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Ancient subducted oceans controlling the positioning of deep mantle plumes

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1 **Seismic imaging of the Earth's interior reveals plumes originating from relatively hot re-**
2 **gions of the lower mantle, surrounded by cooler material thought to be remnants of ancient**
3 **subducted oceans. Based largely on geophysical data, two opposing hypotheses dominate the**
4 **discussion on dynamics at the base of mantle: the large hot anomalies are thermo-chemical in**
5 **nature; or, alternatively, they are purely thermal plume clusters. In previous modelling stud-**
6 **ies, deep chemical heterogenities have been argued to be essential in developing appropriate**
7 **present-day plume positions. Here, we quantify how the chemical composition of large, hot**
8 **regions in the deep mantle influences the location of rising mantle plumes using numerical**
9 **3-D global mantle convection models constrained by 410 million years of palaeo-ocean evolu-**
10 **tion. For the first time, we show that purely thermal convection can reproduce the observed**

11 **positions of present-day hotspots. By demonstrating that a lower mantle without large-scale**
12 **chemical heterogeneities can generate appropriate global dynamics, we illustrate the power**
13 **of sinking ocean plates to stir mantle flow and control the thermal evolution of the mantle.**
14 **Because our models with a thermo-chemical anomaly reproduce the observed hotspot posi-**
15 **tions equally well, we posit that the deep hot anomalies in the mantle are purely passive in**
16 **global dynamics - regardless of their (thermal or chemical) origin.**

17 Large structures of low seismic velocities (the Large Low Shear Velocity Provinces, or
18 LLSVPs) observed in the lower mantle in seismic tomography studies and interpreted as regions
19 of elevated temperatures, appear to play a defining role in the origins and positioning of man-
20 tle plumes¹⁻⁴. Although the fundamentals of composition and origin of LLSVPs are still being
21 debated⁵⁻⁸, there is growing agreement⁹ that plumes may form in some way from these regions
22 (Figure 1). The formation of plumes can influence supercontinent dispersal, leading to a further
23 repositioning of subduction zones and a change in mantle flow¹⁰.

24 The nature of LLSVPs has been at the centre of plume generation discussions for over
25 a decade^{2,6,12}. Geochemical data from surface melts indicate the presence of chemical mantle
26 reservoirs¹³⁻²², with hotspot-derived melts associated with plumes consistently showing a different
27 mantle source composition than mid-ocean ridge basalts. Recent evidence suggests that spatial
28 geochemical patterns at oceanic islands reflect preferential sampling of a distinct source of deep
29 mantle material originating from the LLSVPs²³⁻²⁵. These observations have led to the idea that the
30 LLSVPs may be ‘thermo-chemical’ piles of material with a different composition than the average

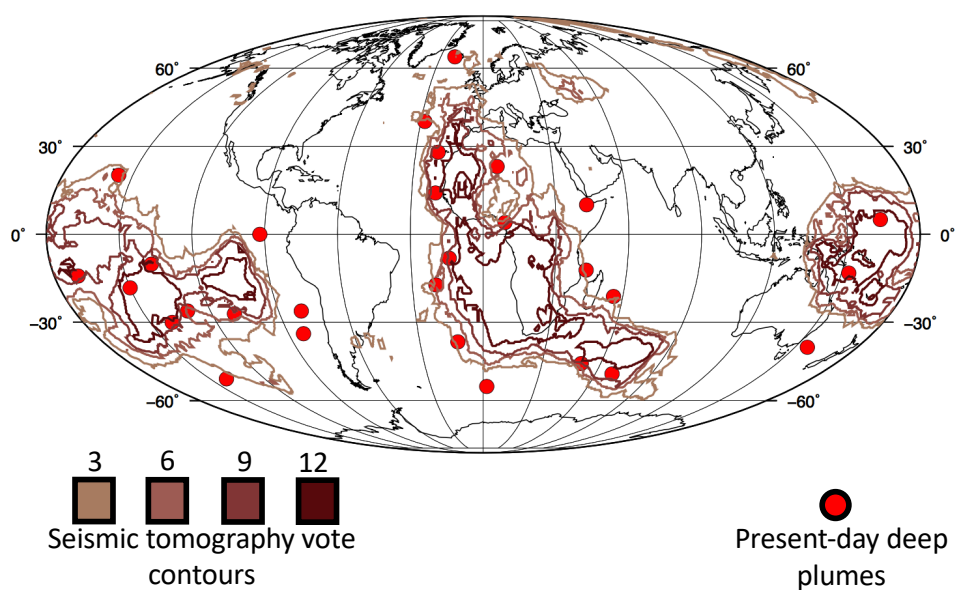


Figure 1: Present-day mantle features. Position of low velocity provinces (LLSVPs) at 2800 km depth alongside present-day deep mantle plume and hotspot locations^{1,4} (SI Table S.1). The low velocity regions are contours where 3, 6, 9, and 12 (out of 14) seismic tomography models agree that there are slow anomalies¹¹.

A divided view is presented by seismology. The two deep anomalies have been characterized by higher than average densities through early tomographic normal mode data and the anti-correlation of bulk sound and shear wave anomalies²⁶⁻³⁰. More recent analysis using tidal tomography has also inferred structures of increased density in the lower mantle³¹. However, work using Stoneley mode data offered an opposing view which indicated the potential for lighter material⁷. Compositional heterogeneities and/or a phase change to post-perovskite could both fit seismic observations in the lowermost mantle^{32,33} – though, the possibility that LLSVPs are purely thermal anomalies has also been put forward³⁴.

Previous geodynamic modelling studies have indicated that a thermo-chemical nature of the deep mantle material is essential in producing appropriate present-day plume positions³⁵⁻³⁷. However, the rheological nature of LLSVPs, and their role in mantle dynamics, is still being debated^{6,8,9,12,38,39}. Important open questions remain: Are LLSVPs thermal plume clusters or dense, stable thermo-chemical piles, or a mixture of the two?; Have LLSVPs been present and stable for most of the Earth's history?; and do they control the location of rising plumes? A key component to answering these questions lies in understanding the mechanisms that cause plumes to rise in the locations where hotspots are observed on Earth today.

To study mantle dynamics, we use 3-D thermo-chemical global mantle convection experiments that take into account palaeo-subduction history from 410 Ma to the present day⁴⁰ (which incorporates plate reconstructions models of the Paleozoic⁴¹ and Mesozoic-Cenozoic⁴²). We then

51 compare the results at the end of the model run to observations of LLSVP shape^{3,4,11} and plume
52 positions^{1,4}. By applying different subduction histories as a time-dependent boundary condition to
53 our models *e.g.*,^{6,35–37,43–48}, we gain insights on the role of subducted slabs in shaping the LLSVPs
54 and triggering the ascent of mantle plumes. In addition, we systematically vary the thickness of
55 compositional heterogeneity in the lowermost mantle, thereby controlling the dynamics of the hot
56 regions that represent the LLSVPs. This allows us to evaluate which of the proposed end-member
57 scenarios – an anomaly that is purely thermal or one that is thermo-chemical – is supported by the
58 observations.

59 A key innovation in our approach compared to previous studies is that we incorporate subduc-
60 tion history encompassing the formation and breakup of the supercontinent Pangea. The life-cycle
61 of a supercontinent has a dramatic impact on mantle dynamics⁴⁹, and we show that the subduction
62 history affects plume and LLSVP formation on different time scales. Here, our models discover the
63 mantle circulation dynamics that connect past plate motions and ocean island volcanism. We show
64 that the two opposing hypotheses – plume clusters and thermo-chemical piles – are both able to
65 reproduce the observed positions of present-day hotspots equally well, highlighting the importance
66 of ancient subducted oceans in stirring the convecting mantle.

67 **Results**

68 **Mantle dynamics.** We computed a suite of geodynamic models with varying amounts of chem-
69 ical heterogeneity in the mantle in order to investigate a wide range of lower-mantle convection

70 regimes: plume clusters, stable thermo-chemical piles, and metastable piles. Our models begin
71 from an initial condition at 410 Ma (Figure 2 and Table 1) that assumes a realistic average mantle
72 temperature and composition. The initial mantle temperature is laterally homogeneous, following
73 an adiabatic profile with thermal boundary layers at the top and bottom⁵⁰, and any anomalous,
74 dense material is initially emplaced as a uniform layer at the core-mantle boundary. We vary the
75 initial thickness of this layer between 0 km (purely thermal model) and 300 km for different model
76 runs (SI Table S.2; Model0, Model100, Model150, Model200, Model250, and Model300). In
77 Model200, the subduction pattern during the formation and dispersal of the supercontinent Pangea
78 shapes the hot thermo-chemical material into two distinct regions in roughly the same locations as
79 the observed LLSVPs (e.g., Figure 3). We herein refer to thermo-chemical piles from the numer-
80 ical models as ‘TCPs’ and reserve the ‘LLSVPs’ for the seismically imaged structures to which
81 TCPs are compared.

82 All simulations where the initial dense layer is at least 150 km thick produce two distinct
83 thermo-chemical piles (TCP, Figure 4). Model100 results in all the thermo-chemical material
84 being entrained into the mantle during the simulation (SI Figure S.2), and Model0 features thermal
85 plume clusters. Correlating the modelled African TCP with the shape of the African LLSVP (as
86 given by the agreement of 8 seismic tomography studies), we find that Model200 produces the
87 best fit (SI Table S.2). The Pacific thermo-chemical pile also forms in the correct hemisphere in
88 Model200, but the agreement with seismic tomography¹¹ is generally not as good as in the African
89 hemisphere. The modelled Pacific pile extends further north and south than the slow anomalies in
90 tomography studies, and it is missing the western limb when compared with the LLSVP outline

91 (Figure 3h).

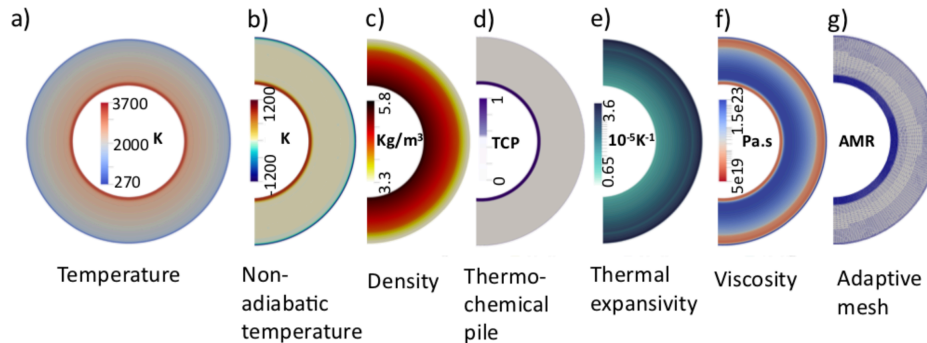


Figure 2: Initial conditions of the models (at 0 Myr/410 Ma), showing cross-sectional views of a) temperature, b) non-adiabatic (excess) temperature, c) density, d) initial thermo-chemical pile thickness (200 km in Model200 shown here), e) variable thermal expansivity, f) viscosity, g) the initial adaptive mesh.

92 Figure 5a shows the mantle under the northern Atlantic and Europe, indicating the excess
 93 temperature anomalies for upwelling plumes (red) and downwelling material (blue) for Model200.
 94 Comparing this model output at the equator to seismic tomography studies³ shows a number of
 95 similarities. Under Africa, our model TCP matches the shape of the LLSVP outline well, high-
 96 lighting the steep western side⁵⁵ and angled eastern flank (A and B, Figure 5b). In the Pacific
 97 hemisphere, our Model200 produces plume positions in the Galapagos (C, Figure 5b) as well as a
 98 TCP in the eastern part of the Pacific LLSVP (D, Figure 5b). However, our model misses the west-
 99 ern flank of the Pacific LLSVP across the equator (E, Figure 5b) as shown in Figure 3h. Further
 100 direct comparisons with another seismic study⁴ are given in SI Figure S.1.

101 **Present-day plume analysis.** Direct comparison between tomography slices (Figure 5b-c,

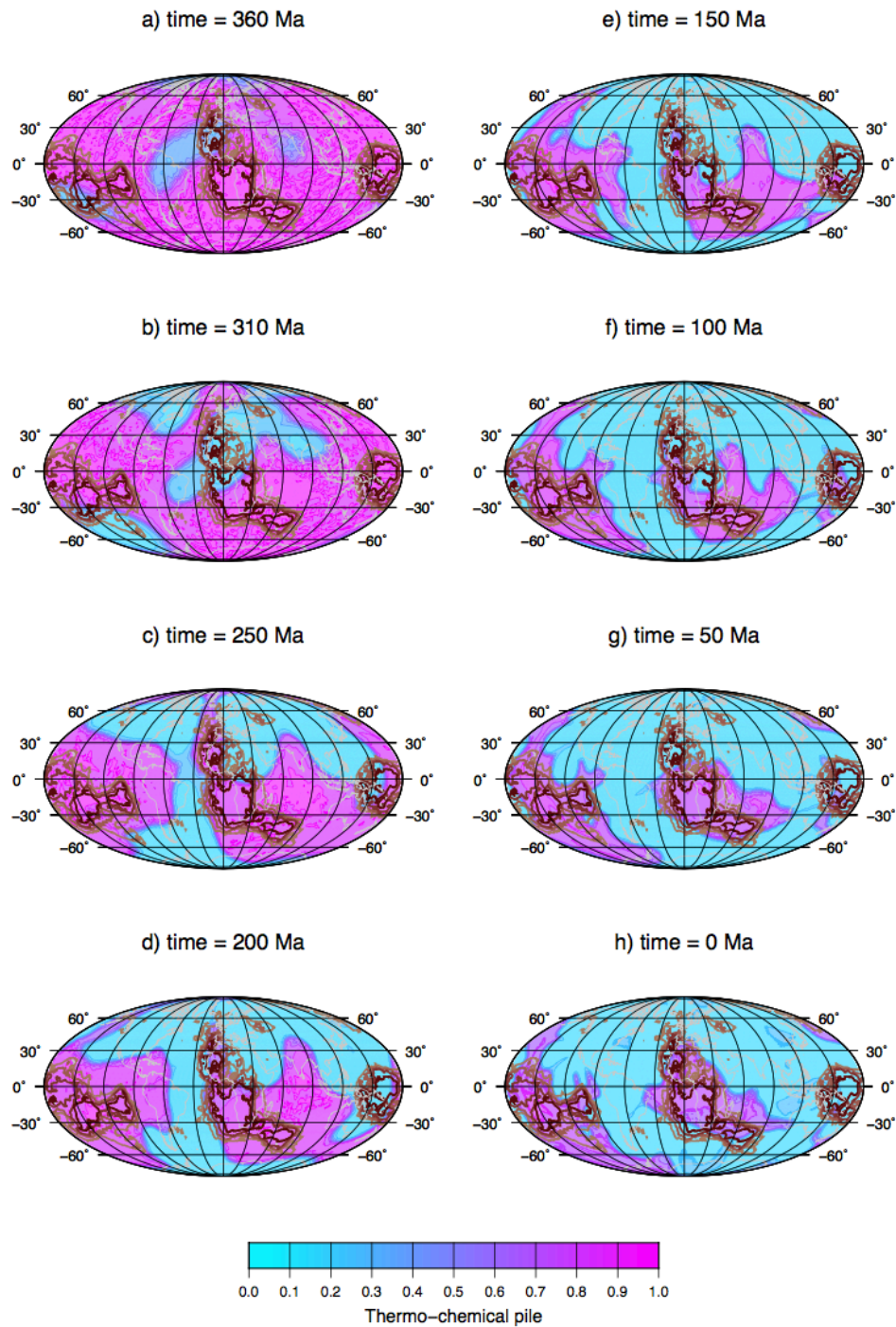


Figure 3: Evolution of chemical heterogeneities in Model200 from 410 Ma to present day.

The snapshots (a)-(h) show the proportion of anomalous, dense material in a spherical slice at 2800 km depth, illustrating how the material evolves from a layer with uniform thickness into a thermo-chemical pile. Brown contours indicate low seismic velocity anomalies as in Figure 1.

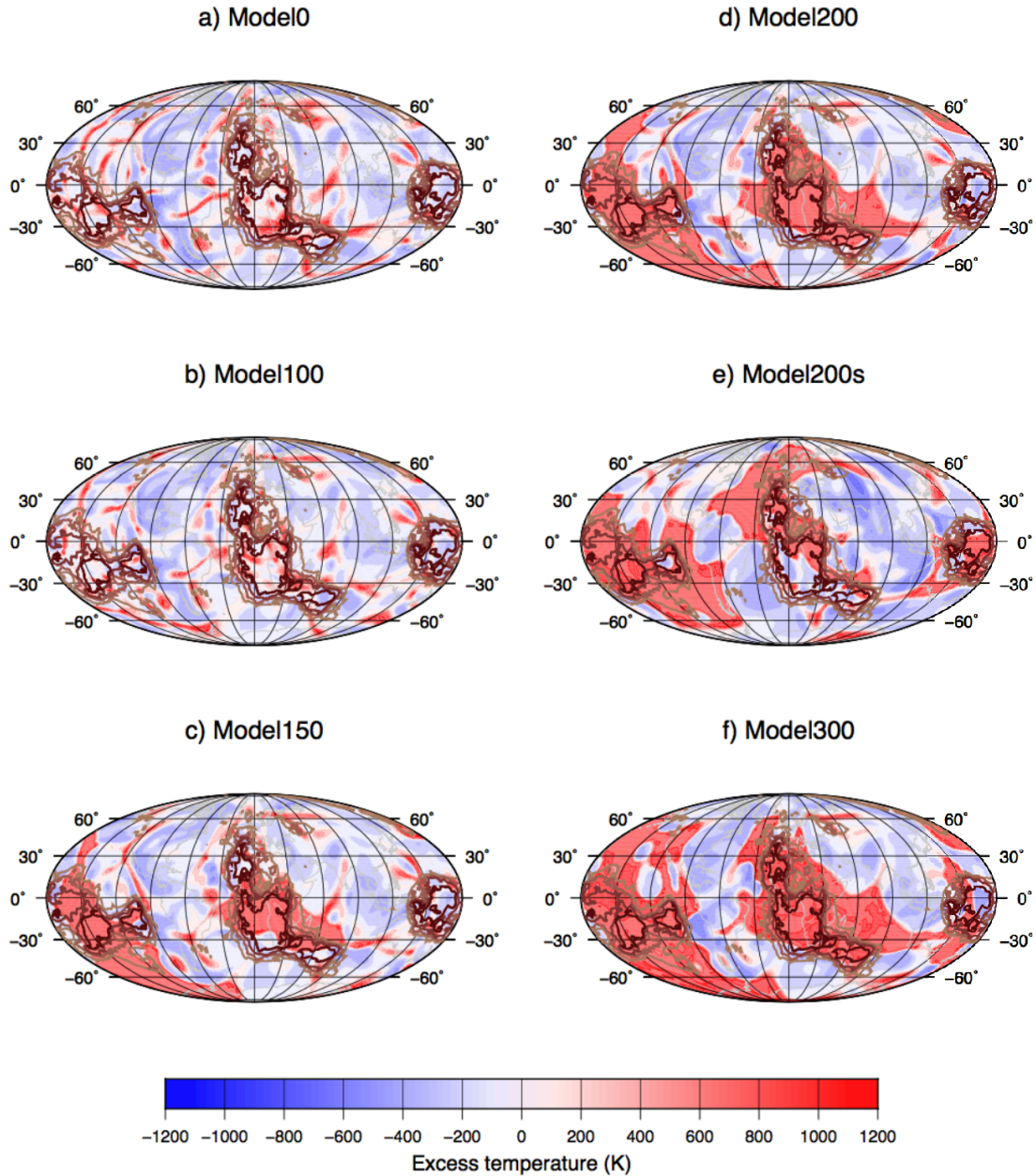


Figure 4: Differences in core-mantle boundary temperature and dynamics between models with varying amount of chemical heterogeneity. Present-day (0 Ma) excess temperature in 2800 km depth for a) Model0, b) Model100, c) Model150, d) Model200, e) Model200s, and f) Model300.

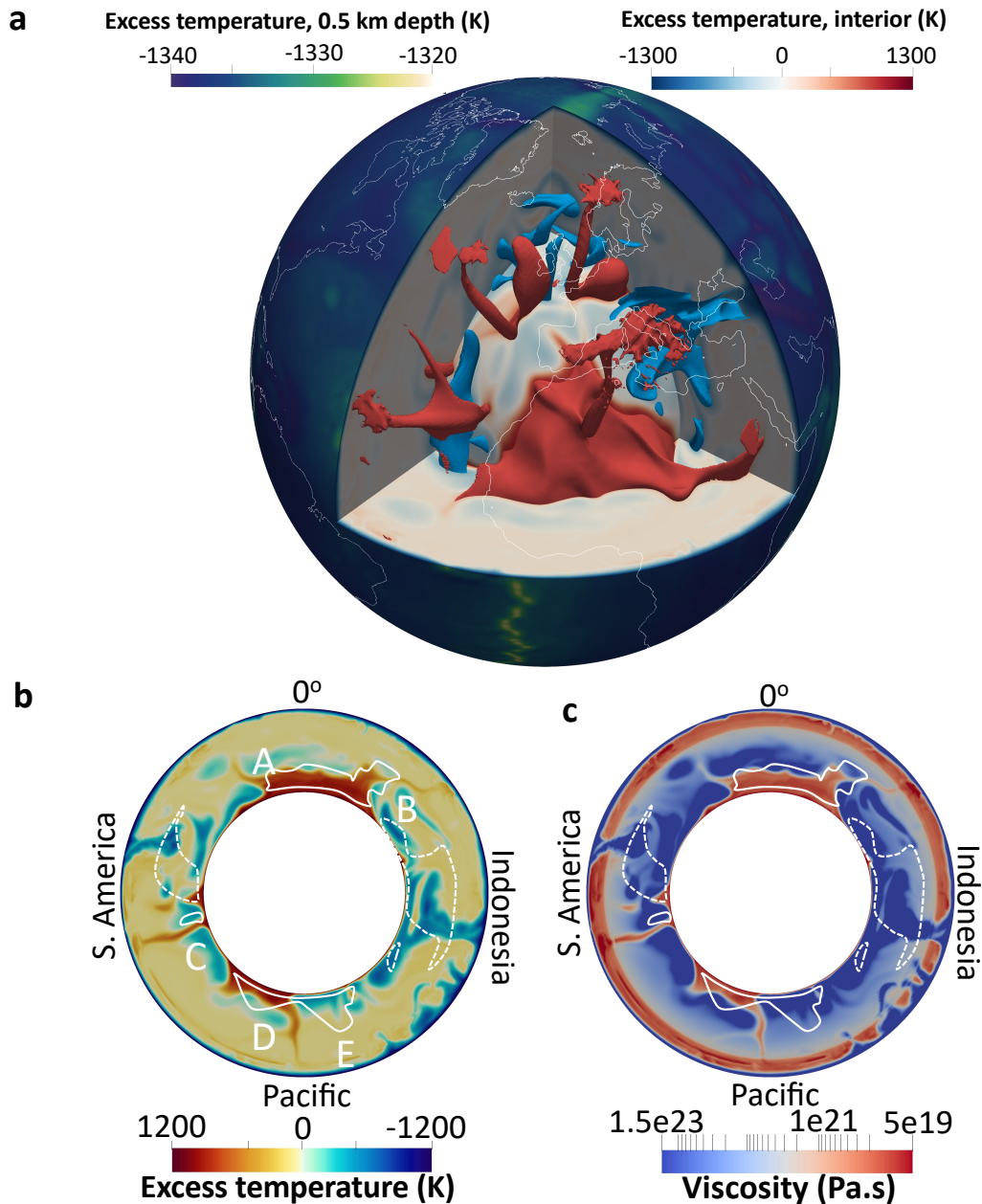


Figure 5: Global mantle dynamics and temperature compared with seismic models. a) A view into the mantle beneath the northern Atlantic and Europe for Model200. Excess mantle temperature anomaly contoured surfaces are given for warm (300 K, red) and cold (-500 K, blue) regions. b) Excess temperature slice at the equator, and the corresponding viscosity distribution (c), with the outline of where four seismic tomography studies agree that there is a shear velocity anomaly³ shown by white solid (slow) and dashed (fast) lines. Key: A: steep western side of the African LLSVP; B: angled eastern flank of the African LLSVP; C: Galapagos plume position; D:

Table 1: Material parameters.

Symbol	Variable	Value
η	Dynamic viscosity	literature ⁵¹
ρ	Density	PerpleX ⁵²⁻⁵⁴
g	Gravitational acceleration	9.81 m.s ⁻¹
C_p	Specific heat capacity	PerpleX ⁵²⁻⁵⁴
k	Thermal conductivity	4.7 W.m ⁻¹ .K ⁻¹
H	Radiogenic heat production	6×10^{-12} W.kg ⁻¹
α	Thermal expansivity	PerpleX ⁵²⁻⁵⁴

102 SI Figure S.1) can give a first order interpretation of how well the models reproduce present-day
 103 mantle dynamics. However, a better representation of model accuracy is to compare the difference
 104 between a present-day deep mantle plume and hotspot database^{1,4} (SI Table S.1) and the position
 105 of hot thermal anomalies in our models.

106 Figure 6a illustrates the position of upwellings of Model200 at 700 km depth (green) along-
 107 side the present-day plume and hotspot database^{1,4} (red). The plume positions in our model do
 108 not appear to be randomly distributed, with upwellings forming close to the upwards projected
 109 margins of the LLSVPs in the African domain. In the Pacific, a number of Model200 plumes form
 110 close to edges of the LLSVPs. Interestingly, this occurs in the western limb of the Pacific LLSVP
 111 even though there is no modelled thermo-chemical pile in that position (Figure 3h). A number
 112 of Model200 plumes form within 1000 km of the Azores, Afar, Comores, St Helena, San Felix,

113 Galapagos, Easter, and Pitcairn plumes (markers G to N in Fig 6a, respectively). Furthermore,
114 a deep mantle plume is positioned at the Louisiade site (marker O, Figure 6a) where a previous
115 study⁴ had imaged a plume despite no corresponding surface hotspot.

116 Figure 6a also shows the positioning of plumes generated from a model that does not feature
117 any thermo-chemical piles (Model0). The blue circles show plume positions at 700 km depth
118 for a purely thermal mantle convection model without chemical heterogeneities. There is a strong
119 similarity between the blue and green plume positions despite different lowermost mantle structure.
120 Indeed, more than half of the Model0 plumes form within a radius of 900 km from the Model200
121 plumes (SI Figure S.4), indicating that the thermo-chemical pile is not the driving force in plume
122 positioning.

123 To test how statistically relevant our results are, we look at the distance each plume in our
124 model forms away (at 700 km depth) from the nearest present-day hotspot in the database^{1,4} (as-
125 suming sub-vertical ascent of the plume in the upper mantle). Figure 7 shows the percentage of
126 database hotspots as a function of their distance from the nearest plume in Model0 (no thermo-
127 chemical pile), Model200 (200 km initial thickness), and Model300 (300 km initial thickness). A
128 lower distance for the database plume percentage indicates that there is a better fit between mod-
129 eled plume positions and the location of hotspots in the database (the model is performing well).
130 For instance, Model300 shows that all of the model plumes fall within 4000 km of a database
131 plume. However, less than 48% of the database plumes fall within 1500 km of a Model300 plume.
132 In comparison, Model0 and Model200 perform much better, with 51% and 59% of the database

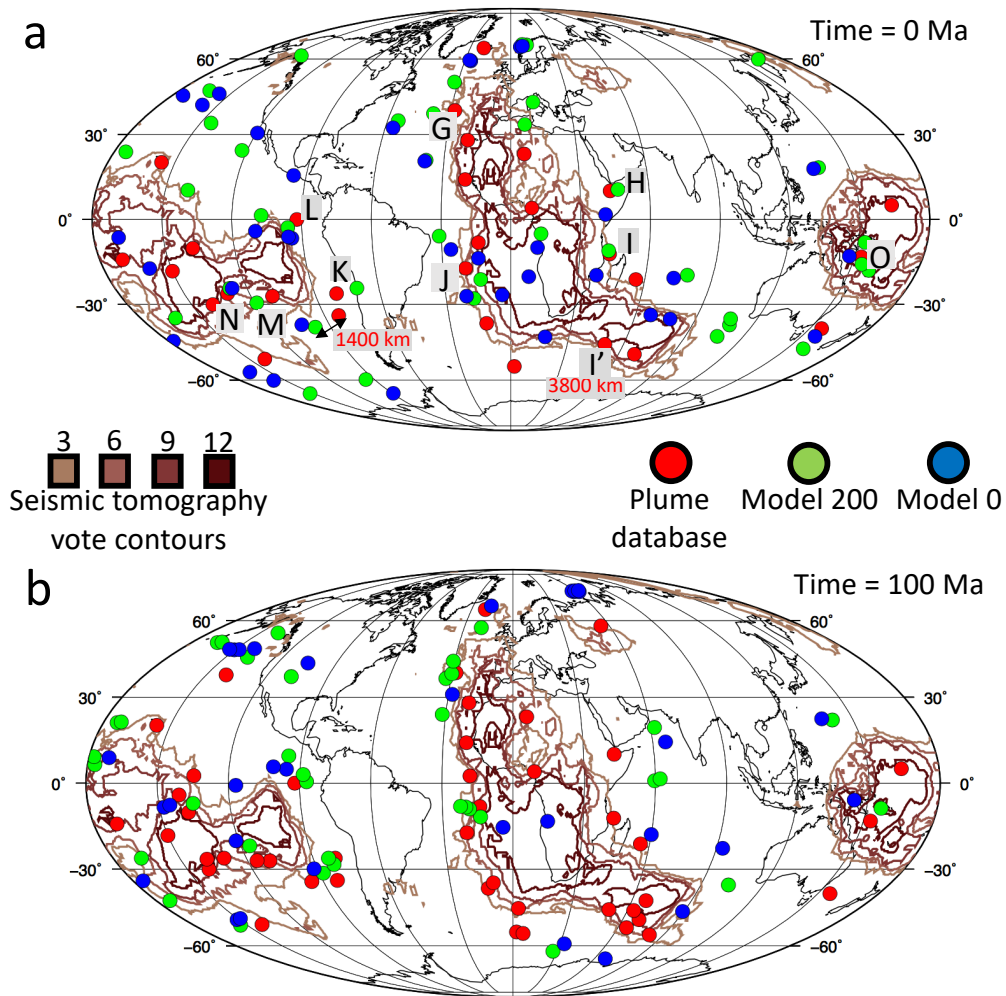


Figure 6: A number of model upwellings form close to the edges of LLSVPs and reach close to the surface near observed present-day hotspots. a) Plume positions at 0 Ma for Model200 (green), Model0 (blue), and the hotspot database^{1,4} (red) as given in SI Table S.1. b) Plume positions at 100 Ma for the above models alongside a large igneous province database for 300 – 100 Ma (red) as given in SI Table S.3². Plume key: G: Azores; H: Afar; I: Comores; J: St Helena; K: San Felix; L: Galapagos; M: Easter; N: Pitcairn; O: Louisiade. Two markers are given to indicate plume distances: 1400 km is shown between the Juan Fernandez plume and the nearest Model200 output; and 3800 km between I and I' (representing the largest distance between a database plume and a Model200 output).

134 Figure 7 also shows 10,000 randomly generated sets of 39 plume positions (grey lines). The
135 percentile distribution of these artificially generated points is given by the areas shaded in grey
136 and shows what percentage of the random models falls within a specific distance. In the case of
137 the lightest grey area, less than 15% of the random models have smaller distances to the database
138 plumes than shown. The profiles of Model0 and Model200 follows this 15% line closely, indicating
139 that these models perform better than 85% of the randomly generated results.

140 Model200s applies a 410 Myr plate reconstruction that uses an extended version of Seton
141 et al.⁵⁶ 250 Ma to 90 Ma surface velocities for 410 Ma to 250 Ma, then repeats from 250 Ma
142 until present-day⁴⁵ (0 Ma) using the rheological setup of Model200. In other words, the velocity
143 boundary conditions applied in the first 160 Myr of model evolution probably do not represent
144 past plate motions well. In applying these surface velocities, we can test the relative importance
145 of supercontinent formation (e.g., 410 Ma to 250 Ma) in developing plume positions. Figure 7
146 indicates that Model200s performs as well as Model200 for the majority of database plumes.
147 However, the surface velocities from the more recent global model of Matthews et al.⁴⁰ that in-
148 corporate supercontinent formation do have an impact on the shape of the thermo-chemical piles
149 in our models. Figure 4 indicates that applying appropriate surface velocities for the final 250 Myr
150 (e.g., Model200s) is not enough time to mould the thermo-chemical pile under Africa into place
151 (or, indeed, that piles are required to be already in position at that time). Although there is good
152 agreement between the TCP and the Pacific LLSVP, the African thermo-chemical pile appears to

153 be at a higher latitude when compared to a model that has undergone supercontinent formation and
154 dispersal surface velocities (e.g., Model200).

155 **Discussion**

156 In contrast to previous studies^{35–37}, we reproduce present-day plume positions both in models
157 that begin (at 410 Ma) with and without thermo-chemical piles (Figure 6a). This result reveals
158 that LLSVPs are a relatively passive feature of mantle dynamics (as shown in Figure 8), and that
159 ancient subducted ocean plates, rather than chemical heterogeneity, are the dominant factor driving
160 global mantle circulation.

161 The work presented here has wider implications on mantle dynamics. The proximity of an-
162 cient large igneous provinces (LIPs) to the edges of present-day LLSVPs has often been cited as a
163 reason why deep mantle heterogeneities may be dynamically stable over 500 Myr time-frames^{2,57}.
164 By analysing the upwelling locations for Model200 and Model0 during the Cretaceous (100 mil-
165 lion years in the past), the position of plumes still often occurs near the margins of present-day
166 LLSVPs (Figure 6b). This is significant given that our thermo-chemical piles are not fixed over
167 time (Figure 3 and 8) or not present at all (e.g., Model0). Consequently, LLSVPs do not need to
168 be laterally fixed on supercontinent timescales^{2,57} to produce appropriate plume positions in the
169 present or (potentially) in the past. Instead, the relative stability of subduction zones over the last
170 100 million years may be the crucial factor for fixing the proximity of past plumes to present-day
171 LLSVP margins. Our models indicate a decoupling of mantle timescales for plume positions (short

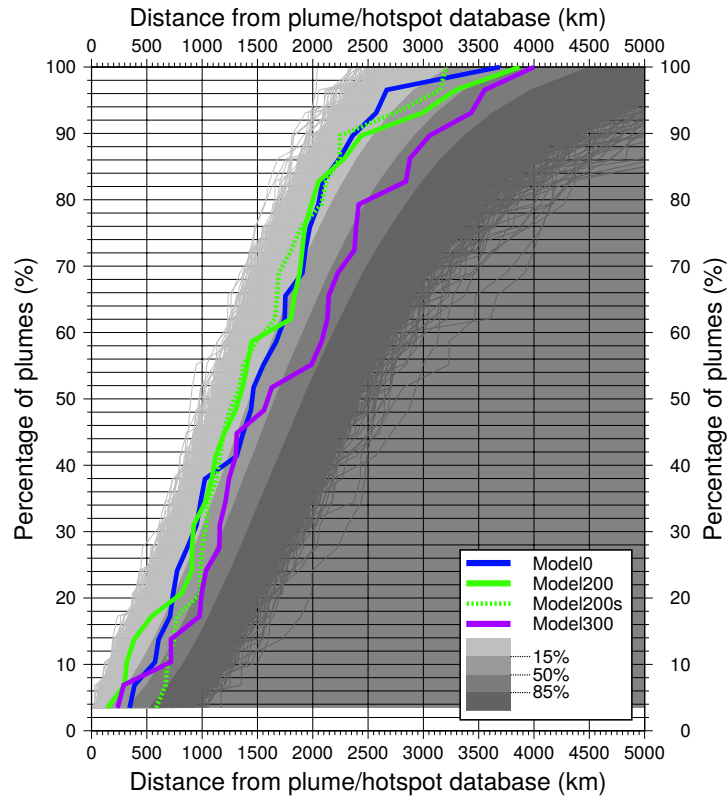


Figure 7: Analysis of model plume positions at 0 Ma as a function of distance away from the nearest present-day hotspot^{1,4}. For any distance given on the x-axis, the blue, green and purple lines show the percentage of model plumes that reach the surface closer to an observed present-day hotspot than this distance (for details on the analysis, see the Methods section). The purely thermal Model0 generates 37 plumes; Model200 (starting from a basal layer of dense material with 200 km thickness) features 39 plumes; and Model300 (with a 300-km basal layer) produces 32 plumes. Model200s utilizes 410 Myr of palaeo-subduction history using a modified reconstruction history (see Table S.2) and produces 29 plumes. Grey lines show 10,000 sets of 39 artificially generated random plume positions. The shaded grey areas indicate the percentile distance distribution 15%, 50%, and 85% of the 10,000 random models (see also SI Figure S.3). For most distances, Model0 and Model200 perform better than 85% of the random model runs.

scale) and LLSVP dynamics (longer scale), controlled by the plate motion history.

Model200s further indicates a decoupling in timescales between the development of plumes and the positioning of LLSVPs beneath Africa and the Pacific. For instance, by only running appropriate plate reconstruction history for the past 250 Ma (instead of 410 Ma in Model200), Model200s does not produce a thermo-chemical pile in the southern hemisphere under Africa (Figure 4e). However, in analysis of the plume positions, Model200s performs as well as Model200 (Figure 7). Although the past 250 Myr of subduction history is sufficient to develop a mantle flow with appropriate present-day upwellings, it may take longer for dense thermo-chemical piles to be moulded into position by subduction.

Although we highlight that LLSVPs do not need to be spatially fixed over supercontinent timescales, our models do not, and cannot, rule out an LLSVP being either purely thermal anomalies (Model0), (meta)stable chemical piles of dense material (Model200), or a combination of the two (Model100) (Figure 9). In contrast to our work, previous studies found that a thermo-chemical simulation, rather than a purely thermal model, is required to produce appropriate plume positions for present-day hotspots³⁵⁻³⁷. However, Davies et al.⁶ previously showed that observed lower mantle shear wave velocity anomalies do not require large-scale chemical piles to obtain appropriate LLSVP dynamics. Although Davies et al.⁶ did not conduct a formal plume analysis as given in our study, they did find that purely thermal LLSVPs (e.g., iso-chemical) could reconcile observed shear wave velocity anomalies and gradients. The work presented in our study highlights that purely thermal geodynamic models can not only produce seismological features of the deep

192 mantle^{e.g.}⁶, but also reproduce mantle dynamics in the form of plume locations (Figure 6a).

193 Our work implies that the decoupling of plume generation and large-scale mantle flow timescales
194 makes it difficult to unravel the composition of LLSVPs from numerical models alone (Figure 9).
195 As a result, we posit that the key to understanding the seismic anomalies in the present-day deep
196 mantle can only come from direct sampling, and therefore additional geochemical data and seismo-
197 logic observations. However, we suggest that the meta-stability of these structures can be cogent
198 with observed mantle dynamics (Figure 5), regardless of whether they are chemically distinct or
199 purely thermal (Figure 6), and signal that the deep mantle may evolve as significantly as our tec-
200 tonic surface (Figure 3 and 8).

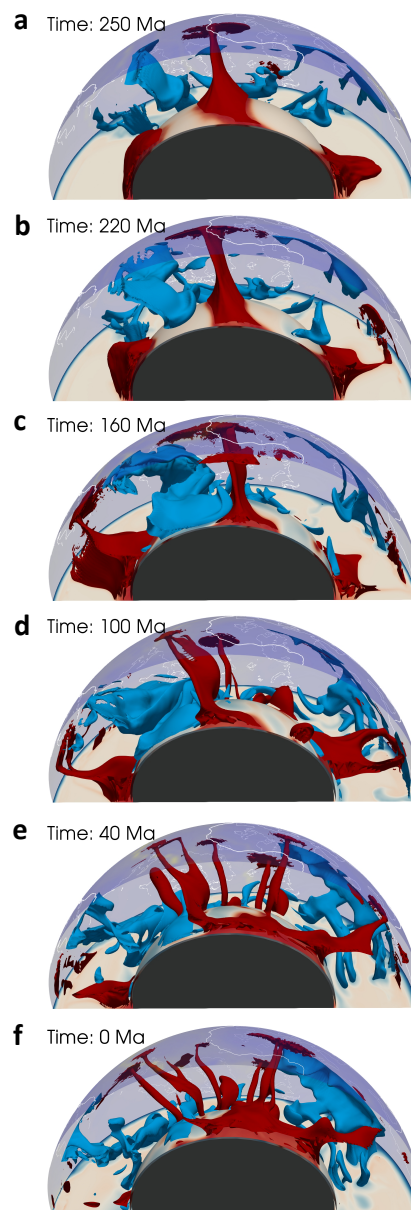


Figure 8: Thermo-chemical piles are passive in mantle dynamics. Evolution of the hot and cold anomalies under the African northern hemisphere from 250 Ma (a) to present-day (f) for Model200. Slice contours as Figure 5.

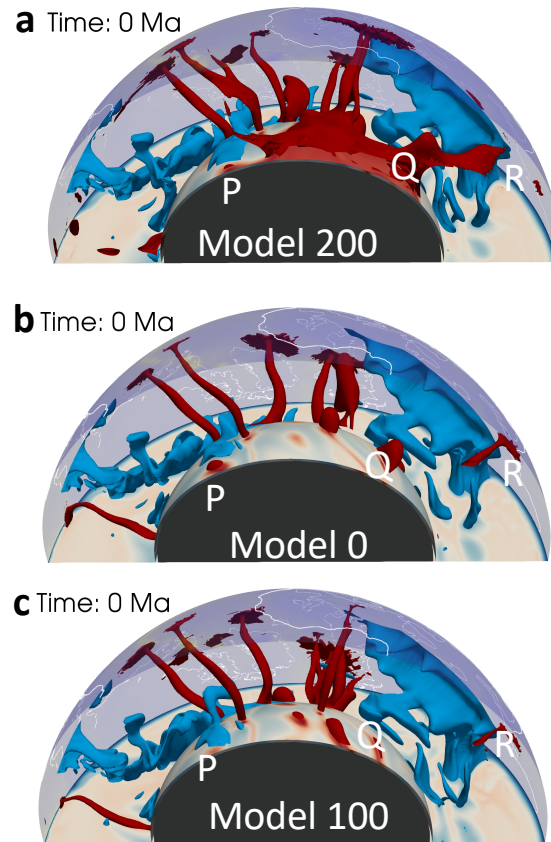


Figure 9: Models with different amount of chemical heterogeneity can produce appropriate plume dynamics. Present-day snapshot of the hot and cold anomalies under the African northern hemisphere for (a) Model200, (b) Model0, and (c) Model100. Slice contours as Figure 5. Markers P, Q, and R indicate equatorial plumes that could be linked to present-day hotspots St Helena, Comores, and Afar, respectively.

201 **Methods**

Model setup. Computations were done using the ASPECT code^{58,59}. ASPECT solves the following set of equations for compressible convection in the Earth’s mantle, describing the mass, force and energy balance (taking into account adiabatic heating, shear heating and radiogenic heat production), and the transport of chemical composition:

$$-\nabla \cdot (2\eta\dot{\epsilon}) + \nabla p = \rho\mathbf{g}, \quad (1)$$

$$\nabla \cdot (\rho\mathbf{u}) = 0, \quad (2)$$

$$\rho C_p \left(\frac{\partial T}{\partial t} + \mathbf{u} \cdot \nabla T \right) - \nabla \cdot k \nabla T = \rho H + 2\eta(\dot{\epsilon} : \dot{\epsilon}) + \alpha T (\mathbf{u} \cdot \nabla p), \quad (3)$$

$$\frac{\partial C}{\partial t} + \mathbf{u} \cdot \nabla C = 0. \quad (4)$$

202 The equations are solved for velocity \mathbf{u} , pressure p , temperature T and chemical composition C . η
 203 is the viscosity, $\dot{\epsilon} = \frac{1}{2}(\nabla\mathbf{u} + \nabla\mathbf{u}^T) - \frac{1}{3}(\nabla \cdot \mathbf{u})\mathbf{1}$ is the deviatoric strain rate, \mathbf{g} is the gravitational
 204 acceleration, ρ is the density, C_p is the specific heat capacity (at constant pressure), k is the thermal
 205 conductivity, H is the radiogenic heat production, and α is the thermal expansivity.

206 We use the global model setup⁵⁹ employed in a previous study⁵⁰, except for the following
 207 modifications:

- 208 • our plate reconstruction⁴⁰, which provides velocity boundary conditions at the surface of the
 209 model, extends back in time 410 million years,
- 210 • we vary the amount of dense recycled oceanic crust material (e.g., the thermo-chemical pile)
 211 that is present in the mantle, changing the thickness of the layer at the core-mantle boundary

212 where the material is distributed at the model start between 0 to 300 km (see SI Table S.2),

- 213 • we fix the core-mantle boundary temperature to 3700 K, corresponding to a buoyancy num-
214 ber $B = \frac{\Delta\rho}{\alpha\rho\Delta T} = 1.9$ (where $\Delta\rho$ is the density contrast between the anomalously dense
215 material and the average mantle composition and ΔT is the difference between the actual
216 core-mantle boundary temperature and the adiabatic profile) at a depth of 2500 km.

217 As in a previous study⁵⁰, we assume an average mantle composition of 82% harzburgite and 18%
218 recycled oceanic crust⁵⁴ and compute the material properties ρ , α and C_p using Perple_X⁵² and
219 a mineral physics database⁵³. Accordingly, all material properties include the effects of phase
220 transitions, with with thermal expansivity and specific heat approximating the corresponding latent
221 heat release⁶⁰. All model parameters are given in SI Table 1, and the initial radial profiles of
222 temperature- and depth-dependent material properties are plotted in Figure 2.

223 **Analysis of plume positions and thermo-chemical pile shape.** To determine the locations of
224 mantle plumes in our models, we define a plume as a positive temperature anomaly of at least 300
225 K excess temperature at a depth of 700 km. This depth was chosen in order to capture a first-order
226 location of the conduit and to avoid complications with plume head and upper mantle entrainment
227 complexities. These regions are then collated and spatially averaged into one data point for each
228 plume. Table S.1 and Fig 1 show the hotspot catalogue^{1,4} we use as a comparison for the output of
229 our geodynamic models. We quantify how well each model output fits this database of present-day
230 hotspots (as illustrated in Figures 7, S.3, S.4 and S.7) by using the following statistical method for
231 comparing two sets of points:

232 We calculate the distance d between the model plumes and the hotspots in the database using
 233 the Spherical Law of Cosines:

$$d = \arccos(\sin(\phi_{db}) \sin(\phi_m) + (\cos(\phi_{db}) \cos(\phi_m) \cos(\lambda_m - \lambda_{db})))r_e \quad (5)$$

234 where ϕ and λ are the latitude and longitude, respectively, of the plume location in the model (m)
 235 and the position of the closest hotspot in the database (db), and r_e is the radius of the Earth. The
 236 results are given in Fig 7, showing the percentage of the 29 hotspots in the database that are in a
 237 given distance from any of the modeled plumes.

238 To quantify how well each model fits the hotspot database, we use Equation 5 to compute the
 239 same statistics for 10,000 random plume positions, generated from a set of 39 random latitudes ϕ_m
 240 and longitudes λ_m . Each result of a random run is plotted as a grey line in Fig 7. These random
 241 runs allow us to compute percentiles (shown as shaded grey areas in Fig 7), indicating if the fit of
 242 a given model is better than, for example, the 15%, 50%, or 85% of the best random sets of plume
 243 positions. Here, the 50% percentile is equivalent to the average fit expected for a random plume
 244 distribution, and the 15% and 85% percentiles indicate approximately one standard deviation of
 245 this distribution. We generate the percentiles in Fig 7 through analysing histograms, with each
 246 histogram corresponding to a horizontal slice through the 10,000 random plume models in the
 247 diagram. In other words: To generate a histograms for a given cumulative fraction of hotspots x ,
 248 we pick the x database hotspots that are closest to any plume in a given random plume distribution.
 249 Each of the 10,000 data points in each of the histograms shows the maximum distance between any
 250 of these hotspots and the closest model plume. For example, a histogram showing the maximum
 251 distance of the 52% database hotspots that are closest to plumes in each of the 10,000 random

252 models is given in Fig S.3. This distribution allows us to determine the cutoff distance for the
253 15%, 50%, and 85% of the best random models. To generate a percentile distribution profile, as
254 given by the shaded areas in Fig 7, these cutoff distances are produced for all the cumulative plume
255 percentages (e.g., from 1 out of 29 plumes (3%) to 29 out of 29 (100%)).

256 To test the robustness of our findings we also computed a higher-resolution convection sim-
257 ulation with the same setup as Model200. Compared to Model200, we increased the resolution
258 by a factor of two, reduced the time step size by a factor of 2, and reduced the nonlinear solver
259 tolerance from 10^{-3} to 10^{-5} . The results of this Model200hr are shown in Figures S.5 to S.7. The
260 position of the thermo-chemical piles on an equatorial slice are similar in both models, but the
261 higher resolution model allows for more accurate (lower) thermal diffusion, reducing the thickness
262 of cold downwellings, as shown in Figure S.5.

263 For the same reason, we also find that in the upper mantle, the upwelling plumes are thinner
264 and not as hot in the higher resolution model. To generate a similar number of model plumes for the
265 statistical analysis as in the other models, we define plumes as regions with an excess temperature
266 of at least 200 K at 700 km depth (rather than 300 K as in Model200). This classification produces
267 29 Model200hr plumes in positions shown by the white circles in Figure S.6. The plume positions
268 are not identical to the lower resolution model (Figure S.6a) and the thermo-chemical pile covers
269 more of the core-mantle boundary than in Model200 (Figure S.6b). However, Figure S.7 does show
270 that the results are robust with regards to the plume positioning consistently performing better than
271 the random models.

272 Vote maps¹¹ are a simple tool that counts how many tomographic models detect the same
273 structure in a given location; a high “vote” count indicates agreement across the given suite of
274 models. In this study, the vote maps detect slow anomalous wavespeeds in order to detect the
275 edges of the LLSVPs. The vote maps were constructed with 14 tomography models (7 P-wave
276 and 7 S-wave) and methodology used in the construction of the vote maps are listed in a previous
277 study¹¹. The correlation between the thermo-chemical pile shape and present-day LLSVP outline
278 under Africa, as outlined in Table S.2, is calculated by the correlation between the area bound by
279 the 0.55 thermo-chemical ratio in the model and the 8 vote contour from seismic data¹¹.

280 The correlation coefficient is given as

$$r = \frac{\sum \sum (A - \bar{A})(B - \bar{B})}{\sqrt{(\sum \sum (A - \bar{A})^2)(\sum \sum (B - \bar{B})^2)}}, \quad (6)$$

281 where A is the model area and B is the vote map area.

282 **Code availability** The modelling code we used is open source and freely available online under
283 the terms of the [GNU General Public License](https://github.com/geodynamics/aspect) at <https://github.com/geodynamics/aspect>. For our
284 study, we have used version 2.0.0-pre (commit [572f967](https://github.com/geodynamics/aspect/commit/572f967)).

285 **Data availability** All input files that are required to reproduce our results are available online at
286 <https://github.com/jdannberg/paper-plume-positions-data>.

287 **References**

- 288 1. Steinberger, B. Plumes in a convecting mantle: Models and observations for individual
290 hotspots. *Journal of Geophysical Research: Solid Earth* **105**, 11127–11152 (2000). URL
291 <http://dx.doi.org/10.1029/1999JB900398>.
- 292 2. Torsvik, T. H., Smethurst, M. A., Burke, K. & Steinberger, B. Large igneous provinces gen-
293 erated from the margins of the large low-velocity provinces in the deep mantle. *Geophysical*
294 *Journal International* **167**, 1447–1460 (2006). URL [http://dx.doi.org/10.1111/
295 j.1365-246X.2006.03158.x](http://dx.doi.org/10.1111/j.1365-246X.2006.03158.x).
- 296 3. Cottaar, S. & Lekic, V. Morphology of seismically slow lower-mantle structures. *Geophysical*
297 *Journal International* **207**, 1122–1136 (2016). URL [http://dx.doi.org/10.1093/
298 gji/ggw324](http://dx.doi.org/10.1093/gji/ggw324).
- 299 4. French, S. W. & Romanowicz, B. Broad plumes rooted at the base of the earth’s mantle
300 beneath major hotspots. *Nature* **525**, 95–99 (2015). URL [http://dx.doi.org/10.
301 1038/nature14876](http://dx.doi.org/10.1038/nature14876).
- 302 5. Deschamps, F., Cobden, L. & Tackley, P. J. The primitive nature of large low shear-wave
303 velocity provinces. *Earth and Planetary Science Letters* **349-350**, 198–208 (2012). URL
304 <http://dx.doi.org/10.1016/j.epsl.2012.07.012>.
- 305 6. Davies, D. R. *et al.* Reconciling dynamic and seismic models of earth’s lower mantle: The
306 dominant role of thermal heterogeneity. *Earth and Planetary Science Letters* **353-354**, 253–

- 307
308 7. Koelemeijer, P., Deuss, A. & Ritsema, J. Density structure of earth's lowermost mantle from
309 stoneley mode splitting observations. *Nature Communications* **8**, 15241 (2017). URL <http://dx.doi.org/10.1038/ncomms15241>.
310
- 311 8. Zaroli, C., Koelemeijer, P. & Lambotte, S. Toward seeing the earth's interior
312 through unbiased tomographic lenses. *Geophysical Research Letters* **44**, 11,399–11,408
313 (2017). URL <https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1002/2017GL074996>.
314 <https://agupubs.onlinelibrary.wiley.com/doi/pdf/10.1002/2017GL074996>.
315
- 316 9. Garnero, E. J., McNamara, A. K. & Shim, S.-H. Continent-sized anomalous zones with low
317 seismic velocity at the base of earth's mantle. *Nature Geoscience* **9**, 481–489 (2016). URL
318 <http://dx.doi.org/10.1038/ngeo2733>.
- 319 10. van Hinsbergen, D. J. J., Steinberger, B., Doubrovine, P. V. & Gassmöller, R. Acceleration
320 and deceleration of india-asia convergence since the cretaceous: Roles of mantle plumes and
321 continental collision. *Journal of Geophysical Research* **116** (2011). URL <http://dx.doi.org/10.1029/2010JB008051>.
322
- 323 11. Shephard, G. E., Matthews, K. J., Hosseini, K. & Domeier, M. On the consistency of seis-
324 mically imaged lower mantle slabs. *Scientific Reports* **7** (2017). URL <http://dx.doi.org/10.1038/s41598-017-11039-w>.
325

- 326 12. Davies, D. & Davies, J. Thermally-driven mantle plumes reconcile multiple hot-spot obser-
327 vations. *Earth and Planetary Science Letters* **278**, 50 – 54 (2009). URL <http://www.sciencedirect.com/science/article/pii/S0012821X08007383>.
328
- 329 13. Allègre, C. J. Chemical geodynamics. *Tectonophysics* **81**, 109 – 132 (1982). URL <http://www.sciencedirect.com/science/article/pii/0040195182901251>. Geo-
330 physics final symposium.
331
- 332 14. White, W. M. Sources of oceanic basalts: Radiogenic isotopic evidence. *Geology* **13**, 115–118
333 (1985). URL [https://doi.org/10.1130/0091-7613\(1985\)13<115:SOOBRI>](https://doi.org/10.1130/0091-7613(1985)13<115:SOOBRI>2.0.CO;2)
334 [https://pubs.geoscienceworld.org/geology/article-pdf/](https://pubs.geoscienceworld.org/geology/article-pdf/13/2/115/3507682/i0091-7613-13-2-115.pdf)
335 [13/2/115/3507682/i0091-7613-13-2-115.pdf](https://pubs.geoscienceworld.org/geology/article-pdf/13/2/115/3507682/i0091-7613-13-2-115.pdf).
- 336 15. Allègre, C. J. & Turcotte, D. L. Implications of a two-component marble-cake mantle. *Nature*
337 **323**, 123–127 (1986). URL <http://dx.doi.org/10.1038/323123a0>.
- 338 16. Carlson, R. W. Mechanisms of earth differentiation: Consequences for the chemical structure
339 of the mantle. *Reviews of Geophysics* **32**, 337 (1994). URL <http://dx.doi.org/10.1029/94RG01874>.
340
- 341 17. Tackley, P. J. Mantle convection and plate tectonics: Toward an integrated physical and chem-
342 ical theory. *Science* **288**, 2002–2007 (2000). URL [http://dx.doi.org/10.1126/](http://dx.doi.org/10.1126/science.288.5473.2002)
343 [science.288.5473.2002](http://dx.doi.org/10.1126/science.288.5473.2002).
- 344 18. Tackley, P. Mantle geochemical geodynamics. *Treatise on Geophysics* 437–505 (2007). URL
345 <http://dx.doi.org/10.1016/B978-044452748-6.00124-3>.

- 346 19. Gonnermann, H. M. & Mukhopadhyay, S. Non-equilibrium degassing and a primordial source
347 for helium in ocean-island volcanism. *Nature* **449**, 1037–1040 (2007). URL [http://dx.
doi.org/10.1038/nature06240](http://dx.
348 doi.org/10.1038/nature06240).
- 349 20. Gonnermann, H. M. & Mukhopadhyay, S. Preserving noble gases in a convecting mantle.
350 *Nature* **459**, 560–563 (2009). URL [http://dx.doi.org/10.1038/
nature08018](http://dx.doi.org/10.1038/nature08018).
- 351 21. Jackson, M. G. *et al.* Helium and lead isotopes reveal the geochemical geometry of the
352 samoan plume. *Nature* **514**, 355–358 (2014). URL [http://dx.doi.org/10.1038/
nature13794](http://dx.doi.org/10.1038/
353 nature13794).
- 354 22. White, W. Probing the earth’s deep interior through geochemistry. *Geochemical Perspectives*
355 **95–251** (2015). URL <http://dx.doi.org/10.7185/geochempersp.4.2>.
- 356 23. Weis, D., Garcia, M. O., Rhodes, J. M., Jellinek, M. & Scoates, J. S. Role of the deep mantle
357 in generating the compositional asymmetry of the hawaiian mantle plume. *Nature Geoscience*
358 **4**, 831–838 (2011). URL <http://dx.doi.org/10.1038/ngeo1328>.
- 359 24. Farnetani, C. G., Hofmann, A. W. & Class, C. How double volcanic chains sample geochem-
360 ical anomalies from the lowermost mantle. *Earth and Planetary Science Letters* **359-360**,
361 240–247 (2012). URL <http://dx.doi.org/10.1016/j.epsl.2012.09.057>.
- 362 25. Williams, C. D., Li, M., McNamara, A. K., Garnero, E. J. & van Soest, M. C. Episodic entrain-
363 ment of deep primordial mantle material into ocean island basalts. *Nature Communications* **6**
364 (2015). URL <http://dx.doi.org/10.1038/ncomms9937>.

- 365 26. Ishii, M. Normal-mode and free-air gravity constraints on lateral variations in velocity and
366 density of earth's mantle. *Science* **285**, 1231–1236 (1999). URL [http://dx.doi.org/
367 10.1126/science.285.5431.1231](http://dx.doi.org/10.1126/science.285.5431.1231).
- 368 27. Masters, G. & Gubbins, D. On the resolution of density within the earth. *Physics of the Earth
369 and Planetary Interiors* **140**, 159–167 (2003). URL [http://dx.doi.org/10.1016/
370 j.pepi.2003.07.008](http://dx.doi.org/10.1016/j.pepi.2003.07.008).
- 371 28. Simmons, N. A., Forte, A. M., Boschi, L. & Grand, S. P. Gypsum: A joint tomographic model
372 of mantle density and seismic wave speeds. *Journal of Geophysical Research* **115** (2010).
373 URL <http://dx.doi.org/10.1029/2010JB007631>.
- 374 29. Trampert, J. Probabilistic tomography maps chemical heterogeneities throughout the
375 lower mantle. *Science* **306**, 853–856 (2004). URL [http://dx.doi.org/10.1126/
376 science.1101996](http://dx.doi.org/10.1126/science.1101996).
- 377 30. Koelemeijer, P. J., Deuss, A. & Trampert, J. Normal mode sensitivity to earth's d layer
378 and topography on the core-mantle boundary: what we can and cannot see. *Geophysical
379 Journal International* **190**, 553–568 (2012). URL [http://dx.doi.org/10.1111/j.
380 1365-246X.2012.05499.x](http://dx.doi.org/10.1111/j.1365-246X.2012.05499.x).
- 381 31. Lau, H. C. P. *et al.* Tidal tomography constrains earth's deep-mantle buoyancy. *Nature* **551**,
382 321 EP – (2017). URL <https://doi.org/10.1038/nature24452>.
- 383 32. Trampert, J., Vacher, P. & Vlaar, N. Sensitivities of seismic velocities to temperature, pressure
384 and composition in the lower mantle. *Physics of the Earth and Planetary Interiors* **124**, 255

- 385 – 267 (2001). URL [http://www.sciencedirect.com/science/article/pii/](http://www.sciencedirect.com/science/article/pii/S0031920101002011)
386 [S0031920101002011](http://www.sciencedirect.com/science/article/pii/S0031920101002011).
- 387 33. Deschamps, F. & Trampert, J. Mantle tomography and its relation to temperature and com-
388 position. *Physics of the Earth and Planetary Interiors* **140**, 277 – 291 (2003). URL [http://](http://www.sciencedirect.com/science/article/pii/S0031920103001894)
389 www.sciencedirect.com/science/article/pii/S0031920103001894.
- 390 34. Schubert, G., Masters, G., Olson, P. & Tackley, P. Superplumes or plume clusters? *Physics of*
391 *the Earth and Planetary Interiors* **146**, 147–162 (2004). URL [http://dx.doi.org/10.](http://dx.doi.org/10.1016/j.pepi.2003.09.025)
392 [1016/j.pepi.2003.09.025](http://dx.doi.org/10.1016/j.pepi.2003.09.025).
- 393 35. Steinberger, B. & Torsvik, T. H. A geodynamic model of plumes from the margins of large
394 low shear velocity provinces. *Geochemistry, Geophysics, Geosystems* **13**, n/a–n/a (2012). URL
395 <http://dx.doi.org/10.1029/2011GC003808>.
- 396 36. Hassan, R., Flament, N., Gurnis, M., Bower, D. J. & Müller, D. Provenance of plumes in
397 global convection models. *Geochemistry, Geophysics, Geosystems* **16**, 1465–1489 (2015).
398 URL <http://dx.doi.org/10.1002/2015GC005751>.
- 399 37. Li, M. & Zhong, S. The source location of mantle plumes from 3d spherical models of mantle
400 convection. *Earth and Planetary Science Letters* **478**, 47 – 57 (2017). URL [http://www.](http://www.sciencedirect.com/science/article/pii/S0012821X17304806)
401 [sciencedirect.com/science/article/pii/S0012821X17304806](http://www.sciencedirect.com/science/article/pii/S0012821X17304806).
- 402 38. Torsvik, T. H., Burke, K., Steinberger, B., Webb, S. J. & Ashwal, L. D. Diamonds sampled by
403 plumes from the core–mantle boundary. *Nature* **466**, 352–355 (2010). URL [http://dx.](http://dx.doi.org/10.1038/nature09216)
404 [doi.org/10.1038/nature09216](http://dx.doi.org/10.1038/nature09216).

- 405 39. Burke, K. Plate tectonics, the wilson cycle, and mantle plumes: Geodynamics from the top.
406 *Annual Review of Earth and Planetary Sciences* **39**, 1–29 (2011). URL <http://dx.doi.org/10.1146/annurev-earth-040809-152521>.
407
- 408 40. Matthews, K. J. *et al.* Global plate boundary evolution and kinematics since the late pa-
409 leozoic. *Global and Planetary Change* **146**, 226 – 250 (2016). URL <http://www.sciencedirect.com/science/article/pii/S0921818116302417>.
410
- 411 41. Domeier, M. & Torsvik, T. H. Plate tectonics in the late paleozoic. *Geoscience Frontiers* **5**,
412 303 – 350 (2014). URL <http://www.sciencedirect.com/science/article/pii/S1674987114000061>.
413
- 414 42. Müller, R. D. *et al.* Ocean basin evolution and global-scale plate reorganization events since
415 pangea breakup. *Annual Review of Earth and Planetary Sciences* **44**, 107–138 (2016). URL
416 <https://doi.org/10.1146/annurev-earth-060115-012211>.
- 417 43. McNamara, A. K. & Zhong, S. Thermochemical structures beneath africa and the pa-
418 cific ocean. *Nature* **437**, 1136–1139 (2005). URL <http://dx.doi.org/10.1038/nature04066>.
419
- 420 44. Zhang, N., Zhong, S., Leng, W. & Li, Z.-X. A model for the evolution of the earth’s mantle
421 structure since the early paleozoic. *Journal of Geophysical Research* **115** (2010). URL <http://dx.doi.org/10.1029/2009JB006896>.
422

- 423 45. Bower, D. J., Gurnis, M. & Seton, M. Lower mantle structure from paleogeographically
424 constrained dynamic earth models. *Geochemistry, Geophysics, Geosystems* **14**, 44–63 (2013).
425 URL <http://dx.doi.org/10.1029/2012GC004267>.
- 426 46. Hassan, R., Müller, R. D., Gurnis, M., Williams, S. E. & Flament, N. A rapid burst in hotspot
427 motion through the interaction of tectonics and deep mantle flow. *Nature* **533**, 239–242 (2016).
428 URL <http://dx.doi.org/10.1038/nature17422>.
- 429 47. Flament, N., Williams, S., Müller, R. D., Gurnis, M. & Bower, D. J. Origin and evolution of
430 the deep thermochemical structure beneath eurasia. *Nature Communications* **8**, 14164 (2017).
431 URL <http://dx.doi.org/10.1038/ncomms14164>.
- 432 48. Zhang, N. & Li, Z.-X. Formation of mantle “lone plumes” in the global downwelling zone
433 — a multiscale modelling of subduction-controlled plume generation beneath the south china
434 sea. *Tectonophysics* **723**, 1 – 13 (2018). URL [http://www.sciencedirect.com/
435 science/article/pii/S0040195117305000](http://www.sciencedirect.com/science/article/pii/S0040195117305000).
- 436 49. Heron, P. J. Mantle plumes and mantle dynamics in the wilson cycle. *Geological Society,*
437 *London, Special Publications* SP470.18 (2018). URL [http://dx.doi.org/10.1144/
438 SP470.18](http://dx.doi.org/10.1144/SP470.18).
- 439 50. Dannberg, J. & Gassmüller, R. Chemical trends in ocean islands explained by plume–slab
440 interaction. *Proceedings of the National Academy of Sciences* **115**, 4351–4356 (2018). URL
441 <http://dx.doi.org/10.1073/pnas.1714125115>.

- 442 51. Steinberger, B. & Calderwood, A. R. Models of large-scale viscous flow in the earth's mantle
443 with constraints from mineral physics and surface observations. *Geophysical Journal Interna-*
444 *tional* **167**, 1461–1481 (2006). URL [http://dx.doi.org/10.1111/j.1365-246X.](http://dx.doi.org/10.1111/j.1365-246X.2006.03131.x)
445 [2006.03131.x](http://dx.doi.org/10.1111/j.1365-246X.2006.03131.x).
- 446 52. Connolly, J. A. D. The geodynamic equation of state: What and how. *Geochemistry,*
447 *Geophysics, Geosystems* **10**, n/a–n/a (2009). URL [http://dx.doi.org/10.1029/](http://dx.doi.org/10.1029/2009GC002540)
448 [2009GC002540](http://dx.doi.org/10.1029/2009GC002540).
- 449 53. Stixrude, L. & Lithgow-Bertelloni, C. Thermodynamics of mantle minerals - ii. phase equi-
450 libria. *Geophysical Journal International* **184**, 1180–1213 (2011). URL [http://dx.doi.](http://dx.doi.org/10.1111/j.1365-246X.2010.04890.x)
451 [org/10.1111/j.1365-246X.2010.04890.x](http://dx.doi.org/10.1111/j.1365-246X.2010.04890.x).
- 452 54. Xu, W., Lithgow-Bertelloni, C., Stixrude, L. & Ritsema, J. The effect of bulk composition
453 and temperature on mantle seismic structure. *Earth and Planetary Science Letters* **275**, 70–79
454 (2008). URL <http://dx.doi.org/10.1016/j.epsl.2008.08.012>.
- 455 55. Ni, S., Tan, E., Gurnis, M. & Helmberger, D. Sharp sides to the african superplume. *Sci-*
456 *ence* **296**, 1850–1852 (2002). URL [https://science.sciencemag.org/content/](https://science.sciencemag.org/content/296/5574/1850)
457 [296/5574/1850](https://science.sciencemag.org/content/296/5574/1850). [https://science.sciencemag.org/content/296/5574/](https://science.sciencemag.org/content/296/5574/1850.full.pdf)
458 [1850.full.pdf](https://science.sciencemag.org/content/296/5574/1850.full.pdf).
- 459 56. Seton, M. *et al.* Global continental and ocean basin reconstructions since 200ma. *Earth-*
460 *Science Reviews* **113**, 212 – 270 (2012). URL [http://www.sciencedirect.com/](http://www.sciencedirect.com/science/article/pii/S0012825212000311)
461 [science/article/pii/S0012825212000311](http://www.sciencedirect.com/science/article/pii/S0012825212000311).

- 462 57. Torsvik, T. H., Steinberger, B., Cocks, L. R. M. & Burke, K. Longitude: Linking earth's
463 ancient surface to its deep interior. *Earth and Planetary Science Letters* **276**, 273 –
464 282 (2008). URL [http://www.sciencedirect.com/science/article/pii/
465 S0012821X08006328](http://www.sciencedirect.com/science/article/pii/S0012821X08006328).
- 466 58. Kronbichler, M., Heister, T. & Bangerth, W. High accuracy mantle convection simulation
467 through modern numerical methods. *Geophysical Journal International* **191**, 12–29 (2012).
468 URL <http://dx.doi.org/10.1111/j.1365-246X.2012.05609.x>.
- 469 59. Heister, T., Dannberg, J., Gassmüller, R. & Bangerth, W. High accuracy mantle convection
470 simulation through modern numerical methods. II: Realistic models and problems. *Geophys-
471 ical Journal International* **210**, 833–851 (2017). URL [https://doi.org/10.1093/
472 gji/ggx195](https://doi.org/10.1093/gji/ggx195).
- 473 60. Nakagawa, T., Tackley, P. J., Deschamps, F. & Connolly, J. A. D. Incorporating self-
474 consistently calculated mineral physics into thermochemical mantle convection simulations in
475 a 3-d spherical shell and its influence on seismic anomalies in earth's mantle. *Geochemistry,
476 Geophysics, Geosystems* **10**, n/a–n/a (2009). URL [http://dx.doi.org/10.1029/
477 2008GC002280](http://dx.doi.org/10.1029/2008GC002280).
- 478 61. Loken, C. *et al.* Scinet: Lessons learned from building a power-efficient top-20 system and
479 data centre. *Journal of Physics: Conference Series* **256**, 012026 (2010). URL [http://dx.
480 doi.org/10.1088/1742-6596/256/1/012026](http://dx.doi.org/10.1088/1742-6596/256/1/012026).

481 62. French, S. W. & Romanowicz, B. A. Whole-mantle radially anisotropic shear velocity struc-
482 ture from spectral-element waveform tomography. *Geophysical Journal International* **199**,
483 1303–1327 (2014). URL <http://dx.doi.org/10.1093/gji/ggu334>.

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501 **Competing Interests.** The authors declare that they have no competing financial interests.

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504 **S Supporting Information**

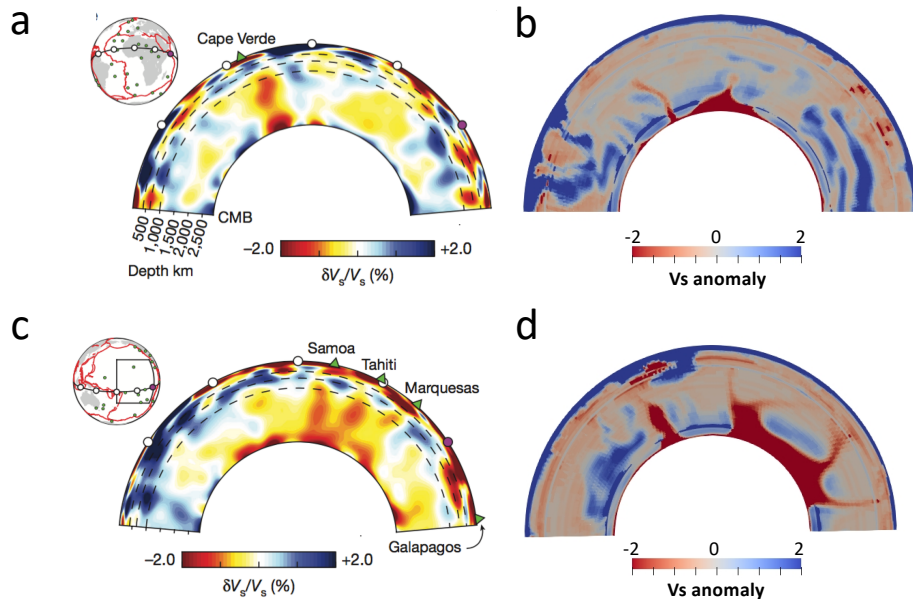


Figure S.1: Whole-mantle depth cross-section modified from⁴ of relative shear-velocity variations using model SEMUCB-WM1⁶² near Cape Verde (a) and Marquesas (c), with the corresponding shear velocity anomaly output from Model200 shown in (b) and (d), respectively.

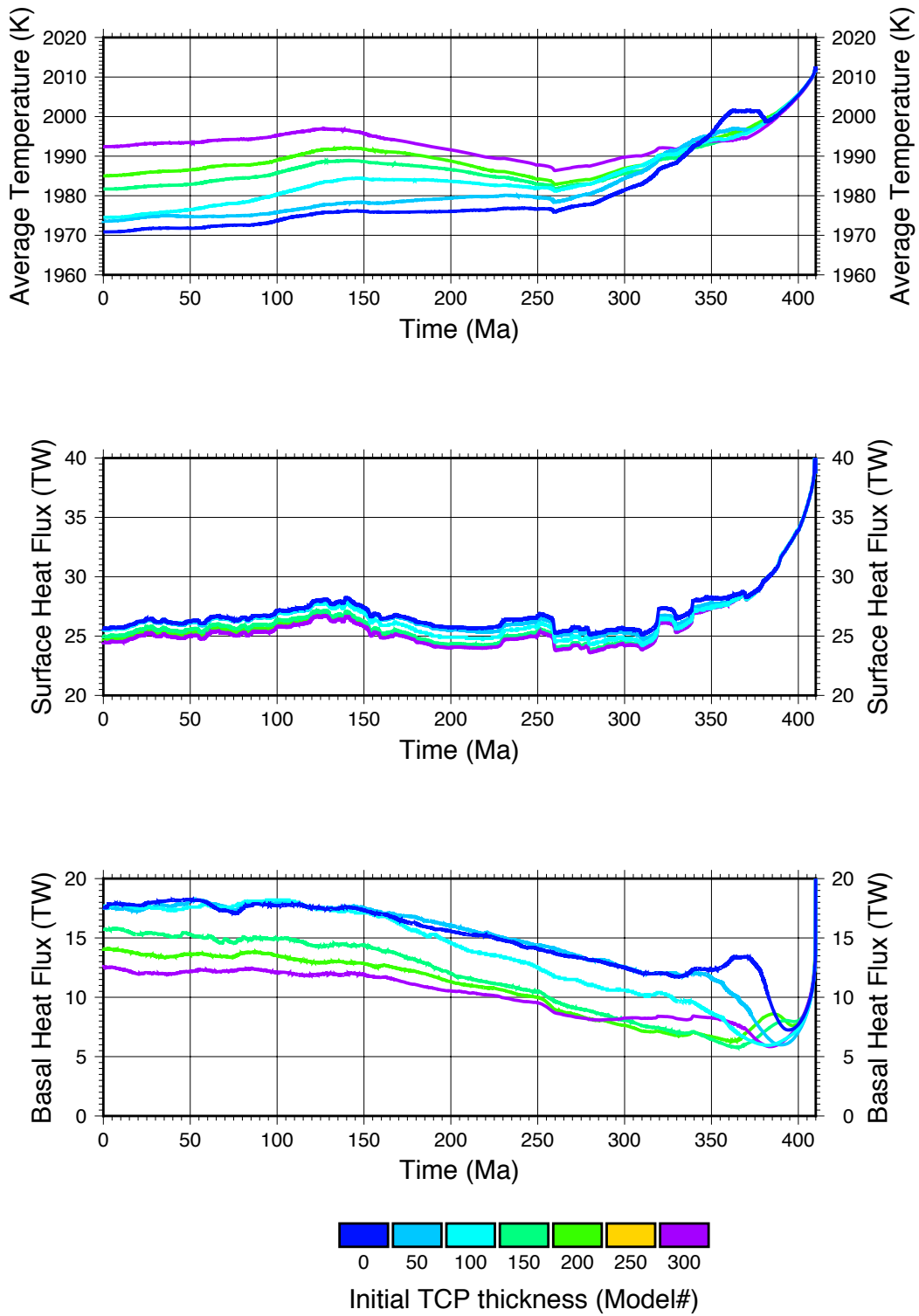


Figure S.2: Thermal evolution of the models. a) Average mantle temperature, b) surface heat flux,

c) basal heat flux.

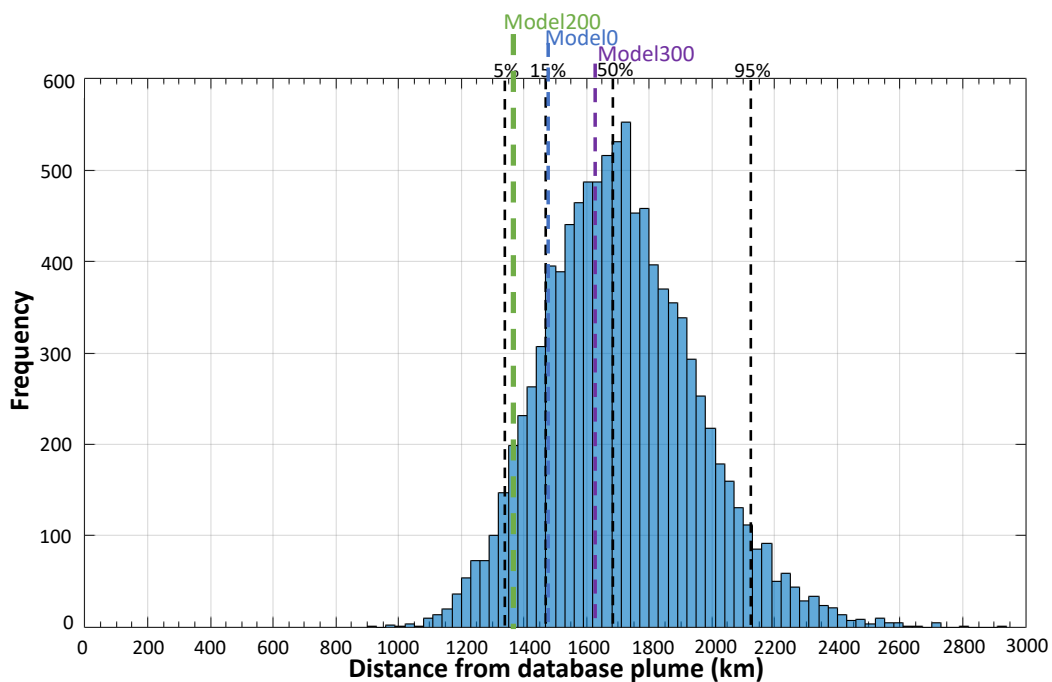


Figure S.3: Model200 performs better than around 95% of the 10,000 random models. Histogram showing the spread of distances from 52% of the database plumes. Annotated are the 5%, 15%, 50%, and 95% distributions of the random models, as well as the Model200, Model0, and Model300 results.

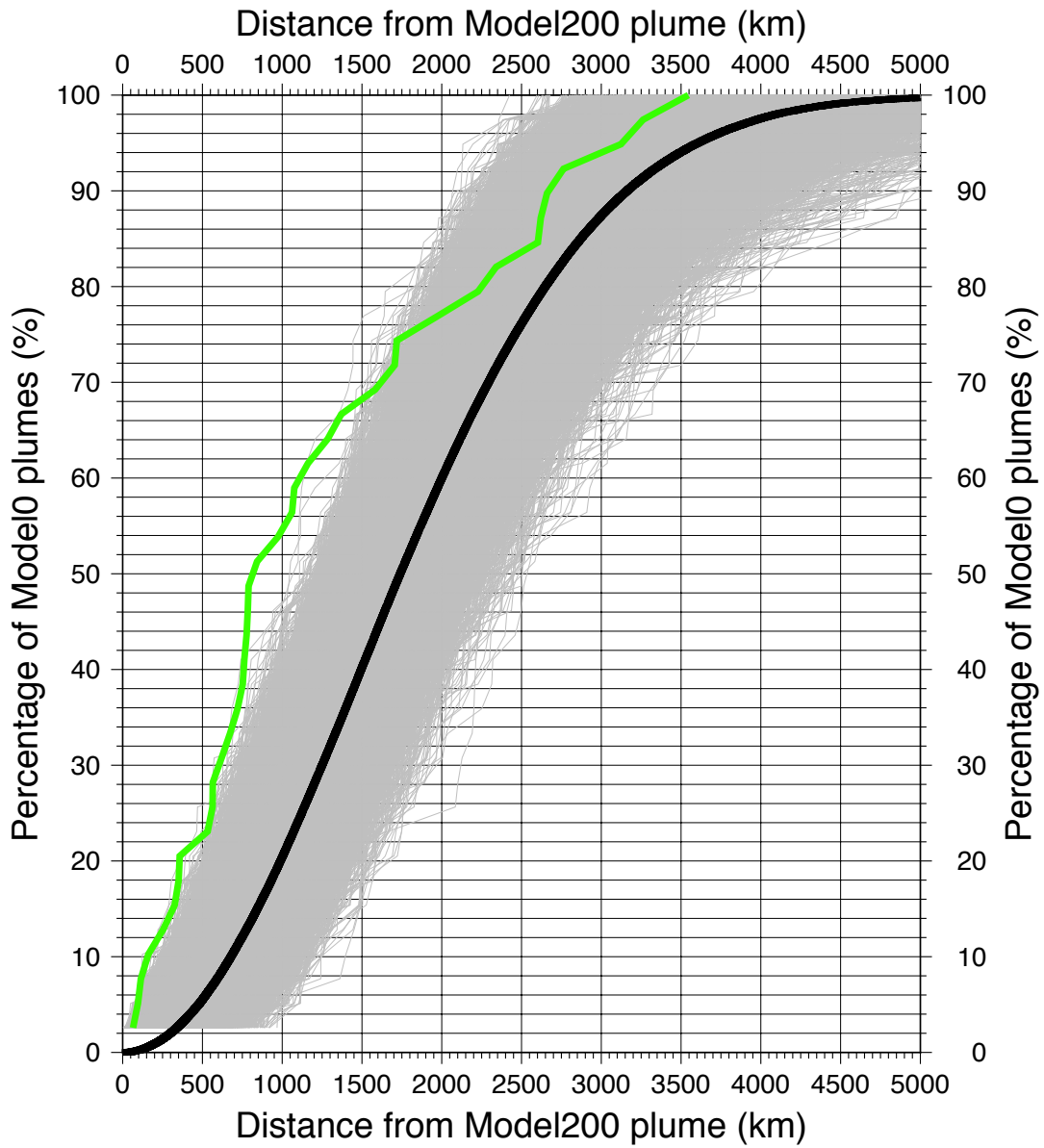


Figure S.4: Analysis of Model0 plume positions at 0 Ma as a function of distance away from nearest Model200 plume. Models have differing number of plumes (see main text). Model0 features no thermo-chemical pile (TCP) and Model200 has a 200 km initial TCP thickness. The black line is the accumulation of 10,000 sets of 39 random plume positions artificially generated (grey lines).

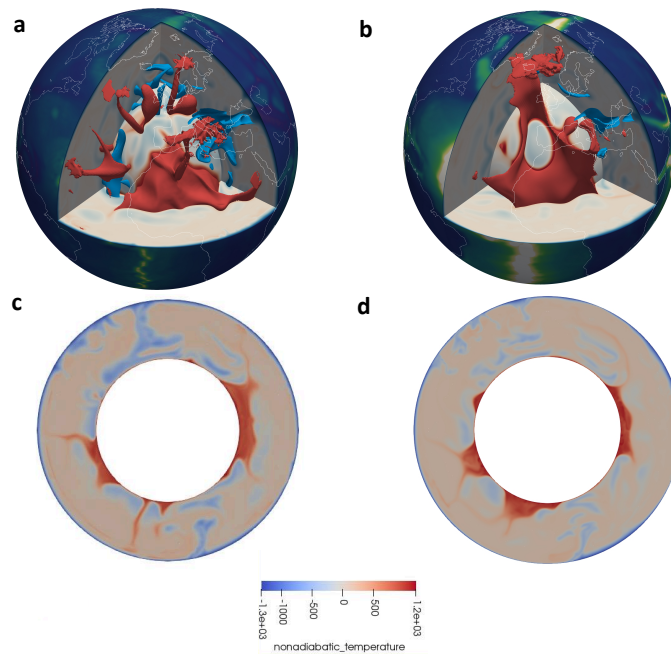


Figure S.5: Comparison between Model200 and Model200hr (a higher resolution simulation). Figure shows a view into the mantle beneath the northern Atlantic and Europe for Model200 (a) and Model200hr (b). Excess temperature anomaly contours are given for warm (300 K, red) and cold (-500 K, blue) regions. Excess temperature at the equator is given for Model200 (c) and Model200hr (d).

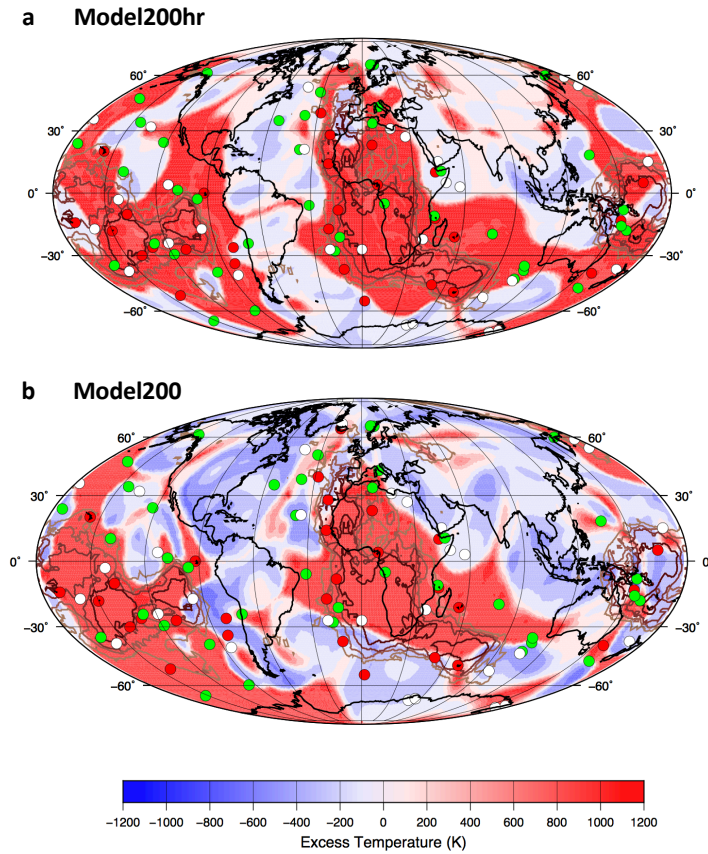


Figure S.6: Comparison between Model200 and Model200hr (a higher resolution simulation). The excess temperature at 2800 km depth for Model200hr is given in (a), with seismic tomography contours as Figure 1. Green circles and white circles show Model200 and Model200hr plume positions, respectively. Red circles are of the present-day plume database^{1,4}. See Methods for Model200hr description. The excess temperature at 2800 km depth for Model200 is given in (b) for comparison. Circles as (a).

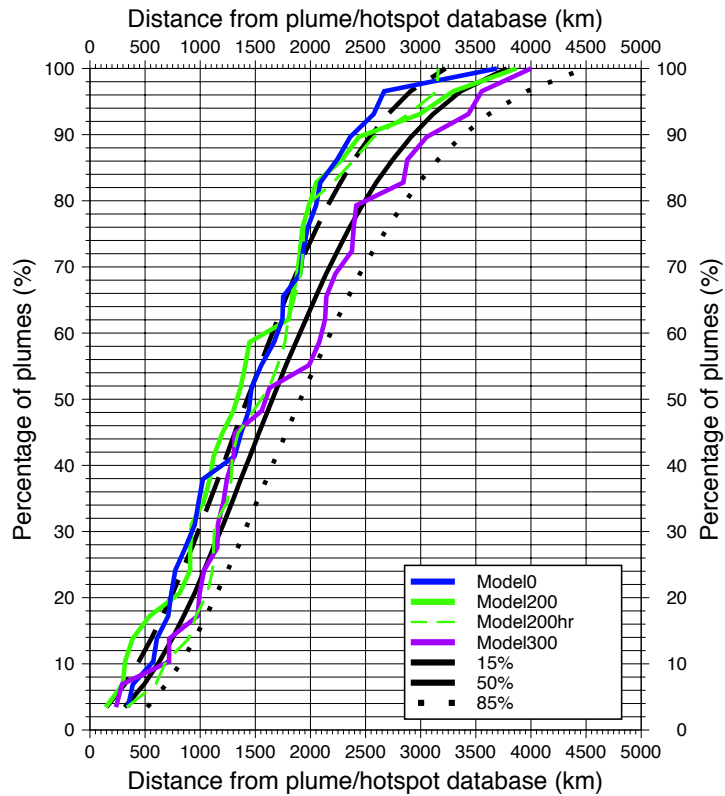


Figure S.7: Analysis of model plume positions at 0 Ma as a function of distance away from nearest present-day plume^{1,4}. This image is as Figure 7 (see caption for details) but with the addition of a high resolution Model200 (Model200hr). See Methods for description of high resolution model.

Table S.1: Database of deep plume and hotspot locations^{1,4}.

#	Latitude	Longitude	Notes
1	-14	-170	Samoa
2	20	-156	Hawaii
3	-18	-150	Tahiti/Society
4	-51	-141	Louisville
5	-30	-140	Macdonald
6	-10	-138	Marquesas
7	-26	-130	Pitcairn
8	-27	-110	Easter
9	0	-92	Galapagos
10	-34	-83	Juan Fernandez
11	-26	-80	San Felix
12	39	-28	Azores
13	28	-20	Mid-Atlantic
14	14	-20	Cape Verde
15	65	-20	Iceland
16	-17	-20	St Helena
17	-8	-14	Ascension
18	-37	-12	Tristan
19	-54	2	Bouvet
20	23	6	Hoggar
21	4	9	Cameroon
22	10	43	Afar
23	-12	43	Comores
24	-45	50	Crozet
25	-21	56	Reunion
26	-49	69	Kerguelen
27	-13	153	Louisade
28	-39	156	Tasmanid
29	5	164	Caroline

Table S.2: List of models with corresponding initial thermo-chemical (TC) thickness, plate reconstruction study, and correlation values. The correlation value is the direct comparison between the African hemisphere slow contour of 8 in the vote map analysis¹¹ and the outline of the model African thermo-chemical pile. 'Modified' refers to a 410 Myr reconstruction that uses Seton et al.⁵⁶ 250 Ma to 90 Ma surface velocities for 410 Ma to 250 Ma, then repeats until present day.

Name	Initial TC	Plate reconstruction	Correlation
Model0	0 km	Matthews et al. ⁴⁰	N/A
Model50	50 km	Matthews et al. ⁴⁰	N/A
Model100	100 km	Matthews et al. ⁴⁰	N/A
Model150	150 km	Matthews et al. ⁴⁰	0.71
Model200	200 km	Matthews et al. ⁴⁰	0.84
Model250	250 km	Matthews et al. ⁴⁰	-
Model300	300 km	Matthews et al. ⁴⁰	0.81
Model200s	200 km	Modified Seton et al. ⁵⁶	-

Table S.3: Database of large igneous province positions^{2,57} from 250 Ma to 100 Ma.

Age (Ma)	Latitude	Longitude	LIP name
251	57.7	54.7	Siberian Traps
200	2.5	341.9	CAMP
182	-44.6	2.8	Karoo Basalts
147	-4	219	Shatsky Rise
145	-26.9	244.3	Megallan Rise
132	-55.3	81.6	Bunbury Basalts
132	-34.9	350.6	Parana-Etendeka
125	-54.7	5.9	Maud Rise
123	-34.5	264.5	Manihiki Plateau
121	38.2	219.7	Ontong Java Plateau
118	-41.6	67	Rajmahal Traps
114	-52.3	64.9	South Kerguelen
111	-26.4	222	Naura Basalts
100	-45.4	63.3	Central Kerguelen
99	2.6	225.4	Hess Rise