; Nakamura & Milligan, 2015; Nakamura, 2016, see Figure 1.c-
 149 d.); ambient noise tomography in the crust (e.g., Young et al., 2013; Pilia et al., 2015,
 150 see Figure 7.d) and MT conductivity models (Aivazpourporgou et al., 2015; Kirkby et
 151 al., 2020; Heinson et al., 2021). Important first-order information about the lithospheric
 152 structure beneath southeast Australia has been obtained from conventional studies, such
 153 as ambient noise and teleseismic tomography (Rawlinson et al., 2016; Davies et al., 2015;
 154 Young et al., 2013), xenolith thermobarometry (e.g., Lu et al., 2018), thermal modeling
 155 (e.g., Tesauro et al., 2020) and a recent 3D conductivity model (Kirkby et al., 2020). This
 156 latter 3D conductivity model showed, for the first time, that some of the crustal struc-
 157 tures associated with the orocline persist below the Moho, providing new insights about
 158 the lithospheric architecture and geodynamic history of the region.

159 Throughout the late Mesozoic and the entire Cenozoic, eastern Australia was sub-
 160 ject to voluminous mafic intraplate volcanism, which formed the extensive Eastern Aus-
 161 tralian Volcanic Province (EAVP, Johnson et al., 1989; Sutherland et al., 2012; Shea et
 162 al., 2022). Several regions across eastern Australia contain recent eruptions; in north-
 163 eastern Australia, the Kinrara vent contains lavas $\sim 7 \text{ ka} \pm 2 \text{ ka}$ (Cohen et al., 2017), while
 164 the Mount Gambier, Newer Volcanics (NV) in southeastern Australia contains lavas \sim
 165 $4\text{--}7.5 \text{ ka}$ (Blackburn et al., 1982; Smith & Prescott, 1987).

166 The EAVP comprises 67 separate volcanic centers with two dominant volcanic cen-
 167 ter compositions: basalt and potassic leucitite (Figure 1.b). While basaltic volcanics erupted
 168 through thinner lithosphere ($< 110 \text{ km}$) along the eastern and south-eastern seaboard,
 169 the leucitite volcanic centers lie on thick lithosphere ($> 125 \text{ km}$) in central New South
 170 Wales and central Victoria (Davies & Rawlinson, 2014; Rawlinson et al., 2017). This leuci-
 171 tite suite represents the most atypical and extraordinarily enriched melt compositions
 172 reported for mafic melts in eastern Australia (Cundari, 1973; Birch, 1978). Particularly,
 173 they represent melts from pervasively metasomatized source assemblages, likely a Ti-bearing
 174 oxide phlogopite websterite \pm apatite (see the recent review by Shea et al., 2022). The
 175 lack of anhydrous peridotite and the abundance of hydrous minerals in their mantle source
 176 assemblages is of particular importance to this work, indicating widespread mantle meta-
 177 somatism beneath eastern Australia.

178 The EAVP is also unique in the sense that about half of the volcanism is age-progressive
 179 and commonly linked to a hot mantle plume (e.g., Sutherland et al., 2012; Davies et al.,
 180 2015), whereas the remaining volcanic centres show no age-progression and no obvious
 181 melting mechanism (Wellman & McDougall, 1974). To further exacerbate this issue, lava
 182 compositions throughout the EAVP (including both age-progressive and non-age-progressive
 183 volcanic centers) argue for low-temperature melting of metasomatized mantle source litholo-
 184 gies. In contradiction to the presence of a hot mantle plume, these compositions suggest
 185 mild, but localized perturbations in mantle temperatures (Shea et al., 2022).

186 3 Methods and data sets

187 3.1 Data

188 The input data used in our joint probabilistic inversion include magnetotelluric (MT)
 189 data from the AusLAMP array (Australian Lithospheric Architecture Magnetotelluric
 190 Project) in southeast Australia and P-wave velocities from the tomography model of Rawl-
 191 inson et al. (2016). The long-period MT data were acquired at 298 AusLAMP stations
 192 (blue triangles in Figure 1.f) across a ~ 55 km spaced array covering an area of $950 \times$
 193 950 km. Details about the data acquisition and processing are given in Kirkby et al. (2020).
 194 The MT data are the full impedance tensor for periods between 6.4 to 40,000s. Error
 195 floors are set to 5% of $\max(|Z_{xx}|, |Z_{xy}|)$ for the components Z_{xx} and Z_{xy} and 5% of $\max(|Z_{yy}|, |Z_{yx}|)$
 196 for the components Z_{yy} and Z_{yx} . We assume uncorrelated data errors that follow a dou-
 197 ble exponential distribution (e.g., Farquharson & Oldenburg (1998); Rosas-Carbajal et
 198 al. (2013); Manassero et al. (2021)).

199 The P-wave velocity model used in this study (Rawlinson et al., 2016) was constructed
 200 from teleseismic tomography using data from the mainland component of the WOMBAT
 201 transportable seismic array (Rawlinson et al., 2015). In order to account for the unre-
 202 solved crustal component of the teleseismic arrival time residuals, the model of Rawl-
 203 inson et al. (2016) includes a detailed crustal model from ambient noise tomography (Young
 204 et al., 2013) and the Moho from AuSREM (Kennett & Salmon, 2012) as starting model
 205 of the tomographic inversion. Despite this additional constraint, uncertainties in abso-
 206 lute velocities within the crust remain relatively high, especially in the lower-crust (Young
 207 et al., 2013; Rawlinson et al., 2016). We obtained seismic velocities by interpolating the
 208 velocity model on a data-point grid of 50×50 km at the surface (shown in red dots in
 209 Figure 1f) and 24 points between the surface and 340 km depth. The velocity data er-
 210 rors are assumed to be uncorrelated and normally distributed with a standard deviation
 211 of 1%, according to Burdick & Lekić (2017). Examples of data and data fits for MT data
 212 and seismic velocities are shown in Figures 2. Additional figures can be found in the Sup-
 213 plementary Material.

214 3.2 Bayesian inversion

215 In the Bayesian or probabilistic approach to the inverse problem, inference about
 216 the model parameters \mathbf{m} given observed data \mathbf{d} is based on the so-called posterior prob-
 217 ability density function (PDF):

$$P(\mathbf{m}|\mathbf{d}) = \frac{P(\mathbf{d}|\mathbf{m})P(\mathbf{m})}{P(\mathbf{d})} \propto \mathcal{L}(\mathbf{m})P(\mathbf{m}) \propto \exp(\phi)P(\mathbf{m}), \quad (1)$$

where $P(\mathbf{m})$ denotes the prior PDF describing all the information on the model's pa-
 rameters prior to the inversion (e.g., prior geological or petrological knowledge in the area
 of study). $\mathcal{L}(\mathbf{m})$ is the likelihood function, which is specified by the statistical distribu-
 tion of the data errors, and ϕ is the *misfit* of model \mathbf{m} . In the case of MT, the data mis-

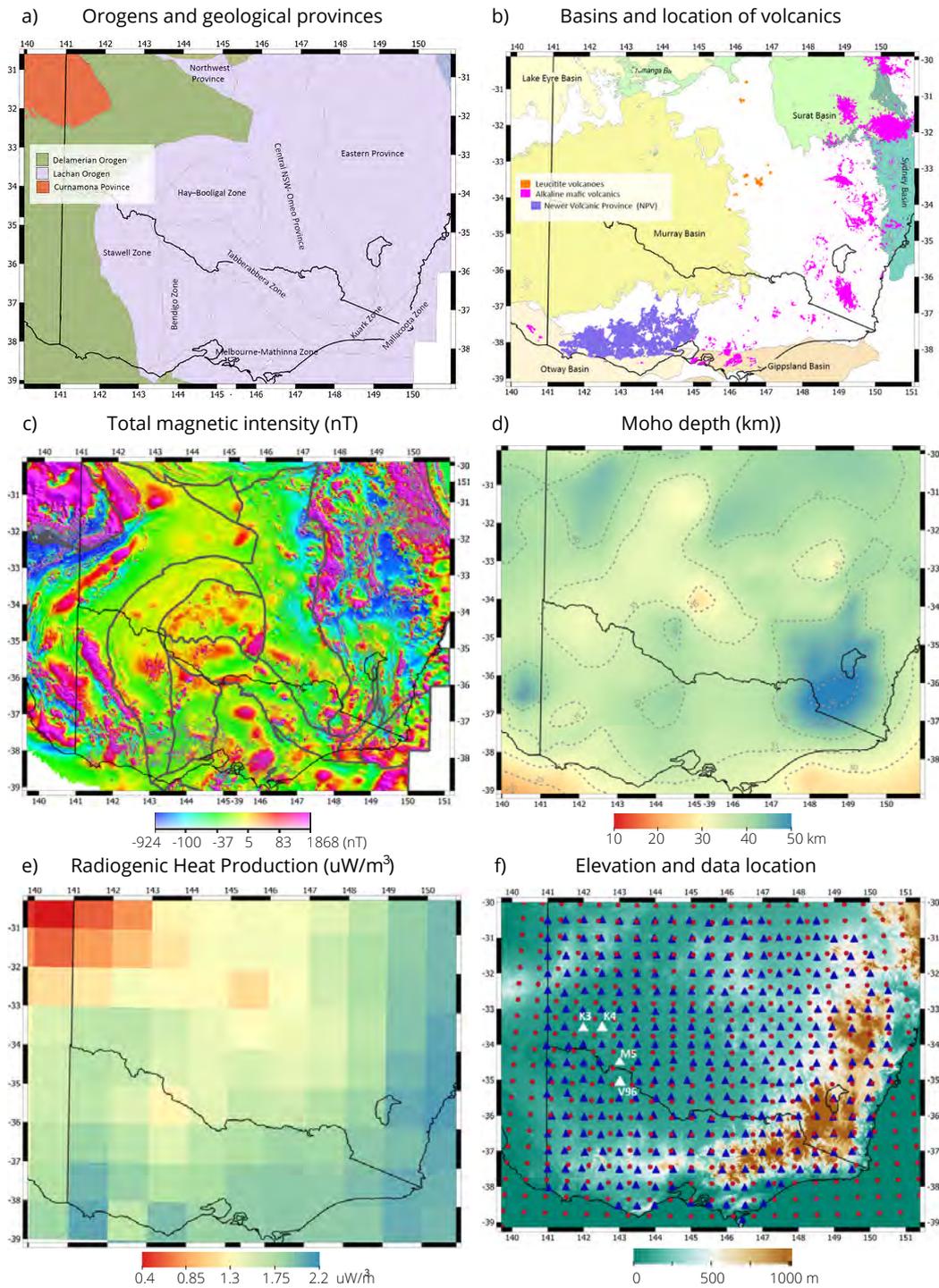
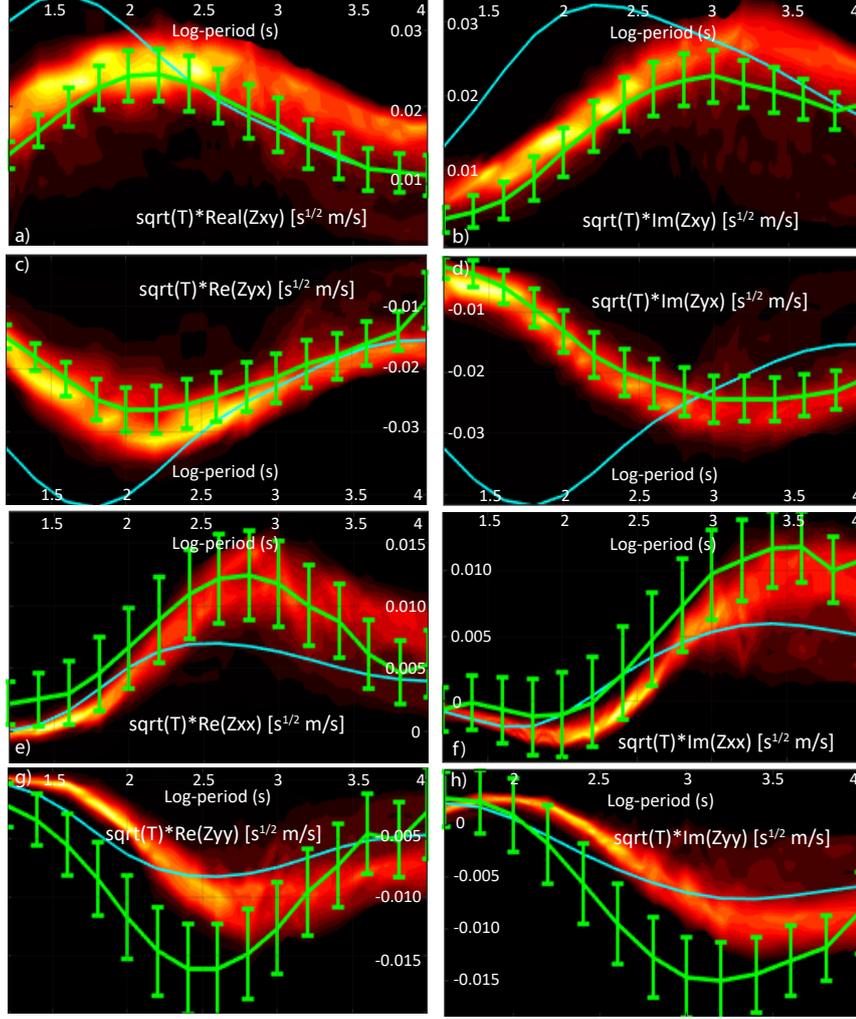


Figure 1. Maps showing (a) orogens that comprise the Tasmanides of southeastern Australia with grey outline denoting geological provinces (Raymond et al., 2018); (b) Mesozoic to Cenozoic sedimentary basins after Raymond et al. (2012), leucitites volcanoes (orange) and basaltic volcanics (pink) after Shea et al. (2022). The basaltic in NV are highlighted in purple. (c) Total magnetic intensity map (TMI) which includes airborne-derived TMI data for onshore and near-offshore continental areas (Nakamura & Milligan, 2015). (d) Moho depth from the AusREM model (Kennett & Salmon, 2012) where 5km-contour lines are shown in dashed-grey. (e) Mean crustal RHP from Haynes et al. (2020) (f) Elevation map of southeast Australia including the AusLAMP MT stations (blue triangles) and the location of the velocity data (red dots). Panel (a) and (c) show major tectonic boundaries are outlined in grey. White triangles indicate stations where data fits are shown.

I. MT data at station M5 (143 E, 34.463 S)



II. a) Vp data at ST15 (139.71 E, 32.74 S)

b) Vp data at ST181 (144.4 E, 33.2 S)

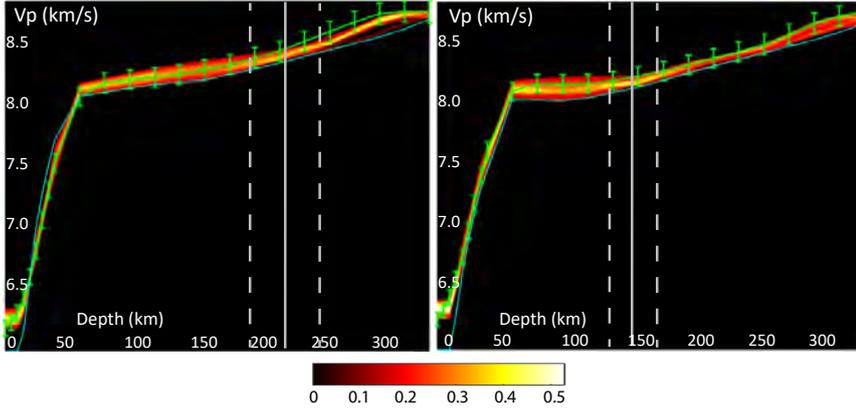


Figure 2. I. Posterior PDFs (refer to next section) of MT data for station M5. Field data and error bars are plotted in green and the computed data for the initial model is plotted in blue. Panels (a), (b), (c) and (d): Posterior PDFs of the real and imaginary parts of the off-diagonal components (Z_{xy} and Z_{yx}). Panels (e), (f), (g) and (h): Posterior PDFs of the real and imaginary parts of the diagonal components (Z_{xx} and Z_{yy}). The data has been scaled by the square-root of the period (T) in all panels. The location of the station is shown in Figure 1f.

II. Posterior PDFs (refer to next section) of P-wave velocity data for stations (a) ST15 located at 139.71 E, 32.74 S and (b) ST182 at 144.4 E, 33.20 S. P-wave velocity data and error bars are plotted in green and the computed data for the initial model is plotted in blue. For those locations, the LAB depths corresponding to the mean, lower and upper bound of the 68% CI models are shown in solid and dashed grey lines, respectively

fit is given by (Tarantola, 2005):

$$\phi = - \sum_{i=1}^N \frac{|g_i(\mathbf{m}) - d_i(\mathbf{m})|}{s_i}, \quad (2)$$

whereas the misfit for the seismic data takes the following form:

$$\phi = - \frac{1}{2} \sum_{i=1}^N \left(\frac{g_i(\mathbf{m}) - d_i(\mathbf{m})}{s_i} \right)^2. \quad (3)$$

218 For each data set, \mathbf{g} is the solution of a particular forward problem for model \mathbf{m} , N is
219 the total number of data and s_i denotes the standard deviation for the i -th data error.

220 The posterior PDF over data and parameters is commonly approximated using sampling-
221 based Markov chain Monte Carlo (MCMC) algorithms (Gilks et al., 1995). In our joint
222 inversions of independent data sets, we use the Delayed Rejection Adaptive Metropol-
223 is (DRAM) scheme of Haario et al. (2006) in combination with the Cascaded Metropol-
224 is (CM) approach (Tarantola, 2005; Hassani & Renaudin, 2013; Manassero et al., 2021).
225 Details about the general inversion framework (RB+MCMC) are given in Manassero et
226 al. (2021) while particular details about the sampling strategy, prior information, and
227 the initial seismic velocity and electrical conductivity models used for the current inver-
228 sion are given in Section S1 of the Supplementary Material.

229 3.3 Model parameters

230 In order to define the model parameters and their interdependence in the joint in-
231 version, we distinguish between primary and secondary parameters (Khan et al., 2006;
232 Afonso et al., 2013a; Manassero et al., 2021). The latter are directly linked to the prop-
233 erties used to solve the forward problems in their classic forms (e.g., V_p and electrical
234 conductivity). The former are the fundamental thermodynamic parameters, namely, tem-
235 perature (T), pressure (P) and bulk composition (C). These control the magnitudes of
236 the secondary parameters in the mantle via equations of state and thermodynamic con-
237 straints (this applies to the mantle only; for crustal parameters see below). A specific
238 configuration of the primary parameters in the 3D space defines what we refer to as the
239 *background state* (or background contribution). In this way, the background P-wave ve-
240 locity and electrical conductivity in the mantle can be written as $V_p(T, P, C)$ and $\sigma_b(T, P, C)$.

241 As shown by Afonso et al. (2013a, 2013b) and Manassero et al. (2021), an efficient
242 way to parameterize the background state is to divide the 3D space into m rectangular
243 columns and use the following model parameters in each column: i) the depth to the ther-
244 mal lithosphere-asthenosphere boundary (LAB), here defined as the 1250°C isotherm (Afonso,
245 Moorkamp, & Fullea, 2016), and ii) n ‘thermodynamic nodes’ distributed throughout
246 the mantle. The LAB depths allow us to solve for a lithospheric conductive temperature
247 profile, while the temperature of the thermodynamic nodes placed in the sub-lithospheric
248 mantle are allowed to vary during the inversion as required by the inverted data. The
249 computation of the pressure (P) and definition of the bulk composition (C) are described
250 below and in Section 3.5, respectively.

251 Since the electrical conductivity is also highly sensitive to factors other than T, P,
252 and C (e.g., hydrogen content, localized fluid/melt pathways, presence of hydrated phases
253 or graphite), we expand the space of secondary parameters and write $\sigma = \sigma_b(T, P, C) +$
254 $\sigma(X)$, where X stands for any factor that cannot be captured by the background. This
255 means that $\sigma(X)$ is a representation of any anomalous conductivity associated with pro-
256 cesses superimposed on the background state (Manassero et al., 2021). This anomalous
257 conductivity contribution and the conductivity in the crust ($\sigma_c(X)$) are parameterized
258 with l conductivity nodes distributed throughout the whole domain (see details in Sec-
259 tion S3 in Supplementary Material and in Manassero et al., 2021).

260 In order to parameterize the rest of the properties of the crust, we divide the crust
 261 into three layers (sediments, upper crust and lower crust) from the surface to the Moho.
 262 Within each column, layers have fixed thicknesses and their own set of physical prop-
 263 erties: thermal conductivity (k), volumetric radiogenic heat production (RHP) and P-
 264 wave (V_p) velocity. During the inversion, only V_p is allow to vary within their assigned
 265 uncertainties (Rawlinson et al., 2016); all remaining parameters are assumed constant.
 266 The thermal conductivity of the crustal layers are set to $k^1=2.8$, $k^2=2.6$ and $k^3=2.3 W/m^\circ C$.
 267 The crustal RHP is obtained using the mean crustal RHP from a previous 1D joint prob-
 268 abilistic inversion (Figure 1.e Haynes et al., 2020) while the Moho depths are taken from
 269 the regional AusREM model (Kennett & Salmon, 2012). We also incorporate one con-
 270 ductivity cell below each MT station as extra parameters to account for the galvanic dis-
 271 tortion effect produced by near-surface inhomogeneities beyond the resolution of our model
 272 (Jones, 2011; Chave & Jones, 2012; Avdeeva et al., 2015). Similarly to the approach used
 273 in ModEM (Kelbert et al., 2014), these cells are placed in the first (thin) layer of the nu-
 274 merical mesh used to solve the MT forward problem.

275 3.4 Forward Problems and Model Discretization

276 The main forward problems solved during the probabilistic inversion are the 3D
 277 MT problem and the conductive heat transfer in the lithosphere. These have been de-
 278 scribed in detail in Manassero et al. (2020, 2021) and in Afonso et al. (2013a, 2013b);
 279 Afonso, Rawlinson, et al. (2016), respectively. In what follows, we focus on the model
 280 discretization and on the derivation of the seismic velocity and background conductiv-
 281 ity models given a realization of the primary parameters.

282 The study area is subdivided into 441 columns of size $0.45^\circ \times 0.45^\circ \times 410$ km. Each
 283 column is discretized at three different scales:

- 284 1. The coarser discretization includes the mantle thermodynamic nodes, placed ev-
 285 ery 50 km in the vertical direction. These nodes are used to obtain stable min-
 286 eral assemblages and physical properties in the mantle by Gibbs free-energy min-
 287 imization (Afonso et al., 2013b).
- 288 2. The intermediate discretization comprises the finite elements (FE) used to solve
 289 the MT forward problem. In each column, we have $3 \times 3 \times 36$ FE (a total of $63 \times$
 290 63×36 FE in the whole domain) of size 17×17 km in the horizontal and vari-
 291 able vertical size with depth. The air comprises four FE cells and a total thick-
 292 ness of 106 km.
- 293 3. A fine mesh (2 km) is used to solve the steady-state heat transfer equation within
 294 the lithosphere (via a FE algorithm), subject to Dirichlet boundary conditions at
 295 the LAB ($T_{LAB} = 1250$ °C) and at the model's surface ($T_S = 10$ °C) (Afonso et
 296 al., 2013b).

297 During the probabilistic inversion, a realization of the background parameters in-
 298 cludes a specific LAB depth and temperatures for all the sub-lithospheric thermodynamic
 299 nodes in the entire domain, both randomly sampled from their prior distributions. Af-
 300 ter solving for the conductive geotherm corresponding to the sampled LAB depth, we
 301 interpolate the temperatures to the lithospheric thermodynamic nodes (i.e. those ther-
 302 modynamic nodes that reside inside the lithosphere). The pressure is computed at all
 303 thermodynamic nodes using a quadratic lithostatic-type approximation (see Section S2
 304 of Supplementary Material). Using these T, P and a pre-defined composition, we retrieve
 305 all thermo-physical properties at the thermodynamic nodes from pre-computed tables.
 306 These tables are calculated by Gibbs free-energy minimization with components of the
 307 software Perple_X (Connolly, 2009; Afonso et al., 2013b) and the database and thermo-
 308 dynamic formalism of Stixrude & Lithgow-Bertelloni (2011), within the CFMAS system
 309 ($CaO, FeO, MgO, Al_2O_3, SiO_2$). All thermophysical properties computed at the ther-
 310 modynamic nodes are linearly interpolated to the fine mesh to create the correspond-

311 ing seismic model and to the FE mesh for the computation of the conductivity model
 312 (see details in Section S2) and the MT forward solution.

313 3.5 Mantle composition

314 The pre-computed tables and their equilibrium assemblages are computed using
 315 a mean bulk mantle composition (i.e., specific CFMAS compositions) of 44.3 wt% SiO_2 ,
 316 2.8 wt% Al_2O_3 , 8.5 wt% Fe_2O_3 , 39.3 wt% MgO and 2.7 wt% CaO . We estimate this
 317 mean composition by averaging eight spinel lherzolites xenoliths (see Table S2 in Sup-
 318plementary material) that were entrained in EAVP lavas. We use major element com-
 319positions from Irving (1980), O'Reilly & Griffin (1987), Griffin et al. (1987) and unre-
 320ported samples from Bokhara River (J. Shea, personal communications), which cover the
 321area of interest. Since this is the most recent volcanism in eastern Australia, these xeno-
 322liths are the most representative samples of current mantle compositions available.

323 The use of an average mantle composition is justified by the fact that V_p and elec-
 324trical conductivity have second-order sensitivity to (dry) bulk mantle composition (see
 325Figure S1, S2 and Özaydın & Selway, 2020; Trampert et al., 2001; Goes et al., 2000). We
 326also assume a dry mantle composition for the background properties. The reasons for
 327this choice are: i) V_p is not significantly affected by the small amounts of water com-
 328monly observed in mantle samples (Yu et al., 2011; Cline Ii et al., 2018, and references
 329therein), and ii) the conductivity nodes can represent any positive anomalies (e.g., wa-
 330ter content) over the background values (which represents the most resistive end-member
 331at the given T-P-C conditions), which reduces the number of parameters by two (see Man-
 332assero et al., 2021).

333 3.6 Mantle water content as a proxy for metasomatism

334 Using outputs of the joint probabilistic inversion (thermal structure, conductivity
 335models, equilibrium assemblages and mineral compositions), we can estimate the bulk
 336water content in the mantle (i.e. hydroxyl or OH^- bound to nominally anhydrous min-
 337erals) that would be required to explain the inversion results. Importantly, the water con-
 338tent as estimated here lumps all unmodeled chemical effects resulting in high conduc-
 339tivity (e.g. connected sulfides, presence of melt) and it is taken as a proxy for mantle meta-
 340samatism (see Discussion)

341 The water content computations are done using the software MATE (Özaydın &
 342Selway, 2020), which includes several experimental models for electrical conductivity, wa-
 343ter partitioning and solubility (based on petrological studies). In particular, we used the
 344electrical conductivity models of Gardés et al. (2014), Dai & Karato (2009a), Liu et al.
 345(2019), and Dai & Karato (2009b) for olivine, orthopyroxene, clinopyroxene and garnet,
 346respectively.

347 The solutions for the water content lie between the bounds defined by the dry litho-
 348sphere (i.e., 0 ppm) and the maximum bulk water content calculated using the olivine
 349water solubility model of Padrón-Navarta & Hermann (2017). The experimental coef-
 350ficients used in the water partitioning are: $D_{opx/ol}^{OH} = 5.6$, $D_{cpx/opx}^{OH} = 1.9$ of Demouchy
 351et al. (2017) and $D_{gt/ol}^{OH} = 0.8$ of Novella et al. (2014); which reflect the sub-solidus con-
 352ditions found in the continental lithospheric mantle in southeastern Australia. Since we
 353aim to portray variations of water content in the mantle rather than fitting the real wa-
 354ter content seen in xenoliths, the choice of the experimental parameters is adequate for
 355our calculations. All water calculations are done using the calibration of Withers et al.
 356(2012) for olivine, and the calibration of Bell et al. (1995) for pyroxenes and garnet.

357 The electrical conductivity of each individual mineral phase is turned into bulk con-
 358ductivity through the Generalised Archie's Law (Glover, 2010) with cementation com-
 359ponents (m) of $m = 2$ for orthopyroxenes, $m = 4$ for clinopyroxenes and garnet, and

360 $m < 1$ for olivine (perfectly connected). The Generalised Archie’s Law is preferred over
 361 the conservative estimates of Hashin-Shtrikman lower-bound since it allows us to incor-
 362 porate the effects of specific minerals in the conductivity values, such as highly-interconnected
 363 phlogopites. The main cementation components used here, however, provide similar val-
 364 ues to the Hashin-Shtrikman lower-bound for a lherzolitic matrix (Özaydın & Selway,
 365 2020).

366 4 Results

367 4.1 Thermal structure of Southeast Australia

368 The depth to the LAB obtained from the joint probabilistic inversion of seismic ve-
 369 locities and MT data is shown in Figure 3; the complete 3D temperature structure is shown
 370 via depth slices in Figure 4 (first three columns). Figure 3 also includes a recent LAB
 371 model obtained from a low-resolution 1D joint probabilistic inversion of elevation, sur-
 372 face heat flow, Rayleigh wave dispersion curves, and geoid anomalies (Haynes et al., 2020;
 373 Afonso et al., 2013b) and the estimated LAB depths from two recent seismic tomogra-
 374 phy models in eastern Australia (Davies et al., 2015; Rawlinson et al., 2017). The max-
 375 imum absolute difference in lithosphere thickness inside the region is ~ 200 km, with shal-
 376 low LAB depths (< 100 km) in the eastern and southern ends of the continental block
 377 and deep LAB depths (> 250 km) beneath the Curnamona Province (CP) and the north-
 378 ern part of the Delamerian Orogen. We observe a clear correlation between lithospheric
 379 structure and the locations of recent volcanism (e.g., Figure 3.a): leucitite volcanic cen-
 380 ters correlate with regions of intermediate lithospheric thickness (125-160km) while the
 381 basaltic volcanoes are located in regions where the LAB is shallower than 120 km. At
 382 least two clear and step-like changes in LAB depth are observed in our model along a
 383 transect from the CP to the southeast corner of the model and across the leucitite vol-
 384 canoes (see also Figures 10).

385 The first order features of our LAB model are in good agreement with the mean
 386 LAB obtained by Haynes et al. (2020), even though the data sets used in each inversion
 387 are different. Many features in our LAB model are also present in those derived from
 388 seismic velocity models (Davies et al., 2015; Rawlinson et al., 2017). In particular, we
 389 observe similar LAB depths beneath the basaltic volcanoes and west of 146°E , where a
 390 wedge-like structure follows the curvature of the Stawell Zone (SZ, see Figure 3.1). All
 391 models show a thickening of the lithosphere towards the northwest part of the region.
 392 However, some significant discrepancies are found in the CP and towards the center of
 393 the model. Beneath the CP, our LAB depths are considerably larger than those of Davies
 394 et al. (2015) and Rawlinson et al. (2017). While one could attribute some of this differ-
 395 ence to the fact that the mantle composition in this area is likely more depleted (and
 396 thus ‘faster’) than the average bulk mantle composition used here (see Section 3.5), ad-
 397 ditional calculations shown in Figure S3.a of the Supplementary Material reject this pos-
 398 sibility as the main cause. Rather, the main reason is the different definitions of LAB
 399 adopted in these works. While in shallow lithospheric environments there is a marked
 400 minimum in V_p near the LAB (and thus easy to pick), deep lithospheric environments
 401 are characterized by smooth V_p profiles, which makes it harder to choose the thermal LAB
 402 unambiguously (Figures S3.b and 2.II) based on V_p profiles only. This therefore explains
 403 why our LAB estimates are similar to those in Davies et al. (2015) and Rawlinson et al.
 404 (2017) in thin lithospheric settings, but they diverge in absolute magnitude when the LAB
 405 gets thicker.

406 The depths to the thermal LABs obtained after an RB+MCMC inversion using MT
 407 data alone (Figure S5 in Supplementary Material) are shown in Figure S6 (Supplemen-
 408 tary Material). Compared to the results from the joint inversion, these figures show large
 409 variability and no clear trend in the thickness of the lithosphere from the CP to the south-
 410 east corner of the model (a feature observed in all other models). This comparison il-

411 illustrates well the facts that i) MT alone has difficulties in discriminating thermal causes
 412 from other factors controlling the electrical conductivity in the mantle (Jones, 1999) and
 413 ii) other types of data (e.g., seismic) need to be included when imaging lithospheric struc-
 414 ture.

415 4.2 Seismic velocity structure

416 Depth slices of the P-wave velocity structure predicted by our model are shown in
 417 Figure 4. The P-wave velocity model of Rawlinson et al. (2016) is also shown for refer-
 418 ence. In all cases, the velocities are plotted relative to the AusREM model at $34.4^{\circ}S, 145^{\circ}E$
 419 (Figure S4 in Supplementary Material). We observe that the inversion succeeded in re-
 420 producing the V_p structure of the input model (Rawlinson et al., 2016). In particular,
 421 the mean P-wave velocity down to 100 km is practically identical in both models. The
 422 Newer Volcanic province stands out as a low-velocity anomaly at depths between 60 and
 423 80 km, whereas the basaltic volcanoes in the middle of the Eastern Province ($\sim 149^{\circ}E,$
 424 $34^{\circ}S$) correlate well with deeper low-velocity anomalies.

425 Some minor discrepancies between our results and the model of Rawlinson et al.
 426 (2016) are observed at depths > 100 km. For instance, we obtain slightly higher seismic
 427 velocities (0.6% higher on average) in the depth range 100-180 km at the eastern end of
 428 the model. Supplementary tests allows us to attribute these differences to the constraints
 429 imposed by the MT in the joint inversion. Similarly, we obtain slightly slower velocities
 430 throughout the whole model at 200-220 km depth (see Figure 4). At these depths, the
 431 local discrepancies are consequence of the different physical parameterizations used in
 432 this work and by Rawlinson et al. (2016). Nevertheless, neither the original tomography
 433 model of Rawlinson et al. (2016) nor our model have sufficient resolution at these depths
 434 to justify further comparisons.

435 4.3 Electrical conductivity structure

436 The conductivity models for the crust and mantle predicted by the joint inversion
 437 are shown in Figures 5 and 6, respectively. For comparison, these figures include the re-
 438 sults obtained from a recent deterministic inversion of MT data (Kirkby et al., 2020),
 439 using the ModEM software (Kelbert et al., 2014). The main structures observed in the
 440 conductivity models are comparable (within model uncertainties) to those in the model
 441 of Kirkby et al. (2020) at all depths.

442 4.3.1 Crust

443 Figure 7 illustrates the agreement between the conductivity structure in the crust
 444 and other sources of information: sedimentary basins (Raymond et al., 2012), magnetic
 445 anomalies (Nakamura & Milligan, 2015) and a shear velocity model (Pilia et al., 2015).
 446 In particular, Panel (a) shows that the extent of the Paleozoic to Cenozoic sedimentary
 447 basins in the region is well outlined by the mean conductivity model at 2 km depth. A
 448 comparison between Panels (c,e,f) and Panel (b) highlights the correlation between to-
 449 tal magnetic anomalies and conductivity features at different depths. Examples of these
 450 are (A) a conductor in the CP; (B) a SW-NE linear structure close to the NW limit of
 451 Murray Basin; conductors (C) and (D) in the Tabberabbera Zone; (E) a N-S conduc-
 452 tor aligned with the western border of the Eastern Province; (F) two resistive structures
 453 west of the Sydney Basin; (G) the Sydney Basin; (H) a conductive region aligned with
 454 the north limit of the Northwest and Central NSW provinces; (I) a highly resistive re-
 455 gion in the Stawell Zone near the NSW-VIC border; (J) a circular structure in the mid-
 456 dle of the model; (K) a high-conductivity anomaly; and (L) a conductor east of CP.

457 The conductor (A) correlates well with the conductor seen in the MT study of Robert-
 458 son et al. (2016), using data from 74 AusLAMP stations placed in the Ikara–Flinders Ranges

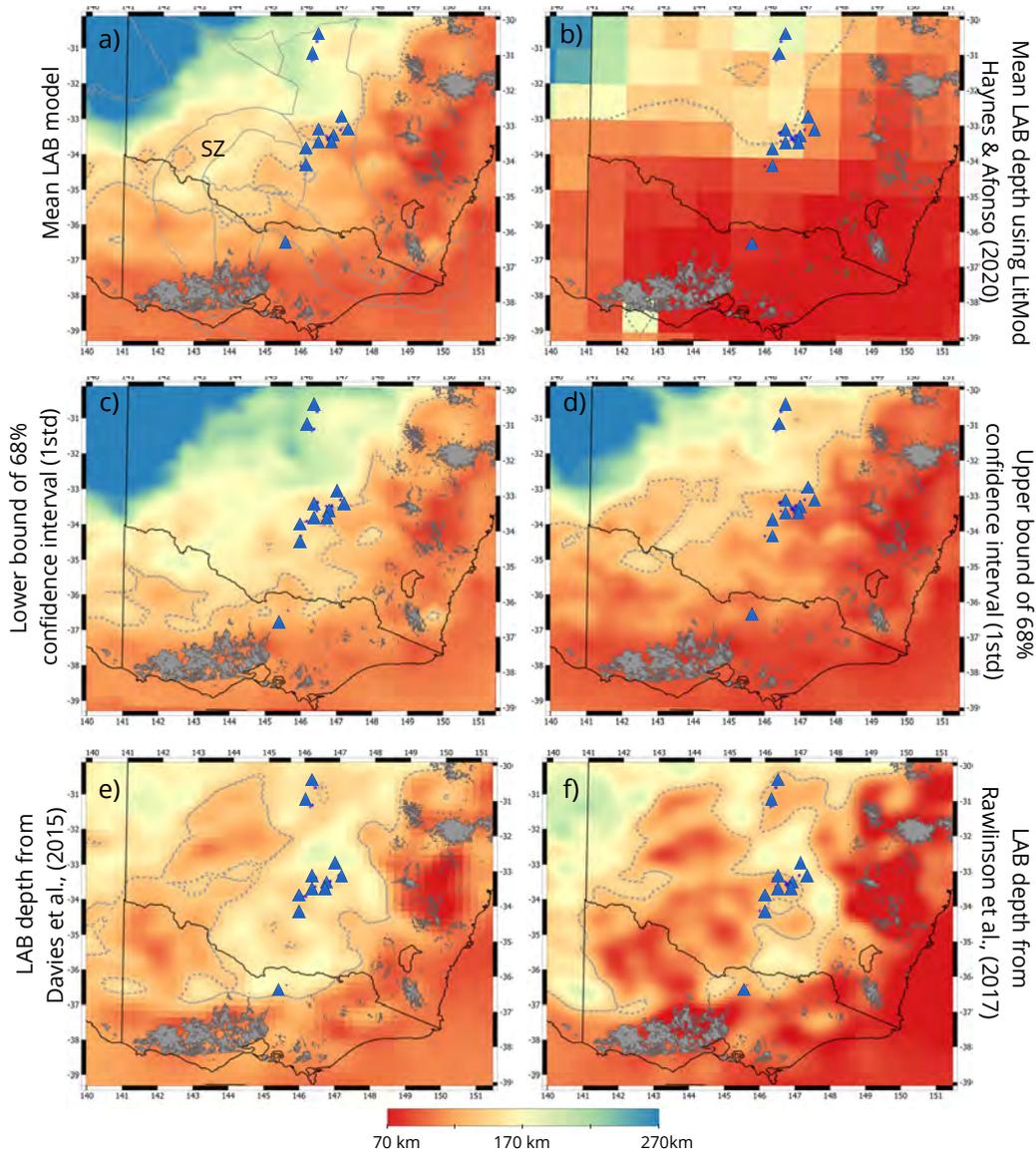


Figure 3. Depth of the thermal LAB. (a) Mean model after the joint probabilistic inversion; (b) mean model obtained after a 1D joint probabilistic inversion (Haynes et al., 2020; Afonso et al., 2013b); (c) and (d) lower and upper bounds of the 68 % confidence interval (1 standard deviation from the mean), respectively. (e) and (f) depth of the LAB after Rawlinson et al. (2017) and Davies et al. (2015), respectively. The location of leucite-bearing volcanism are shown in blue and standard basaltic volcanoes in grey. The 140 km-contour of the LAB depth is shown in dashed-grey line and the outline of the tectonic provinces in solid grey lines. The location of the Stawell Zone (SZ) is marked in panel (a).

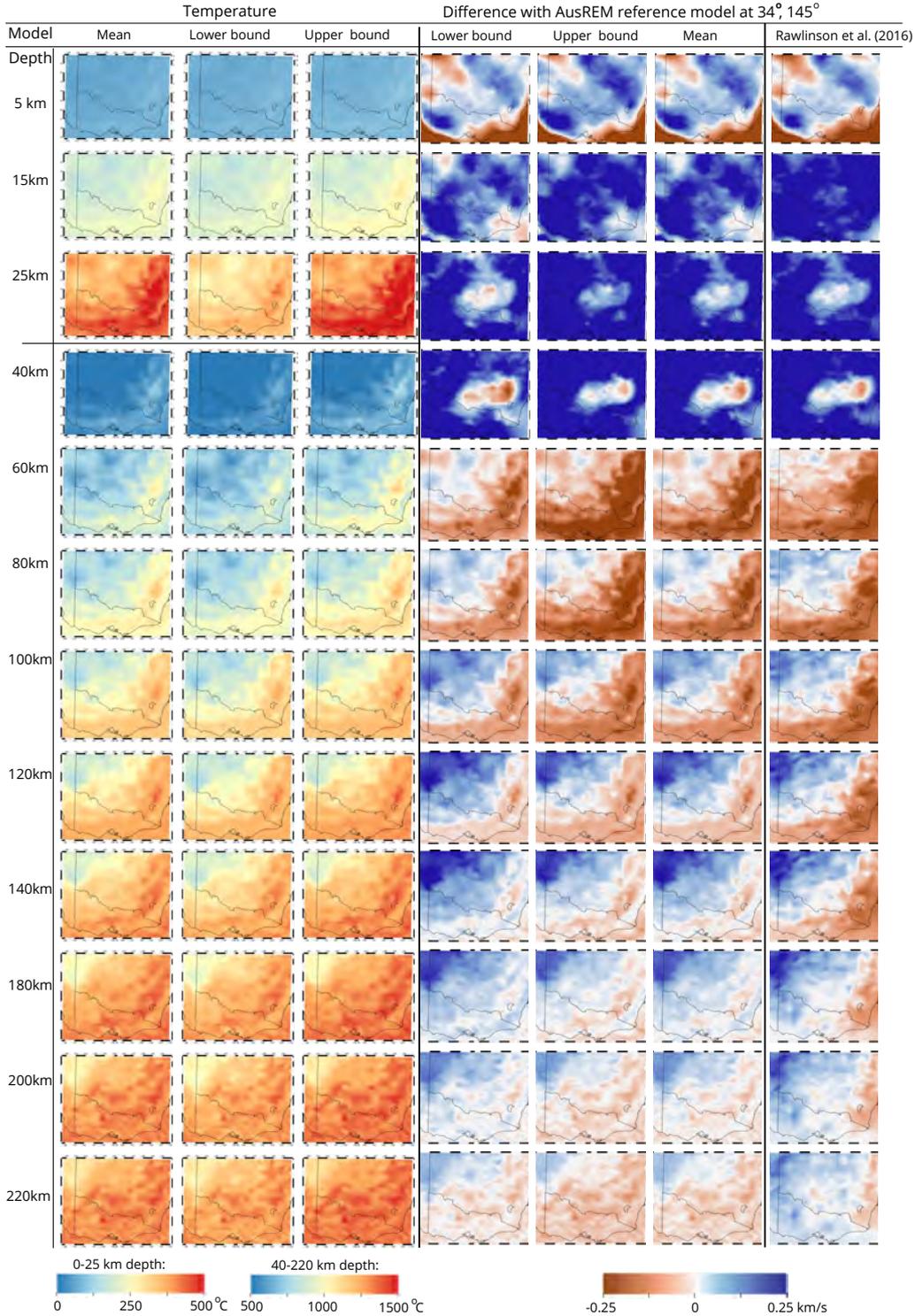


Figure 4. Columns (1)-(3): depth slices from the (1) mean model and those models corresponding to (2) the lower and (3) upper bound of the 68% CI of the posterior PDF for the temperature. Columns (4)-(6): depth slices from the models corresponding to (4) the lower and (5) upper bound of the 68% CI, and (6) mean of the posterior PDF for the P-wave velocity; Column (7): P-wave velocity model of (Rawlinson et al., 2016). Selected depths are shown on the left of the figure. In all cases, velocities are plotted relative to 1-D reference model AusREM at 34.4°S, 145°E shown in Figure S4 of the Supplementary Material.

459 and CP, and in the recent study of Kay et al. (2022), using a densely-spaced MT mod-
 460 eling scheme. Comparing Panels (c) and (d), we observe that the main conductivity struc-
 461 tures correlates well with velocity anomalies imaged by the shear-wave velocity model
 462 of Pilia et al. (2015). We note that the concentric geometries at 29 km depth, such as
 463 conductor (J) and structures on the west of the model, resemble the features of the Lach-
 464 lan Orocline model revealed by potential field and passive seismic data (c.f. Kirkby et
 465 al., 2020).

466 **4.3.2 Mantle**

467 The conductivity models between 80-250 km depth (Figures 6) largely resemble the
 468 ModEM model of Kirkby et al. (2020). In particular, we observe a similar north-eastward
 469 orientation of the conductors in the middle of the model (C_1, C_2, C_3, C_4 and C_5 in Fig-
 470 ures 6 and 8). Comparing the mantle conductivity with the mean LAB structure in Fig-
 471 ures 8, our models suggest that there is some correlation between the LAB topography
 472 and these mantle conductors (cf. Kirkby et al., 2020). In particular, the general NE-SW
 473 trend of the conductors tend to follow the LAB depth structure. Conductor C_5 aligns
 474 well with the LAB wedge northwest of the model, whereas C_{3N} and C_{4S} tend to follow
 475 the 120-140 km LAB depth iso-surfaces. C_1, C_2 and C_{3S} are located in regions where LAB
 476 depths < 120 km. We also observe an intra-lithosphere high-conductivity structure be-
 477 neath the CP (C_6), which agrees well with the structure imaged by previous MT stud-
 478 ies (e.g., Robertson et al., 2016; Thiel & Heinson, 2013). A high-conductivity region (C_4)
 479 is observed below the central-leucitite volcanoes. In our models, the extent of this re-
 480 gion is larger and more connected than in the ModEM model.

481 The main difference between our conductivity model and that of Kirkby et al. (2020)
 482 is that the sub-lithospheric conductivities along the south-east coast and in the middle
 483 of the region are higher in our model (R_1 and R_2 in Figure 6). The same is true when
 484 we compare our model with the results from a probabilistic inversion of MT data only
 485 (Figure S5, Supplementary Material). The high resistivity ($> 10^4 \Omega m$) values in MT-
 486 only inversions are at odds with the mantle resistivity range obtained for sub-lithospheric
 487 temperatures and pressures (Fullea et al., 2011; Naif et al., 2021). We observe that, due
 488 to the constraint imposed by seismic data in the joint probabilistic inversion (via the ther-
 489 mal structure), unrealistically high resistivity values are not present in our model.

490 Another example of the constraint imposed by the seismic data in the conductiv-
 491 ity models is shown in Figures 8. At 140 km depth, we observe that the conductors in
 492 the east (C_1, C_2 and C_{3S} in Figure 8.c) are located within a region defined by a 1250°C -
 493 contour (Figure 8.b). These mantle conductivity structures correlate well with both the
 494 location of the eastern basaltic volcanics and a stripe of low P-wave seismic velocities
 495 (Figure 8.d). At the same time, the stripe of high seismic velocities beneath the east coast
 496 at a depth of 140 km is a clear example of the constraint imposed by the MT data in
 497 the velocity models. This stripe is not seen in the models of Rawlinson et al. (2016) and
 498 correlates with a relatively cold and highly resistive mantle (Figures 8.b-c).

499 The similarities between the mantle conductivity models at ~ 40 -80 km depth and
 500 features found in the gravity anomalies (Nakamura, 2016) are shown in Figures S7 in the
 501 Supplementary Material.

502 **4.4 Joint assessment of bulk water content and temperature maps**

503 The bulk water content maps derived from the mantle conductivity models are shown
 504 in Figures 9. As mentioned in Section 3.6, we emphasize that “bulk water content” is
 505 a lumped proxy for general mantle metasomatism and therefore their absolute values need
 506 to be taken with caution. We observe that most of the localized conductive anomalies
 507 above the background require relatively high bulk water contents. This is the case for

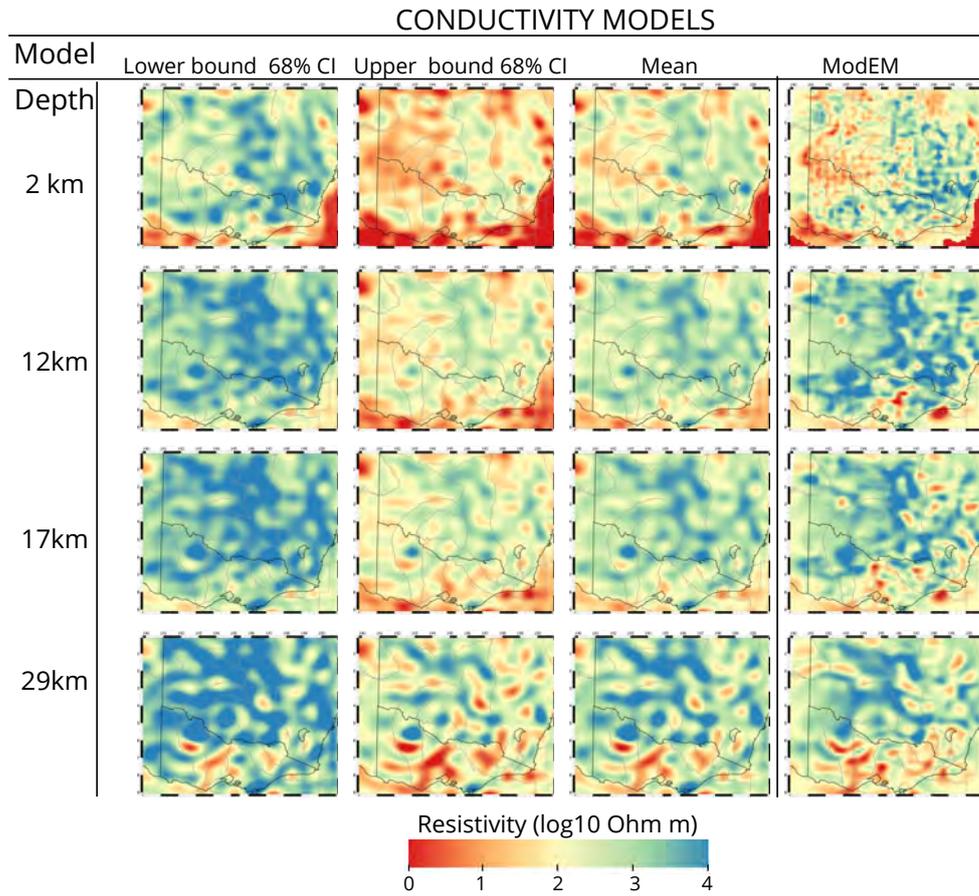


Figure 5. Conductivity in the crust from the joint probabilistic inversion. Columns (1)-(3): depth slices from the (1) the lower, (2) upper bound of the 68% percentile and (3) mean conductivity models of the posterior PDF. Column (4): conductivity model of (Kirkby et al., 2020). Selected depths are shown on the left of the figure and the boundaries of geological provinces are shown in grey lines.

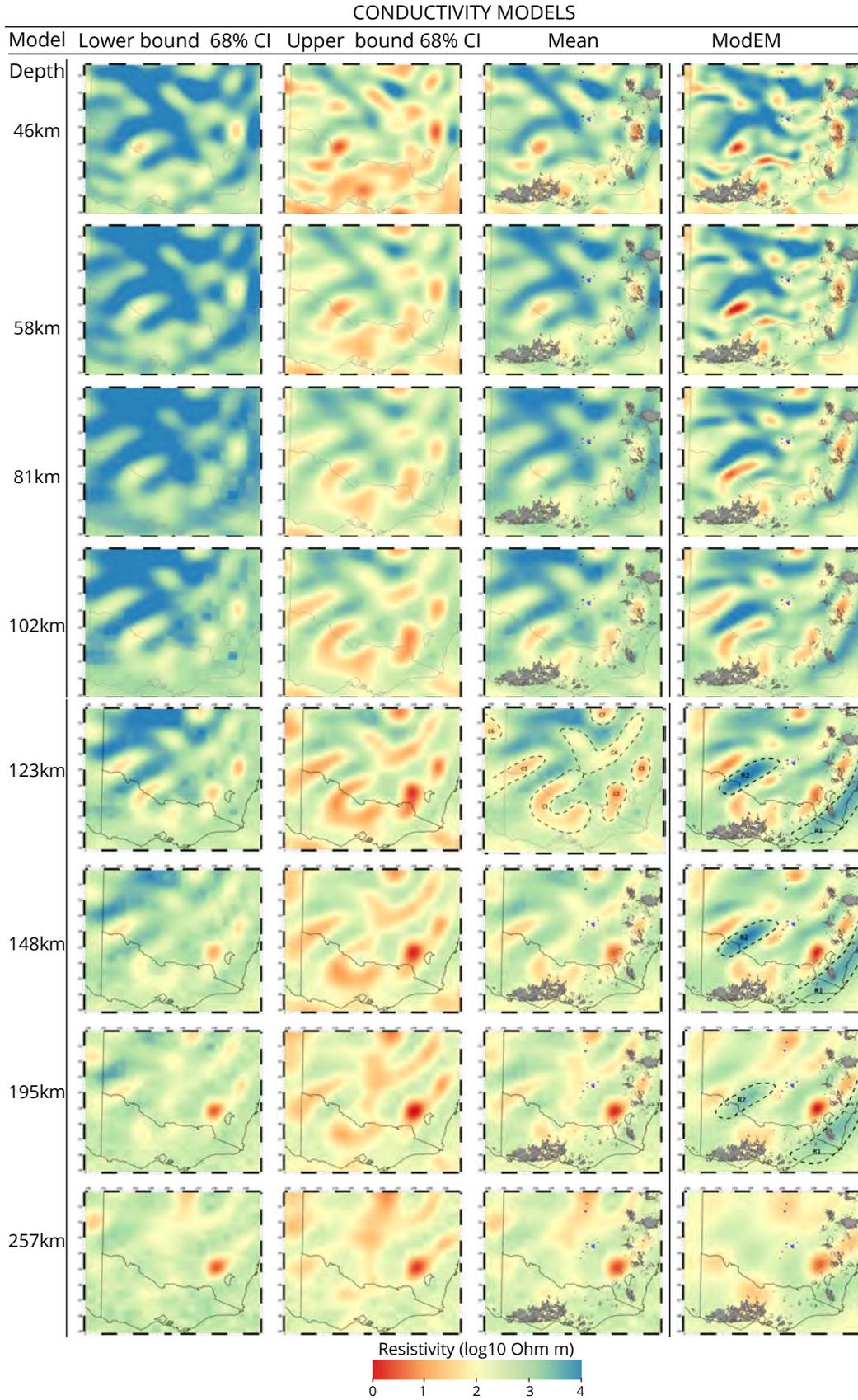


Figure 6. Mantle conductivity from the joint probabilistic inversion. Columns (1)-(3): depth slices from the (1) the lower, (2) upper bound of the 68% percentile and (3) mean conductivity models of the posterior PDF. Column (4): conductivity model of Kirkby et al. (2020). The location of leucitite-bearing volcanism are shown in blue and standard basaltic volcanoes in grey. Selected depths are shown on the left of the figure. Dashed-black lines highlight conductors in the mean model and resistors in the ModEM model below 123 km depth

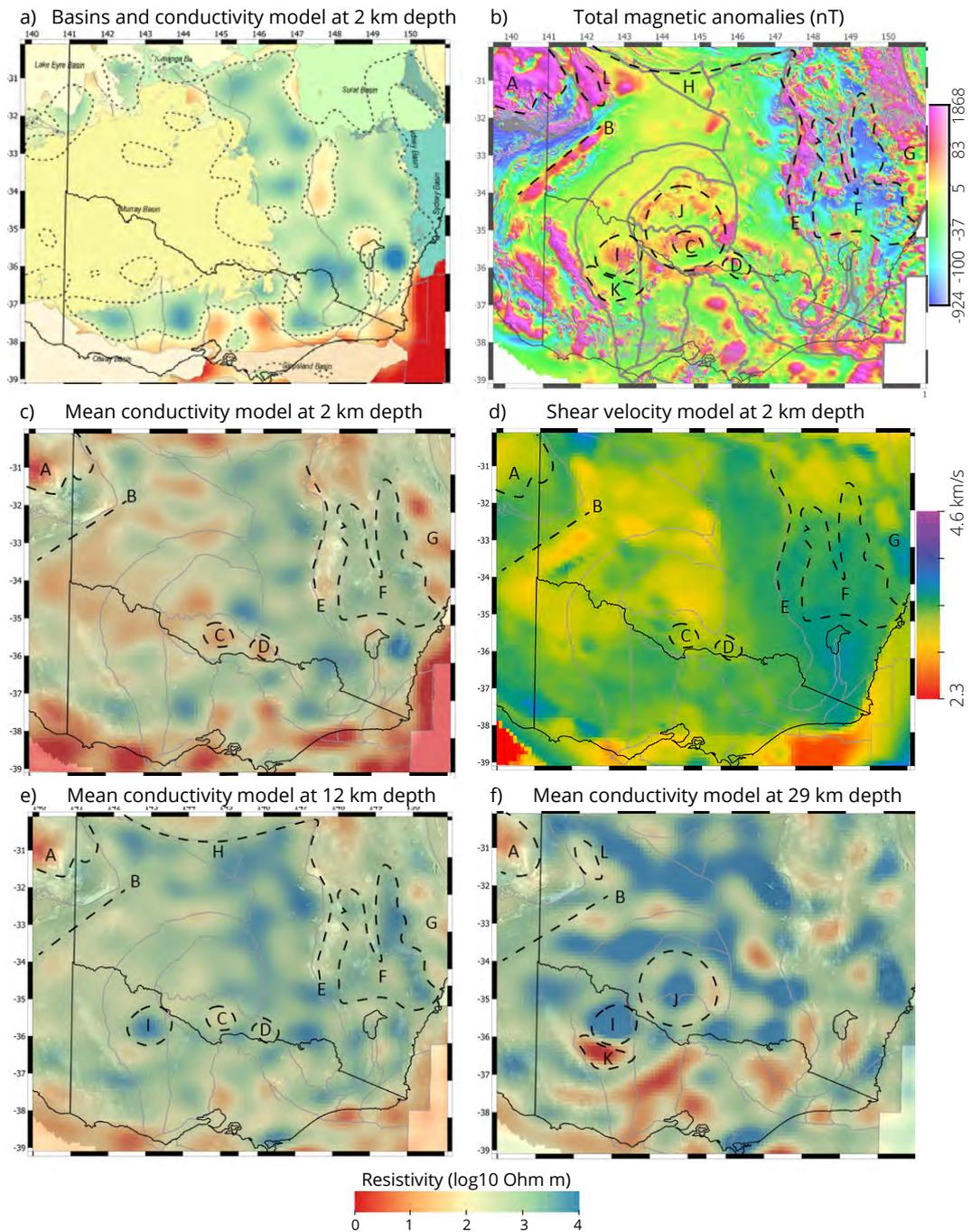


Figure 7. (a) Sedimentary basins overlying mean conductivity model at 2 km depth and 200 Ωm -resistivity contour in dash lines. (b) Total magnetic anomalies after Nakamura & Milligan (2015) (c) Mean conductivity (overlying the magnetic anomalies in grey scale) and (d) shear wave velocity model after Pilia et al. (2015) at 2 km depth. (e-f) Mean conductivity models at 12 and 29 km depth overlying the magnetic anomalies in grey scale. We refer the reader to the main text for a description of structures A-L. Boundaries of geological provinces are shown in grey lines.

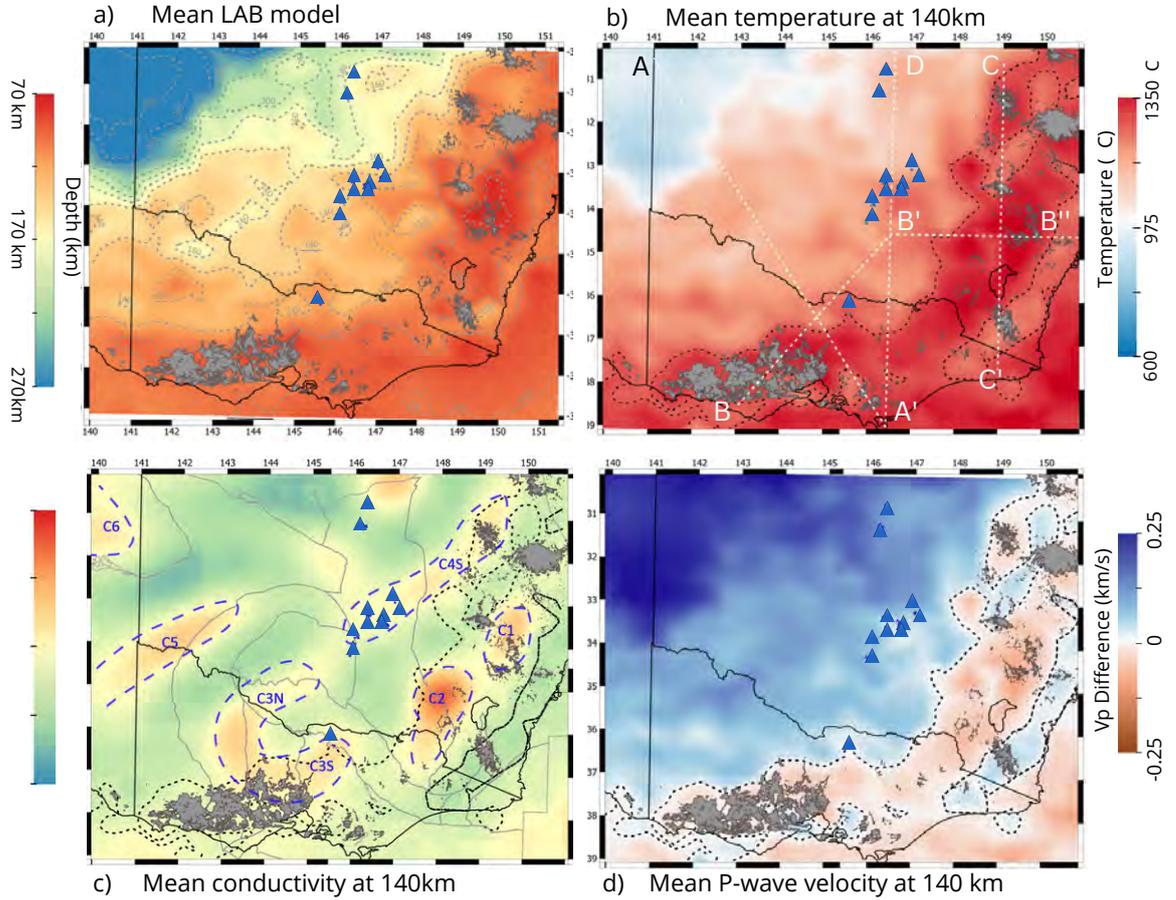


Figure 8. (a) Mean LAB depth. Contours of the LAB depth every 20 km are shown in grey-dashed line. Mean models at 140 km of (b) temperature (c) electrical conductivity and (d) P-wave velocity relative to 1-D reference model AusREM at $34.4^{\circ}S, 145^{\circ}E$. The $1250^{\circ}C$ -contour (corresponding with the thermal LAB) is plotted in dashed-black in (b-d). Panel (c) shows the location of the geological provinces and conductors in dashed blue. The location of leucite volcanoes are shown in blue triangles and the surface outcrop of basaltic volcanics are shown in grey in all panels. Panel (b) shows five transects which are discussed in section 5.1.

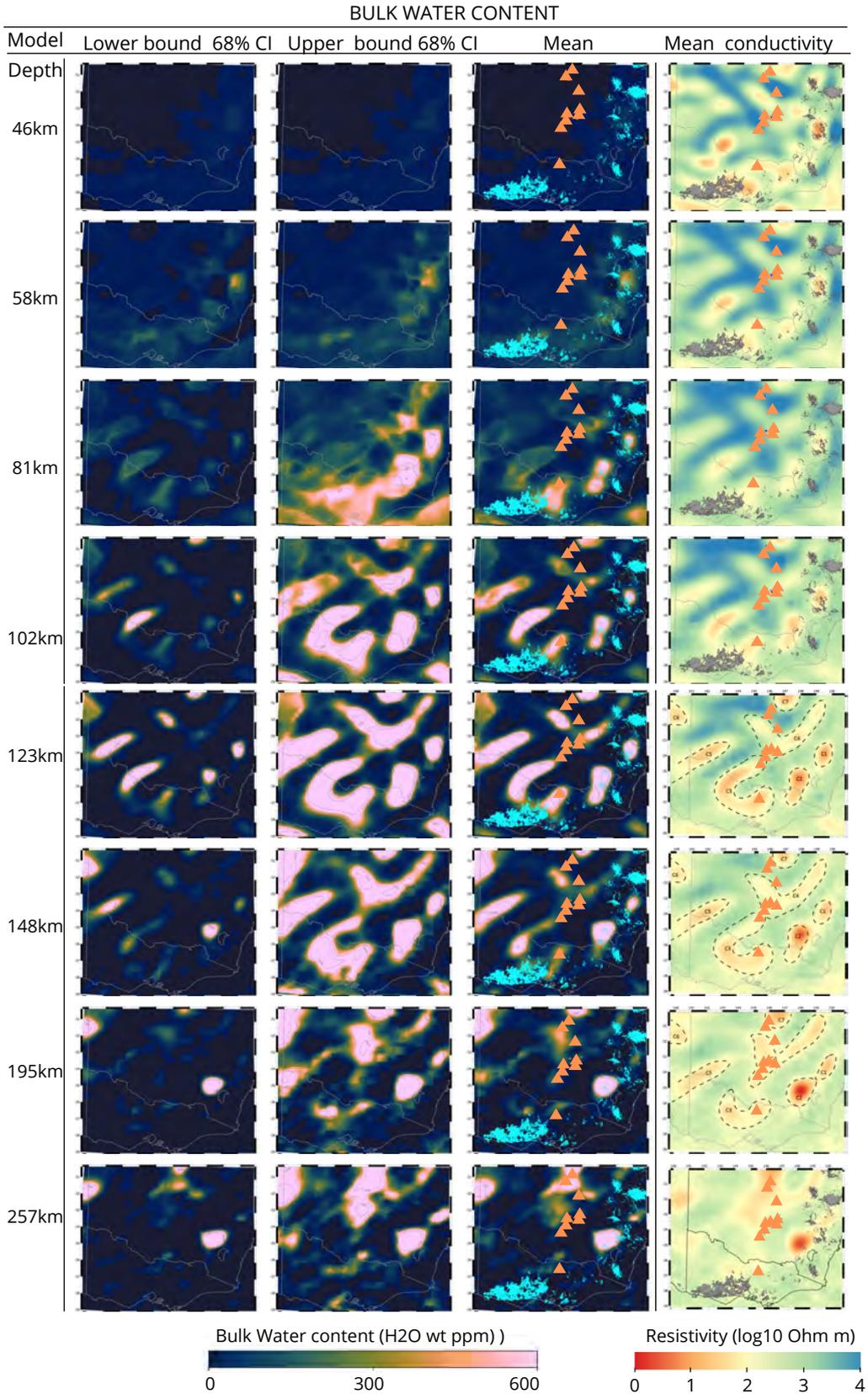


Figure 9. Bulk water content and mantle conductivity models from the joint probabilistic inversion. Columns (1)-(3): water content maps obtained from the (1) the lower, (2) upper bound of the 68% CI and (3) mean conductivity models. Column (4): depth slices from mean conductivity models of the posterior PDF. The location of leucite-bearing volcanism and basaltic volcanoes are shown in orange and turquoise in (3); and blue and in grey in (4). Selected depths are shown on the left of the figure. We refer to the main text for an explanation of structures C1-C7. Selected depths are shown on the left of the figure.

508 the following structures (Figure 9): C_1 , C_2 and C_3 beneath the eastern basaltic volca-
 509 noes; C_4 below the central leucitites; C_5 on the eastern boundary of Delamerian Oro-
 510 gen; C_6 beneath CP; and the deep localized conductor C_7 at $\sim -30.5^\circ N, 147^\circ E$, beneath
 511 the northern leucitites.

512 Figures 10-11.I and Figure S8 (Supplementary material) show vertical slices of the
 513 conductivity, water content and temperature along the four transects depicted in Panel
 514 (b) of Figure 8. The transects in Figure 10.I cross most of the geological provinces on
 515 the west and demonstrate a striking correlation between known geological boundaries
 516 and the alternation between wet/dry portions of the lithosphere. The joint assessment
 517 of these transects clearly shows that the lithospheric mantle beneath CP (C_6) corresponds
 518 to a highly conductive, hydrated, and cold region. We observe a high-conductivity anomaly
 519 (C_5) below the Stawell Zone that crosses the LAB. While the high temperatures found
 520 in this region (T_2 in Figure 10.Ic) can partially explain its conductivity structure, Fig-
 521 ure 10.Ib indicates that a large part of this anomaly is related to metasomatism (or in-
 522 cipient melting?). The high-conductivities observed in region C_{3N} (beneath Tabberra-
 523 berra Zone) and the conductor C_{NV1} (at ~ 90 km depth beneath the NV) can be entirely
 524 explained by a relatively large water content. We observe that while the conductivity
 525 of C_{3S} at ~ 200 km can be explained by the high anomalous temperatures found in that
 526 region (T_1), a substantial part of its conductivity at ~ 150 km is explained by the pres-
 527 ence of water (or melt?). Two shallow conductors C_{NV2} and C_{NV3} are found at ~ 20 -
 528 75 km depth beneath the NV.

529 Along the transects in Figure 10.II, the LAB shows a small perturbation over a large
 530 conductive anomaly at ~ 50 km and two defined steps at ~ 300 and ~ 750 km. These
 531 features correlate with the location of the northern leucitites, central leucitites and basaltic
 532 volcanoes, respectively. They also correlate with the location of high-conductivity regions
 533 (C_7 , C_4 and C_{3S}) in the sub-lithosphere. Comparing the conductivity and bulk water
 534 content along this transect, we observe that while water contents at C_7 and C_4 are rel-
 535 atively large, that of C_{3S} is considerably lower. A series of crustal conductors are also
 536 observed beneath both the leucitites and basaltic volcanoes.

537 Figure 11.I, which transects across the eastern basaltic volcanoes and the NV, shows
 538 the continuation of C_{3S} , C_{NV2} and T_1 beneath the NV. A high-conductivity and wet re-
 539 gion (C_1) is observed in the sub-lithospheric mantle below the Eastern Province. This
 540 deep, wet structure is also seen in Figure S8 and correlates with a high-temperature anomaly
 541 (T_3). A shallow semi-hydrated structure (C_{EP1}) is observed right below the basaltic vol-
 542 canoes. The relationship between these features is hard to reconcile unless we relax the
 543 assumption that the entire conductivity anomaly over the background is due purely to
 544 water content. We discuss this further in the next section.

545 5 Discussion

546 5.1 Mantle metasomatism and volcanism in southeast Australia

547 Mantle metasomatism occurs when incipient melts or fluids react with mantle rocks
 548 (predominantly peridotite). These reactions can i) affect the modal proportions of peri-
 549 dotites, ii) introduce new volatile-bearing phases (phlogopite, amphibole, apatite, and
 550 carbonates) and, in some pervasive cases, iii) create new lithological domains, such as
 551 pyroxenite \pm volatile-bearing phase lithologies (e.g. O'Reilly & Griffin, 1987). The gen-
 552 eration of volatile-bearing phases reduces the solidus temperature (Wallace & Green, 1988;
 553 Foley et al., 2009; Pintér et al., 2021) and increases the electrical conductivity of the man-
 554 tle domain (Selway, 2014).

555 The bulk water content we report in this study acts as a general proxy for meta-
 556 somatism or mantle fertility, i.e., the inclusion of phases (metasomes) that increase the
 557 electrical conductivity of the mantle. Therefore, as mentioned above, this proxy lumps

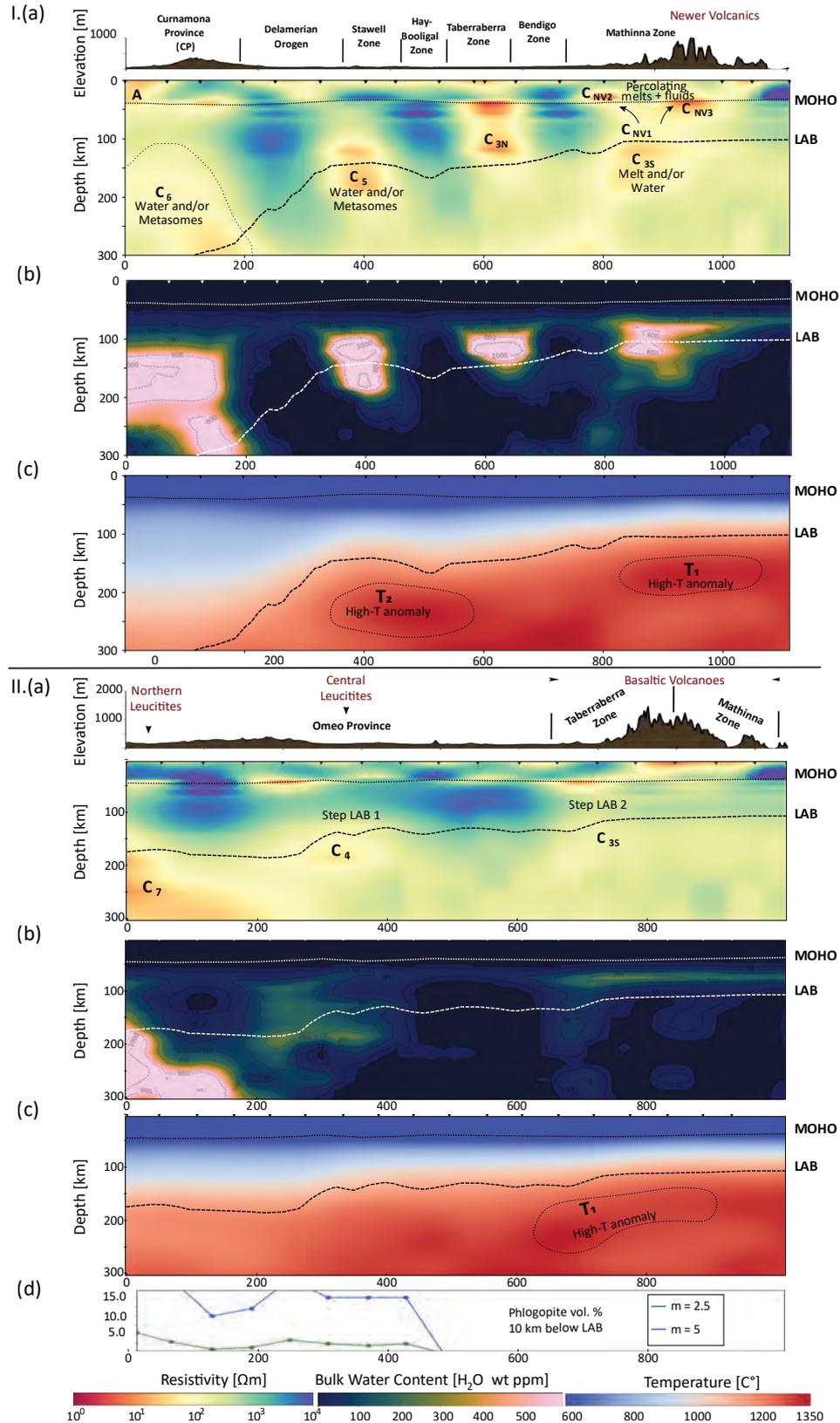


Figure 10. I. Vertical slices along transect A-A' (crossing most of the geological provinces) of a) the mean conductivity model, b) bulk water content derived from the mean conductivity mode and c) temperature. II. Vertical slices along transect D-A' (crossing the leucite volcanics and basaltic volcanics in the south) of a) the mean conductivity model, b) bulk water content derived from the mean conductivity model and c) temperature. d) Intermediately connected ($m=2.5$ in blue) and poorly connected ($m=5$ in green) phlogopite in a dry lherzolitic matrix that fit the observed conductivities along the transect. The Moho and LAB depths along that transects are shown in dashed lines in all panels. The elevation and location of the geological provinces is shown at the top of the figures.

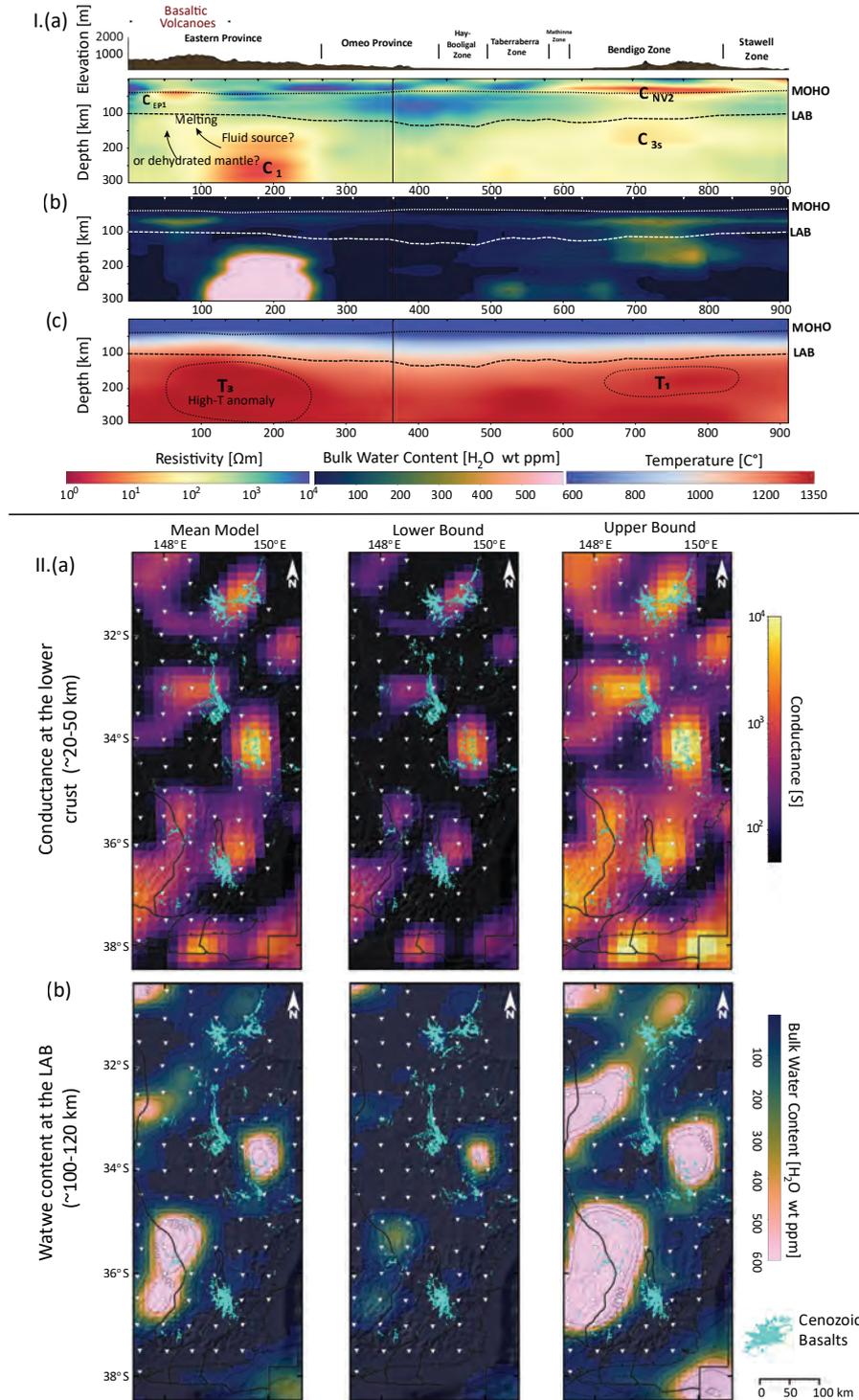


Figure 11. I. Vertical slices along transect B''-B'-B (across the basaltic eastern volcanoes and the NV) of a) the mean conductivity model, b) bulk water content derived from the mean conductivity model and c) temperature. The Moho and LAB depths along that transect are shown in dashed lines in all panels. The LAB depth along that transect is shown in dashed-black line in all panels. The elevation and location of the volcanics and the NV are shown at the top of the figure. II. Relationship between Cenozoic eastern volcanics and the parameters derived from the electrical conductivity model: (a) Conductance of the lower-crust ($\sim 20 - 50$ km), (b) water content calculated around the LAB depth ($\sim 100 - 120$ km).

558 together a number of factors not explicitly modeled in this work, for example, i) the pres-
 559 ence of phases such as graphite or sulphides for depths above 75-120 km (Selway, 2014;
 560 Özaydın & Selway, 2020); ii) co-existing water and phlogopite in cold mantle below 75-
 561 120 km depth; or iii) presence of melt in regions of elevated temperatures.

562 The results of Section 4.4 indicate widespread mantle metasomatism in southeast-
 563 ern Australia (Figures 9-11.I). Clear correlations are observed between the location of
 564 volcanic centers and regions of metasomatized mantle and conductive crust. In partic-
 565 ular, these regions are i) C_{3S} , C_{NV1} , C_{NV2} and C_{NV3} beneath Newer Volcanics; ii) C_1
 566 and C_{EP1} beneath the Eastern Volcanics (also C_4 , C_{EP2} and C_{EP3} Figure S8); iii) C_4
 567 beneath Central Leucitites; and iv) C_7 below Northern Leucitites. Xenoliths entrained
 568 in these lavas show strong evidence for metasomatism in the source (Yaxley et al., 1991,
 569 1998; Shea et al., 2022), further validating the presence of metasomatic agents in these
 570 regions (Frey et al., 1978). Other metasomatized mantle regions show a strong link with
 571 subduction/accretion processes rather than volcanism. These are C_6 , C_5 and C_{3N} be-
 572 neath CP, Stawell Zone and Taberraberra Zone, respectively.

573 One of the key benefit of our inversion is that we can dissociate the effects of the
 574 temperature from other factors controlling the conductivity structures and, ultimately,
 575 map anomalies associated with mantle metasomatism. While the current state of our method-
 576 ology does not allow us to discriminate the different metasomatic factors (presence of
 577 melt, water or phlogopite, for example), these can be inferred via the joint assessment
 578 of the conductivity, temperature and metasomatism models in different regions. Further
 579 analysis with contributions from melt modeling may be required to understand the full-
 580 scope behind metasomatism and the genesis of melts in southeast Australia. This work
 581 is left for a forthcoming publication.

582 **5.1.1 Basaltic volcanoes**

583 The existence of shallow mantle upwellings and thin lithosphere provides a favor-
 584 able setting for extensive decompression melting and mantle metasomatism (Aivazpour-
 585 porgou et al., 2015). The LAB is very shallow in the Newer Volcanics (NV) and East-
 586 ern Volcanics (EV), allowing mantle upwellings to reach depths at which decompression
 587 melting of peridotitic rocks (and the resulting basaltic primitive melts) is possible. In
 588 this context, we also observe that the metasomatized regions C_{3S} and C_1 correlate with
 589 the location of sub-lithospheric high-temperature anomalies (T_1 and T_3 , respectively).
 590 From the low-velocities observed at 60-80 km beneath the NV (Figure 4), Rawlinson et
 591 al. (2017) interpreted T_1 as a mantle upwelling and the source of the NV (see also Rawl-
 592 inson et al., 2015). Similarly, Rawlinson et al. (2015) interpreted the low-velocity anomaly
 593 corresponding with T_3 (Figures 4) as a deep mantle source for the EV. All of the above
 594 point to ideal conditions for the generation of extensive basaltic magmatism in this re-
 595 gion.

596 Figures 10.I and 11.I show a clear conductive pathway from C_{3S} to C_{NV1} - C_{NV3}
 597 and from C_1 to C_{EP1} , while the mantle beneath the volcanics is relatively dry. To fur-
 598 ther illustrate the relationship between the location of volcanics, shallow conductors and
 599 mantle metasomatism, Figure 11.II shows the average conductance at 20-50 km depth
 600 and average water content near the LAB beneath the EV. According to these results,
 601 basalt fields tend to associate with shallow conductors (Figure 11.IIa) and dry mantle
 602 (Figure 11.IIb). This relationship indicates that metasomatic agents percolated from deep
 603 mantle sources and traversed towards the crust. On their ascent, they precipitated con-
 604 ductive minerals forming shallow conductors (C_{NV2} , C_{NV3} , C_{EP1} , C_{EP2} , C_{EP3} in Fig-
 605 ures 10.I, 11.I and S8). We also note that basalts tend to be located in the surround-
 606 ings of the most metasomatized regions of the lithosphere rather than on top of them;
 607 something that has been observed also in kimberlites worldwide (Özaydın & Selway, 2022).

608 The dry mantle beneath the basaltic fields suggests that the melting events exhausted
609 the mantle source in the original metasomes and left behind drier residues.

610 **5.1.2 Leucitite volcanoes**

611 The leucitite lavas have melt compositions comparable to lamproites, and were de-
612 rived from an atypical mantle assemblage of phlogopite-bearing pyroxenite (Shea et al.,
613 2022; Foley et al., 2022). Due to these lava compositions, Kirkby et al. (2020) interpreted
614 the conductors beneath the central leucitites as regions of metasomatized mantle with
615 hydrous minerals such as phlogopite. Given the high conductivities ($< 100\Omega\text{ m}$) and bulk
616 water content ($\sim 200\text{ ppm}$) observed around C_4 (Figures 10.II), our results indicate a
617 high probability for the presence of volatile-bearing minerals, supporting the above in-
618 terpretation.

619 Using the water calculation setup described in Section 3.6 and the phlogopite con-
620 ductivity model of Li et al. (2017), we calculated the electrical conductivities of lherzolite
621 with 5 and 10 % vol. of 0.52 w.t. fluorine-bearing phlogopite (average fluorine value
622 of mantle rocks, Özaydın et al., 2022) for both perfectly connected ($m = 1.1$, modified
623 Archie’s law) and sparsely populated phlogopites ($m = 6$, modified Archie’s law). The
624 results show that perfectly connected cases are 2.5 orders of magnitude more conduc-
625 tive than the observed conductivities in the region, while the conductivity for sparsely
626 populated/disconnected cases lay near the lower bound of the observed conductivities
627 (Figure S9 in Supplementary Material). These results suggest that a lherzolite with 5-
628 10 % vol of partially connected ($6 < m < 1.1$) phlogopite explains the conductivities
629 in C_4 . We interpret the high conductivities in C_4 as a phlogopite-bearing lherzolite with
630 small percentages of partially to sparsely connected phlogopite. These results are also
631 illustrated in Figure 10.II.d, which shows the percentage of partially connected ($m=2.5$)
632 phlogopite that can explain the observed conductivities beneath the leucitites volcanoes.
633 We note that large amounts of phlogopite would drastically lower the seismic velocities
634 (e.g., Selway et al., 2015), something not seen in our models. We therefore reject the hy-
635 pothesis of large percentages ($\geq 15\%$) of poorly connected ($m=5$) phlogopites in this
636 region.

637 The northern leucitites sit above a high-conductivity and metasomatized region be-
638 low the LAB (C_7). The ultrapotassic compositions of these lavas suggest low-degree par-
639 tial melting (Cundari, 1973), which is consistent with the relatively colder temperatures
640 found in the region. Furthermore, potassium-rich magmas are produced by melting of
641 a metasomatized mantle that has been enriched in phlogopite (Xu et al., 2017; Förster
642 et al., 2019). We calculated the effect of phlogopites in this region and found similar re-
643 sults to those for the central leucitites (Figures S9), favoring the scenario with partially
644 connected phlogopites. Compared to the central leucitites, the higher conductivities and
645 colder temperature in this region provide favorable conditions for the presence of exist-
646 ing phlogopites that survived previous melting events.

647 **5.1.3 Curnamona Province**

648 Figure 10.I.a-b shows a successive alternation of conductive/wet and resistive/dry
649 lithospheric domains that resemble the west-to-east subduction-accretion process in east-
650 ern Australia (Glen, 2005; Shea et al., 2022). The joint assessment of the fields depicted
651 in Figure 10.I also suggests that the region C_6 experienced pervasive mantle metasoma-
652 tism. Given the cold temperatures, the lack of present magmatism, and the geological
653 history of southeastern Australia, the metasomatic events C_6 are likely related to accre-
654 tion processes whereby successive subduction and orogenic events introduced metaso-
655 matic agents into the mantle. Overtime, this process preferentially metasomatized the
656 old, thick lithosphere beneath Curnamona province.

657 The crustal conductors (A) below CP (described in Section 4.3.1) can be seen in
 658 Figure 10.Ia. Kay et al. (2022) interpreted these shallow conductors as the deposition
 659 of interconnected graphite. However, given that graphite films are not stable at shallow
 660 depths (Zhang & Yoshino, 2017; Yoshino et al., 2018), we observe that a more feasible
 661 explanation for these crustal conductors are carbon-rich fluids sourced from the deep meta-
 662 somatized region C_6 (Thibault et al., 1992).

663 5.2 Implications for magma generation beneath eastern Australia

664 Age-progressive volcanism in the EAVP, particularly along the Cosgrove track (Davies
 665 et al., 2015), has been widely attributed to long-lived mantle plume activity (Wellman
 666 & McDougall, 1974). However, Shea et al. (2022) has recently argued that primitive melt
 667 compositions throughout the EAVP, including all or most of the age-progressive volcan-
 668 ism, were produced by melting of a metasomatized mantle source at temperatures lower
 669 than those expected in a deep mantle plume. The petrological and geochemical evidence
 670 summarized by Shea et al. (2022) and in the previous sections indicate that primitive
 671 melts, particularly those associated with the leucitites, originated from a pyroxenitic com-
 672 ponent rather than from a peridotitic mantle lithology. Since the solidi of pyroxenites
 673 (Foley et al., 2022) is substantially lower than that of anhydrous peridotites (Walter, 1998),
 674 partial melting of such lithologies is possible with only slight perturbations above am-
 675 bient upper mantle temperatures. The lack of picrites within the EAVP, which are prod-
 676 ucts of high-degrees of partial melting of peridotite, also suggest low-temperature, small
 677 degrees of partial melting dominated volcanism in the EAVP.

678 Our results show a series of steps in the LAB that correlate well with both the lo-
 679 cation of basaltic and leucitite volcanoes (Figures 10.II and 12). Thermomechanical mod-
 680 els have shown that such steps in the LAB constitute areas prone to generating sublitho-
 681 spheric small-scale, edge-driven convection (EDC) instabilities and partial melting (e.g.,
 682 Zlotnik et al., 2008; Van Wijk et al., 2010; Davies & Rawlinson, 2014; Ballmer et al., 2011;
 683 Afonso et al., 2008; Duvernay et al., 2021a). In particular, Duvernay et al. (2021b) showed
 684 that shear-driven upwelling (SDU) and EDC processes that account for water content
 685 in the upper mantle can produce enough melt to explain the total melt volume in the
 686 EAVP.

687 Considering all of the above, a plausible model arises for the melt generation in south-
 688 eastern Australia, which we summarize in Figure 12. The relative motion of the Aus-
 689 tralian plate and the asthenospheric mantle created a favourable asthenospheric flow (pos-
 690 sibly towards the south-southwest) which allowed the generation of EDC/SDU cells be-
 691 neath steps in the LAB (e.g. above C_4 and C_{3S}). Such a convective flow can detach meta-
 692 somatized lithologies from the lowermost portions of the lithosphere (on the thick side
 693 of the step) and drag them into the upwelling limbs of the EDC/SDU convection cell,
 694 where they would preferentially melt (given their lower solidus) and create the primary
 695 metasomatizing melts. The latter would subsequently react with the lower portions of
 696 the lithosphere (in the thin side of the LAB step) and create the source of the leucitite
 697 volcanoes. This scenario is similar to the one presented by Shea et al. (2022) and sup-
 698 ported by our results and by the abundant petrological evidence summarized by these
 699 authors.

700 If EDC/SDU cells are the main driver for the leucitite volcanism, they need to be
 701 able to produce melts for a relatively short period of time (non regenerative) to explain
 702 the punctuated nature of this volcanism. Indeed, numerical simulations clearly show that
 703 ED/SD instabilities are ephemeral in nature and can be hindered by small perturbations
 704 in the sublithospheric flow (e.g., King & Anderson, 1995; Duvernay et al., 2021b, 2022).
 705 The presence of an anomalously hot upwelling or plume can also enhance or shut down
 706 melting near lithospheric steps (Mather et al., 2020; Duvernay et al., 2022), making it
 707 difficult to separate the two mechanisms without further knowledge on the underlying

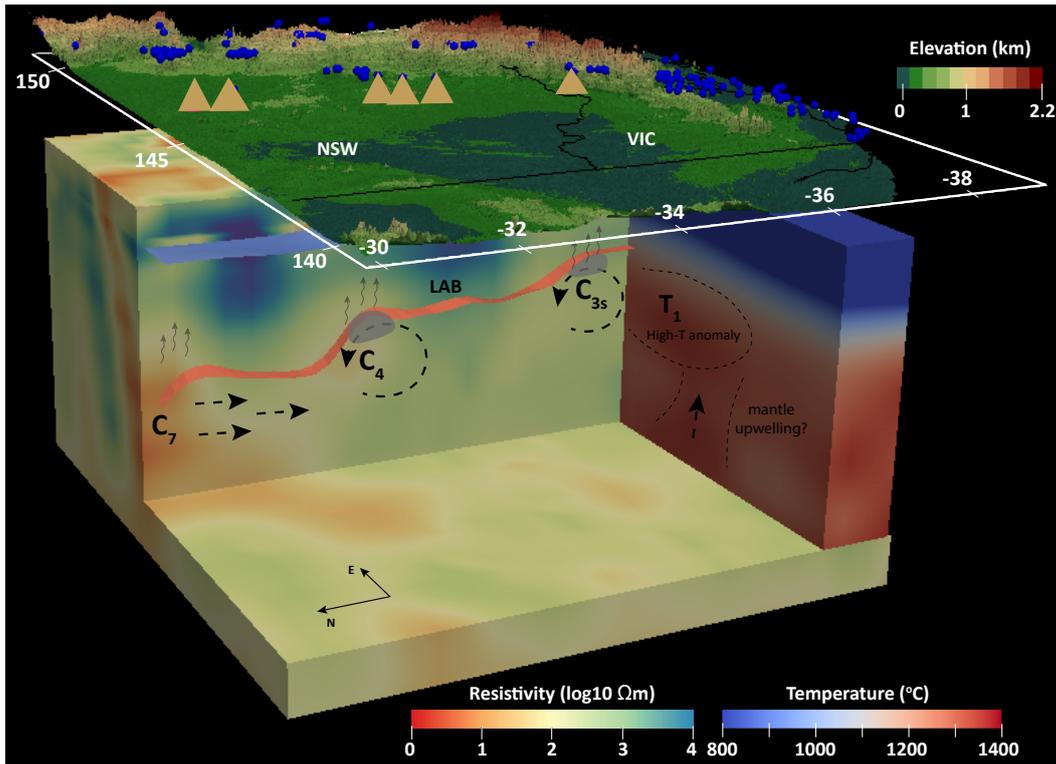


Figure 12. 3D rendering views of the LAB (red surface), mean conductivity and temperature models depicting interactions between mantle metasomatism and steps in the LAB. Dashed black arrows show the flow of the asthenosphere and shearing of enriched mantle material into EDC-cells (circular dashed lines). Local hotspots along the LAB where enriched mantle material crosses its solidus in the sub-lithospheric mantle are indicated by grey blobs. Grey arrows indicate incipient melts that may travel across the LAB from deep to shallow portions of the lithospheric mantle. Leucitites and basaltic volcanic centers are shown in yellow triangles and blue circles, respectively. The Moho is shown in blue over the conductivity model. The figure includes the topography in southeast Australia.

708 mantle dynamics. Thus, while our results provide a firm ground for interpretations on
 709 the origin of the EAVP and the roles that lithospheric structure and composition played,
 710 we cannot assess the deep mantle processes that controlled the large-scale sublithospheric
 711 flow.

712 6 Conclusions

713 We performed a joint probabilistic inversion of 3D magnetotellurics (MT) and seis-
 714 mic velocity data to constrain the lithospheric structure, metasomatic domains and melt-
 715 ing processes in southeast Australia. Our methodology minimizes the non-uniqueness
 716 of the MT problem and provides quantitative information on model uncertainties via full
 717 posterior distributions. This information is crucial for assigning meaningful interpreta-
 718 tions to electrical conductivity anomalies in terms of temperature versus metasomatism/compositional
 719 anomalies.

720 We image a highly heterogeneous lithosphere beneath eastern Australia that we
 721 link to geodynamic and tectono-magmatic processes across multiple scales. In particu-

lar, we detect widespread, but highly irregular mantle metasomatism throughout the region, pointing to complex interactions in the asthenosphere-lithosphere system. We also image alternating conductive/wet and resistive/dry lithospheric domains that correlate with the location of major geological provinces, resembling the west-to-east subduction-accretion process that formed eastern Australia. A series of steps in the present-day thermal structure correlate with the location of intra-plate volcanic centers and moderate thermal anomalies in the sublithospheric mantle, suggesting a genetic link between these three features. Basaltic volcanism is preferentially located in regions of very thin lithosphere and dry mantle, whereas leucitite volcanoes are located in regions of highly metasomatized mantle, intermediate lithospheric thickness and localized lithospheric steps. These results, together with recent petrological and geochemical evidence for relatively low temperatures in the melting region (Shea et al., 2022), suggest that the interaction between a complex, metasomatized lithospheric structure and localized mantle upwellings (e.g. via edge-driven convection or focusing of moderately hot mantle upwellings) are likely responsible for much of the volcanism in the EAVP, rather than a deep, hot mantle plume.

Lastly, it is generally accepted that metasomatized lithospheric mantle plays a critical role in the generation of major ore deposits. The ability to map metasomatized mantle domains paves the way to a new way of exploring for mineral resources.

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