

Estimating the mass eruption rate of volcanic eruptions from the plume height using Bayesian regression with historical data: the MERPH model

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## Highlights

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- A Bayesian linear regression is applied to eruption source parameter datasets
- Inference of mass eruption rate from plume height can be obtained with a quantification of uncertainty
- Analytical forms for the Posterior Predictive distributions facilitate analysis and sampling
- Observational uncertainty can be included in prediction

# Estimating the mass eruption rate of volcanic eruptions from the plume height using Bayesian regression with historical data: the MERPH model

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## **Abstract**

The mass eruption rate (MER) of an explosive volcanic eruption is a commonly used quantifier of the magnitude of the eruption, and estimating it is important in managing volcanic hazards. The physical connection between the MER and the rise height of the eruption column results in a scaling relationship between these quantities, allowing one to be inferred from the other. Eruption source parameter datasets have been used to calibrate the relationship, but the uncertainties in the measurements used in the calibration are typically not accounted for in applications. This can lead to substantial over- or under-estimation. Here we apply a simple Bayesian approach to incorporate uncertainty into the calibration of the scaling relationship using Bayesian linear regression to determine probability density functions for model parameters. This allows probabilistic prediction of mass eruption rate given a plume height observation in a way that is consistent with the data used for calibration. By using non-informative priors, the posterior predictive distribution can be determined analytically. The methods and dataset are collected in a python package, called `merph`, and we illustrate their use in

sampling plausible MER–plume height pairs, and in identifying usual eruptions. We discuss applications to ensemble-based hazard assessments and potential developments of the approach.

*Keywords:* Mass eruption rate, Plume height, Uncertainty quantification,, Bayesian regression

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## 1. Introduction

The mass eruption rate (MER, denoted throughout by  $Q$ ) of an explosive volcanic eruption is a quantity of fundamental importance in volcanology. The MER is often used (along with other quantities) to classify the strength of an eruption (Walker, 1980; Bonadonna and Costa, 2013; Pyle, 2015). During eruptions, rapid estimation of the MER is important for predicting hazards and monitoring evolution of the activity. The MER is an essential input to volcanic tephra transport and deposition models that forecast the dispersion of volcanic ash in the atmosphere (Mastin et al., 2009; Folch, 2012).

A common approach for making rapid estimates of the MER is to use an empirical algebraic expression relating MER to the plume height (denoted here by  $H$ ). Data from historical eruptions, where there are independent estimates of the MER and plume height, have been compiled since the 1978 papers of Settle (1978) and Wilson et al. (1978) that gathered data to support the application of theoretical models of buoyant convection (Morton et al., 1956) to volcanic eruptions. While relatively small datasets ( $n = 6$  eruptions in Settle 1978 and  $n = 8$  eruptions in Wilson et al. 1978), the eruption MER spanned several orders of magnitude, with corresponding plume heights from the low troposphere to the high stratosphere. In the following decades these

20 catalogues have been extended and revised, notably in Sparks (1986) ( $n = 8$   
21 eruptions), Sparks et al. (1997) ( $n = 28$ ), Mastin et al. (2009) ( $n = 35$ ) and  
22 Aubry et al. (2021) ( $n = 130$ ).

23 The eruption datasets can be used to calibrate the parameters in a re-  
24 gression model. In both Sparks et al. (1997) and Mastin et al. (2009) a  
25 power-law relationship is used to model the dependence of the plume height  
26 on the MER, i.e.  $Q = 10^\alpha H^\beta$ , and the fitting parameters  $\alpha$  and  $\beta$  are found  
27 from linear regression of  $x = \log H$  and  $y = \log Q$  with the observational  
28 dataset. This form is consistent with theoretical models of buoyant turbu-  
29 lent convection in a linear stable stratification, where the exponent  $\beta = 4$   
30 (Morton et al., 1956). The curves obtained by regression, shown in figure 1,  
31 describe the leading-order behaviour seen in the data, but there is substan-  
32 tial scatter of the data points around the calibrated relationship. Indeed,  
33 there are eruptions in these datasets that have a measured MER that differs  
34 by in excess of one order-of-magnitude from the power-law prediction. In  
35 applications, an order-of-magnitude under- or over-prediction of the MER  
36 could greatly limit the predictive ability of a dispersion model or inferences  
37 of changes in eruptive behaviour.

38 It is possible to obtain confidence intervals on the fitting parameters in  
39 the model relationship using linear regression of the log-transformed data  
40 (Mastin et al., 2009; Aubry et al., 2021, 2023). While this can be used to  
41 account for some of the scatter in the data, it does not fully account for the  
42 uncertainty in the observational data or in the model relationship used to  
43 describe it. There have recently been attempts to incorporate uncertainties  
44 into the eruption datasets (Aubry et al., 2021) and to consider model un-

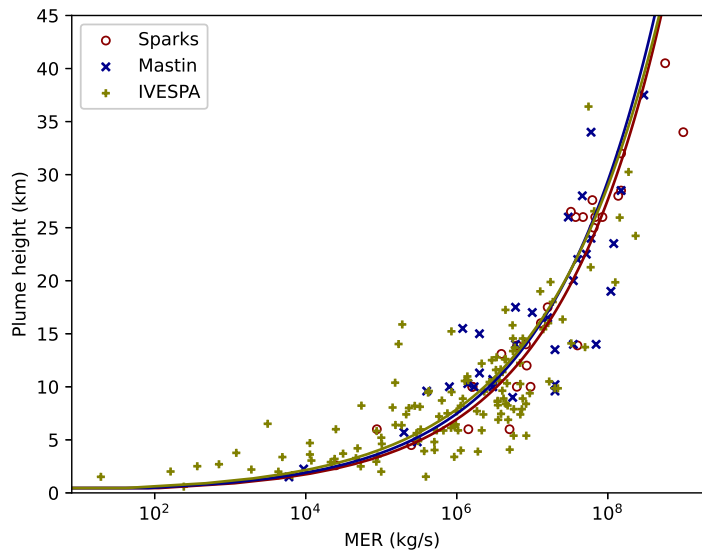


Figure 1: The eruption datasets of Sparks et al. (1997), Mastin et al. (2009) and IVESPA (Aubry et al., 2021), with corresponding power-law regression curves obtained from ordinary least squares linear regression of the log-transformed data.

45 certainties in the application of the relationship using stochastic modelling  
46 (Sparks et al., 2024). Here we present a structured approach using Bayesian  
47 methods to are widely used to quantify uncertainties,

48 Bayesian methods allow us to incorporate a range of uncertainties quan-  
49 titatively into our model, and provide a meaningful way to quantitatively  
50 compare different models (e.g. Gelman et al., 2014). A Bayesian calibration  
51 of a model allows us to make probabilistic predictions, and identify outliers  
52 in the data (and so determine, quantitatively, whether an eruption is un-  
53 usual). Importantly, with probabilistic information built into the predictive  
54 model, we are able to draw samples and build ensembles of plausible MER-  
55 plume height pairs (e.g. a set of plausible MERs corresponding to a height  
56 observation) each with an associated probability. This could allow for en-  
57 semble simulation of atmospheric tephra dispersion with inputs drawn in a  
58 structured way from distributions, facilitating the production of probabilistic  
59 forecasts.

60 In this paper we present a Bayesian approach, Bayesian linear regression,  
61 which is a standard method in statistical modelling (e.g. Gelman et al., 2014).  
62 This contribution illustrates and discusses the application to modelling the  
63 relationship between MER and plume heights using eruption source param-  
64 eter datasets. In our application of the methodology, the statistical model of  
65 the observations has a simple form, assuming the ‘errors’ in the data do not  
66 depend on the observations and are independent and identically distributed.  
67 This allows the calculations needed to calibrate model parameters and to  
68 perform prediction with the fitted model to be performed analytically. The  
69 resulting posterior predictive distribution is a well-known distribution (the

70  $t$ -distribution) and therefore subsequent computations can be performed us-  
71 ing standard numerical libraries. More complicated statistical models could  
72 be applied, and we discuss extensions of the statistical model proposed, but  
73 likely require the use of numerical methods (such as Markov Chain Monte  
74 Carlo) to approximate the posterior distributions of the model parameters  
75 and the posterior predictive distribution.

76 Our contribution is structured as follows. In §2 we introduce three erup-  
77 tion source parameter datasets considered here, and present our statistical  
78 model to relate MER and plume height. We present the analytical formulae  
79 for the posterior distributions of the model parameters and, most impor-  
80 tantly, the posterior predictive distribution. We illustrate the use of the  
81 statistical model in §3, considering inference of unobserved quantities, iden-  
82 tification of unusual events, and incorporation of measurement uncertainty.  
83 In §4 we discuss the limitations and extensions of the statistical model, and  
84 the application of the Bayesian approach to ensemble based modelling, before  
85 presenting our conclusion in §5.

86 The methods developed here have been implemented in the Python pack-  
87 age, MERPH, available from the PyPi package manager (<https://pypi.org/project/merph/>).  
88 This package includes the three datasets, as well as functionality to use alter-  
89 native data sets. Each of the results illustrated below can be easily computed  
90 using MERPH, and the package includes an interactive Jupyter notebook to  
91 illustrate the application.



## 92 2. Methods

### 93 2.1. Eruption source parameter datasets

94 We use the eruption source parameter datasets of Sparks et al. (1997),  
95 Mastin et al. (2009) and Aubry et al. (2021). The datasets do not all contain  
96 the same eruptions, and may distinguish phases of some eruptions. Further-  
97 more, the development over time has increased the number of eruption and  
98 atmospheric '*features*' that are recorded in the datasets. This could facilitate  
99 segregation of the data, but here we retain the complete datasets.

100 The dataset of Sparks et al. (1997) extends the catalogues of Settle (1978)  
101 and Wilson et al. (1978), with 28 eruptions at 18 volcanoes, recording the  
102 source volume flux, column height (above vent elevation) and (in 21 cases)  
103 the duration of the eruption (or phase of eruption). Sparks et al. (1997) notes  
104 substantial uncertainties implicit in both the column height and volume flux,  
105 but also that the studies from which the data derive typically do not quantify  
106 the uncertainties. In four eruptions a range of the volume flux is given, while  
107 11 have a range of values of the column height. In these cases, we adopt  
108 the mid-point of the range. Using a density of  $2500 \text{ kg/m}^3$ , a dense rock  
109 equivalent mass eruption rate is computed from the volume flux.

110 The power-law curve fit obtained from regression using the data of Mastin  
111 et al. (2009), commonly known as 'the Mastin curve', is frequently used in  
112 tephra dispersion modelling. The dataset contains 35 eruptions from 19 vol-  
113 canoes, spanning a period from 1902 to 2005. Mastin et al. (2009) present  
114 their dataset with eruptions separated into '*Silicic and andesitic eruptions*'  
115 (with 28 eruptions) and '*Basaltic eruptions*' (with seven eruptions). Each  
116 eruption has a plume height (above vent elevation), erupted volume, MER

117 and duration. Eight of the plume heights are presented as a range of values;  
118 in these cases we adopt the mid-point value. Four eruptions have a range  
119 for the MER, and we take the mid-point value. Additional features in the  
120 dataset include the magma type and the Volcanic Explosivity Index (VEI) of  
121 an eruption (which is cumulative for eruptions separated into phases). Fur-  
122 thermore, the plume height observations are associated with an observation  
123 method.

124 The IVESPA dataset (Aubry et al., 2021) has substantially increased the  
125 number of events (i.e., eruptions and phases of eruptions) to 137. These  
126 events correspond to 68 eruptions at 45 volcanoes (according to Global Vol-  
127 canism Program 2023). There are also many more features recorded than  
128 in the earlier datasets. For example, IVESPA includes three features that  
129 could be used to quantify the plume height: ‘*Tephra plume top height*’ (which  
130 we adopt here), ‘*Spreading height of the Umbrella Cloud*’, and ‘*SO<sub>2</sub> plume*  
131 *height*’. Aubry et al. (2023) discuss the differences in curve fits obtained  
132 when using these different features to represent the plume height. The plume  
133 heights in IVESPA are given above sea level, so for consistency with the  
134 Sparks et al. (1997) and Mastin et al. (2009) dataset and the basis of the  
135 power-law model, we convert to heights above the vent using vent elevations  
136 contained in IVESPA. A few events do not have a plume height recorded,  
137 and these must therefore be excluded from our analysis, resulting in a dataset  
138 with 130 events. The MER is not recorded in IVESPA, so here is computed  
139 from the ‘*Tephra Erupted Mass*’ and ‘*Duration*’ features. This means that  
140 the MER is an average over the duration of the eruption, which is consistent  
141 with the data in the Sparks et al. (1997) and Mastin et al. (2009) datasets.

142 IVESPA also includes estimates of the uncertainty in the quantities that  
143 could be included in a statistical analysis, but in this study we do not in-  
144 clude these uncertainties.

## 145 2.2. The statistical model

146 The leading-order behaviour in the observational data can be described  
147 by a power-law relationship, and dimensional reasoning supports the power-  
148 law dependence (Morton et al., 1956; Sparks et al., 1997). Therefore, we first  
149 make a logarithmic transformation of the data.

150 Here we present the method assuming that the plume height  $H$  is ob-  
151 served, and we seek to infer MER  $Q$ , as this is the practical use envisaged  
152 for emergency response. However, the method can be used with the roles of  
153 these variables switched (i.e.  $H$  as a function of  $Q$ ), which may be useful for  
154 preparatory modelling and risk analysis, with only changes in the numerical  
155 values that are computed and the interpretation of the results.

156 Taking logarithms of the data, we write  $x_i = \log H_i$  as the ‘explanatory’  
157 variable and  $y_i = \log Q_i$  as the ‘response’ variable, where  $(H_i, Q_i)$  is the pair  
158 of observed plume height (in km above the vent) and mass eruption rate (in  
159 kg/s) for eruption  $i$  in a historical record containing  $n$  eruptions. (Strictly,  
160 the plume height and MER should be non-dimensionalized before taking the  
161 logarithm, and so implicitly we have non-dimensionalized heights using a  
162 length scale of 1 km, and the MER by 1 kg/s.) The leading order power-  
163 law relationship between  $Q$  and  $H$  suggests  $\mathbb{E}(y_i) = \alpha + \beta x_i$ , where  $\mathbb{E}(\cdot)$   
164 denotes the expectation, so that our statistical model for the logarithmically  
165 transformed data is

$$y_i = \alpha + \beta x_i + \epsilon_i, \tag{1}$$

166 where  $\epsilon_i$  is the error in the observation of eruption  $i$  which includes contri-  
167 butions from aleatoric variations, measurement uncertainty and unmodelled  
168 process.

169 Examples of unmodelled processes are varying atmospheric conditions  
170 that are not accounted for in our simple relationship (e.g. wind, which is  
171 known to strongly influence the plume dynamics, e.g. Bursik, 2001; De-  
172 gruyter and Bonadonna, 2012; Woodhouse et al., 2013; Aubry et al., 2023),  
173 and volcanological parameters that are not explicitly included (e.g. physical  
174 properties of the magma, conduit geometry etc.) but which may alter the  
175 relationship. The measurement uncertainty includes both instrumental er-  
176 rors, recording errors, and errors in the derivation of the MER from tephra  
177 deposit volume and eruption duration.

178 As there are several contributions to the error term, and the eruptions  
179 in the databases are located across the world and span several decades, it is  
180 reasonable to assume that the errors for each event are independent. Further-  
181 more, we assume that the error does not depend on the size of the eruption or  
182 the plume height. Therefore, we take the errors to be independent and iden-  
183 tically distributed, with  $\epsilon_i \sim N(0, \sigma^2)$  where  $N(\mu, \sigma^2)$  denotes the Normal  
184 distribution with mean  $\mu$  and variance  $\sigma^2$ . The assumption of homoscedas-  
185 ticity (i.e., an equal error variance for each eruption) is discussed further  
186 below. We note that  $\sigma^2$  is unspecified and must be estimated as part of  
187 model calibration.

188 We are required to determine the three parameters in the statistical  
189 model, with  $\alpha \in \mathbb{R}$ ,  $\beta \in \mathbb{R}$  and  $\sigma^2 > 0$ . These parameters are estimated  
190 using Bayesian linear regression (Gelman et al., 2014). For ease of notation,

191 we define a  $(2 \times n)$  matrix  $X$  as

$$X^T = \begin{pmatrix} 1 & 1 & \dots & 1 \\ x_1 & x_2 & \dots & x_n \end{pmatrix} \quad (2)$$

192 and the vector of fit parameters  $\boldsymbol{\beta} = (\alpha, \beta)^T$ . The mean of  $\mathbf{x}$  is  $\bar{\mathbf{x}} = \frac{1}{n} \sum x_i$ ,  
 193 and  $\text{var}(\mathbf{x})$  denotes the variance of the  $\mathbf{x}$  data (and similarly for  $\mathbf{y}$ ).

194 *2.3. Maximum likelihood estimators*

195 Under the statistical model proposed, the likelihood is

$$\mathbf{y} \mid \boldsymbol{\beta}, \sigma^2, X \sim N(X\boldsymbol{\beta}, \sigma^2 I), \quad (3)$$

196 where  $I$  denotes the  $n \times n$  identity matrix. From this we obtain the maximum  
 197 likelihood estimators

$$\hat{\boldsymbol{\beta}} = VX^T \mathbf{y}, \quad \text{and} \quad \hat{\sigma}^2 = \frac{1}{n-2} (\mathbf{y} - X\hat{\boldsymbol{\beta}})^T (\mathbf{y} - X\hat{\boldsymbol{\beta}}), \quad (4)$$

198 where

$$V = (X^T X)^{-1} = \frac{1}{n^2 \text{var}(\mathbf{x})} \begin{pmatrix} \sum_{i=1}^n x_i^2 & -\sum_{i=1}^n x_i \\ -\sum_{i=1}^n x_i & n \end{pmatrix}. \quad (5)$$

Note, from (4) we find

$$\hat{\beta} = \frac{\text{cov}(\mathbf{x}, \mathbf{y})}{\text{var}(\mathbf{x})}, \quad (6)$$

$$\hat{\alpha} = \bar{\mathbf{y}} - \hat{\beta} \bar{\mathbf{x}}, \quad (7)$$

199 and  $\hat{\sigma}^2$  is the mean square error of the data. These are the same values of  
200 the ‘best fit’ model parameters that would be found from an ordinary linear  
201 regression of the log-transformed data.

#### 202 2.4. Probability distributions for model parameters

203 To incorporate the uncertainty in the observations and the uncertainty  
204 due to limitations of the model, we seek to obtain probability distributions  
205 on the parameters in the model. This is achieved through a Bayesian ap-  
206 proach, where prior probability distributions are assigned to the model vari-  
207 ables based on our beliefs about their values, and these prior probabilities  
208 are updated using Bayes’ theorem with the information contained in the set  
209 of observations. Here we take a standard non-informative prior assuming  $\boldsymbol{\beta}$   
210 and  $\sigma^2$  are independent, known as the Jeffrey’s prior, which has the form

$$p(\boldsymbol{\beta}, \sigma^2 | \mathbf{X}) \propto 1/\sigma^2. \quad (8)$$

211 If  $\sigma^2$  were known, then the posterior distribution of  $\boldsymbol{\beta}$  could be obtained.  
212 However, since  $\sigma^2$  is not known, instead we can write the conditional posterior  
213 distribution of  $\boldsymbol{\beta}$  given  $\sigma^2$ , which is

$$\boldsymbol{\beta} | \sigma^2, \mathbf{y}, \mathbf{X} = N(\hat{\boldsymbol{\beta}}, \mathbf{V}\sigma^2). \quad (9)$$

214 The marginal posterior of  $\sigma^2$  with the non-informative prior is given by

$$\sigma^2 | \mathbf{y}, \mathbf{X} = IG((n-2)/2, (n-2)\hat{\sigma}^2/2), \quad (10)$$

215 where  $IG$  is the inverse-Gamma distribution.

216 If needed, values of the model parameters can be obtained by drawing  
 217 from these distributions, i.e., first a value of the error term variance is drawn  
 218 from the inverse-Gamma distribution using (10), which can then be used in  
 219 (9) and values of  $\alpha$  and  $\beta$  obtained by draws from the multivariate Normal  
 220 distribution. This then gives alternative curve-fits that are feasible given the  
 221 data. A credible interval for the curve fit can be found from the posterior  
 222 distribution of  $y$ , which has the form of a non-standardized  $t$ -distribution,

$$y_j = \alpha + \beta x_j \sim t_{n-2} \left( \hat{\alpha} + \hat{\beta} x_j, \hat{\sigma}^2 \left[ \frac{1}{n} + \frac{(x_j - \bar{x})^2}{n \text{var}(\mathbf{x})} \right] \right). \quad (11)$$

223 Note, for this simple model, the credible interval is numerically the same as  
 224 a confidence interval for the regression line, although the interpretation is  
 225 different (see e.g. Lu et al., 2012).

226 These posterior distributions, and samples or statistics derived from them,  
 227 characterize the uncertainty in the statistical model, given the data. How-  
 228 ever, they are typically not needed in the practical applications. Instead, we  
 229 wish to use the uncertainty in the model parameters, now quantified in the  
 230 posterior distributions, to provide probabilistic predictions of the response  
 231 variable when a new observation of the explanatory variable is made.

### 232 2.5. Posterior prediction

233 The posterior distribution of the model parameters can be used to draw  
 234 sample sets of model parameters that are consistent with the data. If, during  
 235 an eruption, a new observation of the plume height is made, and therefore a  
 236 new explanatory variable  $\tilde{x}$  is given, we can use the posterior distribution of  
 237 the model parameters to predict a distribution of values for the MER that

238 are consistent with the new observation, the data underlying the curve-fit,  
 239 and the uncertainties in the data and the model.

240 Under the statistical model, if the model parameters were known precisely,  
 241 then we have

$$\tilde{y} \sim N(\alpha + \beta\tilde{x}, \sigma^2). \quad (12)$$

242 The uncertainty in the observations and the model is incorporated through  
 243 the posterior distributions of the model parameters  $\alpha$ ,  $\beta$  and  $\sigma^2$ . When in-  
 244 cluding these uncertainties, we arrive at the posterior predictive distribution,  
 245  $p(\tilde{y} | \tilde{x}, \mathbf{y}, \mathbf{X})$ . This can be obtained by simulation, by drawing many sam-  
 246 ples from the posterior predictive distribution by first sampling the model  
 247 parameters from their posterior distributions, and then using these as fixed  
 248 values in equation (12). However, for our simple model and choice of non-  
 249 informative prior, the posterior distribution can be written analytically (see  
 250 Gelman et al., 2014, for details) as

$$\tilde{y} | \tilde{x}, \mathbf{y}, \mathbf{x} \sim t_{n-2} \left( \hat{\alpha} + \hat{\beta}\tilde{x}, \hat{\sigma}^2 \left[ 1 + \frac{1}{n} + \frac{(\tilde{x} - \bar{x})^2}{n\text{var}(\mathbf{x})} \right] \right), \quad (13)$$

251 a  $t$ -distribution with  $n - 2$  degrees-of-freedom, with location at the log-MER  
 252 predicted by the maximum likelihood estimate at the observed plume height,  
 253 and a scale parameter that incorporates the estimated measurement uncer-  
 254 tainty through the maximum likelihood estimate of the error variance. Note,  
 255 the scale parameter for the posterior predictive distribution is larger than the  
 256 scale parameter for the posterior distribution for the curve fit, equation (11),  
 257 which accounts for the uncertainty of a new observation around the curve fit.

258 The logarithmic transformation can be inverted, so that the posterior



259 predictive distribution for the MER given a new plume height observation  
260 and a historical dataset is a log- $t$  distribution.

261 If multiple new height observations are made, so we have a vector of new  
262 explanatory variables  $\tilde{\mathbf{x}}$ , then the posterior predictive distribution of  $\tilde{\mathbf{y}}$  is  
263 found as a multivariate  $t$ -distribution, with

$$\tilde{\mathbf{y}} \mid \tilde{\mathbf{x}}, \mathbf{y}, \mathbf{x} \sim t_{n-2} \left( \tilde{\mathbf{X}} \hat{\boldsymbol{\beta}}, \hat{\sigma}^2 \left[ \mathbf{I} + \tilde{\mathbf{X}} \mathbf{V} \tilde{\mathbf{X}}^T \right] \right). \quad (14)$$

264 However, the typical case will be a single new observation.

265 We note that the analytical form of the posterior predictive distribution  
266 is a major advantage, as it allows samples to be drawn from the distribution  
267 very easily. The  $t$ -distribution is commonly used in statistical analysis, so  
268 there are many software packages that provide algorithms for computing  
269 quantities (such as the CDF and quantiles) from the  $t$ -distribution (e.g.,  
270 the `scipy` package in Python, the Statistics and Machine Learning Toolbox  
271 in MatLab, and in R). Furthermore, the location and scale parameters in  
272 the  $t$ -distribution can be computed directly from the maximum likelihood  
273 estimates for the model parameters and other quantities of the data.

## 274 2.6. Including observational uncertainty in predictions

275 The posterior predictive distribution in (13) assumes that the observation  
276 is exact. In applications where a plume height is observed, there is likely to  
277 be an associated uncertainty. We can characterize this uncertainty through a  
278 probability distribution for the new observation,  $p(\tilde{x})$ , and the joint posterior

279 predictive distribution can be computed using

$$p(\tilde{x}, \tilde{y} | \mathbf{x}, \mathbf{y}) = p(\tilde{y} | \tilde{x}, \mathbf{x}, \mathbf{y}) p(\tilde{x}) \quad (15)$$

280 which allows samples to be drawn using equation (13).

### 281 **3. Results**

282 Here we illustrate the application of the Bayesian regression using the  
283 datasets of Sparks et al. (1997), Mastin et al. (2009) and IVESPA (Aubry  
284 et al., 2021).

#### 285 *3.1. Maximum likelihood estimators and posterior distributions for model pa-* 286 *rameters*

287 The maximum likelihood estimators of the model parameters, computed  
288 from (4) for each dataset, are tabulated in table 1. These give the ‘best-fit’  
289 curves for each dataset, which we refer to as the ‘Sparks curve’, ‘Mastin curve’  
290 and ‘IVESPA curve’, for the Sparks et al. (1997), Mastin et al. (2009) and  
291 IVESPA (Aubry et al., 2021) datasets, respectively. The power-law exponent,  
292  $\hat{\beta}$ , differs only slightly for each of the three datasets (table 1), but there are  
293 more substantial differences between the maximum likelihood estimates of  
294 the pre-factor  $\hat{\alpha}$ . However, these differences do not strongly alter the curves  
295 (figure 1).

296 It is noteworthy that the magnitude of the uncertainty, encapsulated in  
297 the variance of the normally-distributed error term,  $\hat{\sigma}^2$ , increases markedly  
298 as the size of the dataset increases (table 1). This indicates that the in-  
299 creasing size of the historical record is capturing a larger range of eruption

300 and atmospheric conditions, so the scatter of data points around the best-fit  
 301 curve increases, as seen in figure 1.

Dataset	$n$	$\hat{\alpha}$	$10^{\hat{\alpha}}$	$\hat{\beta}$	$\hat{\sigma}^2$
Sparks	28	2.99	981.84	3.47	0.11
Mastin	35	3.07	1164.24	3.36	0.22
IVESPA	130	2.83	668.99	3.54	0.59

Table 1: Maximum likelihood estimates for the model parameters  $\hat{\alpha}$  and  $\hat{\beta}$  for the model  $Q = 10^\alpha H^\beta$ , found from log-transformation of the data of Sparks et al. (1997), Mastin et al. (2009) and IVESPA (Aubry et al., 2021). The size of the dataset,  $n$ , is also given.

302 The posterior distributions for the model parameters are given in equa-  
 303 tions (9) and (10) and can be used to sample curves that fit the data. In  
 304 figure 2 the curves obtained from 100 random samples from the posterior  
 305 distributions of the model parameters using the Mastin et al. (2009) data are  
 306 shown together with data, the Mastin curve obtained from maximum like-  
 307 lihood estimation and the 95% credible interval of the best fit curve. Note  
 308 the 95% credible interval indicates the region where we have high belief in  
 309 the curve fit; there is a probability of 0.95 that the 'true' best-fit curve lies  
 310 within the credible interval.

### 311 3.2. Posterior prediction

312 In applications, the posterior predictive distribution is used to determine  
 313 estimates of the MER that are consistent with a new height observation  
 314 and the data underlying the model. Table 2 presents values of the location  
 315 and scale parameters in the posterior prediction distribution for the log-  
 316 MER given a plume height observation in the range 5–50 km for the Sparks  
 317 et al. (1997), Mastin et al. (2009) and IVESPA (Aubry et al., 2021) datasets.  
 318 The location parameter values (given by  $\mu = \hat{\alpha} + \hat{\beta}\tilde{x}$ , which quantifies the

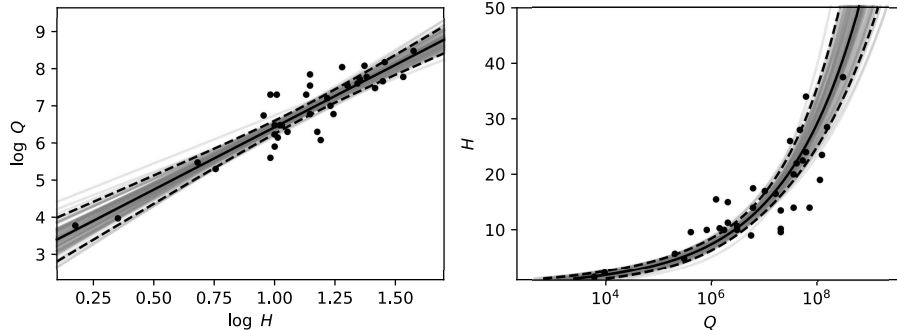


Figure 2: Possible curve fits to the Mastin et al. (2009) dataset obtained from sampling the posterior distributions of the model parameters. The left-hand panel shows the log-transformed data. The right-hand panel show the MER–plume height data as commonly plotted. The black solid lines show the ‘best-fit’ curve obtained from maximum likelihood estimation of the model parameters, with the black dashed lines showing the 95% confidence intervals on the curve fit. The semi-transparent grey lines show 100 alternative curve fits from sampling of the posterior distributions of the model parameters.

319 most likely log-MER) at each height are similar for each dataset, but the  
 320 scale parameter (quantifying the variability around the maximum likelihood  
 321 value) is substantially larger for the IVESPA dataset. For each dataset, the  
 322 scale parameter varies only slightly over the large range of heights, increasing  
 323 gradually with increasing distance of the observed log-plume height from the  
 324 mean of the data,  $|\tilde{x} - \bar{x}|$ .

325 Figure 3 shows the posterior predictive distribution for the log-MER for  
 326 three plume height observations, for each of the Sparks et al. (1997), Mastin  
 327 et al. (2009) and IVESPA (Aubry et al., 2021) datasets. Here the larger  
 328 variability in the IVESPA dataset is apparent. We also note that, for each  
 329 dataset, the posterior predictive distributions for different height observations  
 330 overlap. For example, using the Sparks et al. (1997) data, the most likely  
 331 values of the MER for plume heights of 5, 10 and 20 km differ by a full decade

$\tilde{H}$ (km)	$\tilde{x}$	Sparks		Mastin		IVESPA	
		$\mu$	$\tau$	$\mu$	$\tau$	$\mu$	$\tau$
5	0.699	5.42	0.358	5.41	0.494	5.30	0.772
10	1.00	6.47	0.344	6.43	0.482	6.37	0.771
15	1.18	7.08	0.342	7.02	0.481	6.99	0.774
20	1.30	7.51	0.343	7.44	0.483	7.43	0.777
25	1.40	7.85	0.345	7.76	0.487	7.78	0.780
30	1.48	8.12	0.348	8.03	0.491	8.06	0.783
35	1.54	8.36	0.351	8.25	0.495	8.29	0.786
40	1.60	8.56	0.354	8.45	0.499	8.50	0.788
45	1.65	8.74	0.357	8.62	0.503	8.68	0.791
50	1.70	8.90	0.361	8.78	0.506	8.84	0.793

Table 2: Location ( $\mu$ ) and scale ( $\tau$ ) parameters in the posterior predictive distribution of the logarithm of the MER for specified plume heights,  $\tilde{y} \sim t_{32}(\mu, \tau^2)$ , using data from Sparks et al. (1997), Mastin et al. (2009) and IVESPA (Aubry et al., 2021).

332 (with  $\log Q$  taking values of 5.42, 6.47 and 7.51, respectively), but from the  
333 posterior predictive distribution at these heights we find

$$\begin{aligned}
& P(5.90 < \log Q < 6.00 | H = 5, \text{Sparks}) \\
& \approx P(5.90 < \log Q < 6.00 | H = 10, \text{Sparks}) = 0.037
\end{aligned} \tag{16}$$

334 so that if we observe a plume at height 5 km, there is a probability of 3.7%  
335 that  $\log Q \approx 5.95$ , and the same MER is found with equal probability for a  
336 plume at height 10 km. The larger variability in the IVESPA data results  
337 in more pronounced overlapping of the posterior predictive distributions for  
338 different height observations. In this case,

$$\begin{aligned}
& P(6.32 < \log Q < 6.42 | H = 5, \text{IVESPA}) \\
& \approx P(6.32 < \log Q < 6.42 | H = 20, \text{IVESPA}) = 0.022
\end{aligned} \tag{17}$$

339 so there are appreciable probabilities of a MER with  $\log Q \approx 6.37$  for plumes  
 340 reaching 5 km or 20 km (noting this value for the MER is the most likely for  
 341 plume reaching 15 km). Note this should not be interpreted to mean that  
 342 there are equal probabilities (under the IVESPA model) for a plume height  
 343 of 5 km and 20 km for a MER of  $10^{6.37} = 2.34 \times 10^6$  kg/s; to determine  
 344 these probabilities we need the posterior predictive distribution of the plume  
 345 heights *given a MER*, from which we find that the most likely plume height  
 346 is 9 km.

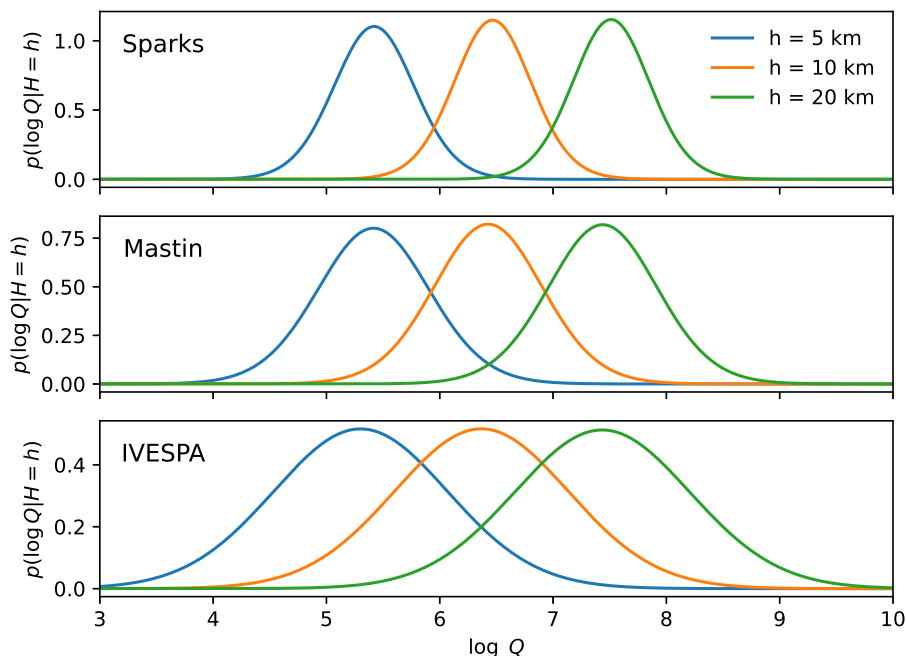


Figure 3: Posterior predictive densities,  $p(\log Q | \log H = \log h, \mathbf{y}, \mathbf{x})$ , for the Sparks et al. (1997), Mastin et al. (2009), and IVESPA (Aubry et al., 2021) datasets, for three plume heights observations:  $h = 5$ km, 10km, and 20km.

347 The posterior predictive distribution can be found for any new height ob-

348 servation. Figure 4 illustrates posterior prediction intervals using the Mastin  
 349 et al. (2009) dataset for  $\log Q | \log H$  for a range of plume heights, together  
 350 with the data, the Mastin curve, and the 95% credible interval for the curve  
 351 fit. The  $(1 - \alpha) \times 100\%$  posterior prediction interval is a centred interval  
 352 satisfying  $P(y_l < Y < y_u | X) = 1 - \alpha$  for a specified  $0 < \alpha < 1$ . Decreasing  
 353  $\alpha$  produces wider bands, indicating a greater probability that the true MER  
 354 lies within the prediction interval.

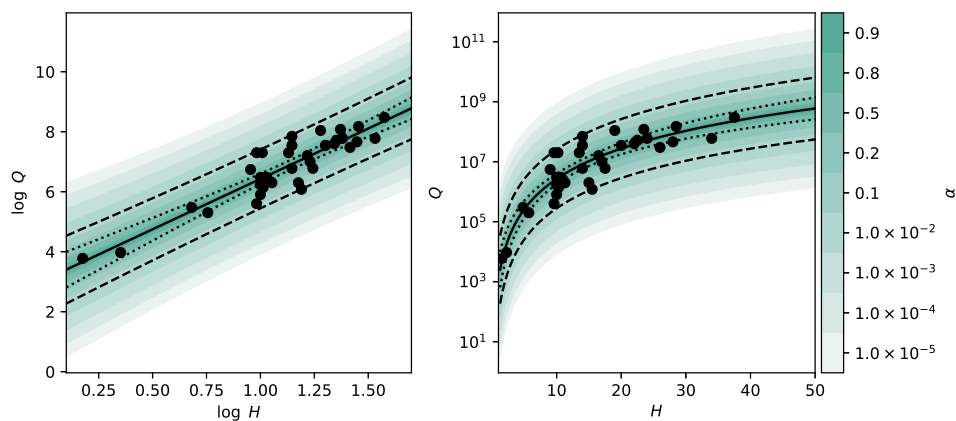


Figure 4: Posterior prediction intervals of the MER given a plume height observation for the Mastin et al. (2009) dataset (data indicated by black circles). The Mastin curve is indicated by the solid black line, with the 95% credible interval on the curve fit shown with dotted lines. The overlapping coloured bands indicate centred prediction intervals (i.e. curves  $y_l(x)$  and  $y_u(x)$  such that  $P(y_l < Y < y_u | X) = 1 - \alpha$ ). The prediction interval for  $1 - \alpha = 0.95$  is shown with dashed lines.

355 Figure 4 shows that the 95% predictive interval is much wider than the  
 356 95% credible interval for the curve fit, which we recall is numerically identical  
 357 to a 95% confidence interval for the curve fit. Therefore, confidence intervals  
 358 on the curve fit by themselves do not adequately capture the uncertainty  
 359 when making predictions. The predictive distribution, with its greater vari-

360 ance, must be used when sampling plausible MER values based on plume  
361 height observations.

362 The contours of the intervals (in log-space) are slightly narrower where  
363 there are more data, although this is less apparent in the for the predictive  
364 intervals than for the credible interval on the curve fit (Figure 4). The illus-  
365 trates that we have a stronger belief that the calibrated curve well represents  
366 the data in regions where the data is clustered, and less belief where there is  
367 sparse data. In contrast, the predictive interval captures the scatter in the  
368 data, so while there is most data for plumes are 11 km, there is consider-  
369 able scatter in the MER for these eruptions, so wide prediction intervals are  
370 needed.

### 371 *3.2.1. Identifying unusual eruptions*

372 The posterior predictive distribution allows us to identify unusual events  
373 quantitatively. Methods have been developed to identify outliers within a  
374 dataset (e.g. Chaloner and Brant, 1988). Here we seek to determine whether  
375 a new eruption, not within a dataset but with a MER and plume height obser-  
376 vation, is unusual. To illustrate this, we select eruptions from the IVESPA  
377 dataset (Aubry et al., 2021) that are not contained in the Mastin et al.  
378 (2009) dataset, and use the posterior predictive distributions derived from  
379 the Mastin et al. (2009) data to characterize these new events.

380 There are 29 eruptions in the IVESPA dataset that occurred since the  
381 publication of the Mastin et al. (2009) dataset. For each of these eruptions,  
382 we use the observations to produce posterior predictive distributions for both  
383 MER and plume height and compute the probability  $P(\log 0.8 + y^* \leq Y \leq$   
384  $\log 1.2 + y^* | x^*, \mathbf{x}, \mathbf{y})$  for each, where  $x^*$  and  $y^*$  denote the observations (either



385  $\log Q$  or  $\log H$ ) from IVESPA, i.e., we determine the probability that the  
 386 posterior prediction lies within  $\pm 20\%$  of the recorded observation. Table 3  
 387 shows the four probabilities for the 2015 eruption of Cotopaxi are notable for the low  
 388 probabilities associated with the observed MER and plume heights, which  
 389 suggest the eruption is unusual with respect to the Mastin et al. (2009)  
 390 dataset. Indeed, for an observed plume height of 6.513 km, the Mastin  
 391 regression would predict  $P(Q \leq 3.14 \times 10^3) \approx 2 \times 10^{-5}$ , and a MER more  
 392 than order-of-magnitude larger than observed is required to reach the 1-  
 393 percentile of the posterior predictive distribution. While the probabilities  
 394 associated with predictions of the plume heights given the MER are larger,  
 395 they remain relatively small.

397 The eruption of Etna, 21 May 2016, is also of interest, as the MER ob-  
 398 servation given the plume height might be considered unusual, with  $P(Q <$   
 399  $7.1 \times 10^2 | H = 2.7) \approx 0.001$ , but the plume height observation given the MER  
 400 is not unlikely (the observed plume height of 4.7 km is less than 2 standard  
 401 deviations from the mean of the posterior predictive distribution). This ex-  
 402 ample illustrates the difference between the posterior predictive distributions  
 403 for  $\log Q | \log H$  and  $\log H | \log Q$ .

### 404 3.2.2. Joint predictive samples

405 Uncertainty in the observation can also be included, by sampling from  
 406 the joint posterior predictive distribution using the decomposition given by  
 407 equation (15). As an example, we consider here uncertain plume height ob-  
 408 servations, so that  $H \in [h_0, h_1]$ , and for illustration take a large range of  
 409 possible values,  $h_0 = 5$  km and  $h_1 = 20$  km. We consider four distributions  
 410 to characterize the uncertainty (referred to below as the ‘measurement dis-

Eruption	$Q^*$ (kg/s)	$H^*$ (km)	$P_{Q^* H^*}$	$P_{H^* Q^*}$
Calbuco 22 April 2015	$1.870 \times 10^7$	17.997	0.144	0.493
Calbuco 23 April 2015	$1.275 \times 10^7$	18.997	0.123	0.417
Puyehue-Cordón Caulle layer A-F, 2011	$4.630 \times 10^6$	10.330	0.134	0.433
Puyehue-Cordón Caulle layer H, 2011	$4.012 \times 10^6$	9.630	0.129	0.406
Puyehue-Cordón Caulle layer K2, 2011	$8.408 \times 10^5$	6.130	0.130	0.344
<b>Cotopaxi 1st phase, 2015</b>	$3.135 \times 10^3$	6.513	<b>0.000</b>	<b>0.004</b>
<b>Cotopaxi 2nd phase, 2015</b>	$3.664 \times 10^2$	2.513	<b>0.001</b>	0.071
<b>Cotopaxi 3rd phase, 2015</b>	$1.624 \times 10^2$	2.013	<b>0.001</b>	0.084
<b>Cotopaxi 4th phase, 2015</b>	$1.939 \times 10^1$	1.513	<b>0.000</b>	0.033
<b>Etna 12 January 2011</b>	$2.510 \times 10^4$	6.000	<b>0.006</b>	0.124
Etna 23 February 2013	$5.051 \times 10^5$	4.800	0.110	0.235
<b>Etna 26 October 2013</b>	$1.136 \times 10^4$	4.700	<b>0.006</b>	0.159
Etna 23 November 2013	$5.824 \times 10^5$	7.200	0.132	0.492
Etna 18 & 19 May 2016	$1.907 \times 10^3$	2.200	0.027	0.391
<b>Etna 21 May 2016</b>	$7.104 \times 10^2$	2.700	<b>0.001</b>	0.104
Eyjafjallajökull 14-16 April 2010	$5.086 \times 10^5$	4.040	0.069	0.108
Eyjafjallajökull 17 April 2010	$3.704 \times 10^5$	3.940	0.085	0.141
Eyjafjallajökull 18 April - 21 May 2010	$8.617 \times 10^4$	2.940	0.116	0.190
Eyjafjallajökull 4-8 May 2010	$8.951 \times 10^4$	3.440	0.136	0.314
Grímsvötn 21 May 2011	$7.500 \times 10^6$	14.550	0.141	0.492
Kelut Unit B, 2014	$5.938 \times 10^7$	21.269	0.127	0.483
<b>Merapi 4 November 2010</b>	$1.698 \times 10^5$	14.032	<b>0.001</b>	0.011
Nakadake - Asosan 14 September 2015	$1.000 \times 10^5$	1.998	0.031	0.017
Ontakesan 27 September 2014	$4.630 \times 10^4$	3.275	0.131	0.402
Shinmoedake - Kirishimayaya Phase SP1+SP2, 2011	$6.667 \times 10^5$	5.879	0.134	0.357
Shinmoedake - Kirishimayaya Phase SP3, 2011	$3.333 \times 10^5$	5.979	0.134	0.477
Tungurahua 14 July 2013	$1.867 \times 10^5$	6.377	0.083	0.450
Tungurahua 1 February 2014	$3.488 \times 10^6$	8.677	0.116	0.345
Villarrica 3 March 2015	$9.392 \times 10^5$	6.253	0.128	0.337

Table 3: Probabilities of eruption source parameters in the IVESPA dataset (Aubry et al., 2021) using the regression model for Mastin et al. (2009). The quantity  $P_{Q^*|H^*}$  gives the probability  $P(\log(0.8Q^*) < \log Q < \log(1.2Q^*) | H^*)$  and similarly for  $P_{H^*|Q^*}$ . Bold entries indicate eruptions where  $P_{Q^*|H^*} < 0.01$  and/or  $P_{H^*|Q^*} < 0.01$ .

411 tributions' of the plume height): a uniform distribution,  $H \sim U(h_0, h_1)$ , and  
412 truncated normal distributions,  $H \sim TN(\mu, \omega^2, h_0, h_1)$  with standard devi-

413 ation ( $\omega = 3$  km) taking mean at the centre of the interval ( $\mu = 12.5$  km),  
 414 a low-skewed mean ( $\mu = 10$  km), and a high-skewed mean ( $\mu = 15$  km).  
 415 Figure 5 shows these probability distributions for the uncertain height ob-  
 416 servation together with the associated joint posterior predictive distributions  
 417  $f_{Q,H}(q, h)$ . Note, these distributions are given for the MER and plume height  
 418 and not the logarithmically transformed variables, obtained by scaling equa-  
 419 tion (15), and this results in small values of the probability density due to  
 420 the large range of values for the MER.

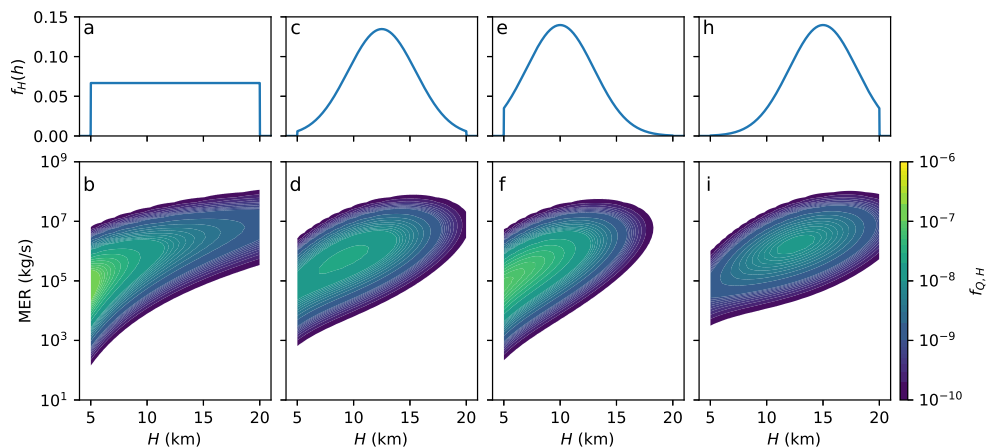


Figure 5: Joint posterior predictive distributions for the MER and plume height, based on uncertain plume height observations using the Mastin et al. (2009) dataset. The upper panels show the probability density function of the height observation,  $f_H(h)$ , and the lower panels show the associated joint posterior density distributions,  $f_{Q,H}(q, h)$ . (a,b) Uniform distribution,  $H \sim U(5, 20)$ . (c,d) Symmetric truncated normal distribution,  $H \sim TN(12.5, 9, 5, 20)$ . (e,f) Low-skewed truncated normal distribution,  $H \sim TN(10, 9, 5, 20)$ . (h,i) High-skewed truncated normal distribution,  $H \sim TN(15, 9, 5, 20)$ .

421 Each of the measurement distributions have finite domain, so curtail the  
 422 joint distribution to this range of plume heights. For the uniform distribu-  
 423 tion (Figure 5a,b) the contours of the joint distribution resembles those of

424 the prediction intervals of Figure 4. However, there is a noticeable asymme-  
425 try, with broader tails of probability density for small MER than large MER,  
426 and higher density for small plume heights than large. This is because the  
427 measurement distribution is not uniform in log-space and the posterior pre-  
428 dictive distribution is not symmetric in linear space. The truncated normal  
429 distributions similarly do not produce symmetric joint distributions. The  
430 location of the maximum joint density moves with the maximum of the ob-  
431 served distribution, but is always offset to lower plume heights.

## 432 4. Discussion

### 433 4.1. Advantages, limitations and extensions

434 Bayesian linear regression of the MER and plume height in eruption  
435 databases provides a valuable methodology to interpret observations and  
436 to predict future eruption conditions. The model proposed here is arguably  
437 the simplest statistical model, but has some notable advantages. Firstly, the  
438 model produces an analytical result in the form of a well-known distribu-  
439 tion (the  $t$ -distribution), so calculations can be performed easily. Secondly,  
440 the inference of model parameters from a dataset is straight-forward. Indeed,  
441 the model parameters of the posterior distribution are functions of quantities  
442 routinely computed in ordinary linear regression, consisting of the curve-fit  
443 parameters (slope and intercept) and the mean square error of the data.  
444 However, the simple approach has some limitations.

445 A key assumption of our statistical model is homoscedasticity (equal error  
446 variance for all observations). While this assumption allows us to obtain the  
447 analytical results, it also causes the inferred variance to be large in datasets

448 where there is increased scatter. For example, in the IVESPA dataset (Aubry  
449 et al., 2021), there is substantially greater scatter of the observations around  
450 the MLE compared to the smaller Sparks et al. (1997) and Mastin et al.  
451 (2009) datasets (see figure 1 and table 1) so the posterior distribution for the  
452 error is wider in order to capture the observations using a normal distribution.

453 A heteroscedastic model, with an observation-specific error variance, may  
454 give improved fit to the observed data. Indeed, based on physical principles,  
455 we may wish to link the error variance to the explanatory variables. For  
456 example, transient weather conditions give highly variable wind and temper-  
457 ature in the troposphere, which can strongly impact plume dynamics (Bursik,  
458 2001; Degruyter and Bonadonna, 2012; Woodhouse et al., 2013), and there-  
459 fore greater aleatoric uncertainty for tropospheric plumes. Therefore, it may  
460 be appropriate to develop a statistical model allowing for larger error vari-  
461 ance at low altitudes. This could be achieved either by using a functional  
462 relationship for the error variance, i.e., letting  $\sigma_i = f(H_i)$ , or by grouping  
463 events in the datasets into tropospheric and stratospheric eruptions. They  
464 may be additional grouping that could be applied, for example considering  
465 tropic, mid-latitude and high-latitude eruptions to assess the effect of erup-  
466 tion location.

467 More sophisticated statistical models could be applied, and there are  
468 likely to be benefits to this, particularly as eruption datasets grow, with  
469 more eruptions and more variables recorded. This would allow other con-  
470 trolling variables to be included in the regression analysis, to reduce epis-  
471 temic uncertainties. For example, atmospheric effects could be included by  
472 incorporating a dependence of the plume height on the atmospheric strat-

473 ification (quantified through the buoyancy frequency,  $N$ ) and/or the wind  
474 speed; these variables and several others are included in the IVESPA dataset  
475 (Aubry et al., 2021).

476 Considering first the buoyancy frequency only, physical reasoning sug-  
477 gests the plume height scales as  $H \sim N^{-3/4}Q^{1/4}$  (Wilson et al., 1978; Settle,  
478 1978; Woods, 1988; Sparks et al., 1997; Degruyter and Bonadonna, 2012;  
479 Woodhouse et al., 2013). We can therefore retain a linear model in logarith-  
480 mically transformed variables, albeit with in a model with two explanatory  
481 variables, so three fitting coefficients. The Bayesian linear regression can  
482 then be applied straight-forwardly (see Gelman et al., 2014).

483 In contrast, when modelling for the effect of wind speed,  $V$ , in inte-  
484 gral plume models, Degruyter and Bonadonna (2012) and Woodhouse et al.  
485 (2013) suggest there is not a simple power-law dependence. Instead, the  
486 plume height scales as  $H \sim N^{-3/4}Q^{1/4}f(\mathcal{W})$ , where  $\mathcal{W} = V/(HN)$  is a di-  
487 mensionless wind speed, and where  $f$  is a decreasing function of  $\mathcal{W}$ . The  
488 functional forms proposed by Degruyter and Bonadonna (2012) and Wood-  
489 house et al. (2013) differ, and we could base a new statistical model on these  
490 forms. However, in neither case do not obtain a linear model in logarithmic  
491 space, so the linear regression cannot be used. However, a nonlinear model  
492 could be used, with Markov Chain Monte Carlo methods used to fit and pre-  
493 dict. This approach would also allow different functional forms to be tested  
494 and quantitatively compared.

495 The extension to model complex models may require computation. There  
496 are now several advanced toolkits for Probabilistic Programming (e.g., Abril-  
497 Pla et al., 2023; Stan Development Team, 2023) that provide easy-to-use

498 interfaces to advanced computation methods for Bayesian inference. These  
499 methods allow non-specialists to implement models by specifying prior and  
500 likelihood functions directly as probability density functions linked to the  
501 data, and perform Markov Chain Monte Carlo integration to fit the model  
502 (i.e., numerically approximate the posterior distribution) and then make  
503 probabilistic predictions. Additionally, non-parametric approaches, such as  
504 Gaussian Process regression (e.g., Rasmussen and Williams, 2006) could be  
505 used to create models that do not rely on a pre-specified form for the relation-  
506 ship. This provides a more ‘data-driven’ approach by removing the need for  
507 an initial physics-based model on which to base the inference, but care must  
508 be taken to ensure physical laws are not violated by the resulting model.

#### 509 *4.2. Applications to tephra dispersion modelling*

510 Estimates of MER and plume height are important inputs for tephra  
511 hazard simulations. For example, operational ash dispersion forecasts require  
512 as input the MER and plume height (Folch, 2012), and typically the plume  
513 height is observed or imposed and a consistent MER estimate is required (e.g.  
514 Folch, 2012; Jenkins et al., 2015; Beckett et al., 2020). In probabilistic tephra  
515 modelling, scenarios are often based on eruption magnitude (e.g. Bonadonna  
516 et al., 2005; Bear-Crozier et al., 2016; Tadini et al., 2022) from which an  
517 MER is imposed and a consistent plume height is determined.

518 In many cases the imposed explanatory variable is given without uncer-  
519 tainty, and the estimate of the response variable is typically derived directly  
520 from the best-fit curve. This approach neglects both uncertainties in the  
521 ‘measured’ variable and the uncertainty in the response variable due to the  
522 use of the observational dataset. Accounting for these uncertainties is impor-

523 tant to ensure that results are not biased by application of the best-fit result.  
524 Tephra dispersion patterns are strongly influenced by the plume height (De-  
525 venish et al., 2012; Dioguardi et al., 2020; Pardini et al., 2022), particularly  
526 where there is significant atmospheric wind shear with altitude, while the  
527 concentrations of airborne tephra, or ground level loadings, depend on the  
528 MER.

529 The Bayesian linear regression used here provides a method to compute  
530 values of the response variable that are consistent with the imposed value  
531 of the explanatory variable, and associate probabilities with these  $(Q, H)$   
532 pairs, by drawing samples from the posterior predictive distribution. As  
533 the statistical model provides an analytical posterior predictive distribution,  
534 we can readily draw random samples directly. Alternatively, we can use a  
535 structured sampling design, for example by drawing values of the response  
536 variable at specified percentiles, which can ensure that unusual  $(Q, H)$  pairs  
537 are included in an ensemble, with knowledge of their probability. Accessing  
538 the tails of the distribution is likely to be important to ensure that ‘rare’  
539 events are included in ensembles. Knowledge of the probabilities of inputs  
540 allows ensemble members to be weighted appropriately when aggregating to  
541 create probabilistic outputs, reducing the need for large ensembles. These  
542 approaches allow quantitative uncertainty to be included in dispersion mod-  
543 elling.

544 In forecasting applications, the uncertainty in both the observed plume  
545 height and inferred MER can be quantified using the Bayesian linear regres-  
546 sion. Specifically, imposing a proper probability distribution for the plume  
547 height observation to represent its measurement uncertainty, and sampling



548 from the joint posterior predictive distribution allows a set of  $(Q, H)$  pairs  
549 to be constructed as inputs for an ensemble of dispersion simulations. This  
550 could be combined with meteorological and other eruption source parameter  
551 uncertainty (Beckett, Witham, Hort, Stevenson, Bonadonna and Millington,  
552 2015; Osman, Beckett, Rust and Snee, 2020, e.g.) to create probabilistic  
553 airborne ash dispersion forecasts. In Williams et al. (2024) we show how the  
554 analytical form of the posterior predictive distribution can be used to effi-  
555 ciently generate probabilistic volcanic ash hazard forecasts that incorporate  
556 uncertainty in eruption source parameters and meteorological fields.

## 557 **5. Conclusion**

558 Eruption source parameter datasets are valuable catalogues of past erup-  
559 tions. Relationships derived from these datasets provide useful tools that can  
560 be rapidly deployed during response to eruption, or used to inform prepared-  
561 ness. However, the aleatoric and epistemic uncertainties captured within the  
562 datasets propagate through curve fits into methods that adopt these expres-  
563 sions. By applying Bayesian linear regression, these uncertainties can be  
564 quantified, providing new capability for probabilistic approaches that adopt  
565 these relationships.

566 The expansion of eruption source parameter dataset provides new insights  
567 into the controls on plume dynamics (Aubry et al., 2021, 2023). With in-  
568 creasing numbers of events in the catalogue, more of nature’s variations are  
569 recorded, so it is not surprising that we observe increased scatter around the  
570 simplest MER–plume height relationship. Including additional explanatory  
571 variables may improve a model’s predictive capability by reducing epistemic

572 uncertainty, but this is tensioned by the increased demand to measure and  
573 specify variables when applying models. The increased complexity of multi-  
574 variable modelling must be considered carefully against the improvement in  
575 prediction. In this regard Bayesian approaches are particularly useful, with  
576 quantitative methods for comparing competing models.

577 Analysis of the eruption source parameter datasets to derive simple re-  
578 lationships facilitates rapid response hazard modelling by allowing variables  
579 that are difficult to determine to be inferred from easily measure quantities.  
580 Bayesian approaches are likely to provide useful tools for this analysis, but  
581 providing meaningful and useful uncertainty quantification.

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591 tol's impact acceleration account.

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