

1 **Synthesizing Sea Surface Temperature and Satellite**
2 **Altimetry Observations Using Deep Learning Improves**
3 **the Accuracy and Resolution of Gridded Sea Surface**
4 **Height Anomalies**

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9 **Key Points:**

- 10 • This paper is a non-peer reviewed preprint submitted to EarthArXiv and has been
11 submitted for publication to the Journal of Advances in Modeling Earth Systems
12 (JAMES) for peer review. Subsequent versions of this manuscript may have slightly
13 different content.
- 14 • We developed a deep learning method that significantly improves the accuracy and
15 resolution of gridded sea surface height anomalies
- 16 • This data-driven method takes advantage of combining sea surface temperature
17 and altimetry observations
- 18 • The inferred surface geostrophic currents are quantitatively and qualitatively more
19 realistic than those from existing sea surface height maps

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Abstract

Gridded sea surface height (SSH) maps estimated from satellite altimetry are widely used for estimating surface ocean geostrophic currents. Satellite altimeters observe SSH along one-dimensional tracks widely spaced in space and time, making accurately reconstructing the two-dimensional (2D) SSH field challenging. Traditionally, SSH is mapped using optimal interpolation (OI). However, OI artificially smooths the SSH field leading to high mapping errors in regions with rapidly-evolving mesoscale features such as western boundary currents. Motivated by the dynamical relation between SSH and sea surface temperature (SST) and the notion that even the chaotic evolution of mesoscale ocean turbulence may contain repeating patterns, we outline a deep learning (DL) approach where a neural network is trained to reconstruct 2D SSH by synthesizing altimetry and SST observations. In the Gulf Stream Extension region, dominated by mesoscale variability, our DL method substantially improves the SSH reconstruction compared to existing methods. Our SSH map has 17% lower root-mean-square error and resolves spatial scales 30% smaller than OI compared against independent altimeter observations. Surface geostrophic currents calculated from our map are closer to surface drifter observations and appear qualitatively more realistic, with stronger currents, a clearer separation between the Gulf Stream and neighboring eddies, and the appearance of smaller coherent eddies missed by other methods. Our map yields significant re-estimations of important dynamical quantities such as eddy kinetic energy, vorticity, and strain rate. Applying our DL method to produce a global SSH product may provide a more accurate and higher resolution product for studying mesoscale ocean turbulence.

Plain Language Summary

Satellites observe small variations in the height of the sea surface but with large gaps in the observations. Having an estimate of the two-dimensional sea surface height field allows one to estimate surface ocean currents, so filling in the gaps between the observations is an important problem. The traditionally-used method for filling in the gaps between sea surface height observations struggles when there are lots of small-scale, rapidly-interacting ocean currents. We developed a deep learning model to estimate the sea surface height field more accurately. We achieved this by combining the sea surface height observations with satellite observations of sea surface temperature. The relationship between sea surface temperature and height is non-trivial, but our deep learning model learned to use information from the sea surface temperature observations in the places where sea surface height wasn't observed to improve the accuracy of the sea surface height estimate. We applied and tested our method in the Gulf Stream and demonstrated that our sea surface height map is more accurate than that from the traditional method. Our method provides a more accurate sea surface height map which could allow us in future to learn new lessons about small-scale surface currents in the ocean.

1 Introduction**1.1 Background**

Sea surface height (SSH) maps - i.e. time-varying gridded maps of the height of the ocean's surface referenced against the geoid - derived from satellite altimetry observations have been a crucial tool for physical oceanographers for many years (Fu et al., 2010; Abdalla et al., 2021). Under the assumption of geostrophy, SSH is proportional to a streamfunction for the surface currents. Hence a SSH map provides easy access to an estimated map of surface geostrophic ocean currents, which would otherwise be challenging to obtain through direct in-situ observations with global coverage. SSH maps have been used to directly observe the inverse kinetic energy cascade in the global oceans (Scott & Wang, 2005), to track and study the behavior of mesoscale eddies (Chelton et al., 2011; Fu et al., 2010), and to estimate surface eddy diffusivities (Abernathey & Marshall, 2013).

To date, the only global SSH observations have come from satellite-borne nadir radar altimeters, which measure SSH along one-dimensional tracks directly beneath the satellite’s orbit resolving along-track scales of O(20-100km) (Dufau et al., 2016). Reconstructing the full two-dimensional (2D) SSH field from altimeter observations requires significant interpolation in space and time. Hence, the length-scales 2D products can accurately resolve are coarser than the along-track resolution of the altimeters (Ballarotta et al., 2019).

The most widely used 2D SSH product is that generated by the ‘Data Unification and Altimeter Combination System’ (DUACS) (Taburet et al., 2019) and distributed by the ‘Copernicus Marine Environment Monitoring Service’ (CMEMS) (*note that this is the same product as was formerly distributed by the ‘Archiving, Validation and Interpretation of Satellite Oceanographic data’ (AVISO) service*). This product is created using the optimal interpolation (OI) method (Bretherton et al., 1976; Le Traon et al., 1998), which is otherwise known as ‘objective mapping’ or ‘objective analysis’. OI provides the best least squares linear estimator of the SSH in the gaps between observations, given an a priori model for the covariance of SSH between different locations and times, and knowledge of the instrument noise covariance. In the DUACS mapping, the SSH covariance is assumed to be a single-scale Gaussian with prescribed decorrelation length- and time-scales that have been tuned empirically and the instrument noise covariance has been chosen to be suitable for the constellation of satellite altimeters (Taburet et al., 2019).

While the DUACS SSH product continues to be valuable to the oceanography community, the OI method has been shown to introduce significant deficiencies. Amores et al. (2018) used an Observer System Simulation Experiment (OSSE) in which a DUACS-like SSH product was generated using pseudo-observations from the output of a high-resolution ocean model to study how accurately DUACS captured the model’s mesoscale SSH field. They demonstrated that the eddy fields inferred from the DUACS SSH map are significantly distorted, with often multiple smaller eddies being aliased into larger eddies. The real-world accuracy of an SSH mapping method can also be studied by generating a map with one satellite altimeter’s observations withheld and examining the errors of the mapped SSH compared to these independent along-track observations. In this way, Ballarotta et al. (2019) showed that the DUACS product only accurately resolves SSH signals down to time-scales of O(30 days) and length-scales O(100km) at high latitudes, increasing to O(800km) in the tropics. OI tends to smooth out small-scale, fast-evolving features and strong peaks/troughs in SSH where observations are scarce and the oceanic mesoscale is energetic.

The surface currents derived from the DUACS SSH product are used extensively in oceanographic studies, so it is essential to ensure these inferred currents are as accurate and high resolution as possible using the existing satellite observing capabilities. There is therefore increasing focus on developing better methods for reconstructing the 2D SSH field from satellite altimetry observations (Ubelmann et al., 2015, 2021; Manucharyan et al., 2021; Fablet, Amar, et al., 2021; Le Guillou et al., 2021). Efforts to improve SSH mapping methods have been aided by the the creation in recent years of community-maintained ‘Ocean Data Challenges’¹. By establishing a common set of evaluation metrics, a common study region and time, and by sharing code and results these data challenges allow for direct comparison between different SSH mapping methods.

OI can only provide the statistical best *linear* estimator of the unobserved SSH for a given a priori covariance model. While the decorrelation length- and time- scales used in this covariance model can be tuned to best fit the data, as is done to create the DUACS product, a single-scale Gaussian covariance model and its resulting linear estimator is ultimately limited in its ability to accurately represent the non-linear dynamics

¹ <https://github.com/ocean-data-challenges>

120 of mesoscale ocean turbulence. An explicitly dynamics-based approach to SSH mapping
 121 would be to use a data assimilation framework to constrain an ocean circulation model
 122 to best match the available observations, then use this constrained model’s SSH field as
 123 the estimate of 2D SSH. However, such an approach requires additional observations to
 124 accurately constrain other essential model variables such as subsurface flow and density
 125 which are not typically available. Thus, only very idealized dynamical models, such as
 126 a one layer quasi-geostrophic (QG) model, have been effectively applied to SSH map-
 127 ping (Le Guillou et al., 2021). Such an idealized model is unable to capture many of the
 128 dynamical effects observed in the real-world ocean.

129 The recent advances in the application of deep learning (DL) methods to the earth
 130 sciences (Sonnewald et al., 2021; Sun et al., 2022) have inspired an increasing focus on
 131 the possibility of using data-driven interpolation methods to improve SSH mapping (Beauchamp
 132 et al., 2020; Manucharyan et al., 2021). DL models can be trained to approximate highly
 133 non-linear mappings from inputs to outputs. Thus it is plausible that a DL model could
 134 be trained to recognize dynamical signatures in partial SSH observations and use these
 135 to estimate SSH in unobserved regions given sufficient training examples. By training
 136 a DL model on a large library of examples of real ocean turbulence the DL model would
 137 be able to implicitly use ocean dynamics to create a more accurate SSH mapping than
 138 would be possible with OI.

139 1.2 Previous work applying DL to SSH mapping

140 A number of attempts to use DL for SSH mapping have been published recently
 141 (Fablet, Amar, et al., 2021; Fablet, Chapron, et al., 2021; Barth et al., 2022; Buongiorno Nardelli
 142 et al., 2022). These studies have mostly focused on the OSSE setting, where the SSH
 143 observations are generated from a circulation model in which the full 2D SSH field is known.
 144 Fablet, Amar, et al. (2021) and Fablet, Chapron, et al. (2021) developed a DL approach
 145 (4DVarNet) for interpolating SSH which formulates SSH mapping as a four-dimensional
 146 variational data assimilation (4DVar) problem but replaces the dynamical model and solver
 147 with trainable neural networks. They demonstrated that 4DVarNet yields a significant
 148 improvement over OI when tested in idealized OSSEs, though the improvement is more
 149 modest when applied to real-world observations.

150 While these results are auspicious, OSSEs can only offer a limited approximation
 151 of the real-world SSH interpolation problem since the full 2D SSH ground truth is un-
 152 known in real-world settings. It therefore cannot be used to calculate the reconstruction
 153 error to be minimized during training. Barth et al. (2020) introduced a Convolutional
 154 Neural Network (CNN) method, Data INterpolating Convolutional Auto-Encoder (DIN-
 155 CAE), for reconstructing sea surface temperature (SST) from partial satellite observa-
 156 tions. In a subsequent paper, Barth et al. (2022) refined and applied the method to re-
 157 construct SSH in the Mediterranean from partial real-world SSH and SST observations.
 158 They found that including SST yielded a lower reconstruction error than SSH observa-
 159 tions alone. However, an error comparison to the operational DUACS method was not
 160 made, so this method’s performance is hard to compare to other SSH mapping meth-
 161 ods. Buongiorno Nardelli et al. (2022) developed another CNN approach for reconstruct-
 162 ing and super-resolving SSH from partial SSH and SST satellite observations in the Mediter-
 163 ranean with promising results. However, to split the observations into independent train-
 164 ing and testing datasets they randomly selected days to withhold for testing, meaning
 165 adjacent days could be used for training and testing. SSH and SST fields contain sig-
 166 nificant temporal autocorrelation over time scales of days to weeks. Hence, such a sam-
 167 pling strategy leads to a testing dataset that could be highly correlated to the training
 168 data, calling into question the independence of the performance metrics calculated.

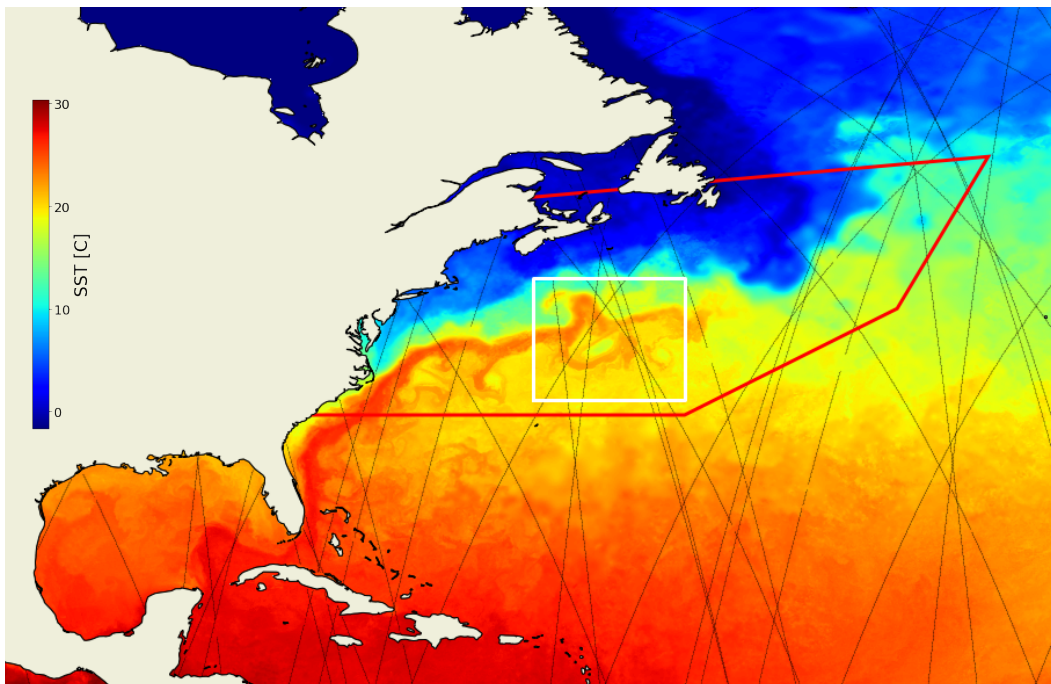


Figure 1. The typical data availability for one day (2017-04-11) during our testing period: the color map shows the GHRSSST MUR gridded SST, and the black tracks indicate the locations of satellite altimeter observations from all satellites available on that date (CryoSat-2, Jason-2, Jason-3, Haiyang-2A, Sentinel-3A, and SARAL/AltiKa). As discussed in section 5, we draw training examples from within the red polygon (the Gulf Stream Extension), and test our method in the white box (where other mapping methods have been applied so we can compare to their accuracy).

169 To date, DL SSH mapping studies remain limited to theoretical OSSE studies or
 170 limited regional experiments and the gridded SSH product of choice for oceanographers
 171 remains the OI-generated DUACS product.

172 **1.3 Problem statement and paper structure**

173 In this study, we seek to demonstrate in practice that DL can provide a more ac-
 174 curate SSH mapping with a more physically realistic eddy evolution than was previously
 175 possible. Further, our DL approach allows us to include co-located observations of other
 176 ocean surface variables in the input to the mapping, something that is challenging with
 177 existing SSH mapping methods. Specifically, we will demonstrate that the inclusion of
 178 SST observations as an additional source of information for the SSH mapping leads to
 179 dramatic improvements in the accuracy and resolution of the reconstructed SSH evolu-
 180 tion. In section 2, we provide a dynamical motivation for why SST observations are ex-
 181 pected to improve SSH mapping and provide the rationale for our DL approach. Our
 182 method is both trained on and tested against real-world satellite observations (the datasets
 183 we use are presented in section 3). We present our mapping method in section 4 and how
 184 it performed during training in section 5. In section 6, we present the results of a regional
 185 experiment to quantitatively compare our method’s accuracy to that of the conventional
 186 and several recently proposed SSH mapping methods in the Gulf Stream Extension. In
 187 section 7, we discuss the implications of these results and outline a road-map to creat-
 188 ing a more accurate global SSH product with DL.

2 Rationale for our proposed methodology

2.1 Dynamical motivation for supplementing SSH observations with SST

In addition to satellite altimetry observations, satellites routinely observe other surface ocean quantities, such as SST, sea surface salinity, and ocean color. The spatiotemporal evolution of these fields is certainly not independent from one another and from that of SSH, since they are all affected and mediated by the dynamics of the ocean (and by biological processes in the case of ocean color). Thus, if one of these other surface variables is observed in the gaps between satellite altimetry observations, and the relationship to SSH is known, then this can inform the mapping of the unobserved SSH.

In this study, we supplement satellite altimetry with SST observations to improve the SSH mapping for reasons both physical and practical. Physically, this is a natural choice since SSH and SST are known to have a close dynamical relationship. Indeed, in high eddy kinetic energy regions (e.g. western boundary currents in the northern hemisphere and the Antarctic Circumpolar Current in the southern hemisphere), temperature explains much of the density field, especially in winter (Isern-Fontanet et al., 2006), indicating that mesoscale SST anomalies can be interpreted as a dynamical part of surface-intensified baroclinic eddies (Smith & Vallis, 2001). Hausmann and Czaja (2012) confirmed these results: SSH and SST anomalies have similar spatial scales and are characterized by a westward shift as expected from baroclinic instability, which is the main source of mesoscale eddies. In addition, a close spectral relationship between mesoscale SSH and SST can be inferred when the large-scale meridional gradient of potential vorticity at depth is linearly related to the large-scale meridional SST gradient through a function only depth-dependent (Lapeyre & Klein, 2006), which is verified in high eddy kinetic energy regions (Lapeyre, 2009). This relationship makes use of the surface quasi-geostrophic (SQG) approximation in which the ocean interior dynamics at mesoscale is implicitly taken into account through an “effective” Brunt-Vaissala frequency, N_{eff} (Lapeyre & Klein, 2006; Isern-Fontanet et al., 2006; LaCasce & Mahadevan, 2006; Klein et al., 2019). The resulting relationship in spectral space is

$$\widehat{SSH}_{SQG}(\mathbf{k}) = -\frac{f_0}{N_{eff}|\mathbf{k}|}\alpha.\widehat{SST}(\mathbf{k}), \quad (1)$$

where a ‘hat’ symbol indicates a horizontal Fourier transform, \mathbf{k} the horizontal wavenumber vector, α the thermal expansion coefficient, and f_0 the Coriolis frequency. Isern-Fontanet et al. (2006) proposed to use this relationship to recover SSH and therefore surface ocean currents from SST observations. Practically, satellite observations of SST have higher spatial resolution and lower missing data rates than those of SSH, which emphasizes the pertinence of using SST observations because of their close dynamical relationship with SSH.

Applying equation 1 to directly reconstruct SSH from SST observations is challenging in that the oceanic conditions described above do not always apply directly in all seasons and ocean regions. Other attempts have used a prescribed analytical relation between SST and SSH like that in equation 1 to reconstruct the unobserved SSH field (Isern-Fontanet et al., 2014; González-Haro & Isern-Fontanet, 2014). All such methods face the challenge that the dynamical relationship between SSH and SST is non-trivial and is likely to change in space and time. Lapeyre (2009) showed that the SSH field is in some regions dominated by the SQG mode, making equation 1 applicable, whereas in others the first baroclinic mode dominates, meaning the SSH estimated from equation 1 would be inaccurate. Thus, methods that rely on the SQG framework (Isern-Fontanet et al., 2006, 2014; González-Haro & Isern-Fontanet, 2014) and methods that rely on a QG framework (Ubelmann et al., 2015, 2021; Le Guillou et al., 2021) will each be geographically and seasonally limited in their ability to reconstruct the SSH field by the pertinence of their respective dynamical assumptions.

2.2 Rationale for deep learning

DL models (specifically, neural networks) have been shown to be universal function approximators, given sufficient depth and training data (Hornik et al., 1989). By employing a DL approach for SSH mapping, we implicitly make two hypotheses. *(i)*: Available satellite observations contain sufficient dynamical information that there exists a function mapping from the observations to an estimate of the dynamical state of the ocean, and by extension to the 2D SSH field. *(ii)*: There exists sufficient training data that we can train a DL model to approximate this function accurately enough to provide better SSH maps than OI. Both are motivated by earlier work applying DL neural networks to spatiotemporal SSH interpolation in a two-layer quasi-geostrophic (QG) model given partial observations of the surface (Manucharyan et al., 2021). George et al. (2021) also showed a DL model could learn to infer subsurface dynamics, which will, in turn, influence SSH evolution, from snapshots of the surface SSH field. The overarching conclusion from the two aforementioned studies is that the surface SSH field contains implicit signatures of the subsurface ocean dynamics and that DL models are sufficiently expressive that they can learn to recognize and exploit these signatures to more accurately reconstruct unobserved ocean variables.

The advantage of a DL approach that combines the SSH and SST fields is that a sufficiently expressive DL model, given enough training data, can, in theory, learn a dynamical relationship of arbitrary complexity. Thus, a DL model could learn to exploit the SST observations in regions and times when the dynamics makes them pertinent, while not being restricted to assuming the SQG approximation is always appropriate. The promise of a DL approach that combines SSH and SST observations to reconstruct the unobserved SSH field more accurately was also noted recently in Fablet and Chapron (2022) and Fablet et al. (2022). However, to date, there are no operationally-used SSH reconstructions that utilize SST observations or DL, and most of the studies on the topic have been theoretical in nature.

3 Datasets

3.1 Sea surface height (SSH) observations

We use Level 3 1Hz along-track satellite altimetry SSH observations from 2010-2020 (distributed by CMEMS). This product consists of observations from multiple satellite altimetry missions merged, calibrated, and corrected for several geophysical effects, including barotropic tides and atmospheric effects. The complete steps for producing this product are described in Taburet et al. (2019) and references therein. Note that at the Level 3 stage of the observation processing chain, the observations have not yet been mapped to a 2D grid, as shown in Figure 1. These along-track observations have been shown to accurately resolve the SSH signal down to length-scales $O(20-100\text{km})$ with root-mean-square instrument noise of 1-4cm depending on the location, season, and satellite altimeter in question (Dufau et al., 2016).

3.2 Sea surface temperature (SST) observations

We use the Level 4 Multiscale Ultrahigh-Resolution (MUR) gridded SST product from the same years provided by the Group for High-Resolution SST (GHRSSST). The MUR SST product is generated by combining observations from multiple satellites and in-situ sensors before interpolating to a regular grid using OI. Clouds cause gaps in the highest resolution SST satellite sensors, meaning that the OI product has high resolution in cloud-free regions but suffers artificial smoothing in cloud-occluded regions. An alternative approach would be to use the high-resolution observations directly with gaps, or to use only the lower resolution cloud-free observations provided for example by the NASA Advanced Microwave Scanning Radiometer-EOS (AMSR-E) which would both

offer a uniform spatial resolution. Until a uniformly high-resolution gridded SST product is developed, there is an inevitable trade-off between using high-resolution SST with gaps, low-resolution SST without gaps, and an OI product like MUR with non-uniform spatial resolution. We here use the MUR product noting that even the along-track spatial scales that are resolved in altimetry observations are typically larger than the scales smoothed by the OI in the MUR SST product.

3.3 Data pre-processing: domain and coordinate system

We pursue a patch-based approach, where SSH is reconstructed in a local square domain using observations from only its surroundings. This approach is motivated by the fact that the dynamics governing SSH evolution is predominantly local in space and time, thus estimating the SSH at a point should only require observations from some finite space-time window surrounding that point. Our method requires the input data to be defined on a regular grid, so we bin-average the available SSH and SST observations within the domain onto regular square grids as shown in Figure 2. For this study, we choose the dimensions of the square domain to have a side-length of 960km with a grid resolution of 7.5km (128x128 grid points). The input SSH and SST observations are first re-scaled by subtracting the mean and dividing by the standard deviation, before being bin-averaged onto a 7.5km grid with zero padding where there are no observations. Our mapping method estimates the 2D SSH field on this same grid. This choice of grid dimensions limits the computational resources required for the reconstruction while ensuring that the domain remains large enough that a significant number of interacting mesoscale eddies fit within the model’s ‘field of view’. Defining a local orthonormal projection allows us to find the latitude and longitude for each grid point given the coordinates of the square domain’s center. This projection ensures that the shape and size of eddies are not distorted at different latitudes, as would be the case if the square grids were defined in latitude-longitude space.

To avoid overfitting to the persistent SSH patterns in one such domain and with a view to creating a location-agnostic SSH mapping method, we generate the training data for our DL method by randomly selecting the coordinates on which to center the square domain for each training example. For most experiments in this study the length of the observation time series used to map SSH was chosen to be 30 days, and the effect of varying this length is explored in section 5.1. The input to our mapping method for a single training example therefore consists of a time series (centered on a randomly chosen date) of bin-averaged SSH and SST observations within a square domain centered on a randomly chosen point in the ocean.

4 Deep learning method for SSH mapping

4.1 Deep learning architecture

Here, we describe the DL neural network architecture we use to map SSH, a schematic diagram of our method is shown in Figure 2. Since the dynamical relationship between SSH and SST is non-trivial, we first use a ‘ResNet’ CNN (He et al., 2016) to encode each variable separately in a learnable latent space representation. Conceptually, each ResNet encoder learns a mapping from the SSH/SST observations to some combined latent space in which the information from the SSH and SST observations can be combined on an equal footing. Note that trying out a more straightforward approach of combining SSH and SST as two different ‘channels’ in the input to a single ResNet encoder yielded lower accuracy SSH mapping.

The ResNet architecture used for each encoder consists of alternating downsampling and residual learning blocks. A downsampling block consists of: a convolution with a stride of 2, followed by a rectified linear unit (ReLU) activation function, and finally,

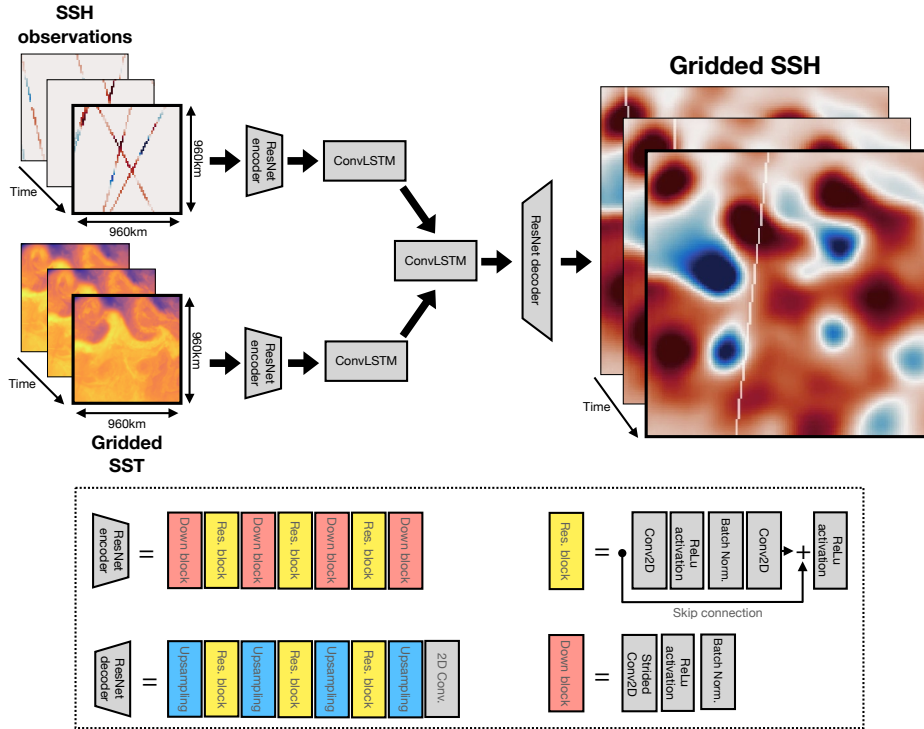


Figure 2. Schematic diagram of our deep learning method for SSH mapping. A time series of SSH and SST observations are each inputted into separate ResNet encoders (whose architecture is depicted in the lower panel) and ConvLSTM layers. Their representations are then combined through concatenation before being passed through another ConvLSTM layer and a ResNet decoder (whose architecture is in the lower panel). The loss function is then minimized along the location of the withheld altimeter observations (overlaid in white).

317 batch normalization. A residual learning block (He et al., 2016) consists of: a convolu-
 318 tion with a stride of 1, ReLU activation, batch normalization, and another convolu-
 319 tion of stride 1. The original input to the block is then combined with the output of the convolu-
 320 tions through addition (the so-called ‘skip connection’ (He et al., 2016)), before a
 321 final ReLU activation function. Note that the same ResNet encoder is applied to each
 322 time step in the time series of input observations. Thus, the resulting latent space rep-
 323 resentation is one in which observations from each time step do not yet inform the rep-
 324 resentation at other time steps.

325 The resulting time series of latent space representations are then each passed through
 326 a bi-directional convolutional long short-term memory (ConvLSTM) layer (Shi et al., 2015).
 327 ConvLSTM is a type of recurrent neural network widely used for problems involving regularly-
 328 spaced spatiotemporal data which has been demonstrated to capture complex dynamical
 329 relationships, for example to perform precipitation nowcasting (Shi et al., 2015). This
 330 layer must learn how partial observations from different times inform the state in the lo-
 331 cations where observations are missing, i.e. the dynamics governing SSH/SST evolution.
 332 Separate ConvLSTM layers are applied to the SSH and SST representations, resulting
 333 in latent space representation time series for SSH and SST where each time step is in-
 334 formed by observations at all other times.

335 The SSH and SST latent space representations are combined through concatena-
 336 tion. This time series is then passed through another bi-directional ConvLSTM layer which
 337 can learn relationships between the SSH and SST representations to give a combined la-
 338 tent space representation of the dynamical state of the ocean informed by both SSH and
 339 SST.

340 From this latent space representation, we finally use a ResNet to decode the rep-
 341 resentation to a gridded SSH map. The decoder is like the ResNet encoder networks, but
 342 with the downsampling blocks replaced by nearest-neighbor upsampling layers, alternat-
 343 ing with residual learning blocks. The final layer has the dimensions of the desired fi-
 344 nal target SSH grid using a linear activation. The same ResNet decoder is applied to each
 345 step in the time series. The full neural network (ResNet encoders, then separate Con-
 346 vLSTMs, then joint ConvLSTM, then ResNet decoder) provides an end-to-end trainable
 347 mapping from a time series of SSH and SST observations to a time series of 2D SSH maps.

348 Although our method produces a time series of SSH maps for the duration of the
 349 input observations, the reconstruction error increases away from the center of the time
 350 series. This is expected, since dynamically informative observations of past (future) states
 351 are missing for days at the beginning (end) of the time series. Thus, to generate our fi-
 352 nal SSH product, we retain only the reconstruction for the central time step, creating
 353 a time series of gridded SSH by successively shifting the input observation time series
 354 by one day.

355 4.2 Formulating a loss function in the absence of gridded SSH data

356 The network is trained by minimizing the mean squared error between the recon-
 357 structed SSH and the ground truth. However, the full 2D ground-truth SSH is not avail-
 358 able when training exclusively on real-world observations. To train using real-world ob-
 359 servations, we implement the following strategy. We first withhold some of the SSH tracks
 360 from the input and then linearly interpolate the network’s 2D reconstruction at the lo-
 361 cations of these withheld observations to calculate the mean squared error between the
 362 reconstruction and the withheld observations. This way the error for any given train-
 363 ing example is calculated only along a few withheld satellite tracks. Upon training, the
 364 network is forced to produce realistic 2D reconstructions throughout the domain since
 365 the withheld tracks appear at random places in the domain, so the network is unaware
 366 of where it will be evaluated.

367 Since there are typically several satellite altimeters operational at any time, for each
 368 training example we randomly select one of the available satellites, withhold its obser-
 369 vations from the input, and use them as the ground truth in the loss function calcula-
 370 tion. For each example, up to five satellite altimeters (depending on mission availabil-
 371 ity) are randomly selected to be used as the input SSH observations. The remaining satel-
 372 lites (or one satellite in the times when fewer than six missions are operational) are with-
 373 held for use as the ground truth when calculating the loss function.

374 To reduce over-smoothing and the appearance of high-frequency artifacts in the re-
 375 construction, we include additional regularization terms in the loss function proportional
 376 to the mean squared error in the first and second along-track derivatives of SSH. Thus
 377 the cost, \mathcal{L} , we seek to minimize during training is given by:

$$\mathcal{L} = \frac{1}{N} \sum_{i=1}^N \left(\frac{1}{\sigma_0^2} (\tilde{\eta}_i - \eta_i)^2 + \frac{\lambda_1}{\sigma_1^2} (\partial_x \tilde{\eta}_i - \partial_x \eta_i)^2 + \frac{\lambda_2}{\sigma_2^2} (\partial_x^2 \tilde{\eta}_i - \partial_x^2 \eta_i)^2 \right), \quad (2)$$

378 where N is the number of observations, η_i is the true observed SSH for the i -th obser-
 379 vation, $\tilde{\eta}_i$ is the corresponding mapped SSH, x is a spatial coordinate following the satel-
 380 lite track, σ_0^2 , σ_1^2 , and σ_2^2 are the variances of η , $\partial_x \eta$, and $\partial_x^2 \eta$ respectively, and λ_1 and
 381 λ_2 are tunable parameters controlling the relative weighting for each regularization term

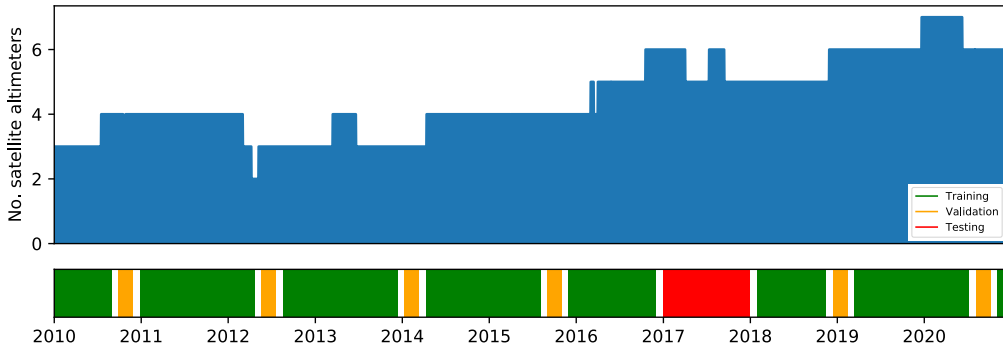


Figure 3. Dates are partitioned into non-overlapping and spaced out training (green), validation (orange), and testing (red) windows (the white bars represent the gaps between partitions when no examples are drawn). Examples are then drawn centered on days randomly drawn from the respective windows. The blue histogram shows the changing availability of satellite altimeters from which we can draw data.

382 in the loss function. In this study, we set $\lambda_1 = \lambda_2 = 0.05$. Along-track derivatives are
 383 estimated using first-order centered difference.

384 4.3 Choice of training, validation, and testing datasets

385 The data are split into training, validation, and testing datasets to prevent over-
 386 fitting. The validation dataset is used to determine at what point the model is trained
 387 to its full potential. In contrast, the testing dataset is withheld and used only at the end
 388 to test the model’s performance on unseen data. Ideally, these datasets should be inde-
 389 pendent. To this end, we divide the available history of observations into time windows
 390 for training, validation, and testing as shown in Figure 3. The year 2017 was kept aside
 391 for testing to coincide with the testing period covered by an Ocean Data Challenge which
 392 allows us to compare our method to other proposed mapping methods (this data chal-
 393 lenge is described in section 6.1). The remaining times were then split into interleaved
 394 training and validation windows. We choose interleaved windows to ensure that the train-
 395 ing and validation datasets both capture various seasons and sample interannual vari-
 396 ability. These windows were chosen such that 80% of the dates fall within the training
 397 dataset and 20% in the validation dataset (specifically, we broke the available dates into
 398 60 day chunks then chose the first 4 chunks to be training, the next 2 to be validation,
 399 then the next 4 to be training, and repeated this procedure until all dates had been as-
 400 signed). Because there is significant autocorrelation in the SSH field over a time scale
 401 of $O(10\text{-}20)$ days, it is important to choose training and validation windows that are widely
 402 enough separated in time to prevent significant autocorrelation between the training and
 403 validation data. We ensure this by leaving a gap of 30 days between training and val-
 404 idation windows, since we observed the autocorrelation in the SSH field of an ocean global
 405 circulation model (GCM) to drop by $\sim 50\%$ in the Gulf Stream Extension over this time.

406 5 Training performance

407 To provide a proof of concept for our method in this study we focus on mapping
 408 SSH in a region with energetic mesoscale dynamics, namely the Gulf Stream Extension.

409 To this aim, we draw training and validation examples from the area shown in red in Fig-
 410 ure 1. This area was chosen such that all examples feature the dynamics characteristic
 411 of the Gulf Stream Extension. In principle, we could draw training and validation ex-
 412 amples from anywhere in the ocean, but since there are a wide range of dynamical condi-
 413 tions in the ocean it is challenging for a single neural network to learn to map SSH in
 414 all regions at once (this will be further discussed in section 7.3). Training and valida-
 415 tion examples are drawn from this region with the square domain centered at a randomly
 416 selected point and time (respecting the training-validation date partition outlined in sec-
 417 tion 4.3). This region features an abundance of mesoscale eddies and complex jet-eddy
 418 interactions between the Gulf Stream and neighboring eddies. It is thus a challenging
 419 region for SSH mapping, where the widely-used DUACS product displays relatively high
 420 mapping errors, with the error becoming comparable to the standard deviation of the
 421 along-track SSH signals.

422 **5.1 Determining the optimal length of the input time series and num-** 423 **ber of training examples**

424 The SSH reconstruction on a given day is generated using a time series of obser-
 425 vations centered on that day. Thus, a key parameter of our approach is the length of this
 426 time series. Including observations from further into the past and future increases the
 427 aggregated spatial coverage of the SSH observations and provides information about the
 428 time variation of the surface flow. To accurately reconstruct the SSH in unobserved re-
 429 gions, the network needs to learn an approximation of the surface dynamics, so that ob-
 430 servations from other days can be used to inform the reconstruction.

431 Physically, we expect there to be some predictability time horizon beyond which
 432 observations become decreasingly helpful in informing the reconstruction. Therefore, we
 433 expect to see diminishing improvement to the reconstruction as the length of the time
 434 series is increased beyond some time period. We tested this hypothesis by varying the
 435 length of the input time series and training our ConvLSTM model, with both SSH and
 436 SST observations for input, and finding the mean value for the cost, \mathcal{L} , achieved on the
 437 validation dataset (Figure 4a). \mathcal{L} reduces with increasing time series length until about
 438 twenty days (i.e. ten days on either side of the reconstruction), beyond which further length-
 439 ening yields minimal improvement. As demonstrated in Figure 4b, we checked that \mathcal{L}
 440 had stopped improving significantly as a function of the number of training examples
 441 for each time series length. This time-scale could be a feature of our neural network ar-
 442 chitecture rather than a physical predictability horizon.

443 The results in the remainder of this study use a time series length of thirty days
 444 (i.e. fifteen days on either side of the date for reconstruction) noting that the significant
 445 computational cost of extending the time series beyond this length is unlikely to yield
 446 significant reductions in \mathcal{L} .

447 **5.2 Quality of mapped SSH field**

448 Figure 5 provides an example time series of SSH mapped using our method from
 449 the independent testing dataset (i.e. during 2017). The SSH field features separated, dis-
 450 tinct mesoscale eddies and a strong Gulf Stream, evident in the surface geostrophic cur-
 451 rents calculated from the SSH map. The time-evolution of the reconstructed field also
 452 looks qualitatively realistic, with a Gulf Stream meander pinching off to form a new mesoscale
 453 eddy to the north of the jet. The capability of our method to produce realistic-looking
 454 SSH fields from observations, unseen during training, gives us confidence that the model
 455 has learned a robust mapping from SSH and SST observations to the 2D SSH field, rather
 456 than overfitting to the training data.

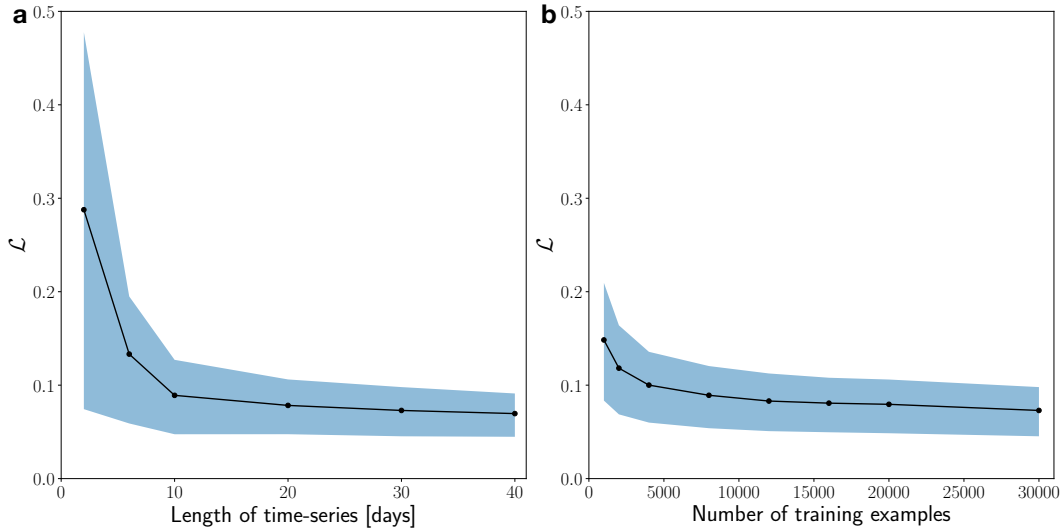


Figure 4. (a): the mean cost, \mathcal{L} , achieved on the validation dataset (containing 2000 samples) for our ConvLSTM model with SSH and SST input observations for increasing time series length. The blue shading shows the 16th and 84th percentiles, thus containing the same fraction of points as $\pm 1\sigma$ for a normal distribution. (b) the reduction in \mathcal{L} with increasing numbers of training examples for our ConvLSTM model with the length of the input observation time series set to 30 days (15 days either side of the reconstruction). Shading as for (a).

6 Inter-comparison of existing SSH mapping methods in the Gulf Stream Extension region

6.1 Mapping methods and evaluation metrics

We use the data distributed through the AVISO Ocean Data Challenge² to provide a direct quantitative comparison of our method’s accuracy to that of five established and experimental SSH mapping methods. The data consist of SSH maps created for the year 2017 (the year set aside for testing) in the region covering ($55^\circ - 65^\circ\text{W}$, $33^\circ - 43^\circ\text{N}$) using five different mapping methods, described below. Each map was created using observations from all available altimeters apart from CryoSat-2, which is withheld for use as a ground-truth for calculating the maps’ errors.

We compare our approach to five other mapping methods:

- *DUACS*: the community-standard SSH product created using OI (Taburet et al., 2019). The full mapping method is not publicly available, but a map was generated excluding CryoSat-2 for the data challenge.
- *DYMOST*: the ‘Dynamic Interpolation’ method proposed by (Ubelmann et al., 2015) and evaluated by Ubelmann et al. (2016) and Ballarotta et al. (2020). This method is a variant of the OI approach. The Gaussian a priori SSH covariance model used in DUACS is replaced by a dynamically informed model based on the forward and backward in time integration of an idealized potential vorticity conservation equation. In this method the SSH evolution is assumed to be governed only by the first baroclinic mode.

² DOI:10.24400/527896/a01-2021.005

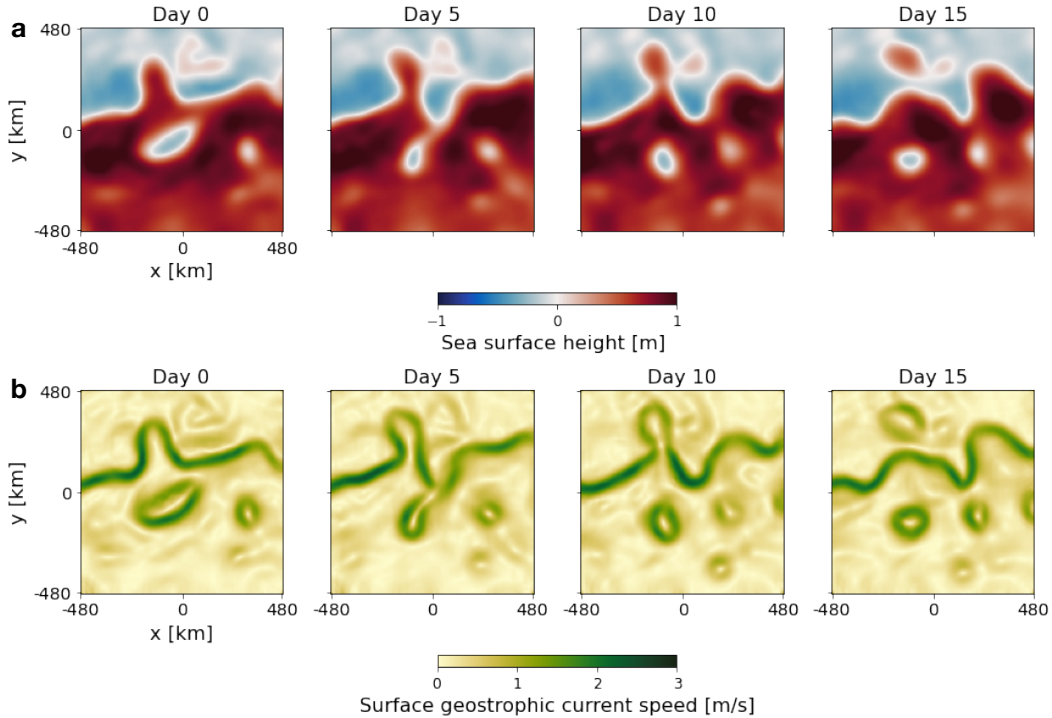


Figure 5. An example of the output from our ConvLSTM SSH+SST mapping method. Time series of (a) SSH (absolute dynamic topography) mapped using ConvLSTM SSH+SST and (b) surface geostrophic current speed calculated from the SSH field.

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- *MIOST*: the ‘Multiscale and Multivariate Interpolation’ method proposed by Ubelmann et al. (2021). This method is another variant of the OI approach where the single-scale Gaussian covariance model for SSH used in DUACS is replaced with a wavelet basis where the amplitude of each scale component is set to match the power spectrum observed in along-track altimetry, thus assuming isotropy.
 - *BFN-QG*: the ‘Back-and-Forth Nudging QG’ method proposed by Le Guillou et al. (2021). This method uses the back-and-forth nudging method for data assimilation to interpolate the altimetry observations while respecting the dynamics of a one-layer QG model of ocean turbulence.
 - *4DVarNet*: the DL method proposed by Fablet, Amar, et al. (2021) and Fablet, Chapron, et al. (2021). This approach poses the SSH interpolation problem as a four-dimensional variational data assimilation (4DVar) problem. It replaces the dynamical model and solver with trainable neural networks. We note that this method has been used in several OSSE studies, so many different versions exist in the literature. Here, we use the map provided through the data challenge which takes only SSH observations as input³.

³Fablet and Chapron (2022) and Fablet et al. (2022) recently showed that including SST in the 4DVarNet approach could significantly improve the SSH reconstruction. However, this has only been demonstrated in the OSSE setting where the full underlying ground-truth SSH is known during training. It has not yet been demonstrated to improve SSH mapping using real-world SST observations which suffer gaps due to cloud cover if the highest resolution observations are used. Hence, we cannot directly compare the two methods with the inclusion of SST but note that the improvements seen for both methods in their

494 For our analysis, we reduce the size of the data challenge’s study region to (55° -
 495 65°W, 34° - 42°N) to ensure the full region fits inside the ‘field of view’ of our neural
 496 network’s 960x960km output. All methods are assessed in this same region. Figure 5 shows
 497 an example of the result of applying our method to map SSH in this study region.

498 We use the withheld CryoSat-2 observations to calculate various metrics to quan-
 499 titatively compare the performance of the SSH mapping methods. These metrics are sum-
 500 marized below:

- 501 • *Mean RMSE (cm)*: The SSH maps are interpolated to the locations of the with-
 502 held observations and the root mean square error (RMSE) of the map is calcu-
 503 lated for each day where observations are available, then the mean of these daily
 504 RMSEs is taken.
- 505 • *Standard deviation of RMSE (cm)*: The standard deviation of the daily RMSE
 506 values.
- 507 • *Effective spatial resolution (km)*: Calculated using the along-track SSH spectra
 508 for each satellite pass. This is the wavelength at which the power spectral den-
 509 sity of the misfit between the map and the observations becomes comparable to
 510 that of the observations, quantifying the spatial scales accurately resolved by the
 511 map. See Ballarotta et al. (2019) for a complete discussion of this metric.
- *RMSE score (no units)*:

$$\text{RMSE score} = 1 - \frac{\text{RMSE}_{\text{map}}}{\text{RMS}_{\text{obs}}}$$

512 A perfect map would have a score of 1, a map that predicted zeroes everywhere
 513 would have a score of 0.

- 514 • *Standard deviation of RMSE score (no units)*: Standard deviation of the daily RMSE
 515 scores.

516 6.2 SSH validation against independent satellite observations

517 6.2.1 Root mean square error (RMSE)

518 All the experimental mapping methods were found to have lower mean RMSE val-
 519 ues than DUACS and the spread in their daily RMSE is also lower, as can be seen in Fig-
 520 ure 6 and Table A1. BFN-QG provides only a marginal improvement in RMSE, whereas
 521 DYMOST and MIOST significantly improve (a 12% reduction in mean RMSE). All the
 522 DL methods tested here (4DVarNet SSH and all configurations of our ConvLSTM method)
 523 yielded lower RMSE than all the other methods. 4DVarNet SSH showed a 14% reduc-
 524 tion in RMSE while our ConvLSTM method respectively gave a 13% and 17% reduc-
 525 tion in RMSE when run with just SSH (ConvLSTM SSH) and with combined SSH and
 526 SST observations (ConvLSTM SSH+SST).

527 To explore the eddy-jet configurations that present the most significant challenge
 528 to traditional altimetry mapping and further illustrate our method’s improvement, in
 529 Figure 7 we present three case studies from the testing period where we contrast the DU-
 530 ACS reconstruction with our ConvLSTM SSH+SST method. These cases correspond
 531 to the three CryoSat-2 satellite tracks of at least 800km in length for which the DUACS
 532 reconstruction shows its highest RMSE. In all three cases, the primary source of error
 533 is the misrepresentation of a sharp kink in the Gulf Stream in the top left corner of the
 534 domain (recall that geostrophic surface currents flow along contours in SSH). In all three
 535 cases ConvLSTM SSH+SST leads to a significantly more accurate along-track SSH pro-
 536 file. Physically, the eddy features in SSH are expected to evolve relatively quickly in shape

respective settings are consistent with the expectation from SQG dynamics that the SST field contains relevant information about the dynamics of the SSH field.

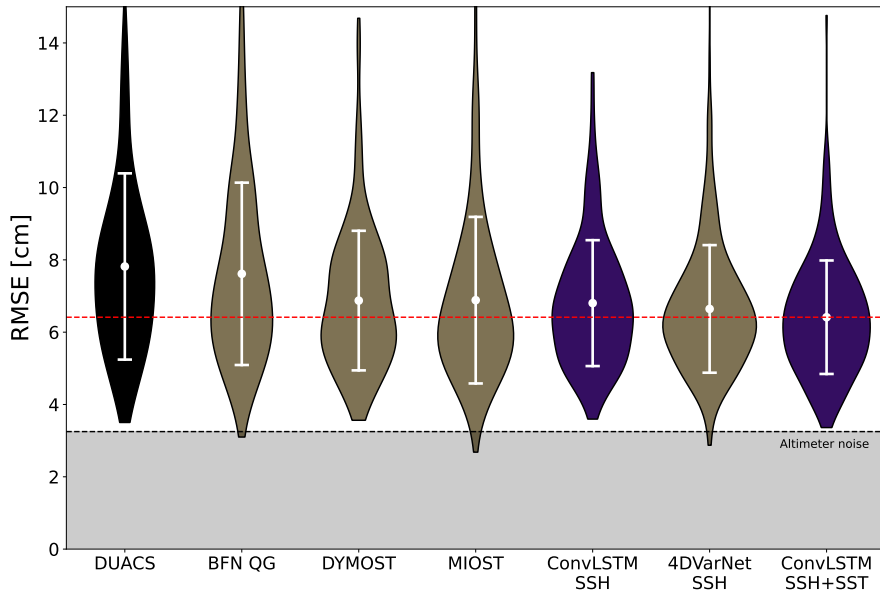


Figure 6. Violin plot showing the distribution of daily reconstruction RMSE for each method when compared to independent satellite altimeter observations for the test period. The black violin shows the community-standard OI product, the gold violins are experimental methods proposed by others in the literature, and the purple violins are using our new method. The width of each violin is proportional to the probability density function, the white circle indicates the mean, and the whiskers show one standard deviation. The shaded region indicates the estimated instrument noise in the observations (as reported by CMEMS).

537 and amplitude in regions where the jet (i.e. the Gulf Stream) is bent at a sharp angle,
 538 thus presenting a challenge to the simple SSH covariance model employed by DUACS.
 539 Our ConvLSTM SSH+SST method overcomes this shortcoming.

540 **6.2.2 Effective along-track spatial resolution**

541 All the experimental mapping methods were able to reconstruct smaller-scale SSH
 542 signals than DUACS, as can be seen in Figure 8. The spectra in Figure 8a show that the
 543 DUACS reconstruction underestimates the strength of the SSH signal at smaller scales
 544 compared to the observations. All other methods show spectra closer to the observations
 545 at these scales to varying degrees. The spectral coherence plot in Figure 8b demonstrates
 546 that the DUACS reconstruction is less accurate than all other methods at scales of 70-
 547 200km. None of the methods accurately resolve SSH below these scales due to the lim-
 548 ited along-track resolution of the observations and the size of the gaps between obser-
 549 vations.

550 The effective spatial resolution (the scale at which each curve in Figure 8b crosses
 551 0.5, also listed in Table A1) for each method reveals significant differences. MIOST and
 552 DYMOST only offer marginal improvements on this metric (a 7% and 12% reduction
 553 in the smallest wavelength accurately resolved), whereas BFN-QG provides a more sub-
 554 stantial improvement (20%). Again, all DL methods yield the largest improvements on

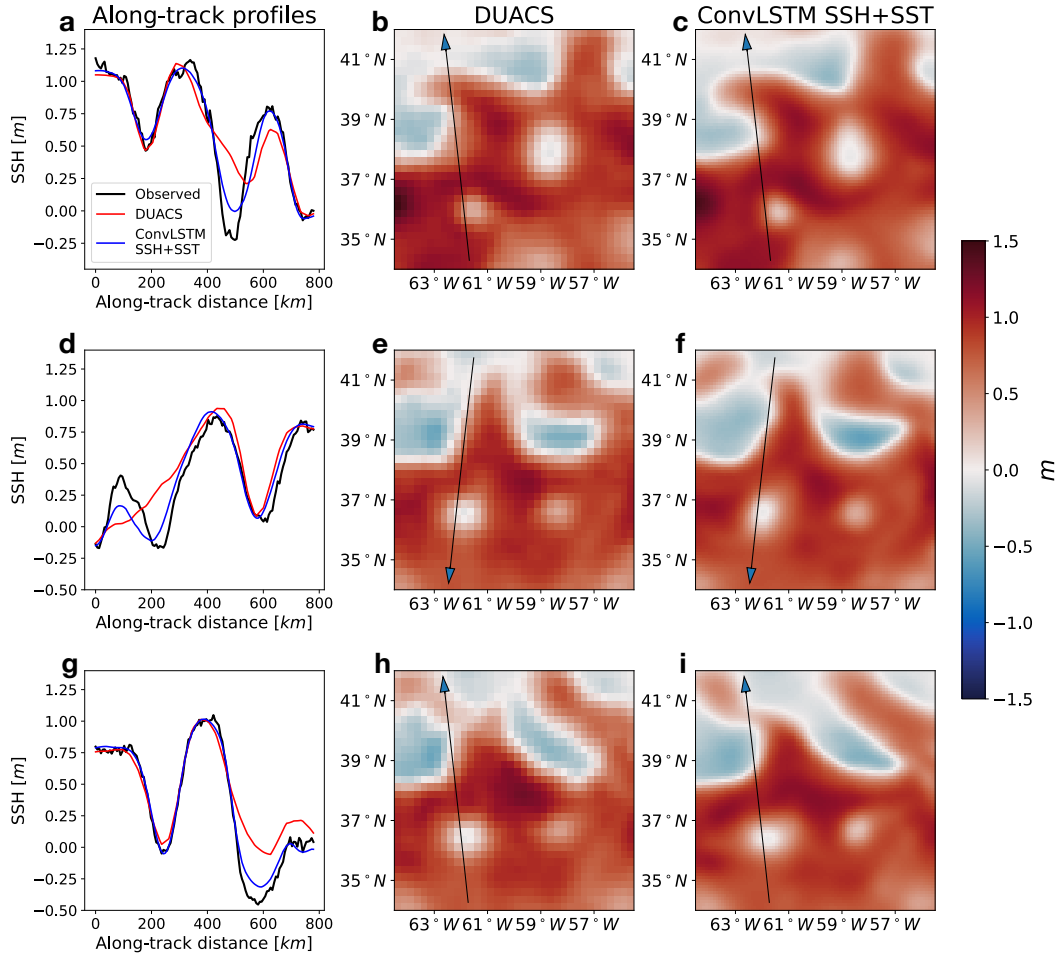


Figure 7. Example along-track SSH profiles. Shown are the three 800km track segments (a, d, g) where DUACS suffers the highest RMSE during 2017 along with the track’s location superimposed over the DUACS reconstruction (b, e, h) and the ConvLSTM SSH-SST reconstruction (c, f, i). In each case, DUACS over-smoothes one or more large peaks/troughs in SSH. In all of these cases ConvLSTM provides a significant improvement.

555 this metric, with ConvLSTM SSH, 4DVarNet SSH, and ConvLSTM SSH+SST provid-
 556 ing 23%, 28%, and 30% improvements respectively. The SSH map produced using our
 557 ConvLSTM SSH+SST method accurately resolved SSH signals with wavelengths as small
 558 as 104km.

559 The significant improvement in all metrics by all the DL methods provides a comp-
 560 pelling case for the use of DL in developing an improved 2D SSH product. The addition
 561 of SST observations also provides a clear improvement in our method’s reconstruction
 562 accuracy and resolution.

563 **6.2.3 Frequency spectra of SSH fields**

564 Another important dynamical property of a mapped SSH field is its frequency spec-
 565 trum. Unlike for the wavenumber spectrum, discussed in section 6.2.2, estimating a ground
 566 truth from the independent altimeter observations is challenging since at any one loca-

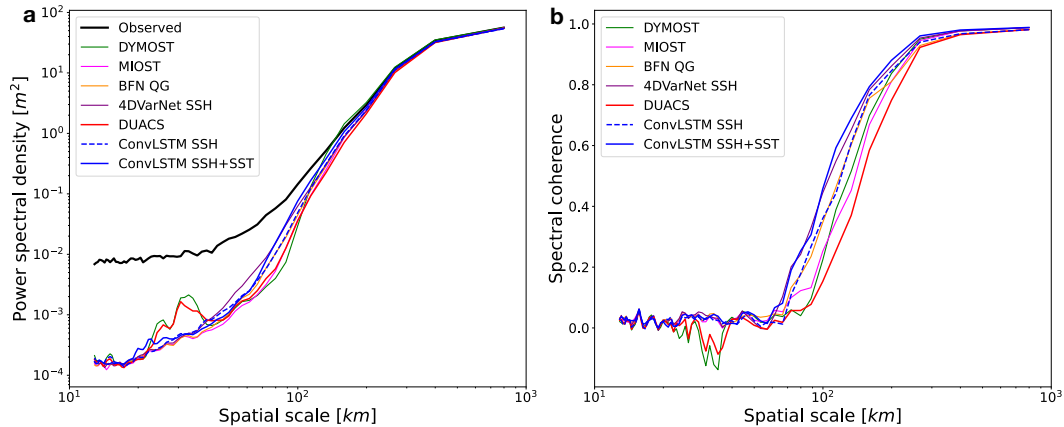


Figure 8. (a): wavenumber spectrum of mapped SSH along the tracks of the withheld CryoSat-2 observations for each method, compared to that of the observations. The observations show a high power spectral density for small scales due to instrument noise. (b): spectral coherence for each mapping method (approaches unity for scales at which the SSH signal is well resolved by the map).

567 tion the temporal sampling is relatively sparse (CryoSat-2 observes the same $1/8^\circ$ by $1/8^\circ$
 568 spatial bin once every 5-6 days on average in this region). Thus, we cannot calculate the
 569 frequency spectral coherence in the manner of section 6.2.2 using just the withheld al-
 570 timeter observations.

571 Nevertheless, we can still characterize the frequency spectrum of each SSH map and
 572 draw comparisons among the maps, as shown in Figure 9a. Compared to all other meth-
 573 ods, DUACS has a substantially lower power spectral density at time scales shorter than
 574 30 days, implying that DUACS may be underestimating the short-time SSH variability,
 575 as has been noted in the literature (Ballarotta et al., 2019). The other maps all show
 576 similar frequency spectra down to time scales of around 5 days, below which 4DVarNet
 577 and BFN QG show significantly higher power spectral density. Without knowing the ground
 578 truth frequency spectrum from observations, it is hard to discern which map’s spectrum
 579 most accurately represents the variability at these time scales. The higher variability at
 580 short temporal scales for 4DVarNet and BFN QG could be due to the appearance of non-
 581 physical artifacts along the location of the input satellite tracks (discussed more in sec-
 582 tion 6.3).

583 To provide some estimate of the SSH variability from the observations, we also cal-
 584 culated the second-order structure function in time, as shown in Figure 9b. To do this,
 585 we used all available satellite altimetry observations from the study region and time, se-
 586 lected the 100 most frequently observed $1/8^\circ$ by $1/8^\circ$ bins, used the observations to con-
 587 struct a SSH time series for each bin, and used these to estimate the second order
 588 structure function. We thus assume that the frequency spectrum of the SSH field is approx-
 589 imately isotropic within this region. We used the bias-corrected and accelerated boot-
 590 strap method to estimate 95% confidence intervals.

591 Calculating the second-order structure function for each of the SSH maps at uniformly-
 592 spaced points throughout the region allows us to compare the maps’ variability to the
 593 observational estimate. We see that DUACS is indeed underestimating the variability
 594 of the SSH field at time scales below 20-30 days as the higher variability shown by the
 595 other maps is replicated in the observations. The differences between the other maps at
 596 short (< 10 days) time scales are small and the confidence interval on the observational

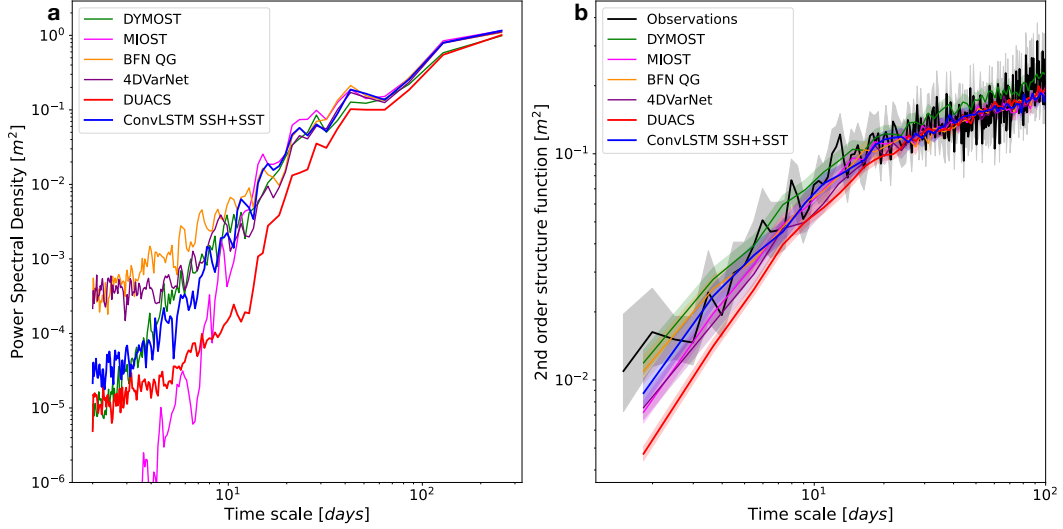


Figure 9. (a): frequency spectrum for each SSH map averaged over points throughout our study region. (b): second order structure function calculated from all available satellite observations (including those used to make the maps) at satellite track crossover points, compared against that of each map average over uniformly-spaced points in the domain. Error bars show the 95% confidence interval estimated using the bias-corrected and accelerated bootstrap method.

597 estimate is large since satellite altimetry provides poor sampling of the SSH field at these
 598 short time scales.

599 6.3 Surface geostrophic currents

The primary use of SSH maps is to infer surface ocean currents through the assumption of geostrophy. The eastward and northward geostrophic surface currents (u_g and v_g respectively) are related to the SSH, η , through

$$(u_g, v_g) = \frac{g}{f} \left(-\frac{\partial \eta}{\partial y}, \frac{\partial \eta}{\partial x} \right), \quad (3)$$

600 where x and y are eastward and northward spatial coordinates respectively, f is the Cori-
 601 olis frequency, and g is the acceleration due to gravity. The currents are thus propor-
 602 tional to the first-order spatial derivatives of η .

From these currents, other physical quantities such as relative vorticity, ω , and strain rate, γ , can be calculated by taking spatial derivatives of the velocity field (corresponding to second-order derivatives of the SSH field):

$$\omega = \frac{\partial v_g}{\partial x} - \frac{\partial u_g}{\partial y} \quad (4)$$

$$\gamma = \sqrt{\left(\frac{\partial u_g}{\partial x} - \frac{\partial v_g}{\partial y} \right)^2 + \left(\frac{\partial v_g}{\partial x} + \frac{\partial u_g}{\partial y} \right)^2}. \quad (5)$$

603 These higher order derivatives are an important diagnostic since they give information
 604 about the dynamics (i.e. the acceleration) of the surface motions (Hua & Klein, 1998).
 605 Relative vorticity (and potential vorticity) are crucial quantities when studying ocean
 606 turbulence because in 2D turbulence vorticity is conserved along streamlines and poten-
 607 tial vorticity is conserved in QG and SQG motion. As such the interaction of eddies with

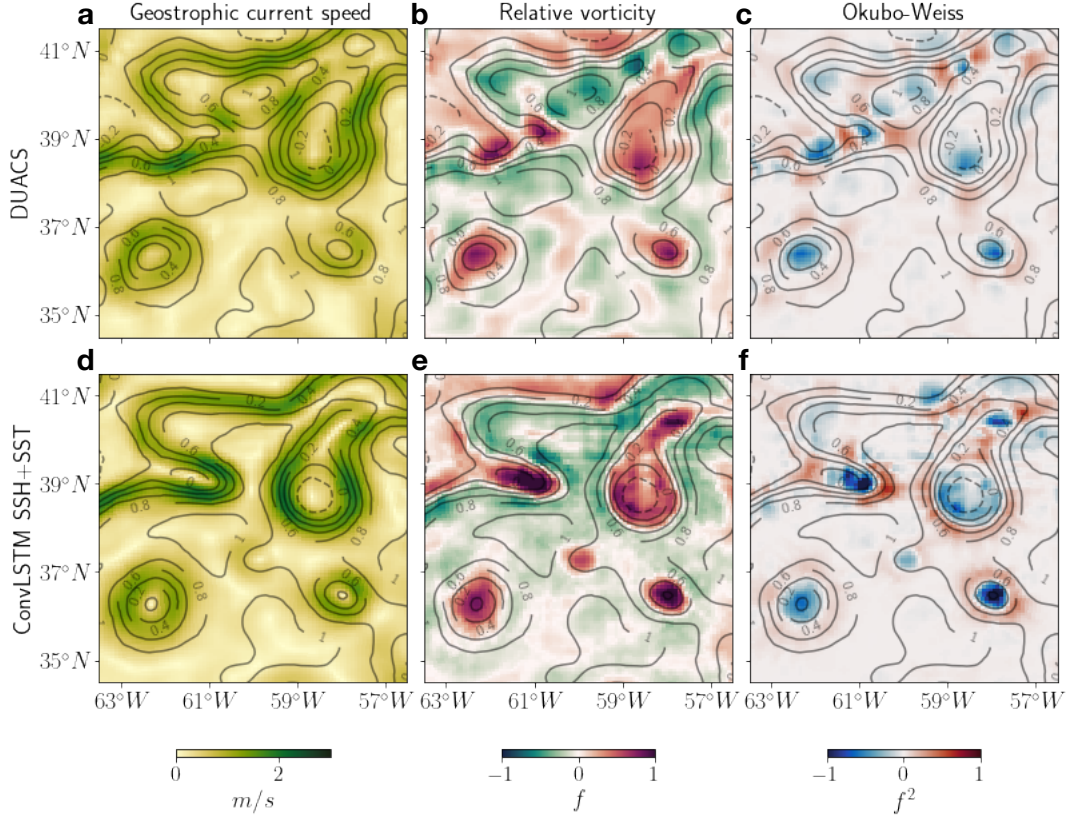


Figure 10. (a-c): surface geostrophic current speed, relative vorticity, and the Okubo-Weiss quantity calculated from the DUACS SSH map. (d-f): the same quantities calculated from the SSH map produced by our ConvLSTM SSH+SST method. Contours of the corresponding SSH maps (absolute dynamic topography) are overlaid in black. Relative vorticity and Okubo-Weiss are normalized by the Coriolis frequency, f , and its square respectively.

608 each other and with the mean flow, as well as their equilibration, filament formation, and
 609 and frontogenesis are all processes that are viewed through the lens of vorticity and strain
 610 evolution.

One such quantity that is important to consider is the Okubo-Weiss quantity, W , defined as the difference between the square of the strain rate and the square of relative vorticity (Okubo, 1970; Weiss, 1991; McWilliams, 1984; Hua & Klein, 1998; Hua et al., 1998)

$$W = \gamma^2 - \omega^2. \quad (6)$$

611 When W is negative, relative vorticity and therefore rotation dominates. This typically
 612 occurs in the regions inside mesoscale eddies where SSH extrema are found. However,
 613 when W is positive or close to zero, the strain rate dominates or is close to vorticity, indicat-
 614 ing an intensification of SST fronts and associated vertical velocity. This typically
 615 occurs in regions between mesoscale eddies and at the eddy edges (Hua et al., 1998; Lapeyre
 616 et al., 1999; Klein et al., 2019).

617 We calculated the surface geostrophic currents, ω , and W from the DUACS and
 618 ConvLSTM SSH+SST maps. To closely match the procedure for calculating current speeds
 619 from the DUACS SSH map used to generate the surface current speeds distributed in

620 the CMEMS Level 4 SSH product⁴, we use a 9-point stencil (as advocated in Arbic et
 621 al. (2012)) for estimating derivatives of the DUACS SSH field. Because the DUACS prod-
 622 uct is distributed on a $1/4^\circ$ ($\sim 30\text{km}$) grid whereas our SSH reconstruction is on a 7.5km
 623 grid, we take all DUACS derivatives on its native grid before linearly interpolating to
 624 our higher-resolution grid.

625 The ConvLSTM SSH+SST reconstruction results in narrower and stronger cur-
 626 rents than those calculated from the DUACS reconstruction, as shown in Figure 10a&d.
 627 The maximum current speed for the date shown in Figure 10 is 40% stronger in the Con-
 628 vLSTM SSH+SST reconstruction than in the DUACS reconstruction. This result is ex-
 629 pected since the covariance model used in DUACS leads to an overly smooth SSH field,
 630 smaller-scale features being blurred out. Hence, the magnitude of the steepest gradients
 631 of the SSH field will, in turn, be weaker.

632 The NOAA Atlantic Oceanographic & Meteorological Laboratory (AOML) global
 633 surface drifter program provides in-situ observations of the total near-surface (15m depth)
 634 ocean currents through satellite-tracked drifters. To quantitatively test the accuracy of
 635 the surface currents inferred from the SSH maps, we here compare the mapped currents
 636 to all available AOML drifters within our study region and time. We use the global AOML
 637 drifter product distributed by CMEMS⁵. Both ConvLSTM and 4DVarNet show a clear
 638 improvement in current reconstruction accuracy relative to DUACS and the other meth-
 639 ods, with all the DL methods showing a 10-12% decrease in current speed RMSE (cur-
 640 rent speed RMSE for each method can be found in Table A1). Surface drifters don't di-
 641 rectly measure geostrophic currents, since Ekman and ageostrophic components will be
 642 aliased onto the drifter-observed currents. However, the fact that we can more accurately
 643 estimate the surface currents through geostrophy from our improved SSH map than from
 644 the DUACS map is compelling since surface drifters are an observational dataset that
 645 is entirely independent of the satellite altimetry methodology.

646 The apparent difference between the DUACS and ConvLSTM SSH+SST recon-
 647 structions becomes starker for the relative vorticity and Okubo-Weiss quantity, W , (shown
 648 in Figure 10b,c,e,f) since deficiencies in the reconstructed SSH field are amplified as deriva-
 649 tives of higher order are taken. The path of the Gulf Stream is hard to identify from the
 650 DUACS relative vorticity field visually and is practically impossible to pick out in the
 651 Okubo-Weiss field. By contrast, the ConvLSTM SSH+SST map results in a relative vor-
 652 ticity and Okubo-Weiss fields that show a clear, physically realistic Gulf Stream trajec-
 653 tory. Two evident mesoscale eddies can also be seen in the Okubo-Weiss field as distinct
 654 blue regions where vorticity dominates the flow surrounded by red regions where the eddy
 655 is inducing high strain rates. One deficiency of our reconstruction is the presence of faint
 656 high frequency grid-like artifacts. Such artifacts are common when using CNNs, and their
 657 intensity decreases with increasing number of training examples. We discuss these more
 658 in section 7.3 but note that they only become visually apparent when taking higher-order
 659 derivatives of the SSH field.

660 Although several of the methods in this study exhibit similar summary statistics,
 661 there are large qualitative differences in the resulting current reconstructions. In Fig-
 662 ure 11, we show the relative vorticity field calculated from each SSH map on the same
 663 date. The DYMOST and MIOST vorticity fields are qualitatively like that of DUACS
 664 so the Gulf Stream's large-scale structure is challenging to discern. By design, BFN-QG
 665 results in a physically realistic-looking vorticity field since the SSH map is the result of
 666 solving an idealized dynamical model of ocean turbulence, albeit an overly-idealized one.
 667 4DVarNet produces a vorticity map with similar large-scale structures to other meth-

⁴ DOI:10.48670/moi-00148

⁵ The results shown here are after we applied the wind slippage correction distributed within the prod-
 uct, however, removing this correction does not change the hierarchy of the methods' errors.

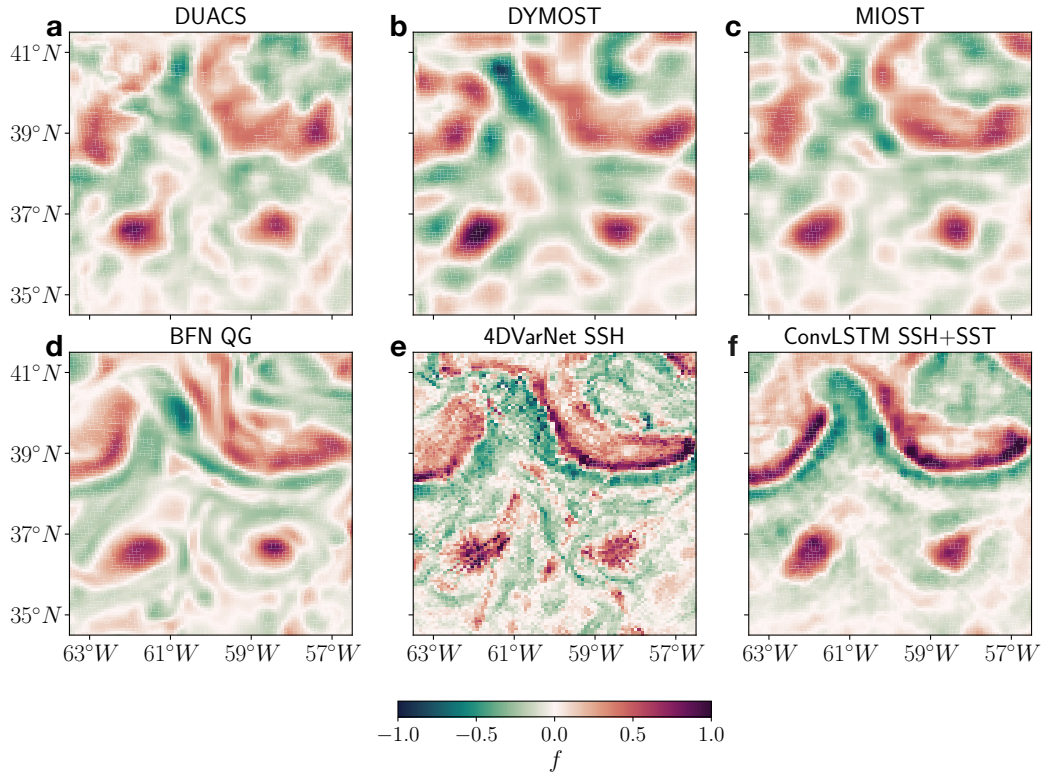


Figure 11. Relative vorticity calculated from the surface geostrophic currents from each method’s SSH map, normalized by the Coriolis frequency, f .

668 ods but with the addition of some smaller-scale features and clear artifacts. The promi-
 669 nent mesoscale eddies (the two distinct red regions in the lower half of the domain) ex-
 670 hibit high-frequency radial features in the 4DVarNet reconstruction that are likely to be
 671 CNN-induced artifacts. In this reconstruction, there are straight-line features along which
 672 the map exhibits smaller-scale features. These lines correspond to the ground tracks of
 673 the satellites providing the input SSH observations (see Supplementary Video for the vor-
 674 ticity evolution).

675 These qualitative differences have impact on two key physical quantities of interest:
 676 eddy kinetic energy (EKE) and enstrophy. Figure 12a shows each map’s time- and
 677 domain-averaged EKE. DUACS has a significantly lower EKE than the other maps while
 678 DYMOST and ConvLSTM SSH+SST have the highest EKE. Larger-scale flows domi-
 679 nate EKE, so the fact that DYMOST and ConvLSTM SSH+SST have a substantially
 680 higher EKE than other methods implies these reconstructions result in faster large-scale
 681 currents (e.g. the Gulf Stream) than the others. Figure 12b shows the time-and domain-
 682 averaged enstrophy for each map. Enstrophy is a quantity that is typically dominated
 683 by smaller-scale flows. DUACS has a significantly lower enstrophy than the other maps,
 684 consistent with the expectation that OI smooths out small-scale features. 4DVarNet has
 685 a very high enstrophy compared to the other methods due to the unphysical high-frequency
 686 artifacts in the relative vorticity field.

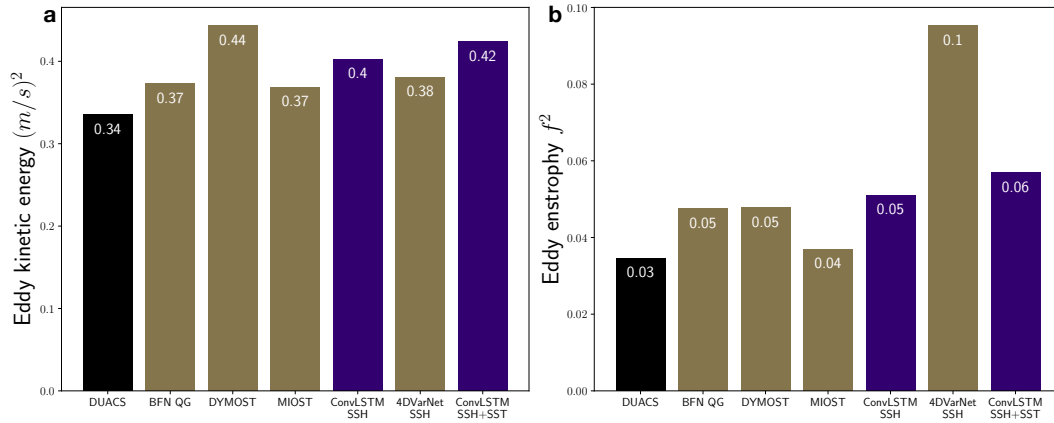


Figure 12. Domain- and time-averaged eddy kinetic energy (a) and enstrophy (b) for the geostrophic currents calculated from each SSH map. Enstrophy is normalized by the square of the Coriolis frequency, f . The black bars shows the community-standard OI product, the gold bars are experimental methods proposed by others in the literature, and the purple bars are using our new method.

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6.4 New mesoscale eddy found by synthesizing SST and SSH observations

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To further emphasize the power of using SST observations to help map SSH, we here present one case study from our testing dataset where ConvLSTM SSH+SST was able to reconstruct a mesoscale eddy that was missed by all the other methods, which only used SSH observations as input. Figure 13 shows the relative vorticity fields for five consecutive days in July 2017 reconstructed using both ConvLSTM SSH and ConvLSTM SSH+SST. On July 7th in the ConvLSTM SSH+SST reconstruction there is a cyclonic mesoscale eddy in the center of the domain (boxed in blue in the figure) which is not present in the ConvLSTM SSH reconstruction or any of the other maps (not shown in the figure). None of the satellite altimeters passed over this eddy (the altimeter tracks are overlaid on the figure), but there is a clear cold signature in the SST coinciding with the reconstructed eddy. The formation process of the eddy can be seen in the SST field for the days leading up to July 7th. On July 3rd, a filament develops on the large eddy in the northeast of the domain (boxed in black in the figure). Over the following days this filament becomes unstable and sheds to form the new eddy. This formation process all happens within just 5 days and no altimeter track passes over the region during this period, making it very challenging for an SSH mapping method that only uses altimetry observations to predict this process from the available observations. Using SST observations to aid SSH mapping is thus essential if we are to capture the formation of short-lived, small-scale mesoscale eddies.

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7 Discussion and conclusions

7.1 Summary

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In this study, we presented a DL framework for mapping SSH from satellite altimetry and SST observations, motivated by the close dynamical relationship between these surface ocean variables. By training and testing our method on real-world satellite observations, we demonstrated that we could reconstruct the mesoscale SSH field with a high degree of physical realism. Including SST as additional observations caused a significant improvement in the accuracy of our reconstruction, in line with expectations from

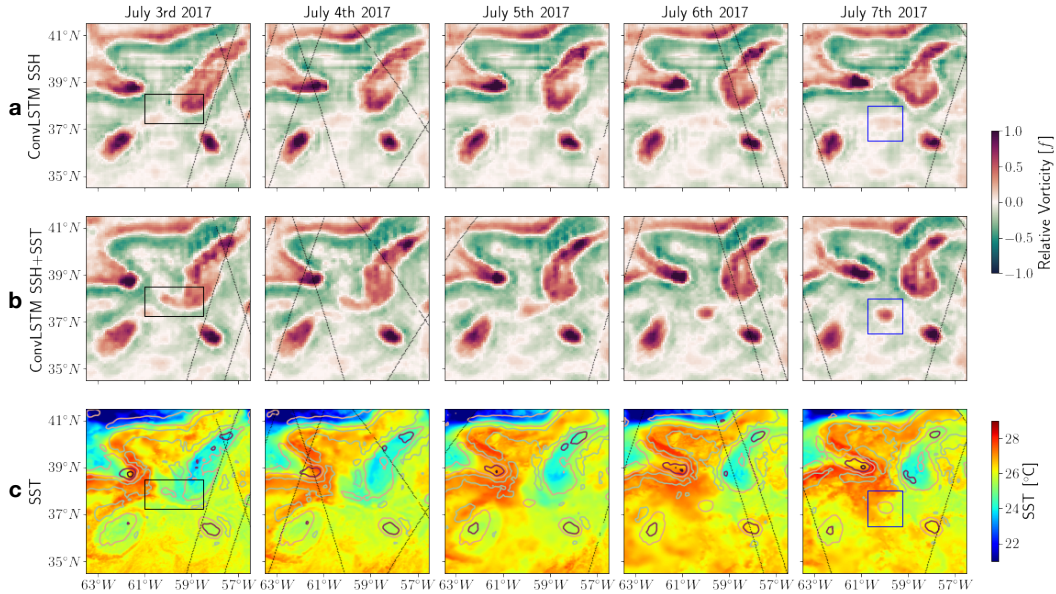


Figure 13. Relative vorticity from (a) the ConvLSTM SSH and (b) the ConvLSTM SSH+SST reconstructions for 5 consecutive days. GHRSSST MUR SST for the same days is shown in row (c) and the locations of the satellite altimetry observations on each day are overlaid in black. ConvLSTM SSH+SST relative vorticity contours in increments of $0.5f$ from $-1.25f$ to $1.25f$ are overlaid over the SST observations. The black box on July 3rd highlights the filament from which this eddy formed.

716 SQG theory. When compared against several other SSH mapping methods proposed in
 717 the literature, our method was able to map SSH in the Gulf Stream Extension with the
 718 highest accuracy and effective spatial resolution. Our study provides a roadmap towards
 719 replacing the decades-old OI method with state-of-the-art deep learning methods to pro-
 720 duce higher-accuracy and higher-resolution global SSH maps.

721 7.2 Physical implications of improved mesoscale SSH field

722 The higher-accuracy and higher-resolution mesoscale SSH field that our DL map-
 723 ping method provides is an exciting new dataset for physical oceanographers. The geostrophic
 724 surface currents calculated from our SSH field resulted in a substantially (25%) higher
 725 EKE than from the DUACS SSH map. This result, if replicated in other ocean regions,
 726 suggests that current satellite altimetry-derived estimates of oceanic mesoscale eddies'
 727 contribution to the global energy budget could be a significant underestimate (Thoppil
 728 et al., 2011; Xu et al., 2014; Martínez-Moreno et al., 2019). An updated global estimate
 729 of EKE using our DL SSH mapping method is the subject of an ongoing study.

730 Our SSH reconstruction also resulted in relative vorticity and Okubo-Weiss quan-
 731 tity, W , fields that appear qualitatively to be more physically realistic, with a clearer sep-
 732 aration between the Gulf Stream and nearby coherent mesoscale eddies. As discussed
 733 in 6.3, W is an important diagnostic for distinguishing between rotation-dominated re-
 734 gions (inside mesoscale eddies) and regions of high strain with SST front intensification
 735 and relatively strong vertical velocities. Thus, W is a prerequisite to get access to the
 736 3D eddy field (Qiu et al., 2020; Siegelman et al., 2020). Another interpretation of W is
 737 that, whereas first-order SSH derivatives only give access to currents, second-order deriva-
 738 tives (through W) give access to accelerations and therefore to the time evolution of the

739 eddy field (Hua & Klein, 1998; Hua et al., 1998; Lapeyre et al., 1999). One consequence
 740 is that improving the diagnosis of second-order SSH derivatives would lead us to bet-
 741 ter estimate the kinetic energy cascade, i.e. the balance between the merging between
 742 eddies leading to larger eddies (inverse KE cascade) and the splitting of eddies by other
 743 eddies leading to smaller eddies (direct KE cascade) (Scott & Wang, 2005; Klein et al.,
 744 2019; Storer et al., 2022). These comments emphasize the importance of accurately di-
 745 agnosing second-order SSH derivatives. Our DL SSH mapping method could allow this
 746 diagnosis to now be made from current satellite observing capabilities.

747 Much of what we know about coherent mesoscale eddies in the ocean comes from
 748 applying eddy tracking algorithms to SSH maps (Chelton et al., 2011; Martínez-Moreno
 749 et al., 2019, 2021). By including SST observations in the input to the SSH mapping, we
 750 demonstrated in section 6.4 that we could reconstruct short-lived, small-scale coherent
 751 eddies that would otherwise not be captured in SSH maps. Thus, applying eddy track-
 752 ing algorithms to our new SSH map could lead to changes in the conclusions drawn in
 753 these past studies, especially about the number of detected coherent eddies, their shapes
 754 and strengths.

755 **7.3 Potential for development of a global DL SSH product**

756 While the results of this study are promising, several challenges must be overcome
 757 before a global DL SSH map can be made available to the oceanography community. The
 758 main challenge we foresee is that different regions of the ocean exhibit a diverse range
 759 of dynamical regimes. It would therefore be challenging for a single DL model to learn
 760 to map SSH accurately in all ocean regions. This is evident in the final row of table A1
 761 where we show the result of training our method on examples drawn from anywhere in
 762 the world (rather than just from the Gulf Stream Extension as for the results in the main
 763 text). While an improvement with respect to traditional SSH mapping methods is still
 764 seen, the accuracy and resolution are significantly worse than when the training data comes
 765 only from the Gulf Stream Extension. This result is unsurprising, since regions with Gulf
 766 Stream-like dynamics make up a small fraction of the global ocean, so the training data
 767 distribution will be dominated by more quiescent, ocean interior regions. An approach
 768 involving an ensemble of regional DL models could be employed to overcome this chal-
 769 lenge where bespoke models are trained for each ocean region before the resulting regional
 770 SSH maps are merged to form a global product.

771 The power of DL models comes from their ability to continue to improve dramati-
 772 cally as they are trained on larger volumes of data. While we limited ourselves here to
 773 using real-world observations for the training dataset, future work could employ ‘transfer
 774 learning’ where the DL model is trained on a large available dataset drawn from a
 775 system with similar characteristics to that under consideration before being fine-tuned
 776 on the desired dataset (Subel et al., 2022). In the context of SSH mapping, the abun-
 777 dant output from high-resolution ocean GCMs might provide a valuable dataset for pre-
 778 training our DL model, though all GCMs contain imperfect representations of ocean physics
 779 and so the resulting SSH product would need to be used with caution if being used for
 780 observational validation of the GCM itself.

781 As highlighted in section 6.3, CNN methods such as ours and 4DVarNet can suf-
 782 fer non-physical high-frequency artifacts. Without enforcing hard physical constraints
 783 on the SSH reconstruction, there is to our knowledge no way to guarantee that no such
 784 artifacts would appear in the SSH map. However, in this study we found adding terms
 785 regularizing the first and second along-track SSH derivatives significantly reduced the
 786 artifacts.

787 For a global DL SSH product to become widely adopted in the oceanographic com-
 788 munity, the product must have a consistent spatial resolution that does not depend on
 789 the observational sampling. We have demonstrated here that this can be achieved by con-

790 structuring a training dataset where the loss function is only calculated at points not in-
 791 cluded in the input to ensure that the DL model does not learn to skillfully reproduce
 792 a high-resolution field near the input observations but with a smoother low-resolution
 793 map in the gaps between observations. The development and validation of a global SSH
 794 product using our method is the subject of an ongoing study.

795 Other ocean observation datasets could be incorporated into a similar framework
 796 to that presented here to improve the mapping of SSH or other quantities of interest.
 797 Surface salinity and ocean color are routinely observed by satellites and may contain dif-
 798 ferent dynamical signatures to the SST observations. ARGO profiles provide regular ver-
 799 tical profiles of temperature and salinity (among other variables) albeit with sparser hor-
 800 izontal sampling density than satellite observations of the surface. Since the vertical struc-
 801 ture of mesoscale eddies plays a significant role in governing their evolution, it is possi-
 802 ble that incorporating ARGO observations into a similar DL framework to that described
 803 here would lead to more accurate SSH mapping. Although described here in the con-
 804 text of satellite altimetry, our DL framework could similarly be applied to other remote
 805 sensing interpolation problems.

806 Finally, the recent launch of the Surface Water Ocean Topography (SWOT) mis-
 807 sion will soon offer, for the first time, high-resolution 2D snapshots of the SSH field through
 808 its wide swath. The development of a method to use SWOT observations to map high-
 809 resolution SSH is an area of active research (Beauchamp et al., 2020; Fablet, Amar, et
 810 al., 2021; Le Guillou et al., 2021) and our framework could also be adapted to suit the
 811 nature of SWOT observations. The launch of the SWOT mission certainly does not re-
 812 move the need for a more accurate method for mapping SSH from traditional nadir alti-
 813 metry. The nadir observations are available over a longer period than the likely dura-
 814 tion of the SWOT mission, and so are an important dataset for studies of inter-annual
 815 variability and climate change. SSH maps and the tracked eddies inferred from them are
 816 also widely used to aid the interpretation of in-situ observations across all oceanographic
 817 disciplines, so improving the SSH maps for the pre-SWOT years would add value to many
 818 years of already-collected in-situ observations. To this end, SWOT observations in fu-
 819 ture could be used as a ground truth to further validate our nadir altimeter-derived DL
 820 SSH map or calculate the loss function during training.

821 The effective use of a DL framework such as that outlined in this study to map the
 822 global ocean’s SSH field from existing and future satellite observations would represent
 823 an exciting new paradigm for studying surface ocean dynamics and, by extension, global
 824 climate.

825 Appendix A Full performance metric comparison

Mapping methods’ summary statistics						
Mapping method	RMSE [cm]	std. of RMSE [cm]	RMSE score [no units]	std. of RMSE score [no units]	Eff. resolution [km]	Drifter RMSE [m/s]
DUACS	7.7	2.6	0.88	0.055	149	0.213
DYMOST	6.8	2.0	0.89	0.047	131	0.208
MIOST	6.8	2.3	0.89	0.057	139	0.203
BFN-QG	7.5	2.6	0.88	0.053	119	0.200
4DVarNet SSH	6.6	1.8	0.90	0.046	107	0.188
ConvLSTM SSH	6.7	1.8	0.89	0.050	115	0.192

ConvLSTM SSH+SST (Gulf Stream)	6.4	1.6	0.90	0.048	104	0.190
ConvLSTM SSH+SST (global)	6.6	1.8	0.90	0.050	115	0.195

Table A1: Summary of the performance metrics calculated for the SSH mapping methods compared in this study. Methods above the double horizontal line are from the literature (see the text for references) and those below are those proposed in this study. See section 6.1 for a description of each metric.

826 Appendix B Open Research

827 The Level 3 satellite altimetry (<https://doi.org/10.48670/moi-00146>) and AOML
828 global drifter (<https://doi.org/10.17882/86236>) data used in this study are freely
829 publicly available from the CMEMS data store. The GHRSSST MUR Level 4 SST prod-
830 uct is freely publicly available from the NASA Earthdata PODAAC ([https://doi.org/](https://doi.org/10.5067/GHGMR-4FJ04)
831 [10.5067/GHGMR-4FJ04](https://doi.org/10.5067/GHGMR-4FJ04)). The Ocean Data Challenge data used in this study to compare
832 our method to the other mapping methods ([https://doi.org/10.24400/527896/a01](https://doi.org/10.24400/527896/a01-2021.005)
833 [-2021.005](https://doi.org/10.24400/527896/a01-2021.005)) were developed, validated by CLS and MEOM Team from IGE (CNRS-UGA-
834 IRD-G-INP), France and distributed by Aviso+. Code to generate training examples from
835 these public datasets, to define and train our ConvLSTM SSH and ConvLSTM SSH+SST
836 mapping methods, and to reproduce the results figures in this manuscript is publicly avail-
837 able here: <https://github.com/smartin98/deep-learning-ssh-mapping-JAMES-paper>.
838 Also provided in this repository are the underlying data for Table A1 and .nc files con-
839 taining the SSH maps we generated for the Gulf Stream study region for the year 2017.

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