

The Fate of Carbon during Earth's Core–Mantle Differentiation

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Abstract

Carbon is an essential element for the existence and evolution of life on Earth, constitutes up to 50% of dry biomass, and is likely a requirement for all life in the universe. Its high abundance in Earth’s crust and mantle (the Bulk Silicate Earth, BSE) is surprising because carbon is strongly siderophile (metal-loving) and should have segregated almost completely into Earth’s core during accretion^{1–3}. Estimates of the concentration of carbon in the mantle lie mostly in the range of 80–120 ppm^{4,5}, which is much higher than expected based on simple models of core–mantle differentiation^{2,6,7}. Here we show through experiments at 49–71 GPa and 3600–4000 K that carbon is significantly less siderophile at such conditions than at the low pressures (≤ 3 GPa) of previous studies^{1–3,8}. We derive a new parameterization of the pressure–temperature dependence of the metal–silicate partitioning of carbon and apply this in a state-of-the-art model of planet formation and differentiation^{9,10} that is based on astrophysical N-body accretion simulations. Results show that BSE carbon concentrations increase strongly starting at a very early stage of Earth’s accretion and, depending on the concentration of carbon in accreting bodies, can easily reach or exceed estimated BSE values. In contrast, simple models of “continuous core formation”^{11–13} require all BSE carbon to be accreted after core formation ended, but this is not consistent with astrophysical models of accretion.

The extents to which the volatile contents of terrestrial planets were determined by the processes of accretion and core formation, impact-driven volatile loss, and addition of a late veneer subsequent to core formation are highly controversial^{9,14–17}. The Earth is considered to have formed from bodies of chondritic composition, although it is unlikely that any particular chondrite group constituted the sole building blocks of the planet¹⁸. Different classes of chondrites have varying carbon concentrations, but importantly, such concentrations are all

high, at the weight percent level¹⁹. This is far in excess of the concentration of carbon present in the Bulk Silicate Earth (BSE) today, which is currently estimated to be 80–120 ppm^{4,5} or possibly as high as 765 ± 300 ppm²⁰. Additionally, the metal–silicate partition coefficient of carbon at low pressures and temperatures (≤ 13 GPa and ≤ 2000 K) is very high ($D > 1000$; see D defined in Eq. 2 and ref.^{1–3,21}). Consequently, almost all carbon should have partitioned into the core during its formation, leaving the mantle almost completely depleted in this element. Several explanations have been proposed for the BSE C concentration, including crystallization of C-rich phases (e.g. diamond) from a C-saturated magma ocean, imperfect metal–silicate equilibration during core formation, accretion of a “late veneer” after the end of core formation, and carbon outgassing from the core^{1,3,22}.

Here we present the results of experiments that determine the metal–silicate partition coefficient of carbon at conditions that are likely similar to that of metal–silicate equilibration/segregation during core formation²³, and that are at much higher pressures and temperatures than those of previous studies. We use the results in a state-of-the art model that combines accretion and core–mantle differentiation^{9,10}, and so we show that the present BSE carbon concentration is the direct consequence of core formation.

We performed laser-heated diamond anvil cell (LH-DAC) experiments to reproduce the pressures and temperatures of the putative conditions of Earth’s core-mantle differentiation in the laboratory. We performed experiments at 49–71 GPa and 3600–4000 K for a few tens of seconds using a double-sided laser-heated system (see Methods section for details). After melting and quenching at high pressure, our recovered samples consisted of a central metallic sphere surrounded by quenched silicate glass, as observed in previous similar studies^{24–26} (Fig. 1a). Samples were prepared for analysis using a Focused Ion Beam system and analysed for major elements using Electron Probe Microanalysis (EPMA see Methods section). To measure the concentrations of carbon in the quenched silicate melt, we used NanoSIMS, for which we

synthesized appropriate standards (the carbon contents of which were analysed by FTIR; see the Methods section).

We calculated the oxygen fugacities of our experiments relative to the iron-wüstite redox buffer (ΔIW), based on the assumption of ideal mixing of both phases (as justified by the high temperatures):

$$\Delta IW = 2 \log \frac{X_{\text{FeO}}}{X_{\text{Fe}}}. \quad (1)$$

Oxygen fugacities of our experiments lie between -0.9 and -1.5 log units relative to the IW buffer (Table 1). The concentration of carbon in the metal lies in the range from 5 to 9.4 wt.% and in the quenched silicate melt varies from 545 to 2800 ppm. The concentrations of carbon in the quenched silicate melt are much higher than concentrations observed in low P–T experiments which are typically 10–200 ppm^{1–3,8}. Thus, metal–silicate partition coefficients for carbon are 24–166 (Table 1) and are 1–2 orders of magnitude lower than those determined in previous low-pressure studies.

In order to understand the metallic phase better, we obtained images with SEM and TEM of the sample BAS C 39 (Fig. 1), and we performed electron diffraction analyses that show the presence of stoichiometric Fe₇C₃ inclusions in the quenched metal (Fig. 1b). To further investigate the quench textures (Fig. 1c), we performed scanning TEM-energy dispersive spectroscopy (EDS) mapping to identify Si-O-rich exsolved inclusions in the metal phase (Fig. 1d) along with a C-rich domain that could be an exsolved diamond, as also observed in previous similar experiments²⁷.

To understand the effects of pressure and temperature on the metal–silicate partitioning of carbon, we assembled a dataset for our model from previous studies to supplement our new results (Fig. 2). However, only some published experiments are relevant to core formation because prior studies were often focussed on the effects of silicate melt composition², fO_2 ^{28,29}, and the interaction of carbon with elements in the metal⁷. The partitioning of carbon is affected

by the ratio of non-bridging oxygens to tetrahedral cations² (NBO/T), so we only included data from the literature for which NBO/T lies in a narrow range (0.5–1.5); the range for our samples is 0.56–1.34. We also only included data from previous low P–T studies with low concentrations (< 1 wt.%) of light elements (S, N, O, Si) in the metal because the interaction between carbon and other light elements may be significant⁶ but is poorly constrained. As a consequence, we also excluded very low fO_2 data ($\Delta IW < -3$), as they are often associated with high concentrations of Si in the metal. On the other hand, our high P–T samples contain significant concentrations of oxygen and silicon in the metallic phase, as is typical for high temperature experiments^{24,25}. By combining these two datasets, we thereby empirically include the interaction effects of these light elements in the regression presented below, because the concentrations of O and Si in metal are generally low during core formation at low P–T conditions^{11,31} consistent with the data^{1-3,8} selected from previous studies, whereas at high P–T conditions such concentrations become high^{30,31}, which is consistent with the present study. Interactions of carbon with oxygen and silicon in the metal of our experiments might be at least partly responsible for the scatter of our data (Fig. 2). Finally, we did not use literature data for which the silicate phase did not quench to a glass, because the quality of SIMS analyses may be compromised¹. In Fig. 2, we present a comparison of our data together with the selected data from the literature as a function of temperature and pressure.

The metal–silicate partition coefficient of carbon (D_C) is calculated as:

$$D_C = \frac{X_C^{\text{metal}}}{X_{(CO_3)_2}^{\text{silicate}}}, \quad (2)$$

where X is the mole fraction of the element in the phase of interest. D is a function of several parameters, including pressure (P), temperature (T), and oxygen fugacity.

To remove the effect of oxygen fugacity on D values, we use the distribution coefficient K_D :

$$K_D = \frac{D_C}{(D_{Fe})^{n/2}}, \quad (3)$$

where n is the valence of C when dissolved in silicate liquid, which is 4 for carbon in this case³². Based on the aforementioned data selection, we fitted previously–published data^{1–3,8} along with our new results to derive the following expression:

$$\log K_D = -1(\pm 0.5) + \frac{4842(\pm 920)}{T} + 31(\pm 19)\frac{P}{T}. \quad (4)$$

This relation shows that carbon becomes less siderophile with increasing depth along a magma ocean geotherm, and predicts K_D values fairly well for both low and high P–T experiments (see Extended Data Fig. S3).

We used the parameterization of Eq. 4 in a combined model of accretion and core–mantle differentiation¹⁰. This model is based on N-body accretion simulations that typically start with 80–220 Mars-size embryos (initially located between 0.7 and 3.0 AU) that are embedded in a protoplanetary disk consisting of a few thousand much smaller planetesimals that are initially distributed over a heliocentric distance of 0.7 to ~9.5 AU. In the accretion simulations, planets typically grow from the starting embryos through accretional collisions with other embryos and planetesimals. Each collision potentially involves a core formation event during which accreted metal equilibrates with silicate liquid in a magma ocean and then segregates to the proto-core. By modelling metal–silicate equilibration using a mass balance approach, the evolving compositions of the mantles and cores of all accreting bodies are tracked. The approach requires that the compositions of all starting bodies are defined, with refractory element compositions matching CI chondrites. An oxidation gradient in the protoplanetary disk defines the oxygen content of starting bodies, which is the main compositional variable. Embryos and planetesimals that formed within ~1 AU of the Sun are highly reduced but with increasing heliocentric distance are increasingly oxidized^{9,33}. Planetesimals originating from beyond ~4.5 AU are fully oxidized, contain no metal, and have an H₂O content of ~20 wt.%. An important

feature of the model is that accreted metal only equilibrates with a small fraction of the silicate mantle – which is quantified using a hydrodynamic model of the interaction of metal and silicate as material from the impactor’s core sinks in a magma ocean³⁴.

Here we use the “Grand Tack”^{35,36} accretion simulation “4:1-0.5-8”⁹ with the same model parameters as in our previous study¹⁰. We first assume that carbon is accreted to Earth only in bodies that originate beyond 4.5 AU and that their average C concentration is 1.6 wt.% (an approximately mean carbonaceous chondrite (CI, CM, CO and CV) value). The mantle C abundance increases throughout accretion from a very early stage and finally reaches the BSE value of 80–120 ppm (Fig. 3). However, in this model the final concentration of C in the core is 47 ppm which means that the carbon concentration of the bulk Earth is only 83 ppm, i.e. far below the lowest estimate of 520 ppm²⁰. We have therefore developed a scenario that results in a bulk concentration of ~520 ppm. We assume that bodies from beyond 4.5 AU contain on average 1.6 wt.% carbon and that C concentrations in bodies originating at < 4.5 AU decrease along a linear gradient to reach a value of zero at a heliocentric distance of 1.88 AU. This is based on the condensation of volatile elements such C increasing with heliocentric distance as temperatures decrease and is analogous to a concentration gradient derived for sulfur¹⁰. The model then predicts a final mantle C concentration of 314 ppm (Fig. 3), a core concentration of 950 ppm and a bulk Earth concentration of 518 ppm. This calculation is not unique because the concentration (1.6 wt.% C) and distance (1.88 AU) parameters are of course correlated. Furthermore, if water-rich outer solar system planetesimals have concentrations of C exceeding 1.6 wt.%, then it is possible to obtain mantle concentrations consistent with the estimate of 765 ± 300 ppm²⁰ because C delivered to Earth’s mantle is inefficiently extracted by subsequent core formation similar to S and the HSEs¹⁰, as explained below. In summary, the final mantle abundance of Earth exceeds 200 ppm for all plausible C concentrations of the initial bodies in the protoplanetary disk.

A BSE carbon concentration of > 200 ppm is considerably higher than the estimate of 100 ± 20 ppm^{4,5}. If the 100 ± 20 ppm value is correct, an additional process, that is not included in our model, may have transported excess carbon to the core. There is experimental evidence that the pressure-induced disproportionation of $\text{Fe}^{2+} \rightarrow \text{Fe}^{3+} + \text{Fe}$ in deep magma oceans causes the formation of dispersed Fe metal droplets^{37,38}. The gravitational segregation of these metal droplets could have transported the excess carbon to the core, leaving ~ 100 ppm in the mantle and resulting in 0.14 wt.% C in the core.

There are two critical features of the accretion/differentiation model that explain why the BSE carbon concentration is not efficiently extracted during core formation in spite of its siderophile behaviour. First, bodies originating from the outer solar system are water ice-bearing and so fully oxidize and contain no metal. Thus, there is no core formation event when they are accreted and the delivered C remains in the magma ocean/mantle. Second, when differentiated metal-bearing bodies are accreted, the metal of the impactor only equilibrates with a fraction of the target's mantle, of which the value depends on the size of the impactor's core and the depth of the magma ocean^{9,34}. The fractions of the mantle that equilibrate with metal in the present simulation are 0.16 to 2.5% for planetesimal impacts and 2.6 to 9.9% for embryo impacts. Consequently, core formation events are extremely inefficient at removing carbon from the bulk of the mantle and transferring it to the core. These two critical features are absent in the models of "continuous core formation" that are currently applied in many studies of Earth's differentiation^{12,13}; the result of a continuous core formation model for carbon is presented in the supplementary information for comparison. In such models all accreted material is assumed to contain metal and it is assumed (unrealistically) that all accreted metal equilibrates chemically with the entire Earth's mantle. Consequently, such models predict that essentially the entire BSE carbon budget has to be delivered during late accretion, i.e. after core formation has ended (see supplementary Information and Extended Data Fig. S3). We know of no

astrophysical accretion simulations that would support this scenario which requires the accretion of carbon-poor material during core formation and carbon-rich material afterwards during late accretion.

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Table

Table 1: Experimental conditions of carbon partitioning experiments.

Run #	BAS C 50	BAS C 60	BAS C 39	BAS C 42
P (GPa)	60	71	49	56
T (K)	3800	4000	3600	3700
ΔIW	-1.06	-1.2	-0.89	-1.51
D_c	149 ± 7	69 ± 14	24 ± 1	166 ± 6
K_D	12.9 ± 1	5 ± 1.4	3.0 ± 1.0	2.9 ± 1.1
nbo/t	0.95	0.58	1.34	0.56

ΔIW is the calculated fO_2 relative to the IW buffer (see text for details).

D_c stands for the metal–silicate partitioning of carbon (mol.% of carbon in the metal / mol.% carbon in the silicate).

K_D is the distribution coefficient where $K_D = \frac{D_c}{(D_{Fe})^{n/2}}$, nbo/t is the number of non-bridging oxygen atoms over tetrahedrally coordinated cations.

The uncertainties on temperature and pressure measurements are estimated to be about 300 K and 5 GPa respectively.

Figures

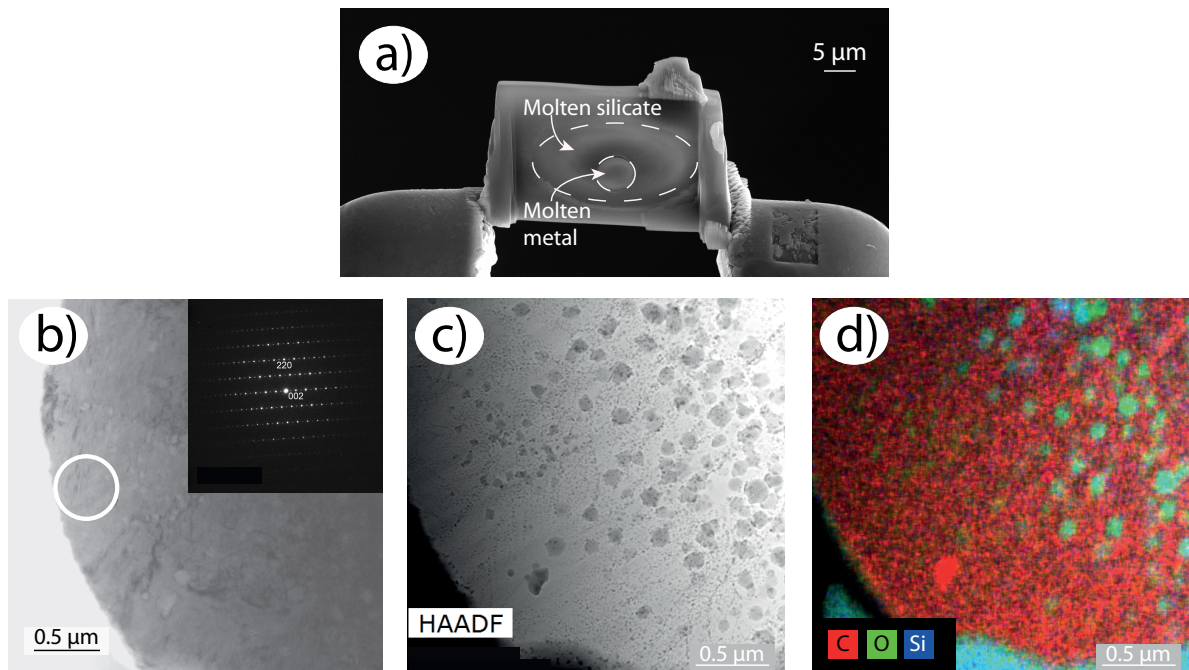


Figure 1: Images of run BAS C 39 synthesized at 49 GPa and 3600 K. a) Secondary electron image where two different phases can be observed in the sample: the molten silicate and the molten metal that were equilibrated at high pressure and high temperature. b) Bright field TEM image and selected area electron diffraction pattern (Inset) of the metallic phase showing the presence of stoichiometric Fe_7C_3 (also confirmed with EELS). c) High-angle annular dark field of the same sample showing exsolutions that developed during quenching along with d) Scanning TEM-EDS map of the same area which shows exsolved inclusions that are oxygen rich (green) and silicon-rich (blue). The bright red inclusion in image c) is believed to be an exsolved diamond²⁷.

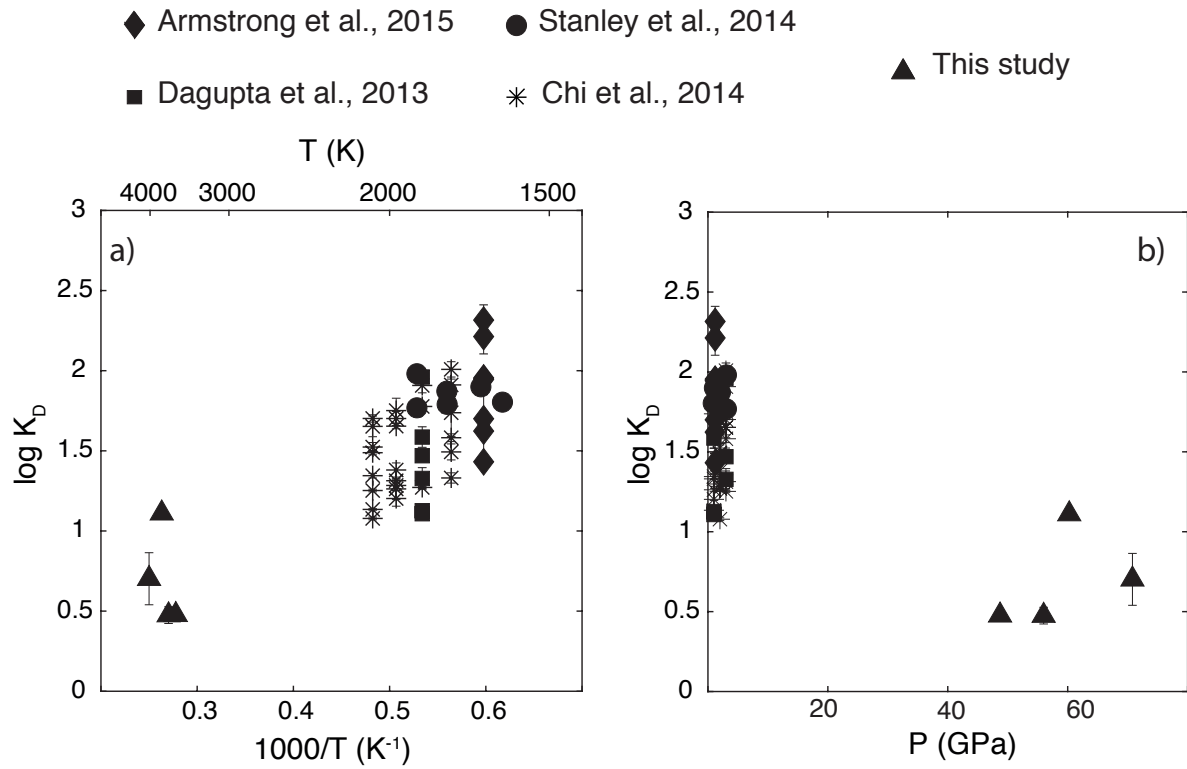


Figure 2: Comparison of our data with results of previous studies with respect to temperature (a) and pressure (b). Uncertainties on temperature and pressure measurements for diamond anvil cell experiments are of 300 K and 5 GPa respectively.

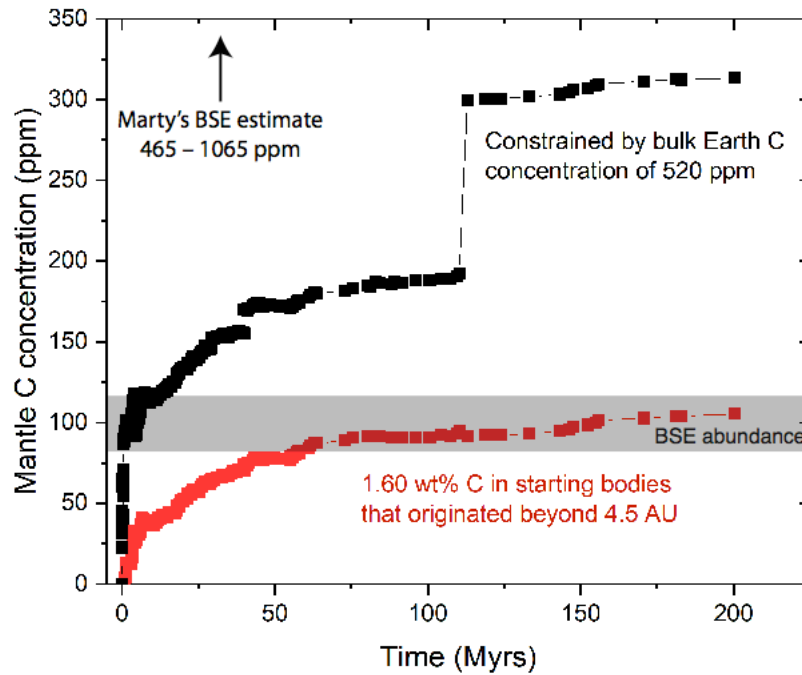


Figure 3: *Mantle C concentration during accretion as a function of time based on a combined accretion/differentiation model^{9,10}. The red symbols show the evolution of mantle carbon concentration assuming that all starting planetesimals that originated from beyond 4.5 AU contained an average of 1.6 wt.% C; bodies originating at heliocentric distances < 4.5 AU are assumed to contain zero C. The black symbols show the evolution when the average carbon concentration in starting bodies originating from beyond 4.5 AU is 1.6 wt.% and C concentrations in bodies originating at < 4.5 AU decrease along a linear gradient to reach a value of zero at a heliocentric distance of 1.88 AU. The prominent jump at 110 Myrs is caused by the final giant impact. Estimates of the BSE concentration range from 100 ± 20 ppm⁴, as shown, to 765 ± 300 ppm²⁰.*

1 **Methods**

2 **a) Starting materials**

3 For the silicate starting glass, we mixed oxide and carbonate compounds in stoichiometric
4 proportions, decarbonated over night at 900°C, then added FeO. Powders were ground under
5 acetone for some time to ensure a homogeneous composition. We created pellets that were
6 subsequently fused in argon flux for several seconds at about 1400°C using a levitation furnace
7 device in Orléans, France³⁹. Chemical composition of the glass was checked using EPMA
8 (“Basalt” in table S1) and we could verify that no compositional variation was detectable. The
9 glass sphere was subsequently polished down to a thickness of 20 µm and machined in IPGP,
10 France, to obtain small disks of 80 microns diameter²⁵.

11 The goal of this study is to understand the effect of core formation on the partitioning of carbon
12 between Earth’s most important reservoirs: the silicate mantle and the metallic core. Carbon is
13 a prominent contaminant in the laboratory, with multiple sources of contamination, from the
14 use of ethanol and acetone, various type of glue or ambient contamination. It has also been
15 suggested that diamonds used as pressure transmitting tool can also diffuse carbon to the
16 samples⁴⁰. To better understand the origin of carbon, we used ¹³C as the source of carbon in our
17 experiments. We synthesized the metallic phase using piston-cylinder apparatus. In order to
18 assess the extent of carbon contamination during LH-DAC experiments, we only used ¹³C. We
19 mixed 5 wt.% of ¹³C powder (97% pure from Cambridge Isotope Laboratories, Inc.) with 95
20 wt.% of Fe. We melted this alloy at 2 GPa and 1873 K in MgO capsule for 10 minutes using a
21 ½” piston-cylinder assembly at BGI. This ¹³C-doped carbide was then crushed to be used for
22 the metallic part of the for laser-heated diamond anvil cell experiments. The composition was
23 checked by EPMA (see table S1, “Metal”) using an Fe₃C standard (see sections d) and e)).

24

25 **b) Laser-heated diamond anvil cell experiments (LH-DAC)**

26 We used diamonds with 250 μm culets and rhenium gaskets. Gaskets were pre-indented to
27 obtain a thickness of about 40 to 50 μm , and subsequently laser-drilled to have an experimental
28 chamber of about 90 μm in diameter. We then loaded two silicates discs encapsulating a flake
29 of carbide²⁵. Samples were compressed to the target pressure, and then laser-heated at the
30 desired temperature using a doubled sided laser heating system at BGI. Temperature was
31 generated by two fiber continuous-wave (CW) YAG lasers (SPi©) with wavelength of 1064
32 nm and delivering 100 watts each. The laser beam was focused onto the sample using two NIR-
33 Mitutoyo objectives lenses with x20 magnification. Temperature was measured simultaneously
34 and continuously on both sides of the diamond cell using a spectro-radiometric technique⁴¹
35 using a 2500i spectrometer and pixis 400 CCD camera from Princeton-Instrument© from the
36 light collected through both objective lenses. The peak temperature, above the liquidus
37 temperature of the material, was maintained for few tens of seconds, before switching off the
38 electric power of the lasers to quench the experiment. The samples quenched extremely quickly
39 due to the high thermal conductivity of diamonds and the cells were subsequently decompressed
40 slowly over several hours. As observed in previous studies using LH-DAC, the silicate part
41 quenched to a glass, and the Fe-rich metal presents Si- and O-rich inclusions^{24,25}.

42

43 **c) Sample recovery**

44 We used the Focused Ion Beam facility (FIB, dual beam Scios FEI©) in Bayerisches Geoinstitut
45 (BGI) to recover small lamellae of 3–5 μm thickness, about 40 μm in length and 20 μm high^{24,25}.

46 In the samples, we observed two quenched phases that had been molten at high pressure and
47 high temperature: a silicate and metallic one (see Fig. 1-a in the main text).

48

49 **d) Fabrication of analytical standards**

50 For EPMA analysis of the carbon contents of metals, we synthesized a Fe_3C (cementite) carbide
51 standard at BGI⁴², by inserting a 1 mm diameter, 10 mm long, 99.99 % purity Fe wire in a thick
52 graphite sleeve and heating it to 1423 K at 15 kbar in a piston cylinder for one week. The
53 reaction product was confirmed to have the Fe_3C cementite structure by XRD and was assumed
54 to be stoichiometric.

55 For the nanoSIMS analysis of quenched glasses, we synthesized ^{13}C -doped glass standards⁴³.
56 We first mixed oxides in basaltic proportions and made a glass at 1600°C in a furnace for two
57 hours. This glass was analyzed by EPMA for major elements, and by FTIR to confirm that it
58 was carbon-free (“B1” in table S1). Subsequently, we performed piston-cylinder experiments
59 to incorporate ^{13}C into the basaltic glass. The source of ^{13}C was chosen as oxalic acid enriched
60 in ^{13}C (Cambridge Isotope Laboratories, Inc.), that was loaded along with the basaltic glass in
61 a Pt capsule and pressurized to 2 GPa and heated to 1600°C for 10 minutes⁴³. We created two
62 glass standards (B1145 and B1147) containing 785 ppm and 1263 ppm of carbon respectively,
63 as measured by FTIR at BGI. For the FTIR measurements, we used an extinction coefficient⁴⁴
64 of 69500 $\text{L}\cdot\text{mol}^{-1}\cdot\text{cm}^{-1}$. We report the FTIR measurements along with EPMA analyses on table
65 S1 for the carbon-free (B1) and the two carbon-doped glasses.

66

67 **e) EPMA analysis**

68 Major elements from both the metallic and silicate phases were analyzed by Electron Probe
69 Microanalysis (EPMA). We calibrated using wollastonite, spinel and olivine for Ca, Al and Si,
70 Mg, Fe in the silicate respectively with current of 15 keV and 15 nA, a focused beam and a phi-
71 rho-Z correction. For the metallic phase, we used a pure Fe wire, FeSi, Fe_2O_3 and Fe_3C
72 standards to calibrate for Fe, Si, O and C respectively, at 15 keV with a 25 nA beam current,
73 again using a focused beam and phi-rho-Z correction. The counting time was 10 seconds on the

74 background, and 20 seconds on the peak for all elements, and the samples and Fe₃C standard
75 were not carbon coated.

76 Analyzing carbon using the microprobe requires special consideration⁴⁵. In order to quantify
77 the carbon content of both the experimental metals and the metal starting material, we
78 synthesized a Fe₃C carbide primary standard (see section d)). During C analysis, we took
79 particular care in repeating the measurements on the metallic standards to quantify background
80 levels (contamination) inside the EPMA. We measured the count rate from the C α peak on a
81 pure Fe standard continuously for ten minutes, and found no statistically-significant change in
82 peak height over that timescale⁴². We performed repeated measurements of pure iron wire to
83 check for carbon contamination, and observed a background for carbon of 0.49 ± 0.07 wt.%
84 which is well below the amount of carbon present in the metallic phase of our samples. The
85 same was done with pure FeSi, for which we measured 0.36 ± 0.05 wt.% of carbon in the
86 standard. As for oxygen contamination, we measured a background of 0.04 ± 0.04 wt.% in the
87 Fe wire and 0.15 ± 0.05 wt.% in FeSi. Those values are again well below the amount measured
88 in our samples. Standard analyses were consistent and stable over the whole duration of the
89 analytical sessions.

90 Analyses performed on the quenched silicate highlighted high concentrations of FeO (from 14
91 to 27 wt.%), as observed previously in similar experiments^{24,25}. In the metallic phase, we
92 observe a range of concentrations of carbon, from 5 to 9.5 wt.%, depending on the sample.
93 Oxygen and silicon are also present in the metallic phase, from 2.8 to 7.5 wt.% oxygen and 0.3
94 to 5.6 wt.% silicon. We report on tables S2 and S3 the full analyses of our recovered runs for
95 silicate and metal respectively. Whilst some analytical totals are low, these low totals are not
96 explained by the thin samples, as a 3 μ m deep lamella is thick enough to not lose electrons, nor
97 significant X-ray flux, through the base⁴⁶.

98

99 **f) nanoSIMS analysis**

100 The abundance of carbon in the silicate phase was expected to be extremely small, due to its
101 high siderophilicity¹. Laser heated diamond anvil cell experiments produce very small samples,
102 with silicates phases that are only few micrometers wide. Therefore, NanoSIMS was used in
103 this study to achieve required analytical precision and spatial resolution and quantitatively
104 analyze the low carbon contents in the silicate phase.

105 NanoSIMS measurements are sensitive to matrix effects, so we took particular care in
106 synthesizing and using relevant standard glasses (see section d)). Along with those two
107 standards, we used two standards of natural rocks containing natural and known amounts of
108 carbon: a basaltic glass from the East Pacific Rise^{47,48} ALV 981-R23 containing 405 ppm of
109 CO₂; and a piece of glass from D'Orbigny meteorite^{49,50} containing 40 ppm of C. Composition
110 of those standards and of the natural samples are given in table S1. During the NanoSIMS
111 session, we analyzed the standards each day to check for consistency and demonstrate the
112 reproductivity of the measurements.

113 NanoSIMS analyses were performed using the CAMECA NanoSIMS 50L at the Open
114 University (Milton Keynes, UK). Prior to NanoSIMS analysis, a large area of each sample was
115 pre-sputtered using a focused primary beam of 16 KeV Cs⁺ ions with probe current of 100 pA
116 to remove surface contamination. The sizes of the pre-sputtered areas vary from 5×5 to 7×7
117 μm². Analyses were then carried out in spot mode first by scanning a 50 pA Cs⁺ beam onto the
118 inner 3×3-5×5 μm² areas, with secondary ions of ¹²C⁻, ¹³C⁻, ³⁰Si⁻, ²⁴Mg¹⁶O⁻, ²⁷Al¹⁶O⁻ and
119 ⁵⁶Fe¹⁶O⁻ collected in electron multipliers simultaneously. Only the data from the inner 60% of
120 the area were collected using electron gating to avoid carbon contamination from the
121 surrounding areas. Each spot analysis consisted of 200 measurements, with a total analysis
122 time of ~2 minutes.

123 After spot analyses, selected sample areas were also measured in imaging mode. A 50 pA Cs⁺
124 beam was also used in imaging mode, with five secondary ions (¹²C⁻, ¹³C⁻, ³⁰Si⁻, ⁵⁶Fe¹³C⁻ and
125 ⁵⁶Fe¹⁶O⁻) monitored simultaneously. Each imaging analysis consisted of 10 frames. A frame
126 size of 256×256 pixels was used for all images with an integration time of 500 ms per pixel,
127 leading to total analysis time of 5-6 minutes for each image.

128 The mass resolving power (m/Δm) for both spot and imaging mode was set to 9000 (CAMECA
129 definition), sufficient to resolve all interferences from neighbouring mass peaks, such as ¹³C
130 from ¹²CH. We could extract from the region of interest (ROIs) the ratios of ¹³C/³⁰Si and
131 ¹²C/³⁰Si. We derived a precise calibration line using those standards, which was used to infer
132 the amount of carbon (both ¹²C and ¹³C, see figures S1-a and S1-b) present in each of our
133 samples. During analyses, we carefully avoided metal blebs that are sometimes present in the
134 silicate phase in this style of experiments, interpreting them to be entrained rather than exsolved
135 on quench. We present in Fig. S2 NanoSIMS ion images of the metallic phase of sample BAS-
136 C-39, to highlight the level of detail that can be reached with NanoSIMS.

137 Some uncertainty of the NanoSIMS measurements may derive from the difference between the
138 composition of the samples and the one of the standards. We calculated the amount of carbon
139 in each of our samples using the C/Si ratio of the standards and the experimental runs for both
140 ¹²C and ¹³C. The standards we used contained between 40 and 50 wt.% of SiO₂ whereas the
141 runs contained between 35 and 40 wt.% SiO₂. To encompass the range of potential sources of
142 uncertainty, we have assigned an uncertainty of ± 15% to the nanoSIMS measurements.

143 Analysing our samples and standards with the NanoSIMS revealed that the carbon present in
144 both cases was not only ¹³C but also ¹²C, with a lower ¹³C/¹²C ratio in experimental silicates
145 than in the metal (values from 2.8 to 6.9 for metal versus 0.37 to 1.7 for silicate, see Fig. S1-c).
146 The presence of significant concentrations of ¹²C was unexpected, since we carefully prepared
147 the samples using only ¹³C, but demonstrates that carbon contamination is probably ubiquitous

148 in high-temperature DAC experiments. The silicate starting powders were heated overnight at
149 900 °C and then converted to glass using an aerodynamic levitation system at temperatures of
150 about 1400 °C. Thus, despite several hours at elevated temperatures, it is extremely hard to
151 eliminate the presence of environmental carbon. ^{12}C could originate from the diamond anvils,
152 the carbonate powders being not fully decarbonated, by sorption of carbon from the atmosphere
153 or from residues of organic carbon contamination during sample preparation.

154 The $^{13}\text{C}/^{12}\text{C}$ measured in the experimental silicates was lower than in the corresponding metals
155 (Fig. S1-c). This indicates that experiments did not reach isotopic equilibrium, and implies that
156 the most significant source of carbon contamination was the silicate starting material rather than
157 the diamond anvils. This observation is contrary to what has been proposed in previous LH-
158 DAC studies^{26,40}. We also measured silicate in unmelted regions in DAC samples BAS C 39
159 and BAS C 42 at the NanoSIMS, that contained 386 and 347 ppm of carbon ($^{12}\text{C} + ^{13}\text{C}$)
160 respectively, highlighting the presence of carbon in the starting material. We demonstrate here,
161 for the first time with support of NanoSIMS measurements, that carbon contamination in our
162 DAC experiments is not due to carbon diffusion from the diamond anvils to the samples, but is
163 from the samples themselves.

164

165 **h) TEM analysis**

166 We performed transmission electron microscope (TEM) analyses on a FEI Titan G2 80-200
167 S/TEM equipped with X-ray energy dispersive spectrometer (EDS) and electron energy-loss
168 spectrometer (EELS) in Bayerisches Geoinstitut in order to observe and further analyze the
169 experimental metals. To do so, we further thinned down one of our samples to a thickness of
170 60-80 nm using FIB. The FIB lamella was plasma cleaned prior to TEM analysis in order to
171 remove surface contamination. We determined Fe/C ratios in our sample and also perform EDS
172 mapping in the metallic phase. The quantification of the EELS analyses followed the procedure

173 previously described^{40,51}, using experimentally-determined ratios of partial cross sections of C
174 K and O K edges against the Fe L edge versus sample thickness, which were calibrated with
175 synthetic Fe_{0.94}O and Fe₃C samples. Note that due to the inhomogeneity of the metallic phase
176 of the sample (see Fig. 1-a,-b,-c), EELS technic cannot provide a measurement of the bulk
177 carbon contents of this phase.

178

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- 215

216 **Extended Data Tables and Figures**

217

218 *Table S1: Composition of the synthetic basaltic starting silicate used for the LH-DAC*
 219 *experiments (“Basalt”), the standards for the NanoSIMS synthesized with piston-cylinder at 2*
 220 *GPa and 1600 °C (“B1145” and “B1147”), the carbon-free starting material for piston-*
 221 *cylinder experiments (“B1”) and the starting metal used in LH-DAC experiments. We also*
 222 *report the composition of the two natural samples that were used at the NanoSIMS as*
 223 *standards. b.d. stands for below detection limit.*

Samples	Basalt	B1145	B1147	B1	D ^o Orbigny glass ^a	ALV 981-R23 ^c	Metal
wt.%	N=10	N=7	N=7	N=7	-	-	N=14
SiO ₂	52.18 (0.15)	48.05 (0.18)	50.89 (0.27)	53.56 (0.13)	41.7	49.63	C 5.67 (1)
Na ₂ O	1.90 (0.03)	1.80 (0.06)	1.86 (0.06)	2.03 (0.03)	-	2.91	Fe 91.61 (1.31)
CaO	10.92 (0.07)	9.47 (0.10)	10.03 (0.08)	10.70 (0.05)	23.3	11.61	
FeO	11.35 (0.04)	13.65 (0.20)	10.71 (0.06)	10.04 (0.06)	11.0	8.16	
Al ₂ O ₃	14.66 (0.07)	16.37 (0.11)	15.22 (0.10)	15.48 (0.05)	19.6	16.60	
MgO	7.45 (0.07)	7.22 (0.08)	7.57 (0.10)	7.73 (0.06)	1.87	8.36	
C (ppm)	-	785 (8)*	1263 (13)*	b.d.*	40	111	
Total	98.46 (0.20)	96.66 (0.19)	96.41 (0.33)	99.55 (0.15)	99.3 ^b	98.96 ^d	97.28 (0.76)

224

*Measured using FTIR (see text for details)

225

^aComposition from ref.⁵⁰

226

^bTotal includes also 2 wt.% of TiO₂, Cr₂O₃, MnO

227

^cComposition from refs.^{47,48}

228

^dTotal includes TiO₂, MnO, K₂O and P₂O₅

229

230 *Table S2: Compositions of quenched silicate liquid. Major elements were analysed using*

231 *EPMA and carbon was analysed using NanoSIMS. See text for details.*

Run #	BAS C 50	BAS C 60	BAS C 39	BAS C 42
wt.%	N=4	N=5	N=8	N=6
SiO ₂	35.84 (0.38)	39.97 (1.98)	34.09 (0.83)	39.66 (0.99)
MgO	8.85 (0.20)	8.60 (0.28)	9.53 (0.51)	9.02 (0.26)
Al ₂ O ₃	19.46 (0.12)	25.93 (0.50)	16.34 (1.36)	24.48 (0.64)
CaO	4.02 (0.16)	4.89 (0.12)	5.67 (0.20)	5.68 (0.31)
Na ₂ O	2.13 (0.07)	2.22 (0.10)	2.32 (0.15)	2.53 (0.11)
FeO	23.82 (1.40)	18.60 (0.60)	27.35 (2.55)	13.71 (1.72)
C (ppm)*	312 (47)	545 (82)	2805 (420)	272 (41)
Total	94.12 (1.85)	100.2 (1.96)	95.30 (2.34)	95.07 (1.55)

232

*15% uncertainties estimated on NanoSIMS measurements (see text for more details).

233

234 *Table S3: Composition of the metal phase of our recovered runs analysed using EPMA. See*

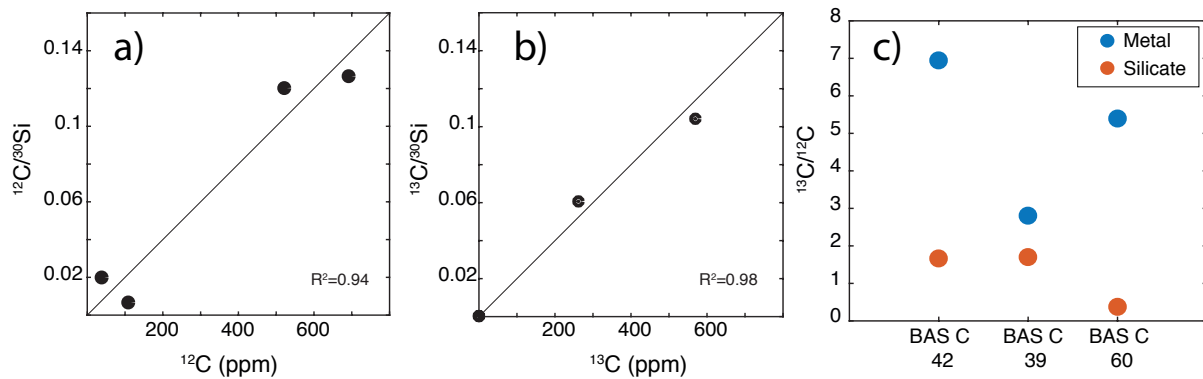
235 *text for details.*

Run #	BAS C 50	BAS C 60	BAS C 39	BAS C 42
EPMA (wt.%)	N=3	N=3	N=4	N=2
Fe	81.68 (1.25)	79.56 (0.9)	85.96 (0.58)	78.01 (0.78)
Si	0.46 (0.16)	5.56 (0.16)	0.33 (0.11)	2.65 (0.40)
O	3.79 (0.89)	7.56 (0.37)	2.85 (0.48)	3.42 (0.10)
C	5.90 (0.18)	4.99 (0.28)	9.44 (0.24)	5.35 (0.20)
Total	91.83 (0.41)	97.66 (0.38)	98.57 (0.29)	89.43 (0.48)

236

237

238



239

240 *Figure S1: Results from NanoSIMS measurements. a) and b) calibration lines obtained for*

241 *^{12}C and ^{13}C respectively using our standards and two natural samples (d'Orbigny and ALV*

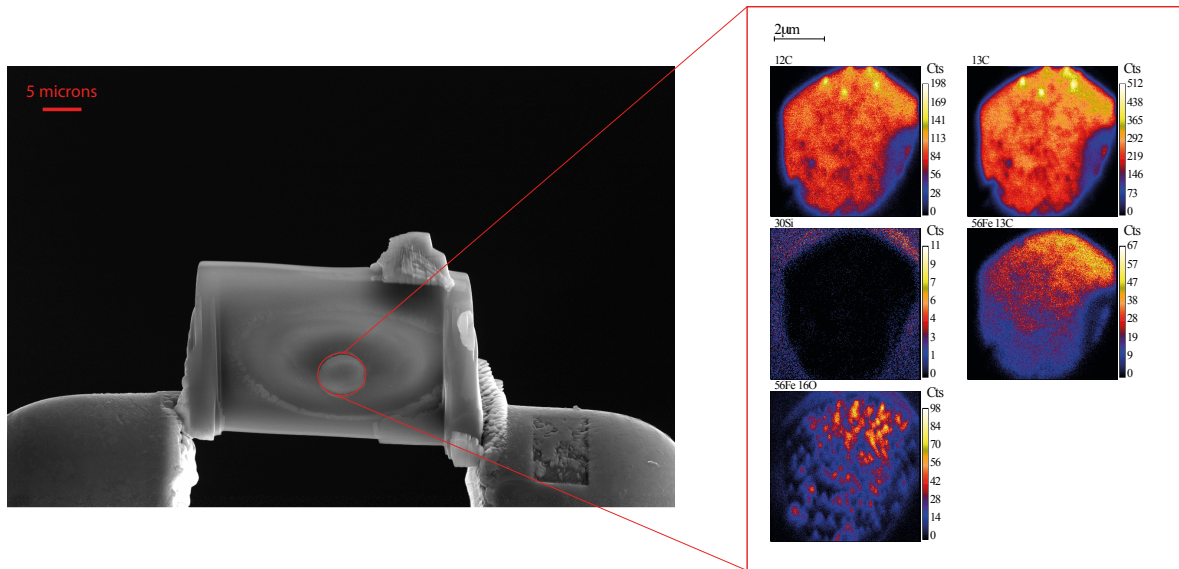
242 *981-R23) and both FTIR and NanoSIMS measurements. Note that for ^{13}C measurements, the*

243 *two natural samples plot at 0. c) Illustration of the isotopic disequilibrium between the*

244 *metallic and the silicate phase for three of our recovered runs which illustrates the presence*

245 *of ^{12}C in our reacted samples, despite the use of pure ^{13}C in the starting material.*

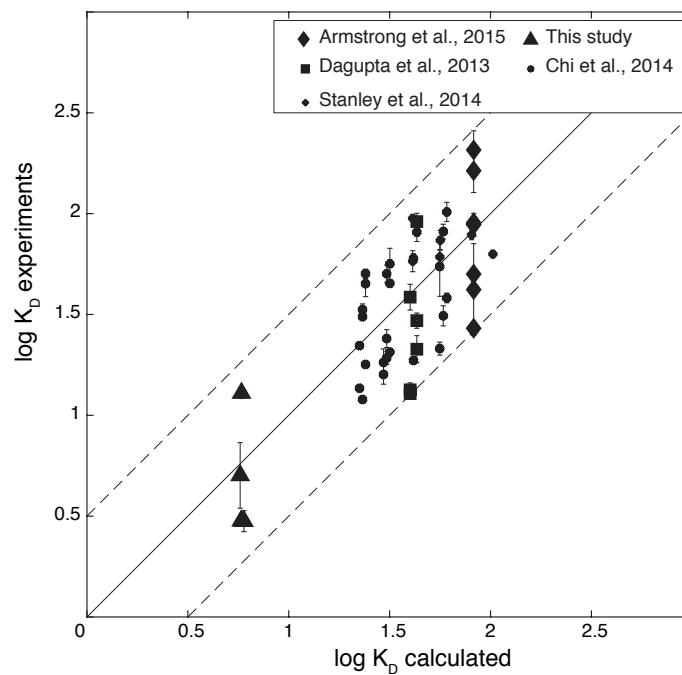
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247

248 *Figure S2: Secondary electron image of sample BAS-C-39 along with NanoSIMS mapping of*
 249 *the metallic phase.*

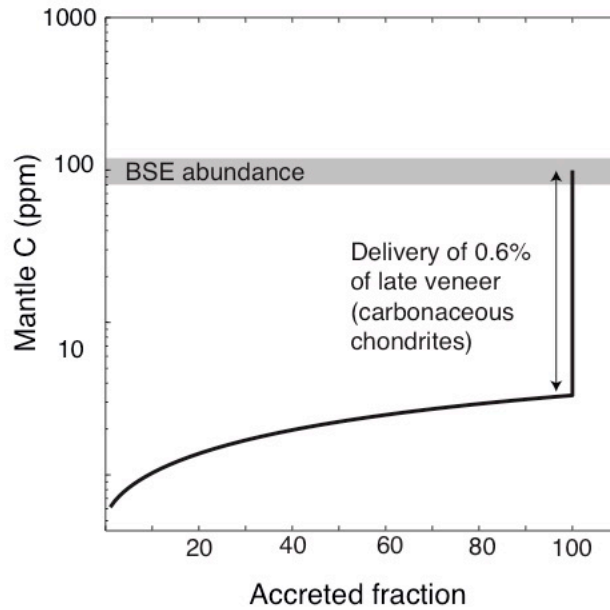
250



251

252 *Figure S3: Validation of our partitioning model. Comparison of the K_D values (see eq. 3 in*
 253 *the main text) calculated from the experiments and from our model (eq. 4 in the main text).*

254



255

256 *Figure S4: Continuous core formation model. Evolution of the concentration of carbon in*
 257 *the mantle during Earth's core–mantle differentiation and the late accretion stage based on a*
 258 *continuous core formation model. The estimated BSE concentration of 80–120 ppm is shown*
 259 *in both cases⁴.*

260

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273

274 **Authors contributions**

275 I.B. conceived the project, performed experiments, analyzed data and derived the
276 thermodynamic model. E.S.J. analyzed data. I.A.F., X.Z. performed NanoSIMS measurements.
277 S.P. performed laser heating experiments and helped with FIB. N.M. performed TEM
278 measurements. D.C.R. and S.A.J. built the accretion model. I.B., E.S.J. and D.C.R wrote the
279 manuscript.

280

281 **Additional information**

282 **Supplementary Information** is available for this paper. Correspondence and requests for
283 materials should be addressed to Ingrid Blanchard (ingrid.blanchard@uni-bayreuth.de).

284

285 **Supplementary Information**

286 Using our high pressure and high temperature experiments along with the ones of the literature,
287 we built a new parameterization for the metal–silicate partitioning of carbon (Eq. 4 in the main
288 text). We compare in Fig. S3 the agreement between the experimental K_D and the one predicted
289 by our regression. Then, we utilized two styles of modelling in this study. The first one is the
290 one presented in the main text, and the second one is a more simple continuous core formation
291 model that has been applied in many geochemical studies^{1–3}. Here, Earth is assumed to accrete
292 material in small (1%) increments and each batch of accreted metal is assumed to equilibrate
293 chemically with the entire mantle. The equilibration pressure is assumed to correspond to
294 conditions the base of the magma ocean and increases as the planet grows. The temperature at
295 the base of the magma ocean, which is constrained by the peridotite solidus–liquidus, also

296 increases as the Earth grows. For simplicity, we set the fO_2 during accretion at IW-2.3 as given
297 by the current mantle/core FeO/Fe ratio⁴.

298 The concentration of carbon in the mantle today is estimated to lie between 80 and 120 ppm
299 (Dasgupta and Hirschman 2010; Hirschman 2016). On the other hand, the bulk Earth carbon
300 content has been proposed to lie between 520 ppm⁵ and 750 ppm⁶. Here, we assume that each
301 stepwise-accretion step delivers material with a bulk C content that matches the bulk Earth
302 value⁵. Using the higher value of 750 ppm⁶ would not significantly change the results. The
303 partition coefficient of carbon is calculated after each equilibration step following eq. 4, and
304 the resulting amount of carbon present in both mantle and core is derived from this calculation.
305 Hence, carbon is present throughout accretion.

306 The Earth is believed to have been bombarded by bodies comprised of chondritic material after
307 core formation ceased (the so-called late accretion stage which resulted in addition of the “late
308 veneer”) in order to explain the concentrations of highly siderophile elements present in the
309 mantle⁷⁻⁹. The amount of material that was delivered to Earth during late accretion was about
310 0.6% of Earth’s mass¹⁰, and was not incorporated into the core, but rather remained in the
311 mantle. Chondrites are carbon-rich, with up to 3.2 wt.% of carbon reported in CI chondrites¹¹,
312 so their delivery to Earth brought a substantial amount of carbon. We incorporated the effect of
313 such late accretion in our model by assuming that the carbon content of the late veneer matches
314 the mean carbonaceous chondrite (CI, CM, CO and CV) value of 1.6 wt.%.

315 Results of the continuous core formation model are presented in Fig. S4. The combined effects
316 of core formation and later addition of late veneer are needed in order to reproduce the carbon
317 content of the BSE. Core–mantle differentiation effectively strips the mantle of all carbon and
318 the addition of C during late accretion is required to achieve the mantle carbon concentration.
319 According to this model, the resulting concentration of carbon in the core after the combined
320 effect of differentiation and late veneer is 0.16 wt.%. This result is consistent with previous

321 suggestions, based on first principles molecular dynamics¹² and experiments¹³, that carbon is
322 not a dominant element in Earth's core. As discussed in the main text, this simple scenario is
323 not coherent with astrophysical accretion scenarios. This model assumes that the bulk Earth C-
324 content is known *a priori*, and the assumption that all accreted metal equilibrate with all the
325 mantle is also unrealistic.

326

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