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# **Towards a widely applicable earthquake detection algorithm for fibreoptic and hybrid fibreoptic-seismometer networks**

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## **SUMMARY**

Distributed Acoustic Sensing (DAS) is a promising technology for providing dense (metre-scale) sampling of the seismic wavefield. However, harnessing this potential for earthquake detection with accurate phase picking and associated localisation remains challenging. Single-channel algorithms are limited by individual channel noise, while machine learning and semblance methods are typically limited to specific geological settings, have no physically-constrained phase association and/or require specific fibre geometries. Here, we present a method that seeks to detect seismicity for any geological setting, applicable for any fibre geometry, and combining both fibreoptic and conventional seismometer data to maximise the information used for detection and source localisation. This method adapts a proven back-migration detection method to also include DAS observations, propagating energy from many receivers back in time to search for localised peaks in energy, corresponding to seismic sources. The strengths of this method are capitalising on coherency over many channels to enhance detection sensitivity even in high-noise environments compared to single-channel algorithms, applicability to arbitrary fibre geometries, as well as built-in, physics-informed phase association and source

21 localisation. We explore the performance of the method using three geologically and geometri-  
22 cally diverse settings: a glacier, a volcanic eruption and a geothermal borehole. Our results ev-  
23 idence the effect of spatial-sampling extent and non-optimal fibreoptic geometries, accounting  
24 for P and S wave sensitivity, coupling effects, and how the sensitivity of native fibreoptic strain  
25 measurements to shallow subsurface heterogeneities can affect detection. Finally, we attempt  
26 to also present a method-ambivalent overview of key challenges facing fibreoptic earthquake  
27 detection and possible avenues of future work to address them.

28 **Key words:** seismology, distributed acoustic sensing, earthquake detection, network seismol-  
29 ogy, computational seismology

## 30 1 INTRODUCTION

31 Earthquakes are essential monitoring various natural hazards, imaging subsurface structure and  
32 interrogating various Earth system processes. In order to harness the potential of earthquakes for  
33 either monitoring or insight into fundamental processes, one first has to detect and locate any  
34 seismicity. Typically, earthquake detection has been performed using conventional seismometers  
35 sensitive to the seismic velocity wavefield, but recent new optical instrumentation now allows one  
36 to use fibreoptic cables to measure the seismic strain wavefield with far denser spatial sampling  
37 over a wide bandwidth [Hartog et al., 2018, Lindsey and Martin, 2021, Lindsey et al., 2020a, Paitz  
38 et al., 2021]. This technology is often referred to as Distributed Acoustic Sensing (DAS). How-  
39 ever, currently these fibreoptic technologies only provide single-component measurements, with  
40 deployed fibre geometries often not favourable for earthquake detection and/or location. There-  
41 fore, there is a need for earthquake detection algorithms that can: firstly, be applied for arbitrary  
42 fibre geometries; secondly, maximise the spatial sampling extent of the seismic wavefield by com-  
43 bining fibreoptic and conventional seismometer data; and thirdly, capitalise on the spatio-temporal  
44 coherency of the earthquake wavefield.

45 Earthquake detection methods can be broadly separated into two categories: (1) receiver-by-  
46 receiver detection, with earthquake arrivals triggered on each seismogram in isolation and then  
47 combined afterwards using some form of phase association; or (2) multi-receiver detection, where

48 earthquake arrivals at multiple receivers are combined together in a physics-constrained frame-  
 49 work. Common receiver-by-receiver algorithms are applied in the time-domain using short-term-  
 50 average to long-term-average (STA/LTA) methods, for example [Allen, 1978, Withers et al., 1998],  
 51 or in the frequency-domain by looking for energy peaks within a certain frequency band [Helm-  
 52 stetter et al., 2015, O’Neel et al., 2007]. Machine-learning techniques, such as convolutional neural  
 53 networks, have also been applied to receiver-by-receiver seismic phase detection [Mousavi et al.,  
 54 2020, Zhu and Beroza, 2019, Zhu et al., 2023, Hernandez et al., 2022]. Recently, novel hybrid  
 55 algorithms combining multiple STA/LTA functions with machine-learning have also been devel-  
 56 oped [Latto et al., 2024]. Perpetual limitations of any receiver-by-receiver method are not using  
 57 spatio-temporal coherency information and the challenge of associating phase arrivals from each  
 58 receiver with one another, especially for multiple wave types [Ross et al., 2019a]. Multi-receiver  
 59 detection methods overcome these limitations by identifying coherent signals arriving at multiple  
 60 receivers. Methods that do not explicitly require knowledge of the medium’s velocity structure,  
 61 designed specifically for DAS measurements include: semblance-based coherence [Porrás et al.,  
 62 2024]; and machine-learning based image recognition methods [Stork et al., 2020, Huot et al.,  
 63 2022]. More general approaches include array-processing techniques, such as beamforming [Hud-  
 64 son et al., 2023] or covariance matrix analysis [Seydoux et al., 2016]. However, although these  
 65 methods do not explicitly require subsurface velocity structure information, they either implicitly  
 66 assume, learn or are sensitive to the local velocity structure. A final multi-receiver method of note  
 67 is back-migration, which explicitly requires an estimate of velocity structure in the region of inter-  
 68 est, using this information to effectively perform physics-informed stacking of energy arriving at  
 69 multiple receivers [Drew et al., 2013, Hudson et al., 2019, Smith et al., 2020, Winder et al., 2021,  
 70 Hudson et al., 2021, Guidarelli et al., 2020, Wagner et al., 2017].

71 When discussing earthquake detection, we deem it helpful to consider the following key in-  
 72 gredients for optimal earthquake detection algorithm performance:

- 73 (i) Maximise spatial coverage and sampling density
- 74 (ii) Exploit signal coherency
- 75 (iii) Maximise sensitivity to multiple seismic phases



76 (iv) Quantify event origin-time and phase arrival-time uncertainty

77 (v) Optimise computational efficiency

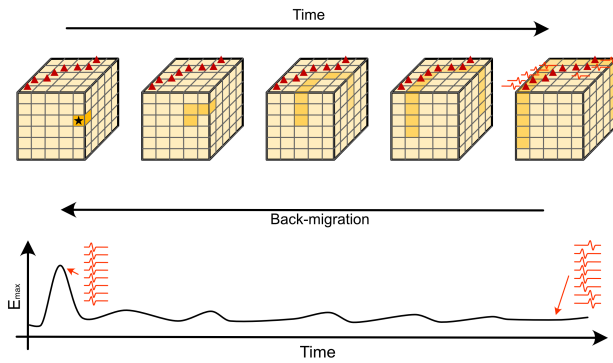
78 (vi) (*Bonus: Towards universal applicability, while considering the trade-off with computa-*  
79 *tional efficiency*)

80 In this work we quantify how important these ingredients are for earthquake detection and  
81 evidence the reasons why, relating them to the modifications required to apply an existing back-  
82 migration earthquake detection method to fibreoptic and hybrid fibreoptic-seismometer datasets.  
83 This builds on previous work, where back-migration was applied to a DAS dataset without any  
84 DAS-specific adaptations [Hudson et al., 2021]. While we focus here on local to regional micro-  
85 seismicity applications, these ingredients should generally be transferable to global earthquake  
86 detection [Selby, 2011] and noise localisation applications [Igel et al., 2023]. We favour a back-  
87 migration method because it includes all the key ingredients specified above in the recipe. In par-  
88 ticular, the adapted back-migration method presented here allows one to use arbitrary fibreoptic  
89 deployment geometries, maximise spatial coverage and seismic phase sensitivity by also including  
90 conventional seismometer data, and requires no further modification or retraining when applied to  
91 new datasets. It also provides earthquake location estimates without additional cost. Using fibre-  
92 optic (DAS) datasets for earthquake detection requires various specific considerations, including:  
93 in-axis fibre directional sensitivity, the use of native strain/strain-rate measurements, coupling of  
94 the fibre to the medium; and weighting the relative contribution of 1000s of DAS channels with  
95 far fewer conventional seismometer measurements. As we detail the method and its performance,  
96 we also identify and discuss remaining challenges of using fibreoptic datasets for earthquake de-  
97 tection.

## 98 **2 THE EARTHQUAKE DETECTION RECIPE**

### 99 **2.1 Back-migration at a glance**

100 The back-migration earthquake detection method converts continuous seismograms at each re-  
101 ceiver into onset functions that represent the energy from a particular seismic phase arriving at



**Figure 1.** Schematic of back-migration of energy in space through time. Top shows the 3D volume at different points in time, with darker shading corresponding to higher back-migrated energy amplitude at particular grid cells. Red triangles denote receivers and black star denotes location of peak back-migrated energy, corresponding to a hypothetical earthquake. Bottom shows energy corresponding to a maximum amplitude grid cell at each point in time. As energy is back-migrated through the grid in time, the energy coalesces towards a singularity in space.

102 each receiver through time [Drew et al., 2013]. These characteristic onset functions from all re-  
 103 ceivers are then back-migrated in space through time, effectively stacking the data with physically-  
 104 meaningful time-shifts. Potential events are detected by identifying coalescence peaks in the back-  
 105 migrated energy through time (see Figure 1). The key strength of this method lies in events only  
 106 being triggered by coherent source singularities rather than incoherent noise.

## 107 2.2 The QuakeMigrate algorithm

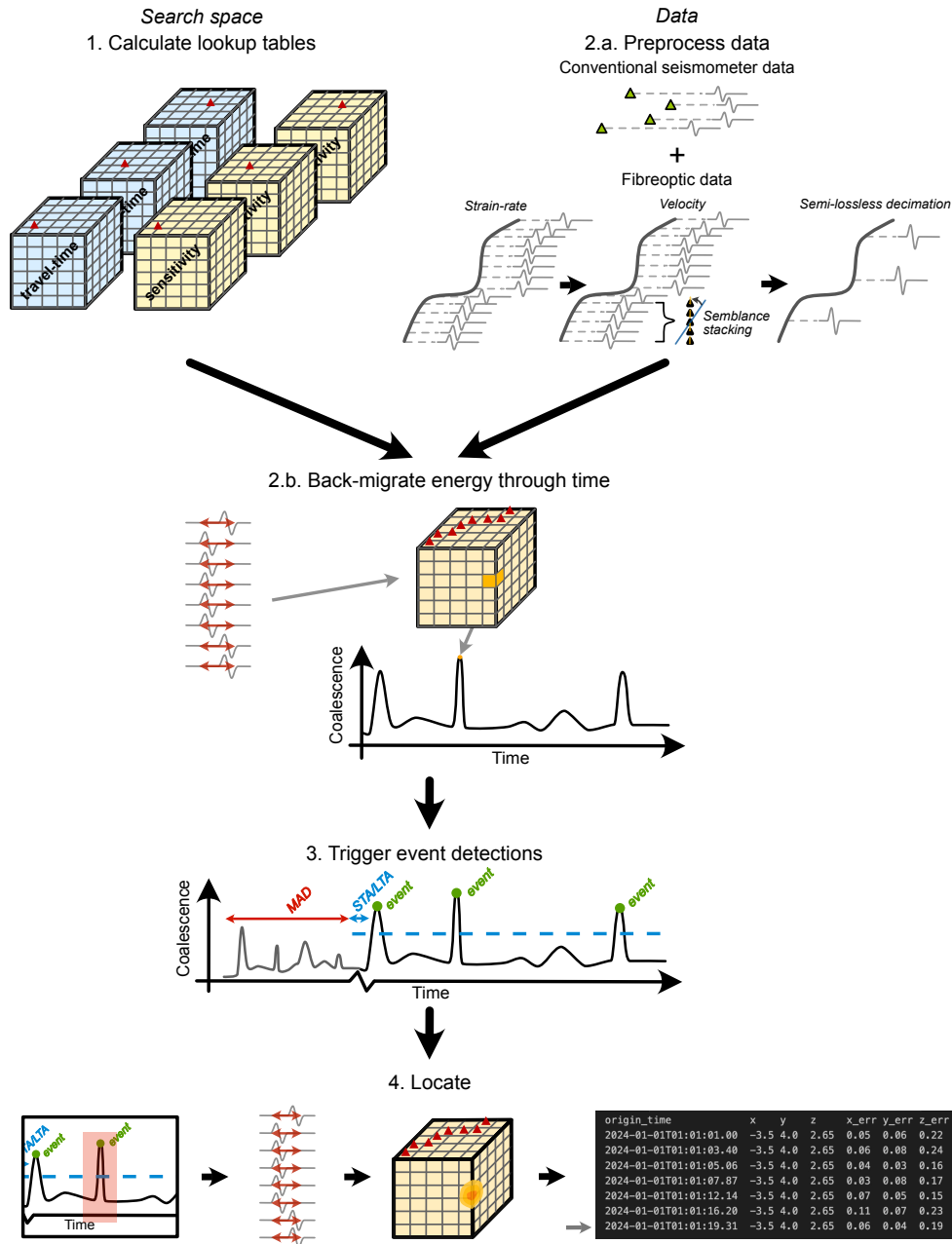
108 The specific back-migration algorithm used in this work is a modified version of the open source  
 109 software QuakeMigrate [Hudson et al., 2019, Smith et al., 2020, Winder et al., 2021]. The specific  
 110 steps of the modified QuakeMigrate algorithm are summarised in Figure 2. First, 3D travel-time  
 111 and fibre-sensitivity lookup tables are generated for each receiver and each seismic phase (e.g. P,  
 112 S), corresponding to the time-shifts required to back-migrate the characteristic onset function to a  
 113 particular point in space. This computationally expensive step is only performed once for a given  
 114 network and velocity model. One should note that this requires a velocity model, which theoret-  
 115 ically limits the universal applicability of the method. However, typically one can make an ap-  
 116 proximate yet sufficient guess at an initial model. Second, continuous seismograms are read in for  
 117 every receiver. Characteristic onset functions representing the energy arriving at each receiver are

118 calculated, for example by using an STA/LTA ratio. Typically for conventional three-component  
119 seismic data, vertical components are used for P wave arrivals and horizontal components for S  
120 wave arrivals, although this obviously differs for DAS data (see Section 2.3). These characteristic  
121 onset functions are shifted and stacked in time for each point in the 3D search space, for all seismic  
122 phases (e.g. P and S). For each point in time, the value of the grid cell with the maximum coales-  
123 cence of energy is recorded, producing a maximum coalescence time-series. The third step is to  
124 trigger possible event detections based on finding peaks in the coalescence time-series. Typically,  
125 we find the best trigger threshold is dynamic, using a multiple of the Median Absolute Deviation  
126 (MAD) value from a moving window of several hours in duration. Finally, one refines the event  
127 location by repeating the back-migration for each triggered candidate event individually. Event lo-  
128 cation can be refined by using data at a higher sampling-rate, a more spatially-dense lookup table,  
129 or different frequency filters, for example.

130 The outputs are an event catalogue, including arrival time picks for each seismic phase at  
131 each receiver, the earthquake origin time, and an estimate of earthquake location. Uncertainties  
132 are quantified for all parameters. Arrival time uncertainties are defined as the standard deviation  
133 of a Gaussian fit to the characteristic onset function for a given receiver and phase. Similarly, the  
134 earthquake origin time uncertainty is approximated as the standard deviation in time of the peak  
135 coalescence function. The earthquake location uncertainty is estimated from the standard deviation  
136 of a Gaussian fit to the marginalised coalescence in space at the earthquake origin time, which is  
137 assumed to be a proxy for the probability density function in space. If desired, one can also output  
138 additional information such as plots of the coalescence in space and time or plots of arrival time  
139 pick labelled waveform data. These are useful for initial refinement of the detection parameters for  
140 a particular dataset, especially regarding STA/LTA values, bandpass filters and lookup table grid  
141 resolution.

### 142 **2.3 Fibreoptic-specific modifications**

143 Although back-migration earthquake detection methods already exist, a number of modifications  
144 are required to include fibreoptic datasets and optimise the algorithm. Modifications include com-



**Figure 2.** Overview of the modified back-migration detection method. 1. Travel-time and DAS sensitivity lookup tables calculated for each channel. 2.a. Data is preprocessed, including conversion and stacking of DAS data. 2.b. Onset functions are calculated and back-migrated through time. 3. Potential events are triggered from peaks in the time-shifted, stacked onset functions (corresponding to coalescence (or a measure of energy)). 4. Back-migration is performed again, just for candidate events. Uncertainties are estimated in this step and a final earthquake catalogue is generated.

binning datasets with different units (velocity, strain/strain-rate), lossless spatial decimation of DAS channels, considering fibre-medium coupling effects, accounting for in-axis fibre sensitivity, and both incoherent and coherent noise reduction via spatio-temporal filtering. All modifications are implemented into the QuakeMigrate algorithm. Below, we describe how these modifications are implemented, organised in relation to the aforementioned key earthquake detection ingredients.

### 2.3.1 *Maximising spatial sampling extent and density*

Including fibreoptic data for earthquake detection is an obvious way to enhance spatial sampling density. However, as we show in the results, enhancing spatial sampling density alone does not necessarily equate to enhanced earthquake detection performance. Where receivers are placed geographically may be of similar importance as number of receivers deployed [Strutz and Curtis, 2024, Toledo et al., 2020]. Fibreoptic deployments fall into two categories: fibreoptic cables deployed specifically for a seismological application versus interrogation of existing dark-fibre telecommunication networks. While the first category allows one to design an optimal network geometry, one has no influence of the geometry for the second category. Fibreoptic geometries can severely limit back-migration-based detection methods (e.g. it is impossible to uniquely back-migrate energy from a linear fibre geometry). Overcoming such issues is only possible by including other data, for example conventional seismometers, which will enhance spatial coverage in almost any scenario.

However, combining data from different instrument types is non-trivial. Fibreoptic interrogators measure strain-rate (or strain) whereas seismometers typically measure velocity. Strain-rate,  $\dot{\epsilon}$ , is the spatial derivative of velocity,  $v$ . One might assume that we are ambivalent to receiver units, since we back-migrate an approximation of the normalised energy arriving at each receiver through time. However, spatially or temporally differentiating or integrating a periodic time-series leads to a systematic change in frequency-amplitude content. For example, let us assume that an earthquake has a simple sinusoidal signal,

$$v = \sin(\mathbf{k} \cdot \mathbf{x} - \omega \mathbf{t}),$$

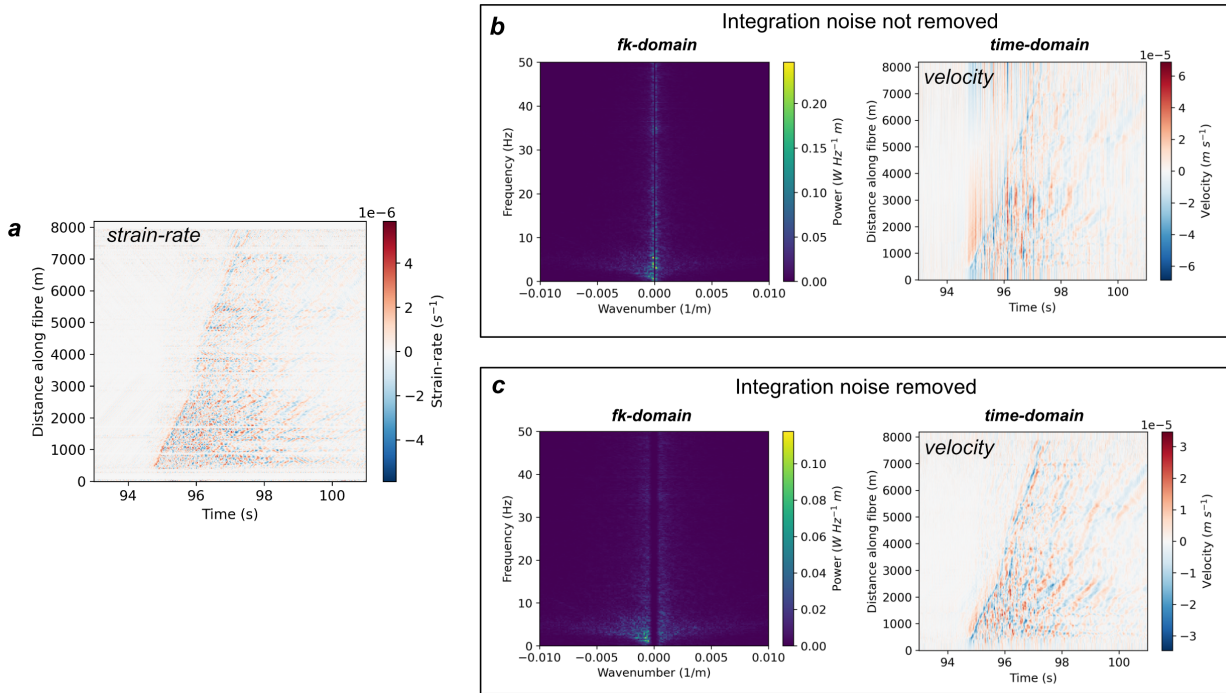
171 where  $\mathbf{k}$  is the wave number,  $\mathbf{x}$  is the direction vector,  $\omega$  is angular frequency and  $t$  is time. Then  
 172 conversion to  $\dot{\epsilon}$  gives,

$$173 \quad \dot{\epsilon} = \frac{\partial v}{\partial x} = k \cdot \sin(\mathbf{k} \cdot \mathbf{x} - \omega t).$$

174 Since  $k = 1/\lambda$ , where  $\lambda$  is the wavelength, and the wavelength is proportional to the frequency  
 175 (since  $c = \lambda f$ ), the amplitude of  $\dot{\epsilon}$  is dependent on frequency. Such a relative change in frequency  
 176 content in seismograms from different receivers could cause issues when pre-processing data in-  
 177 puts (e.g. bandpass filtering), and affect onset function amplitudes and hence overall coalescence  
 178 amplitudes.

179 We therefore opt for converting all data into the same units before calculating and back-  
 180 migrating onset functions if possible. An example of conversion from strain-rate to velocity is  
 181 shown in Figure 3. This is only possible over approximately linear sections of fibre. If the fibre  
 182 does not have significant sections that fulfil this criteria then we leave the data in native strain-rate.  
 183 We choose to convert DAS strain-rate to velocity rather than converting seismometer velocity to  
 184 strain-rate for several reasons. Firstly, converting seismometer velocities into strain-rate is highly  
 185 challenging as one would have to combine all seismometer data, reconstruct the wavefield, and  
 186 then take the spatial derivative of that wavefield at seismometer locations [Muir and Zhan, 2022].  
 187 This is impossible unless one has many conventional receivers. Secondly, integrating DAS strain-  
 188 rate to velocity is not only more practical, but the resultant integration noise has infinite apparent  
 189 velocity that can be removed using an  $fk$ -filter, increasing SNR. Thirdly, the integration also acts  
 190 as a spatial low-pass filter, removing some incoherent noise. Fourthly, strain-rate is inherently  
 191 more sensitive to local velocity structure than velocity [Capdeville and Sladen, 2024], with con-  
 192 version to velocity removing this sensitivity. If an earthquake is far from the receiver and one does  
 193 not know the local velocity structure adequately, then again one gains an improvement in back-  
 194 migration performance. As we have already hinted, we convert from strain-rate to velocity via  
 195 direct integration followed by an  $fk$ -filter to remove near-infinite apparent velocities.

196 In summary, maximising spatial coverage requires using data from as many receivers as possi-  
 197 ble. Back-migration algorithm performance is then improved by converting all data into the same



**Figure 3.** Example of strain-rate to velocity conversion. a. An earthquake distance along fibre vs. time plot for an earthquake from the Reykjanes Peninsula, Iceland in units of strain-rate. b. The same earthquake as in (a), but converted to velocity without removing infinite apparent velocity integration noise, both in  $fk$ -space and distance-time space. c. Same as (b) but with integration noise removed.

198 units (velocity), which both homogenises frequency content as well as reducing noise and hetero-  
 199 geneous local velocity structure effects.

### 200 2.3.2 Exploit signal coherency

201 The back-migration method inherently exploits coherency, which introduces both benefits and  
 202 challenges when including fibreoptic datasets.

203 One benefit of exploiting coherency is associated with coupling of fibre to the medium. Ide-  
 204 ally, one would quantify coupling and remove poorly-coupled channels from any analysis. How-  
 205 ever, quantifying coupling in fibreoptic deployments remains challenging. Instead, here we sim-  
 206 ply capitalise on the assumption that poor coupling results in incoherent noise that cannot be  
 207 back-migrated. Poorly coupled channels will reduce the overall maximum theoretical coalescence  
 208 value, since the channels may not contribute but crucially poor coupling will not contribute to false  
 209 triggers.

210 The primary challenge associated with including fibreoptic data is the large number of chan-  
211 nels compared to conventional receivers. The challenge lies in how to balance information from  
212 typically lower SNR, single-component fibre channels and higher SNR conventional seismometer  
213 data. This is not an issue if using only fibreoptic data. However, in the most extreme case one might  
214 have many ( $n$ ) fibreoptic receivers and a single seismometer receiver. Theoretically one could as-  
215 sign receiver types different weights, but such weighting is limited because as one preferentially  
216 weights data from one receiver, it also preferentially weights noise. Hypothetically if one equally  
217 weighted a single seismometer with the same weight as  $n$  fibre channels, seismometer receiver  
218 noise would contribute  $n$  times that of noise on a single DAS channel. This would potentially  
219 negate any coherency gain.

220 We instead opt for weighting the contribution of fibreoptic receivers relative to conventional  
221 instrumentation via semi-lossless decimation. This semi-lossless decimation refers to performing  
222 semblance-based stacking on every  $n$  fibreoptic receivers, similar to [Porras et al. \[2024\]](#). Specifi-  
223 cally, semblance-stacking comprises time-shifting every channel relative to every other channel in  
224 order to maximise the stacked amplitude. Time-shifts are limited by a maximum permissible ap-  
225 parent velocity. Theoretically this preserves both amplitude and directional information. However,  
226 we refer to the method as semi-lossless rather than lossless as we discard the directional infor-  
227 mation but preserve coherency post decimation. This semblance-based decimation improves deci-  
228 mated fibreoptic receiver signal quality, which at some point would theoretically have sufficiently  
229 high SNR to provide equal constraint/information in the detection algorithm to any conventional  
230 receivers. Weighting therefore remains subjective, controlled by the number of DAS channels  
231 stacked. In practice, one is limited by the fibre geometry, since semblance-based stacking does not  
232 work on arbitrary orientations of fibre (e.g. it would fail if applied to two channels orthogonal to  
233 one another). Therefore, in practice we currently employ the philosophy of decimating as much as  
234 possible while preserving both semblance-stack performance (e.g. only over linear segment scales  
235 or the gauge-length) and spatial coverage.



### 2.3.3 *Maximise sensitivity to multiple seismic phases*

A strength of back-migration detection methods is that one can use multiple seismic phases to constrain earthquake location better, enhancing coalescence and therefore improving detection performance. QuakeMigrate traditionally does this for P- and S- waves by back-migrating vertical and horizontal receiver component onset functions through P-wave and S-wave velocity models, respectively. Fibreoptic channels only measure signals in the fibre axis, making it non-trivial to isolate different seismic phases. Instead, we modify the algorithm to allow one to use fibreoptic channels for both P and/or S phases, as well as surface waves. This can be specified wholesale or individually for every channel, depending on whether a channel is horizontally deployed on the surface or vertically in a borehole, for example. Currently, the surface wave implementation involves specifying a single group velocity for the medium and assuming the energy migrates within the near-surface. However, more rigorous inclusion of surface waves could be achieved by modifying the method to back-migrate energy through different phase-velocity models for different frequency bands, approximately simulating surface wave dispersion and Rayleigh vs. Love waves.

Above, we hint that fibre channel orientation and surface vs. subsurface deployment play a role in the sensitivity of a particular channel to a particular seismic phase. While this topic could be the subject of numerous studies, one can typically assume that subsurface channels deployed vertically are sensitive to both P and S waves, whereas horizontal surface channels are dominantly sensitive to S-waves because of steep near-surface velocity gradients resulting in near-vertical ray incident angles [Hudson et al., 2021]. An exception to this is if the medium has an approximately homogeneous velocity structure, for example if deployed on ice [Walter et al., 2020]. However, for the surface DAS examples in this work, one has a homogeneous velocity structure and so has similar sensitivity for both P and S waves, and although the other is dominantly sensitive to S-waves, we still observe some P-wave energy in that case too.

In reality, the sensitivity of fibreoptic cables to different seismic waves is not binary. A more sophisticated approach that we implement here is to calculate fibreoptic channel sensitivity based on ray takeoff angle derived from the same velocity model used to calculate travel-time lookup tables. We calculate the takeoff angle for a ray propagating to each receiver from every grid cell

264 in the search volume, for both P and S waves. Fibreoptic sensitivity to strain-rate and velocity  
 265 differ, so for flexibility we implement both. For velocity as measured by fibreoptic channels, the  
 266 sensitivity of a fibreoptic channel to P, SV and SH phases is given by [Martin, 2018],

$$267 \zeta_P = \cos(\phi_1 - \theta)\cos(\phi_2),$$

$$268 \zeta_{SV} = \cos(\phi_1 - \theta)\sin(\phi_2),$$

$$269 \zeta_{SH} = \sin(\phi_1 - \theta),$$

270 where  $\theta$  is the angle of the fibre on a plane relative to a reference direction (e.g. north),  $\phi_1$  is the  
 271 in-plane angle of a plane wave propagation direction relative to the same reference direction, and  
 272  $\phi_2$  is the angle of the plane wave relative to the plane-perpendicular angle (e.g. angle from vertical  
 273 down for a horizontal fibre channel). Alternatively, for strain-rate the sensitivity is,

$$274 \zeta_P = \cos^2(\phi_1 - \theta)\cos^2(\phi_2),$$

$$275 \zeta_{SV} = \cos^2(\phi_1 - \theta)\sin(2\phi_2),$$

$$276 \zeta_{SH} = \sin(2(\phi_1 - \theta))\cos(\phi_2).$$

277 Note that these are for point strain and we drop any wave amplitude factors and frequency or phase  
 278 dependence since we do not know the amplitude, frequency content or phase of any prospective  
 279 arrivals prior to detecting them. SV and SH sensitivities then have to be somehow combined. Since  
 280 the proportion of SV to SH energy incident at a receiver prior to detection is unknown, maximum  
 281 sensitivity to any given S-wave polarisation is assumed. We therefore define the S wave sensitivity  
 282 as,

$$283 \zeta_S = \max \left( \sqrt{\frac{\zeta_{SV}^2 + \zeta_{SH}^2}{2}}, \zeta_{SV}, \zeta_{SH} \right).$$

284 These equations are now implemented in QuakeMigrate. We then define a sensitivity threshold be-  
 285 low which we deem that a particular fibre channel is insensitive to that location within the search  
 286 volume. Such a threshold could be selected, for example, by determining at what sensitivity value  
 287 the amplitude of a fibre channel would fall below the noise level of another channel with perfect  
 288 sensitivity. We then mask these regions for the associated seismic phase at that particular channel.  
 289 We do this for all fibreoptic channels. This provides greater constraint over where potential events

290 coalesce. This approach is dependent on knowledge of the approximate velocity structure, espe-  
291 cially the shallow velocity gradient. It should therefore be used with caution, with the user able to  
292 choose whether to implement it or not. While it is computationally intensive, it is only performed  
293 once when the travel-time lookup tables are generated, so subsequent runtime on continuous seis-  
294 mic data is unaffected.

295 While we only address body-wave sensitivity here, theoretically, one can also surface wave  
296 sensitivity in a similar way.

#### 297 *2.3.4 Quantify uncertainty*

298 No explicit uncertainty quantification modifications are made to the QuakeMigrate back-propagation  
299 method. However, several modifications will affect uncertainty estimates. Arrival time uncertainty  
300 estimates for fibreoptic receivers will depend on the number of channels stacked, with stacking im-  
301 proving the SNR, but potentially reducing the first break resolution. Applying semblance-stacking  
302 minimises this issue by time-shifting each trace so that first breaks should be aligned. Here, we typ-  
303 ically perform semblance stacking on upsampled data ( $\times 10$ ), preserving time precision. Including  
304 fibreoptic data also affects hypocentre uncertainty estimation. Dense fibreoptic channel spacing  
305 (typically order of metres) may be far smaller than the size of the grid cells in travel-time lookup  
306 tables, which means that these fibreoptic channels may not optimally contribute to hypocentre  
307 constraints and hence not minimise uncertainty. Furthermore, using fibreoptic channel sensitivity  
308 to mask out regions of the search space could change hypocentre uncertainty estimates, either re-  
309 ducing uncertainty if the velocity model structure used to calculate sensitivity is sufficiently close  
310 to the true structure, or artificially perturbing the uncertainty if not.

#### 311 *2.3.5 Maximise computational efficiency*

312 Back-migration detection methods are inherently computationally expensive compared to simpler,  
313 receiver-by-receiver detection methods. However, QuakeMigrate runs in approximately real-time  
314 for the experiments presented here (using 8 processors on an Apple M3 Pro CPU). These effi-  
315 ciencies are primarily driven by three contributions: computing lookup tables only once for entire

316 datasets; reading in continuous seismic data in blocks rather than entire files, minimising read-  
317 write operations; and implementing the core back-migration step in the pre-compiled C language.  
318 To minimise the additional computational expense of including fibreoptic data, we only compute  
319 sensitivity lookup tables once, and perform lookup table masking without the need to subsequently  
320 store sensitivity information independently in memory. Secondly, we support reading of a num-  
321 ber of native DAS data formats (hdf5, segy etc) directly, which are typically split into small, one  
322 minute duration files. One can then simply run QuakeMigrate over minute long time-windows,  
323 optimising costly read-write processes and memory usage.

### 324 **3 DATA**

325 Three datasets are used to investigate performance of the new method. These datasets approxi-  
326 mately represent end-members of current fibreoptic deployments: (1) a dense 2D fibreoptic grid  
327 deployed coincident with nodes on a glacier; (2) a dark-fibre located within the vicinity of a vol-  
328 canic eruption; and (3) a downhole fibre at a geothermal field. The specific detection algorithm  
329 settings used in each case are given in Table 1.

330 The glacier dataset comprises of 1.2 km fibre deployed at Gornergletscher in the Swiss Alps,  
331 in October 2023. The network has a  $\sim 100$  m aperture, with 29 single vertical component Sercel  
332 WiNG nodes deployed in the same area. The interrogator used is a Sintela Onyx, measuring strain,  
333 with a gauge-length of 6 m and a channel spacing of 1.6 m. All data were acquired with a sampling  
334 rate of 1000 Hz. The majority of microseismicity at the study site is thought to be caused by near-  
335 surface crevassing [[Walter et al., 2009](#), [Hudson et al., 2020](#)].

336 The volcanic eruption dataset comprises of an 8 km dark fibre, interrogated during the first  
337 Svartsengi volcanic eruption on the Reykjanes Peninsula, Iceland, in December 2023. We also  
338 include data from a broadband seismometer operated by the Icelandic Meteorological Office. The  
339 fibreoptic interrogator used is a Silixa iDAS, measuring strain-rate, with a gauge-length of 10  
340 m and a channel spacing of 16 m. All data are sampled at 100 Hz. Seismicity detected here is  
341 attributed to one intrusion episode on the 18<sup>th</sup> December 2023.

342 The downhole geothermal dataset is from the Utah Frontier Observatory for Research in

**Table 1.** QuakeMigrate detection settings for each dataset in this study.

Parameter	Glacier	Volcanic eruption	Geothermal borehole
Phases used	P, surface	P, S	P, S
Sampling rate	1000 Hz	100 Hz	1000 Hz
Frequency filter, P	10-250 Hz	1.2-20 Hz	2-300 Hz
Frequency filter, S	n/a	1.2-20 Hz	2-300 Hz
Frequency filter, surface	5-150 Hz	n/a	n/a
Grid resolution, x	8 m	150 m	400 m
Grid resolution, y	8 m	150 m	400 m
Grid resolution, z	10 m	300 m	100 m
STA/LTA P	0.01/0.2	0.2/1	0.01/0.5
STA/LTA S	0.02/0.2	0.2/1	0.01/0.5
Coalescence detection threshold	1.15	1.125	1.7
Marginal window	0.25 s	2 s	0.25 s
DAS specific settings			
Spatial downsamp. factor	1	5	10
Channel spacing	1.6 m	16 m	1 m
Gauge-length	6.38 m	10 m	10 m
Semblance-stacking	no	yes	yes

343 Geothermal Energy (FORGE) 2019 experiment, consisting of 1.2 km of fibre cemented into a  
344 vertical monitoring borehole [Lellouch et al., 2020]. A network of seismometers was deployed at  
345 the surface. The fibre is interrogated using a Silixa iDAS interrogator with a gauge length of 10 m,  
346 a channel spacing of 1 m and a sampling rate of 100 Hz. Here, we focus on one particularly active  
347 hour of seismicity during a well stimulation at 17:00 to 18:00 on 27 April 2019.

348 While these example datasets are not comprehensive, the majority of fibreoptic deployments  
349 for studying seismicity are likely similar to at least one of these examples, perhaps with the excep-  
350 tion of subsea and urban deployments.

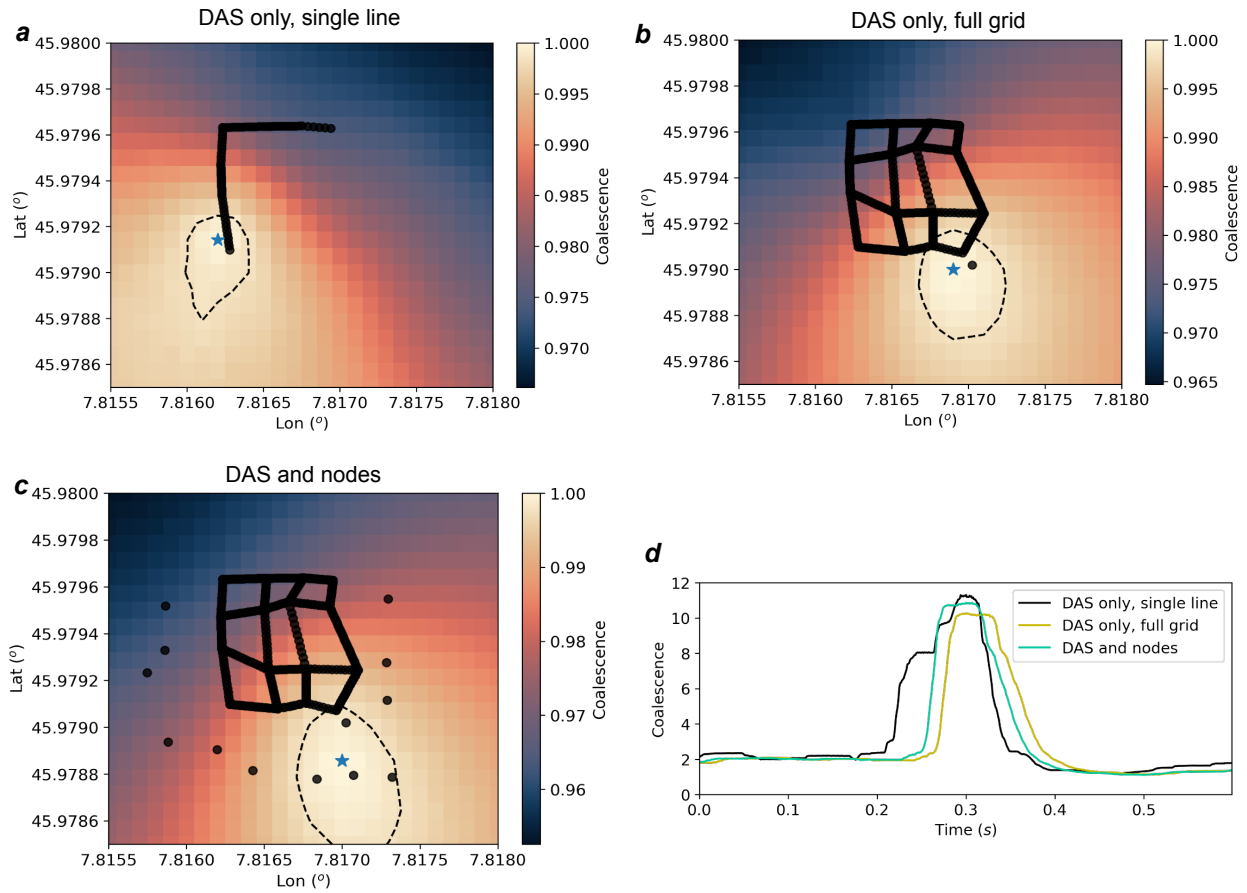
## 4 RESULTS

### 4.1 Algorithm performance

Before showing earthquake catalogues for each example dataset, we first investigate various specific aspects of the detection algorithm performance in more detail. In particular, we focus on the importance of spatial coverage when including both fibreoptics and seismometers, stacking fibreoptic channels, fibreoptic coupling, and fibreoptic sensitivity.

#### 4.1.1 Spatial coverage and fibreoptics-only vs. combined fibreoptics and seismometers

Adequate source localisation is essential for any back-migration detection method, such as the one described in this study. Theoretically, the better the source localisation, the higher the peak energy observed in the space-time search space and hence the more likely an event is to be detected above the ambient noise level. Figure 4 exemplifies how important spatial extent of sampling of the seismic wavefield is for source localisation and hence detection. Here, P and surface waves are used to detect an example icequake. Figure 4a shows the example earthquake detected by QuakeMigrate using only two vertices of the Gornergletscher fibre. Such a geometry is typical of many fibre deployments, for example that of the Iceland volcanic eruption dataset in this study. The dashed line in Figure 4a shows the 95% contour of the peak coalescence, with the blue star indicating the location of the peak corresponding to the inferred icequake location. The result of Figure 4a can be compared to that of Figure 4b, where the whole 2D fibre is used to detect and locate the same icequake. The icequake hypocentre moves considerably, with no overlap in the 95% contours for the two solutions. When one also includes the 29 nodes in the detection (Figure 4c) then although the icequake hypocentre moves again, this time it is within the 95% contour. The additional constraint provided by the nodes further constrains event location and hence the ability to detect the event above the noise level. The results of Figure 4 show how important maximising spatial sampling coverage is for source localisation and hence detection. For glacier deployments, it is conceivable to deploy dense 2D geometries, but in other situations it is likely that maximising spatial coverage will often require the combination of conventional seismometers and/or nodes in addition to fibre.



**Figure 4.** Example of how important spatial coverage is for back-migration earthquake detection. Data shown is from an icequake at Gornergletscher, Switzerland. a. Map view of the 3D coalescence space at the earthquake depth and origin time, for a simple fibre geometry consisting of only two vertices of the entire fibre deployment, typical of many current or dark fibre deployments. Blue star indicates the peak coalescence corresponding to the inferred icequake hypocentre. Black dashed line indicates the corresponding 95% contour. b. Same as (a) but using all fibre channels to detect the icequake. c. Same as (a,b) but also including the co-deployed vertical single-component nodes. d. Comparison of maximum single-pixel coalescence values through time for each setup in (a) to (c).

#### 378 4.1.2 Stacking and coupling

379 Two immediate challenges of processing fibreoptic datasets are processing large data volumes  
 380 resulting from inherently dense spatial sampling and fibre-medium coupling issues. Figure 5 sum-  
 381 marises the effects of stacking fibre channels to reduce data volumes and accounting for coupling  
 382 effects by removing poorly coupled channels, for the glacier dataset.

383 The effect of stacking multiple fibre channels is shown in Figure 5b. The motivation for stack-

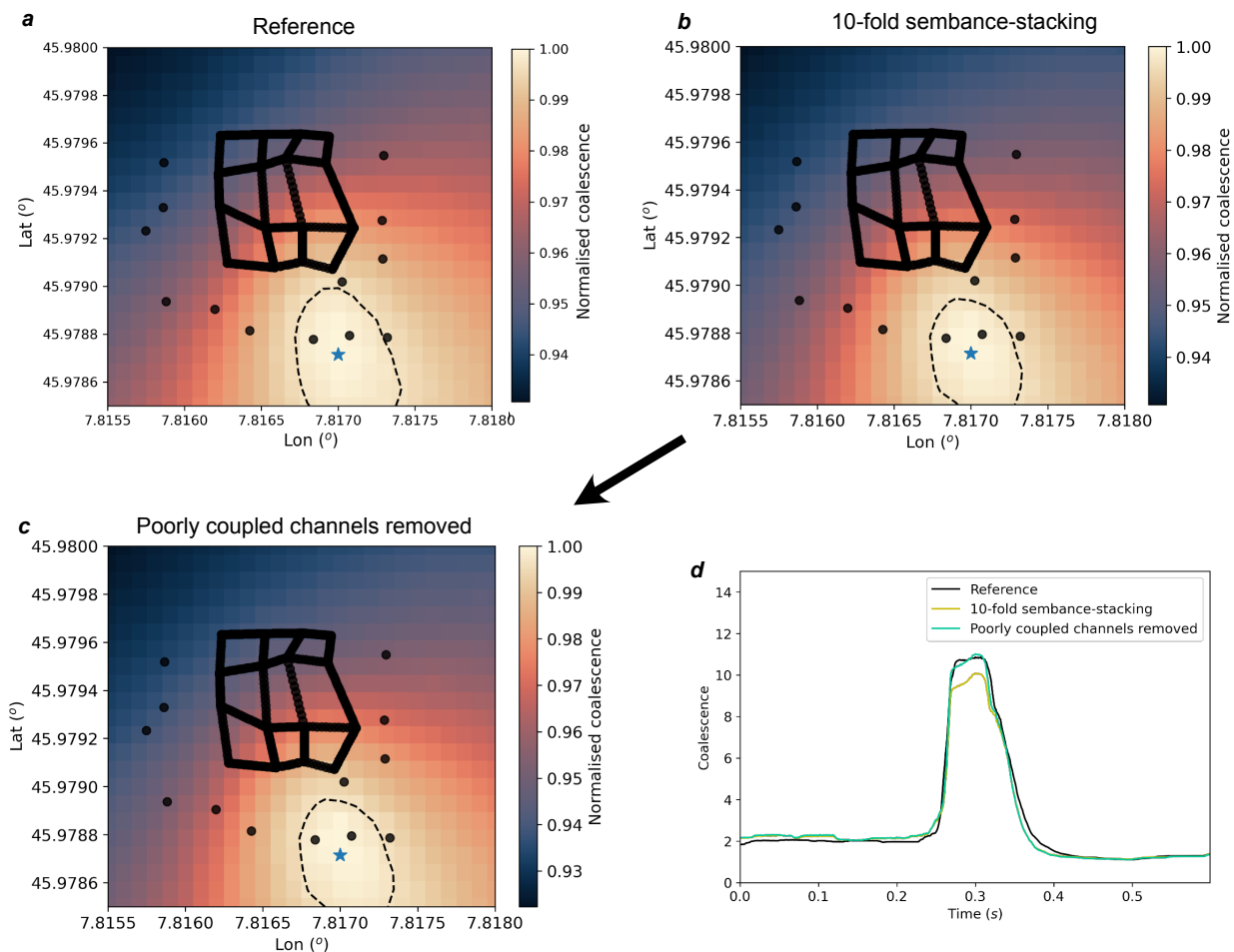
ing is two-fold: firstly to reduce data volumes and hence increase computational efficiency; and secondly to move towards balancing contributions from fibre channels and seismometers/nodes. The result in Figure 5b applies semblance-stacking to every 10 channels (16 m), aiming to preserve spatio-temporal coherency information while spatially downsampling the data. The results show that downsampling the fibre via stacking does not have a significant effect on the hypocentral location but does result in a spatially more constrained peak in the coalescence. This result is interesting since it shows that the semblance-stacking not only preserves the coherency information but acts to reduce noise effects between individual channels, enhancing the overall coherency of the solution in space and hence improving detection performance. Semblance-stacking therefore not only gives a computational performance gain (both in terms of efficiency and memory usage), but also enhances detection algorithm performance, at least in this instance.

The effect of removing poorly coupled channels on detection performance is shown in Figure 5c. Here, we remove channels that traverse crevasses and are hence poorly coupled to the ice and comparatively well coupled to the atmosphere. Not only are these channels therefore approximately insensitive to subsurface seismic energy, but actually have higher noise amplitudes due to atmospheric effects such as wind, for example. The results are remarkably similar to those of Figure 5b, with no distinguishable difference in the location of the coalescence peak, but the 95% contour becoming better constrained spatially. One might expect removing poorly coupled channels to have a greater effect. However, we attribute the relatively insignificant change in performance to be a result of a key strength of back-migration, in that poorly-coupled channels represent incoherent noise that theoretically should not contribute significantly to event detection. At least in this example, accounting for coupling appears to not be of first-order importance.

#### 4.1.3 Including fibreoptic sensitivity

The influence of accounting for the effects of fibreoptic measurement sensitivity are shown in Figure 6 for both a glacier icequake (Figure 6a-d) and a volcano-tectonic earthquake (Figure 6e-i). For the glacier icequake, accounting for sensitivity affects the coalescence, including both the peak and extent of the 95% contour, moving the peak and tightening the spatial constraint (see Figure 6b





**Figure 5.** Effect of stacking and coupling on detection algorithm performance. a. Reference coalescence map, using all available fibre channels for detection of an example icequake. b. 10-fold (10 channel) semblance-stacking result for the same icequake as in (a). c. Detection result for same icequake as in (a) except with poorly coupled channels removed. d. Comparison of maximum single-pixel coalescence values through time for each setup in (a) to (c).

411 vs. Figure 6(a). The effect is more extreme than the effects of stacking or removing poorly coupled  
 412 channels (Figure 5). Accounting for sensitivity for the volcanic earthquake example has a smaller  
 413 effect. This is despite the geometry being far more linear than in the icequake example. While the  
 414 95% contour encloses a smaller spatial extent, the earthquake hypocentre moves only  $\sim 100$  m,  
 415 relative to the 8 km fibre. We attribute this behaviour to the fact that the somewhat linear fibre ge-  
 416 ometry here has a highly non-uniform sensitivity to both P and S waves, and so the network is only  
 417 sensitive to seismic energy from these regions already, so the additional sensitivity constraint we  
 418 impose has little effect. This is contrary to the icequake example, where the overall network has an

419 approximately homogeneous sensitivity to incoming energy from any source location. We argue  
420 that it is therefore debatable whether one should impose any sensitivity constraint in practice, since  
421 it is already incorporated into the analysis and if the velocity structure is poorly constrained then  
422 the sensitivity maps will also be poorly constrained. However, in any case, plotting up the sensi-  
423 tivity maps for various seismic waves (Figure 6c,g,h) is insightful and should always be compared  
424 to the distribution of earthquake hypocentres output from a detection and location algorithm.

## 425 4.2 Generating earthquake catalogues

### 426 4.2.1 Glacier

427 Figure 7 shows the icequake catalogue from the fibreoptic deployment at Gornergletscher in the  
428 Swiss Alps. The majority of these icequakes are likely generated by near-surface crevassing [Wal-  
429 ter et al., 2009, Hudson et al., 2020]. Firstly, the 2D fibre geometry results in no apparent bias in the  
430 spatial distribution of seismicity. In this example, we use both P- and surface- waves to contribute  
431 to event detection. The surface fibre deployment is sensitive to P-waves due to the approximately  
432 homogeneous velocity structure of ice, with no shallow slow velocity firm layer present. Such slow  
433 velocity layers at glaciers have proven problematic for P-wave detection previously [Hudson et al.,  
434 2021]. S-wave energy generated by crevassing is expected to be minimal [Hudson et al., 2020], so  
435 is not used here. Since the ice column is assumed to be of  $< 100$  m thickness and the dominant  
436 icequake generation mechanism is expected to be near-surface crevassing, we also use surface-  
437 waves for detection. The site has several sources of noise, including wind and subsurface fluids  
438 that flow through some of the fractures. Coupling directly to glacier ice in such conditions can  
439 be challenging, but in this experiment the fibre is generally well-coupled to the ice since the fibre  
440 froze in within the first 12 hours of the deployment. The quality of coupling can be seen in Figure  
441 7b, with most channels showing clear P- and surface- wave arrivals for an icequake, but with a  
442 number of channels showing only noise where they are suspended above a crevasse (at  $\sim 470$  m,  
443  $\sim 750$  m and  $\sim 970$  m, for example). Here, all channels are included for detection. This is based  
444 on the finding that the algorithm is unaffected, at least to first-order, by poorly coupled channels  
445 with spatially uncorrelated noise (see Figure 5c).

446 173 icequakes are detected in six hours. While there are likely more icequakes in the dataset,  
447 we opt to lower the detection threshold only to a point where we are confident that we minimise  
448 false detections. Figure 7b-d shows results for one icequake. The detection algorithm generally  
449 picks both P- and surface- wave first breaks where one would manually identify them, after ac-  
450 counting for uncertainty, although uncertainties are large (of the order of the dominant surface-  
451 wave period). Seismic energy arriving at one of the nodes is also shown in Figure 7d, evidencing  
452 that the detection algorithm performs adequately on both conventional and fiberoptic data in com-  
453 bination, one of the key aims of this work.

#### 454 4.2.2 *Volcanic eruption*

455 Figure 8 shows an earthquake catalogue during one episode of the ongoing Sundhnúkaíggar erup-  
456 tion, Iceland. The dark fibre interrogated during this experiment has a somewhat linear geometry,  
457 typical of many dark fibre geometries, following a road from a geothermal power plant to the  
458 coast. Both the geothermal plant (at the fibre origin) and the coast (at the far end of the fibre) gen-  
459 erate coherent noise. Coherent noise sources can affect the performance of the detection algorithm  
460 if they contain energy within the bandwidth of interest. These noise sources detrimentally affect  
461 phase arrival identification from 0 m to 300 m and beyond 7900 m along the fibre. The presence  
462 of both natural and anthropogenic noise, along with the numerous earthquakes that occur over the  
463 time period make this dataset an ideal case study.

464 The fiberoptic data combined with a single seismometer detects 886 earthquakes within the  
465 region shown in Figure 8a on the 18<sup>th</sup> December 2023, compared to 826 earthquakes detected  
466 within the same region by the permanent regional monitoring network (operated by the Icelandic  
467 Met Office). However, although the energy from the earthquakes coalesces sufficiently to make de-  
468 tections, the locations typically remain poor. We expect the majority of seismicity to align with the  
469 opening rift [Sigmundsson et al., 2024], but instead find that the seismicity clusters near one end  
470 of the fibre. This is likely for several reasons. Firstly, many of the earthquakes detected are close  
471 to the noise level, affecting the accuracy of individual channel arrival time picks. Secondly, and  
472 likely more significantly, the geometry of the fibre provides poor location constraint (see Figure

473 [6c,d](#)). This is partly due to poor azimuthal coverage but also likely the result of low sensitivity to  
474 P-waves from certain regions of the seismically active rift. Poor azimuthal constraint likely results  
475 in poor locations, primarily because the regional velocity model used is likely faster than the true  
476 shallow velocity structure.

477 Even though locations are poorly constrained, phase arrival times for some events are promis-  
478 ing. For example, the earthquake shown in Figure [6b-d](#) shows P and S arrival time picks close to  
479 the first break, with realistic uncertainties, even though the P-wave amplitude is close to the noise  
480 level. This is particularly clear in Figure [6c](#). It is surprising that one can even observe P-waves on  
481 the fibre at all, given the crustal setting. However, we attribute this to be due to the fibre being de-  
482 ployed within metres of the bedrock, removing much of the effect of a steeply varying near surface  
483 velocity gradient that would otherwise refract P-waves towards vertical incidence. The earthquake  
484 arriving at a conventional seismometer also included for detection is shown in Figure [6d](#), again  
485 confirming the promising performance when processing hybrid fibreoptic-seismometer datasets.

486 While it is encouraging that the detection algorithm works well even when locations are poorly  
487 constrained, our findings from this dataset illustrate the challenges associated with using fibreoptic  
488 data from dark fibres for source localisation.

#### 489 4.2.3 Geothermal borehole

490 Figure [9](#) shows an earthquake catalogue for one hour during a stimulation test at the Utah FORGE  
491 experiment. The deployment is typical of many borehole DAS deployments, with only a vertical  
492 fibre cemented into a well, with seismometers deployed at the surface. For the hour of data we  
493 analyse, the closest surface seismometers to the injection well have too high noise to observe any  
494 subsurface seismicity. We therefore use only one seismometer, FORU, in combination with the  
495 downhole fibreoptic data.

496 We detect 135 earthquakes, compared to 125 earthquakes detected using a combination of stan-  
497 dard methods and matched filter processing with a string of 12 borehole geophones [\[Dzubay et al.,  
498 2022\]](#). Most earthquakes are present in both catalogues. One should note that comparing numbers  
499 of earthquake detections can be misleading, since one can lower detection thresholds and detect

500 many more earthquakes or vice versa. Obviously one can check candidate events individually, but  
501 this is not feasible for datasets with  $\sim 1000$ s earthquakes. Here, we try to avoid this by setting the  
502 detection threshold to a level where we minimise false detections while still detecting as many  
503 real events as possible. However, it is still worth emphasising that the detection algorithm of this  
504 study detects a comparable number of earthquakes to the most sensitive detection method possible,  
505 a matched filter method. Although there is no ground truth, we deem earthquake locations to be  
506 poor, with a portion of the earthquakes detected locating far from the stimulation well located at  
507 approximately  $38.5^\circ N$ ,  $112.9^\circ W$ . Similar to the volcanic eruption example, we attribute poor lo-  
508 cation constraint to be dominantly caused by network geometry. If more surface instruments were  
509 usable, locations would likely be constrained better.

510 Although locations are interpreted to be poor, this does not translate to poor earthquake de-  
511 tections, consistent with our findings from the volcanic eruption results. This is demonstrated for  
512 the example earthquake shown in Figure 9b. Both P- and S- wave phases are identified with small  
513 uncertainties compared to the other datasets presented. This is likely due to a combination of good  
514 coupling of the cemented fibre, as well as perhaps lower noise levels in the subsurface and a rela-  
515 tively homogeneous velocity structure in the vicinity of the monitoring well. S-wave arrivals may  
516 be incorrectly identified in some cases (see Figure 9c), but this is to be expected given the noise  
517 levels and/or P-wave coda.

## 518 **5 DISCUSSION**

### 519 **5.1 Practical considerations influencing earthquake detection with fibreoptic sensing**

520 The findings of this work emphasise a number of practical considerations affecting earthquake  
521 detection using fibreoptic sensing. While we focus on local microseismicity detection using a back-  
522 migration method, most of the points below also hold for earthquake detection using fibreoptic  
523 sensing more generally.

524 (i) The spatial sampling extent of the seismic wavefield plays the most important role in earth-  
525 quake localisation, yet back-migration based detection still performs successfully in practice, even

526 if the coherent source of energy is poorly constrained spatially. Specifically, the geometry of the  
527 fibreoptic cable plays an important role, evidenced by the glacier 2D grid deployment of fibre de-  
528 tecting and locating events better than highly linear fibre geometries such as the volcanic eruption  
529 dataset, for example. Extrapolating from this, we expect 3D fibre geometries to further improve  
530 performance. In practice, many fibreoptic cable geometries are limited by practical constraints,  
531 and so combining fibreoptic data with conventional receivers provides an optimal alternative.

532 (ii) Fibreoptic sensing inherently provides dense sampling of the seismic wavefield. Decima-  
533 tion of this data is important for increasing the computational efficiency and minimising memory  
534 usage in practice. We find that using semblance-stacking [Porras et al., 2024] to decimate the data  
535 preserves wavefield coherency information. Such stacking is applicable where fibre geometries  
536 have linear sections at least of the order of the spatial decimation/stacking length.

537 (iii) Conversion from native strain or strain-rate to velocity enhances performance, for fibre  
538 geometries with substantial linear sections or low curvature. For example, we find a phase pick  
539 accuracy gain for the volcanic eruption example but not the glacier example. This gain is primarily  
540 attributed to dampening sensitivity to subsurface heterogeneities [Capdeville and Sladen, 2024] but  
541 also the isolation of integration noise.

542 (iv) Compensating for fibreoptic sensitivity can provide additional spatial constraint of coales-  
543 cence peaks, but the benefit is limited. This finding is important because it implies that sensitivity  
544 may not be a major practical concern for earthquake detection for most deployments.

545 (v) There exists an optimal spatial resolution of the 3D search grid. We introduce this spe-  
546 cific back-migration concept here, since it can play an important role in our detection algorithm  
547 performance. Specifically, too high a resolution grid results distributes coherent energy between  
548 more grid cells, due to velocity structure uncertainty, resulting in lower peaks in coalescence. Con-  
549 versely, too low a resolution grid results in high coalescence values of a single grid, but risks not  
550 isolating coherent earthquake signals from coherent or even incoherent noise. Optimal grid cell  
551 resolutions for each example in this work are given in Table 1.

552 (vi) There also exists an optimal moving time-window with which to process the data. Again,  
553 this point is a specific back-migration concept, but important for users of the detection algorithm

554 presented here. Specifically, the time-window over which one processes phase arrivals must be  
555 sufficiently long that the uncertainty in phase arrivals can be adequately estimated, while being  
556 sufficiently short to eliminate phase mis-association.

## 557 **5.2 Comparison to alternative approaches**

558 The greatest strength of fibreoptic sensing is the orders-of-magnitude denser spatial sampling over  
559 conventional instrumentation. Coherency-based back-migration methods such as that described in  
560 this study by definition capitalise on coherency between these densely sampled channels in a way  
561 that receiver-by-receiver detection methods cannot. Receiver-by-receiver STA/LTA algorithms are  
562 susceptible to triggering off incoherent noise. Receiver-by-receiver machine-learning-based phase  
563 arrival detection methods are theoretically less sensitive to incoherent noise since they approxi-  
564 mately learn to identify the noise field, but this requires the availability of a training dataset of  
565 existing detected earthquakes. Furthermore, all receiver-by-receiver methods struggle with phase  
566 association. Although machine-learning derived phase associators exist [Ross et al., 2019b], we  
567 find that they do not always perform adequately for local seismicity (within 5 km of a network).  
568 The only benefit of receiver-by-receiver methods over back-migration is computational efficiency.  
569 However, the all the examples presented in this work run faster-than-real-time on a standard com-  
570 puter (8 processors on a Apple M3 Pro CPU). Surpassing the faster-than-real-time benchmark is  
571 key for any seismic monitoring application, with any additional gains only a bonus.

572 Some promising advances in detecting earthquakes using fibreoptics have been made by as-  
573 sessing coherent energy arriving at many fibre channels simultaneously. These include explicit  
574 methods such as assessing the curvature of arrivals on linear fibre using semblance methods [Porras  
575 et al., 2024], and implicit methods that use machine-learning image recognition algorithms [Stork  
576 et al., 2020, Huot et al., 2022] to identify similar features. Although these methods are promising,  
577 typically detecting events close to the noise-level, there are a number of limitations. Firstly, explicit  
578 semblance-based methods require specific, linear fibre geometries, or at least substantial lengths  
579 of linear fibre. It is unclear how sensitive machine-learning image recognition algorithms are to  
580 fibre geometry, but current implementations would have to be retrained for every new deployment

581 or new fibre geometry. A further limitation is that it is challenging to conceive of ways to include  
582 both fibreoptic and conventional receivers in such detection frameworks. [Huot et al., 2024] lay  
583 foundations for combining fibreoptic and conventional receivers in a machine learning algorithm  
584 by running two binary logistic regression models in parallel and identifying events detected by  
585 both. They find that this significantly reduces false triggers compared to only using the fibreoptic  
586 data, emphasising the importance of combining data from all available receivers where possible.  
587 However, unlike the back-migration method we present here, such methods do not yet fully cou-  
588 ple the information provided by multiple receiver types, and therefore do not yet fully realise the  
589 associated performance gain.

### 590 5.3 Perpetual challenges

591 A number of challenges remain for using fibreoptics for earthquake detection. Particular challenges  
592 that are difficult or even impossible to overcome are as follows. Firstly, fibreoptic cables are sen-  
593 sitive only to in-axis strain. This limits sensitivity to multiple seismic phases, especially in the  
594 presence steeply varying shallow subsurface velocity gradients [Hudson et al., 2021]. Helically-  
595 wound fibre has the potential to overcome this issue [Baird, 2020], but still represents a pseudo-1D  
596 measurement, is costly, and is never deployed in telecommunications networks. Secondly, fibre-  
597 medium coupling remains generally both poorly constrained and for some experiments uncontrol-  
598 lable [Paitz et al., 2021]. Recent work allows one to quantify expected coupling response [Celli  
599 et al., 2024] and if a fibre is buried, frozen-in or cemented in-situ then coupling is approximately  
600 perfect. However, for fibres deployed on the surface or deployed in conduits, for example dark  
601 fibres, coupling is challenging to quantify. Thirdly, our results emphasise the importance of fibre  
602 geometries and spatial extent, yet at least for dark fibre deployments one has little or no control  
603 over the deployed geometry. Overcoming this issue may be possible by interrogating many fibres  
604 in dense, urban environments, but will remain a challenge in rural or subsea regions. Fourthly,  
605 data volumes remain challenging. Experiments can generate 100s GB to TBs of data per day.  
606 Downsampling data will therefore become essential with increasing deployment duration and/or  
607 increased deployment spatial scales. The semblance stacking component described in this study



608 may allow on-the-fly decimation while preserving some directivity information. A final challenge  
609 of note is near-field coherent noise, for example roads or train lines [Dou et al., 2017, Lindsey  
610 et al., 2020b]. The back-migration method we present would be sensitive to coherent vehicle noise  
611 that would likely prevent simultaneous earthquake detection. A key benefit of fiberoptic sensing  
612 is measuring the seismic wavefield in urban environments, so going forward this challenge would  
613 have to be addressed for urban earthquake detection.

#### 614 **5.4 Future directions**

615 The aforementioned challenges inspire directions for future work. Recording and processing large  
616 data volumes (TBs) is non-trivial and limits real-time earthquake detection using back-migration  
617 style algorithms. A possible avenue for reducing data volumes could be applying compressive-  
618 sensing based approaches [Muir and Zhan, 2021] to effectively reduce the number of fibre chan-  
619 nels used while retaining sufficient information to detect seismicity. In a similar vain, using non-  
620 uniform or cascading coalescence search grids, inspired by those used for earthquake location and  
621 tomography [Lomax and Curtis, 2001, Thrastarson et al., 2024], could optimise memory usage and  
622 compute expense while refining peak coalescence values and hence detection. Another avenue of  
623 future work is reducing or removing noise. Recent advances applying machine learning to min-  
624 imise instrument noise [Lapins et al., 2023] could readily be applied to our detection workflow,  
625 improving detection performance. A remaining endeavour is to perform masking of coherent noise  
626 sources. For example, certain regions of the coalescence search grid may correspond to roads that  
627 act as temporally-varying noise sources. One could envisage adapting our detection method to re-  
628 move parts of the wavefield corresponding to a coalescence of energy from these locations, in a  
629 similar approach to that taken to mitigate fiberoptic sensitivity in this work. Simultaneously arriv-  
630 ing lower SNR earthquakes may then be detectable even in the presence of higher local coherent  
631 noise sources.

## 632 **6 CONCLUSIONS**

633 Here, we describe adaptations made to a back-migration earthquake detection method required  
634 to process fibreoptic (DAS) receiver measurements. We show the inherent strengths of back-  
635 migration detection methods for harnessing the dense spatial sampling of the seismic wavefield  
636 that fibreoptic sensing provides. Such methods can be deployed for processing fibreoptic cables  
637 of arbitrary geometries. They also enable the combination of fibreoptic and conventional receivers  
638 to maximise the spatial extent and information used for detection. We also attempt to provide an  
639 overview of the current status of earthquake detection using fibreoptic sensing more generally, es-  
640 pecially regarding identifying challenges provided by fibreoptic strain measurements and possible  
641 ways to overcome them. Finally, although the detection method presented here appears robust for  
642 earthquake detection in diverse geological settings, we briefly identify key remaining challenges  
643 and future directions that these challenges inspire.

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**DATA AVAILABILITY**

Earthquake catalogues for each dataset are provided, along with a full working example of how to run the modified QuakeMigrate package, via an archived online repository [Hudson, 2024a]. The exact DAS modified version of the QuakeMigrate software is also archived online [Hudson, 2024b], which will be merged with the main QuakeMigrate repository in due course.. Data for the example earthquakes from the Gornergletscher glacier and the Iceland volcanic dataset are also provided in the same repository. Also provided in this repository are jupyter-notebook examples of running the modified version of QuakeMigrate. FORGE DAS data are available from US DOE Geothermal Data Repository (<https://doi.org/10.15121/1603679>) and seismometer data from the University of Utah (<https://doi.org/10.7914/SN/UU>), accessed from Incorporated Research Institutions for Seismology (<https://ds.iris.edu/mda/UU/>).

**References**

- R. V. Allen. Automatic earthquake recognition and timing from single traces. *Bulletin of the Seismological Society of America*, 68(5):1521–1532, 1978.
- A. F. Baird. Modelling the Response of Helically Wound DAS Cables to Microseismic Arrivals. In *First EAGE Workshop on Fibre Optic Sensing*, pages 1–5. European Association of Geoscientists Engineers, 2020. doi: 10.3997/2214-4609.202030019. URL <https://www.earthdoc.org/content/papers/10.3997/2214-4609.202030019>.
- Y. Capdeville and A. Sladen. DAS sensitivity to heterogeneity scales much smaller than the minimum wavelength. *Seismica*, 3(1), jan 2024. ISSN 2816-9387. doi: 10.26443/seismica.v3i1.1007. URL <https://seismica.library.mcgill.ca/article/view/1007>.
- N. L. Celli, C. J. Bean, and G. S. O’Brien. Full-waveform simulation of DAS records, response and cable-ground coupling. *Geophysical Journal International*, 236(1):659–674, jan 2024. ISSN 1365246X. doi: 10.1093/gji/ggad449.
- S. Dou, N. Lindsey, A. M. Wagner, T. M. Daley, B. Freifeld, M. Robertson, J. Peterson, C. Ulrich, E. R. Martin, and J. B. Ajo-Franklin. Distributed Acoustic Sensing for Seismic Monitoring of the Near Surface: A Traffic-Noise Interferometry Case Study. *Scientific Reports*, 7(1):1–12,

- 682 2017. ISSN 20452322. doi: 10.1038/s41598-017-11986-4. URL <http://dx.doi.org/10.1038/s41598-017-11986-4>.
- 683
- 684 J. Drew, R. S. White, F. Tilmann, and J. Tarasewicz. Coalescence microseismic mapping. *Geophysical Journal International*, 195(3):1773–1785, 2013. ISSN 0956540X. doi: 10.1093/gji/ggt331.
- 685
- 686
- 687 A. Dzubay, M. Mesimeri, K. M. Whidden, D. Wells, and K. Pankow. Developing a comprehensive seismic catalog using a matched-filter detector during a 2019 stimulation at Utah FORGE. *47th Workshop on Geothermal Reservoir Engineering*, pages 1–9, 2022.
- 688
- 689
- 690 M. Guidarelli, P. Klin, and E. Priolo. Migration-based near real-time detection and location of microearthquakes with parallel computing. *Geophysical Journal International*, 221(3):1941–1958, 2020. ISSN 1365246X. doi: 10.1093/gji/ggaa111.
- 691
- 692
- 693 A. H. Hartog, L. B. Liokumovich, N. A. Ushakov, O. I. Kotov, T. Dean, T. Cuny, A. Constantinou, and F. V. English. The use of multi-frequency acquisition to significantly improve the quality of fibre-optic-distributed vibration sensing. *Geophysical Prospecting*, 66:192–202, 2018. ISSN 13652478. doi: 10.1111/1365-2478.12612.
- 694
- 695
- 696
- 697 A. Helmstetter, L. Moreau, B. Nicolas, P. Comon, and M. Gay. Intermediate-depth icequakes and harmonic tremor in an Alpine glacier (Glacier d’Argentière, France): Evidence for hydraulic fracturing. *Journal of Geophysical Research: Earth Surface*, 120(3):402–416, 2015. ISSN 21699011. doi: 10.1002/2014JF003289.
- 698
- 699
- 700
- 701 P. D. Hernandez, J. A. Ramirez, and M. A. Soto. Deep-Learning-Based Earthquake Detection for Fiber-Optic Distributed Acoustic Sensing. *Journal of Lightwave Technology*, 40(8):2639–2650, 2022. ISSN 15582213. doi: 10.1109/JLT.2021.3138724.
- 702
- 703
- 704 T. S. Hudson. Supplementary data for the article: ”Towards an widely applicable earthquake detection algorithm for fibreoptic and hybrid fibreoptic-seismometer networks”. *Zenodo*, 2024a. doi: 10.5281/zenodo.13355664.
- 705
- 706
- 707 T. S. Hudson. QuakeMigrate with DAS integration. *Zenodo*, 2024b. doi: 10.5281/zenodo.13355775.
- 708
- 709 T. S. Hudson, J. Smith, A. M. Brisbourne, and R. S. White. Automated detection of basal ice-

- 710 quakes and discrimination from surface crevassing. *Annals of Glaciology*, 60(79):167–181, sep  
711 2019. ISSN 0260-3055. doi: 10.1017/aog.2019.18. URL [https://www.cambridge.org/  
712 core/product/identifier/S0260305519000181/type/journal\\_article](https://www.cambridge.org/core/product/identifier/S0260305519000181/type/journal_article).
- 713 T. S. Hudson, A. M. Brisbourne, R. S. White, J. M. Kendall, R. Arthern, and A. M. Smith.  
714 Breaking the Ice: Identifying Hydraulically Forced Crevassing. *Geophysical Research Let-  
715 ters*, 47(21), nov 2020. ISSN 0094-8276. doi: 10.1029/2020GL090597. URL [https:  
716 //onlinelibrary.wiley.com/doi/10.1029/2020GL090597](https://onlinelibrary.wiley.com/doi/10.1029/2020GL090597).
- 717 T. S. Hudson, A. F. Baird, J. M. Kendall, S. K. Kufner, A. M. Brisbourne, A. M. Smith,  
718 A. Butcher, A. Chalari, and A. Clarke. Distributed Acoustic Sensing (DAS) for Natural Mi-  
719 croseismicity Studies: A Case Study From Antarctica. *Journal of Geophysical Research: Solid  
720 Earth*, 126(7):1–19, 2021. ISSN 2169-9313. doi: 10.1029/2020jb021493.
- 721 T. S. Hudson, A. M. Brisbourne, S.-k. Kufner, J.-m. Kendall, and M. Smith. Array processing in  
722 cryoseismology. *The Cryosphere*, 2023. doi: 10.5194/egusphere-2023-657.
- 723 F. Huot, A. Lellouch, P. Given, B. Luo, R. G. Clapp, T. Nemeth, K. T. Nihei, and B. L. Biondi.  
724 Detection and Characterization of Microseismic Events from Fiber-Optic das Data Using Deep  
725 Learning. *Seismological Research Letters*, 93(5):2543–2553, 2022. ISSN 19382057. doi:  
726 10.1785/0220220037.
- 727 F. Huot, R. G. Clapp, and B. L. Biondi. Detecting local earthquakes via fiber-optic cables in  
728 telecommunication conduits under Stanford University campus using deep learning. *Computers  
729 and Geosciences*, 190(May), 2024. ISSN 00983004. doi: 10.1016/j.cageo.2024.105625.
- 730 J. K. H. Igel, D. C. Bowden, and A. Fichtner. SANS: Publicly Available Daily Multi-Scale  
731 Seismic Ambient Noise Source Maps. *Journal of Geophysical Research: Solid Earth*, 128(1):1–  
732 22, jan 2023. ISSN 2169-9313. doi: 10.1029/2022JB025114. URL [https://onlinelibrary.  
733 wiley.com/doi/10.1029/2022JB025114](https://onlinelibrary.wiley.com/doi/10.1029/2022JB025114).
- 734 S. Lapins, A. Butcher, J.-M. Kendall, T. S. Hudson, A. L. Stork, M. J. Werner, J. Gunning,  
735 and A. M. Brisbourne. DAS-N2N: machine learning distributed acoustic sensing (DAS) sig-  
736 nal denoising without clean data. *Geophysical Journal International*, 236(2):1026–1041, dec  
737 2023. ISSN 0956-540X. doi: 10.1093/gji/ggad460. URL <https://academic.oup.com/gji/>

[article/236/2/1026/7453669](https://doi.org/10.1026/7453669).

R. B. Latto, R. J. Turner, A. M. Reading, and J. P. Winberry. Towards the systematic reconnaissance of seismic signals from glaciers and ice sheets - Part 1: Event detection for cryoseismology. *Cryosphere*, 18(4):2061–2079, 2024. ISSN 19940424. doi: 10.5194/tc-18-2061-2024.

A. Lellouch, N. J. Lindsey, W. L. Ellsworth, and B. L. Biondi. Comparison between distributed acoustic sensing and geophones: Downhole microseismic monitoring of the FORGE geothermal experiment. *Seismological Research Letters*, 91(6):3256–3268, 2020. ISSN 19382057. doi: 10.1785/0220200149.

N. J. Lindsey and E. R. Martin. Fiber-Optic Seismology. *Annual Review of Earth and Planetary Sciences*, pages 309–336, 2021.

N. J. Lindsey, H. Rademacher, and J. B. Ajo-Franklin. On the Broadband Instrument Response of Fiber-Optic DAS Arrays. *Journal of Geophysical Research: Solid Earth*, 125(2):1–16, 2020a. ISSN 21699356. doi: 10.1029/2019JB018145.

N. J. Lindsey, S. Yuan, A. Lellouch, L. Gualtieri, T. Lecocq, and B. Biondi. City-Scale Dark Fiber DAS Measurements of Infrastructure Use During the COVID-19 Pandemic. *Geophysical Research Letters*, 47(16):1–8, 2020b. ISSN 19448007. doi: 10.1029/2020GL089931.

A. Lomax and A. Curtis. Fast, probabilistic earthquake location in 3D models using oct-tree importance sampling. In *Geophys. Res. Abstr*, volume 3, 2001.

E. R. Martin. Passive imaging and characterization of the subsurface with DAS. *PhD thesis*, 2018.

S. M. Mousavi, W. L. Ellsworth, W. Zhu, L. Y. Chuang, and G. C. Beroza. Earthquake transformer—an attentive deep-learning model for simultaneous earthquake detection and phase picking. *Nature Communications*, 11(1):3952, dec 2020. ISSN 2041-1723. doi: 10.1038/s41467-020-17591-w. URL <http://dx.doi.org/10.1038/s41467-020-17591-w><http://www.nature.com/articles/s41467-020-17591-w>.

J. B. Muir and Z. Zhan. Seismic wavefield reconstruction using a pre-conditioned wavelet-curvelet compressive sensing approach. *Geophysical Journal International*, 227(1):303–315, 2021. ISSN 1365246X. doi: 10.1093/gji/ggab222.

- 766 J. B. Muir and Z. Zhan. Wavefield-based evaluation of DAS instrument response and array  
767 design. *Geophysical Journal International*, 229(1):21–34, 2022. ISSN 1365246X. doi: 10.  
768 1093/gji/ggab439.
- 769 S. O’Neel, H. P. Marshall, D. E. McNamara, and W. T. Pfeffer. Seismic detection and analysis of  
770 icequakes at Columbia Glacier, Alaska. *Journal of Geophysical Research: Earth Surface*, 112  
771 (3):1–14, 2007. ISSN 21699011. doi: 10.1029/2006JF000595.
- 772 P. Paitz, P. Edme, D. Gräff, F. Walter, J. Doetsch, A. Chalari, C. Schmelzbach, and A. Ficht-  
773 ner. Empirical investigations of the instrument response for distributed acoustic sensing (Das)  
774 across 17 octaves. *Bulletin of the Seismological Society of America*, 111(1):1–10, 2021. ISSN  
775 19433573. doi: 10.1785/0120200185.
- 776 J. Porras, D. Pecci, G. M. Bocchini, S. Gaviano, M. De Solda, K. Tuinstra, F. Lanza,  
777 A. Tognarelli, E. Stucchi, and F. Grigoli. A semblance-based microseismic event detector for  
778 DAS data. *Geophysical Journal International*, 236(3):1716–1727, jan 2024. ISSN 0956-540X.  
779 doi: 10.1093/gji/ggae016.
- 780 C. Porter. ArcticDEM, Version 4.1. *Harvard Dataverse*, 2023. doi: 10.7910/DVN/3VDC4W,.
- 781 Z. E. Ross, Y. Yue, M. A. Meier, E. Hauksson, and T. H. Heaton. PhaseLink: A Deep Learning  
782 Approach to Seismic Phase Association. *Journal of Geophysical Research: Solid Earth*, 124  
783 (1):856–869, 2019a. ISSN 21699356. doi: 10.1029/2018JB016674.
- 784 Z. E. Ross, Y. Yue, M. A. Meier, E. Hauksson, and T. H. Heaton. PhaseLink: A Deep Learning  
785 Approach to Seismic Phase Association. *Journal of Geophysical Research: Solid Earth*, 124  
786 (1):856–869, 2019b. ISSN 21699356. doi: 10.1029/2018JB016674.
- 787 N. D. Selby. Improved Teleseismic Signal Detection at Small-Aperture Arrays. *Bulletin of*  
788 *the Seismological Society of America*, 101(4):1563–1575, aug 2011. ISSN 0037-1106. doi:  
789 10.1785/0120100253. URL [https://pubs.geoscienceworld.org/bssa/article/101/4/  
790 1563-1575/349509](https://pubs.geoscienceworld.org/bssa/article/101/4/1563-1575/349509).
- 791 L. Seydoux, N. M. Shapiro, J. De Rosny, F. Brenguier, and M. Landès. Detecting seismic activity  
792 with a covariance matrix analysis of data recorded on seismic arrays. *Geophysical Journal*  
793 *International*, 204(3):1430–1442, 2016. ISSN 1365246X. doi: 10.1093/gji/ggv531.

- 794 F. Sigmundsson, M. Parks, H. Geirsson, A. Hooper, V. Drouin, K. S. Vogfjörð, B. G. Ófeigsson,  
795 S. H. M. Greiner, Y. Yang, C. Lanzi, G. P. De Pascale, K. Jónsdóttir, S. Hreinsdóttir, V. Tolpekin,  
796 H. M. Fririksdóttir, P. Einarsson, and S. Barsotti. Fracturing and tectonic stress drive ultra-  
797 rapid magma flow into dikes. *Science*, 383(6688):1228–1235, mar 2024. ISSN 0036-8075.  
798 doi: 10.1126/science.adn2838. URL [https://www.science.org/doi/10.1126/science.  
799 adn2838](https://www.science.org/doi/10.1126/science.adn2838).
- 800 J. D. Smith, R. S. White, J.-P. Avouac, and S. Bourne. Probabilistic earthquake locations of in-  
801 duced seismicity in the Groningen region, the Netherlands. *Geophysical Journal International*,  
802 222(1):507–516, 2020. ISSN 0956-540X. doi: 10.1093/gji/ggaa179.
- 803 A. L. Stork, A. F. Baird, S. A. Horne, G. Naldrett, S. Lapins, J.-M. Kendall, J. Wookey, J. P.  
804 Verdon, A. Clarke, and A. Williams. Application of Machine Learning To Microseismic Event  
805 Detection in Distributed Acoustic Sensing (Das) Data. *Geophysics*, pages 1–53, 2020. ISSN  
806 0016-8033. doi: 10.1190/geo2019-0774.1.
- 807 D. Strutz and A. Curtis. Variational Bayesian experimental design for geophysical applications:  
808 seismic source location, amplitude versus offset inversion, and estimating CO<sub>2</sub> saturations in  
809 a subsurface reservoir. *Geophysical Journal International*, 236(3):1309–1331, 2024. ISSN  
810 1365246X. doi: 10.1093/gji/ggad492.
- 811 S. Thrastarson, D. P. van Herwaarden, S. Noe, C. J. Schiller, and A. Fichtner. REVEAL: A  
812 Global Full-Waveform Inversion Model. *Bulletin of the Seismological Society of America*, 114  
813 (3):1392–1406, 2024. ISSN 19433573. doi: 10.1785/0120230273.
- 814 T. Toledo, P. Jousset, H. Maurer, and C. Krawczyk. Optimized experimental network design for  
815 earthquake location problems: Applications to geothermal and volcanic field seismic networks.  
816 *Journal of Volcanology and Geothermal Research*, 391:106433, 2020. ISSN 03770273. doi:  
817 [10.1016/j.jvolgeores.2018.08.011](https://doi.org/10.1016/j.jvolgeores.2018.08.011). URL [https://doi.org/10.1016/j.jvolgeores.2018.  
818 08.011](https://doi.org/10.1016/j.jvolgeores.2018.08.011).
- 819 F. Wagner, A. Tryggvason, R. Roberts, B. Lund, and Gudmundsson. Automatic seismic event  
820 detection using migration and stacking: A performance and parameter study in Hengill, south-  
821 west Iceland. *Geophysical Journal International*, 209(3):1866–1877, 2017. ISSN 1365246X.



822 doi: 10.1093/gji/ggx127.

823 F. Walter, J. F. Clinton, N. Deichmann, D. S. Dreger, S. E. Minson, and M. Funk. Moment tensor  
824 inversions of icequakes on Gornergletscher, Switzerland. *Bulletin of the Seismological Society  
825 of America*, 99(2):852–870, 2009. doi: 10.1785/0120080110.

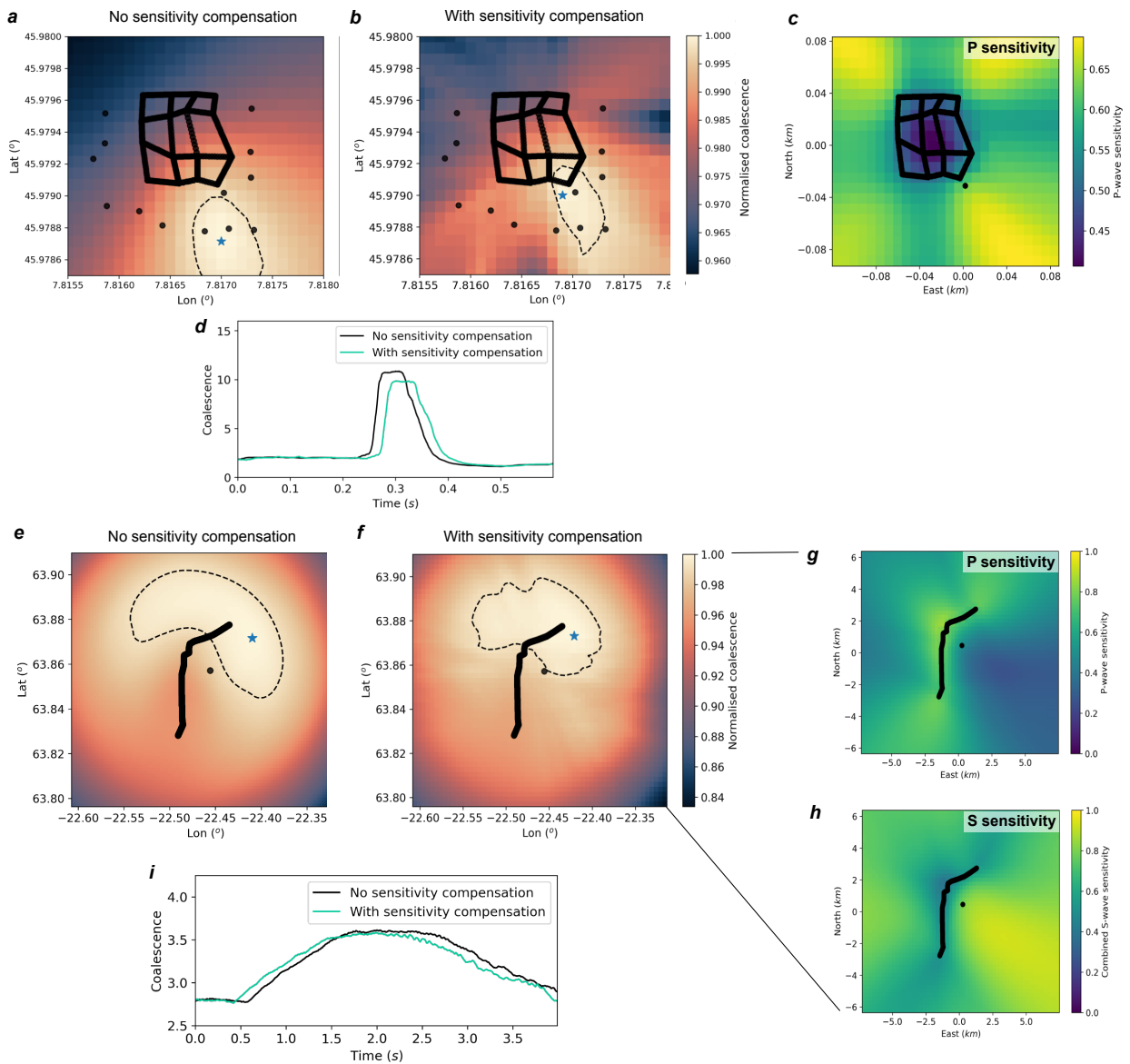
826 F. Walter, D. Gräff, F. Lindner, P. Paitz, M. Köpfl, M. Chmiel, and A. Fichtner. Dis-  
827 tributed Acoustic Sensing of Microseismic Sources and Wave Propagation in Glaciated Ter-  
828 rain. *Nature Communications*, 53(9):1689–1699, 2020. ISSN 1098-6596. doi: 10.1017/  
829 CBO9781107415324.004.

830 T. Winder, C. Bacon, J. D. Smith, T. S. Hudson, J. Drew, and R. S. White. QuakeMigrate v1.0.0  
831 (v1.0.0). *Zenodo*, 2021. doi: 10.5281/zenodo.4442748.

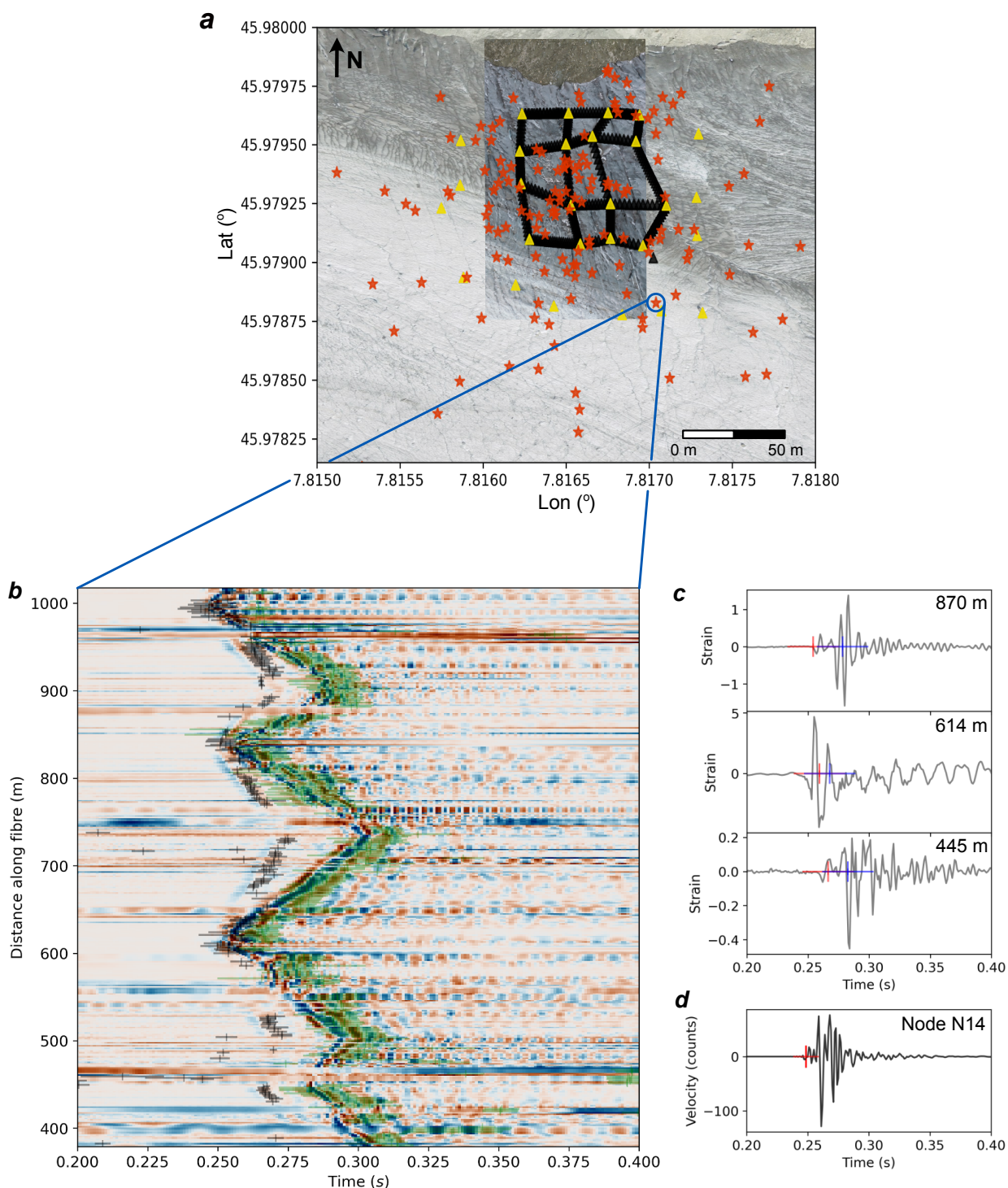
832 M. Withers, R. Aster, C. Young, J. Beiriger, M. Harris, S. Moore, and J. Trujillo. A comparison  
833 of select trigger algorithms for automated global seismic phase and event detection. *Bulletin of  
834 the Seismological Society of America*, 88(1):95–106, 1998. ISSN 00371106.

835 W. Zhu and G. C. Beroza. PhaseNet: A deep-neural-network-based seismic arrival-time picking  
836 method. *Geophysical Journal International*, 216(1):261–273, 2019. ISSN 1365246X. doi:  
837 10.1093/gji/ggy423.

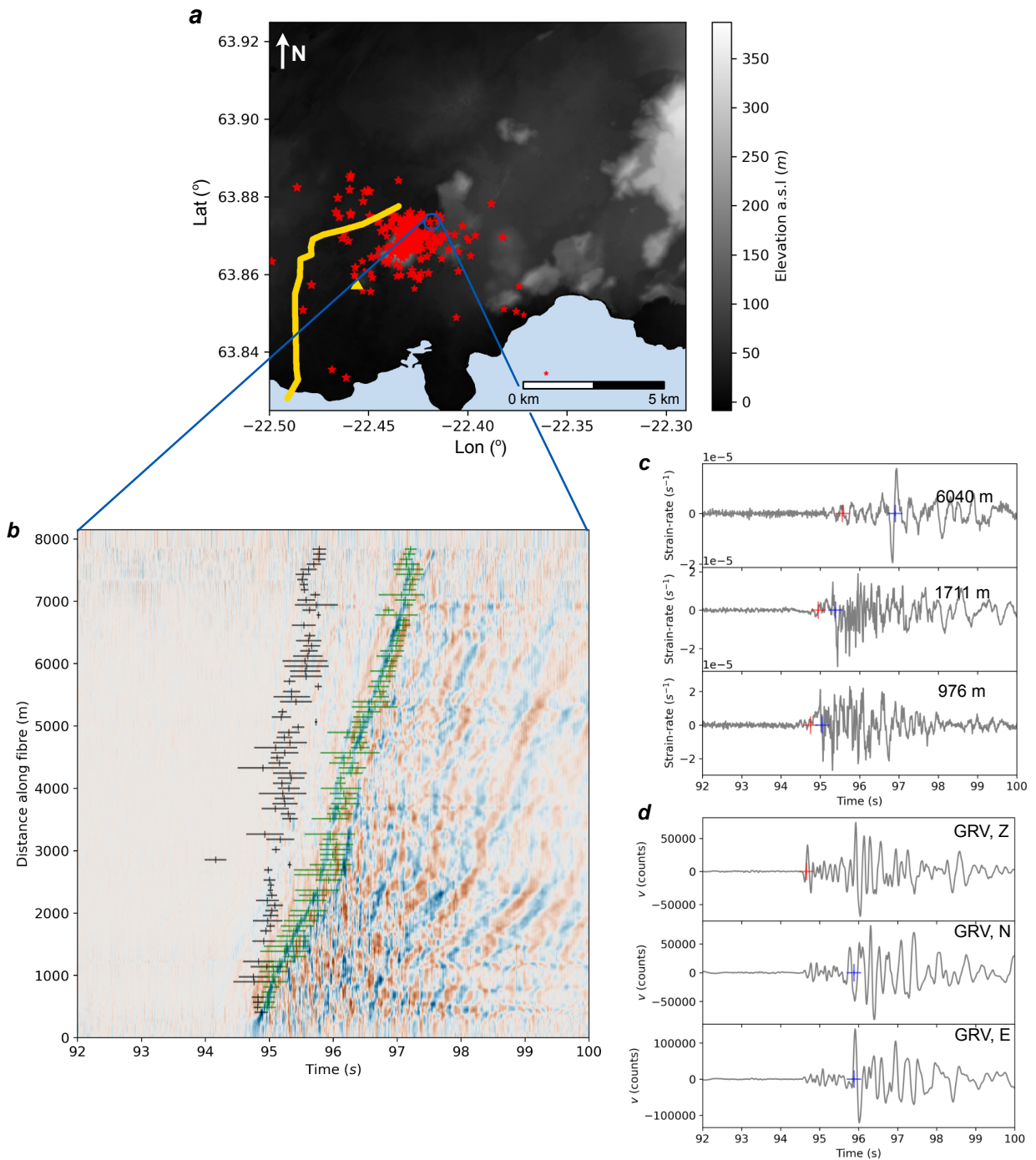
838 W. Zhu, E. Biondi, J. Li, J. Yin, Z. E. Ross, and Z. Zhan. Seismic arrival-time picking on  
839 distributed acoustic sensing data using semi-supervised learning. *Nature Communications*, 14  
840 (1):1–11, 2023. ISSN 20411723. doi: 10.1038/s41467-023-43355-3.



**Figure 6.** Effect of fibreoptic sensitivity on detection performance. a,b. Coalescence maps for a glacier icequake example, without and with sensitivity compensation, respectively. c. P-wave sensitivity map for the Gornergletscher network corresponding to the results of (b). d. Comparison of maximum single-pixel coalescence values through time with and without sensitivity compensation. e,f. Coalescence maps for a volcano-tectonic example, without and with sensitivity compensation, respectively. g,h. P- and S- wave sensitivity maps, respectively, for the Reykjanes Peninsula network corresponding to the results of (e). For each channel and seismic phase, regions with sensitivities  $< 0.1$  are masked for individual receivers. i. Same as (d) but for the volcano-tectonic example.

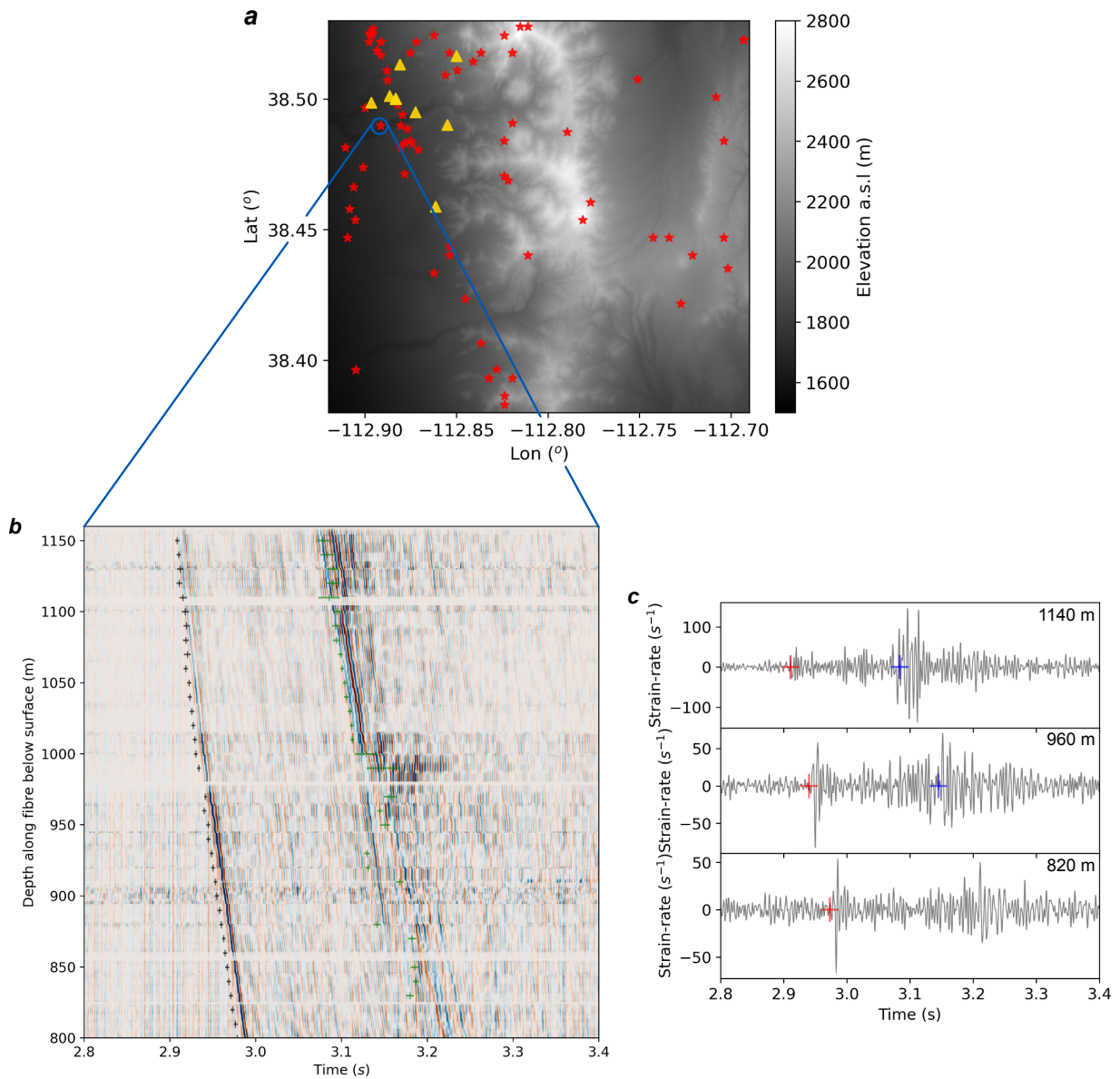


**Figure 7.** Glacier example from a crevasse field at Gornergletscher, Switzerland. a. Gornergletscher icequake catalogue (2023-10-22T04:00:00 to 2023-10-22T10:00:00). Red stars are icequakes, black triangles are DAS channels and gold triangles are seismic nodes. Background images are from an unmanned aerial vehicle during field deployment and from Swiss Topo (accessed: 10<sup>th</sup> June 2024). b. Example icequake signal measured on the fibre. Colours are normalised strain. Black and green scatter points show P-wave and surface-wave arrival time picks, respectively, by QuakeMigrate. c. Phase arrival picks for P-waves (red) and surface-waves (blue) for three channels in more detail. d. Example of P-wave arrival detected on a seismic vertical single-component node for comparison.



**Figure 8.** Volcanic eruption example from Sundhnúkagíggar on the Reykjanes Peninsula, Iceland. a. Earthquake catalogue from one eruptive period (2023-12-18T00:00:00.0 to 2023-12-20T00:00:00.0). Fibre receivers are shown by gold line, seismometer receiver is shown by gold triangle, earthquakes detected and located in this study are red stars and matched Icelandic Met Office events are shown by blue stars. Digital elevation model is from the ArcticDEM [Porter, 2023]. b. Example earthquake signal measured on the fibre. Colours are normalised strain-rate. Black and green scatter points show P-wave and S-wave arrival time picks, respectively, by QuakeMigrate. c. Phase arrival picks for P-waves (red) and S-waves (blue) for three channels in more detail. d. Example of P-wave and S-wave arrivals detected on a conventional seismometer operated by IMO, for comparison.





**Figure 9.** Geothermal borehole example from the Utah FORGE experiment, US. a. Earthquake catalogue from a one hour period during stimulation (17:00 to 18:00 on 27<sup>th</sup> April 2019). Receivers are shown by gold triangles and earthquakes detected and located in this study are red stars. Digital elevation model is from the NASA Shuttle Radar Topography Mission (SRTM). b. Example earthquake signal measured on the fibre. Colours are normalised strain-rate. Black and green scatter points show P-wave and S-wave arrival time picks, respectively, by QuakeMigrate. c. Phase arrival picks for P-waves (red) and S-waves (blue) for three channels in more detail.