

The Fate of Carbon during Earth's Core–Mantle Differentiation

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Abstract

Carbon is an essential element for the existence and evolution of life on Earth, constitutes up to 50% of dry biomass, and is likely a requirement for all life in the universe. Its high abundance in Earth's crust and mantle (the Bulk Silicate Earth, BSE) is surprising because carbon is strongly siderophile (metal-loving) and should have segregated almost completely into Earth's core during accretion^{1–3}. Estimates of the concentration of carbon in the mantle lie mostly in the range of 80–120 ppm^{4,5}, which is much higher than expected based on simple models of core–mantle differentiation^{2,6,7}. Here we show through experiments at 49–71 GPa and 3600–4000 K that carbon is significantly less siderophile at such conditions than at the low pressures (≤ 3 GPa) of previous studies^{1–3,8}. We derive a new parameterization of the pressure–temperature dependence of the metal–silicate partitioning of carbon and apply this in a state-of-the-art model of planet formation and differentiation^{9,10} that is based on astrophysical N-body accretion simulations. Results show that BSE carbon concentrations increase strongly starting at a very early stage of

26 **Earth’s accretion and, depending on the concentration of carbon in accreting bodies, can**
27 **easily reach or exceed estimated BSE values. In contrast, simple models of “continuous**
28 **core formation”^{11–13} require all BSE carbon to be accreted after core formation ended,**
29 **but this is not consistent with astrophysical models of accretion.**

30

31 The extents to which the volatile contents of terrestrial planets were determined by the
32 processes of accretion and core formation, impact-driven volatile loss, and addition of a late
33 veneer subsequent to core formation are highly controversial^{9,14–17}. The Earth is considered to
34 have formed from bodies of chondritic composition, although it is unlikely that any particular
35 chondrite group constituted the sole building blocks of the planet¹⁸. Different classes of
36 chondrites have varying carbon concentrations, but importantly, such concentrations are all
37 high, at the weight percent level¹⁹. This is far in excess of the concentration of carbon present
38 in the Bulk Silicate Earth (BSE) today, which is currently estimated to be 80–120 ppm^{4,5} or
39 possibly as high as 765 ± 300 ppm²⁰. Additionally, the metal–silicate partition coefficient of
40 carbon at low pressures and temperatures (≤ 13 GPa and ≤ 2000 K) is very high ($D > 1000$; see
41 D defined in Eq. 2 and ref.^{1–3,21}). Consequently, almost all carbon should have partitioned into
42 the core during its formation, leaving the mantle almost completely depleted in this element.
43 Several explanations have been proposed for the BSE C concentration, including crystallization
44 of C-rich phases (e.g. diamond) from a C-saturated magma ocean, imperfect metal–silicate
45 equilibration during core formation, accretion of a “late veneer” after the end of core formation,
46 and carbon outgassing from the core^{1,3,22}.

47 Here we present the results of experiments that determine the metal–silicate partition coefficient
48 of carbon at conditions that are likely similar to that of metal–silicate equilibration/segregation
49 during core formation²³, and that are at much higher pressures and temperatures than those of
50 previous studies. We use the results in a state-of-the art model that combines accretion and

51 core–mantle differentiation^{9,10}, and so we show that the present BSE carbon concentration is
52 the direct consequence of core formation.

53 We performed laser-heated diamond anvil cell (LH-DAC) experiments to reproduce the
54 pressures and temperatures of the putative conditions of Earth’s core-mantle differentiation in
55 the laboratory. We performed experiments at 49–71 GPa and 3600–4000 K for a few tens of
56 seconds using a double-sided laser-heated system (see Methods section for details). After
57 melting and quenching at high pressure, our recovered samples consisted of a central metallic
58 sphere surrounded by quenched silicate glass, as observed in previous similar studies^{24–26} (Fig.
59 1a). Samples were prepared for analysis using a Focused Ion Beam system and analysed for
60 major elements using Electron Probe Microanalysis (EPMA see Methods section). To measure
61 the concentrations of carbon in the quenched silicate melt, we used NanoSIMS, for which we
62 synthesized appropriate standards (the carbon contents of which were analysed by FTIR; see
63 the Methods section).

64 We calculated the oxygen fugacities of our experiments relative to the iron-wüstite redox buffer
65 (ΔIW), based on the assumption of ideal mixing of both phases (as justified by the high
66 temperatures):

$$67 \quad \Delta IW = 2 \log \frac{X_{\text{FeO}}}{X_{\text{Fe}}}. \quad (1)$$

68 Oxygen fugacities of our experiments lie between -0.9 and -1.5 log units relative to the IW
69 buffer (Table 1). The concentration of carbon in the metal lies in the range from 5 to 9.4 wt.%
70 and in the quenched silicate melt varies from 545 to 2800 ppm. The concentrations of carbon
71 in the quenched silicate melt are much higher than concentrations observed in low P–T
72 experiments which are typically 10–200 ppm^{1–3,8}. Thus, metal–silicate partition coefficients for
73 carbon are 24–166 (Table 1) and are 1–2 orders of magnitude lower than those determined in
74 previous low-pressure studies.

75 In order to understand the metallic phase better, we obtained images with SEM and TEM of the
76 sample BAS C 39 (Fig. 1), and we performed electron diffraction analyses that show the
77 presence of stoichiometric Fe_7C_3 inclusions in the quenched metal (Fig. 1b). To further
78 investigate the quench textures (Fig. 1c), we performed scanning TEM-energy dispersive
79 spectroscopy (EDS) mapping to identify Si-O-rich exsolved inclusions in the metal phase (Fig.
80 1d) along with a C-rich domain that could be an exsolved diamond, as also observed in previous
81 similar experiments²⁷.

82 To understand the effects of pressure and temperature on the metal–silicate partitioning of
83 carbon, we assembled a dataset for our model from previous studies to supplement our new
84 results (Fig. 2). However, only some published experiments are relevant to core formation
85 because prior studies were often focussed on the effects of silicate melt composition², $f\text{O}_2$ ^{28,29},
86 and the interaction of carbon with elements in the metal⁷. The partitioning of carbon is affected
87 by the ratio of non-bridging oxygens to tetrahedral cations² (NBO/T), so we only included data
88 from the literature for which NBO/T lies in a narrow range (0.5–1.5); the range for our samples
89 is 0.56–1.34. We also only included data from previous low P–T studies with low
90 concentrations (< 1 wt.%) of light elements (S, N, O, Si) in the metal because the interaction
91 between carbon and other light elements may be significant⁶ but is poorly constrained. As a
92 consequence, we also excluded very low $f\text{O}_2$ data ($\Delta\text{IW} < -3$), as they are often associated with
93 high concentrations of Si in the metal. On the other hand, our high P–T samples contain
94 significant concentrations of oxygen and silicon in the metallic phase, as is typical for high
95 temperature experiments^{24,25}. By combining these two datasets, we thereby empirically include
96 the interaction effects of these light elements in the regression presented below, because the
97 concentrations of O and Si in metal are generally low during core formation at low P–T
98 conditions^{11,31} consistent with the data^{1-3,8} selected from previous studies, whereas at high P–T
99 conditions such concentrations become high^{30,31}, which is consistent with the present study.

100 Interactions of carbon with oxygen and silicon in the metal of our experiments might be at least
 101 partly responsible for the scatter of our data (Fig. 2). Finally, we did not use literature data for
 102 which the silicate phase did not quench to a glass, because the quality of SIMS analyses may
 103 be compromised¹. In Fig. 2, we present a comparison of our data together with the selected data
 104 from the literature as a function of temperature and pressure.

105 The metal–silicate partition coefficient of carbon (D_C) is calculated as:

$$106 \quad D_C = \frac{X_C^{\text{metal}}}{X_{(\text{CO}_3)_2}^{\text{silicate}}}, \quad (2)$$

107 where X is the mole fraction of the element in the phase of interest. D is a function of several
 108 parameters, including pressure (P), temperature (T), and oxygen fugacity.

109 To remove the effect of oxygen fugacity on D values, we use the distribution coefficient K_D :

$$110 \quad K_D = \frac{D_C}{(D_{\text{Fe}})^{n/2}}, \quad (3)$$

111 where n is the valence of C when dissolved in silicate liquid, which is 4 for carbon in this case³².

112 Based on the aforementioned data selection, we fitted previously– published data^{1–3,8} along with
 113 our new results to derive the following expression:

114

$$115 \quad \log K_D = -1(\pm 0.5) + \frac{4842(\pm 920)}{T} + 31(\pm 19) \frac{P}{T}. \quad (4)$$

116

117 This relation shows that carbon becomes less siderophile with increasing depth along a magma
 118 ocean geotherm, and predicts K_D values fairly well for both low and high P–T experiments (see
 119 Extended Data Fig. S3).

120 We used the parameterization of Eq. 4 in a combined model of accretion and core–mantle
 121 differentiation¹⁰. This model is based on N-body accretion simulations that typically start with
 122 80–220 Mars-size embryos (initially located between 0.7 and 3.0 AU) that are embedded in a
 123 protoplanetary disk consisting of a few thousand much smaller planetesimals that are initially

124 distributed over a heliocentric distance of 0.7 to ~9.5 AU. In the accretion simulations, planets
125 typically grow from the starting embryos through accretional collisions with other embryos and
126 planetesimals. Each collision potentially involves a core formation event during which accreted
127 metal equilibrates with silicate liquid in a magma ocean and then segregates to the proto-core.
128 By modelling metal–silicate equilibration using a mass balance approach, the evolving
129 compositions of the mantles and cores of all accreting bodies are tracked. The approach requires
130 that the compositions of all starting bodies are defined, with refractory element compositions
131 matching CI chondrites. An oxidation gradient in the protoplanetary disk defines the oxygen
132 content of starting bodies, which is the main compositional variable. Embryos and
133 planetesimals that formed within ~1 AU of the Sun are highly reduced but with increasing
134 heliocentric distance are increasingly oxidized^{9,33}. Planetesimals originating from beyond ~4.5
135 AU are fully oxidized, contain no metal, and have an H₂O content of ~20 wt.%. An important
136 feature of the model is that accreted metal only equilibrates with a small fraction of the silicate
137 mantle – which is quantified using a hydrodynamic model of the interaction of metal and silicate
138 as material from the impactor’s core sinks in a magma ocean³⁴.

139 Here we use the “Grand Tack”^{35,36} accretion simulation “4:1-0.5-8”⁹ with the same model
140 parameters as in our previous study¹⁰. We first assume that carbon is accreted to Earth only in
141 bodies that originate beyond 4.5 AU and that their average C concentration is 1.6 wt.% (an
142 approximately mean carbonaceous chondrite (CI, CM, CO and CV) value). The mantle C
143 abundance increases throughout accretion from a very early stage and finally reaches the BSE
144 value of 80–120 ppm (Fig. 3). However, in this model the final concentration of C in the core
145 is 47 ppm which means that the carbon concentration of the bulk Earth is only 83 ppm, i.e. far
146 below the lowest estimate of 520 ppm²⁰. We have therefore developed a scenario that results in
147 a bulk concentration of ~520 ppm. We assume that bodies from beyond 4.5 AU contain on
148 average 1.6 wt.% carbon and that C concentrations in bodies originating at < 4.5 AU decrease

149 along a linear gradient to reach a value of zero at a heliocentric distance of 1.88 AU. This is
150 based on the condensation of volatile elements such C increasing with heliocentric distance as
151 temperatures decrease and is analogous to a concentration gradient derived for sulfur¹⁰. The
152 model then predicts a final mantle C concentration of 314 ppm (Fig. 3), a core concentration of
153 950 ppm and a bulk Earth concentration of 518 ppm. This calculation is not unique because the
154 concentration (1.6 wt.% C) and distance (1.88 AU) parameters are of course correlated.
155 Furthermore, if water-rich outer solar system planetesimals have concentrations of C exceeding
156 1.6 wt.%, then it is possible to obtain mantle concentrations consistent with the estimate of 765
157 ± 300 ppm²⁰ because C delivered to Earth's mantle is inefficiently extracted by subsequent core
158 formation similar to S and the HSEs¹⁰, as explained below. In summary, the final mantle
159 abundance of Earth exceeds 200 ppm for all plausible C concentrations of the initial bodies in
160 the protoplanetary disk.

161 A BSE carbon concentration of > 200 ppm is considerably higher than the estimate of 100 ± 20
162 ppm^{4,5}. If the 100 ± 20 ppm value is correct, an additional process, that is not included in our
163 model, may have transported excess carbon to the core. There is experimental evidence that the
164 pressure-induced disproportionation of $\text{Fe}^{2+} \rightarrow \text{Fe}^{3+} + \text{Fe}$ in deep magma oceans causes the
165 formation of dispersed Fe metal droplets^{37,38}. The gravitational segregation of these metal
166 droplets could have transported the excess carbon to the core, leaving ~ 100 ppm in the mantle
167 and resulting in 0.14 wt.% C in the core.

168 There are two critical features of the accretion/differentiation model that explain why the BSE
169 carbon concentration is not efficiently extracted during core formation in spite of its siderophile
170 behaviour. First, bodies originating from the outer solar system are water ice-bearing and so
171 fully oxidize and contain no metal. Thus, there is no core formation event when they are
172 accreted and the delivered C remains in the magma ocean/mantle. Second, when differentiated
173 metal-bearing bodies are accreted, the metal of the impactor only equilibrates with a fraction of

174 the target's mantle, of which the value depends on the size of the impactor's core and the depth
175 of the magma ocean^{9,34}. The fractions of the mantle that equilibrate with metal in the present
176 simulation are 0.16 to 2.5% for planetesimal impacts and 2.6 to 9.9% for embryo impacts.
177 Consequently, core formation events are extremely inefficient at removing carbon from the
178 bulk of the mantle and transferring it to the core. These two critical features are absent in the
179 models of "continuous core formation" that are currently applied in many studies of Earth's
180 differentiation^{12,13}; the result of a continuous core formation model for carbon is presented in
181 the supplementary information for comparison. In such models all accreted material is assumed
182 to contain metal and it is assumed (unrealistically) that all accreted metal equilibrates
183 chemically with the entire Earth's mantle. Consequently, such models predict that essentially
184 the entire BSE carbon budget has to be delivered during late accretion, i.e. after core formation
185 has ended (see supplementary Information and Extended Data Fig. S3). We know of no
186 astrophysical accretion simulations that would support this scenario which requires the
187 accretion of carbon-poor material during core formation and carbon-rich material afterwards
188 during late accretion.

189

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291 **Table**

292 *Table 1: Experimental conditions of carbon partitioning experiments.*

Run #	BAS C 50	BAS C 60	BAS C 39	BAS C 42
P (GPa)	60	71	49	56
T (K)	3800	4000	3600	3700
ΔIW	-1.06	-1.2	-0.89	-1.51
D_c	149 \pm 7	69 \pm 14	24 \pm 1	166 \pm 6
K_D	12.9 \pm 1	5 \pm 1.4	3.0 \pm 1.0	2.9 \pm 1.1
nbo/t	0.95	0.58	1.34	0.56

293 ΔIW is the calculated fO_2 relative to the IW buffer (see text for details).

294 D_c stands for the metal–silicate partitioning of carbon (mol.% of carbon in the metal / mol.%
 295 carbon in the silicate).

296 K_D is the distribution coefficient where $K_D = \frac{D_c}{(D_{Fe})^{n/2}}$, nbo/t is the number of non-bridging

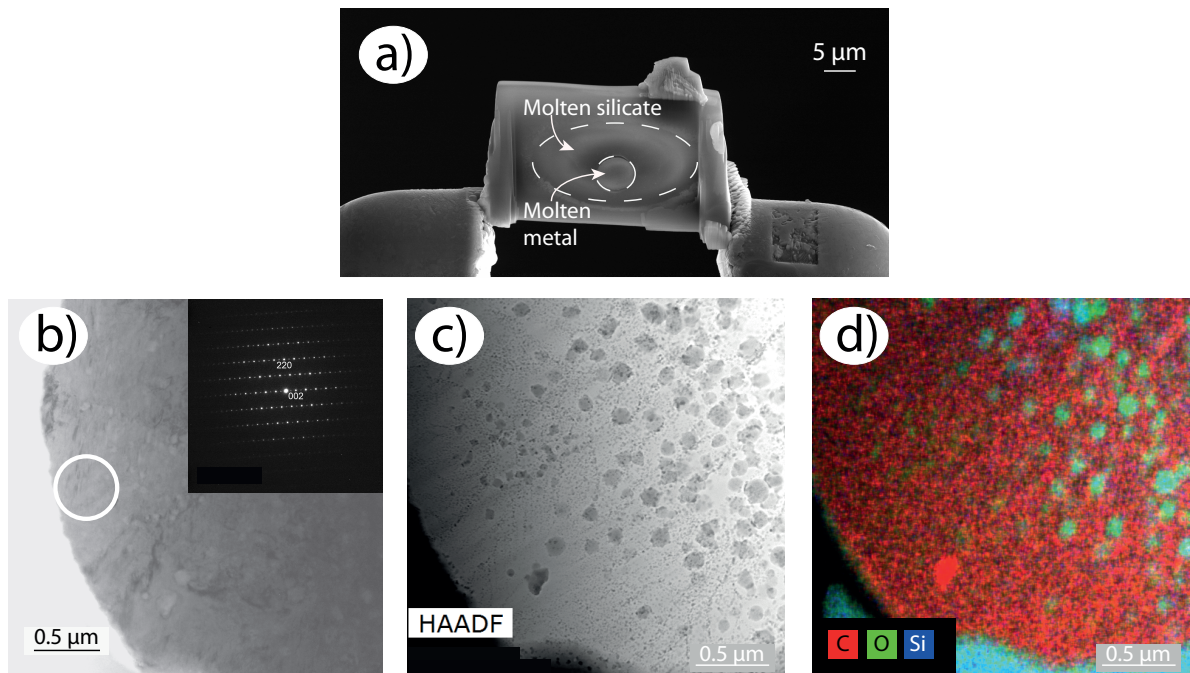
297 oxygen atoms over tetrahedrally coordinated cations.

298 The uncertainties on temperature and pressure measurements are estimated to be about 300 K
 299 and 5 GPa respectively.

300

301 **Figures**

302



303

304 *Figure 1: Images of run BAS C 39 synthesized at 49 GPa and 3600 K. a) Secondary electron*

305 *image where two different phases can be observed in the sample: the molten silicate and the*

306 *molten metal that were equilibrated at high pressure and high temperature. b) Bright field TEM*

307 *image and selected area electron diffraction pattern (Inset) of the metallic phase showing the*

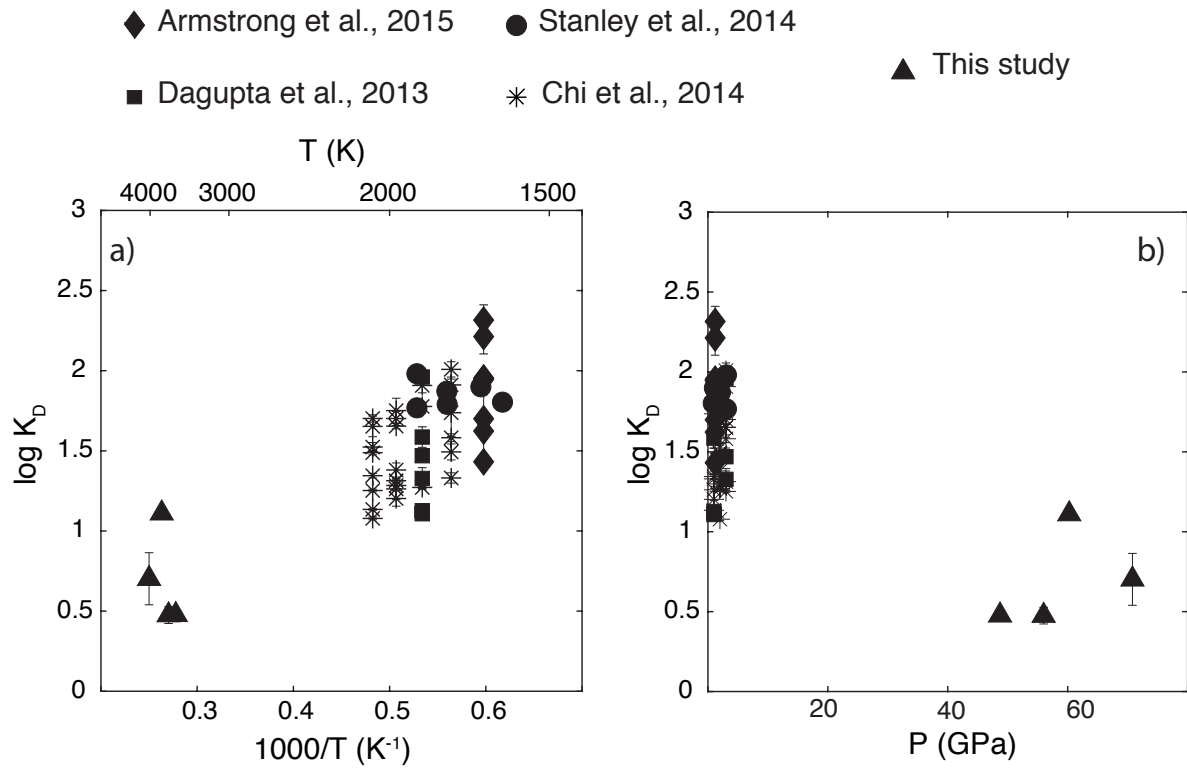
308 *presence of stoichiometric Fe_7C_3 (also confirmed with EELS). c) High-angle annular dark field*

309 *of the same sample showing exsolutions that developed during quenching along with d)*

310 *Scanning TEM-EDS map of the same area which shows exsolved inclusions that are oxygen*

311 *rich (green) and silicon-rich (blue). The bright red inclusion in image c) is believed to be an*

312 *exsolved diamond²⁷.*



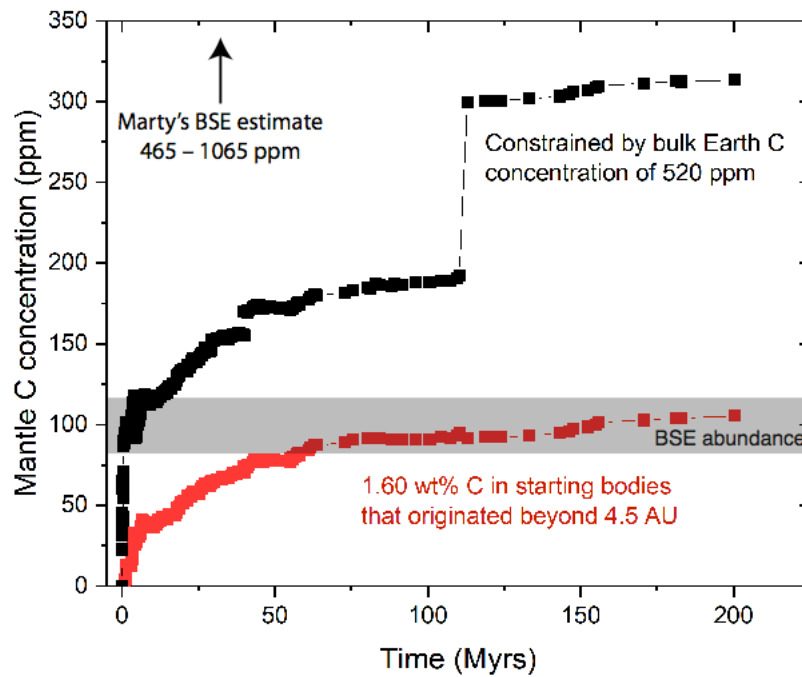
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314 *Figure 2: Comparison of our data with results of previous studies with respect to*

315 *temperature (a) and pressure (b). Uncertainties on temperature and pressure measurements*

316 *for diamond anvil cell experiments are of 300 K and 5 GPa respectively.*

317



318

319 *Figure 3: Mantle C concentration during accretion as a function of time based on a*
 320 *combined accretion/differentiation model^{9,10}. The red symbols show the evolution of mantle*
 321 *carbon concentration assuming that all starting planetesimals that originated from beyond*
 322 *4.5 AU contained an average of 1.6 wt.% C; bodies originating at heliocentric distances <*
 323 *4.5 AU are assumed to contain zero C. The black symbols show the evolution when the*
 324 *average carbon concentration in starting bodies originating from beyond 4.5 AU is 1.6 wt.%*
 325 *and C concentrations in bodies originating at < 4.5 AU decrease along a linear gradient to*
 326 *reach a value of zero at a heliocentric distance of 1.88 AU. The prominent jump at 110 Myrs*
 327 *is caused by the final giant impact. Estimates of the BSE concentration range from 100±20*
 328 *ppm⁴, as shown, to 765 ± 300 ppm²⁰.*

329 **Methods**

330 **a) Starting materials**

331 For the silicate starting glass, we mixed oxide and carbonate compounds in stoichiometric
332 proportions, decarbonated over night at 900°C, then added FeO. Powders were ground under
333 acetone for some time to ensure a homogeneous composition. We created pellets that were
334 subsequently fused in argon flux for several seconds at about 1400°C using a levitation furnace
335 device in Orléans, France³⁹. Chemical composition of the glass was checked using EPMA
336 (“Basalt” in table S1) and we could verify that no compositional variation was detectable. The
337 glass sphere was subsequently polished down to a thickness of 20 µm and machined in IPGP,
338 France, to obtain small disks of 80 microns diameter²⁵.

339 The goal of this study is to understand the effect of core formation on the partitioning of carbon
340 between Earth’s most important reservoirs: the silicate mantle and the metallic core. Carbon is
341 a prominent contaminant in the laboratory, with multiple sources of contamination, from the
342 use of ethanol and acetone, various type of glue or ambient contamination. It has also been
343 suggested that diamonds used as pressure transmitting tool can also diffuse carbon to the
344 samples⁴⁰. To better understand the origin of carbon, we used ¹³C as the source of carbon in our
345 experiments. We synthesized the metallic phase using piston-cylinder apparatus. In order to
346 assess the extent of carbon contamination during LH-DAC experiments, we only used ¹³C. We
347 mixed 5 wt.% of ¹³C powder (97% pure from Cambridge Isotope Laboratories, Inc.) with 95
348 wt.% of Fe. We melted this alloy at 2 GPa and 1873 K in MgO capsule for 10 minutes using a
349 ½” piston-cylinder assembly at BGI. This ¹³C-doped carbide was then crushed to be used for
350 the metallic part of the for laser-heated diamond anvil cell experiments. The composition was
351 checked by EPMA (see table S1, “Metal”) using an Fe₃C standard (see sections d) and e)).

352

353 **b) Laser-heated diamond anvil cell experiments (LH-DAC)**

354 We used diamonds with 250 μm culets and rhenium gaskets. Gaskets were pre-indented to
355 obtain a thickness of about 40 to 50 μm , and subsequently laser-drilled to have an experimental
356 chamber of about 90 μm in diameter. We then loaded two silicates discs encapsulating a flake
357 of carbide²⁵. Samples were compressed to the target pressure, and then laser-heated at the
358 desired temperature using a doubled sided laser heating system at BGI. Temperature was
359 generated by two fiber continuous-wave (CW) YAG lasers (SPi©) with wavelength of 1064
360 nm and delivering 100 watts each. The laser beam was focused onto the sample using two NIR-
361 Mitutoyo objectives lenses with x20 magnification. Temperature was measured simultaneously
362 and continuously on both sides of the diamond cell using a spectro-radiometric technique⁴¹
363 using a 2500i spectrometer and pixis 400 CCD camera from Princeton-Instrument© from the
364 light collected through both objective lenses. The peak temperature, above the liquidus
365 temperature of the material, was maintained for few tens of seconds, before switching off the
366 electric power of the lasers to quench the experiment. The samples quenched extremely quickly
367 due to the high thermal conductivity of diamonds and the cells were subsequently decompressed
368 slowly over several hours. As observed in previous studies using LH-DAC, the silicate part
369 quenched to a glass, and the Fe-rich metal presents Si- and O-rich inclusions^{24,25}.

370

371 **c) Sample recovery**

372 We used the Focused Ion Beam facility (FIB, dual beam Scios FEI©) in Bayerisches Geoinstitut
373 (BGI) to recover small lamellae of 3–5 μm thickness, about 40 μm in length and 20 μm high^{24,25}.

374 In the samples, we observed two quenched phases that had been molten at high pressure and
375 high temperature: a silicate and metallic one (see Fig. 1-a in the main text).

376

377 **d) Fabrication of analytical standards**

378 For EPMA analysis of the carbon contents of metals, we synthesized a Fe₃C (cementite) carbide
379 standard at BGI⁴², by inserting a 1 mm diameter, 10 mm long, 99.99 % purity Fe wire in a thick
380 graphite sleeve and heating it to 1423 K at 15 kbar in a piston cylinder for one week. The
381 reaction product was confirmed to have the Fe₃C cementite structure by XRD and was assumed
382 to be stoichiometric.

383 For the nanoSIMS analysis of quenched glasses, we synthesized ¹³C-doped glass standards⁴³.
384 We first mixed oxides in basaltic proportions and made a glass at 1600°C in a furnace for two
385 hours. This glass was analyzed by EPMA for major elements, and by FTIR to confirm that it
386 was carbon-free (“B1” in table S1). Subsequently, we performed piston-cylinder experiments
387 to incorporate ¹³C into the basaltic glass. The source of ¹³C was chosen as oxalic acid enriched
388 in ¹³C (Cambridge Isotope Laboratories, Inc.), that was loaded along with the basaltic glass in
389 a Pt capsule and pressurized to 2 GPa and heated to 1600°C for 10 minutes⁴³. We created two
390 glass standards (B1145 and B1147) containing 785 ppm and 1263 ppm of carbon respectively,
391 as measured by FTIR at BGI. For the FTIR measurements, we used an extinction coefficient⁴⁴
392 of 69500 L.mol⁻¹.cm⁻¹. We report the FTIR measurements along with EPMA analyses on table
393 S1 for the carbon-free (B1) and the two carbon-doped glasses.

394

395 **e) EPMA analysis**

396 Major elements from both the metallic and silicate phases were analyzed by Electron Probe
397 Microanalysis (EPMA). We calibrated using wollastonite, spinel and olivine for Ca, Al and Si,
398 Mg, Fe in the silicate respectively with current of 15 keV and 15 nA, a focused beam and a phi-
399 rho-Z correction. For the metallic phase, we used a pure Fe wire, FeSi, Fe₂O₃ and Fe₃C
400 standards to calibrate for Fe, Si, O and C respectively, at 15 keV with a 25 nA beam current,
401 again using a focused beam and phi-rho-Z correction. The counting time was 10 seconds on the

402 background, and 20 seconds on the peak for all elements, and the samples and Fe₃C standard
403 were not carbon coated.

404 Analyzing carbon using the microprobe requires special consideration⁴⁵. In order to quantify
405 the carbon content of both the experimental metals and the metal starting material, we
406 synthesized a Fe₃C carbide primary standard (see section d)). During C analysis, we took
407 particular care in repeating the measurements on the metallic standards to quantify background
408 levels (contamination) inside the EPMA. We measured the count rate from the C α peak on a
409 pure Fe standard continuously for ten minutes, and found no statistically-significant change in
410 peak height over that timescale⁴². We performed repeated measurements of pure iron wire to
411 check for carbon contamination, and observed a background for carbon of 0.49 ± 0.07 wt.%
412 which is well below the amount of carbon present in the metallic phase of our samples. The
413 same was done with pure FeSi, for which we measured 0.36 ± 0.05 wt.% of carbon in the
414 standard. As for oxygen contamination, we measured a background of 0.04 ± 0.04 wt.% in the
415 Fe wire and 0.15 ± 0.05 wt.% in FeSi. Those values are again well below the amount measured
416 in our samples. Standard analyses were consistent and stable over the whole duration of the
417 analytical sessions.

418 Analyses performed on the quenched silicate highlighted high concentrations of FeO (from 14
419 to 27 wt.%), as observed previously in similar experiments^{24,25}. In the metallic phase, we
420 observe a range of concentrations of carbon, from 5 to 9.5 wt.%, depending on the sample.
421 Oxygen and silicon are also present in the metallic phase, from 2.8 to 7.5 wt.% oxygen and 0.3
422 to 5.6 wt.% silicon. We report on tables S2 and S3 the full analyses of our recovered runs for
423 silicate and metal respectively. Whilst some analytical totals are low, these low totals are not
424 explained by the thin samples, as a 3 μ m deep lamella is thick enough to not lose electrons, nor
425 significant X-ray flux, through the base⁴⁶.

426

427 **f) nanoSIMS analysis**

428 The abundance of carbon in the silicate phase was expected to be extremely small, due to its
429 high siderophilicity¹. Laser heated diamond anvil cell experiments produce very small samples,
430 with silicates phases that are only few micrometers wide. Therefore, NanoSIMS was used in
431 this study to achieve required analytical precision and spatial resolution and quantitatively
432 analyze the low carbon contents in the silicate phase.

433 NanoSIMS measurements are sensitive to matrix effects, so we took particular care in
434 synthesizing and using relevant standard glasses (see section d)). Along with those two
435 standards, we used two standards of natural rocks containing natural and known amounts of
436 carbon: a basaltic glass from the East Pacific Rise^{47,48} ALV 981-R23 containing 405 ppm of
437 CO₂; and a piece of glass from D'Orbigny meteorite^{49,50} containing 40 ppm of C. Composition
438 of those standards and of the natural samples are given in table S1. During the NanoSIMS
439 session, we analyzed the standards each day to check for consistency and demonstrate the
440 reproductivity of the measurements.

441 NanoSIMS analyses were performed using the CAMECA NanoSIMS 50L at the Open
442 University (Milton Keynes, UK). Prior to NanoSIMS analysis, a large area of each sample was
443 pre-sputtered using a focused primary beam of 16 KeV Cs⁺ ions with probe current of 100 pA
444 to remove surface contamination. The sizes of the pre-sputtered areas vary from 5×5 to 7×7
445 μm². Analyses were then carried out in spot mode first by scanning a 50 pA Cs⁺ beam onto the
446 inner 3×3-5×5 μm² areas, with secondary ions of ¹²C⁻, ¹³C⁻, ³⁰Si⁻, ²⁴Mg¹⁶O⁻, ²⁷Al¹⁶O⁻ and
447 ⁵⁶Fe¹⁶O⁻ collected in electron multipliers simultaneously. Only the data from the inner 60% of
448 the area were collected using electron gating to avoid carbon contamination from the
449 surrounding areas. Each spot analysis consisted of 200 measurements, with a total analysis
450 time of ~2 minutes.

451 After spot analyses, selected sample areas were also measured in imaging mode. A 50 pA Cs⁺
452 beam was also used in imaging mode, with five secondary ions (¹²C⁻, ¹³C⁻, ³⁰Si⁻, ⁵⁶Fe¹³C⁻ and
453 ⁵⁶Fe¹⁶O⁻) monitored simultaneously. Each imaging analysis consisted of 10 frames. A frame
454 size of 256×256 pixels was used for all images with an integration time of 500 ms per pixel,
455 leading to total analysis time of 5-6 minutes for each image.

456 The mass resolving power (m/Δm) for both spot and imaging mode was set to 9000 (CAMECA
457 definition), sufficient to resolve all interferences from neighbouring mass peaks, such as ¹³C
458 from ¹²CH. We could extract from the region of interest (ROIs) the ratios of ¹³C/³⁰Si and
459 ¹²C/³⁰Si. We derived a precise calibration line using those standards, which was used to infer
460 the amount of carbon (both ¹²C and ¹³C, see figures S1-a and S1-b) present in each of our
461 samples. During analyses, we carefully avoided metal blebs that are sometimes present in the
462 silicate phase in this style of experiments, interpreting them to be entrained rather than exsolved
463 on quench. We present in Fig. S2 NanoSIMS ion images of the metallic phase of sample BAS-
464 C-39, to highlight the level of detail that can be reached with NanoSIMS.

465 Some uncertainty of the NanoSIMS measurements may derive from the difference between the
466 composition of the samples and the one of the standards. We calculated the amount of carbon
467 in each of our samples using the C/Si ratio of the standards and the experimental runs for both
468 ¹²C and ¹³C. The standards we used contained between 40 and 50 wt.% of SiO₂ whereas the
469 runs contained between 35 and 40 wt.% SiO₂. To encompass the range of potential sources of
470 uncertainty, we have assigned an uncertainty of ± 15% to the nanoSIMS measurements.

471 Analysing our samples and standards with the NanoSIMS revealed that the carbon present in
472 both cases was not only ¹³C but also ¹²C, with a lower ¹³C/¹²C ratio in experimental silicates
473 than in the metal (values from 2.8 to 6.9 for metal versus 0.37 to 1.7 for silicate, see Fig. S1-c).
474 The presence of significant concentrations of ¹²C was unexpected, since we carefully prepared
475 the samples using only ¹³C, but demonstrates that carbon contamination is probably ubiquitous

476 in high-temperature DAC experiments. The silicate starting powders were heated overnight at
477 900 °C and then converted to glass using an aerodynamic levitation system at temperatures of
478 about 1400 °C. Thus, despite several hours at elevated temperatures, it is extremely hard to
479 eliminate the presence of environmental carbon. ¹²C could originate from the diamond anvils,
480 the carbonate powders being not fully decarbonated, by sorption of carbon from the atmosphere
481 or from residues of organic carbon contamination during sample preparation.

482 The ¹³C/¹²C measured in the experimental silicates was lower than in the corresponding metals
483 (Fig. S1-c). This indicates that experiments did not reach isotopic equilibrium, and implies that
484 the most significant source of carbon contamination was the silicate starting material rather than
485 the diamond anvils. This observation is contrary to what has been proposed in previous LH-
486 DAC studies^{26,40}. We also measured silicate in unmelted regions in DAC samples BAS C 39
487 and BAS C 42 at the NanoSIMS, that contained 386 and 347 ppm of carbon (¹²C + ¹³C)
488 respectively, highlighting the presence of carbon in the starting material. We demonstrate here,
489 for the first time with support of NanoSIMS measurements, that carbon contamination in our
490 DAC experiments is not due to carbon diffusion from the diamond anvils to the samples, but is
491 from the samples themselves.

492

493 **h) TEM analysis**

494 We performed transmission electron microscope (TEM) analyses on a FEI Titan G2 80-200
495 S/TEM equipped with X-ray energy dispersive spectrometer (EDS) and electron energy-loss
496 spectrometer (EELS) in Bayerisches Geoinstitut in order to observe and further analyze the
497 experimental metals. To do so, we further thinned down one of our samples to a thickness of
498 60-80 nm using FIB. The FIB lamella was plasma cleaned prior to TEM analysis in order to
499 remove surface contamination. We determined Fe/C ratios in our sample and also perform EDS
500 mapping in the metallic phase. The quantification of the EELS analyses followed the procedure

501 previously described^{40,51}, using experimentally-determined ratios of partial cross sections of C
502 K and O K edges against the Fe L edge versus sample thickness, which were calibrated with
503 synthetic Fe_{0.94}O and Fe₃C samples. Note that due to the inhomogeneity of the metallic phase
504 of the sample (see Fig. 1-a,-b,-c), EELS technic cannot provide a measurement of the bulk
505 carbon contents of this phase.

506

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544 **Extended Data Tables and Figures**

545

546 *Table S1: Composition of the synthetic basaltic starting silicate used for the LH-DAC*
 547 *experiments (“Basalt”), the standards for the NanoSIMS synthesized with piston-cylinder at 2*
 548 *GPa and 1600 °C (“B1145” and “B1147”), the carbon-free starting material for piston-*
 549 *cylinder experiments (“B1”) and the starting metal used in LH-DAC experiments. We also*
 550 *report the composition of the two natural samples that were used at the NanoSIMS as*
 551 *standards. b.d. stands for below detection limit.*

Samples	Basalt	B1145	B1147	B1	D ^o Orbigny glass ^a	ALV 981-R23 ^c	Metal
wt.%	N=10	N=7	N=7	N=7	-	-	N=14
SiO ₂	52.18 (0.15)	48.05 (0.18)	50.89 (0.27)	53.56 (0.13)	41.7	49.63	C 5.67 (1)
Na ₂ O	1.90 (0.03)	1.80 (0.06)	1.86 (0.06)	2.03 (0.03)	-	2.91	Fe 91.61 (1.31)
CaO	10.92 (0.07)	9.47 (0.10)	10.03 (0.08)	10.70 (0.05)	23.3	11.61	
FeO	11.35 (0.04)	13.65 (0.20)	10.71 (0.06)	10.04 (0.06)	11.0	8.16	
Al ₂ O ₃	14.66 (0.07)	16.37 (0.11)	15.22 (0.10)	15.48 (0.05)	19.6	16.60	
MgO	7.45 (0.07)	7.22 (0.08)	7.57 (0.10)	7.73 (0.06)	1.87	8.36	
C (ppm)	-	785 (8)*	1263 (13)*	b.d.*	40	111	
Total	98.46 (0.20)	96.66 (0.19)	96.41 (0.33)	99.55 (0.15)	99.3 ^b	98.96 ^d	97.28 (0.76)

552 *Measured using FTIR (see text for details)

553 ^aComposition from ref.⁵⁰

554 ^bTotal includes also 2 wt.% of TiO₂, Cr₂O₃, MnO

555 ^cComposition from refs.^{47,48}

556 ^dTotal includes TiO₂, MnO, K₂O and P₂O₅

557

558 *Table S2: Compositions of quenched silicate liquid. Major elements were analysed using*

559 *EPMA and carbon was analysed using NanoSIMS. See text for details.*

Run #	BAS C 50	BAS C 60	BAS C 39	BAS C 42
wt.%	N=4	N=5	N=8	N=6
SiO ₂	35.84 (0.38)	39.97 (1.98)	34.09 (0.83)	39.66 (0.99)
MgO	8.85 (0.20)	8.60 (0.28)	9.53 (0.51)	9.02 (0.26)
Al ₂ O ₃	19.46 (0.12)	25.93 (0.50)	16.34 (1.36)	24.48 (0.64)
CaO	4.02 (0.16)	4.89 (0.12)	5.67 (0.20)	5.68 (0.31)
Na ₂ O	2.13 (0.07)	2.22 (0.10)	2.32 (0.15)	2.53 (0.11)
FeO	23.82 (1.40)	18.60 (0.60)	27.35 (2.55)	13.71 (1.72)
C (ppm)*	312 (47)	545 (82)	2805 (420)	272 (41)
Total	94.12 (1.85)	100.2 (1.96)	95.30 (2.34)	95.07 (1.55)

560 *15% uncertainties estimated on NanoSIMS measurements (see text for more details).

561

562 *Table S3: Composition of the metal phase of our recovered runs analysed using EPMA. See*

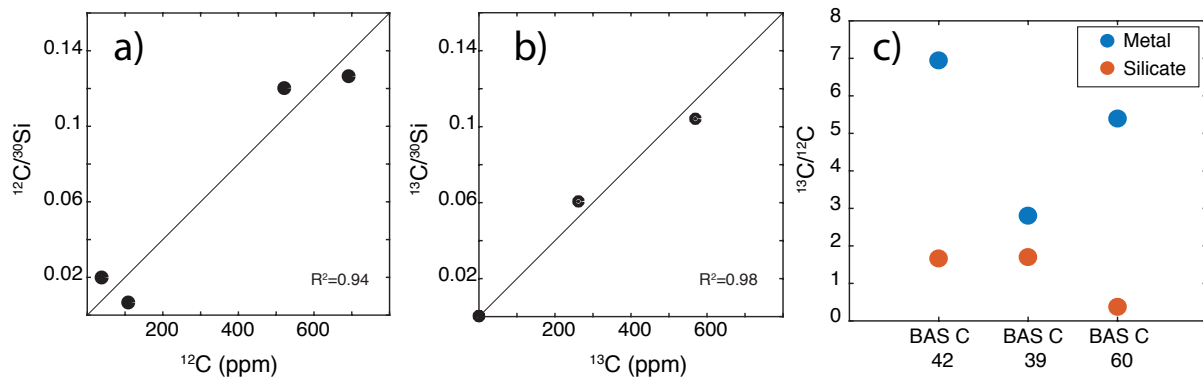
563 *text for details.*

Run #	BAS C 50	BAS C 60	BAS C 39	BAS C 42
EPMA (wt.%)	N=3	N=3	N=4	N=2
Fe	81.68 (1.25)	79.56 (0.9)	85.96 (0.58)	78.01 (0.78)
Si	0.46 (0.16)	5.56 (0.16)	0.33 (0.11)	2.65 (0.40)
O	3.79 (0.89)	7.56 (0.37)	2.85 (0.48)	3.42 (0.10)
C	5.90 (0.18)	4.99 (0.28)	9.44 (0.24)	5.35 (0.20)
Total	91.83 (0.41)	97.66 (0.38)	98.57 (0.29)	89.43 (0.48)

564

565

566



567

568 *Figure S1: Results from NanoSIMS measurements. a) and b) calibration lines obtained for*

569 *^{12}C and ^{13}C respectively using our standards and two natural samples (d'Orbigny and ALV*

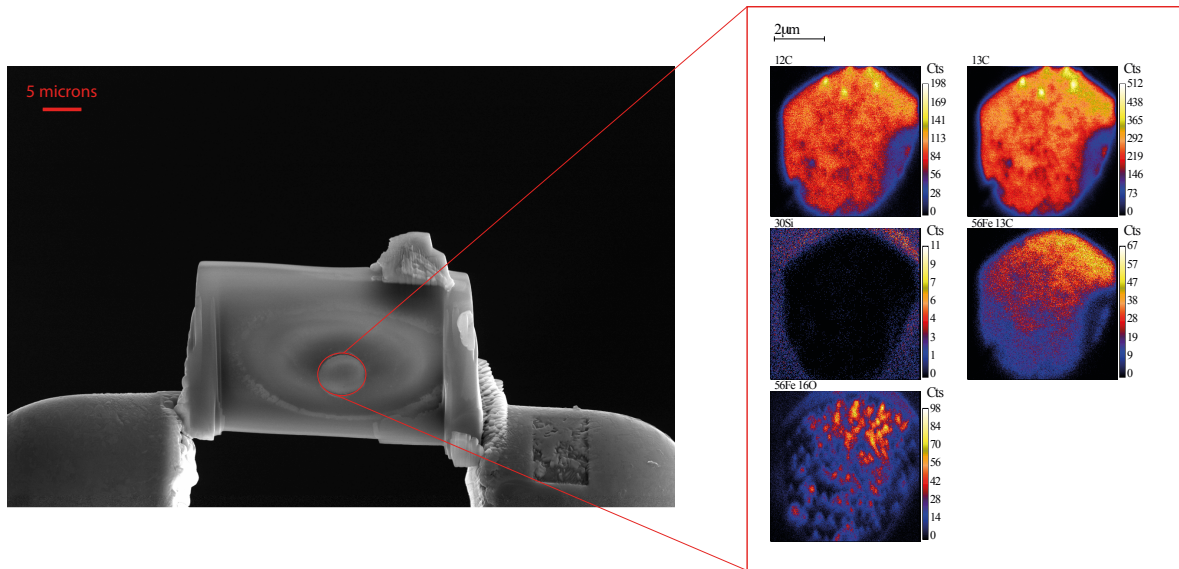
570 *981-R23) and both FTIR and NanoSIMS measurements. Note that for ^{13}C measurements, the*

571 *two natural samples plot at 0. c) Illustration of the isotopic disequilibrium between the*

572 *metallic and the silicate phase for three of our recovered runs which illustrates the presence*

573 *of ^{12}C in our reacted samples, despite the use of pure ^{13}C in the starting material.*

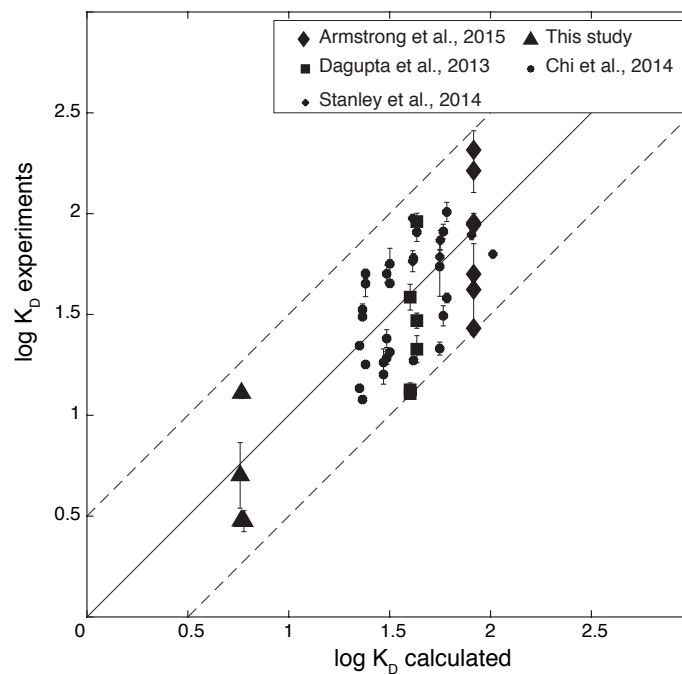
574



575

576 *Figure S2: Secondary electron image of sample BAS-C-39 along with NanoSIMS mapping of*
 577 *the metallic phase.*

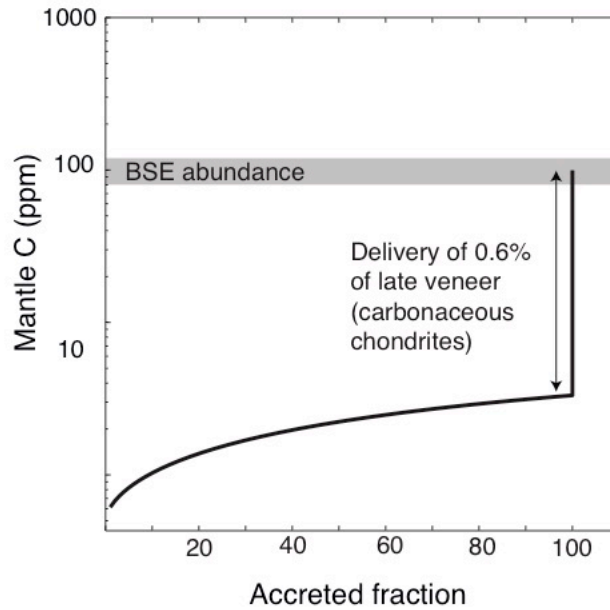
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579

580 *Figure S3: Validation of our partitioning model. Comparison of the K_D values (see eq. 3 in*
 581 *the main text) calculated from the experiments and from our model (eq. 4 in the main text).*

582



583

584 *Figure S4: Continuous core formation model. Evolution of the concentration of carbon in*
 585 *the mantle during Earth's core–mantle differentiation and the late accretion stage based on a*
 586 *continuous core formation model. The estimated BSE concentration of 80–120 ppm is shown*
 587 *in both cases⁴.*

588

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601

602 **Authors contributions**

603 I.B. conceived the project, performed experiments, analyzed data and derived the
604 thermodynamic model. E.S.J. analyzed data. I.A.F., X.Z. performed NanoSIMS measurements.
605 S.P. performed laser heating experiments and helped with FIB. N.M. performed TEM
606 measurements. D.C.R. and S.A.J. built the accretion model. I.B., E.S.J. and D.C.R wrote the
607 manuscript.

608

609 **Additional information**

610 **Supplementary Information** is available for this paper. Correspondence and requests for
611 materials should be addressed to Ingrid Blanchard (ingrid.blanchard@uni-bayreuth.de).

612

613 **Supplementary Information**

614 Using our high pressure and high temperature experiments along with the ones of the literature,
615 we built a new parameterization for the metal–silicate partitioning of carbon (Eq. 4 in the main
616 text). We compare in Fig. S3 the agreement between the experimental K_D and the one predicted
617 by our regression. Then, we utilized two styles of modelling in this study. The first one is the
618 one presented in the main text, and the second one is a more simple continuous core formation
619 model that has been applied in many geochemical studies^{1–3}. Here, Earth is assumed to accrete
620 material in small (1%) increments and each batch of accreted metal is assumed to equilibrate
621 chemically with the entire mantle. The equilibration pressure is assumed to correspond to
622 conditions the base of the magma ocean and increases as the planet grows. The temperature at
623 the base of the magma ocean, which is constrained by the peridotite solidus–liquidus, also

624 increases as the Earth grows. For simplicity, we set the fO_2 during accretion at IW-2.3 as given
625 by the current mantle/core FeO/Fe ratio⁴.

626 The concentration of carbon in the mantle today is estimated to lie between 80 and 120 ppm
627 (Dasgupta and Hirschman 2010; Hirschman 2016). On the other hand, the bulk Earth carbon
628 content has been proposed to lie between 520 ppm⁵ and 750 ppm⁶. Here, we assume that each
629 stepwise-accretion step delivers material with a bulk C content that matches the bulk Earth
630 value⁵. Using the higher value of 750 ppm⁶ would not significantly change the results. The
631 partition coefficient of carbon is calculated after each equilibration step following eq. 4, and
632 the resulting amount of carbon present in both mantle and core is derived from this calculation.
633 Hence, carbon is present throughout accretion.

634 The Earth is believed to have been bombarded by bodies comprised of chondritic material after
635 core formation ceased (the so-called late accretion stage which resulted in addition of the “late
636 veneer”) in order to explain the concentrations of highly siderophile elements present in the
637 mantle⁷⁻⁹. The amount of material that was delivered to Earth during late accretion was about
638 0.6% of Earth’s mass¹⁰, and was not incorporated into the core, but rather remained in the
639 mantle. Chondrites are carbon-rich, with up to 3.2 wt.% of carbon reported in CI chondrites¹¹,
640 so their delivery to Earth brought a substantial amount of carbon. We incorporated the effect of
641 such late accretion in our model by assuming that the carbon content of the late veneer matches
642 the mean carbonaceous chondrite (CI, CM, CO and CV) value of 1.6 wt.%.

643 Results of the continuous core formation model are presented in Fig. S4. The combined effects
644 of core formation and later addition of late veneer are needed in order to reproduce the carbon
645 content of the BSE. Core–mantle differentiation effectively strips the mantle of all carbon and
646 the addition of C during late accretion is required to achieve the mantle carbon concentration.
647 According to this model, the resulting concentration of carbon in the core after the combined
648 effect of differentiation and late veneer is 0.16 wt.%. This result is consistent with previous

649 suggestions, based on first principles molecular dynamics¹² and experiments¹³, that carbon is
650 not a dominant element in Earth's core. As discussed in the main text, this simple scenario is
651 not coherent with astrophysical accretion scenarios. This model assumes that the bulk Earth C-
652 content is known *a priori*, and the assumption that all accreted metal equilibrate with all the
653 mantle is also unrealistic.

654

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