

# Soil moisture modulation of midlatitude heat waves

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# Soil moisture modulation of midlatitude heat waves

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## Abstract

Heat waves are broadly expected to increase in severity and frequency under climate change. Case studies highlight a number of physical mechanisms that play a role in present-day heat waves, which typically occur during a coalescence of anomalous atmospheric and land surface conditions. However, a unified model of heat wave physics is lacking, primarily owing to difficulty in disentangling the forcing versus feedback roles of soil moisture and atmospheric variability. Here, we provide observational evidence that soil moisture modulation of heat waves is a generic feature of midlatitude continental climates, and develop a theoretical framework to understand this modulation. Using separation of timescales we derive a diagnostic equation for the nonlinear response of temperature to soil moisture variations, and a dynamical Hasselmann-like model for the soil moisture variations themselves. We find that soil moisture fluctuations control the frequency and intensity of temperature extremes by slowly altering the background state

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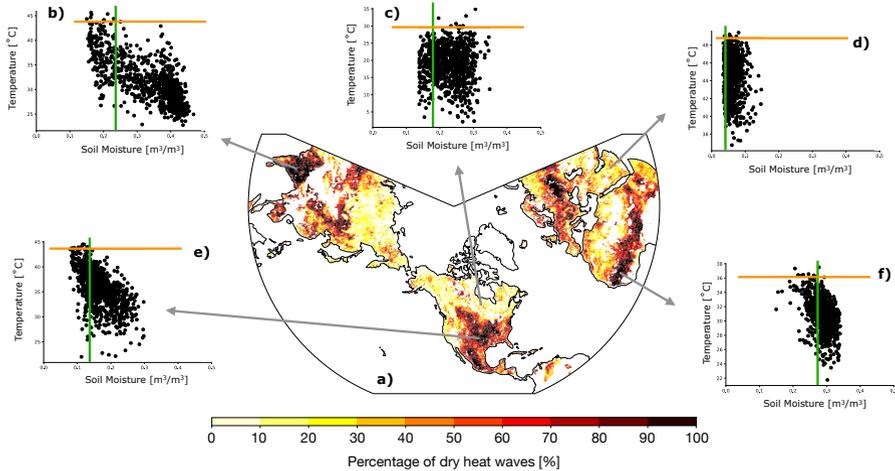
047 on which rapid atmospheric variability acts, rather than by altering  
048 atmospheric variability itself. We also find the slow soil moisture vari-  
049 ations are well-approximated as being primarily driven by stochastic  
050 precipitation variability. Our framework provides a first-principles under-  
051 standing of soil moisture’s role in midlatitude temperature extremes.

052 **Keywords:** heat waves, soil moisture, climate extremes, stochastic models  
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056 Heat waves inflict immense damage to economies [1], ecosystems [2], and  
057 human health [3]; the European heat wave of 2003 alone is estimated to have  
058 killed over seventy thousand people. As average temperatures increase, mean  
059 shifts of daily temperature distributions inexorably lead to increased frequency  
060 and severity of heat waves [4, 5]. The precise impact of a warming climate on  
061 extreme events, however, is highly sensitive to the details of the underlying  
062 distribution [6]. A short and nonstationary observational record makes it dif-  
063 ficult to robustly estimate how often a given historical temperature threshold  
064 should be exceeded. Statistical approaches to attribution depend on fitting the  
065 limited number of observations to a generalized probability distribution with  
066 many free parameters, which makes it difficult to assign a likelihood to even  
067 well-characterized events such as the 2021 Pacific Northwest Heat Wave [7].  
068 At the same time, climate models show no consistent pattern between mean  
069 warming and changes in temperature distributions [8], and are biased in their  
070 representation of temperature variability [9]. These limitations of both the  
071 empirical record and numerical models highlight a growing need for a general  
072 theory of present-day heat waves, especially outside the tropics [10, 11].

073 A unified understanding of heat waves has been hampered by difficulty in  
074 disentangling contributing mechanisms that act across a breadth of timescales:  
075 atmospheric blocking [12], cloud variability [13], and thermal advection near  
076 the surface [14, 15] contribute to fast variations in near-surface temperature  
077 relative to soil moisture dynamics [16–20]. Throughout this paper, we will refer  
078 to the relatively fast processes (variations in blocking, thermal advection, and  
079 cloud radiative effects) as the atmospheric forcing in contrast to variations  
080 driven by (comparatively) slower fluctuations in soil moisture. Capturing the  
081 interaction of these mechanisms is a key component of any physical theory of  
082 heat waves, as it is a general property of stochastic systems that the shape  
083 of distributional tails is controlled by the interaction of mechanisms across  
084 timescales [21].

085 A generic picture of how atmospheric circulation contributes to generat-  
086 ing a midlatitude summer “heat dome” involves the spatial stagnation of a  
087 high pressure system, known as a “blocking” event. Anticyclonic circulation  
088 leads to enhanced surface warming through a combination of subsidence-driven  
089 adiabatic heating [12], advection of warm tropical air [22], and suppression  
090 of convection and clouds [23]. Each of these mechanisms leads to increased  
091 energetic forcing of the surface boundary layer, and thus, of the surface.  
092



**Fig. 1** Dry heat waves. **a)** shows the percentage of days where the maximum temperature exceeds the local 99<sup>th</sup> percentile while soil moisture is less than the local first quantile. **b) – f)** show daily maximum temperature–soil moisture phase space plots with the horizontal orange line indicating the local 99<sup>th</sup> temperature percentile and the vertical green line indicating the local first soil moisture quantile.

While this excess of energy inevitably warms the surface, it does not necessarily lead to the onset of a heat wave. Instead, the extent to which increased surface forcing leads to extreme temperatures is modulated by the thermodynamic state of the surface. The land surface thermodynamic state controls the partitioning of incoming atmospheric forcing into two channels: an ET channel, where incoming energy evaporates moisture, and a heating channel, where incoming energy raises temperatures. The distraction of energy by ET away from heating the surface is commonly referred to as “latent cooling.” The strength of latent cooling is controlled by two factors: (1) the atmospheric vapor pressure deficit (VPD) and (2) the available moisture in the soil for ET [24–26]. The influence of soil moisture on individual heat waves has been shown in a number of case studies [27]. However, a theoretical model of this influence is lacking, as is quantification of its generality without using climate model or reanalysis output.

## Soil moisture modulation of heat wave frequency

Here, we quantitatively describe how each of the contributing factors to temperature extremes – atmospheric forcing and soil moisture variability – influence the intensity and frequency of heat waves. We begin by using observational data to show the generality of soil moisture’s influence on heat waves in midlatitude climates. To our knowledge, this is the first demonstration of soil moisture’s general influence on heat waves using observational data in

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139 the literature. We then use a much longer reanalysis product to demonstrate  
140 the fast-slow influence of atmospheric forcing and soil moisture dynamics on  
141 temperature distributions, as well as statistically quantify the impact of soil  
142 moisture on near-surface temperature. We interpret these data by building  
143 a conceptual model that quantifies the impact of soil moisture on the inten-  
144 sity and frequency of temperature extremes. We do this by first quantifying  
145 how a soil moisture anomaly modifies the probability of a heat wave occur-  
146 ring under favorable atmospheric forcing, and then quantifying the frequency  
147 of soil moisture anomalies.

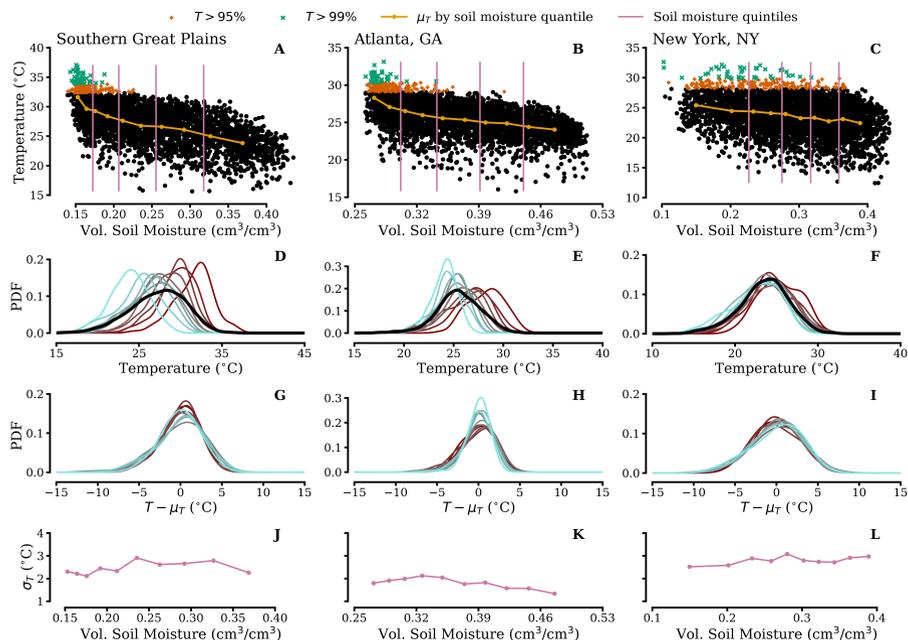
148 Observations in Figure 1a show a clear connection between extremely hot  
149 days and soil moisture deficits in the midlatitudes. In many regions, more  
150 than 90% of days where temperatures exceed the local 99<sup>th</sup> percentile occur  
151 when the soil is dry (see *Methods*). The influence is strongest in continental  
152 midlatitude regions such as the central United States and central Europe, as  
153 well as subtropical regions with strong precipitation seasonality such as South  
154 Asia and the Sahel. The influence is weakest in high latitudes regions such as  
155 Canada and Siberia, and in very arid regions such as the US Southwest and  
156 the Sahara.

157 The difference between areas where soil moisture’s influence is sizable  
158 and areas where its influence is moderate to minimal is best visualized in  
159 temperature-soil moisture phase space (Figs. 1b–f). In regions where extremely  
160 high temperatures are clustered around periods where the soil is dry, a nonlin-  
161 ear relationship emerges where soil moisture depletion leads to a rapid upward  
162 shift in temperature.

163 To understand this relationship, we focus on a small number of continental  
164 locations in the United States that span the potential influence of soil mois-  
165 ture on near-surface temperature (see Table 2) and use the ERA5 reanalysis  
166 product that provides a longer record of surface temperature and soil moisture  
167 (see *Methods* for details) [28]. The reanalysis record shows the same emer-  
168 gent relationship between temperature and soil moisture (Fig. 2A–C) found  
169 in observations, but provides a long enough record to allow us to estimate  
170 conditional means and probability distribution functions of temperature for  
171 individual soil moisture quantiles (Fig. 2A–I). We find that soil moisture has  
172 a potentially large impact on mean temperature across soil moisture quantiles  
173 (up to  $\sim 5 - 10$  °C), but has a minimal impact on temperature variabil-  
174 ity within each quantile regardless of its impact on mean temperature (see  
175 Fig. 2G–L).

176 Based on our analysis in Figures 1 and 2, we highlight two conclusions  
177 that inform and motivate our theoretical model. The first is that soil moisture  
178 influences heat waves primarily by modulating the background state on top  
179 of which atmospheric variability acts, rather than by changing atmospheric  
180 variability itself. However, the static scatter-plot analysis carried out in Figs. 1  
181 and 2 cannot answer the question of whether these soil moisture variations are  
182 themselves a response to – and thus a feedback on – synoptic forcing. We will  
183 explore this ambiguity below.

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**Fig. 2 Decoupling soil moisture and atmospheric variability.** A–C shows scatter plots of the temperature-soil moisture phase space for three locations, with extreme temperature days colored orange ( $T > 95^{\text{th}}$  percentile) and green ( $T > 99^{\text{th}}$ ), as well as the mean temperature averaged within each soil moisture quantile (gold curve), see the legend above A–C. D–F show the estimated probability distribution function (PDF) of near surface temperature for a given quantile of soil moisture using ERA5 reanalysis output. The cyan (maroon, resp.) PDFs are for the high (low, resp.) quantiles of soil moisture. The mean PDF is in black. G–I are as D–F, after subtracting the mean. J–L show the standard deviation of the temperature PDF in each soil moisture quantile.

The second implication is that the degree to which soil moisture impacts heat wave frequency and intensity is related to the strength of the nonlinear correspondence between soil moisture and temperature. Indeed, soil moisture has a sizable impact in the SGP which has a strongly nonlinear relationship between soil moisture and temperature, whereas soil moisture has a minimal impact in New York which has a weak nonlinearity strength (see the gold lines in Fig. 2A–C). Atlanta serves as a “middle of the road” example between the SGP and New York.

## Results

We now turn to building a conceptual *soil moisture—atmosphere coupled model* (SMACM) that will provide insight into both the nonlinear coupling of soil moisture and temperature, as well as the mechanisms that drive soil moisture variations. SMACM is formulated as a one-dimensional box model of surface temperature and soil moisture, where the land surface is forced by the atmosphere. A full derivation of model equations and our model evaluation is

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231 **Table 1 SMACM parameters.** List of SMACM parameter names, their symbols,  
 232 average order of magnitude (OM), and units.

Parameter	Symbol	OM	Units
Effective heat capacity of surface layer	$C$	10	$\text{J K}^{-1} \text{m}^{-2}$
Two meter air temperature	$T$	$10^2$	K
Time	$t$	1	s
Radiative forcing	$\mathcal{F}$	$10^2$	$\text{W m}^{-2}$
Dry feedback strength	$\alpha$	10	$\text{W m}^{-2} \text{K}^{-1}$
Latent heat of vaporization	$\lambda$	$10^6$	$\text{J kg H}_2\text{O}^{-1}$
Surface evaporative conductance	$\nu$	$10^{-2}$	$\text{kg H}_2\text{O m}^{-2} \text{s}^{-1}$
Clausius-Clapeyron derivative at mean $T_d$	$\gamma$	$10^{-4}$	$\text{K}^{-1}$
Soil moisture fraction	$m$	$10^{-1}$	–
Constant deep soil moisture fraction	$m_0$	$10^{-1}$	–
Two meter dew point temperature	$T_d$	$10^2$	K
Soil holding capacity	$\mu$	10	$\text{kg H}_2\text{O m}^{-2}$
Precipitation	$\mathcal{P}$	$10^{-1}$	$\text{kg H}_2\text{O m}^{-2} \text{s}^{-1}$

246 provided in the *Methods*; see Table 1 for parameter descriptions, units, and  
 247 average order of magnitudes used in this study.

248 Assuming the timescale of thermal adjustment is much smaller than that  
 249 of moisture adjustment (we prove this explicitly in the *Methods*, but it is a  
 250 common assumption in various studies, e.g., [18–20]) and using the govern-  
 251 ing equations of SMACM (Eqns. (12) and (13) in the *Methods*), a diagnostic  
 252 relationship between temperature,  $T$ , and soil moisture,  $m$ , can be derived as  
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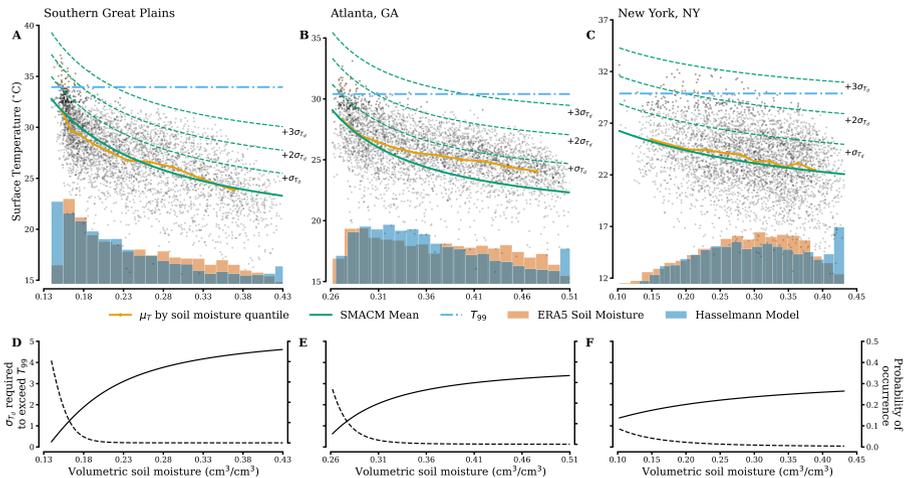
$$254 \quad T(m) = T_0 \left( \frac{1}{1 + \eta m} \right) + T_d, \quad (1)$$

255 where we have defined  $T_0 := \mathcal{F}/(\alpha + \lambda\nu\gamma m_0)$  and  $\eta := \lambda\gamma\nu/(\alpha + \lambda\gamma\nu m_0)$ .  $T_0$   
 256 and  $\eta$  have salient physical interpretations.  $T_0$  is the maximum soil moisture  
 257 induced temperature departure from the dew point temperature (occurring  
 258 when the surface is completely dry).  $\eta$  measures the coupling strength between  
 259 soil moisture and temperature, and is determined by the relative strength of the  
 260 latent heat feedback associated with a saturated *surface* layer to the strength  
 261 of the dry feedback and the recalcitrant latent heat feedback associated with  
 262 deep soil moisture.  
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266 This procedure separates the two aforementioned modes of temperature  
 267 variability; the dew point temperature is a purely additive component to the  
 268 soil moisture dependent piece. In the language of dynamical systems, (1)  
 269 is referred to as a *nullcline*, defined as a curve in phase space where the dynamics  
 270 of one or more variables is neglected [29, 30], and we will henceforth refer to (1)  
 271 as the temperature nullcline, or just “the nullcline.”

## 272 Heat wave frequency

273 The nullcline can be used to understand the role of soil moisture in heat  
 274 wave frequency. In Figure 3A–C, we show the daily temperature-soil moisture  
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**Fig. 3 Quantifying the impact of soil moisture on heat wave frequency and intensity.** **A–C** shows the temperature-soil moisture phase space for the SGP, Atlanta and New York, respectively, with the mean temperature by soil moisture quantile (gold line) and the temperature nullcline (Eqn. (1), green solid line). The dashed lines are the nullcline after anomalies of varying size in two meter dew point temperature have been applied. We also show the soil moisture histograms for ERA5 (salmon) and our Hasselmann model (3) (light blue). In **D–F**, the solid lines show the required anomaly size in the two meter dew point temperature for two meter temperature given by (1) to exceed the 99<sup>th</sup> percentile in each location as a function of soil moisture. The dashed lines represent the probability of such an anomaly to occur, assuming a normally distributed dew point temperature.

phase space of the SGP, Atlanta, and New York with the temperature nullcline overlaid in green, assuming mean values for the net shortwave radiation and dew point temperature. We apply dew point temperature anomalies to the nullcline in each location to show the impact of atmospheric variability on the orientation of the nullcline in phase space. (An analogous figure for radiative forcing anomalies is provided in the *Supplementary Information*.)

When the soil is dry, less extreme dew point temperature anomalies are needed to exceed the local 99<sup>th</sup> temperature percentile than when the soil is wet. We calculate the anomaly size needed to exceed the local 99<sup>th</sup> percentile, as well as the probability of such an event occurring, as a function of soil moisture in panel 3D–F. For the SGP, the anomaly size required grows nonlinearly with soil moisture, whereas in New York the anomaly size is relatively flat and linear, showing that the higher the degree of nonlinearity between temperature and soil moisture, the more dry soils load the dice in favor of heat waves.

This can be explained by the coupling parameter  $\eta$  determining the curvature of the nullcline, i.e.,

$$\Delta T := \frac{T(m=1) - T(m=0)}{T_0} = \frac{\eta}{\eta + 1}. \quad (2)$$

323 If  $\eta$  is small (weak coupling regime), then  $\Delta T \sim 0$ , and soil moisture’s influence  
 324 on temperature disappears by virtue of the nullcline “flattening out.” If  $\eta$   
 325 is large (strong coupling regime), then  $\Delta T \sim 1$ , resulting in the maximum  
 326 difference between temperatures when the soil is dry or wet.

327 This showcases how the nonlinear coupling between soil moisture and tem-  
 328 perature influences heat wave frequency: in locations with high  $\eta$  parameters,  
 329 relatively common atmospheric events can cause temperature extremes when  
 330 the soil is dry, implying that the frequency at which the soil dries out is a  
 331 leading order indicator of heat waves in these locations. In locations with low  
 332  $\eta$  parameters, soil moisture matters little for heat wave frequency, if at all.

### 334 Heat wave intensity

335 Soil moisture’s impact on heat wave intensity can be determined by assum-  
 336 ing the soil is dry (if the soil is wet, heat waves are either just as or far  
 337 less likely than when the soil is dry, depending on the coupling strength, see  
 338 Figure 3D–F). In this regime, the departure of temperature away from the  
 339 dew point temperature is determined by  $T_0$ , or “dry-out” temperature. The  
 340 dry-out temperature depends on root-level soil moisture content, highlighting  
 341 the role of deep soil moisture content on surface temperatures. Atmospheric  
 342 dynamics impact the nullcline by altering the magnitude of the dry-out tem-  
 343 perature. (Recall  $T_0$  linearly depends on  $\mathcal{F}$ .) Thus, when surface forcing is high,  
 344 the impact of the land thermodynamics on near surface temperature is exac-  
 345 erbated, particularly when the soil is dry. Our framework therefore suggests  
 346 an important feedback mechanism between the atmospheric forcing and local  
 347 thermodynamics, where strong atmospheric forcing makes the land response  
 348 more skewed towards extremely high temperatures. We find across a variety  
 349 of nullcline configurations,  $T_0$  captures at least 52% of the temperature depar-  
 350 ture from the mean dew point temperature on the hottest day in the reanalysis  
 351 product (see Table 2).

### 354 Soil moisture variability

355 So far, we have consistently invoked dry soil conditions to demonstrate the  
 356 role of soil moisture in temperature extremes. But how often does the soil  
 357 dry out? Indeed, the nonlinearity between temperature and soil moisture  
 358 described above is a necessary, but not sufficient condition for soil moisture  
 359 dynamics being a leading order factor in heat wave frequency. Soil moisture is  
 360 impacted by both precipitation frequency and intensity that ultimately con-  
 361 trol its influence on temperature. For example, if the temperature nullcline is  
 362 highly nonlinear, but the soil is perpetually dry or wet, then the nonlinearity  
 363 matters little because soil moisture is too restricted in phase space to cause  
 364 significant temperature variance (panel 1D is an example of this scenario).  
 365 Hence, the soil moisture distribution also plays a crucial role in assessing the  
 366 amount of temperature variance owing to soil moisture.

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**Table 2 SMACM parameter value impacts on heat wave intensity and frequency.** We calculate the ratio of temperature variance owing to soil moisture to total temperature variance for five locations in the continental United States. We also show the ratio of the dry-out temperature,  $T_0$ , to the maximum departure from the mean dew point temperature in the reanalysis product,  $T_X := \max(T) - \bar{T}_d$ , for five locations in the continental United States. Provided as well are the nonlinear strength parameter,  $\eta$ , the dry-out temperature,  $T_0$ , and the mean and variance of the soil moisture distribution in each location.

Location	$\eta$ [-]	$T_0$ [°C]	$\bar{m}$ [-]	$\sigma_m^2$ [-]	$\sigma_{T_m}^2/\sigma_T^2$ [-]	$T_0/T_X$ [-]
Southern Great Plains	3.43	12.4	0.36	0.06	0.49	0.66
Atlanta, GA	2.49	7.61	0.42	0.07	0.39	0.56
Wichita, KS	2.14	11.8	0.41	0.07	0.39	0.70
Dallas, TX	1.70	9.24	0.31	0.07	0.39	0.56
New York, NY	0.94	7.80	0.58	0.05	0.06	0.52

We use SMACM to derive a nonlinear Hasselmann-like model [31] for the VPD-driven soil moisture response to stochastic precipitation forcing. Again invoking slow timescales, we write

$$\mu \frac{dm}{dt} = \mathcal{P} - \nu\gamma T_0 \left( \frac{m}{1 + \eta m} \right). \quad (3)$$

(Derivation in *Methods*.) We see that  $\eta$  again determines the degree of non-linearity. We show the distribution of soil moisture when (3) is forced with ERA5 precipitation time series in Figure 3A–C and see that our model largely reproduces the distribution of soil moisture found in ERA5, despite neglecting atmospheric forcing other than precipitation.

We use (3) to compute the ratio of temperature variance owing to soil moisture to total temperature variance in five locations in Table 2 (see *Methods* for calculation details). We find that in locations where  $\eta$  and  $T_0$  are large, and soil moisture is generally dry while still maintaining non-negligible variance, the ratio of soil moisture induced temperature variance to total temperature variance is significant ( $\gtrsim 40\%$ ). Meanwhile, New York has a low  $\eta$  and a relatively wet soil, leading to a small fraction of temperature variance owing to soil moisture. These findings corroborate the trends seen in Figure 1, where higher fractions of dry heat waves are found where the ratio of temperature variances seen in Table 2 are highest, and the lower fractions are found where variances are lowest.

## Synthesis

We have explained soil moisture modulation of heat waves found in observations (see Fig. 1) by partitioning near-surface temperature fluctuations into a component driven by rapid atmospheric variability and another driven by slow soil moisture dynamics. We used a theoretical framework to encode these components into the fast and slow modes of a dynamical system. We find slow

415 soil moisture variability has a nonlinear impact on the background tempera-  
416 ture state, while rapid atmospheric variability is well-approximated as linearly  
417 additive on top of background state variations (Fig. 2G–L). The slow soil mois-  
418 ture variability itself is well approximated as primarily driven by stochastic  
419 precipitation, with other atmospheric forcing factors likely playing a secondary  
420 role (Fig. 3A–C).

421 Both a strong nonlinear coupling and significant soil moisture variability  
422 are required for soil moisture to be a strong control on heat waves. The degree  
423 of nonlinearity is encoded in the temperature nullcline (Eqn. (1)) and varies  
424 across space depending on land surface properties (such as the surface con-  
425 ductance and root-level soil moisture content). Differences in the strength of  
426 the nullcline partially explain the heterogeneity of soil moisture’s influence on  
427 temperature extremes found in observations and in reanalysis. We quantified  
428 these differences in nullcline slope by introducing the coupling parameter  $\eta$ .  
429 In strongly coupled locales (where  $\eta$  is large), moisture depletion increases the  
430 potential for a heat wave, and atmospheric anomalies turn this potential into  
431 actuality. In weakly coupled locales (where  $\eta$  is small), soil moisture matters  
432 little for temperature extremes, as the nullcline “flattens out” and temperature  
433 extremes are determined primarily by atmospheric forcing.

434 By combining the nullcline (1) with the Hasselmann-like model for soil  
435 moisture (3), we can calculate the fraction of temperature variance attributable  
436 to soil moisture’s influence on temperature. We find that the fraction of tem-  
437 perature variance attributable to soil moisture fluctuations is highest in the  
438 same locations where heat waves disproportionately favor dry soils (see Fig. 1).  
439 Over a majority of the continental midlatitude regions heat waves are condi-  
440 tioned on precipitation deficits depleting soil moisture and priming the surface  
441 for temperature extremes when a favorable atmospheric anomaly occurs.

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## 443 **Methods**

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### 445 **Observational data**

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447 Daily temperature data are from the [CPC Global Unified Temperature](#)  
448 [dataset](#), while daily soil moisture data are from the [ESA-CCI v06.1](#) [32–34].  
449 June, July, and August values of both quantities were used from 2010–2019. A  
450 value of 100 in Fig. 1a indicates that all days above the local 99<sup>th</sup> percentile  
451 occurred when the soil was drier than the local first quantile (i.e. all points  
452 above the orange line in Figs. 1b)–e) are also left of the green line).

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### 454 **Reanalysis data**

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456 ERA5 daily mean temperature and soil moisture data [28] for the southern  
457 great plains (we use the Department of Energy Atmospheric Radiation Mea-  
458 surement site for the precise coordinates for this analysis), Atlanta and New  
459 York are used to create Figure 2. In Figure 4, we use hourly precipitation and

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solar radiative forcing data from the SGP to force our model equations, as well as the climatological two meter dew point temperature. 461  
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## Derivation of SMACM model equations 463 464

Consider a land surface with a dynamic temperature,  $T$ , and soil moisture content,  $m$ .  $m$  is a value between zero and unity that denotes the fractional saturation of the land surface. The land surface is coupled to a non-dynamic atmosphere with dew point temperature  $T_d$  and reference specific humidity  $q_r$ . 465  
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In the energy sector, the land is forced by radiation from the Sun, denoted as  $\mathcal{F}$ . The surface cools by releasing heat through latent, sensible, and longwave channels. Thus, conservation of energy implies 472  
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$$C \frac{dT}{dt} = \mathcal{F} - F_{LW} - G - \mathcal{H} - \lambda \mathcal{E}, \quad (4) \quad 474$$

where  $F_{LW}$  is the net longwave heat flux,  $G$  is the ground heat flux,  $\mathcal{H}$  is the sensible heat flux,  $C$  is the effective surface heat capacity,  $\lambda$  is the latent heat of vaporization of water, and  $\mathcal{E}$  is evapotranspiration. We parameterize the sum of the longwave, ground, and sensible heat fluxes as linear departures from the atmospheric dew point, such that 475  
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$$G + \mathcal{H} + F_{LW} = \alpha(T - T_d). \quad (5) \quad 482$$

We parameterize evapotranspiration as, 483  
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$$\mathcal{E} = \nu m (q_s(T) - q_r), \quad (6) \quad 485$$

where  $q_s(T)$  is the saturation humidity given by the Clausius-Clapeyron relationship and  $\nu$  is the surface evaporative conductance. Linearizing (6) around the dew point temperature yields 486  
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$$\mathcal{E} = \nu \gamma m (T - T_d), \quad (7) \quad 492$$

where we have used that  $q_s(T = T_d) = q_r$  and  $\gamma$  is the temperature derivative of the Clausius-Clapeyron relationship evaluated at the dew point temperature. 493  
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Using (5) and (7) in (4) results in 498

$$C \frac{dT}{dt} = \mathcal{F} - (\alpha + \lambda \gamma \nu m)(T - T_d). \quad (8) \quad 499$$

So far, this model is identical to the one presented in [19], but one final alteration is required to (8). Figure 4 shows that when the surface soil is depleted, the latent heat flux does not approach zero. This owes to the fact that deep, root-level soil moisture is present, and contributes to the latent heat flux. We 500  
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507 therefore add an additional soil moisture fraction to (8) to capture this effect,  
 508 denoted as  $m_0$ , resulting in the final temperature evolution equation,

$$510 \quad C \frac{dT}{dt} = \mathcal{F} - (\alpha + \lambda\gamma\nu(m + m_0))(T - T_d), \quad (9)$$

511  
 512 as desired. Note that deep soil is assumed to act on longer timescales than are  
 513 relevant for our analysis, and therefore is treated as constant.

514 In the moisture sector, we consider the conservation of water mass flux to  
 515 write

$$516 \quad \mu \frac{dm}{dt} = \mathcal{P} - \mathcal{E} - \mathcal{R} - \mathcal{I}, \quad (10)$$

517 where  $\mu$  is the holding capacity of the soil,  $\mathcal{P}$  is stochastic precipitation forcing,  
 518  $\mathcal{R}$  is runoff and  $\mathcal{I}$  is infiltration. In this model, runoff and infiltration act only  
 519 to keep soil moisture within restricted bounds set by the field capacity,  $\Theta$ ;  
 520 therefore,  $\mathcal{R} = \mathcal{I} = 0$ . At all other times, we have

$$521 \quad \mu \frac{dm}{dt} = \mathcal{P} - \nu\gamma m(T - T_d), \quad (11)$$

522 Taken together, the governing equations of SMACM are

$$527 \quad C \frac{dT}{dt} = \mathcal{F} - (\alpha + \lambda\gamma\nu(m + m_0))(T - T_d), \quad (12)$$

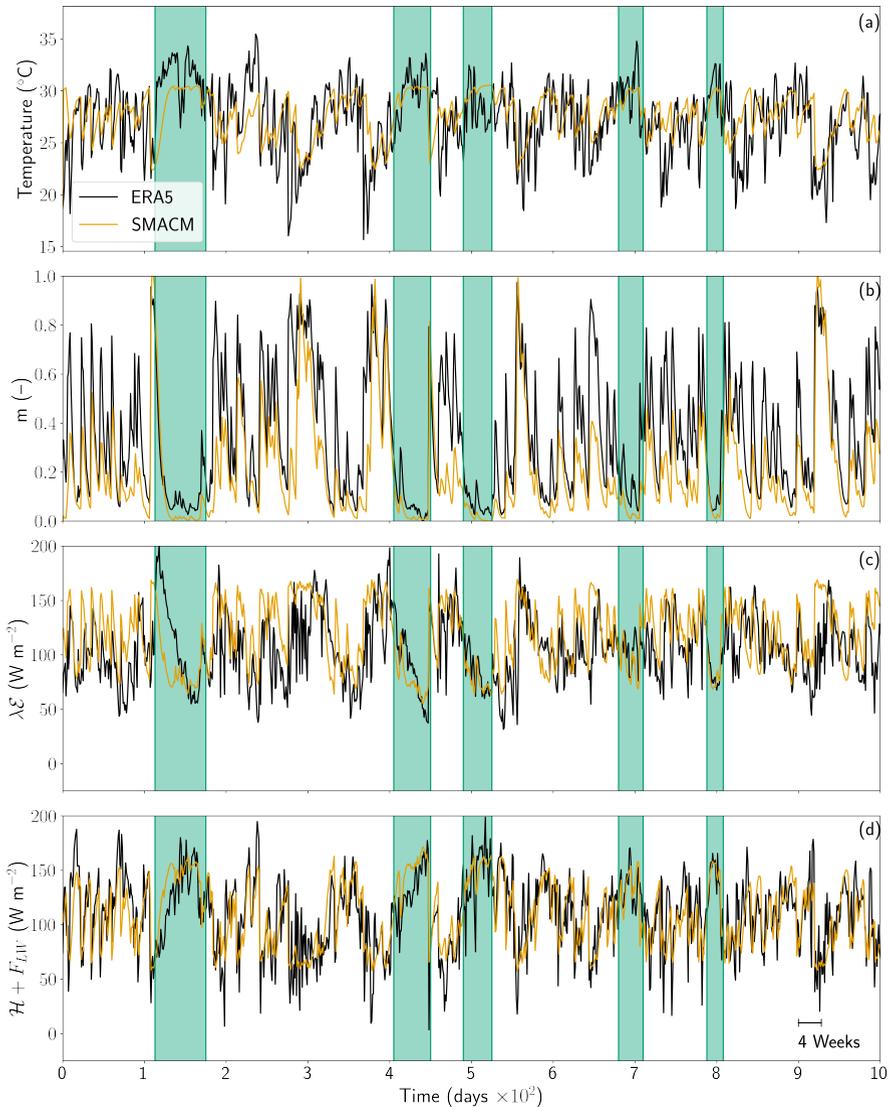
$$529 \quad \mu \frac{dm}{dt} = \mathcal{P} - \nu\gamma m(T - T_d). \quad (13)$$

530  
 531 To derive the nulleline equation (1), we simply assume that the temperature  
 532 derivative in (12) can be set to zero, owing to the assumption that the timescale  
 533 of thermal adjustment is much smaller than that of soil moisture adjustment  
 534 (as is evidenced by Fig. 2 as well as other studies, e.g., [18–20]). Setting (12)  
 535 to zero and solving for  $T$  in terms of soil moisture results in (1) after some  
 536 rearrangement of terms.

537 To derive the Hasselmann-like model for soil moisture, we use (1) for  
 538 temperature in (13). Simplifying results in (3).

## 541 Model evaluation

542 We use reanalysis output from ERA5 1979–2021 over the SGP [28] to eval-  
 543 uate SMACM performance. Using hourly output for net shortwave radiation  
 544 and precipitation as model forcings, taking  $T_d$  to be the average climatological  
 545 two meter dew point temperature, and tuning the remaining free parameters,  
 546 we arrive at the time series shown in Figure 4A–D. Of particular interest are  
 547 periods where soil moisture is low (green shading in Figure 4), where a reduc-  
 548 tion in the latent heat flux (panel 4C) is met with an increase in the sensible  
 549 and longwave heat fluxes (panel 4D), as well as an increase in near surface air  
 550 temperature (panel 4A). Some high-frequency variability in SMACM is miss-  
 551 ing because we have not included fast fluctuations in atmospheric circulation;



**Fig. 4 Model evaluation.** A comparison of SMACM to ERA5 output at the SGP. **A–D** show the near surface air temperature, soil moisture fraction, latent heat flux, and dry heat flux, respectively.

the only rapid fluctuations considered are fluctuations in shortwave radiation (the  $\mathcal{F}$  term). Despite the lack of rapid atmospheric forcing, SMACM adequately reproduces soil moisture variations, as well as the partition of sensible and latent heat. SMACM is also able to reproduce slow variations in temperature, particularly elevated mean temperatures during times of dry soil moisture conditions.

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599 **Proof of timescale separation in SMACM**

600 Throughout this paper, we have argued that all temperature variability can  
 601 be decomposed into a “slow” mode driven by soil moisture fluctuations and a  
 602 “fast” mode driven by fluctuations in shortwave radiation and thermal advec-  
 603 tion. Here we prove that, within SMACM, the timescale of thermal adjustment  
 604 is fast relative to moisture adjustment, thus justifying our separation of  
 605 timescales approach.

606 Assume in all that follows  $\mathcal{P} = 0$ , as we are analyzing the *response*  
 607 of SMACM’s model equations to precipitation forcing. Thus, the relevant  
 608 equations are

$$610 \quad \frac{dT}{dt} = \frac{1}{C} (\mathcal{F} - (\alpha + \lambda\nu\gamma(m + m_0))(T - T_d)) =: f_T(m, T), \quad (14)$$

$$611 \quad \frac{dm}{dt} = -\frac{\nu\gamma m(T - T_d)}{\mu} =: f_m(m, T), \quad (15)$$

612 where we have defined the right-hand-sides of each equation for simplicity in  
 613 what follows.

614 **Step 1: Equilibrium analysis**

615 The equilibria of (14)–(15) correspond to the set of solutions of the algebraic  
 616 system  $f_T(m, T) = f_m(m, T) = 0$ .  $f_m(m, T) = 0$  implies there are two equi-  
 617 librium, namely  $m = 0$  and  $T = T_d$ . However,  $T = T_d$  does not satisfy the  
 618 condition  $f_T(m, T = T_d) = 0$  for any  $m$ , as  $\mathcal{F} \neq 0$  by assumption and  $C$  is  
 619 finite. Therefore, the sole equilibrium is given by  $(m^*, T^*) = (0, T^*)$ , where  
 620  $T^* := \mathcal{F}/(\alpha + \lambda\nu\gamma m_0) + T_d$ . This makes intuitive sense: the only equilib-  
 621 rium is a totally dry soil column, and the temperature is prescribed solely by  
 622 atmospheric conditions and thermodynamic attributes of the surface.

623 **Step 2: Determine stability of equilibrium**

624 The stability of the equilibrium above is determined by the Jacobian of (14)–  
 625 (15) [29]. Therefore, we evaluate

$$626 \quad \mathcal{J}(m, T) = \begin{pmatrix} \partial_T f_T(m, T) & \partial_m f_T(m, T) \\ \partial_T f_m(m, T) & \partial_m f_m(m, T) \end{pmatrix}, \quad (16)$$

$$627 \quad = \begin{pmatrix} -(\alpha + \lambda\nu\gamma(m + m_0))/C & -\lambda\nu\gamma(T - T_d)/C \\ -\nu\gamma m/\mu & -\nu\gamma(T - T_d)/\mu \end{pmatrix}, \quad (17)$$

628 where  $\partial_x := \partial/\partial x$ . Evaluating (17) at  $(m^*, T^*)$  results in

$$629 \quad \mathcal{J}(m^*, T^*) = \begin{pmatrix} -(\alpha + \lambda\nu\gamma m_0)/C & -\lambda\nu\gamma\mathcal{F}/(C(\alpha + \lambda\nu\gamma m_0)) \\ 0 & -\nu\gamma\mathcal{F}/(\mu(\alpha + \lambda\nu\gamma m_0)) \end{pmatrix}. \quad (18)$$

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Equation (18) is an upper triangular matrix, and therefore the eigenvalues lie on the diagonals, and the corresponding eigenvectors are the columns [35]. Therefore, the eigenvalues are,

$$E_1 := -\frac{\alpha + \lambda\nu\gamma m_0}{C}, \quad E_2 := -\frac{\nu\gamma\mathcal{F}}{\mu(\alpha + \lambda\nu\gamma m_0)}. \quad (19)$$

As both eigenvalues are negative, the equilibrium  $(m^*, T^*)$  is a stable node [29].

### Step 3: Solve for eigenvectors

As mentioned above, the eigenvectors are the columns of (18). Therefore, we have

$$\vec{\xi}_1 := \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad \vec{\xi}_2 := \begin{pmatrix} \frac{\mathcal{F}\lambda\nu\gamma\mu}{C\mathcal{F}\nu\gamma - (\alpha + \lambda\nu\gamma m_0)^2\mu} \\ 1 \end{pmatrix}, \quad (20)$$

where we've simplified each eigenvector to illustrate the point that  $\vec{\xi}_1$  points solely along the  $T$ -axis, and therefore corresponds to thermal adjustments only. It follows that the timescale associated along this direction of decay corresponds to the timescale of thermal adjustment, whereas the timescale of decay along  $\vec{\xi}_2$  corresponds to moisture adjustment (and the resulting temperature adjustment as the soil dries). This proves that temperature has two modes that evolve on different timescales, as represented in (1).

### Step 4: Compare timescales of decay

The timescales of decay along  $\vec{\xi}_1$  and  $\vec{\xi}_2$  are found by taking the absolute value and inverse of their corresponding eigenvalues [29], such that

$$\tau_1 = \frac{C}{\alpha + \lambda\nu\gamma m_0}, \quad \tau_2 = \frac{\mu(\alpha + \lambda\nu\gamma m_0)}{\nu\gamma\mathcal{F}}. \quad (21)$$

Using the average order of magnitude of each parameter from Table 1, we find that,

$$\frac{\tau_2}{\tau_1} \sim 10^6. \quad (22)$$

Thus, we have shown that the timescale of purely thermal adjustment,  $\tau_1$ , is fast relative to the timescale of moisture adjustment and its influence on temperature,  $\tau_2$ , as desired.

## Calculation the temperature variance owing to soil moisture

To calculate the ratios of temperature variance owing to soil moisture to dew point temperature variance, we utilize the ERA5 soil moisture time series at each location and use it to force (1). The variance is then taken, and compared with the dew point temperature variance from ERA5 in each location. The

691 parameter values at each location (i.e.,  $\eta$  and  $T_0$ ) are calculated by fitting the  
 692 model to ERA5 time series, as was done in Figure 4. Soil moisture mean and  
 693 variance is calculated from ERA5.

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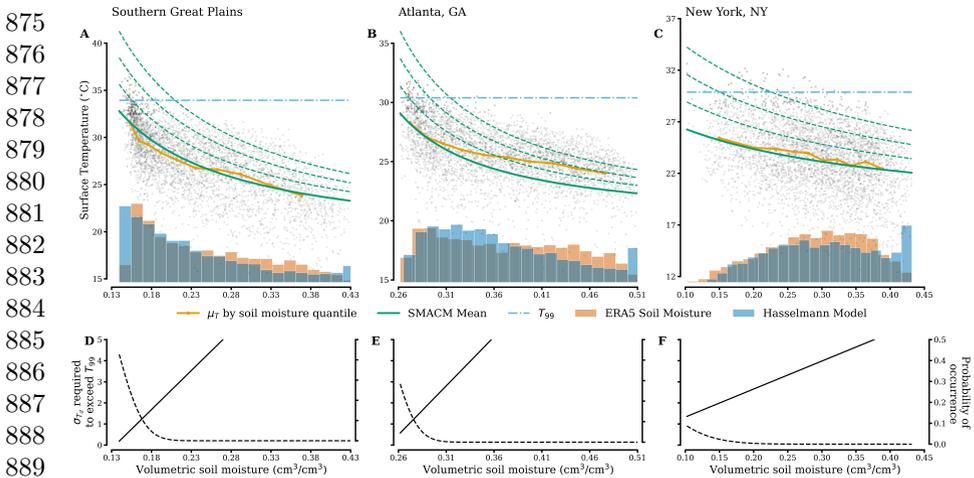
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890 **Fig. 5** As Figure 3, but with anomalies being applied in the radiative forcing,  $\mathcal{F}$ , rather  
891 than dew point temperature.

892  
893 **Supplementary information.**

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905 **Author contributions.** AMB carried out research tasks, did analytic cal-  
906 culations, and prepared figures. LRVZ created Figure 1. CP contributed to  
907 Figure 2. LRVZ and CP provided advising and designed the project. AMB,  
908 LRVZ, and CP wrote the manuscript.

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