

Gravity affects magma-induced crustal deformation: comparing laccoliths on the Moon, Mars, and Earth

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2 **comparing the Moon, Mars, and Earth**

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9 **Key Points:**

- 10 • Inflation of a magmatic laccolith at gravity of the Moon, Mars, and Earth was
11 implemented in the two-dimensional Discrete Element Method.
- 12 • Simulations show that the magma-induced fracturing accumulates in narrower zones in
13 weaker rock at lower gravity.
- 14 • Simulations show that laccolith inflation induces more surface displacement over a wider
15 area at lower gravity, and opens longer and wider cracks in the crust.

16 **Abstract**

17 Dome-shaped, uplifted surface areas and associated fractures on Mars and the Moon are inferred
18 to result from the shallow emplacement of magma intrusions. This inference originates from
19 analogue observations at partially eroded or active volcanic systems on Earth. Computational
20 models help estimating the geometry and emplacement depth of those inferred magma bodies.
21 Models often do not consider that the gravitational acceleration is different on planetary bodies
22 with different masses, however, and have not simulated large concentrations of magma-induced
23 strain and dynamic fracturing of the host rocks. We used the two-dimensional Discrete Element
24 Method (2D DEM) to simulate the inflation of a laccolith-shaped magma intrusion in particle-
25 based assemblages of different mechanical strength, under gravitational acceleration of the
26 Moon, Mars, and Earth. The 2D DEM model simulates the magma-induced displacements,
27 principal stresses, and dynamic fracturing, and allows deriving shear strains in the crust. For weak
28 rocks, the vertical surface displacement is nearly twice as high on the Moon, compared to Earth.
29 For stronger rocks, the amount of magma-induced cracks on the Moon is half of the amount of
30 cracks induced on Earth. Our 2D DEM simulations show, for the first time, that gravity specific to
31 a rocky planetary body affects the pattern and amount of fracturing and surface displacement
32 above inflating laccoliths. This calls for a careful reevaluation, and future modelling, of
33 differences seen in the morphology of intrusive domes found on Earth, Mars and the Moon.

34

35 **Plain Language Summary**

36 Domed surface areas have been observed on the Moon and on Mars, but how they formed
37 remains cryptic. By analogy with similar features in volcanic areas on Earth, at least some of these
38 features were likely caused by stalling and inflation in the shallowest few kilometers of crust of
39 laccoliths, which are magma intrusions with a convex roof and horizontal base. Gravitational
40 acceleration is lower on the Moon than on Mars, and is lower on Mars than on Earth. With those
41 differences in specific gravities, can we expect that similar amounts and patterns of displacement
42 and fracturing at the lunar and martian surface were caused by laccoliths of similar dimensions?
43 To address this question, we used the Discrete Element Method to simulate the displacement
44 and, for the first time, progressive fracturing of host rocks around an inflating laccolith intrusion
45 at specific gravities of the Moon, Mars, and Earth. Our results show that laccolith intrusions of
46 similar volume cause more surface displacement on a planetary body with weaker gravity,
47 compared to one with stronger gravity. The amount of magma-induced fracturing, however, is
48 more affected by the host rock's strength than by gravity. We conclude that magma-induced
49 deformation patterns depend on a planetary body's specific gravity and the strength of its crust.
50 Both features need to be accounted for when inferring magma intrusion properties from surface
51 observations on other planets and moons.

52

53

54 **1 Introduction**

55 Dome-shaped surface features in volcanic terrains have been related by geological observations
56 to the emplacement of thick magma intrusions in the shallowest several kilometers of the Earth's
57 crust (Bunger & Cruden, 2011; Gilbert, 1877; Morgan, 2018; Pollard & Johnson, 1973). Such thick
58 intrusions include sills (tabular, horizontal intrusions between preexisting layers of rock) and
59 laccoliths (intrusions with a horizontal base and a convex-upward roof). The crustal stresses
60 associated with the inflation of those intrusions induce bending of the overburden rocks that
61 translates in a dome-shaped uplift patterns at the surface, but also host rock compaction,
62 fracturing and faulting (Mattsson et al., 2018; Wilson et al., 2016). The scarcity of recent and well-
63 monitored laccolith intrusion events makes understanding their intrusion dynamics challenging.
64 The 2011-2012 emplacement and explosive eruption of a viscous rhyolite laccolith at Puyehue
65 Cordón Caulle, in Chile, is the only recent monitored event where dome-shaped surface uplift
66 and tensile surface fracturing were observed through satellite imagery and sparse seismic
67 monitoring (Castro et al., 2016). To understand laccolith- and sill-induced structural deformation
68 thus relies on geological and geophysical observations at now-solidified and exposed volcanic
69 and igneous plumbing systems, and modelling (Bunger & Cruden, 2011; Galland et al., 2018;
70 Magee et al., 2018; Morgan, 2018).

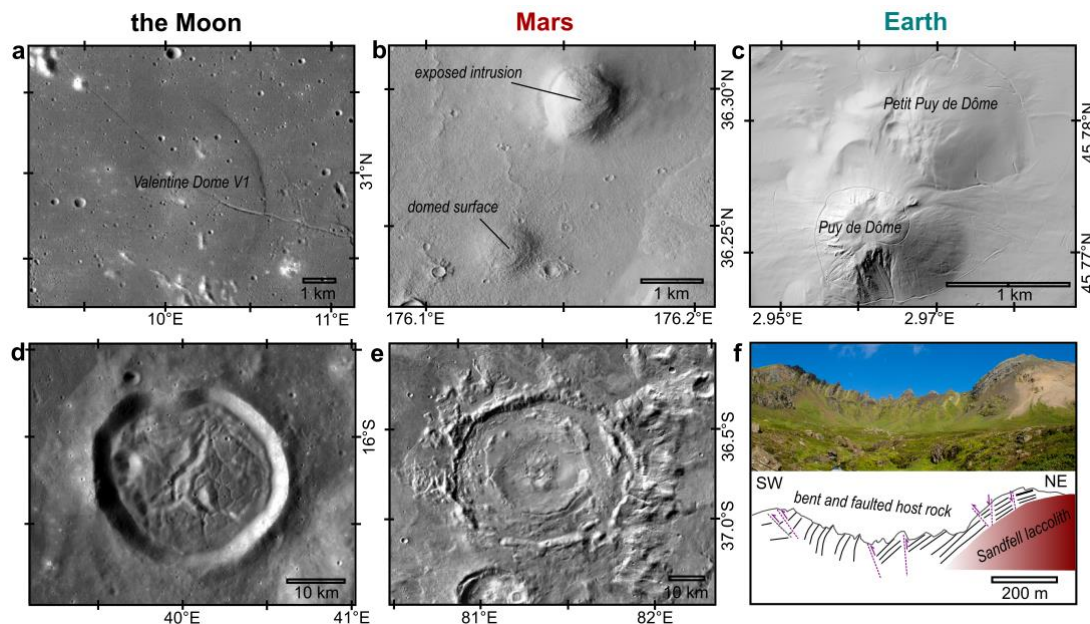
71
72 Areas with a dome-shaped positive topography have been observed at the surface of the Moon
73 since the Apollo missions in the early 1970's. Most of those domes have rough surfaces and
74 summit pits that suggest they are the result of the extrusion of viscous lava (Ivanov et al., 2016;
75 Lena et al., 2013; Wilhelms, 1987). Some, however, have a gentle slope, open fractures and no
76 eruptive vents; these are likely induced by magma emplacement in the shallowest few kilometers
77 of the lunar crust without eruption (Head & Wilson, 2017; Lena et al., 2013; Wöhler & Lena,
78 2009). One of the most known intrusive domes are the Valentine domes in Mare Serenitatis (Fig.
79 1a). Lunar intrusive domes range in diameter from less than one kilometer to more than 30
80 kilometers (Lena et al., 2013). Their relative scarcity is explained by the negative buoyancy of
81 dense mafic magma in the shallowest kilometres of the less dense lunar crust, which is
82 dominantly composed of anorthosite rocks and highly porous due to a long impact history
83 (Wieczorek et al., 2013; Wöhler & Lena, 2009). Lunar intrusive domes also often lie in areas
84 between proximal mare-filled impact basins where crustal extensional stresses might favor
85 magmatic buoyancy and ascent into shallower crustal levels (McGovern et al., 2014; Thomas et
86 al., 2015).

87
88 On Mars, high-resolution orbital imagery has allowed the detection of dome-shaped terrain in
89 regions with distributed volcanism away from the largest Martian volcanoes, but no global-scale
90 overview exists to date (Farrand et al., 2011; Platz et al., 2015; Rampey et al., 2007). Most of the
91 magmatism on both Mars and the Moon is considered dominantly mafic with low viscosity (Head
92 & Wilson, 2017; Platz et al., 2015; Vaucher et al., 2009). Some of the intrusive domes on Mars,
93 however, lie adjacent to intrusions exposed by erosion (Fig. 1b), eruptive cones, apparent
94 pyroclastic deposits and thick flows of felsic, viscous lava (Brož et al., 2015; Farrand et al., 2021).
95 This association makes it likely that thick sills or laccoliths may have intruded and inflated in the
96 shallowest few kilometers of the Martian crust. Models have confirmed the possibility of

97 cryptodome and laccolith formation on Arcadia Planitia, albeit in interaction with now-molten
 98 ice sheets (Farrand et al., 2011; Michaut et al., 2013).

99

100 Other sites of inferred magma-induced surface doming and fracturing on Mars and the Moon are
 101 found at ‘floor-fractured craters’ (FFCs) (Fig. 1d, 1e) (Jozwiak et al., 2012; Michaut, 2011; Schultz,
 102 1976). Floor-fractured crater morphology has been explained in the past by post-impact viscous
 103 relaxation of the crust (Bamberg et al., 2014; Hall & Solomon, 1981; Montigny et al., 2022).
 104 Models have shown, however, that impact cratering provides a lithospheric stress deficit and
 105 driving overpressure for magma to ascend below the crater floor (Michaut & Pinel, 2018). There,
 106 the magma may then have inflated a sill or laccolith and further fractured and uplifted the
 107 overburden rocks (Michaut et al., 2020; Michaut & Pinel, 2018; Wöhler & Lena, 2009).
 108 Observations of gravity anomalies by NASA’s Gravity Recovery and Interior Laboratory (GRAIL)
 109 have confirmed the shallow presence of dense magma intrusions below floor-fractured craters
 110 on the Moon (Wieczorek et al., 2013).



111 **Fig. 1** Magma-induced dome-shaped uplift features on three planetary bodies of different mass:
 112 **a.** Valentine 1 dome in Mare Serenitatis on the Moon (Lunar Reconnaissance Orbiter (LRO), NASA);
 113 **b.** domed surface feature and exposed intrusion East of Phlegra Montes on Mars (Context Camera
 114 (CTX), NASA); **c.** hillshade relief image of the intrusive dome Petit Puy de Dôme and extrusive lava
 115 dome Puy de Dôme, Chaine-des-Pûys, France (LIDAR dataset www.craig.fr); **d.** lunar floor-
 116 fractured crater Bohnenberger (LRO, NASA); **e.** martian floor-fractured crater 28-000072 (Themis,
 117 NASA); **f.** image and interpretive sketch of the exposed Sandfell laccolith and bent and faulted
 118 host rocks in East-Iceland (after Mattsson et al., 2018).

119
 120 Existing analytical and numerical models assume a linearly-elastic response of the lunar and
 121 martian crust to magma-induced stresses (Bunger & Cruden, 2011; Grosfils et al., 2015; Michaut,
 122 2011; Michaut & Pinel, 2018; Pollard & Johnson, 1973; Thorey & Michaut, 2016; Wöhler & Lena,
 123 2009). Linearly-elastic behavior implies that rocks deform instantaneously and proportionally to

124 stress, and may reverse to its initial state once the stress source is removed. Once the critical
125 stress is exceeded, tensile fractures open and the rock ruptures (Jaeger et al., 2007; Segall, 2010).
126 The geological observations at terrestrial laccoliths, however, such as at Sandfell in Iceland (Fig.
127 1f), show the importance of fracturing and faulting, besides elastic bending, of the overlying rocks
128 (Mattsson et al., 2018; Wilson et al., 2016). The dome-shaped, uplifted surfaces
129 to enter text. are often heavily dissected by faults and fractures, such as at the Petit Puy de Dôme
130 in the Chaîne des Puys volcanic field in France (Fig. 1c) (Petronis et al., 2019).

131
132 The role of non-elastic deformation during intrusion inflation has been documented in scaled
133 laboratory experiments where surface deformation can be directly related to the intrusion of
134 analog magma as thick sills, laccoliths and cryptodomes in granular, cohesive materials (Currier
135 & Marsh, 2015; Montanari et al., 2017; Poppe et al., 2019; Schmiedel et al., 2017). Furthermore,
136 magnetic crystal fabrics in now-solidified and exposed viscous magma in Iceland and Argentina
137 show that laccoliths can grow through repetitive magma pulses of high ascent rates and week-
138 to month-long pauses during which the magma partially cools and solidifies (Burchardt et al.,
139 2019; Mattsson et al., 2018). Such crustal accommodation through displacement, straining and
140 fracturing often allows for more than doubling of the intruded magma volume without eruption.

141
142 The geological, geophysical and experimental observations outlined above have motivated the
143 implementation of more complex stress responses though, such as elastoplasticity (Daniels et al.,
144 2012; Scheibert et al., 2017). High concentrations of strain, and dynamic fracturing, can be
145 simulated in the Discrete Element Method (DEM) (Cundall & Strack, 1979). The DEM discretizes
146 a medium into an assemblage of spherical disks or spheres of which the position at each timestep
147 is calculated according to Newton's laws of motion (Cundall & Strack, 1979). Unlike Finite Element
148 Models (FEM), the DEM allows for the concentrations of large strains, discontinuous deformation
149 and the dynamic opening and propagation of tensile fractures and shear bands. The particle
150 assembly can thus be modelled to respond mechanically similarly to natural rock (Cundall &
151 Strack, 1979; Potyondy & Cundall, 2004; Schöpfer et al., 2009). DEM has been used to simulate,
152 amongst others, lava dome effusion and stability, hydraulic fracture propagation, or caldera
153 collapse (Harnett et al., 2020; Harnett & Heap, 2021; Holohan et al., 2015; Huang et al., 2022;
154 Morgan & McGovern, 2005; Woodell et al., 2023). Using two-dimensional (2D) DEM simulations,
155 Morand et al. (2024) recently showed how mechanical properties control the surface
156 displacement magnitudes and fracture network development during laccolith growth on Earth.

157
158 Accounting for nonelastic deformation and dynamic fracturing during laccolith inflation may be
159 especially impactful for planets and moons smaller than Earth, where lower global mass results
160 in lower gravitational acceleration compared to that on Earth. Compilations of rock deformation
161 experiments have shown that reduced surface gravity leads to increased porosity and,
162 consequentially, reduced brittle and tensile strength for similar depths (Heap et al., 2017).
163 Models that assume linear elasticity during laccolith inflation (Bunger & Cruden, 2011; Michaut
164 & Pinel, 2018; Walwer et al., 2021; Wöhler & Lena, 2009), have thus not clarified what the effect
165 is of specific gravity on the development of crustal fracturing and surface displacements during
166 the inflation of a laccolith.

167

168 We present simulations of the inflation of a laccolith at one kilometer depth under gravitational
169 acceleration for Earth, Mars and the Moon with a 2D DEM model. We vary the strength of the
170 host particle assemblage to represent a range of planetary crustal strengths. In this way, our 2D
171 DEM implementation allows, for the first time, to discuss the effect of gravity as a function of
172 planetary body mass and crustal strength on the magma-induced crustal deformation, dynamic
173 fracturing, and surface displacement patterns. Click or tap here to enter text.

174 **2 Method**

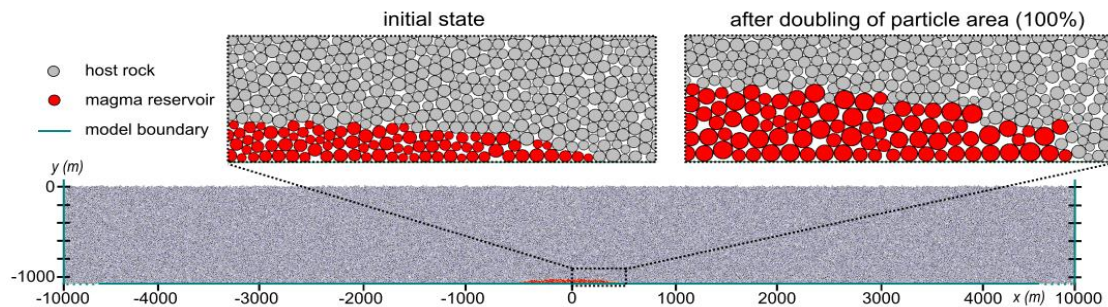
175 2.1 Model set-up

176 We implemented the inflation of a laccolith-shaped magma intrusion in the commercial two-
177 dimensional (2D) DEM Particle Flow Code (PFC2D 7.0) from Itasca Consulting Group, Ltd.
178 (<https://www.itascacg.com/software/PFC>). Spherical particles with radii normally distributed
179 around a mean of 6.65 ± 1.65 m were generated in a rectangular model domain of 20 km wide
180 and 1 km high, constrained by lateral and bottom walls and with a free upper surface (Fig. 2). We
181 assigned particles either to a class of rock, or magma. Contacts between two rock particles and
182 between rock particles and the model walls were governed by the soft-bond model (Jiang et al.,
183 2015; Ma & Huang, 2018a). This contact model represents a nonelastic solid medium defined by
184 a bond effective modulus E^* , bond tensile strength, bond cohesion (kept 10 times the bond
185 tensile strength in our model such that contacts only break in tension), and two softening
186 parameters (Ma & Huang, 2018b). All other numerical parameters were kept constant in our
187 model (see Supplementary Table S1). The two softening parameters were kept constant to
188 approach a 10:1 ratio of unconfined compressive strength to tensile strength in accordance with
189 laboratory experiments on intact rock samples (Heap et al., 2021). We further followed the
190 method of initialization of a stable particle assemblage as described by Harnett & Heap (2021)
191 and Morand et al. (2024).

192
193 The absolute values of the rock particle bond strength parameters differ from the bulk strength
194 of the simulated rock. Here, we simulated a range of realistic strengths of planetary crust, not
195 one particular rock strength. To define the particle bond strength parameters needed to obtain
196 the desired bulk properties, we performed numerical uniaxial compressive strength (UCS) and
197 tensile strength (TS) tests (Harnett & Heap, 2021; Ma & Huang, 2018a; Morand et al., 2024). The
198 obtained UCS, TS and Young's modulus (E) values correspond to values measured on intact
199 samples of natural rocks (Table 1), and are used to label figures 3 – 6. The DEM bond strength
200 parameter values affect two properties of the rock strength. The toughness relates to the
201 resistance of the rock to cracking (mainly controlled by the bond cohesion and bond tensile
202 strength at a 10:1 ratio); tougher rocks can accommodate more strain before breaking. The
203 stiffness relates to the resistance of the rock to elastic bending (mainly controlled by the bond
204 effective modulus and ratio of normal to shear stiffness); stiffer rocks require more force to
205 obtain a given amount of displacement and strain. We simulated six combinations of three rock
206 toughnesses (TS of 0.8 ± 0.1 MPa, 1.9 ± 0.2 MPa, or 3.7 ± 0.1 MPa) for two rock stiffnesses
207 (Young's modulus of 2.5 ± 0.1 GPa or ~ 8.4 GPa) (Table 1). These ranges represent ranges of
208 realistic elastic moduli, tensile strength and unconfined compressive strength obtained for intact

209 rock samples scaled up to that of the weaker crust of a rocky planetary body, that is fractured
 210 and layered (Heap et al., 2017, 2020, 2021).

211
 212 Magma particles were assigned within a body with a horizontal base with a width of 1,000 m,
 213 and a half-ellipse shaped convex roof with a top height of 50 m (Fig. 2). This laccolith was
 214 positioned at the bottom center of the model, and with its central top 1,000 m deep below the
 215 upper free surface. Contacts between two magma particles, between a rock and magma particle,
 216 and between magma particles and the lower model interface were defined by the linear parallel
 217 bond model, with zero friction, zero cohesion, zero tensile strength but a high effective modulus
 218 E^* . This approach approximates the behavior of incompressible magma (Morand et al., 2024).



219 **Fig. 2** Model set-up implemented in the two-dimensional Discrete Element Method (2D DEM) in
 220 PFC2D (Itasca Ltd.) with an assemblage of host rock particles (grey) and a particle-based, half-
 221 ellipsoid-shaped magma body (red) with its top at 1 km depth. The subset images display the
 222 initial state and final state of the magma body after doubling of the magma particle area, i.e.
 223 after 100% of laccolith inflation.

224 2.2 Model exploitation

225 We aimed to investigate the effect of the inflation of a laccolith on the shallow crust in function
 226 of gravity. We thus ran 2D DEM simulations that are subjected to the gravitational acceleration
 227 constant g of Earth (9.81 m.s^{-2}), Mars (3.71 m.s^{-2}), or the Moon (1.62 m.s^{-2}). We incrementally
 228 increased the area of the circular magma particles in the laccolith in steps of 1% until an area
 229 inflation of 100% was obtained (Fig. 2). Assuming an axial intrusion symmetry, this represents the
 230 doubling of laccolith volume from $\sim 0.1 \text{ km}^3$ to $\sim 0.2 \text{ km}^3$. This approach of particle area change is
 231 similar to that used to simulate the deflation of a magma body (Holohan et al., 2011, 2015), but
 232 differs from that of Harnett et al. (2020) and Morand et al. (2024), wherein new particles were
 233 added from below to inflate the lava dome or laccolith.

234
 235 The inflating laccolith exerted stress on the surrounding rock particles, which were displaced. As
 236 a consequence, rock-rock contacts were strained. Particles resisted relative rotation which
 237 simulated a rigid interface between them. Once strained rock particle bonds failed, the soft-bond
 238 model was replaced by a rolling-resistance model that was only governed by a friction coefficient
 239 of 0.5 but with zero bond cohesion or bond tensile strength. This is analogue to strain-induced
 240 opening of tensile cracks and displacement of the detached rock masses along the crack plane in
 241 natural rock.

242 2.3 Model visualization and analysis

243 Particle position, velocity, displacement, bond stresses and bond cracking are tracked throughout
244 the model. By convention, upward and rightward displacements are positive; downward and
245 leftward displacements are negative. Cracks are formed by bond breakage (i.e. when the tensile
246 bond strength is exceeded). In post-processing, the displacement gradient tensor is used to
247 determine the finite shear strain from the Cauchy-Green deformation tensor (Schöpfer et al.,
248 2006), further described by Harnett et al. (2018) and Morand et al. (2024). To emphasize the
249 relative shear strain distribution within a simulation, we normalize all finite shear strain values to
250 the simulation's maximum. We further apply a cutoff criterion to all plots wherein all normalized
251 finite shear strain values higher than 0.05% are set to 0.05%, because more than 90 % of finite
252 shear strain values lies below that value. We express the number of cracked bonds as a
253 percentage of the initial amount of rock particle bonds. To determine the maximum lateral extent
254 of the vertically displaced surface in Table 1, all surface particles are included that are vertically
255 displaced by more than 0.5% of the maximum vertical displacement value detected in the final
256 time step of the simulation.

257 3 Results

258 3.1 Influence of gravity on magma-induced deformation

259 We systematically analyzed the patterns of magma-induced displacement (Fig. 3), and of vertical
260 and horizontal displacement at the surface (Fig. 4). We provide figures displaying the final model
261 outcome, the total displacements, stress, strain, and crack patterns for each simulation in the
262 Online Supplementary materials.

263 We first discuss the displacement results irrespective of the value of gravitational acceleration.
264 All displaced overburden rock areas are the narrowest near the magma-rock contact and the
265 widest at the surface, and are laterally delimited by one or a series of en-echelon fractures. The
266 inward dip of those fractures is the steepest for weak rocks (Fig. 3a-c, 3j-l) and the most gentle
267 for the toughest rocks (Fig. 3g-i, 3p-r). Consequently, the displaced area is the widest in the
268 toughest rock (Fig. 3g-i, 3p-r). In those toughest rocks, the displacement magnitude decreases
269 continuously laterally outwards, we observe the longest and widest fractures, and also the widest
270 displacement zones at the surface (Fig. 4i-l). Weaker rocks display discontinuities along which the
271 displacement magnitude abruptly decreases laterally outwards from the model centre (Fig. 3a-
272 i). Under these conditions, the vertical and horizontal displacement varies stepwise (Fig. 4a-h), in
273 contrast with the continuous decrease laterally outwards in the toughest rocks (Fig. 4i-l). The
274 locations of those stepwise variations in surface displacement magnitude correspond to open
275 fractures at the model surfaces seen in Figure 3.

276 We then discuss the displacement results in function of specific gravity. The inward dip of the
277 fractures that laterally delimit the displaced rock is the steepest for the highest gravity (Earth)
278 (e.g., Fig. 3i) and the most gentle for the lowest gravity (the Moon) (e.g., Fig. 3g). Consequently,
279 the displaced area is the widest at the lowest gravity where the bounding fractures do not reach
280 the surface any longer but propagate horizontally away from the inflating laccolith (Fig. 3g, 3p).
281 At that lowest gravity, the maximum of the displacement magnitude is the highest at 50 m, the

282 extent of displacement at the surface is the widest and the slope of the vertical surface
 283 displacement curve is declining the steepest away from the center (Fig. 4, black curves). Under
 284 these lower-gravity conditions, the stepwise variations in horizontal displacement are also the
 285 largest across open surface fractures (Fig. 4b, 4d, 4f, 4h). In weaker rocks, the maxima of vertical
 286 displacement are 10-15 m higher at lower gravity on the Moon compared to those at Earth's
 287 higher gravity (Fig. 4a, 4e); the maxima of horizontal displacement are 4-6 m lower (Fig. 4b, 4f).
 288 There is less difference in absolute values at the domes' crests for stronger rocks across gravities
 289 (Fig. 4g-l). Displacement magnitudes and surface displacements for the intermediate gravity
 290 specific to Mars are overall intermediary between those on the Moon and on Earth. Only for
 291 Earth's higher gravity we observe a central block near the surface bound by normal faults with
 292 less positive vertical displacement than the dome flanks and near-zero horizontal displacement
 293 (Fig. 4a-f).

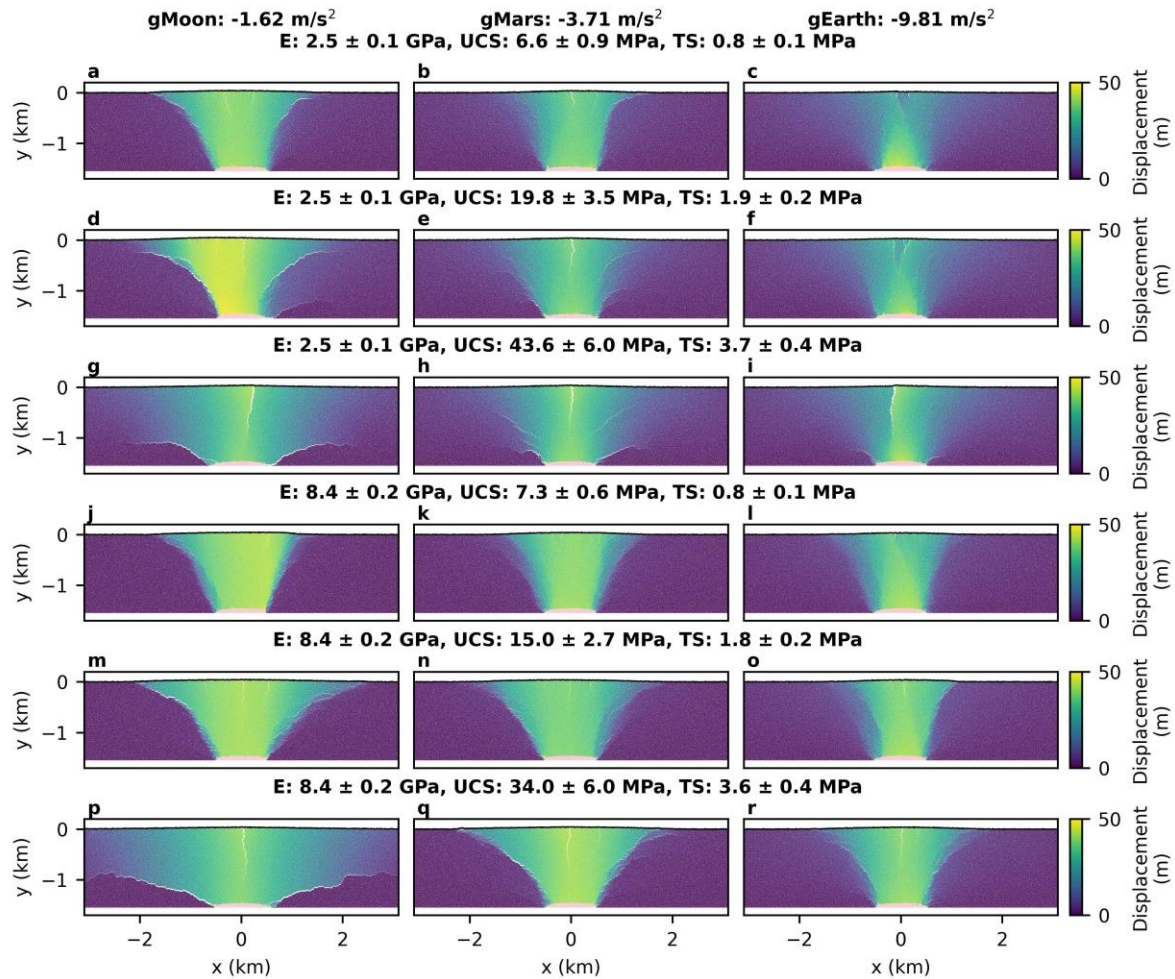
294

295 **Table 1.** Summary of bond strength and bulk strength values and key results for the 18 performed
 296 2D DEM simulations.

297

| Model setup | Bond parameters | | | Mechanical properties | | | | Model results | | |
|-------------|--------------------------------|----------|-----------|-----------------------|------------|-----------|------------|---------------|---|---------------------------------------|
| | Gravity g (m.s ⁻²) | E* (GPa) | Coh (MPa) | Ten (MPa) | E (GPa) | UCS (MPa) | TS (MPa) | UCS/ TS | Final U _y ^{max} (m) | Lateral extent of U _y (km) |
| -1.62 | 3.0 | 10.0 | 1.0 | 2.5 ± 0.1 | 6.6 ± 0.9 | 0.8 ± 0.1 | 8.3 ± 2.0 | 40.0 | 3.699 | Type 1 |
| -3.71 | 3.0 | 10.0 | 1.0 | 2.5 ± 0.1 | 6.6 ± 0.9 | 0.8 ± 0.1 | 8.3 ± 2.0 | 36.9 | 3.903 | Type 1 |
| -9.81 | 3.0 | 10.0 | 1.0 | 2.5 ± 0.1 | 6.6 ± 0.9 | 0.8 ± 0.1 | 8.3 ± 2.0 | 26.1 | 4.330 | Type 1 |
| -1.62 | 3.0 | 25.0 | 2.5 | 2.5 ± 0.1 | 19.8 ± 3.5 | 1.9 ± 0.2 | 10.4 ± 2.6 | 46.3 | 4.735 | Type 1 |
| -3.71 | 3.0 | 25.0 | 2.5 | 2.5 ± 0.1 | 19.8 ± 3.5 | 1.9 ± 0.2 | 10.4 ± 2.6 | 37.2 | 4.328 | Type 1 |
| -9.81 | 3.0 | 25.0 | 2.5 | 2.5 ± 0.1 | 19.8 ± 3.5 | 1.9 ± 0.2 | 10.4 ± 2.6 | 28.0 | 4.211 | Type 1 |
| -1.62 | 3.0 | 50.0 | 5.0 | 2.5 ± 0.1 | 43.6 ± 6.0 | 3.7 ± 0.4 | 11.8 ± 2.6 | 40.4 | 5.723 | Type 2 |
| -3.71 | 3.0 | 50.0 | 5.0 | 2.5 ± 0.1 | 43.6 ± 6.0 | 3.7 ± 0.4 | 11.8 ± 2.6 | 37.6 | 4.670 | Mixed |
| -9.81 | 3.0 | 50.0 | 5.0 | 2.5 ± 0.1 | 43.6 ± 6.0 | 3.7 ± 0.4 | 11.8 ± 2.6 | 37.2 | 4.191 | Mixed |
| -1.62 | 10.0 | 10.0 | 1.0 | 8.4 ± 0.2 | 6.6 ± 0.9 | 0.8 ± 0.1 | 8.3 ± 2.0 | 43.5 | 3.537 | Type 1 |
| -3.71 | 10.0 | 10.0 | 1.0 | 8.4 ± 0.2 | 6.6 ± 0.9 | 0.8 ± 0.1 | 8.3 ± 2.0 | 40.0 | 3.942 | Type 1 |
| -9.81 | 10.0 | 10.0 | 1.0 | 8.4 ± 0.2 | 6.6 ± 0.9 | 0.8 ± 0.1 | 8.3 ± 2.0 | 35.6 | 3.821 | Type 1 |
| -1.62 | 10.0 | 25.0 | 2.5 | 8.4 ± 0.2 | 19.8 ± 3.5 | 1.9 ± 0.2 | 10.4 ± 2.6 | 42.5 | 4.923 | Type 1 |
| -3.71 | 10.0 | 25.0 | 2.5 | 8.4 ± 0.2 | 19.8 ± 3.5 | 1.9 ± 0.2 | 10.4 ± 2.6 | 38.0 | 4.242 | Type 1 |
| -9.81 | 10.0 | 25.0 | 2.5 | 8.4 ± 0.2 | 19.8 ± 3.5 | 1.9 ± 0.2 | 10.4 ± 2.6 | 39.5 | 3.791 | Type 1 |
| -1.62 | 10.0 | 50.0 | 5.0 | 8.4 ± 0.2 | 43.6 ± 6.0 | 3.7 ± 0.4 | 11.8 ± 2.6 | 39.5 | 7.813 | Type 2 |
| -3.71 | 10.0 | 50.0 | 5.0 | 8.4 ± 0.2 | 43.6 ± 6.0 | 3.7 ± 0.4 | 11.8 ± 2.6 | 42.3 | 4.702 | Type 1 |
| -9.81 | 10.0 | 50.0 | 5.0 | 8.4 ± 0.2 | 43.6 ± 6.0 | 3.7 ± 0.4 | 11.8 ± 2.6 | 39.1 | 4.045 | Type 1 |

298



299

300 **Fig. 3** Displacement magnitude after doubling of the magma body area in the 2D DEM simulations
 301 subjected to gravity specific to the Moon, Mars, or Earth for a range of host rock Young's moduli
 302 (E), unconfined compressive strengths (UCS) and tensile strengths (TS). Specific gravity increases
 303 from left to right, rock strength increases from top to bottom. Rock stiffness is the lowest in the
 304 three top rows and the highest in the three bottom rows.

305

3.2 Influence of gravity on magma-induced strain and cracking

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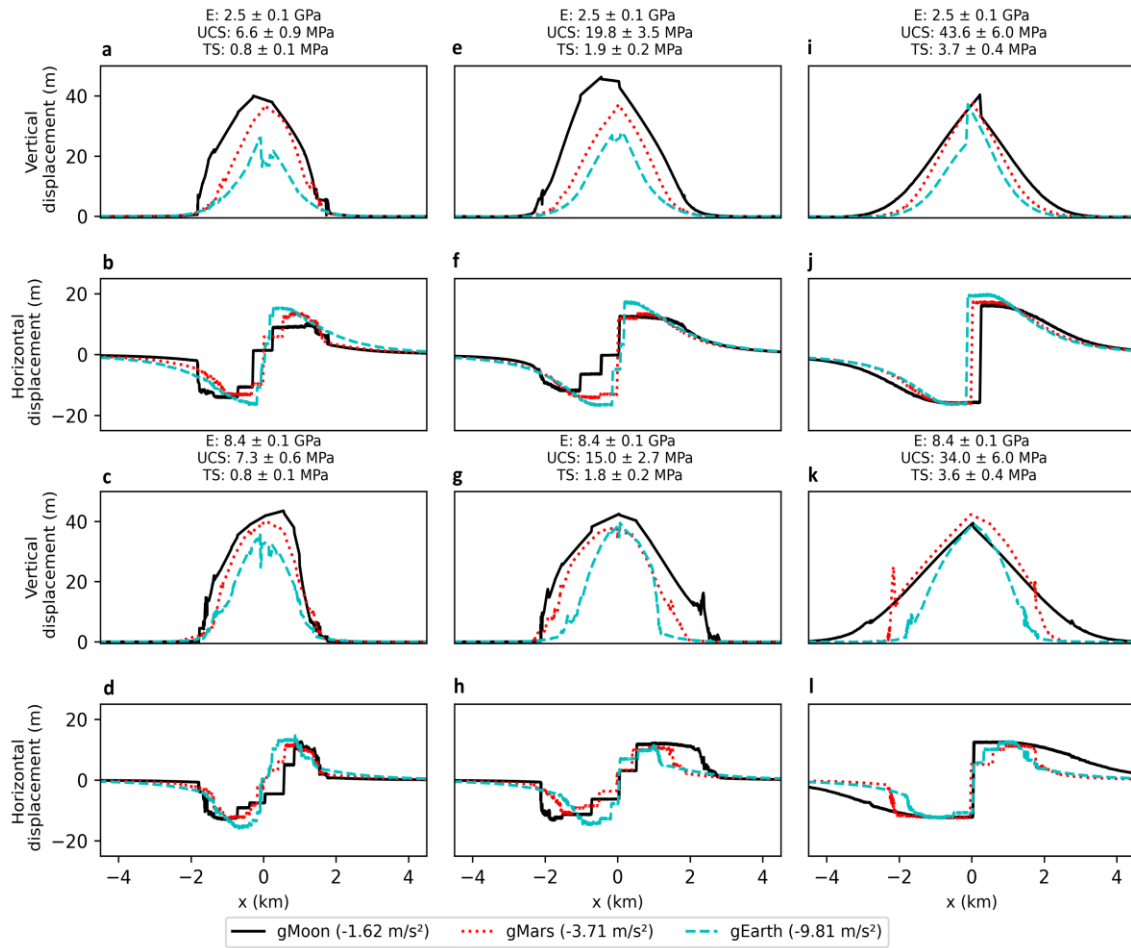
313

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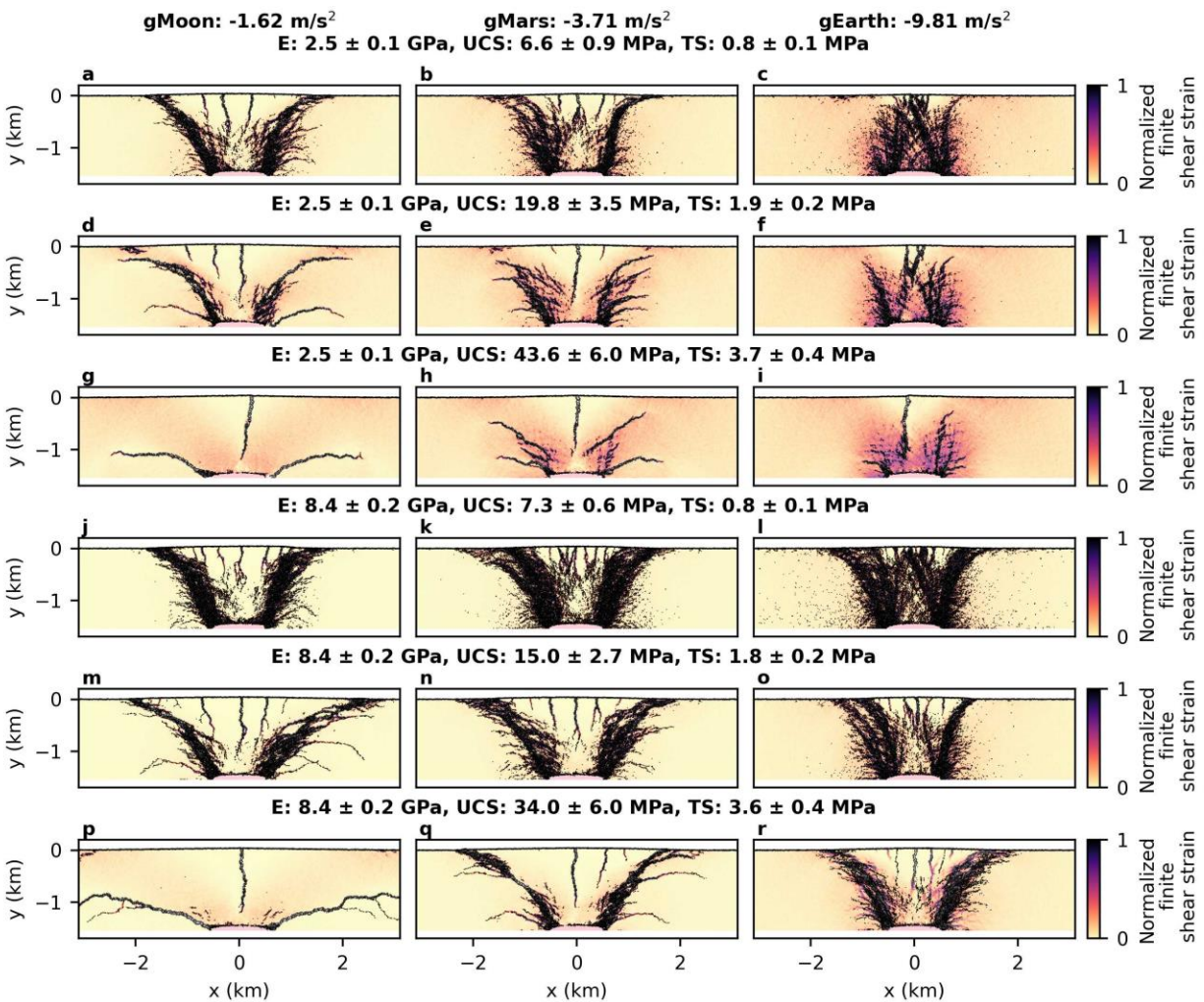
We first discuss the calculated normalized finite shear strain and cracking irrespective of the value of gravitational acceleration. For the same stiffness, less tough rocks accumulate the least shear strain and are the most cracked (Fig. 5a-c, 5j-l), whereas tougher rocks accumulate the most shear strain and crack the least (Fig. 5g-i, 5p-r). The fracture zones in which cracks concentrate are the widest in less tough rocks and the narrowest in tougher rocks. The zones of highest shear strain are the narrowest in less tough rocks and the narrowest in tougher rocks. As observed above, the fracture zones that delimit the deformed rock zone above the inflating laccolith are the steepest in less tough rocks (Fig. 5a-c, 5j-l), and the most gentle to subhorizontal in the tougher rocks (Fig. 5g-i, 5p-r). Quantitatively, less cracks open in less stiff rocks (up to 3.5% of initial bonds broken) (Fig. 6a-c), compared to more cracks in stiffer rocks (up to 6.5% of initial

316 bonds broken) (Fig. 6d-f). For the same stiffness, more cracks open in less tough rocks (4.5% –
 317 6.5%) (Fig. 6a, 6d), compared to less cracks in tougher rocks (0.5% – 2.5%) (Fig. 6c, 6f).
 318



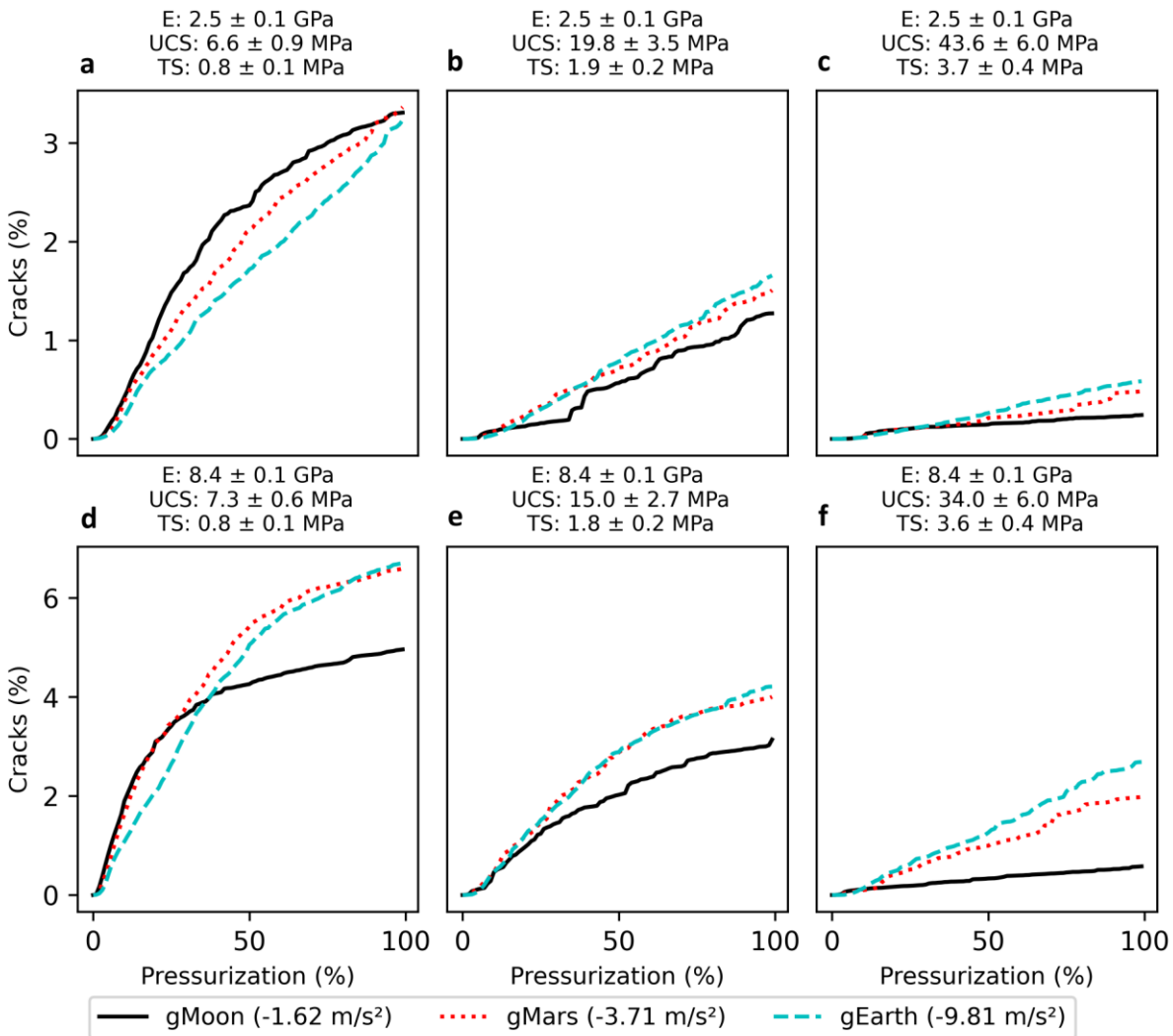
319
 320 **Fig. 4** Horizontal and vertical surface displacement after doubling of the magma body area in 2D
 321 DEM simulations subjected to gravity specific to the Moon (black), Mars (red), or Earth (blue) for
 322 a range of host rock Young's moduli (E), unconfined compressive strengths (UCS) and tensile
 323 strengths (TS). Toughness increases from the left column to the right column, stiffness increases
 324 from the top two rows to the bottom two rows. Gravity specific to each planetary body is indicated
 325 by line color and dashed.

326 We then discuss the calculated normalized finite shear strain and cracking in function of specific
 327 gravity. For the same stiffness, rocks accumulate the least shear strain and are the least cracked
 328 at lower gravity specific to the Moon (Fig. 5a, 5d, 5g, 5j, 5m, 5p), whereas rocks accumulate the
 329 most shear strain and crack the most at higher gravity specific to Earth (Fig. 5c, 5f, 5i, 5l, 5o, 5r).
 330 The zones that delimit the deformed rock area in which shear strain and cracks concentrate are
 331 the narrowest at lower gravity and the widest at higher gravity. As observed above, the fracture
 332 zones that delimit the deformed rock zone above the inflating laccolith are the most gentle to
 333 subhorizontal at the lowest gravity and the steepest at the highest gravity. Time series of finite
 334 shear strain and crack pattern development (see Supplementary videos) show that the tensile
 335 fracture that opens at the crest of the uplifted surface propagates the deepest with the largest
 336 width at the lowest gravity (compare Fig. 5g-i and Fig. 5p-r). At the higher gravity and stiffer rocks,
 337 narrow bands of high finite shear strain remain uncracked (Fig. 5r).



338 **Fig. 5** Normalized finite shear strain (see color bars) and cracked particle bonds (black lines) after
 339 doubling of the magma body area in the 2D DEM simulations subjected to gravity specific to the
 340 Moon, Mars, or Earth for a range of host rock Young's moduli (E), unconfined compressive
 341 strengths (UCS) and tensile strengths (TS).

342 Quantitatively, less cracks open at the lower gravity (Fig. 6, black curves), compared to more
 343 cracks opening at higher gravity (Fig. 6, blue curves). Crack amounts increases faster – the rock
 344 cracks more rapidly when stressed – at lower gravity than at high gravity in weak rocks (Fig. 6a),
 345 but in contract crack amounts increase faster at higher gravity in tougher and stiffer rocks (Fig.
 346 6b-c, 6e-f). The difference between the amount of cracks in function of gravity is the highest with
 347 2% for the strongest rocks (Fig. 6f). Notably, there is no significant difference in amount of cracks
 348 in function of specific gravity in the weakest rocks (Fig. 6a), but the spatial fracture and shear
 349 strain pattern is markedly different (Fig. 5a-c). Overall, Martian rocks respond in an intermediate
 350 manner to the inflation of the laccolith, compared to those on the Moon and Earth.
 351



352 **Fig. 6** Development of cracks (failed particle bonds) during the inflation of the magma body area
 353 in 2D DEM simulations subjected to gravity specific to the Moon, Mars, or Earth for a range of
 354 host rock Young's moduli (E), unconfined compressive strengths (UCS) and tensile strengths (TS).

355 **4 Discussion**

356 The results from our 18 2D DEM simulations show that the spatial distribution and intensity of
357 cracking of host rocks induced by the inflation of a laccolith of similar volume at similar depth lies
358 along a spectrum between two end-members. Morand et al. (2024) attributed this spectrum to
359 differences in the host rock toughness and stiffness. We show that the end-member spectrum is
360 also controlled by the difference in gravitational acceleration between planetary bodies of
361 different mass, here the Moon, Mars, and Earth. The first end-member develops a highly
362 fractured host rock, with multiple tensile fractures that open from the surface downward into
363 the host rock and two well-developed, inward-dipping shear fracture zones that develop from
364 the laccolith edges and surface as thrust faults that bound the uplifted surface area (e.g., Fig. 3r,
365 Fig. 5r). The second end-member develops a poorly fractured host rock, with one tensile fracture
366 that opens from the surface downward at the crest of the uplifted surface area, and two inward-
367 dipping to near-horizontal fractures that open from the intrusion edges outward (e.g.; Fig. 3p,
368 Fig. 5p). When similarly tough crust is assumed (TS of ~ 3.7 MPa in our models), the deformation
369 pattern approaches the first, highly fractured end-member rather for higher gravity on Earth and
370 Mars, and the second, poorly fractured end-member for lower gravity on the Moon (Fig. 3p-r,
371 Fig. 5p-r). For less tough crust (TS below 2.0 MPa in our models), we mainly see the most
372 displacements and a highly fractured first end-member especially for lower gravity on the Moon.
373

374 Our model results show, however, that the linearly-elastic assumption is not compatible with
375 syn-intrusive dynamic fracturing dynamics, wherein narrower bounding zones concentrate shear
376 strain and close-spaced cracking at lower gravity on the Moon than on Earth (Fig. 5). The fracture
377 and shear strain zones that develop from the edges of the inflating laccolith, grow upward and
378 truncate the surface as reverse faults, except in the second end-member. These bounding
379 fracture zones effectively laterally constrain the displaced and strained overburden rock, and
380 once that lateral constraint has developed, the overburden can be displaced upward in coherent
381 zones separated by one or more tensile fractures that propagate from the surface downward.
382 The width of the displaced surface area thus becomes delimited by the fracturing dynamics in
383 the host rock, which in turn is controlled by the gravity, and thus lithostatic loads at similar depth,
384 specific to the planetary body the laccolith is inflating in. We observe less cracking when laccoliths
385 inflate at lower gravity on the Moon, where the magma-induced stress is rather accommodated
386 by upwards crustal displacement than by fracturing. As a result, we observe more surface
387 displacement over a wider area on the Moon, compared to Earth, with Mars intermediate
388 between the two.
389

390 Standard models that invert the surface displacements to obtain the diameter, but also the depth
391 and opening of sills and laccoliths, simplify the crustal response to that of a homogeneous,
392 linearly-elastic half-space (Battaglia et al., 2013; Mogi, 1958; Okada, 1985). Those analytical
393 solutions do not involve a gravitational constant, whereas our model results demonstrate a
394 significant effect of gravity on the surface displacements. More recent analytical and finite
395 element models do include a gravitational constant (Fernández et al., 1997; Galland & Scheibert,
396 2013; Got et al., 2019; Michaut, 2011). Models at a similar Young's modulus of 10 GPa than ours
397 find more displacement over a wider area above laccoliths on the Moon compared to Earth

398 (Michaut, 2011). The slope of the domed surface is overall more gentle in those models for lower
399 gravity, however, than in our models. All these models ignore fracturing dynamics, or only include
400 a damage criterion (Got et al., 2019). It has been inferred that, because lunar intrusive surface
401 doming features are wider than terrestrial ones, magmatic laccoliths and sills in the shallow lunar
402 crust may be wider on the Moon than they are on Earth (Wöhler & Lena 2009). Our 2D DEM
403 results nevertheless show that concentrations of high strains and dynamic fracturing can narrow
404 the width of the displaced surface area, especially for less tough rocks. This finding raise the
405 expectation that those models may overestimate the width and depth of laccoliths on the Moon
406 and Mars as compared to Earth. Direct comparisons between existing models and our DEM
407 model, using the same strength properties, are now required to directly determine the difference
408 between model outcomes that include, or ignore, dynamic fracturing.

409
410 We noted in the results that only our models for Earth's gravity show a central area at the crest
411 of the domed surface that is bound by opposing normal faults and that has been uplifted less
412 than the surrounding area, i.e., a graben (Figs. 3, 4). The subsurface crack distribution is much
413 more distributed in these conditions compared to those at lower gravities (Fig. 5). Geological
414 observations at exposed igneous plumbing systems, seismic reflection data, and recent volcanic
415 events on Earth, have shown that such graben systems can form as elongated topographic
416 depressions above vertical, magma-filled fractures (dykes) (Magee & Jackson, 2021; Mastin &
417 Pollard, 1988; Sigmundsson et al., 2014; Smittarello et al., 2022). Existing analytical models have
418 used the geometry of such grabens observed on orbital imagery of the Moon or Mars to postulate
419 that the lower gravity there would favor wider opening of higher dykes (Ernst et al., 2001; Head
420 & Wilson, 1993; Klimczak, 2014). The height and opening of those tensile fractures in the crust of
421 Mars or the Moon, versus those on Earth, is not well constrained though (Klimczak, 2015). Tensile
422 fractures in our simulations open wider and propagate deeper from the surface downward on
423 the Moon, and to lesser on Mars, than they can on Earth, for similar amounts of laccolith inflation
424 (Fig. 3, Fig. 5). This progressive fracturing mechanism was not simulated previously, but the
425 relative easiness of propagating open fractures deeper into the crust, in combination with
426 impact-induced surface unloading (Michaut & Pinel, 2018; Wöhler & Lena, 2009), may have
427 favored the ascent of dense mafic magma to the surface despite a negative buoyancy contrast
428 with the porous shallow lunar crust.

429
430 Our new 2D DEM implementation will allow future simulations to investigate the interaction
431 between extensional crustal stresses and impact-induced unloading, and laccolith-induced
432 stresses. Simulations of broad ranges of laccolith thickness and width are also necessary to
433 investigate if our conclusions hold for the widths of tens of kilometers of intrusive domes and
434 floor-fractured craters found on the Moon and Mars. Our simulations furthermore do not
435 account for the effects of magma viscosity and cooling on flow dynamics and geometric
436 development of laccoliths found previously (Burchardt et al., 2019; Mattsson et al., 2018;
437 Michaut, 2011; Thorey & Michaut, 2016). To investigate such effects on the dynamics of host
438 rock fracturing and displacements will require a different approach for simulating magma flow
439 and laccolith growth that goes beyond this work and that of (Morand et al., 2024). Constraints
440 on the porosity, and thus strength, of the lunar and martian crust, remain a source of uncertainty,
441 however, because they are either theoretical, or based on orbital observations, modeling of

442 limited planetary seismic data, or laboratory experiments on terrestrial rocks (Heap et al., 2017;
443 Q. Huang & Wieczorek, 2012; Knapmeyer-Endrun et al., 2021; Wieczorek et al., 2022). It remains
444 necessary, therefore, as we have done here, to include a wide range of rock toughnesses and
445 stiffnesses, in numerical simulations of tectonic and magma-induced deformation of the crust of
446 Mars, the Moon and other rocky planetary bodies.

447
448 Finally, our 2D DEM models are the first to directly simulate dynamic fracturing and the
449 concentration of high strains induced by an inflating laccolith under gravitational forces on
450 planetary bodies of different mass. The results only do so along a vertical 2D profile. Three-
451 dimensional DEM simulations of caldera and sinkhole collapse have shown how asymmetrical
452 and anisotropic stresses and fracture patterns can develop (Hardy, 2021; Wang et al., 2022). This
453 limits the interpretation of fracture patterns and displacement patterns observed from orbit on
454 the Moon and Mars. Future 3D DEM modeling will be required to better understand the
455 development of asymmetric surface doming (Lena et al., 2013), hierarchical crack patterns at
456 floor-fractured craters (Montigny et al., 2022), and the accuracy of using elastic rheology theory
457 to infer intrusion depth from observed surface fracture patterns (Walwer et al., 2021).

458 **5 Conclusions**

459 We ran unprecedented 2D DEM simulations of displacement, straining and fracturing of
460 planetary crust of a range of stiffness and toughness, by inflating a laccolith at one kilometer
461 depth. Our simulations show that the different gravitational acceleration on the Moon, Mars,
462 and Earth influences the magma-induced displacement patterns, as well as the concentration of
463 shear strain and cracking in the subsurface. A same amount of magma emplaced at the same
464 depth in the lunar or martian crust must have induced higher surface displacement compared to
465 what is expected on Earth. This influence of the gravitational acceleration impacts estimates of
466 volumes of magma intruded into the shallowest few kilometers of crust on the Moon and Mars.
467 The amount of laccolith-induced cracking is, however, more dominantly controlled by the
468 stiffness and toughness of the planetary crust, as found previously for Earth (Morand et al., 2024).
469 Our model results show that inferences on magma intrusion properties in the often already
470 fractured and heterogeneous crusts of planetary bodies smaller than Earth, based on existing
471 models that assume simplified crustal rheologies, should be assessed with caution.

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485

486 **Open Research**

487 The authors declare no conflicts of interest relevant to this study. The simulations were produced under an
488 academic license in the commercial software PFC2D (Itasca Ltd.) of which the code cannot be shared publicly.
489 Python scripts used to plot the figures and process the simulation output, as well as the simulation output itself
490 such as particles' positions, displacements, stresses, finite shear strain, radius and group of particles, and
491 cracks' positions, orientations, length and aperture, will be made available in ASCII (.txt) format in a Zenodo
492 data set (Poppe et al., in prep.).

493

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Supplementary Table S1. Varied computational parameters and corresponding rock properties in the 2D DEM model.

| | |
|---|---|
| <i>Model geometry</i> | |
| Model width | 20,000 m |
| Depth of laccolith top | 1,000 m |
| Laccolith width | 1,000 m |
| Laccolith height | 50 m |
| Particle radii | 6.65 ± 1.65 m |
| Number of particles | ~120 × 10 ³ |
| <i>PFC2D7 contact model</i> | |
| Rock-rock | Soft-bond |
| Rock-wall | Soft-Bond |
| Magma-magma | Linear parallel bond |
| Magma-wall | Linear |
| Magma-rock | Linear parallel bond |
| Along broken bonds (cracks) | Rolling resistance linear |
| <i>Bond parameters that control stiffness</i> | |
| Effective modulus (E*) | 2.5 × 10 ⁹ (Pa) ^α |
| Ratio between normal and shear stiffness | 2.5 |
| <i>Bond parameters that control toughness</i> | |
| Bond tensile strength (ten) | 0.0 (Pa) ^α |
| Bond cohesion (coh) | 10 × bond tensile strength (Pa) |
| <i>Other contact parameters</i> | |
| Friction angle | Rock-rock/rock-wall: 30.0 (°) Rock-magma: 26.6 (°) Other bonds: 0.0 (°) |
| Friction coefficient between unbonded particles | Magma-magma/magma-wall: 0.0 Other bonds: 0.5 |
| Softening factor (γ) | Rock-rock/rock-wall: 13 |
| Softening tensile strength factor (ζ) | Rock-rock/rock-wall: 0.4 |
| Radius multiplier | 1.0 |
| Gap between bonded particles | 6.65 × 10 ⁻³ (m) |

756 ^α for rock-rock bonds and rock-wall bonds, see Table 1 in the main text.

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