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# A Review of the Migration of Hydrogen from the Planetary to Basin Scale

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# <sup>14</sup> Key Points:

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15	•	History of planetary formation and overview of natural hydrogen cycle of Earth.
16	•	Summary of mechanisms for hydrogen migration (diffusion and advection), gen-
17		eration and consumption by microbial reactions.
18	•	Summary of indicative timescales of hydrogen migration within crystalline and sed-
19		imentary rocks.

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#### 20 Abstract

The occurrence of natural hydrogen and its sources have been reviewed extensively in 21 the literature over the last few years, with current research across both academia and 22 industry focused on assessing the feasibility of utilising natural hydrogen as an energy 23 resource. However, gaps remain in our understanding of the mechanisms responsible for 24 the large-scale transport of hydrogen and migration through the deep and shallow Earth 25 and within geological basins. Due to the unique chemical and physical properties of hy-26 drogen, the timescales of migration within different areas of Earth vary from billions to 27 thousands of years. Within the shallow Earth, diffusive and advective transport mech-28 anisms are dependent on a wide range of parameters including geological structure, mi-29 crobial activity, and subsurface environmental factors. Hydrogen migration through dif-30 ferent media may occur from geological timescales to days and hours. We review the na-31 ture and timescale of hydrogen migration from the planetary to basin-scale, and within 32 both the deep and shallow Earth. We explore the role of planetary accretion in setting 33 the hydrogen budget of the lower mantle, discuss conceptual frameworks for primordial 34 or deep mantle hydrogen migration to the Earth's surface and evaluate the literature on 35 the lower mantle's potential role in setting the hydrogen budget of rocks delivered from 36 the deep Earth. We also review the mechanisms and timescales of hydrogen within dif-37 fusive and advective, fossil versus generative and within biologically moderated systems 38 within the shallow Earth. Finally, we summarise timescales of hydrogen migration through 39 different regions within sedimentary basins. 40

#### 41 Plain Language Summary

Over the last several years, naturally-occurring hydrogen has emerged as a poten-42 tial game changer in the energy transition. However, the vast majority of current research 43 focuses on understanding hydrogen generation and underground storage. Important ques-44 tions remain - how did hydrogen come to be within Earth? What are the timescales of 45 hydrogen movement through different regions of Earth and within different rock types? 46 In this review, we describe the processes responsible for entraining hydrogen into Earth's 47 mantle during the time of planetary formation and timescale of evolution towards the 48 present-day hydrogen cycle. We summarise the importance of environmental factors and 49 mineralogy for hydrogen movement and the timescale of generative and destructive pro-50 cesses in the shallow Earth. Finally, we summarise the timescale of hydrogen movement 51 within different regions of Earth from the planetary to basin-scale and within different 52 minerals and rock types. 53

#### 54 1 Introduction

Hydrogen is regarded as an important component of the world's transition towards 55 a low emission, net-zero future (IEA, 2021). Significant efforts are currently being made 56 across academia and industry to improve our understanding of hydrogen subsurface mo-57 bility, especially in the context of natural hydrogen occurrence and underground stor-58 age (e.g., Zgonnik (2020); Muhammed et al. (2022); L. Wang, Jin, et al. (2023)). Within 59 the literature, the term 'natural hydrogen' describes the hydrogen which is not manu-60 factured and is directly found in the subsurface. Natural Hydrogen is encountered as a 61 free gas (i.e., surface seeps), dissolved in groundwater and within fluid inclusions in rocks. 62 Natural hydrogen migrates through a wide range of mechanisms, including diffusion (e.g., 63 through crystalline lattices) and advection (e.g., dissolution in groundwater, migration 64 along faults) (e.g. Farver (2010); Lefeuvre et al. (2021, 2022); Strauch et al. (2023); Truche 65 et al. (2024)). Whilst laboratory experiments and assessments of specific case studies have shed some light on the complex nature of natural hydrogen migration, large gaps remain 67 in our understanding of the mechanisms responsible for the large-scale migration of hy-68 drogen through the deep Earth and within sedimentary basins. In this review, we refer 69

to the regions beneath the lithosphere, i.e. mantle, as the 'deep' Earth, and subsurface
 regions within geological and sedimentary basins as the 'shallow' Earth.

This review overviews the processes responsible for setting Earth's hydrogen budget and the timescales of hydrogen migration across all length-scales within Earth. We review (i) the origin of primordial or deep mantle natural hydrogen supply to the deep Earth, the role of water on mantle mixing and Earth's hydrogen cycle over geological timescales and (ii) the dynamics of diffusive and advective hydrogen migration within the shallow Earth on geological to human timescales.

#### <sup>78</sup> 2 Hydrogen in the deep Earth

Hydrogen is unlike any other molecule. Whilst it is the most abundant element in 79 the universe, molecular hydrogen is scarce on Earth. Estimates of the hydrogen abun-80 dance in Earth's interior have spanned a range from less than the equivalent of the cur-81 rent hydrosphere to on the order of 100 hydrospheres if hydrogen is the dominant light 82 alloying component in Earth's outer core (Williams & Hemley, 2001). There is a limited 83 understanding of its sources, migration through rocks and whether hydrogen can accu-84 mulate in geological formations for significant time periods. To understand the nature 85 of hydrogen migration within Earth, we must first consider its origins and distribution 86 on a planetary scale. Within the deep Earth, there exist reservoirs of primordial or deep 87 mantle hydrogen trapped during the period of planetary formation, which are proposed 88 by some models to be transported by advection within the mantle on timescales of bil-89 lions of years (Peslier et al., 2017; Loewen et al., 2019). For a detailed description of the 90 mineralogical composition and hydrogen content of the deep Earth, see Williams and Hem-91 ley (2001). Following planetary accretion, the stabilisation of liquid water on the sur-92 face and onset of plate tectonics had a profound impact on the dynamics of mantle flow 93 and Earth's hydrogen cycle. This includes the contamination of non-native material into Earth's mantle at subduction zones and its heterogeneity in mantle hydrogen contents. 95 Conversely, the preservation of isotopic signatures indicate that mantle material deliv-96 ered to Earth's surface at hotspot settings preserve their deep mantle or primordial iso-97 topic signatures and do not mix with surrounding mantle (Mangenot et al., 2023). The 98 timescale of hydrogen transport from within the deep Earth to different geological set-99 tings at the surface vary across several orders of magnitude. In this section, we describe 100 the discrepencies between hydrogen and helium isotopic ratios encountered in rocks from 101 different geological settings and compare frameworks for hydrogen migration to the Earth's 102 surface from the deep mantle. 103

#### 104 Global hydrogen cycle

Whilst up to 90% of the proto-solar nebula comprised of hydrogen, <sup>1</sup>H, the isotopes 105 deuterium, <sup>2</sup>H, and Helium-3, <sup>3</sup>He, were also created during the Big Bang. Unlike ter-106 restrial  ${}^{4}$ He, which is mainly produced by decay of uranium and thorium, terrestrial  ${}^{3}$ He 107 is largely of primordial origin, synthesised in the aftermath of the Big Bang (Bania et 108 al., 2002) and incorporated into the Earth primarily during its formation (Lupton & Craig, 109 1975). In spite of its primordial status and 4.56 Ga of planetary evolution, up to  $\sim 2$ 110 kg <sup>3</sup>He continues to leak from Earth's interior and mainly along mid ocean ridges (Olson 111 & Sharp, 2022). The reference proto-solar D/H ratio is  $\sim 2.1 - 2.5 \times 10^{-5}$ , which is 112 close to the Big Bang value. Due to its mass, hydrogen is lost through diffusion prefer-113 entially over deuterium and the D/H ratio increases with geologic time. As <sup>4</sup>He is a de-114 cay product of U-Th-Pb  $\alpha$ -decay systems, <sup>3</sup>He/<sup>4</sup>He ratios with Earth decrease mono-115 tonically with time, with high  ${}^{3}\text{He}/{}^{4}\text{He}$  rocks indicating preservation in mantle domains 116 that are not modified by convective mixing or diffusive homogenisation since early Earth 117 history (Porcelli & Elliott, 2008; Huang et al., 2014; Cooke et al., 2014; Lis et al., 2019). 118 The reaction of deuterium, hydrogen and water,  $HD + H_2O \rightleftharpoons H_2 + HDO$ , is an impor-119

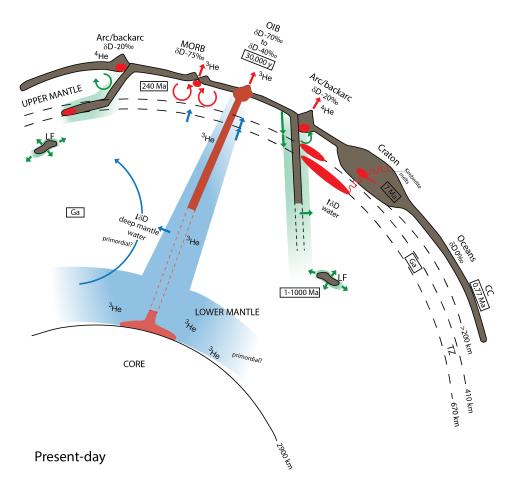


Figure 1. Conceptual model for primordial or deep mantle hydrogen and helium migration through Earth and present-day subduction patterns. Blue shading indicates regions of depleted mantle that sample reservoirs of primordial or deep mantle water, H and He. Green shading indicates infiltration of non-mantle water as a consequence of tectonic process on Earth's surface. Lithospheric fragments (LF) which break from subducting slabs may constitute a source of non-mantle material to the deep mantle (e.g. (Sperner et al., 2001; Zahirovic et al., 2016; van der Meer et al., 2018; Kufner et al., 2021)). The asthenosphere is a region within the upper mantle and beneath the lithosphere in which there is relatively low resistance to plastic deformation due to the partial melting of rocks initiated by the infiltration of hydrous mineral phases entrained on subducting lithosphere. The depth of the asthenosphere varies throughout Earth, however generally lies between 80 - 200 km depth. TZ = transition zone, red ovals = regions of partial melting in the upper mantle, lithosphere and TZ. Boxes = hypothesised residence times of water within different layers of Earth; Ga = billion years, Ma = million years (Bodnar et al., 2013). Continental and lithospheric thicknesses are not to scale. Modified from (Peslier et al., 2017; Loewen et al., 2019).

tant measure of the thermal history of water molecules since the formation of Earth. The D/H ratio is expressed in delta notation as  $\delta D$ , whereby  $\delta D = [(D/H)_{sample}/(D/H)_{V_{SMOW}} - (D/H)_{V_{SMOW}}$ 

<sup>122 1]×10&</sup>lt;sup>3</sup> with  $V_{SMOW}$  the Vienna Standard Mean Ocean Water. Hence, low  $\delta D$  values

within high  ${}^{3}\text{He}/{}^{4}\text{He}$  rocks may be used as a diagnostic isotopic signature to determine

<sup>the origins of water and mantle material within Earth (Geiss & Gloeckler, 1998; Loewen
et al., 2019; Pinti, 2021).</sup> 

Mid Ocean Ridges are transform margins where upwelling mantle is extruded at 126 Earth's surface to form new oceanic crust. Mid-Ocean Ridge Basalts (MORBs) are mafic 127 rocks derived from larger mantle domains that appear to sample deep mantle hydrogen 128 transported to the melting domain in the upper mantle by large-scale mantle convection 129 and typically have low  $\delta D$  values of  $\sim -70\%$  (Table 1) (Craig & Lupton, 1976; Rison 130 & Craig, 1983; Poreda et al., 1986; Graham, 2002; Jackson et al., 2017; Loewen et al., 131 2019). <sup>3</sup>He/<sup>4</sup>He ratios in MORBs are typically homogeneous and have a narrow range 132 of 7 – 9  $R_A$ , whereby  $R_A$  = atmospheric ratio (C. J. Allègre et al., 1995; Gautheron 133 & Moreira, 2002). Helium ratios analysed in MORB glasses by C. J. Allègre et al. (1995) 134 from individual ridge segments show a linear correlation with the ridge spreading rate. 135 Average helium ratios of MORBs are predominantly uniform and distinct from lower man-136 tle and transition zone values, which is interpreted by C. J. Allègre et al. (1995) to in-137 dicate the existence of two scale upper mantle convection. Rapid convection within the 138 uppermost mantle that feeds mid ocean ridges (see red arrows on Figure 1) is respon-139 sible for the homogenisation of helium in this layer and is calculated to have a mixing 140 time of  $\sim 250$  Ma, which is distinct from the  $\sim 1$  Ga residence time of upper mantle 141 rocks (C. Allègre et al., 1983; C. J. Allègre et al., 1995) (Figure 1). 142

Whilst enriched compared to non-mantle rocks, helium ratios in MORBs are sig-143 nificantly less than those found at Ocean Island Basalt (OIB) and Continental Hotspot 144 (CH) settings, which can reach values  $< 40 R_A$  (Table 1). Although the nature of plume 145 development over geological time remains an active topic of research, it is widely accepted 146 that mantle plumes which feed OIB and CH settings may extend as deep as the core/mantle 147 boundary ( $\sim 2900$  km) (Figure 1). Plumes that transport lower mantle rocks with high 148  ${}^{3}\text{He}/{}^{4}\text{He}$  and low  $\delta D$  signatures that are enriched in deep mantle water, i.e. water that 149 has never been present at Earth's surface, are significant as they are conduits of deep 150 mantle hydrogen and helium entrained in ultramafic rocks that are extruded at Earth's 151 surface. Recent investigations into the architecture of hotspot-plume systems suggest that 152 material transported via hotter and more buoyant mantle plumes have increased resilience 153 to mixing with surrounding mantle and increased preservation of deep mantle geochem-154 ical signatures compared to colder and less buoyant plumes (Samuel & Farnetani, 2003; 155 S. C. Lin & Keken, 2006; Garnero et al., 2016; Jackson et al., 2017; Jimenez-Rodriguez 156 et al., 2023). It is noteworthy that CH settings exhibit  $\delta D$  values close to MORB, how-157 ever  ${}^{3}\text{He}/{}^{4}\text{He}$  ratios are significantly higher (< 20 R<sub>A</sub>). Some, like the African hotspots, 158 also have  $\delta D$  values as low as -89% (Table 1) (Jimenez-Rodriguez et al., 2023). Further-159 more, experimental results from Mangenot et al. (2023) indicate that H<sub>2</sub> is sensitive to 160 isotope re-equilibration (e.g. between  $H_2$  and water at its source) during the ascent and 161 cooling of high-temperature crustal, magmatic, and mantle fluids. These observations 162 indicate that material transport from the lower mantle must be fast enough to prevent 163 mixing with both upper mantle and the surrounding continental cratonic rocks and sig-164 nificantly faster than the timescale upper mantle mixing. Hence, we hypothesise that the 165 degree of enrichment of deep mantle hydrogen and helium within rocks delivered to OIB 166 and CH settings, and therefore their surrounding terrestrial environments through sub-167 surface processes e.g. serpentinisation, is controlled by the rate at which rocks are sup-168 plied from lower mantle reservoirs. This must be on the timescale of, at least, millions 169 of years instead of hundreds of millions of years. Thus, the timescale of migration of rocks 170 enriched in hydrogen from within the deep Earth varies over several orders of magnitude 171 between different geological regimes and is distinct from water residence times. 172

Hence, it is likely that whilst the background flux of primordial or deep mantle hydrogen to upper mantle regions which supply melt to Mid Ocean Ridges is set by largescale mantle convection over billions of years, the mixing of primordial or deep mantle
hydrogen and helium-enriched material into the upper mantle is an order of magnitude
faster over hundreds of millions of years.

#### Primordial versus deep mantle hydrogen and helium ratios

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Although the primordial origins of  ${}^{3}$ He are generally accepted, there is controversy 179 within the scientific community regarding application of <sup>3</sup>He/<sup>4</sup>He ratios as a diagnos-180 tic marker for primordial material. The value of the  ${}^{3}\text{He}/{}^{4}\text{He}$  ratio from the local inter-181 stellar medium (LISM) is  $1.7\pm0.8\times10^{-4}$ , and is around two orders of magnitude greater 182 than the present-day atmospheric value,  $R_A = 1.4 \times 10^{-6}$  (Graham (2002); Salerno et 183 al. (2003), see Table 1). LISM values are consistent with protosolar ratios obtained from 184 meteorites and Jupiter's atmosphere, supporting the hypothesis that negligible changes 185 of the abundance of <sup>3</sup>He occurred in the galaxy during the past 4.5 Ga (Salerno et al., 186 2003). High  ${}^{3}\text{He}/{}^{4}\text{He}$  ratios in Ocean Island Basalts (OIBs) were traditionally interpreted 187 as indicators of a primitive, undegassed mantle source that has been trapped within the 188 Earth to the present day (Bouhifd et al., 2013). This interpretation was predicated on 189 the assumption that these ratios reflect the composition of material sampled by rapidly 190 ascending thermally buoyant plumes arising from deep within the mantle (Kurz et al., 191 1982; Morgan, 1971; Kellogg & Wasserburg, 1990). Despite this established view, pri-192 mordial models fail to account for observed discrepancies in helium concentrations and 193 the elemental ratios of He and other noble gases (e.g. Ar and Ne) between OIBs and Mid-194 Ocean Ridge Basalts (MORBs), with the former displaying values an order of magni-195 tude lower than those found in MORBs (Gonnermann & Mukhopadhyay, 2007). 196

The debate over the origins of high  ${}^{3}\text{He}/{}^{4}\text{He}$  ratios in OIBs and how these ratios 197 address the helium paradoxes has been extensively explored in the literature. Gonnermann 198 and Mukhopadhyay (2007) present a model in which the helium concentration paradox, 199 as well as the variance in noble-gas concentrations observed in MORB and OIB glasses, 200 can be explained by disequilibrium open-system degassing of erupting magma. Their work 201 suggests that higher  $CO_2$  content in OIBs leads to more extensive helium degassing in 202 OIB magmas compared to MORBs, thus deriving noble gases in OIB lavas from a largely 203 undegassed primitive mantle source. This interpretation aligns with the conventional view 204 that high <sup>3</sup>He/<sup>4</sup>H ratios in OIBs indicate parts of the deep mantle have remained iso-205 lated from outgassing and the convective upper mantle over Earth's history (Gonnermann 206 & Mukhopadhyay, 2007). 207

Bouhifd et al. (2013) assess helium partitioning in experiments between molten sil-208 icates and iron-rich metal liquids at conditions representative of Earth's lower mantle 209 and core. Their results and estimated concentrations of primordial helium suggest that 210 significant quantities of helium may reside in the core and that the early core could have 211 incorporated enough helium to supply deep-rooted plumes enriched in <sup>3</sup>He throughout 212 Earth history. Bouhifd et al. (2013) therefore suggest that two variations in the  ${}^{3}\text{He}/{}^{4}\text{He}$ 213 ratio observed at the surface in OIBs and MORBs may be explained by two distinct reser-214 voirs in the Earth's interior (e.g. Hopp and Trieloff (2008)). These are a conventional 215 depleted mantle source and a deep, still enigmatic, source that must have been isolated 216 from processing throughout Earth history. However, modelling of helium ingassing into 217 a silicate magma ocean and iron-rich proto-core coupled to a nebular atmosphere of so-218 lar composition and outgassing into a coupled core-mantle system after accretion by Olson 219 and Sharp (2022) indicates that Earth's core may be a substantial and long-lived reser-220 voir of primordial helium. 221

Zhu et al. (2020), however, offer a contrasting perspective by proposing that he-222 lium contents and  ${}^{3}\text{He}/{}^{4}\text{He}$  isotopic ratios can be fractionated by thermal diffusion in 223 the lower mantle, driven by an adiabatic or convective temperature gradient. Their model 224 suggests that the lower mantle is helium stratified due to thermal diffusion, resulting from 225 226 a of  $\sim 400$  K temperature contrast across the lower mantle. Hence, Zhu et al. (2020) argue that helium fractionation, rather than the lower mantle being a primordial and 227 undegassed reservoir, explains the observed high <sup>3</sup>He/<sup>4</sup>H isotopic ratios and lower he-228 lium contents in OIBs. Zhu et al. (2020) argue that OIBs derived from the deepest lower 229 mantle, which display high <sup>3</sup>He/<sup>4</sup>H isotopic ratios and less helium content, can be ex-230

plained by their model, effectively addressing the long-standing helium concentration para dox without necessitating a primordial undegassed lower mantle reservoir.

These differing viewpoints illustrate the complexity of mantle dynamics and the 233 origins of helium isotopic variations and their implications for characterising truly pri-234 mordial source material. It is beyond the scope of this review to explore this subject in 235 greater detail, however it is widely accepted that isotopic signatures may be used to di-236 agnose the preservation of deep lower mantle material over geologically significant time 237 periods transported rapidly over planetary length scales (e.g., Mackintosh and Ballen-238 239 tine (2012)). However, whilst the influence of primordial material on rocks that preserve high  ${}^{3}\text{He}/{}^{4}/\text{He}$  isotopic signatures rocks are debated, we acknowledge that this cannot 240 be excluded entirely. 241

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#### The impacts of water on the deep Earth and early tectonics

As isotopic signatures of igneous rocks may be used to identify material sourced 243 from deep lower mantle reservoirs, they offer an insight into the hydrogen cycle of the 244 deep Earth. Whilst the diffusion of water and hydrogen in silicates is fast compared to 245 other elements, it cannot explain the heterogeneity of lower mantle material or enrich-246 ment within rocks delivered by mantle plumes or upwelling. In the absence of water in-247 filtration into the deep mantle, mantle convection alone would lead to homogeneous wa-248 ter contents among regions of more than  $\sim 100$  km size (Peslier et al., 2017). For typ-249 ical asthenospheric conditions, diffusion of hydrogen over a distance of  $\sim 10$  km takes 250  $\sim 1$  Ga (Karato, 2007; Peslier & Bizimis, 2015; Peslier et al., 2017). Estimates of the 251 residence time of water within different layers in Earth are calculated by Bodnar et al. 252 (2013) as < 3000 years for the hydrosphere, 0.77 - 7 Ma (million years) for the litho-253 sphere and  $\sim$  Ga (billion years) for the transition zone and lower mantle (Figure 1). 254

Combined evidence from several radionuclide systems (Pd-Ag, Mn-Cr, Rb-Sr, U-255 Pb) suggests that water was not incorporated in Earth in significant quantities until the 256 planet had grown to  $\sim 60 - 90\%$  of its current size, while core formation was still on-257 going (Peslier et al., 2017). Prior to the onset of plate tectonics, some models propose 258 that the early Earth lost heat generated from planetary accretion and radioactive de-259 cay of isotopes in the metallic core through degassing and volcanism according to a stag-260 nant or mobile lid regime during the late Hadean (Solomatov & Moresi, 2000; Capitanio 261 et al., 2022). The fractionation of Earth into the core, mantle and early crust is proposed 262 by those models to have created a stratified water structure within the planet. During 263 this period, photolysis from solar radiation and late crust-forming events led to significant loss of water from the early crust and upper mantle, leading to the lower mantle 265 becoming relatively enriched in primordial or lower mantle water, hydrogen and helium 266 (Peslier et al., 2017). The gradual accumulation of liquid water oceans is generally ac-267 cepted to be a result of the impact of chondritic material from the asteroid belt follow-268 ing the period of heavy bombardment and stabilisation of Earth's late veneer. Depend-269 ing on which estimates are used for the water and carbon contents of the bulk silicate 270 Earth, 20 - 100% of the early mantle's hydrogen and carbon may have been brought 271 to Earth by carbonaceous chondrites during this late stage of planetary formation (Marty 272 & Yokochi, 2006; Z. Wang & Becker, 2013; Marty, 2012; Peslier et al., 2017; Loewen et 273 al., 2019). 274

The earliest known evidence for liquid water present on Earth's surface includes the Isua Greenstone Belt, where pillow-lava structures consistent volcanic eruption in submarine conditions occur as early as  $\sim 3.8$  billion years ago (Polat & Hofmann, 2003). The presence of stable liquid water on Earth's surface marks a significant point in geological history, as the infiltration of water-rich mineral phases into the upper mantle reduced melting temperatures and led to the formation of the mechanically weak Astheno-

Setting	$\delta \mathrm{D}~[\%]$	$^{3}\mathrm{He}/^{4}\mathrm{He}\left[\mathrm{R}_{A} ight]$	Example	Reference
OIB	< -75	< 40	Iceland, Hawaii Samoa, Galapagos Easter	(Jackson et al., 2017) (French & Romanowicz, 2015) (Boschi et al., 2007) (Poreda et al., 1986) (Rison & Craig, 1983)
MORB	$-70^{*}$	7 - 9	Iceland Mid Atlantic Ridge	(C. J. Allègre et al., 1995) (Gautheron & Moreira, 2002)
Arc/back-arc	$-20$ to $-40^{**}$	$\sim 1^{***}$	Valu Fa ridge Andes	(Hilton et al., 1993)
СН	-23 to $-89$	8 - 20	Afar, Darfur Hoggar	(Jimenez-Rodriguez et al., 2023) (Jackson et al., 2017)

Table 1. Typical isotope ratios of various tectonic settings. OIB = Ocean Island Basalt, MORB = Mid Ocean Ridge Basalt,  $CH = Continental hotspot. \delta D$  values are measured relative to Standard Mean Ocean Water (see text).  ${}^{3}\text{He}/{}^{4}\text{He}$  ratios are measured relative to present-day atmospheric values,  $R_A$ , where  $R_A = 1.4 \times 10^{-6}$  (Graham, 2002). OIB settings supported by hotter mantle plumes with increased buoyancy are hypothesised to support increased transport of primordial or deep mantle material from the lower mantle to Earth's surface (Jackson et al., 2017). \*Extremely low  $\delta D$  values of < -90 ‰ are observed in some MORBs associated with multi-stage melting and anhydrous minerals isotopic fractionation (see text) (Loewen et al., 2019). \*\*High  $\delta D$  values in typical arc and back-arc settings (< -20%) indicate the recycling of non-mantle water due to the subduction of hydrous minerals contained in the mantle wedge from beneath arc lavas, with lower  $\delta D$  values (> -40\%) associated with mixing with depleted mantle sources (Shaw et al., 2008). \*\*\*Igneous rocks from arc/back-arc settings not associated mantle mixing (i.e. differentiated silicic rocks) have negligible  ${}^{3}\text{He}/{}^{4}\text{He}$  ratios ~  $1R_A$  (Hilton et al., 1993). Continental hotspots, most notably African examples, vary significantly from continental lithoshpere and differentiated back-arc basin igneous rocks. Mantle-supported continental hotspot have  ${}^{3}\text{He}/{}^{4}\text{He}$  ratios ranging from typical MORB values to  $20R_{A}$  (Jackson et al., 2017).

sphere, and the onset of plate tectonics around  $\sim 3$  billion years ago (Farquhar et al., 2002; Shirey & S.H., 2011; Debaille et al., 2013).

Since both helium and hydrogen are incompatible during mantle melting (i.e. both 283 partition into a melt as soon as melting begins), high  ${}^{3}\text{He}/{}^{4}\text{He}$  ratios characterise a man-284 tle that has been isolated from melting and degassing since the earliest stages of Earth 285 history (C. Allègre et al., 1983; Mukhopadhyay, 2012; Loewen et al., 2019). Hence, rocks 286 which contain low  $\delta D$  and high  ${}^{3}\text{He}/{}^{4}\text{He}$  signatures are a prime target for understand-287 ing sources of deep mantle water that has survived significant mixing and transport over 288 planetary lengthscales (Craig & Lupton, 1976; Rison & Craig, 1983; Poreda et al., 1986; 289 Loewen et al., 2019; Mackintosh & Ballentine, 2012). 290

Subduction zones represent the primary regions of terrestrial water exchange be-291 tween Earth's interior and hydrosphere. It is estimated that  $\sim 25\%$  of the water enter-292 ing subduction zones reaches the transition zone and  $\sim 3\%$  reaches the lower mantle through 293 transport via lithosphere fragments (LF) that separate from subducted slabs (Figure 1) 294 (Bodnar et al., 2013). Slab break-off and the transport of lithosphere fragments to the 295 deep mantle has been investigated extensively, and is supported by plate tectonic recon-296 structions and geophysical data (e.g. (Williams & Hemley, 2001; Sperner et al., 2001; 297 Zahirovic et al., 2016; Kufner et al., 2021)). It is impossible to know the precise amount 298

of lithospheric material that has been returned to the mantle over geologic time, however it is reasonable to imagine the presence of graveyards of remnant fossil lithosphere distributed heterogeneously throughout the mantle (e.g., van der Meer et al. (2018)).

The subduction of hydrous minerals and recycling of water into the mantle over 302 geologic time and to the present day thus led to an increase in  $\delta D$  in Earth's crust, up-303 per mantle, oceans and atmosphere compared to primitive or lower mantle materials. Es-304 timates of the water content within Earth range from  $7-14 M_{oceans}$  within the man-305 tle and  $< 12 M_{oceans}$  in the core (Bodnar et al., 2013; Nestola & Smyth, 2016; Peslier 306 et al., 2017). Hence, subuduction and lithosphere fragments have a profound impact on 307 the dynamics of mantle convection and Earth's hydrogen cycle the through mixing of 308 non-mantle material into deep mantle reservoirs. 309

As Earth's crust and tectonic processes evolved, the planet's atmosphere also un-310 derwent significant changes that impacted the stability of water and hydrogen genera-311 tive processes on Earth's surface. Dodd et al. (2022) investigated hydrogen dynamics be-312 fore and after the Great Oxidation Event (GOE) at  $\sim 2.5 - 2.0$  Ga. Initially, abiotic 313 reactions in anoxic conditions led to hydrogen generation from banded iron formations 314 (BIFs), with free hydrogen escaping due to low oxygen levels. As the concentration of 315 biologically generated O<sub>2</sub> within Earth's atmosphere gradually increased, the atmosphere 316 changed from weakly reducing conditions and practically devoid of oxygen into oxidis-317 ing conditions, and containing abundant free oxygen Torres et al. (2015). Post-GOE, el-318 evated oxygen facilitated water formation by reacting with hydrogen, reducing hydro-319 gen escape and transitioning Earth to a more oxidised state supportive of aerobic life and 320 altering geochemical dynamics significantly (Dodd et al., 2022). 321

Present-day patterns of high-angle subduction and mantle wedge hydration (i.e. 322 as shown on Figure 1) were not dominant during the early Earth, as most present-day 323 subduction initiation mechanisms require acting plate forces and existing zones of litho-324 spheric weakness, which are both consequences of plate tectonics. In the absence of plate 325 tectonic-related subduction, mechanisms responsible for the initiation of tectonics dur-326 ing early history are theorised to be plume-induced subduction, which is only feasible 327 in the hotter early Earth for old oceanic plates. In contrast, younger plates favoured episodic 328 lithospheric drips rather than self-sustained subduction and global plate tectonics (Gerva 329 et al., 2015). It is possible that the development of the modern, globally interconnected 330 plate network and subduction-related tectonics did not arise until billions of years af-331 ter the formation of the earliest crust and as late as Proterozoic times (Wan et al., 2020). 332 This assertion is consistent with modeling studies that demonstrate that much of the con-333 tinental crust of Archean cratons could have been generated in the absence of subduc-334 tion (Capitanio et al., 2019; Johnson et al., 2017). 335

However, some evidence indicates localised infiltration of hydrated mantle wedges 336 into the mantle occurred as early as 3.1 Ga, there is a consensus in the literature that 337 higher mantle temperatures, lower mantle viscosity and the subduction or infiltration 338 of oceanic crust at an unusually low angle was responsible for the growth of continen-339 tal crust older than  $\sim 2.5$  Ga (e.g., Perchuk et al. (2023); Smithies et al. (2003)). Al-340 though, it is important to note that a hotter mantle would lead to lower viscosity and 341 thus more melt, with the lower viscosity leading to more frequent slab breakoff, and to 342 increased crustal separation from the mantle lithosphere (van Hunen & van den Berg, 343 2008). Crustal and lithospheric fragments which break off, contaminate, and sink into 344 the mantle raise  $\delta D$  and lower  ${}^{3}He/{}^{4}He$  ratios over geologic time away from primordial 345 or lower mantle values and towards their present-day values. Therefore, it is our view 346 347 that whilst localised tectonic processes early in Earth history will have had some impact on the distribution of hydrogen within the mantle, the onset of global subduction marked 348 the turning point of large-scale mantle mixing leading to present-day heterogeneous man-349 tle hydrogen contents. Throughout Earth history, plume-related tectonics will have been 350

responsible for transporting material enriched in primordial or lower mantle hydrogen and helium to the surface on geologically short timescales.

#### 353 **3** Hydrogen in the shallow Earth

Whilst the migration of hydrogen in the deep Earth is dependent on large-scale man-354 the convective and tectonic processes that operate from billions to millions of years, the 355 enrichment of hydrogen within near-surface systems and ongoing emission from surface 356 seeps across the globe represents an intriguing duality of length-scales and timescales. 357 Shallow Earth processes, such as the migration in porous media, migration along faults 358 and fractures and microbial reactions may operate over timescales of thousands of years 359 to hours. To understand hydrogen migration in the shallow Earth, we must acknowledge 360 the relationship between hydrogen sources and transport mechanisms at the crustal, basin 361 and outcrop scale, i.e.  $10^3 - 1$  m. The primary mechanisms of natural hydrogen gen-362 eration are thought to be: (1) serpentinisation of mafic rocks, (2) radiolysis of water, (3) 363 rock fracturing and (4) volcanic degassing, (5) maturation of organic matter and (6) weath-364 ering of iron-rich rocks (Takai et al., 2004; Klein et al., 2013; Zgonnik, 2020; Lefeuvre 365 et al., 2021, 2022; Geymond et al., 2022; Horsfield et al., 2022; Mahlstedt et al., 2022; Boreham et al., 2023; L. Wang, Jin, et al., 2023). 367

However, to date a distinction between processes that release primordial or deep 368 mantle (i.e. fossil) hydrogen and the chemical and biological production/destruction of 369 'new' hydrogen is seldom made in the literature. Hydrogen gas concentrations of > 10%370 have been encountered in various locations across different tectonic regimes (see ? (?)) 371 with one documented case of a successful resource discovery in Mali (Prinzhofer et al., 372 2018). The mineralogies of Archaen - Proterozoic basement rocks in continental cratonic 373 regions (e.g. Africa, Brazil, Russia) and mantle-derived rocks (e.g. MORBs, OIBs) are 374 enriched in hydrogen. Hence, within the continental realm, regions of high hydrogen con-375 centration coincide with sedimentary basins underlain by Archean - Proterozoic cratonic 376 rocks enriched in hydrogen (Zgonnik, 2020; Moretti, Brouilly, et al., 2021). Early esti-377 mates of global hydrogen production rates via both radiolysis and hydration reactions 378 from the Precambrian continental lithosphere were reported at  $0.36-2.27\times10^{11}$  moles/year 379 and are comparable to estimates from marine systems (Lollar et al., 2014). As recent lit-380 erature has reviewed the topic of natural hydrogen generation extensively (e.g. Zgonnik 381 (2020); Moretti, Brouilly, et al. (2021); L. Wang, Jin, et al. (2023)), we limit our cov-382 erage of this topic and focus on the relationship between natural hydrogen generation 383 and its migration pathways to Earth's surface. In this section, we review the mechanisms 384 and timescales of transport of hydrogen within diffusive and advective systems, includ-385 ing transport along faults and microbial reactions, within the shallow Earth. 386

Diffusion

# 387 388

#### Diffusion in crystalline rocks and minerals

Diffusive mechanisms transport hydrogen without any motion of a material's bulk 389 (e.g. rock, crystalline matrix or fluid). Experimental results of hydrogen diffusivity within 390 crystalline rocks are reviewed extensively by Farver (2010); Demouchy (2010); Li and 391 Chou (2015) (see references therein) and summarised in Figure 2. Within the primary 392 mafic rock-forming minerals olivine, pyroxene and amphibole, the Arrhenius plots of Farver 393 (2010) indicate a pattern of decreasing hydrogen diffusivity from  $10^1 \text{ cm}^2/\text{year}$  to  $10^1$ 394  $\mathrm{mm}^2/10$  ka with decreasing Mg content. Hydrogen diffusivity within quartz and feldspar 395 vary between  $10^1 \text{ cm}^2/\text{year}$  to  $10^1 \text{ cm}^2/\text{ka}$ , however have been measured up to  $1 \text{ m}^2 \text{ yr}^{-1}$ 396 - 1 m<sup>2</sup> day<sup>-1</sup> in the case of fused quartz at temperatures > 1200 K (Li & Chou, 2015). 397 Oxide minerals, which are significant components of soils and regolith, along with meta-398 morphic minerals (e.g. garnet) have hydrogen diffusivities of  $10^1 \text{ cm}^2/\text{year} - 10^1 \text{ cm}^2/\text{ka}$ 300 (Figure 2B). A strong relationship between mineral structure and hydrogen diffusivity 400

<sup>401</sup> is also seen (Figure 2C). Experiments by Kohlstedt and Mackwell (1998); Demouchy and <sup>402</sup> Mackwell (2006); Demouchy (2010) show a clear distinction between hydrogen diffusiv-<sup>403</sup> ity within crystalline aggregate and at grain boundaries (Figure 2D). For olivine, hydro-<sup>404</sup> gen diffusivity at grain boundaries is measured at ~ 1 m<sup>2</sup> year<sup>-1</sup>, which can be consid-<sup>405</sup> ered instantaneous given an average grain boundary thickness of 0.75 nm (Demouchy, <sup>406</sup> 2010).

Whilst such experimental results show that hydrogen diffusivity at grain bound-407 aries and within mineral aggregates (e.g. fused quartz) may be significant, this is out-408 weighed by exponentially decreasing diffusivity with increasing grain size (Figure 2E). 409 Grain sizes for crystalline rocks are a function of their cooling histories, with cratonic 410 crystalline basement and mantle xenoliths exhibiting average grain sizes from millime-411 tres to several centimeters. However, individual crystals can reach up to 30 cm in size 412 in some ultramafic mantle xenoliths with prolonged cooling histories (Hoskin & Sundeen, 413 1985; Speciale et al., 2020; Sharapov et al., 2022)). Furthermore, the maximum temper-414 atures of diffusivity experiments (< 1600 K) are representative of mantle conditions and 415 not encountered within the continental realm and sedimentary basins, which typically 416 vary between  $\sim 300-500$  K (Hantschel & Kauerauf, 2009). The exponential relation-417 ship of Demouchy (2010) indicates a decrease in hydrogen diffusivity of  $\sim 3$  orders of 418 magnitude between grain sizes of 10 mm and 0.1 mm, at which point the relationship 419 flatlines. Given these experiments were conducted at a pressure and temperature rep-420 resentative of upper mantle conditions, it is reasonable to assume that the diffusivity of 421 olivine (and other minerals) at typical continental and sedimentary basin conditions will 422 be many orders of magnitude smaller than the measurements of Demouchy (2010). Given 423 these factors, experimental results indicate that native hydrogen entrained within the 424 mineral structure of crystalline rocks within the shallow Earth may diffuse on geolog-425 ical timescales from the most common rock forming minerals. Geochemical data obtained 426 by Parnell and Blamey (2017) indicate that common felsic lithologies, such as granites, 427 gneiss and conglomerates of Archean - Proterozoic (> 1600 Ma) age consistently con-428 tain an order of magnitude greater hydrogen concentration in their entrained fluid than 429 very young (< 200 Ma) granites. Parnell and Blamey (2017) found that sedimentary 430 rocks containing clasts of old basement also included a greater proportion of hydrogen 431 than young granites and hypothesise that a signature of hydrogen in the basement could 432 be conferred to the sediment and that modern sediment derived from old and young base-433 ment retains the signature of more or less hydrogen, respectively. It should be noted, how-434 ever, that the experimental results summarised by Parnell and Blamey (2017) refer to 435 bulk lithologies whereas those of Farver (2010) refer to individual minerals (e.g. olivine 436 and quartz). Furthermore, the preservation of high hydrogen abundances within fluid 437 inclusions and mineralised veins in ancient granites has been observed (Bourdet et al., 438 2023). Hence, diffusion from enriched Archean - Proterozoic crystalline basement may 439 supply a 'background' hydrogen flux to overlying sedimentary basin rocks on geological 440 timescales. Hydrogen diffusion from coarse grained crystalline rocks, e.g. crystalline base-441 ment, granites and their derived sedimentary products, e.g., conglomerates, must oper-442 ate on timescales of Ma - Ga in order to explain the provenance of high hydrogen sig-443 natures in sedimentary rocks that contain material sourced from enriched Archean - Pro-444 terozoic basement. This is consistent with the widely documented observation of higher 445 hydrogen fluxes in sedimentary basins in continental cratonic regions underlain by Archean 446 - Proterozoic basement (e.g. Zgonnik (2020); Moretti, Brouilly, et al. (2021)). In the case 447 of rapidly cooled upper mantle rocks, e.g. MORBs, volcanic glasses, pillow lavas, how-448 ever, grain sizes may be many orders of magnitude smaller than their continental coun-449 terparts and within the nanometre scale (e.g., Schlinger et al. (1988)). Hence, hydrogen 450 451 diffusivity in rapidly cooled crystalline rocks and at MOR settings will be significantly faster than in continental settings and potentially only a few orders of magnitude slower 452 than the lower temperature ranges of diffusivity experiments, i.e. 100 Ka - Ma or faster. 453 This is significant, since the age of most oceanic crustal rocks is < 60 Ma (Seton et al., 454 2020), hydrogen diffusion within oceanic crustal rocks will operate on the same timescale 455

as the age of rocks themselves and provide a mechanism for the degassing of mantle hy drogen to the surface and oceans.

458

#### Diffusion in sedimentary rocks

The rate of hydrogen diffusivity in sedimentary rocks is dependent on a wide range of factors, including lithology, porosity, permeability, temperature, pressure, salinity and water content. Typical values for hydrogen diffusivity in different sedimentary rocks, water and air are summarised in Table 2

Unlike for crystalline rocks and minerals, hydrogen diffusivity experiments for sed-463 imentary rocks are carried out at temperature and pressure conditions representative of 464 sedimentary basins. These are typically at temperatures between 288–413 K and pres-465 sures < 40 MPa. A strong positive temperature dependence is seen in silt, clay, coal, 466 shale and salt, whereby hydrogen diffusivity values vary by > 50% over a narrow range 467 of  $\sim 40$  K (C. Wang et al., 2024; J. Liu et al., 2022; Keshavarz et al., 2022; Vinsot et 468 al., 2014). Hydrogen diffusivity decreases with increasing pressure, owing to increased 469 fluid density at the same temperature and fixed space causing gas diffusion to be restrained 470 (J. Liu et al., 2022). Measurements by J. Liu et al. (2022) of density profiles of hydro-471 gen in silt (montmorillonite) corresponding to different pressures indicate that the min-472 eral adsorption layer is not influenced by increasing pressure, thus causing pressure to 473 have an important but albeit lesser, effect on hydrogen diffusivity than temperature. At 474 experimental conditions of 353 K and 10 MPa, a reduction in hydrogen diffusivity of  $\sim$ 475 50% at a threshold water content of  $\rho_{H_2O}^{ave} = 0.568 \text{ gcm}^{-3}$ , and  $\sim 12\%$  as salinity in-476 creases from 8-12 wt% is observed in silt (J. Liu et al., 2022). This may be explained by the effects of increasing water content and salinity on the geometry of brine-rock and 478 brine-hydrogen molecule contacts. When water content and salinity are low, hydrogen 479 and brine form a stratified structure and the diffusivity of hydrogen is similar to that 480 of confined pure gas at the same pressure and temperature conditions. Increasing wa-481 ter content and salinty leads to increased connectivity of brine molecules, the formation 482 of water bridges between brine molecules and mixing of hydrogen and brine as a new phase, 483 leading to a decrease in hydrogen diffusivity by up to 5 orders of magnitude (Figure 3, 484 J. Liu et al. (2022)). Measurements of both dry and water-wet samples of various rock 485 types by (Strauch et al., 2023) indicate that hydrogen diffusivity decreases further by 486 an order of magnitude due to the effect of fracture healing with increased water content 487 (Strauch et al., 2023). Hydrogen breakthrough times vary significantly with water con-488 tent, most notably for salt rocks whereby values increase from 1 - 843 hours. 489

Measurements of  $H_2$ -brine contact angles versus pressure and total organic car-490 bon (TOC) in various water-wet Australian antracite shales by Al-Yaseri et al. (2022) 491 indicate that capillary entry pressure decreases with increasing pressure and TOC, thus 492 leading to a reduction in sealing capacity with depth and TOC. Within salt, hydrogen 493 diffusion is strongly dependent on mineralogy, crystal shape and size. Diffusion through 494 intact halite crystals with no discontinuities is negligible, whilst salt crystal boundaries 495 and fractures within grains are the preferential flow paths for gas diffusion (Yuan et al., 496 2023).497

Whilst hydrogen diffusivity experiments offer an insight into the absolute timescale 498 of hydrogen migration through different materials, a more useful measure of sealing abil-499 ity is the breakthrough time. The breakthrough time,  $t_b$  is defined as the time interval 500 between the start of gas purging of the feed chamber and the first detection of hydro-501 gen at the sensor in the permeate chamber. Whilst there are various methods to calcu-502 late  $t_b$ , the time-lag method of Frisch (1957) is widely applied due to its practicality and 503 ability to allow for the transient definition of a microstructure dependent correlation of 504 breakthrough time and sample thickness (Rhode et al., 2022). From Frisch (1957),  $t_b \sim$ 505

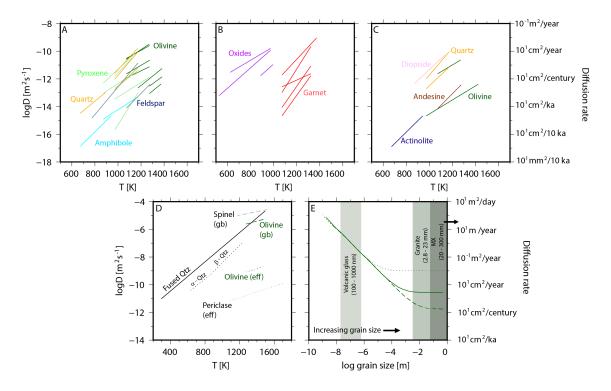


Figure 2. Hydrogen diffusivity as a function of mineralogy and lithology. A = Hydrogen diffusivity in the primary rock-forming minerals. These minerals constitute a significant component of Archean - Proterozoic crystalline basement (both mafic and felsic). B = Hydrogendiffusivity in oxides and garnets. C = Hydrogen diffusivity in different mineral structures. Whilst hydrogen diffusivity varies significantly with mineral structure, there is no obvious relationship between the two. D = Hydrogen diffusivity in fused,  $\alpha$  and  $\beta$  quartz from (Li & Chou, 2015). Effective diffusivity in and diffusivity at grain boundaries (gb) for olivine, spinel and periclase from (Demouchy, 2010). Diffusivity at gb are several orders of magnitude greater than within crystal lattices. E = Hydrogen effective diffusivity in olivine aggregate at 1473 K and 300 MPa 0.75 nm (Demouchy, 2010). Hydrogen diffusivities are calculated and grain boundary width = using the gb diffusion from (Demouchy, 2010) and the "proton-vacancy" mechanism for lattice diffusion in olivine along [001] (dashed line) and along [100] and [010] (dotted line, Demouchy and Mackwell (2006) and "proton-polaron" mechanism for lattice diffusion in olivine along [100] (solid line, Kohlstedt and Mackwell (1998)). Diffusivity decreases exponentially with grain size. Grain size ranges for plutonic granites and mantle xenoliths (MX) are shown in medium and dark grey (Hoskin & Sundeen, 1985; Speciale et al., 2020). Black horizontal arrow indicates MX grain sizes beyond the axes range (e.g., Sharapov et al. (2022)). Typical grain sizes for microcrystalline volcanic glasses vary from 100 - 1000 nm and are shown in light grey (Schlinger et al., 1988). Diffusivity data from Farver (2010) (A-C), citetDemouchy2010, Li2015 (D-E). Panel E modified from Demouchy (2010).

	Rock type	Diffusivity $[m^2 s^{-1}]$	$\mid t_{b_1} \text{ (dry) [years]} \mid t_{b_1} \text{ (wet) [years]}$	$t_{b_1}$ (wet) [years]	Reference
	Bentheimer sandstone	$1.6 - 2.1  imes 10^{-9}$	2.5	3.3	Strauch et al. (2023)
	Silt (montmorillonite)	$4.25-8.27\times 10^{-8}$	0.06	0.12	(J. Liu et al., 2022)
	) Clay	$1.2-5.13 imes 10^{-9}$	1.0	4.4	Vinsot et al. (2014) Strauch et al. (2023)
	Coal (anthracite)	$1.3-7 imes 10^{-8}$	0.08*	0.52*	Bagreev et al. (2004), Keshavarz et al. (2022)
	Shale	$1.3-2.4\times10^{-8}$	2.2	3.7	Al-Yaseri et al. (2022)
	Werra rock salt (halite)	$1.4 \times 10^{-9} - 1.3 \times 10^{-8}$	0.40	3.7	Strauch et al. (2023)
	Water (pure)	$3.9-6.1\times10^{-9}$			de Blok et al. (1982) Ferrell and Himmelblau (1967) Jacops et al. (2017) Hamme and van Bark (2018)
	Air	$0.756 - 1.604 \times 10^{-4}$	0.3 hours	0.6 hours	ToolBox (2018) ToolBox (2018)
	Stainless steel	$1.5  imes 10^{-11}$	352		Owczarek and Zakroczymski (2000)
<b>Table 2.</b> Hyd: time-lag method	rogen diffusivity value (see body text). Hyd	s for various sedimentary roc rogen diffusivity experiments	k types and materia to redimentary roc	ls and breakthrough ks are conducted at	Table 2.       Hydrogen diffusivity values for various sedimentary rock types and materials and breakthrough times through 1 m of material, $t_{b_1}$ , estimated using the time-lag method (see body text). Hydrogen diffusivity experiments for sedimentary rocks are conducted at temperatures and pressures representative of sedimen-
tary basin condi	tions and are much lo	wer than experimental condit	tions used for crystal	lline rocks and miner	tary basin conditions and are much lower than experimental conditions used for crystalline rocks and minerals. Variation in diffusivity values are due to dry and
wet samples. The The second se	the temperatures and p	ressures applied to the exper- tor air and stainlass stad an	imental results sumr	narised in this table	wet samples. The temperatures and pressures applied to the experimental results summarised in this table vary between 288 - 413 K and < 40 MPa, respectively. The diffusivity of hydrocon in price with and stainless studied for information with the literations values for cool are octi-
T THE MITMENT A	u mannden m bure we	aver, all allu scallless sveel al	a monant of the left of	TICE. Due to lack of	available uate III tile Interature, values for coal are

mated using H<sub>2</sub>S diffusivity measurements (Bagreev et al., 2004).

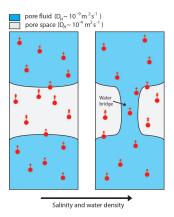


Figure 3. Effects of increasing salinity and water density on hydrogen diffusivity. Increasing salinity and density of pore fluid within sedimentary rocks leads to the formation of water bridges and increased connectivity between brine molecules, inhibiting the diffusion of hydrogen. Red circles = hydrogen molecules. Diffusion rates of hydrogen ( $D_H$ ) in air are up to five orders of magnitude greater than in water and are indicated by double and single arrows, respectively. Modified from (J. Liu et al., 2022)

 $L^2/6D_H$ , whereby L = length of sample thickness. Hence, the breakthrough time through 1 m of rock may be approximated as  $t_{b_1} \sim 1/6D_H$ .

Using published values, estimated  $t_{b_1}$  values for sedimentary rocks are shown in Table 2.  $t_{b_1}$  values are on the scale of years for dry and wetted rock samples, with breakthrough times for wetted rocks typically several times greater than for dry rocks (Strauch et al., 2023; J. Liu et al., 2022; Vinsot et al., 2014; Keshavarz et al., 2022; C. Wang et al., 2024; Al-Yaseri et al., 2022; de Blok et al., 1982; Ferrell & Himmelblau, 1967; Jacops et al., 2017; Hemme & van Berk, 2018; Mostinsky, 2011; ToolBox, 2018; Owczarek & Zakroczymski, 2000).

#### 515 Advection

Advection is the transport of a substance due to the motion of a carrier, e.g. wa-516 ter. In the case of hydrogen, gaseous  $H_2$  may also be carrier for itself. Advective trans-517 port systems within the subsurface include fluid migration along faults and discontinu-518 ities and groundwater flow through aquifers and are summarised on Figure 7. By its na-519 ture, the advection of hydrogen through sedimentary basins is a complex process that 520 cannot be described in a single step. A hydrogen molecule may migrate from depth to 521 the surface through a combination of the mechanisms shown on Figure 7 and diffusion 522 during these processes. 523

#### 524 Migration in porous media

Migration through porous sedimentary rocks is one of the most significant mech-525 anisms of fluid transport within sedimentary basins across the world. Fluids, such as hy-526 drocarbons and  $CO_2$ , displace and interact with formation fluids present within the pore 527 spaces in sedimentary rocks, with flow directed along pressure gradients exerted on the 528 system. In the case of hydrogen, transport via dissolution in groundwater within aquifers 529 is one of the primary mechanisms of advective migration in the shallow Earth. Natural 530 hydrogen migration via transport within sedimentary rocks and groundwater has not re-531 ceived much attention in literature and is not yet well understood. It is however known 532

to be affected by factors such as salinity, temperature, pressure, formation fluid composition and fluid-rock interactions.

Shallow aguifers have typical depths from near-surface to  $\sim 100$  m, whilst deep 535 aquifers may reach up to 9000 m. Temperature and salinity ranges in aquifers vary be-536 tween  $7-174^{\circ}C$  and 5-52000 ppm, respectively (Dopffel et al., 2021). These factors 537 exert a primary control on both the solubility of hydrogen in groundwater and micro-538 bial abundance and activity, which have important implications for hydrogen migration 539 (e.g. Berta et al. (2018); Koproch et al. (2019)). Similar to other gases, the solubility 540 of hydrogen in water decreases proportionally with salt (NaCl) concentration, until the 541 solution is saturated (Chabab et al., 2020). Recent solubility data of hydrogen in water-542 brine under geological conditions typical of aquifers were used by Chabab et al. (2020) 543 to determine empirical relationships between hydrogen solubility and salt concentration 544 in pure water and brine with 0.5-2% Average Absolute Deviation to observed values. 545 These are: 546

$$x_{H_2}^0 = b_1 P T + \frac{b_2 P}{T} + b_3 P + b_4 P_2 \tag{1}$$

547

for pure water at 273.15 < T < 323.15 K and 0.1 < P < 20.3 MPa, and

$$ln\left(\frac{x_{H_2}}{x_{H_s}^0}\right) = a_1 m_{NaCl}^2 + a_2 m_{NaCl} \tag{2}$$

for brine at 323.15 < T < 373.15 K, P 1 < P < 23 MPa and  $m_{NaCl} < 5$  mol/kg, 548 whereby T = temperature, P = pressure,  $X_{H_2} =$  hydrogen solubility,  $x_{H_2}^0 =$  hydro-549 gen solubility in pure water,  $m_{NaCl}$  = molality of salt and  $a_i, b_i$  = empirical coefficients 550 listed in (Chabab et al., 2020). The amount of dissolved hydrogen decreases with increas-551 ing salinity (see Figure 9 Chabab et al. (2020)). New measurements of hydrogen solu-552 bility carried out under conditions of high pressure (< 20 MPa), temperatures ranging 553 from 298-373 K and salinities < 4 mol/kgw of NaCl by Chabab et al. (2024) indicate 554 that H<sub>2</sub> solubility in water and brine increases with pressure and follows Henry's law in 555 a quasi-linear trend. Models optimised on experimental data predict a minimum solu-556 bility temperature  $(T_{xH2,min})$  of around ~ 326 K in pure water, decreasing with salin-557 ity ( $T_{xH2,min} = 315$ K at 2 molal NaCl and  $T_{xH2,min} = 288$ K at 4 molal NaCl) (Chabab 558 et al., 2024). 559

The interaction between natural hydrogen and resident formation fluid, e.g. wa-560 ter, present within the pore spaces of sedimentary rocks is also of significance as this in-561 fluences the displacement of the former over the latter within geological porous media. 562 Measurements of hydrogen-brine interfacial tension,  $\gamma$ , over a wide range of T, P and 563  $m_{NaCl}$  ranges by Hosseini et al. (2022) indicate an inverse linear relationship between 564  $\gamma$ , P and T.  $\gamma$  was found to have a strong dependence on temperature, and decreased 565 linearly at constant pressures and salinity.  $\gamma$  increased significantly and linearly with in-566 creasing salinity at constant temperatures and pressures, whilst decreasing at a lower rate 567 with increasing pressure at constant temperature and salinity due to the increasing inter-568 molecular forces between hydrogen and water at elevated pressures (Iglauer et al., 2012; 569 Hosseini et al., 2022). Hence, temperature and salinity have a greater effect than pres-570 sure on the solubility of hydrogen and its ability to displace formation fluids within porous 571 media (Hosseini et al., 2022; J. Liu et al., 2022). With regards to diffusivity, simulations 572 of molecular dynamics at subsurface conditions conducted by Kalati et al. (2024) indi-573 cate lower diffusivity at higher salinity and lower temperature, the value being  $7.29 \times$ 574  $10^{-9}$  m<sup>2</sup>/s at 323 K, increasing to  $10.2 \times 10^{-9}$  m<sup>-2</sup>/s at 353 K for 1 molal NaCl solu-575 tion. The diffusion coefficient decreases up to 38% as the salinity increases from 1 to 5 576 molal (Kalati et al., 2024). The results of Kalati et al. (2024) correspond to simulation 577

results by Lopez-Lazaro et al. (2019), which measure the temperature minimum hydrogen solubility close to 326 K.

#### Migration along faults and fractures

580

The episodic circulation of fluids and gases along geological faults is intricately con-581 trolled by mechanisms governing fault opening and sealing, as well as the timing of these 582 processes. Fluid advection along crustal faults is a well-documented phenomenon (Cox 583 & Etheridge, 1989), and it is a recurrent occurrence in the Earth's crust (Marques et al., 584 2018). These faults play a pivotal role in the migration of gas such as hydrogen, serv-585 ing as both conduits and barriers. Studies at both reservoir and basin scales have demon-586 strated that fault transmissivity is primarily influenced by (i) the fault's type, geome-587 try, and displacement; (ii) the internal architecture of the fault zone; (iii) the surround-588 ing stratigraphy and lithology; and (iv) the geomechanical stress (Faulkner et al., 2010; 589 Solum et al., 2010; Massiot et al., 2019). Due to the variability of these parameters, fault 590 transmissivity evolves both temporally and spatially (Frery et al., 2015; Frery, Fryer, et 591 al., 2021). The opening of faults can be triggered by seismic events, fluid overpressures, 592 or localised dissolution (Gratier & Gueydan, 2007), while their closure can be attributed 593 to progressive sealing resulting from mechanical (Hancock et al., 1999; Eichhubl & Boles, 594 2000), chemical processes and fault roughness (Renard et al., 2013). 595

Fault zones are complex features that can be effectively modelled as damaged zones 596 and gouges with heterogeneous porosity and permeability architectures. For instance, 597 in the North Perth Basin (Australia), these fault zones exhibit a highly compartmen-598 talised nature, primarily acting as barriers to crossflow while driving upward fluid mi-599 gration. This structural configuration provides an ideal setting for structurally controlled 600 hydrogen migration (Frery, Langhi, et al., 2021). For instance, above a natural subsur-601 face CO2 reservoir, a causal relationship between  $CO_2$  pulsing and fault opening have 602 been demonstrated using isotopic analysis (C, O, Sr ratios, Ba/Ca, and Sr/Ca elemen-603 tal ratios)(Kampman et al., 2012). The opening of veins is associated with a pulse of  $CO_2$ 604 within the system, followed by a degassing phase that occurs simultaneously with vein 605 growth. Consequently, abrupt events may have triggered the opening of fractures, im-606 mediately followed by episodes of fluid circulation. Evidence of fault opening events can 607 be observed with durations ranging from millennia to centuries (Burnside et al., 2013; 608 Frery et al., 2015; Gratier & Gueydan, 2007). 609

In regions of active faulting, stress cycling and the creation and destruction of per-610 meability and fluid flow are closely linked. Both large (km) and small-scale (i m) faults 611 are capable of influencing fluid migration pathways within sedimentary basins. The ad-612 vection of hydrogen-enriched fluids along large-scale faults are attributed to natural hy-613 drogen fluxes recorded in several well-known case studies, including Mali, Brazil and the 614 north Pyrenees (France) (Prinzhofer et al., 2018; Myagkiy, Brunet, et al., 2020; Donzé 615 et al., 2020; Lefeuvre et al., 2022). Common factors include the intersection of deep, crustal-616 scale faults with Archean-Proterozoic crystalline basement or ultramafic mantle bodies 617 that are serpentinised, hydraulic or elevated temperature and pressure gradients that trig-618 ger fluid migration. Measured daily flow rates of gaseous  $H_2$  flux within fault zones by 619 Lefeuvre et al. (2022) ranges from  $0.07-0.15 \text{ m}^3\text{m}^{-2}\text{d}^{-1}$  in the north Pyrenees. These 620 values are comparable to measurements of gaseous  $H_2$  flux within soils from the Sao Fran-621 cisco basin in Brazil and the Semail ophiolite, Oman (Prinzhofer et al., 2019; Zgonnik, 622 2020; Moretti, Prinzhofer, et al., 2021). The measurements of Lefeuvre et al. (2022) equate 623 to a timescale of  $\sim 128 - 274$  years for hydrogen migration over a distance of  $\sim 7$  km 624 from its serpentinite source to trap beneath a clay-rich seal. Templeton et al. (2024) note 625 that whilst low temperature water-rock reactions produces net H<sub>2</sub> from the oxidation 626 of Fe(II)-bearing minerals within the Semail Ophiolite, biological activity is predicted 627 to be stimulated by fluxes of  $H_2$ , giving rise to net  $H_2$  consumption. The most proba-628 ble detection of  $H_2$  at the surface is at hyperalkaline seeps sourced by deep faults, rather 629

than in most soils and peridotite outcrops, due to efficient microbial  $H_2$  scavenging of the available  $H_2$  flux in the upper aquifer, where measured  $H_2(aq)$  levels drop below detection (Templeton et al., 2024). Hence, from a migration perspective it is reasonable to hypothesise that the migration of hydrogen-rich fluids along faults within the upper several kilometers of the subsurface in this instance must occur over a timescale of hours to days, and must be faster the timescales of biogenic reactions responsible for hydrogen consumption.

Aside from large-scale faulting, complex networks of small-scale (10-100 m throw)637 faults restricted to individual sedimentary layers, known as polygonal fault systems (PFS) 638 have been identified as having important impacts on basin-scale fluid flow. The impact 639 of PFS on fluid migration is debated within the literature, with authors attributing PFS 640 for both enhancing (e.g., Ireland et al. (2021); Cartwright (2011); Cartwright et al. (2003)) 641 and restricting (e.g., Xia et al. (2022); Andresen and Huuse (2011)) fluid flow. Whilst 642 PFS have been identified as a mechanism for seal bypass (Ireland et al., 2021), perme-643 ability may be effectively destroyed by clay smearing along fault planes and thus increase 644 the sealing capacity of PFS (Xia et al., 2022; Andresen & Huuse, 2011). Basin inversion, 645 fault reactivation or dewatering of host sediments may lead to the periodic opening of 646 fluid, and thus hydrogen migration pathways along impermeable PFS (Xia et al., 2022). 647 Whilst PFS must inevitably enhance the passage of fluids during their diagensis, it is 648 unlikely that PFS and microfractures provide substantial hydrogen migration pathways 649 over geological timescales during periods of tectonic quiescence. However, PFS may be 650 capable of both providing a mechanism for fluid communication and opening hydrogen 651 migration pathways on short geological timescales during periods of tectonic activity or 652 fault slip. 653

Whilst faulting and fluid flow have been extensively reviewed in the literature, their 654 impact on hydrogen migration have only recently gained significant attention. Early work 655 by by Wakita et al. (1980) hypothesised the production of hydrogen by fault movement, 656 based on measurements of elevated hydrogen concentrations (> 3% by volume) around 657 active fault zones in southwestern Japan compared to background measurements of  $\sim$ 658 0.5 ppm. Su et al. (1992) identified the potential reduction in strength of crystalline min-659 erals (e.g., calcite, dolomite, antigorite) due to hydrogen infiltration at low pressures, lead-660 ing to the weakening of rocks and initiation of faulting. Hydrogen gas measurements and 661 particle size distribution analyses by Niwa et al. (2011) within an active fault zone in-662 dicate that hydrogen gas mostly migrated in permeable fracture zones by advection with 663 groundwater. Firstov and Shirokov (2005) measured seven pulses of hydrogen discharge 664 against background levels in a fault zone trending parallel the Kuril - Kamchatka geostruc-665 tural zone, Russia, from 1999–2003. Hydrogen pulses preceding seismic events lasted 666 from 1.5-6 hours and were 2-14 times higher than measured background levels. Firstov 667 and Shirokov (2005) found that < 80% earthquakes with  $M_W \geq 5.6$  in the southern 668 Kamchatka region occurred within one month of measured hydrogen pulses and consid-669 ered such events as short-term earthquake precursors. In recent years, the migration of 670 natural hydrogen from deep crustal sources along kilometre-scale faults which penetrate 671 crystalline basement have been recorded in several locations across the world (e.g., Brazil, 672 France, Mali) (Prinzhofer et al., 2018; Deronzier & Giouse, 2020; Donzé et al., 2020; Rezaee, 673 2021; Frery, Langhi, et al., 2021; Lefeuvre et al., 2021, 2022). 674

#### 675 Surface seeps

Surface hydrogen emissions are associated with a wide range of geological conditions, including serpentinised mafic rocks, rift zones, Precambrian rocks, volcanic rocks, volcanic gases, geysers, hot springs, mud volcanoes and isolated seeps. The emission of natural hydrogen and gas from surface seeps has been recognised for millennia, e.g., the continuously burning Olympic flame at Mount Olympus, Turkey, dating back 2500 years and comprised of 7.5-11.3% H<sub>2</sub> (Hosgörmez, 2007). Other examples include 'Los Fue-

Type	Example	Characteristics	Reference
H <sub>2</sub> -rich	'Los Fuegos Eternos', Philippines	Noble gas signatures similar to air, $pH < 10$	Vacquand et al. (2018) Abrajano et al. (1990)
$N_2$ -rich	Semail, Oman	Ophiolitic, ${}^{3}\text{He}/{}^{4}\text{He} < 1R_{A}$	Sano et al. $(1993)$
$N_2-H_2-CH_4$	New Caledonia	Ophiolitic, mixing of $CH_4-H_2$ and $N_2$ -rich components	Deville and Prinzhofer (2016)
$\rm H_2-CH_4$	Mount Olympus, Turkey Zambales, Philippines	Ophiolitic, active subduction, ${}^{3}\text{He}/{}^{4}\text{He} > 1R_{A}$	Abrajano et al. (1988) Hosgörmez (2007)

 Table 3.
 Characteristics of surface abiogenic hydrogen seeps compiled from various references.

 $_{682}$  gos Eternos' (the eternal flames), discovered in the Philipines over two centuries ago with  $H_2$  concentrations of 41.4–44.5% (Vacquand et al., 2018; Abrajano et al., 1990). The distribution of surface hydrogen seeps across the globe are reviewed extensively by Zgonnik (2020), and can be separated into four broad categories (Table 3).

A common characteristic of most abiogenic natural hydrogen seeps is an associa-686 tion with ultramafic rocks, ophiolites and serpentinisation. The hydrolysis and oxida-687 tion of primary ferromagnesian minerals, such as olivine and pyroxenes, produces  $H_2$  over 688 a wide range of environmental conditions (Holm et al., 2015). Elevated isotopic signa-689 tures, i.e.  ${}^{3}\text{He}/{}^{4}\text{He} < 1R_{A}$ , within hydrogen-rich fluids encountered at such surface seeps 690 owe their provenance to primordial or deep mantle enrichment within ultramafic hydro-691 gen source rocks. However, whilst the advection of hydrogen entrained within ultramafic 692 rocks occurs over undoubtedly geological timescales, its liberation and transport to the 693 surface must depend on the parameters driving the  $H_2$ -forming serpentinisation reac-694 tion. Measurement of  $H_2$  degassing using in-situ gas chromatography and analysis of ex-695 perimental products using XRD, Raman and X-ray absorption spectroscopy under ser-696 pentinisation conditions ( $300^{\circ}C$  and 30 MPa) show a three stage process during the ser-697 pentinisation reaction: early (0-18 days), intermediate (18-34 days) and late (34-34 days)698 70 days). At the earliest stage, hydrogen is generated due to the crystallisation of mag-699 netite, with Fe-rich serpentine also formed as a reaction product of olivine, enstatite, clinopy-700 roxene and water. As the reaction progresses during the intermediate phase, hydrogen 701 is generated due to the formation of serpentine and clinopyroxene is absent from the re-702 action. During the final stages of serpentinisation, the serpentinisation front has effec-703 tively disappeared and hydrogen is generated due to the oxidation of Fe-rich serpentine 704 (Figure 4A, Table 4, Marcaillou et al. (2011)). The results of Marcaillou et al. (2011) 705 are further supported by Greenberger et al. (2015), who investigated the progression of 706 serpentinisation by mapping Fe oxidation states and analysing stable isotopes of carbon 707 and oxygen in carbonates to constrain the conditions of water-rock interaction during 708 serpentinisation. As groundwater migrates through a rock volume, the area of contact 709 between ultramafic source rocks (e.g., harzburgite, olivine) and migrating fluids, i.e. ser-710 pentinisation front, is greatest at early stages of serpentinisation. The serpentinisation 711 front is reduced as the reaction progresses, with maximum  $H_2$  generation at the earli-712 est stages of serpentinisation (Figure 5B Greenberger et al. (2015)). 713

Within subduction-related and ophiolitic terrains, all four types of seeps listed on 714 Table 3 are encountered. However, the nature of gases and gaseous mixtures emitted at 715 surface seeps is dependant on geodynamic context and the proportionality of mixing be-716 tween different fluids. Noble gases display signatures close to the value of air in  $H_2$ -rich 717 seeps, indicating that hydrogen gas emitted from ophiolitic settings is generated at shal-718 low depths within Earth's crust.  $N_2$ -rich seeps are notably associated with relatively high 719 contents of crustal  ${}^{4}$ He, and the source of N<sub>2</sub> is interpreted as derived mainly from meta-720 morphosed sediments located on the subducted crustal slab, below the ophiolitic units 721

Stage	Reaction	Reaction time
Early	Olivine + enstatite + cpx + $H_2O$ + $H^+ \rightarrow$ Serpentine + magnetite + $Mg^{2+}$ + $Ca^{2+}$ + $SiO_2$ + $H_2$	0-18 days
Intermediate	Olivine + enstatite + $H_2O$ + $SiO_2$ + $Mg^{2+} \rightarrow$ Serpentine + magnetite + $H^+$ + $H_2$	18 - 34 days
Late	$ \begin{array}{l} {\rm Enstatite} + {\rm cpx} + {\rm Fe-rich\ serpentine} + {\rm H}_2{\rm O} \rightarrow \\ {\rm Mg-rich\ serpentine} + {\rm magnetite} + {\rm calcic\ silicate} + {\rm H}_2 \end{array} $	34-70 days

Table 4. Stages of the serpentinsation reaction (Marcaillou et al., 2011)

(Deville & Prinzhofer, 2016; Vacquand et al., 2018).  $H_2$ -CH<sub>4</sub>-rich gas seeps are typi-722 cally characterised by mantle-like C and noble gas characteristics, as evidenced by mea-723 surements from several locations including the Zambales ophiolite, Phillipines and New 724 Caledonia. Fluid communication within fracture networks and mixing between  $N_2$ -rich 725  $H_2-CH_4$ -rich end members is the most likely cause for  $N_2-H_2-CH_4$ -rich gas seeps (Abrajano 726 et al., 1990; Deville & Prinzhofer, 2016). The flux of deep gas into a shallow aquifer iso-727 lated from direct equilibrium with the atmosphere and fractionation during subsequent 728 degassing is suggested by Deville and Prinzhofer (2016) as the simplest explanation for 729 observed <sup>20</sup>Ne and <sup>36</sup>Ar concentrations. A schematic diagram of the distribution of dif-730 ferent gas seeps with an ophiolitic terrain is shown Figure 5A. 731

The timescale of subsurface gas migration is also dependent on parameters that con-732 trol groundwater transport properties. Experimental work by Lamadrid et al. (2017) used 733 synthetic fluid inclusions as micro-reactors in olivine crystals to monitor serpentinisa-734 tion rates in-situ and at serpentinisation conditions  $(280^{\circ}C)$ . Serpentinisation rates were 735 strongly influenced by aqueous fluid salinity, with evidence of reaction after 5 days de-736 creasing from 50% to zero as aqueous fluid salinity of synthetic inclusions increased from 737 1-10 wt%. The time taken to observe the first evidence of reaction for salinity exper-738 iments at 10 wt% was 120 days, with the average rate of reaction being two orders of 739 magnitude lower than experiments conducted with 1-3.5 wt% salinity (Lamadrid et 740 al., 2017). The results of Lamadrid et al. (2017) support those of Rouméjon and Can-741 nat (2014), and indicate that the forward reaction requires continual influx of a lower 742 salinity aqueous fluid (seawater) to dilute the serpentinisation fluid and allow serpen-743 tinisation of olivine to continue. Hydrogen migration and emission from surface seeps 744 is dependent on the interplay between generative processes (e.g. serpentinisation) and 745 destructive processes (e.g. microbial consumption). Important factors which influence 746 the balance between generation and consumption include the water/rock ratio of migra-747 tion or injected fluids (e.g. stimulated hydrogen generation, see next section), fluid chem-748 istry and the formation of Fe(III)-bearing secondary phases (Templeton et al., 2024). 749

The results of Lamadrid et al. (2017) and Rouméjon and Cannat (2014) compliment experimental results of hydrogen dissolution by Iglauer et al. (2012) and Hosseini et al. (2022) (see earlier sections), and lead to an overall consensus that the salinity of the carrier fluid (i.e., groundwater) is a major controlling factor in both the amount of hydrogen gas generated and its rate of transport within sedimentary basins and ophiolitic terranes.

#### Stimulated hydrogen generation

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Over the last few years, the concept of stimulated geological hydrogen, also known
 as 'orange' hydrogen, has gained significant momentum (e.g., Osselin et al. (2022); Tem pleton et al. (2024)). Hydrogen generation may be stimulated by the injection fluids into
 target rock formations rich in reactive Fe(II)-bearing minerals to promote the overar-

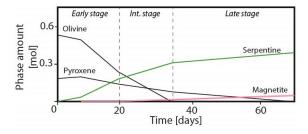


Figure 4. Evolution of mineral phases during serpentinisation reactions shown on Table 4. Serpentinisation products are shown in red (serpentinite) and red (magnetite). Magnetitie crystallises first and is responsible for  $H_2$  generation during the early stage. During the intermediate phase, clinopyroxene is absent due to a lack of enrichment of calcium. During the late stage, olivine is no longer the reactive species and is replaced by serpentine formed during the early phase as the reactant. Late stage serpentine is Mg-rich and distinct from early stage serpentine. Modified from (Marcaillou et al., 2011).

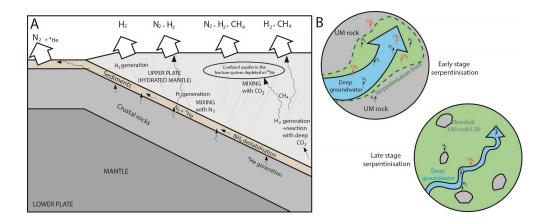


Figure 5. A = Conceptual diagram explaining gas compositions found at ophiolite-related seeps. B = Conceptual diagram showing progression of serpentinisation reaction and fluid migration pathways. UM rock = ultramafic rock. As serpentinisation of a rock volume progresses, water/rock ratios decrease significantly due to volume expansion and reduced permeability of serpentine (O'hanley, 1992). During late stage serpentinisation,  $H_2$  generation decreases significantly due to reduced contact between serpentinisation fluids (e.g. deep groundwater) and reactants (unreacted ultramafic rock and early stage Fe-rich serpentine). Modified from (Marcaillou et al., 2011; Greenberger et al., 2015; Deville & Prinzhofer, 2016; Vacquand et al., 2018).

ching reaction of  $2\text{FeO}(\text{rock}) + \text{H}_2\text{O} \rightarrow \text{Fe}_2\text{O}_3(\text{rock}) + \text{H}_2$ . Most commonly, these are found within ultramafic rocks or ophiolites. H<sub>2</sub> is then extracted by recirculating injected fluids to the surface, if the required (bio) geochemical conditions for rapid hydrogen production have been met (Templeton et al., 2024). Development of this relatively new concept has been spurred by the recent U.S Department of Energy announcement of funding to support research into the production of geologic hydrogen through stimulated mineralogical processes (of Energy, 2024).

However, it is worth noting that the term 'orange' hydrogen has a broader mean-768 ing within the literature, and includes chemical processes that break down traditional 769 hydrocarbons and biofuels (e.g., crude oil, natural gas, gasoline, biogas, etc.) into hy-770 drogen with no carbon dioxide byproduct (Neelameggham et al., 2022). A notable re-771 cent study that explores this includes the generation of  $H_2$  from  $H_2S$  by thermal split-772 ting, a process which is potentially  $\sim 38\%$  more economically viable than green hydro-773 gen production (Nova et al., 2023). As this review is focussed on subsurface hydrogen 774 and migration, we do not explore industrial and chemical processes capable of produc-775 ing hydrogen and avoid the term 'orange' hydrogen to prevent confusion with stimulated 776 hydrogen generation, i.e. due to the injection of fluids within subsurface geological for-777 mations. 778

Osselin et al. (2022) assessed the reactive percolation of a NaHCO<sub>3</sub>-rich brine at 779 160°C and 280°C in natural serpentinite cores in order to study the dynamic competi-780 tion between serpentinisation and carbonation of ultramafic formations. Their results 781 are used to suggest that up to 100 trillion tonnes of  $H_2$  could be produced from Fe(II)-782 bearing rocks near Earth's surface. Similar to the serpentinisation reactions shown on 783 Table 3, the experiments of Osselin et al. (2022) were completed over a time period of 784 days. However, it is important to note that permeability reductions of several orders of 785 magnitude were observed and attributed to the precipitation of carbonates in the main 786 percolation paths. In mafic rocks, permeability is produced by fractures and the reduc-787 tion in fracture permeability with time can be approximated by an exponential function: 788  $k_f = k_{f0} e^{-t/t^*}$ , where  $k_{f0}$  is the starting permeability of the modelled interval and  $t^*$ 789 is the characteristic decay time of the best fit exponential function in hours (Farough 790 et al., 2016). Precipitation of carbonates over silicates is favoured due to fast reaction 791 kinetics in comparission to flow rates. Osselin et al. (2022) conclude that the spatio-temporal 792 lengthscales associated with the different chemical reactions are directly linked to the 793 ratio of chemical reaction rate and transport (also known as the Damkokhler number) 794 and to the type of reaction regime (transport-limited vs reaction-limited). They inter-795 pret their results as a dynamic interplay between dissolution and precipitation, controlled 796 by the local flow rate and the local pore geometry. Osselin et al. (2022) note that the 797 complex pore size distribution in natural rocks leads to very different behavior even for 798 a homogeneous mineralogy. Hence, the work of Osselin et al. (2022) demonstrates the 799 significance of changes in material properties of host rocks due to the flow of fluids used 800 to liberate hydrogen from mafic minerals, which will occur on a timescale of days. 801

More recently, Templeton et al. (2024) discusses factors that influence the timescales 802 of fluid migration through the Semail Ophiolite, Oman, such as the hydraulic conduc-803 tivity of partially-hydrated peridotites, the extent of fracturing, and the geochemical dy-804 namics of the subsurface environment. Fluid flow in the most shallow and fissured rocks 805 occurs in transmissive zones located within 50 m of the surface, indicating that fluid dy-806 namics are highly heterogeneous. Templeton et al. (2024) also describes that some zones 807 are most sensitive to conductive channels, such as partially mineralised fractures, whereas 808 other zones are supplied from rocks above and below. This complexity, including the pres-809 ence of fractures partially filled with secondary minerals produced from both modern and 810 ancient water/rock reactions, suggests that fluid migration rates and, by extension, hy-811 drogen migration rates through the crust are highly variable and dependent on local ge-812 ological conditions. In contrast to Osselin et al. (2022), the analyses of Templeton et al. 813

(2024) were conducted at low temperature (25°C), demonstrating that the sensitivity
 of fluid flowpaths are not restricted to high-temperature systems only or within deeper
 parts of geological basins.

#### Adsorption

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The migration of hydrogen may be prevented by the physical adsorption of  $H_2$  molecules 818 to the surface of minerals, particularly clays. This is demonstrated by Truche et al. (2018), 819 who demonstrate hydrogen enrichment of < 500 ppm (0.25 mol/kg of rock) within organic-820 poor (< 0.5 wt % total organic carbon) clays composed of illite, chlorite and kaolinite 821 from the Cigar Lake uranium ore deposit, Canada. Furthermore, recent experiments by 822 L. Wang, Cheng, et al. (2023) demonstrate that hydrogen adsorption is significantly in-823 fluenced by the pore structure and specific surface area of the clay minerals, with a no-824 table increase in hydrogen adsorption capacity under high pressure and a decrease at higher 825 temperatures, independent of the clay mineral type. L. Wang, Cheng, et al. (2023) show 826 that whilst montmorillonite and chlorite only adsorb hydrogen on their external surface, 827 palygorskite and sepiolite can adsorb hydrogen on both the bulk phase and the exter-828 nal surface. However, adsorption capacity may be compromised by a range of factors, 829 including the presence of water and other adsorbed gases (e.g., CO<sub>2</sub>, CH<sub>4</sub>, He) which 830 may compete for adsorption sites. 831

Whilst significant volumes of hydrogen may accumulate within clay-rich rocks due 832 to adsorption, the degree to which this affects overall hydrogen migration on the basin-833 scale is, however, an open question. Truche et al. (2018) estimate that 4 - 17% of H<sub>2</sub> 834 produced by water radiolysis over the 1.4 Ga lifetime of the Cigar Lake uranium ore de-835 posit is trapped in the surrounding clay alteration halos, thus leaving 83-96% hydro-836 gen unaccounted for. Despite adsorption, hydrogen will migrate due to advection along 837 fractures and by diffusion, with breakthrough times varying on the scale of years per me-838 tre for most rocks, including clays (Table 2). Hence, it is reasonable to assume the timescale 839 over which adsorption sites become occupied must be fairly rapid and on the scale of years 840 to thousands of years for volumes of clay-rich rocks typically within sedimentary basins. 841 In reality, it is likely that whilst adsorption may trap significant and potentially economic 842 volumes of hydrogen within clay-rich rocks and remain stable over geological timescales, 843 the majority of hydrogen within sedimentary basins remains transient and mobile. 844

#### 845 Microbes

Microbial reactions within host rocks and sediments are important moderators of 846 hydrogen flow in the subsurface. Microbial activity may lead to both the generation and 847 loss of hydrogen as it migrates through a reservoir, as summarised on Table 5. Micro-848 bial reactions are dependent on many factors, such as environment (e.g., pH, salinity). 849 iron ( $Fe^{3+}$ ) content of host rocks, groundwater recharge and the presence of other re-850 duced gases from deeper in Earth (Anderson et al., 1998; Stevens & McKinley, 2000). 851 There is a growing consensus that subsurface microbial communities are independent of 852 photosynthesis for carbon and hydrogen supply, and are primarily or completely depen-853 dent on abiotc hydrogen sources in various geological settings as an energy source (Kotelnikova 854 & Pedersen, 1998; Takai et al., 2004; L. H. Lin et al., 2005; McCollom & Amend, 2005; 855 Nealson et al., 2005; Escudero et al., 2018; Gregory et al., 2019). These microorganisms 856 consist of Bacteria and Archaea and exist in great abundances within the subsurface, with 857  $\sim 10^4 - 10^8$  cells/gram of rock up to several km deep and  $\sim 2 - 6 \times 10^{29}$  cells within 858 the continental subsurface (Dutta et al., 2018; Magnabosco et al., 2018). For microbial 859 life to survive, temperature limits must lie between  $-15^{\circ}$ C to  $+121^{\circ}$ C, corresponding 860 to depths of up to 3.5 - 4.5 km beneath Earth's surface at normal geothermal gradi-861 ents of  $30\pm5^{\circ}$  C/km. Temperature, pressure and salinity are important factors for the 862 prevalence of single-celled microorganisms within the subsurface that are responsible for 863 using hydrogen in their metabolism. Whilst there is little to no information about pres-864

sure or brine salinity thresholds, neutral pH (pH = 6 - 7) conditions generally correspond to the greatest abundance and diversity of microbial life. However, microbial life may exist within the pH range 0 - 11 (Dopffel et al., 2021).

In terms of migration, literature on the impact of microbial reactions on hydrogen 868 is scant, with the overwhelming majority of recent and legacy research focused on as-869 sessing microbial hydrogen generation, consumption and associated environmental risk 870 (e.g.  $H_2S$  generation). As generative and destructive processes alter the amount of hy-871 drogen within a subsurface system, it is conceptually reasonable to consider their role 872 873 in migration as a moderator of hydrogen flow, whereby the rate of hydrogen transport through a medium influenced by microbes will depend on the kinetics and rates of microbial-874 hydrogen reactions. Thus, microbial reactions represent an important sink in the migra-875 tion pathway of hydrogen from depth to the surface. Harris et al. (2007) present one of 876 the few experimental assessments of microbial community metabolism directly within 877 a groundwater environment, and estimate hydrogen consumption rates in-situ injection/withdrawal 878 tests conducted in two geochemically varying, contaminated aquifers. The results of Harris 879 et al. (2007) show that first-order hydrogen consumption rates varied from  $0.002 \text{ nM h}^{-1}$ 880 for an uncontaminated, aerobic site to  $2.5 \text{ nM} \text{ h}^{-1}$  for a contaminated site where sul-881 phate reduction was a predominant process. Notably, the hydrogen consumption rate 882 reduced to zero within a denitrifying zone and in the presence of air or an antibiotic mix-883 ture, thus highlighting potential sensitivity to environmental perturbations on field mi-884 crobial activities on the timescale of several hours (Harris et al., 2007). These results may 885 be interpreted as meaning the degree to which subsurface microbial activity moderates 886 hydrogen flow may vary on timescales relevant to groundwater flow through the host rock 887 or sediment, e.g. the acidity and salinity of pore fluids in top soils may vary on timescales 888 of hours - days, whereas deeper rocks and aquifers may attenuate environmental signals 889 over thousands of years or longer. Interestingly, Templeton et al. (2024) demonstrate that 890 the alteration of the chemical composition of fluids introduced into geological formations 891 during stimulated hydrogen generation is pivotal for the optimal generation of Fe(III)-892 enriched secondary mineral phases. Templeton et al. (2024) argue that modifications in 893 fluid chemistry should be strategically engineered to concurrently diminish the micro-894 bial uptake of H<sub>2</sub> within the stimulated region, whilst maintaining elevated capacities 895 for biogenic hydrogen assimilation in the shallow groundwater systems. The assimila-896 tion of biogenic hydrogen into shallow groundwater is essential as this will be saturated 897 with oxidising agents such as nitrate, sulfate, and dissolved inorganic carbon. The rec-898 ommendations of Templeton et al. (2024) serves to mitigate the risk of unintentional hy-899 drogen emissions into the atmosphere, where it contributes as an indirect greenhouse gas 900

The ability of host rocks to sustain microbial activity on geological timescales may 901 also be dependent on whether the rocks contain sufficient reduced iron and other depen-902 dent nutrients (Gregory et al., 2019). The compilation of experimental results by Roden 903 and Jin (2011) show that the relationship between microbial yield and the free energy 904 of aerobic and anaerobic metabolism of hydrogen in soils and sediments follow the same 905 linear trend as other compounds, such as glucose, ethanol, formate, acetate, lactate, pro-906 pionate, butyrate. Roden and Jin (2011)'s results indicate that it is possible to estimate 907 microbial yield values within a factor of 2 (i.e. error  $= \pm 100\%$ ) using a simple linear re-908 lationship, although it should be noted that errors are greatest for hydrogen metabolism. 909 The results of Harris et al. (2007) and Roden and Jin (2011) indicate that it may be pos-910 sible to quantify the role played by subsurface microbial activity as a moderator of hy-911 drogen transport, however further research is required in this area to determine the re-912 lationship between subsurface environmental change and the timescale of hydrogen mi-913 gration. 914

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Process	Reaction	$\Delta G^{0'}$ kj mol <sup>-1</sup> H <sub>2</sub>	Reference
Fermentation	Multiple pathways that breakdown large organics into smaller organics, e.g., mixed acid fermentation e.g., $C_6H_{1,2}O_6 + 4H_5O \rightarrow 2CH_3COO_7^- + 2HCO_7^- + 4H^2 + 4H_5$	-54.0 - 24.2	(Thauer et al., 1977; Conrad, 1999) (Schink et al., 2002; Gregory et al., 2019)
Nitrogen fixation (nitrogenase activity)	$N_2 + 8H + 8e^- (Fd_{red}) \rightarrow 2NH_3 + H_2 (+Fd_{ox})$	-18.1	(Thauer et al., 1977; Conrad, 1999) (Schink et al., 2002; Gregory et al., 2019)
Anaerobic carbon monoxide oxidation	$\rm CO+H_2O\rightarrow CO_2+H_2$	-19.9	(Thauer et al., 1977; Conrad, 1999) (Schink et al., 2002; Gregory et al., 2019)
Phosphite oxidation	$\rm H_3PO_3+H_2O\rightarrow H_3PO_4+H_2$	-46.3	(Thauer et al., 1977; Conrad, 1999) (Schink et al., 2002; Gregory et al., 2019)
Acetate oxidation	$\frac{1}{4}$ CH3COO <sup>-</sup> + 14 H <sup>+</sup> + $\frac{1}{2}$ H <sub>2</sub> O $\rightarrow$ H <sub>2</sub> + $\frac{1}{2}$ CO <sub>2</sub>	23.7	(Thauer et al., 1977; Conrad, 1999) (Schink et al., 2002; Gregory et al., 2019)
	Microbial hydrogen consumption	n	
Hydrogenotrophic methanogenesis	$HCO3^{-} + H_{2} + \frac{1}{4}H^{+} \rightarrow \frac{1}{4}CH_{4} + \frac{3}{4}H_{2}O$	-33.9	(Schink et al., 2002; Greening et al., 2016) (Gregory et al., 2019; Dopffel et al., 2021)
Acetogenesis	$\frac{1}{2}HCO_3^- + H_2 + \frac{1}{4}H^+ \rightarrow \frac{1}{4}CH_3COO^- + 2H_2O$	-26.1	(Schink et al., 2002; Greening et al., 2016) (Gregory et al., 2019; Dopffel et al., 2021)
Sulphate reduction	$\frac{1}{4}SO_4^{2-} + H_2 + \frac{1}{4}H^+ \rightarrow \frac{1}{4}HS^- + H_2O$	-38.0	(Schink et al., 2002; Greening et al., 2016) (Gregory et al., 2019; Dopffel et al., 2021)
Sulphur reduction	${ m H_2+S}  ightarrow { m H_2S}$	-33.1	(Schink et al., 2002; Greening et al., 2016) (Gregory et al., 2019; Dopffel et al., 2021)
Iron(III) reduction	$2 FeOOH + H_2 + 4 H^+ \rightarrow 2 Fe^{2+} + 4 H_2 O$	-228.3	(Schink et al., 2002; Greening et al., 2016) (Gregory et al., 2019: Donffel et al., 2021)
Aerobic hvdrogen oxidation	$H_2 + \frac{1}{2}O_2 \rightarrow H_2O$	-237	
Dehalorespiration	Halogenated compounds + $H_2 \rightarrow$ dehalogenated compounds + 2HCl	-230 to -187	
Fumarate respiration	$H_2 + fumarate \rightarrow succinate$	-86.2	(Schink et al., 2002; Greening et al., 2016) (Gregory et al., 2019; Dopffel et al., 2021)
Denitrification	$rac{2}{5} m NO3^{-1}+~H_2+rac{2}{5} m H^+ ightarrow rac{1}{5} m N2+rac{6}{5} m H_2O$	-240.1	(Schink et al., 2002; Greening et al., 2016) (Gregory et al., 2019; Dopffel et al., 2021)

#### <sup>915</sup> Diffusion versus advection

A property that distinguishes hydrogen from other fluids within geological basins 916 is its ability to migrate through rocks occurs via diffusion or advective processes from 917 human to geological timescales. In reality, it is reasonable to assume that hydrogen flow-918 paths will be a function of both diffusive and advective processes (e.g., Figure 7) and over-919 all fluid chemistry. However, the ability to predict the behaviour of hydrogen in the sub-920 surface similar to other fluids, such as hydrocarbons and groundwater, remains largely 921 unresolved. The results of Mathiesen et al. (2023); Hutchinson et al. (2024) indicate that 922 923 pore throat diameter, and therefore capillary entry pressure, exert a primary control on the mode of hydrogen migration, with increasing advective dominance at larger pore throat 924 sizes (i.e. sandstones) and increasing diffusive dominance at smaller pore throat sizes (i.e. 925 shales, evaporites). Lodhia and Clark (2022) approximate hydrogen mobility and buoy-926 ancy by solving the Darcy flow equation using a series of steps and use lithological pa-927 rameters representative of general rock types. Their method may be applied to estimate 928 the basin-scale maximum vertical velocity,  $v_{max}$ , of pure advective H<sub>2</sub> gas-flow as the 929 product of mobility and buoyancy, as it accounts for upscaling (due to macro-scale fea-930 tures such as faults and fractures), geological regime (e.g., normal, overpressured or hy-931 drostatic) and geothermal gradient (Lodhia, 2023). However, there is no clear relation-932 ship between  $v_{max}$ , permeability and porosity (Lodhia & Peeters, 2024). The transition 933 between diffusive and advective flow for pure and multiphase  $H_2$ , known as the trans-934 slip flow boundary, may be calculated using a charecteristic Knudsen number value of 935 0.1 (Hutchinson et al., 2024; Sakhaee-Pour & Alessa, 2022; Roy et al., 2003). We apply 936 data from Strauch et al. (2023) and Fick's first law to calculate diffusion velocities,  $v_{diff}$ , 937 for dry and wetted sandstone, evaporties and clay. Due to a lack of data in the litera-938 ture,  $v_{diff}$  is not calculated for hydrogen migration in carbonates. Calculated basin-scale 939  $v_{max}$  and  $v_{diff}$  values are shown on Figure 6A–B and indicate diffusive velocities are 940 several orders of magnitude smaller than advective velocites and that  $v_{max}$  decreases ex-941 ponentially with increasing clay content across all rock types. Advective flow of  $H_2$  be-942 comes less effective at shallow depths (< 400 m) due to the rapid increase in mean free 943 paths. Under multiphase subsurface conditions, advective flow will be impaired due to 944 water occupying and restricting pore space, causing the diffusion-advection boundary 945 to be displaced to larger pore throat sizes (Hutchinson et al., 2024). It is also reason-946 able to assume that diffusion velocities for hydrogen in carbonates will follow a similar 947 trend to other rock types and be several orders of magnitude smaller than  $v_{max}$  values. 948 Hence, calculations of  $v_{max}$ ,  $v_{diff}$  and the position of the diffusion-advection boundary 949 provide an estimate for the timescale and mode of hydrogen migration at different depths 950 in shallow sedimentary basins for a range of rock types. We hypothesise that hydrogen 951 migration is dominated by diffusion at shallow depths and operates on a timescale of <952 0.5 cm year - 1 for clastic rocks, shale, evaporites and probably carbonates, which decrease 953 by an order of magnitude for water-saturated rocks or sediments. At intermediate depths, 954 the boundary between diffusive and advective flow marks a peak in migration velocity 955 on the timescale of > 0.1 and < 10 m year<sup>-1</sup> for most clastic and carbonate rocks, with 956 the exception of dolomites which have advective velocities from  $\sim 100-1000$  m year<sup>-1</sup>. 957 The boundary between diffusive and advective hydrogen flow is not uniform across dif-958 ferent rock types, such that multiphase conditions cause displacement towards more coarsely 959 grained rocks and an increased depth envelope for diffusive migration. Advective veloc-960 ities decrease with depth due to the reduction in pore throat sizes and corresponding in-961 crease in capillary entry pressures required for hydrogen flow. Furthermore, whilst in-962 creased clay content will affect the timescale of advective flow, the effect on diffusion is 963 minimal. 964

Future experimental research should focus on improving understanding of porosity and water saturation relationships relevant for hydrogen, as due to a lack of data in the literature, Lodhia and Clark (2022) apply estimates of oil-water and gas-water sys-

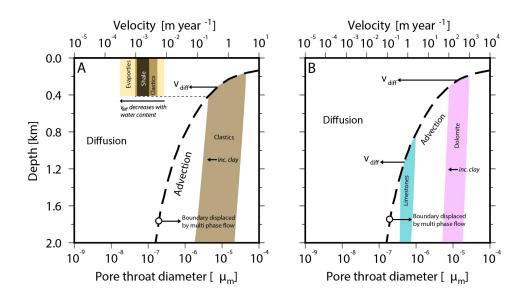


Figure 6. Maximum vertical velocity  $(v_{max})$  and diffusion velocity  $(v_{diff})$  calculated for various rock types using the method of Lodhia and Clark (2022) and data from Hantschel and Kauerauf (2009) and Strauch et al. (2023), respectively. A = velocity  $v_{max}$  and  $v_{diff}$  for clastic rocks. Dashed horizontal line indicates the maximum depth of transition between advective and diffusive migration for pure H<sub>2</sub>.  $v_{diff}$  decreases with increasing water content due to a rapid increase in mean free paths (i.e., increased collisions between H<sub>2</sub> and water molecules) whilst  $v_{max}$ decreases with increasing clay content due to increased capillary entry pressure associated with decreasing pore throat diameter (e.g. Hutchinson et al. (2024)). B =  $v_{max}$  for carbonate rocks. We do not calculate  $v_{diff}$  values for carbonates due to a lack of data within the literature, however assume these to follow a similar trend to clastic rocks and be several orders of magnitude smaller than corresponding  $v_{max}$  values (see labelled arrows). Curved dashed line = diffusionadvection boundary for pure H<sub>2</sub> calculated using a Knudsen number of 0.1 (Hutchinson et al., 2024). This boundary is displaced to larger pore throat sizes for multiphase subsurface flow as indicated by circled arrows.

tems. Future experimental studies may also focus on testing both the robustness of Lodhia and Clark (2022)'s approximations and our hypotheses.

#### 970 Discussion

The migration of hydrogen through the subsurface is a topic seldom addressed directly, yet is critical for exploration and geological storage investigations. To understand the dynamics of the subsurface hydrogen cycle within sedimentary basins and Earth's surface, we must take a holistic view of its supply, emission and intermediate processes.

Within geological basins, long-term hydrogen supply from the radiolysis of water 975 within crystalline basement, Archean - Proterozoic cratonic rocks and other hydrogen 976 abundant mafic igenous rocks will remain steady over geological timescales. However, 977 hydrogen migration pathways will be disproportionately affected by specific processes 978 operating within small regions within sedimentary basins and Earth's crust, such as mi-979 crobial reactions in soil or regolith, advection of fluids along faults and 'trapping' on timescales 980 relative to humans by wet or evaporitic sediments. Environmental factors, such as salin-981 ity and temperature may change the dynamics of subsurface hydrogen systems rapidly. 982 for example a saline aquifer changing from a barrier to a carrier due to an influx of fresh 983 meteoric water following heavy rainfall. Hydrogen supply rates within generative sys-984 tems will be primarily controlled by the availability of fresh water, such as rainfall on 985 ophiolitic systems or groundwater contact with buried igenous rocks. 986

The rate of diffusive migration of hydrogen from crystalline rocks into surround-987 ing sediments will operate on timescales of 1-1000 years (Figure 2), and be controlled 988 primarily by grain size and temperature. Experimental results from the literature indi-989 cate that native hydrogen entrained within the mineral structure of crystalline rocks within 990 the shallow Earth may diffuse on geological timescales from the most common rock form-991 ing minerals. Geochemical data obtained by Parnell and Blamey (2017) indicate that common felsic lithologies, such as granites, gneiss and conglomerates of Archean - Pro-993 terozoic (> 1600 Ma) age consistently contain an order of magnitude greater hydrogen 994 in their entrained fluid than very young (< 200 Ma) granites. Parnell and Blamey (2017) 995 found that sedimentary rocks containing clasts of old basement also included a greater 996 proportion of hydrogen than young granites and hypothesise that a signature of hydro-997 gen in the basement could be conferred to the sediment and that modern sediment de-998 rived from old and young basement retains the signature of more or less hydrogen, re-999 spectively (Figure 7). It should be noted however that the experimental results summarised 1000 by Parnell and Blamey (2017) refer to bulk lithologies whereas those of Farver (2010)1001 refer to individual minerals (e.g. olivine and quartz). Furthermore, the preservation of 1002 high hydrogen abundances within fluid inclusions and mineralised veins in ancient gran-1003 ites has been observed (Bourdet et al., 2023). Hence, diffusion from enriched Archean 1004 - Proterozoic crystalline basement and their derived sedimentary products, e.g., conglom-1005 erates, may supply a 'background' hydrogen flux to overlying sedimentary basin rocks 1006 on geological timescales. This is consistent with the widely documented observation of 1007 higher hydrogen fluxes in sedimentary basins in continental cratonic regions underlain 1008 by Archean - Proterozoic basement (e.g. Zgonnik (2020); ? (?)). In the case of rapidly 1009 cooled upper mantle rocks, e.g. MORBs, volcanic glasses, pillow lavas, however, grain 1010 sizes may be many orders of magnitude smaller than their continental counterparts and 1011 within the nanometre scale (e.g., Schlinger et al. (1988)). Hence, hydrogen diffusivity 1012 in rapidly cooled crystalline rocks and at MOR settings will be significantly faster than 1013 in continental settings and potentially only a few orders of magnitude slower than the 1014 lower temperature ranges of diffusivity experiments, i.e. 100 Ka - Ma or faster. This is 1015 significant, since the age of most oceanic crustal rocks is < 60 Ma (Seton et al., 2020), 1016 hydrogen diffusion within oceanic crustal rocks will operate on the same timescale as the 1017 age of rocks themselves and provide a mechanism for the degassing of mantle hydrogen 1018 to the surface and oceans. Figure 7 summarises the characteristic migration timescales 1019

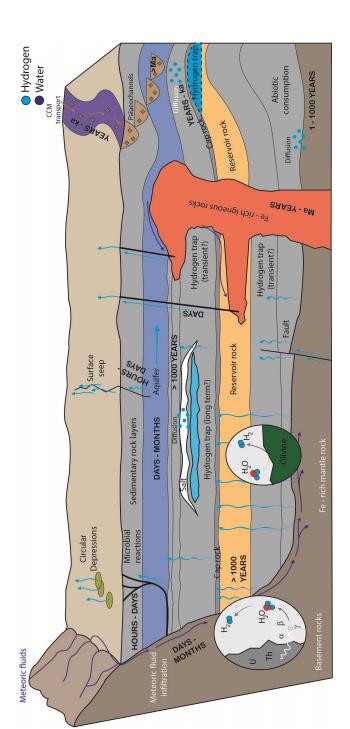


Figure 7. Conceptual method for modelling workflow, migration and indicative timescales through sedimentary basins. Blue arrows indicate advective transport rocks. Diffusion through geological barriers to hydrogen migration, e.g. salt and igneous intrusions may limit hydrogen migration by thousands to millions of years, (e.g. Archean - Proterozoic basement) is believed to be responsible for the provenance of high hydrogen concentrations in some sedimentary basin rocks (Parnell &Blamey, 2017). The transport of CCM (brown diamonds and arrows) by fluvial systems and accumulation in palaeochannel deposits is shown for conceptual purof hydrogen. Fe-rich lithologies (e.g. granites) characterise Archean - Proterozoic continental crystalline basement within cratonic regions and igneous intrusive depending on mineralogy and lithology (Farver, 2010; Parnell & Blamey, 2017). Continental clastic material (CCM) derived from terrains enriched in hydrogen poses. Modified from Hand (2023)

for hydrogen transport through different parts of a sedimentary basin as described in thisarticle.

Whilst faulting and fluid flow have been extensively reviewed in the literature, their 1022 impact on hydrogen migration have only recently gained significant attention. Obser-1023 vations of hydrogen pulses prior to seismic activity are not well documented, but can be 1024 analogous to increased  $CO_2$  emissions possibly due to enhanced porosity of the soil due 1025 to faulting, and accelerated water rock interactions and soil gas emission within the fault 1026 zone (e.g., Z. Liu et al. (2023)). Early work by by Wakita et al. (1980) hypothesised the 1027 production of hydrogen by fault movement, based on measurements of elevated hydro-1028 gen concentrations (> 3% by volume) around active fault zones in southwestern Japan 1029 compared to background measurements of  $\sim 0.5$  ppm. Su et al. (1992) identified the po-1030 tential reduction in strength of crystalline minerals (e.g., calcite, dolomite, antigorite) 1031 due to hydrogen infiltration at low pressures, leading to the weakening of rocks and ini-1032 tiation of faulting. Hydrogen gas measurements and particle size distribution analyses 1033 by Niwa et al. (2011) within an active fault zone indicate that hydrogen gas mostly mi-1034 grated in permeable fracture zones by advection with groundwater. Firstov and Shirokov 1035 (2005) measured seven pulses of hydrogen discharge against background levels in a fault 1036 zone trending parallel the Kuril - Kamchatka geostructural zone, Russia, from 1999– 1037 2003. Hydrogen pulses preceding seismic events lasted from 1.5-6 hours and were 2-1038 14 times higher than measured background levels. Firstov and Shirokov (2005) found 1039 that < 80% earthquakes with  $M_W \ge 5.6$  in the southern Kamchatka region occurred 1040 within one month of measured hydrogen pulses and considered such events as short-term 1041 earthquake precursors. In recent years, the migration of natural hydrogen from deep crustal 1042 sources along kilometre-scale faults which penetrate crystalline basement have been recorded 1043 in several locations across the world (e.g., Brazil, France, Mali) (Prinzhofer et al., 2018; 1044 Deronzier & Giouse, 2020; Donzé et al., 2020; Rezaee, 2021; Frery, Langhi, et al., 2021; 1045 Lefeuvre et al., 2021, 2022). 1046

# Recent hydrogen discoveries and possible importance to migration path ways

Following the landmark discovery of natural hydrogen at Bourakebougou in Mali by Prinzhofer et al. (2018), subsequent research has uncovered numerous hydrogen deposits worldwide. In this section, we review several recent discoveries and examine potential relationships between gas composition and fluid migration pathways.

A regional geochemistry study by Lévy et al. (2023) in Albania and Kosovo focused 1053 on natural springs, revealing a site in northern Kosovo with a hydrogen concentration 1054 of 16%, pointing to serpentinisation of peridotites as a hydrogen source. Notably, this 1055 study found no correlation between hydrogen and helium concentrations but did observe 1056 substantial organic and crustal contributions ( $CH_4$  and  $N_2$ ). Contrastingly, recent data 1057 from South Australia and direct measurements from the Bulqizë chromite mine in Al-1058 bania report significant hydrogen outgassing, with  $H_2$  concentrations > 80% and vary-1059 ing but minor N<sub>2</sub>, CO<sub>2</sub> and CH<sub>4</sub> components (Goh, 2023; Hydrogen, 2023; Truche et al., 1060 2024; Yeo, 2023). Interestingly, whilst the results of Lévy et al. (2023) indicated no a clear 1061 connection between hydrogen and helium in the context of serpentinisation, Karolytė et 1062 al. (2022) documented He-rich hydrocarbon gases in South Africa's Witwatersrand Basin, 1063 where the scarcity of mafic and ultramafic minerals capable of serpentinisation suggests 1064 radiolytic hydrogen production as the predominant mechanism. Despite the detection 1065 of hydrogen alongside  $CH_4$  and  $N_2$  in both scenarios, the system described by Karolytė 1066 et al. (2022) is not associated with hydrocarbon source rocks. Transportation over vast 1067 distances and a significant degree of interaction with groundwater dilute He and lower H/He ratios (Ballentine & Lollar, 2002; Gilfillan et al., 2008). Hence, observed He con-1069 centrations in the Witwatersand Basin by Karolytė et al. (2022) may be explained by 1070

long periods of quiescence for He accumulation within a closed system, characterised by
 its hydrogeological systems being isolated and He preservation.

These examples highlight the rapid conversion of hydrogen to methane in surface 1073 environments and the influence of gas composition on migration pathways, particularly 1074 the impact of helium content on gas mixture compositions over geological timeframes. 1075 Prinzhofer and Cacas-Stentz (2023) present theoretical analyses suggesting that advec-1076 tive leakage of hydrogen-bearing gases out of subsurface reservoirs affects their overall 1077 composition, leading to an increase in nitrogen and methane content at the expense of 1078 1079 hydrogen. Their findings suggest a dichotomy where hydrogen, though renewable on human scales, is diluted on geological timescales, while helium, due to its inert nature, ac-1080 cumulates over similar periods. Their models demonstrate the rapid formation of nat-1081 ural hydrogen deposits, with instances of H2-dominant gas accumulations  $\sim 500$  years 1082 old in Mali, evolving to CH4-dominant mixtures  $\sim 40$  ka old in Turkey, and N<sub>2</sub>-rich vari-1083 ants within timescales of millions of years, exemplified by the Amadeus Basin in Aus-1084 tralia (Boreham et al., 2021). 1085

These analyses indicate that H<sub>2</sub>-rich gas subsurface accumulations are dependent 1086 on recent or ongoing hydrogen generation, whilst fossil accumulations are characterised 1087 by lower  $H_2$  abundances and greater organic and crustal component abundancies. The 1088 relationship between fluid migration and helium is paradoxical, given that for high He 1089 concentrations to be preserved alongside high H2 concentrations (e.g. Hydrogen (2023), 1090 migration must be rapid enough to prevent dilution whilst the isolation of fluids over bil-1091 lions of years could allow He to accumulate whilst  $H_2$  is lost (e.g. Karolytė et al. (2022)). 1092 The relationship between surface gas seep compositions and migration pathways is more 1093 nuanced, given that microbial methanogenesis (Table 5) is depth-dependent (e.g., Truche 1094 et al. (2024)). Hydrogen consumption within the shallowest levels of the subsurface (< 1095 1 m) may imprint a diurnal variation onto otherwise long-lived high-concentration (> 1096 50%) H<sub>2</sub> signals (e.g. Myagkiy, Moretti, and Brunet (2020)). For high H<sub>2</sub> concentrations 1097 to be preserved to the surface, we hypothesise that fluid migration pathways operate on 1098 both short timescales and lengthscales such that the opportunity for environmental hy-1099 drogen uptake is severely limited. However, we propose that seeps characterised by low 1100 hydrogen concentrations (< 20%, e.g. Lévy et al. (2023)) typify migration pathways over 1101 greater distances and timescales, such as basin-scale transport along faults and fractures, 1102 whereby substantial amounts of hydrogen are lost due to microbial consumption and other 1103 processes. 1104

### 1105 Conclusion

Hydrogen within the subsurface remains elusive. While entrenched into Earth dur-1106 ing planetary formation, the exchange of hydrogen between materials is prevalent dur-1107 ing subsurface processes at all depths. Significant advances in understanding the distri-1108 bution and generation of natural hydrogen have been made in recent literature, however 1109 large gaps remain in our understanding of large-scale hydrogen migration. The timescale 1110 of hydrogen migration throughout Earth varies from billions of years to days, and is de-1111 pendent on a wide range of lithological and environmental factors. Grain size, temper-1112 ature and fluid salinity exert important controls on hydrogen diffusivity in crystalline 1113 and sedimentary rocks. Diffusive and advective migration of hydrogen vary by several 1114 orders of magnitude, however operate on timescales of < 0.5 cm to m per > 1000 m 1115 per year, respectively. Fluid migration along faults and fractures is controlled by rock 1116 properties, subsurface stress regimes and groundwater properties. The phenomena of gas-1117 induced fault opening and hydrogen pulses associated with seismic activity require fur-1118 ther research. Microbial reactions moderate subsurface hydrogen flow by altering mass 1119 balance on differing timescales related to depth and environmental factors. Understand-1120 ing the transition between diffusive and advective flow of hydrogen and multiphase flu-1121 ids within different rock types in the subsurface remains a key challenge. 1122

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#### 1130 Data availability statement

The authors confirm that the data supporting the findings of this study are available within the article. Code used to calculate fluid properties (Figure 6) using the method of Lodhia and Clark (2022) is published at https://github.com/lodhia92/hydrogen .mobility (Lodhia, 2023).

#### 1135 **Declarations**

<sup>1136</sup> The authors declare no conflicts of interest.

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