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1 A new mechanism for brittle failure in garnets

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40 **Highlights**

- 41 • High-spatial resolution microscopy is used to study garnet deformation mechanisms
- 42 • Atom probe tomography reveals Fe-rich nanoclusters at high-angle grain boundaries
- 43 • Fe-rich nanoclusters lead to strain hardening and mechanical failure in garnet

44

45 **Abstract**

46 Garnet is a high-strength mineral and preserves structures that can consequently be used to
47 understand the flow strength and evolution of stress within the lower crust. Yet, the deformation
48 mechanisms at the brittle-ductile transition of garnet remain ambiguous. Here, we study garnet
49 porphyroclasts from an eclogite facies mylonite (central Australia) to investigate the mechanisms
50 by which garnet is deformed under relatively dry, lower crustal conditions. Electron backscatter
51 diffraction analysis reveals bands of small, relatively strain-free garnet with scattered orientations,
52 outlined by polygonal to lobate high-angle grain boundaries cross-cutting the garnet
53 porphyroclasts. Atom probe tomography of a high-angle grain boundary shows Fe enrichment in
54 the form of planar and equally spaced arrays of Fe-rich nanoclusters. Our experiments demonstrate
55 Fe segregation along grain boundaries of garnet, resulting in the nucleation of Fe-rich nanoclusters
56 that can act as barriers for migrating dislocations which leads to strain-hardening that facilitates
57 mechanical failure. The occurrence of strain-hardening in garnet potentially contributes to crustal
58 strengthening that can lead to seismicity at depths.

59

60 **Keywords**

61 garnet, atom probe tomography, strain-hardening, nanostructures, crystal-plasticity

62

63 **1. Introduction**

64 The mechanical behaviour of the lithosphere and the variation of strength with depth has been an
65 enduring topic in tectonics. It is generally accepted that the yield strength of the lithosphere at
66 geological timescales can be simplified using the "Christmas tree" envelope model (Goetze and
67 Evans, 1979). In this model, the lithosphere is assumed to be a single uniform layer with an upper
68 brittle part that follows a law of static friction (i.e., Byerlee's law) and a lower ductile part that
69 follows a viscous flow law (Byerlee, 1978; Wang, 2021). In truth, the lithosphere consists of
70 multiple layers where brittle and ductile behaviour alternate (e.g., Bürgmann and Dresen, 2008;
71 Kohlstedt et al., 1995; Ellis and Wang, 2022). Other more sophisticated forms of the strength
72 profile for the continental lithosphere have been proposed including the well-known "jelly
73 sandwich" and "crème brûlée" models (Burov and Watts, 2006; Jackson, 2002). Such models are
74 essential for geodynamic modelling and reconstructing the evolution and stabilization of
75 continents (e.g., Audet and Bürgmann, 2011). Most of these models, however, assume a brittle
76 upper layer and a ductile middle to lower layer separated by a steady-state transition zone.
77 Consequently, the lower continental crust is thought to be too hot to deform by brittle fracturing,
78 and frictional sliding at great depth is commonly considered to be a paradox (Green and Houston,
79 1995; White, 2012). Despite these assumptions, seismicity indicators such as pseudotachylytes in
80 deeply exhumed rocks demonstrate earthquakes do occur at depths (>50 km; e.g., Austrheim and
81 Boundy, 1994). Since the macroscopic rheology of the lower crust is governed by the microscopic
82 mechanical behaviour of its component minerals, microstructures within these minerals are often
83 used to infer its mechanical behaviour (e.g., Kavner et al., 2007; Mainprice et al., 2004; Phillips
84 and Ji, 2021). Under lower crustal conditions, most rock-forming minerals such as quartz and
85 feldspars deform by ductile creep, which can obscure microstructures associated with seismic

86 events (Kirkpatrick and Rowe, 2013). Garnet, notably, is considered mechanically and chemically
87 stable under elevated temperatures and commonly preserves microstructures related to seismicity
88 (e.g., Mancktelow et al., 2022; Trepmann and Stöckhert, 2002). Although garnet is not the most
89 predominant mineral in the lower crust, it becomes increasingly important with depth constituting
90 a large modal proportion in metamorphosed rocks (up to 50%; Villaseca et al., 1999) and is
91 assumed to play a significant role in crustal strengthening (e.g., Ingrin and Madon, 1995; Jin et al.,
92 2001; Ji et al., 2003). Therefore, understanding the style and nature of garnet deformation under
93 natural conditions is important for characterizing the flow strength and the stress history of the
94 lower crust (Mainprice et al., 2004).

95 Although regarded as a high-strength mineral, the pressure-temperature conditions in the
96 lower crust often exceed the conditions necessary for the onset of crystal-plastic behavior in garnet.
97 Experimental deformation studies have showed that under typical strain rates for actively
98 deforming regions (10^{-12} – 10^{-15} s⁻¹; Behr and Platt, 2011), the onset for crystal-plasticity in garnet
99 occurs at differential stresses of a few gigapascals and temperatures of 640–750°C (Karato et al.,
100 1995; Wang and Ji, 1999). Some studies have consequently attempted to link the brittle behaviour
101 of garnet at elevated temperatures to transient high fluid or melt pressures and fluid-assisted
102 eclogitization of granulites (Austrheim and Boundy, 1994; Hacker et al., 2003; Rempel et al.,
103 2006). Fracturing in garnet adjacent to pseudotachylites in deep-seated rocks has also been
104 interpreted to be related to thermal shock due to frictional heating (Papa et al., 2018). Under dry,
105 lower crustal conditions, fracturing may be related to transient high differential stresses from
106 seismic rupture (Giuntoli et al., 2018; Hawemann et al., 2018, 2019a). Despite these efforts, the
107 mechanisms of brittle fracture in garnet at elevated temperatures remain equivocal.

108 To determine the specific mechanisms by which garnet may be weakened under relatively
109 dry, lower crustal conditions, we applied correlated microscopy using 2D and 3D structural and
110 geochemical investigation at the near-atomic scale on a sample from a high-strain zone that formed
111 within the base of the crust (35-40 km depths) of the Proterozoic Musgrave Province, central
112 Australia (Gray, 1978; Camacho and Fanning, 1995). The well constrained pressure, temperature,
113 and time history of the Musgrave Province make it an excellent field location to investigate the
114 deformation mechanisms at the brittle-ductile transition for garnet (Gray, 1978; Camacho and
115 Fanning, 1995; Camacho et al. 1997). We document enhanced diffusion of Fe along grain
116 boundaries of recrystallized garnet, which apparently led to the nucleation of Fe-rich nanoclusters
117 that may have acted as barriers for migrating dislocations. Once these barriers were established,
118 the garnet underwent localized hardening and mechanical failure. Our data thus unequivocally
119 establish the brittle behaviour of garnet during deformation in the lower crust.

120

121 **2. Materials and Methods**

122 *2.1 Geological setting and sample description*

123 The Musgrave Province is a Mesoproterozoic–Neoproterozoic granulite to amphibolite facies
124 terrain located in central Australia that was formed by the amalgamation of different cratonic
125 blocks during the Musgravian Orogeny (ca. 1120–1200 Ma; **Figure 1**; Gray, 1978; Camacho and
126 Fanning, 1995). Granulite facies metamorphism at this time produced garnet-bearing
127 quartzofeldspathic gneisses of peraluminous and intermediate composition (Camacho et al., 1997).
128 The region was later heterogeneously overprinted during the Petermann Orogeny (ca. 550 Ma), by
129 crustal-scale faults and high-strain shear zones including the Woodroffe Thrust, the Mann Fault,
130 and the Davenport Shear Zone (DSZ; **Figure 1**; Collerson et al., 1972; Major, 1973; Camacho and

131 Fanning, 1995; Hawemann et al., 2018, 2019b; Wex et al., 2019). The eclogite facies DSZ is a W–
132 E-striking strike-slip shear zone, with a near horizontal stretching lineation and strongly localized
133 deformation (Camacho et al., 1997; Hawemann et al., 2019b). Peak conditions for the eclogite
134 facies deformation were estimated at ~650°C and ~1.2 GPa (Ellis and Maboko, 1992; Camacho et
135 al., 1997) under relatively dry conditions (<0.002 wt% H₂O; Wex et al. 2019). The sample studied
136 herein is a quartzofeldspathic mylonite from the DSZ with relict granulite facies garnet
137 porphyroclasts (**Figure 1, 2**). Within the sample, granulite facies garnet occurs as bands of
138 millimetre-scale porphyroclasts (0.5–3.0 mm) that are parallel to the main planar foliation and as
139 smaller (5–40 µm) neocrystallized eclogite facies porphyroblasts around relict garnets (**Figure 2a,**
140 **b**; Camacho et al., 1997). Garnet porphyroclasts are strongly fractured, which occur as regularly
141 spaced structures that are nearly perpendicular to the main foliation (**Figure 2a–c**). These fractures
142 do not propagate into the host rock matrix. Retrograde biotite (<5%) occurs between garnet clasts
143 and within fractures (**Figure 2a–c**). Quartz occurs either as fine grains (<50 µm) with serrated
144 boundaries (**Figure S1a**) or slightly larger elongated grains (up to 100 × 200 µm) exhibiting
145 undulose extinction, deformation bands, and new subgrains forming around the edges (**Figure**
146 **S1a**). Quartz also occurs as ribbons, defining the foliation (30 × 100 µm), with undulose extinction
147 and polygonal subgrain development (**Figure S1b**). Comparatively, feldspar occurs as medium to
148 coarse (0.2–0.5 mm × 0.5–2.5 mm) elongated lenticular porphyroclasts (**Figure S1c–d**). Feldspar
149 porphyroclasts are oriented parallel to the main foliation and exhibit undulose extinction, folded
150 and tapered twinning, and new subgrain development around the edges (**Figure S1c–d**).

151

152 *2.2 2D microscopy*

153 Primary characterization of the garnet porphyroclasts was performed by combining electron
154 backscatter diffraction (EBSD) mapping and high-contrast backscattered electron (BSE) imaging
155 to obtain quantitative and qualitative information of the crystalline structure and deformation state
156 of the sample. All crystallographic data were collected at the Max-Planck-Institut für
157 Eisenforschung GmbH (Düsseldorf, Germany) using a high-resolution field emission SEM Sigma
158 (Carl Zeiss Microscopy, Germany) equipped with a Hikari (EDAX, USA) EBSD detector. Prior
159 to analysis, the thin section was chemo-mechanically polished with an alkaline colloidal silica
160 suspension and carbon-coated to establish conductivity. EBSD mapping was conducted first with
161 analytical conditions set to an accelerating voltage of 15 kV and probe currents of 7.4 nA. The
162 sample was tilted to 70° and brought to a working distance of 17.4 mm. Mapping was performed
163 by using a hexagonal grid and a step size of 1.5 µm; the camera was operated by applying a 4 × 4
164 binning. The EBSD data was cleaned using the OIM Data Collection and Analysis software by
165 combining the neighbour orientation correlation (clean-up level: 3) and confidence index
166 standardization (CI: 0.1) functions. Misorientation deviation angle (misorientation of each pixel
167 relative to grain average) and grain boundary maps were then created using the OIM software.
168 Based on the EBSD data, highly-strained garnet porphyroclasts were selected for further analysis
169 and are the main subject of this study.

170 High-contrast BSE imaging was subsequently conducted to further investigate the targeted
171 microstructures. This imaging was performed on the Zeiss SEM Merlin (Carl Zeiss Microscopy,
172 Germany) instrument at the Max-Planck-Institut für Eisenforschung GmbH. Analytical conditions
173 for imaging were set to a 30 kV accelerating voltage, a 2.0 nA probe current and a working distance
174 of 6.5 mm, allowing for the direct observation of crystal defects.

175 Energy-dispersive X-ray spectroscopy (EDX) was also performed on targeted grain
176 boundaries on the Zeiss SEM Sigma using analytical conditions of 15 kV, probe currents of 7.4
177 nA, and a working distance of ~7.8 mm. EDX maps were collected by setting the counting time to
178 100 s and capturing 128-256 frames. Acquired EDX maps were subsequently processed using the
179 EDAX APEX Data Collection and Analysis software.

180 An electron microprobe (EMP) was used to chemically characterize the garnet
181 porphyroclasts. Qualitative X-ray mapping was performed using the Cameca SX-100 at the
182 University of Manitoba (Canada). Mapping for Fe K α , Mg K α , Mn K α , and Ca K α were done at a
183 20 kV accelerating voltage, a 50 nA beam current, 1 μ m beam size and 50 ms dwell time.

184

185 *2.3 Atom probe tomography*

186 A specimen was prepared from a high-angle grain boundary within the garnet by in situ lift-out
187 (Thompson et al., 2007) for atom probe tomography (APT) in an attempt to investigate the detailed
188 composition of some of the microstructural features at the near-atomic scale. The needle-shaped
189 specimen was sharpened by annular milling at 30 kV on a dual-beam scanning electron
190 microscope/focused ion beam FEI Helios Nanolab 600i at the Max-Planck-Institut für
191 Eisenforschung GmbH and subsequently cleaned using the Ga-beam at 5 kV to remove regions
192 potentially severely damaged by the implantation of energetic Ga ions. The specimen was then
193 analysed by APT in a Cameca LEAP 5000 XR (**Table S1**) fitted with a reflectron-lens with a
194 detector efficiency of ~52%. The specimen was analysed in laser pulsing mode at 50 K with a laser
195 pulse energy of 80 pJ focused on an area estimated to be <3 μ m in diameter, a detection rate of 0.5
196 ion detected for 100 pulses and a laser pulse repetition rate of 125 kHz. The data processing and
197 reconstruction were done with the commercial software package AP Suite 6.1. The compositional

198 profile presented herein was generated from a cylindrical region of interest using the 1D
199 concentration profile function and computed using the fixed bin width profile, atomic, and
200 decompose complex ions options. The bin size was adjusted to 0.05 nm to maximize the ratio of
201 the peak signal to the statistical fluctuations. The ranged mass spectrum used for the specimens is
202 shown in **Figure S2**.

203

204 **3. Results**

205 EMP X-ray maps of the garnet porphyroclasts reveal compositional variations across the grains
206 (**Figure 3**). Relative to the core, the rims of the porphyroclasts have a Ca-rich and Fe- and Mg-
207 depleted composition, similar to that of the neocrystallized garnet grains (**Figure 3**).

208 EBSD analysis of the garnet porphyroclasts reveals minor evidence for crystal-plastic
209 deformation within the brittle-dominated grain fragments. Inverse pole figure (IPF) maps of two
210 large ($220 \times 225 \mu\text{m}$; $175 \times 195 \mu\text{m}$) porphyroclasts reveal a cross-cutting band of smaller grains
211 ($1\text{--}10 \mu\text{m}$ in diameter) exhibiting a wide distribution of crystal orientations (**Figure 4a, 5**). The
212 IPFs for the large grains reveal minor dispersion relative to the main orientation (**Figure 5a–c**)
213 whereas small grains show scattered and random orientations, and therefore no obvious systematic
214 misorientation relationship to the orientation of the adjacent larger grains (**Figure 5d–e**). Deviation
215 angle and grain boundary maps reveal lattice distortions within the large grains as heterogeneous
216 misorientation patterns with a maximum misorientation of 36° relative to the grain average (**Figure**
217 **4b**), and low-angle grain boundary (LAGB) development at the large grain rims (**Figure 4c**).
218 Within the larger grains, the misorientation angle of the LAGBs gradually increase from $2\text{--}5^\circ$ in
219 the grain interior to $5\text{--}15^\circ$ near the rim (**Figure 4c**). This contrasts with the smaller grains that
220 exhibit almost no evidence for intragranular lattice distortion and are primarily bordered by high-

221 angle grain boundaries (HAGB; **Figure 4b–c**). Note, regularly spaced fractures completely
222 overprint the misorientation patterns within the garnet clasts and do not appear related to crystal-
223 plastic microstructures.

224 High-contrast BSE imaging also shows the nearly parallel and regularly spaced fractures
225 overprinting the porphyroclasts (**Figure 4d**), yet these fractures do not propagate into the host
226 quartzofeldspathic matrix. Smaller (1–10 μm) subgrains within the larger clasts are apparent as
227 greyscale variations at the rims of the porphyroclasts (**Figure 4d–f**). Small grains within the band
228 display polygonal to slightly lobate grain boundaries (**Figure 4e–f**) with localized microfracturing
229 along the HAGBs (**Figure 3d–f**). HAGBs display noticeably brighter contrasts, indicative of
230 segregation of a heavier element, confirmed by EDX mapping revealing Fe enrichment and Mg
231 depletion (**Figure 6**).

232 APT analysis from a HAGB (**Table S1**) reveals a homogeneous distribution of the major
233 garnet components (8.28 at% Mg; 2.16 at% Ca; 11.04 at% Al; 11.30 at% Si, 48.68 at% O), except
234 for Fe (12.87 at%; **Table 1**). Iron is homogeneously distributed in region 1, with a composition of
235 8.40 at%, whereas in region 2 it makes up 18.79 at% of the bulk composition. Moreover, it forms
236 regular planar arrays of near-spherical nanoclusters with diameters of 3.0–7.5 nm and spacing of
237 2.0–7.5 nm between the clusters. The distance between the planar arrays is 4.5–6.5 nm. The
238 average composition of the clusters, extracted from Fe isosurfaces (isovalue: 0.34), reveals mostly
239 Fe enrichment (49.69 at%) with the other garnet components making up the remaining
240 composition (**Figure 7, Table 1**). The Fe-rich nanoclusters are also highlighted by the distribution
241 of Ga ions within the specimen (**Figure S3**).

242

243 **4. Discussion**

244 *4.1 Deformation mechanisms of garnet*

245 The microstructures documented by EBSD and BSE analyses suggest the garnet porphyroclasts
246 deformed dominantly by brittle to brittle-plastic mechanisms. Plastically deformed regions with
247 the highest degree of misorientation and the development of LAGBs within the garnet are spatially
248 related to the rims of porphyroclasts (**Figure 4**). In the semi-brittle regime of deformation,
249 fracturing can generate dislocations; therefore, fracturing and dislocation glide processes can be
250 superimposed and act reciprocally (Rogowitz et al., 2018; McLaren and Pryer, 2001; FitzGerald
251 et al., 1991). Since plastic deformation in our sample is concentrated at the rims of porphyroclasts,
252 fracturing during cataclasis of the granulite facies porphyroblasts likely played an important role
253 in generating dislocations and facilitating the onset of crystal-plasticity during eclogite facies
254 deformation. Although assumed to be a high-strength mineral in crustal strength profiles (Karato
255 et al., 1995; Wang and Ji, 1999), both brittle and crystal-plastic behaviour have been reported in
256 naturally deformed garnet (e.g., Austrheim et al., 2017; Hawemann et al., 2019a). Experimental
257 deformation studies on garnet conclude that differential stresses on the order of a few gigapascals
258 are required to produce shear fractures in the mineral, and the onset of crystal-plastic deformation
259 for natural strain rates (10^{-12} – 10^{-15} s⁻¹; Behr and Platt, 2011) should only occur at temperatures
260 above 640–750°C (Karato et al., 1995; Wang and Ji, 1999). The estimated peak temperature for
261 our sample (~650°C) falls within the range of the proposed transition from brittle to crystal-plastic
262 deformation. Our micro- and nanostructural data are consistent with previous studies evincing
263 brittle and crystal-plastic deformation of garnet under such temperature conditions within
264 mylonitic shear zones related to seismic stresses (Hawemann et al., 2019a).

265 EBSD IPF maps and plots (**Figure 4a, 5**) reveal a band of 1–10 µm, relatively strain-free
266 grains with scattered orientations. Progressive subgrain rotation due to dislocation glide and creep

267 processes is an effective mechanism to generate new, small grains during dynamic recrystallization
268 at moderate homologous temperatures (Hobbs, 1968). During recrystallization, there is a
269 systematic relationship, termed host-control, between host grains, progressively rotated subgrains,
270 and recrystallized grains (Halfpenny et al. 2006; Stünitz et al., 2003). In our sample, the LAGBs
271 within the larger grains reveal gradually increasing misorientation angles towards the rims. The
272 same systematic relationship can be observed in their respective IPF plots which show minor
273 dispersion around the main grain orientation. However, within the fine-grained band, there is no
274 systematic relationship between the larger (host) grains and the smaller grains. Therefore, subgrain
275 rotation recrystallization alone cannot completely explain the band of small, strain-free grains in
276 our sample.

277 Instead, we propose that at eclogite facies conditions the small grains were derived from
278 fragments of the host granulite grains due to comminution during cataclastic deformation, and
279 underwent various degrees of rotation. Fragments generated by cataclasis, typically have irregular
280 shapes (Stünitz and FitzGerald, 1993), yet in our sample, these small grains display polygonal to
281 slightly lobate grain boundaries, suggesting that the grain boundaries were mobile. Grain boundary
282 mobility can occur due to either the high surface energy of small grains, which can drive normal
283 grain growth and solid-state recrystallization, or the difference in dislocation density between the
284 neighbouring fractured and non-fractured zones within the grains, leading to dynamic
285 recrystallization by bulging or grain boundary migration (GBM; Stünitz et al., 2003; Tullis and
286 Yund, 1992). Polygonal grain boundaries within the fine-grained band (**Figure 4e, f**) corroborate
287 that normal grain growth occurred during metamorphism. Although GBM could explain the lobate
288 grain boundaries, the low homologous temperature experienced by the sample combined with the

289 lack of crystal-plasticity unrelated to brittle structures, make GBM highly unlikely. If dynamic
290 recrystallization did occur, then it most likely did so by bulging.

291 Based on these combined observations, we suggest that the small, relatively strain-free
292 grains represent former fragments of the host porphyroblast, which acted as nuclei for grain growth
293 at during the eclogite facies deformation. The nucleation of new grains by cataclasis and
294 subsequent growth of very small fragments leading to the non-host control orientation of
295 recrystallized grains has also been documented in plagioclase and quartz deformed at low
296 homologous temperatures (Hirth and Tullis, 1992; Stünitz et al., 2003).

297 Moreover, the large set of fractures overprinting the garnet porphyroclasts suggest that the
298 garnet underwent at least two stages of brittle deformation (**Figure 1, 2**). The first being the
299 cataclastic deformation of granulite facies porphyroblasts, which is related to crystal-plasticity
300 indicating it occurred at the peak temperatures. The second stage formed the regularly spaced
301 fractures, which occur nearly perpendicular to the host rock foliation. The lack of crystal-plastic
302 microstructures associated with these fractures, and the presence of retrograde biotite infilling the
303 open space, suggest this stage of brittle deformation occurred after the peak temperatures had
304 dropped below the brittle to crystal-plastic transition for garnet (640–750°C; Karato et al., 1995;
305 Wang and Ji, 1999). The development of regularly spaced opening-mode fractures occurring
306 perpendicular to the garnet layer and terminating at the interface between garnet and the softer
307 quartzofeldspathic host rock is consistent with micro-jointing reported in layered materials (Bai et
308 al., 2000). A study on garnet from the same location in the Musgrave Province has interpreted this
309 second generation of extensional fractures to be related to exhumation of the eclogite facies rocks
310 (Hawemann et al., 2019a).

311

312 *4.2 Nature and composition of nanoclusters*

313 The Fe-rich composition of the nanoclusters was confirmed by estimating the field-strength
314 conditions during field evaporation through the Kingham analysis method (**Figure S4**; Kingham,
315 1981). Based on the ratio of Al^{3+} to Al^{2+} (0.08) in the mass spectrum, the field conditions for the
316 nanoclusters can be estimated at 40 Vnm^{-1} . Under such field conditions, and even though the
317 Kingham analysis can only provide an estimate, the field evaporation and post-ionization of Si
318 causes the ions to be detected only in the double charge state, therefore minimizing any overlap at
319 28 Da and confirming the Fe-rich composition of the nanoclusters.

320 The Fe-rich nanoclusters are also highlighted by the distribution of Ga ions in the specimen
321 (**Figure S3a**). In the 3D reconstruction, Ga (peak at 69 Da) is concentrated both at the tip of the
322 specimen and along the Fe nanoclusters. No other combination of elements could explain the peak
323 observed at 69 Da apart from the implantation of the monoisotopic Ga used in the FIB preparation.
324 This observation is consistent with previous studies showing that when Ga is mobile in the
325 structure, it tends to segregate along pre-existing defects during FIB-based specimen preparation
326 (Gault et al., 2018). Intentional segregation of Ga has even been used in the past to reveal specific
327 phases inside a specimen, including far from the specimen's topmost surface (Ruan et al. 2011).
328 In addition, the 3D reconstruction of another specimen from the same lift-out also shows Ga
329 enrichment at the specimen surface, but no Fe nanoclusters are observed despite the penetration of
330 the ions (**Figure S3b**). The presence of Ga implantation and the absence of clusters in the second
331 specimen suggests the Fe nanoclusters did not result from ion implantation during specimen
332 preparation. Moreover, the diffusion of Fe in garnet at room temperature is not expected to be fast
333 enough to form such clusters as a result of ion implantation or the associated change in the
334 material's composition (Li et al. 2018). Finally, the accelerated diffusion of Ga on the half where

335 clusters are present could be explained by a combination of ion channeling and the Ga affinity for
336 Fe.

337

338 *4.3 Origin of nanocluster arrays*

339 EDX mapping of the HAGBs within the fine-grained band shows Fe enrichment and Mg depletion
340 (**Figure 6**), which is confirmed in the APT data revealing Fe enrichment in the form of planar and
341 nearly equally spaced arrays of Fe-rich nanoclusters (**Figure 7**). Although many studies have
342 previously reported elemental differences at LAGBs and HAGBs in garnets (e.g., Chapman et al.
343 2017, 2019a, 2019b; Konrad-Schmolke et al. 2007; Tacchetto et al. 2022), none have documented
344 nanocluster arrays. Such a regular arrangement of nanoclusters is strikingly similar to those
345 documented in some metallic alloys (e.g., Yen et al., 2011; Seol et al., 2017). In those instances,
346 cooling of the designed alloy transforms it from a high-temperature face-centered cubic phase to
347 a low-temperature body-centred cubic phase. The phase transformation leads to nanoparticles
348 continuously nucleating and subsequently detaching from the migrating heterophase interface,
349 resulting in regular planar arrays of nanoparticles within the bulk material. The thermal treatment
350 of the metal allows for precise control of the dispersion of these property-enhancing precipitates,
351 and hence of the alloy's beneficial mechanical properties (e.g., yield strength, high-temperature
352 strength, and creep resistance; Ashby, 1958; Devaraj et al., 2016; Orowan, 1968; Vogel et al.,
353 2013). When densely arranged in the metallic matrix, these nanoparticles indeed act as effective
354 obstacles for dislocation movement apportioning only small volumes for dislocations to move
355 freely. Known as the Orowan effect, dislocations must either loop around or sweep through the
356 nanoparticles to accommodate strain, requiring relatively higher amounts of energy and therefore
357 stresses (Ashby, 1958; Orowan, 1968).

358 We propose that a very similar nanostructuring has occurred in the garnet at a mobile grain
359 boundary. The mobility of the HAGBs within the fine-grained band is driven by normal growth
360 and solid-state recrystallization during metamorphism (**Figure 8a**), as illustrated by the EBSD
361 data. Grain boundaries in crystalline materials have a high transport coefficient for diffusing atoms
362 and are the loci of solute segregation driven by the reduction in the system's free energy (Raabe
363 et al., 2014). Therefore, we suggest that Fe was segregated into the core of grain boundaries during
364 growth, leading to solute drag (**Figure 8b**; e.g., Luo et al., 2022), which slows down the grain
365 boundary, further facilitating high degrees of segregation. Upon reaching a critical Fe-
366 concentration at the grain boundary, phase separation occurred and Fe-rich nanoclusters nucleated
367 on the grain boundary, similar to localized spinodal decomposition (Kwiatkowski et al., 2018).
368 The continuous migration of the grain boundaries during growth leaves behind planar alignments
369 of equally sized and spaced nanoclusters (**Figure 8c**), forming regular planar arrays in the vicinity
370 of HAGBs (**Figure 8d**). Although the precipitates that form within doped alloys segregate due to
371 the inability of the trace elements to reside within the transformed phase, Fe is a major component
372 of garnet and should be stable within the crystal structure. EPMA X-ray maps of the garnet
373 porphyroclasts reveal Ca-enriched rims that are depleted in Fe and Mg (**Figure 3**). Camacho et al.
374 (2009) studied these same rocks and interpreted the chemical variations as a result of Ca diffusion
375 from the enriched host-rock into the garnet rims. The diffusion of Ca therefore reduced the
376 solubility of Fe and Mg in the crystal structure, leading to their diffusion out of the garnet.
377 Therefore, the uptake of Ca and the reduced solubility of Fe within the growing garnet grains can
378 explain the segregation of Fe into nanoclusters rather than its homogeneous distribution.

379 These regular arrays of Fe-rich nanoclusters in the garnet may have acted as barriers for
380 migrating dislocations. This process can lead to hardening, dislocation tangles, and facilitates

381 intragranular mechanical failure of the garnet porphyroclast through microfracture nucleation at
382 and propagation along grain boundaries (**Figure 8e, f**). The propagation of fractures along grain
383 boundaries can in turn promote the localization of crystal-plasticity creating a positive feedback
384 loop. The microfractures observed at HAGBs (**Figure 4e**) in conjunction with Fe-segregation
385 (**Figure 4b**) in our sample are consistent with this model.

386

387 **Conclusion**

388 We have combined 2D and 3D structural and chemical analyses on garnet porphyroclasts from the
389 relatively dry lower crust of the Musgrave Province (central Australia) to identify the micro- and
390 nanoscale processes associated with deformation in the brittle- to brittle-plastic regime of garnet.
391 EBSD and EDX data reveal a band of randomly oriented small grains outlined by Fe-rich HAGBs
392 formed by cataclasis and subsequent growth. APT data from one of these HAGBs reveals regular
393 planar arrays of Fe-rich nanoclusters. The combined data led us to propose a new strain hardening
394 model in garnet whereby the mobility of grain boundaries during growth led to the formation of
395 regular planar arrays of Fe-rich nanoclusters that act as barriers for migrating dislocations. The
396 arrest of the dislocations resulted in defect tangles and microfracture nucleation and propagation
397 at grain boundaries. Localized hardening in garnet in the vicinity of grain boundaries, as
398 documented herein, can therefore potentially contribute to the mechanisms of mechanical failure
399 within the presumed high-strength mineral. The occurrence of strain-hardening in garnet
400 potentially contributes to crustal strengthening in the lower crust, which could lead to deep
401 seismicity.

402

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410

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Table 1. Bulk composition analysis of garnet specimen

| Ion Type | Whole specimen | | | Region 1 | | | Region 2 | | | Clusters isosurfaces | | |
|--------------|----------------|----------|----------|----------|----------|----------|----------|----------|----------|----------------------|----------|----------|
| | Count | Atomic % | σ | Count | Atomic % | σ | Count | Atomic % | σ | Count | Atomic % | σ |
| H | 104727 | 5.24 | 0.06 | 53861 | 4.94 | 0.10 | 50299 | 5.49 | 0.07 | 3485 | 2.64 | 0.13 |
| O | 972809 | 48.68 | 0.13 | 575857 | 52.84 | 0.24 | 404072 | 44.07 | 0.15 | 47299 | 35.85 | 0.38 |
| Si | 225859 | 11.30 | 0.09 | 136093 | 12.49 | 0.17 | 88695 | 9.67 | 0.11 | 2698 | 2.04 | 0.21 |
| Ca | 43189 | 2.16 | 0.03 | 23826 | 2.19 | 0.06 | 19148 | 2.09 | 0.04 | 1536 | 1.16 | 0.08 |
| Mg | 165546 | 8.28 | 0.07 | 95885 | 8.80 | 0.13 | 68594 | 7.48 | 0.07 | 4867 | 3.69 | 0.13 |
| Al | 220509 | 11.04 | 0.08 | 109755 | 10.07 | 0.14 | 108385 | 11.82 | 0.09 | 6581 | 4.99 | 0.16 |
| Mn | 7910 | 0.40 | 0.02 | 2669 | 0.24 | 0.04 | 5243 | 0.57 | 0.03 | 0 | 0.00 | 0.00 |
| Fe | 257151 | 12.87 | 0.09 | 91576 | 8.40 | 0.15 | 172258 | 18.79 | 0.12 | 65558 | 49.69 | 0.40 |
| In | 390 | 0.02 | 0.00 | 290 | 0.03 | 0.01 | 91 | 0.01 | 0.01 | 0 | 0.00 | 0.00 |
| Cr | 174 | 0.01 | 0.00 | 67 | 0.01 | 0.01 | 99 | 0.01 | 0.00 | 3 | 0.00 | 0.01 |
| Total | 1998264 | 100 | | 1089879 | 100 | | 916884 | 100.00 | | 132027 | 100 | |

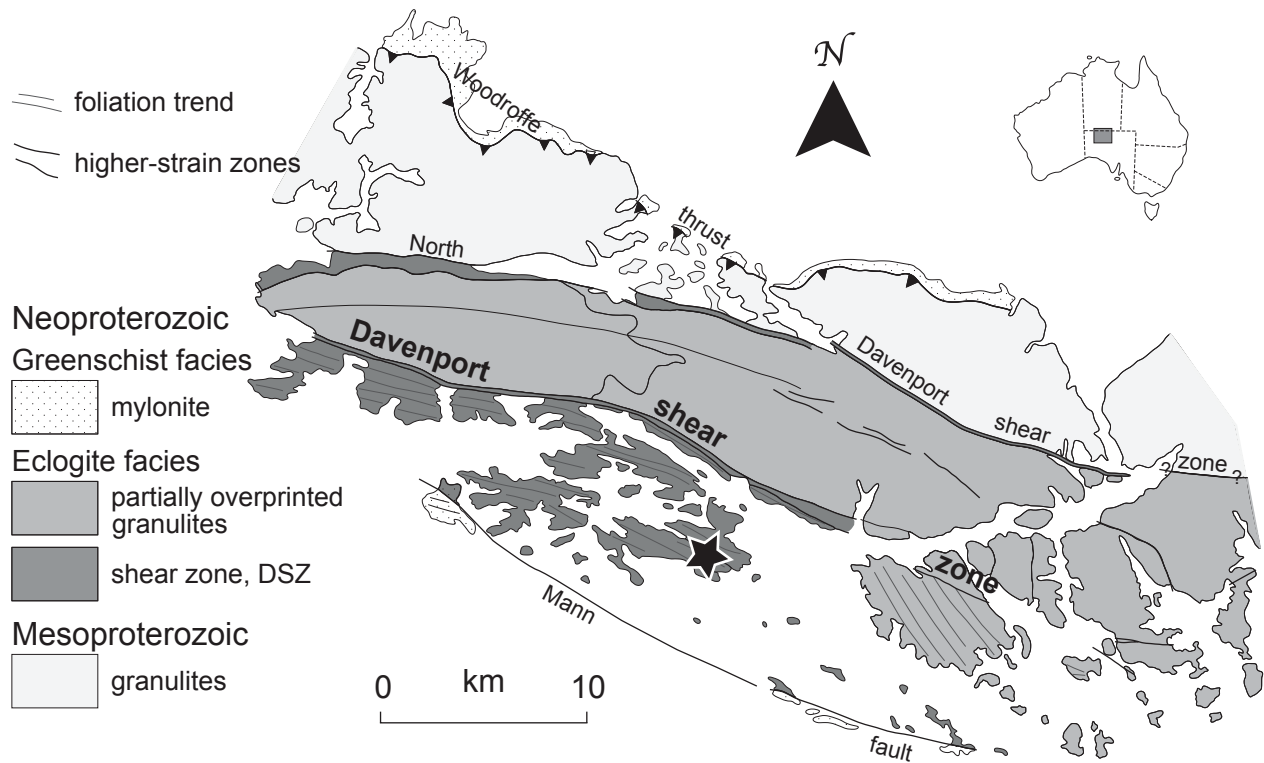


Figure 1. Simplified geological map of Mount Woodroffe region, South Australia, showing the sample location (black star; S26° 24.386' E131° 39.171' Datum WGS84) between the Davenport shear zone (DSZ) and the Mann fault. Modified after Camacho et al. (2009).

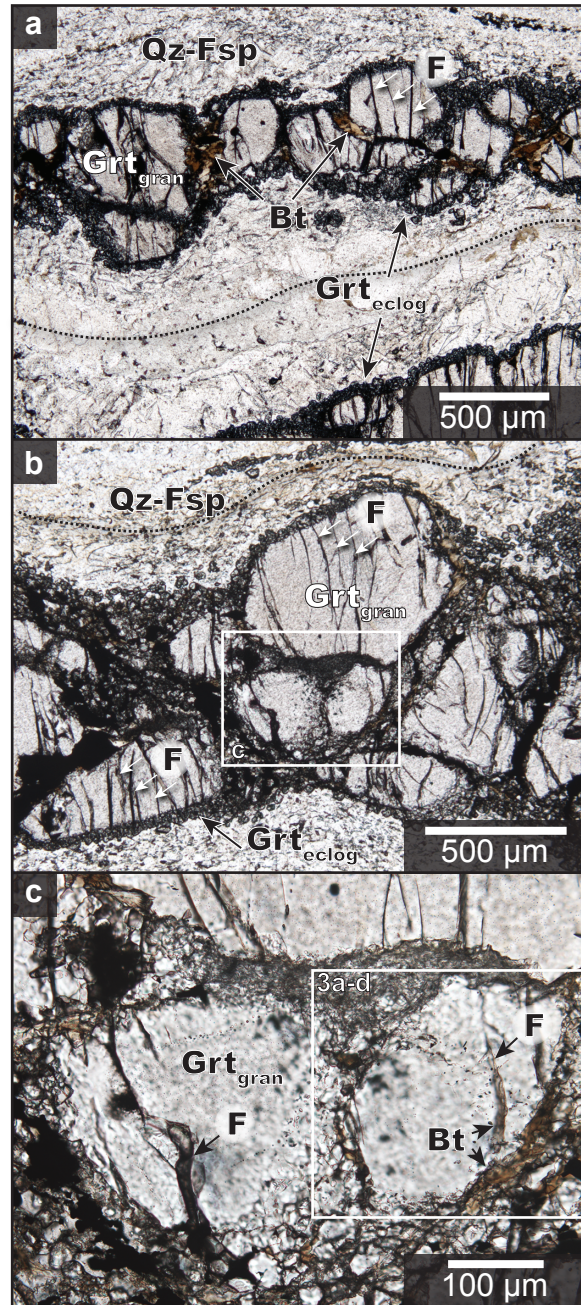


Figure 2. a-c Photomicrographs (PPL) of the investigated garnet porphyroclasts hosted within a mylonitic Qz-Fsp matrix, Musgrave Province, central Australia. Granulite facies garnet porphyroclasts (Grt_{gran}) form foliation parallel bands and are overgrown by small neocrystallized eclogite facies garnet porphyroblasts (Grt_{eclog}). Both generations of garnet are overprinted by regularly-spaced and foliation perpendicular fractures (F). Foliation is indicated by the dashed black lines. Biotite occurs between garnet porphyroclasts and fractures. White box in b shows the location of image c. White box in c shows the location of Figure 3a-d.

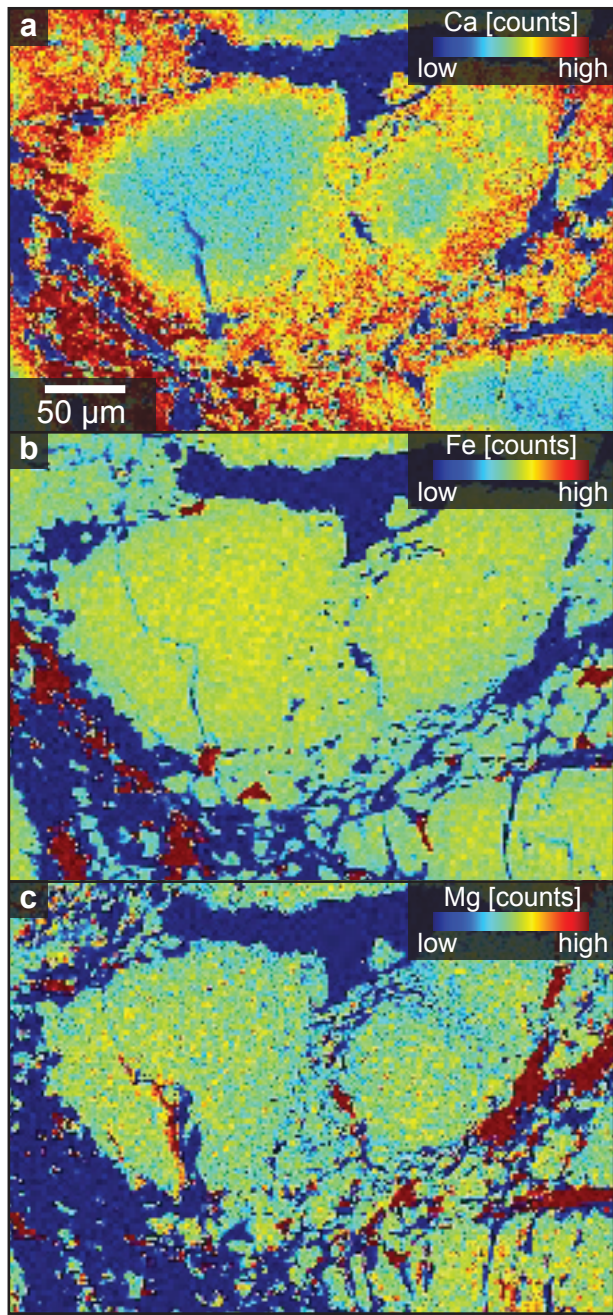


Figure 3. EPM X-ray map for Ca (a), Fe (b) and Mg (c) showing compositional variations along the rims the garnet prophyroclasts shown in Figure 2c. Clast rims are enriched in Ca and depleted in Fe and Mg relative to cores.

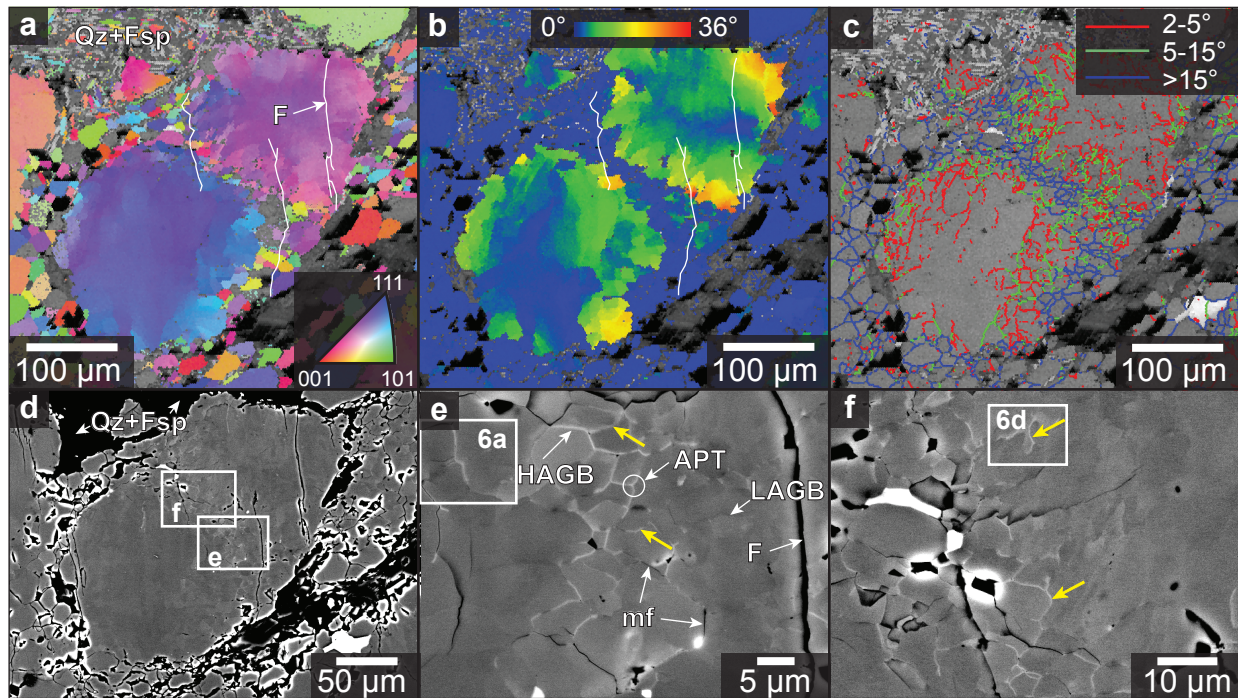


Figure 4. **a** IPF map of garnet porphyroclasts showing the band of smaller grains with scattered orientations between two larger grains. Fractures from BSE image are traced in white (F). **b** Orientation deviation angle map revealing a heterogeneous misorientation pattern within the larger grains with higher misorientation along the rims (max. misorientation of 36°) and relatively no misorientation within the smaller grains. **c** Grain boundary map evincing the development of LAGBs with increasing misorientation angle towards the rims of the large grains and HAGBs within the fine-grained band. **d-f** High-contrast BSE images revealing LAGBs within the larger grains and HAGBs within the band of smaller grains. Small grains within the band exhibit polygonal to slightly lobate grain boundaries (yellow arrows). Note: microfracturing (mf) along HAGBs. Images (e) and (f) show the location of the APT specimen (e) and the EDX maps (e, f).

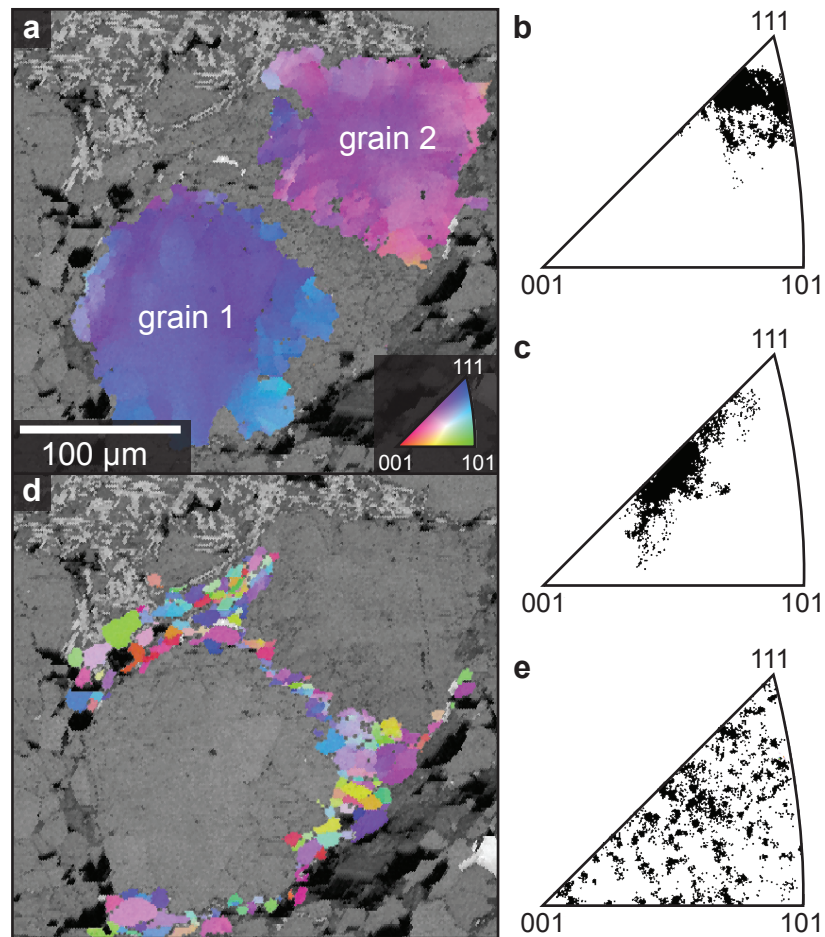


Figure 5. **a** IPF map of larger grains 1 and 2. **b–c** Lower hemisphere IPF plots of grains 1 (c) and 2 (d) showing minor dispersion relative to the main orientation of the grains. **d** IPF map of fine-grained band. **e** Lower hemisphere IPF plot of the fine-grained band showing the large distribution of orientations within the smaller grains with no relationship to the larger grains' orientations.

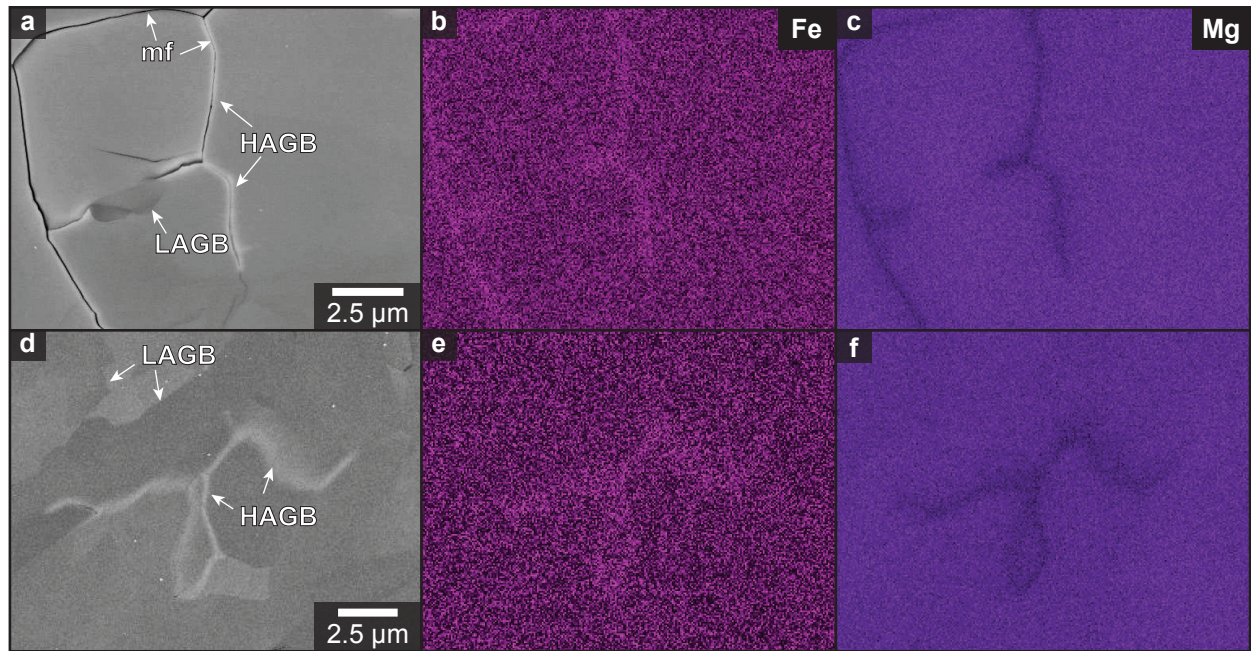


Figure 6. BSE images (**a, d**) and EDX maps of Fe (**b, e**) and Mg (**c, f**) within the garnet porphyroclast revealing Fe enrichment and Mg depletion at fractured (mf) and intact HAGBs.

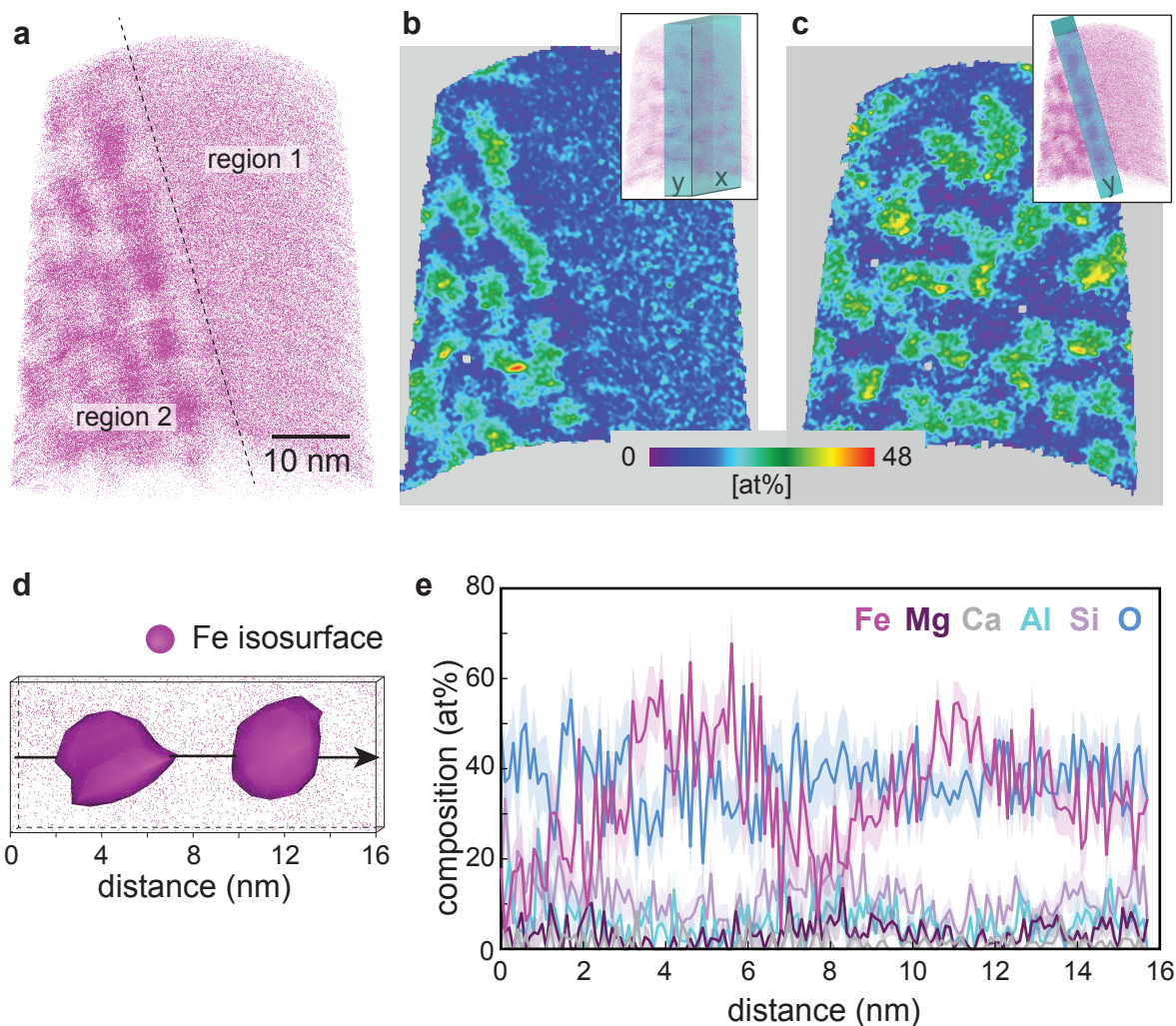


Figure 7. **a** 3D reconstruction of APT specimen revealing Fe-rich nanoclusters aligned into individual planes on the left side of the specimen (region 2). Note: Fe peaks at 28 and 56 Da are plotted in the figure. Details on ranging the peak at 28 Da can be found in Figure S2. **b** and **c** 2D contour plot showing the projection of Fe compositions along the y-axis of a region of interest perpendicular (**b**) and parallel (**c**) to the Fe-nanocluster planes. The regions of interest for (**b**) and (**c**) are shown as turquoise boxes in the inset APT reconstructions on the upper right corner of each image. **d** Clipping of region of interest surrounding two Fe-rich nanoclusters (pink isosurfaces; 34 at% isovalue). **e** Composition profile across the two clusters in (**d**) confirming Fe enrichment (54.7–67.7 at%) and slightly depleted Si, Mg, Al, and O compositions within the matrix between clusters.

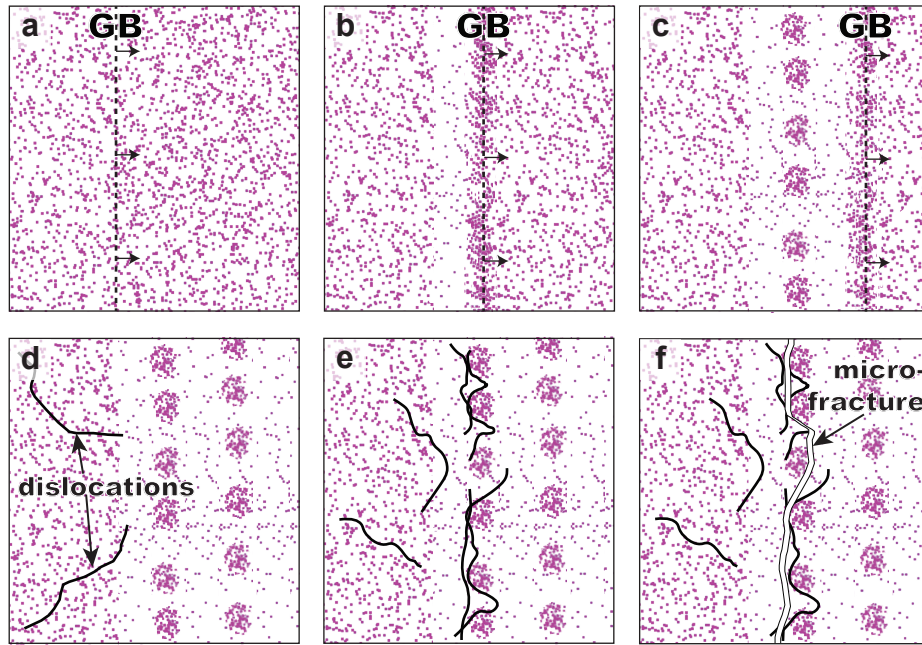


Figure 8. Conceptual model of strain hardening in garnet. **a** Initial state of grain boundary (GB) prior to migration during crystal growth. **b** Accumulation of Fe within the core of the migrating grain boundary leaving behind an Fe-depleted wake. **c** Nucleation and detachment of Fe nanoclusters after reaching saturation and continued crystal growth. **d** Introduction and migration of dislocations (thick black lines) during deformation. **e** Continued migration of dislocations and pinning at Fe nanoclusters, leading to entanglement. **f** Nucleation of microfracture at dislocation tangles.

Table S1. Atom probe tomography data acquisition settings and data summary

| Specimen/Data set | Specimen i |
|---|---|
| Instrument Model | LEAP 5000 XR |
| Instrument settings | |
| Laser wavelength (nm) | 355 |
| Laser pulse energy (pJ) | 80 |
| Voltage pulse fraction (%) | n.a. |
| Pulse frequency (kHz) | 125 |
| Evaporation control | detection rate |
| Target detection rate (ions/pulse) | 0.005 |
| Nominal flight path (mm) | 382 |
| Set point temperature (K) | 50 |
| Chamber pressure (Torr) | 4.66E-11 |
| Data summary | |
| LAS root version | 18.46.533c |
| CAMECAROOT version | 18.46.533c |
| Analysis software | AP Suite 6.1 |
| Total ions: | 3070985 |
| Single (%) | 63.7 |
| Multiple (%) | 34.8 |
| Partial (%) | 1.5 |
| Reconstructed ions: | 2371612 |
| Ranged (%) | 87.5 |
| Unranged (%) | 12.5 |
| Volt./bowl corr. peak (Da) | 32 |
| Mass calib. (peaks/interp.) | 5/Lin. |
| $(M/\Delta M)$ for ^{28}Fe | 933.3 |
| $(M/\Delta M_{10})$ | 50.9 |
| Time-independent background (ppm/ns) | 24 |
| Reconstruction | |
| Final specimen state | fractured |
| Pre-/post-analysis imaging | SEM/n.a |
| Radius evolution model | voltage |
| Field factor (k) | 3.3 |
| Image compression factor | 1.65 |
| Assumed E-field (V/nm) | 25 |
| Detector efficiency (%) | 52 |
| Avg. atomic volume (nm ³) | species specific volume used (see range file) |
| V_{initial} ; V_{final} (V) | 3140; 4125 |

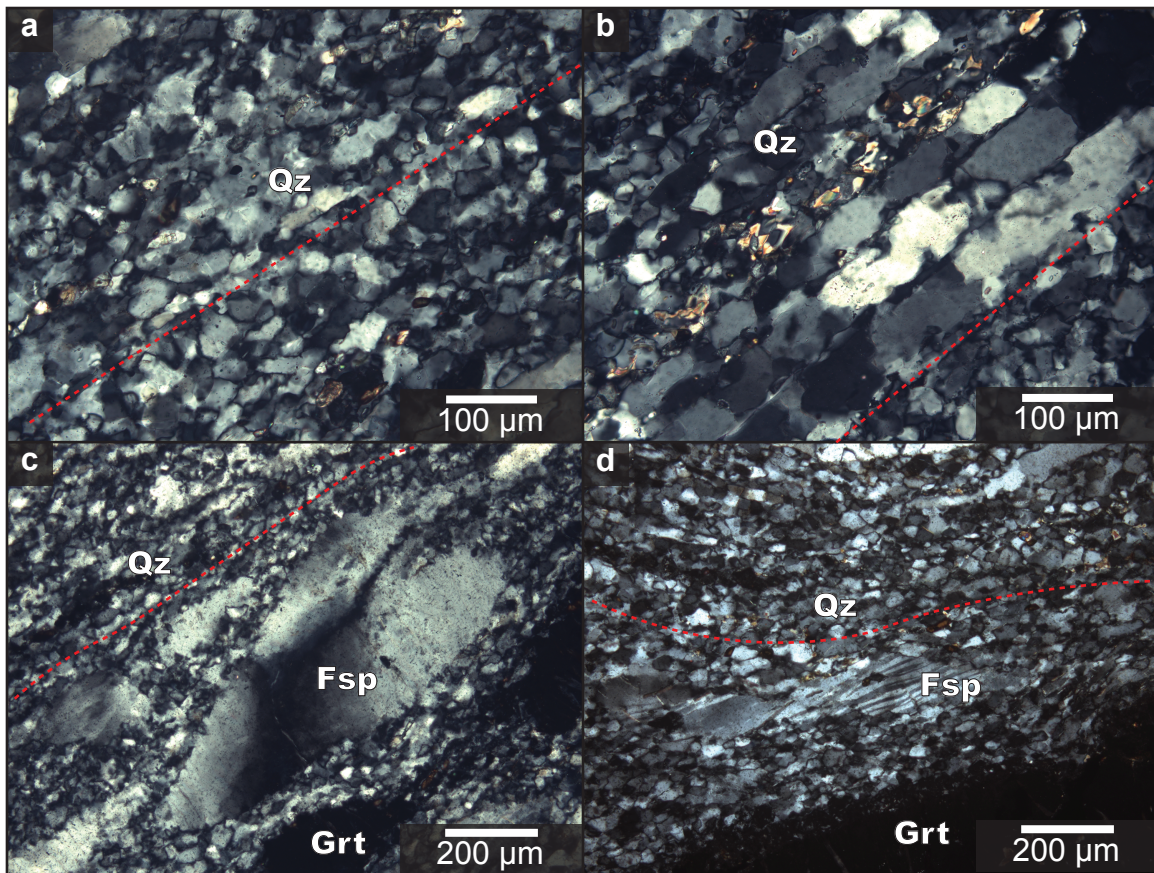


Figure S1. a-d Photomicrographs (XPL) of deformed quartzofeldspathic mylonite matrix. Main foliation is traced by dashed red lines. a Fine-grained ($<50 \mu\text{m}$) quartz with serrated grain boundaries and slightly larger elongated grains (up to $100-200 \mu\text{m}$) exhibiting undulose extinction, deformation bands, and new subgrains forming around the edges. b foliation parallel quartz ribbons ($30-100 \mu\text{m}$) with undulose extinction and polygonal subgrain development. c-d medium to coarse grained ($0.2-0.5 \text{ mm}$ to $0.5-2.5 \text{ mm}$) elongated lenticular porphyroclasts oriented parallel to the main foliation with undulose extinction, folded and tapered twinning, and new subgrain development around the edges.

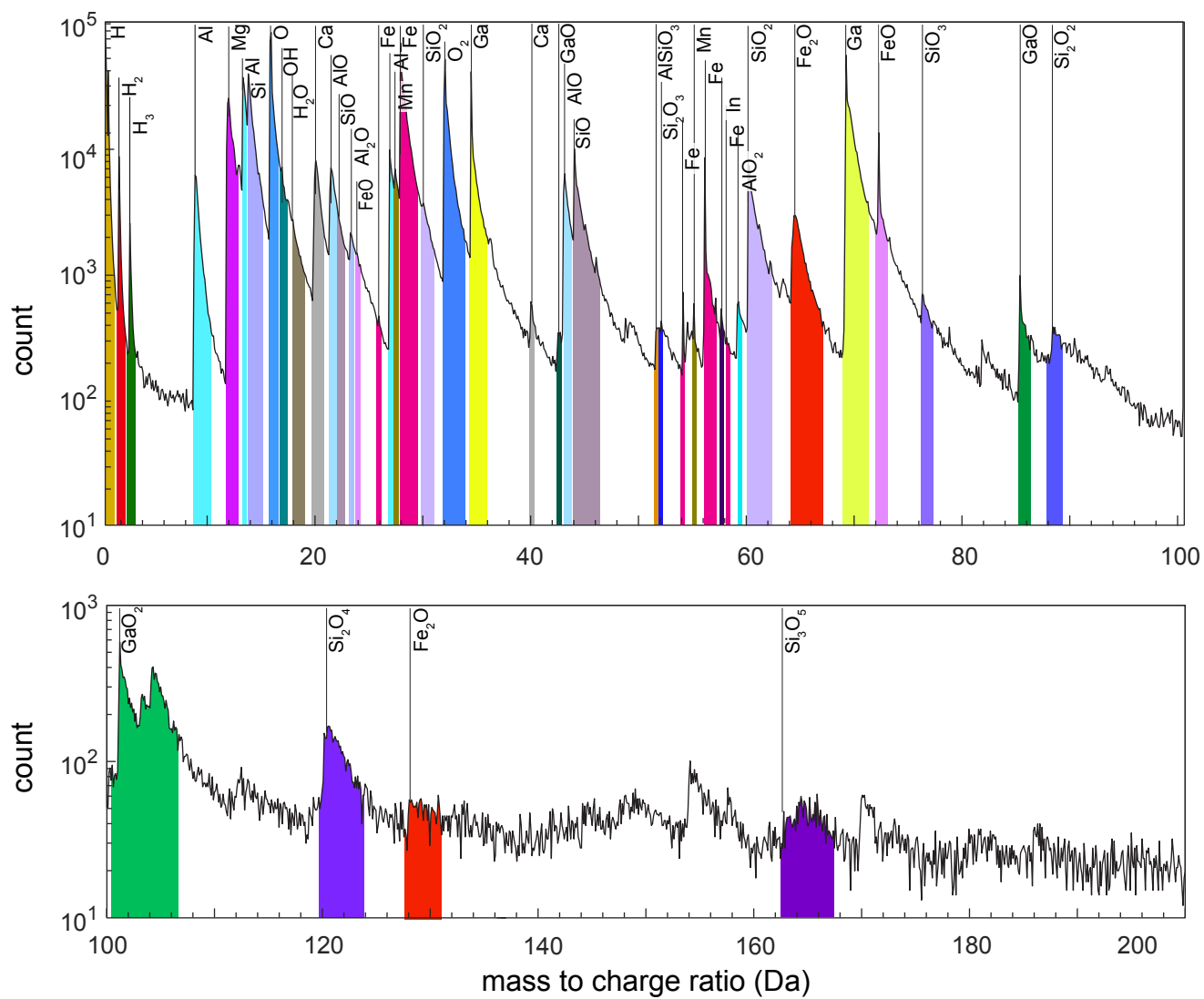


Figure S2. Atom probe mass spectrum of garnet specimen

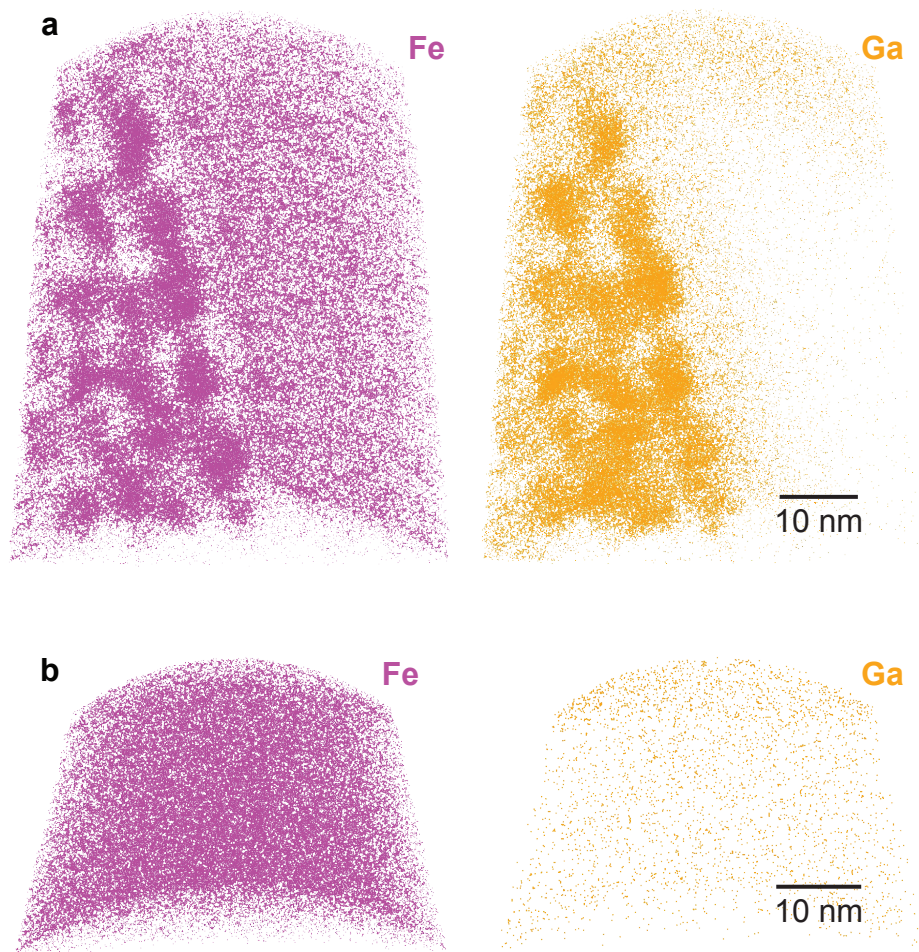


Figure S3. **a** 3D reconstruction of the investigated APT specimen showing the distribution of Fe and Ga. Note: Ga is concentrated both at the tip of the specimen and the within the Fe nanoclusters. **b** 3D reconstruction of another APT specimen from the same lift-out showing the distribution of Fe and Ga. Note: Ga is also enriched at specimen the surface, but no Fe nanoclusters are observed despite the penetration of the ions. The presence of Ga implantation and the absence of clusters in the second specimen (b) suggests the Fe nanoclusters did not result from ion implantation during specimen preparation.

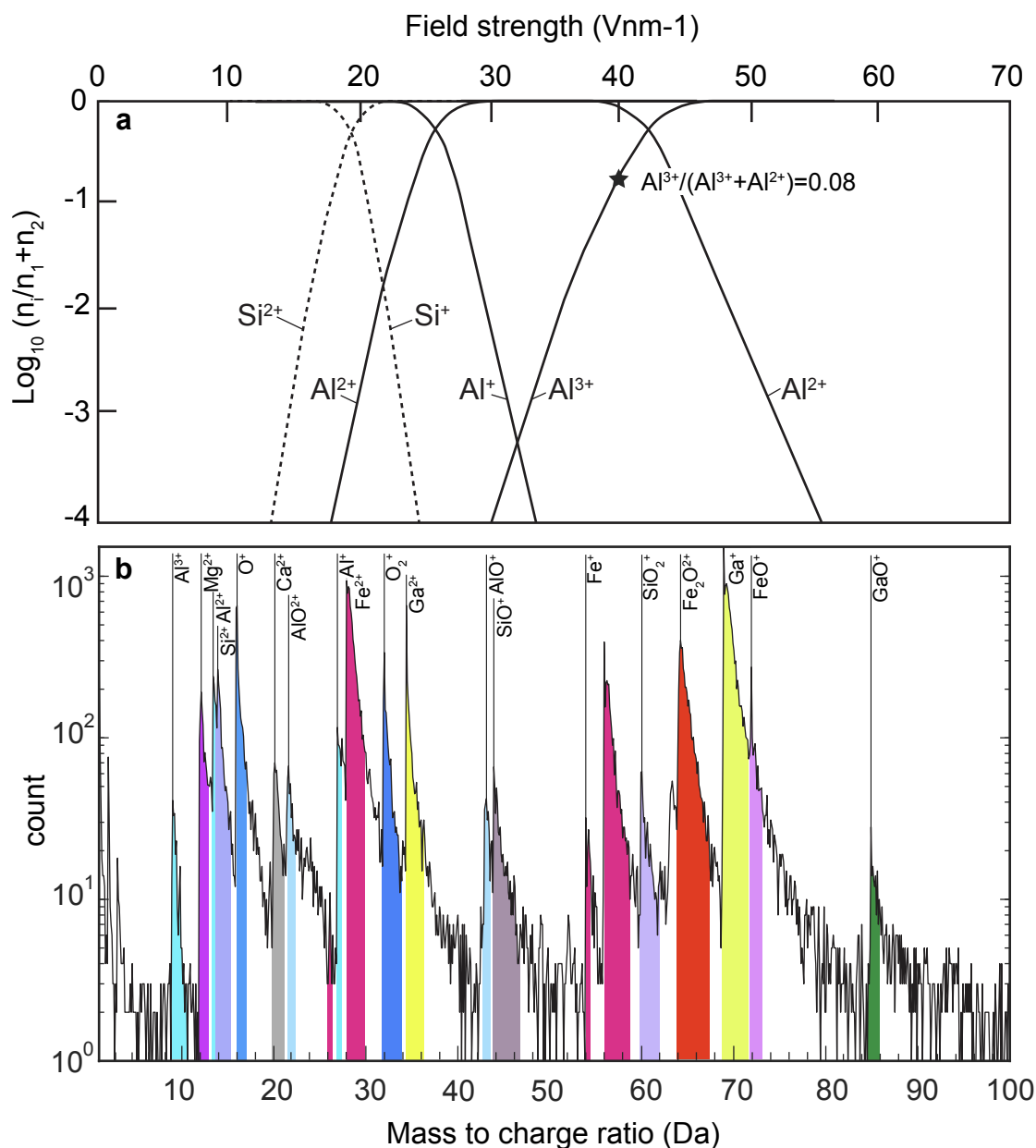


Figure S4. **a** Relative abundance of different charge states for Al and Si during field evaporation plotted logarithmically against field strength (modified after Kingham, 1981). **b** Mass spectrum of Fe nanoclusters extracted from isosurfaces (34 at% isovalue). The ratio of Al³⁺ to Al²⁺ (0.08) in the mass spectrum estimates the field conditions for the clusters at 40 Vnm⁻¹, where Si is expected to field evaporate and post-ionise to be detected only in the double charge state.