

1 **Considering fault interaction in estimates of absolute stress along**
2 **faults in the San Gorgonio Pass region, southern California**

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7
8 **ABSTRACT**

9 Present-day shear tractions along faults of the San Gorgonio Pass region can be estimated
10 from stressing rates provided by three-dimensional forward crustal deformation models.
11 Modeled dextral shear stressing rates on the San Andreas and San Jacinto faults differ from rates
12 resolved from the regional loading due to fault interaction. In particular, fault patches with
13 similar orientations and depths on the two faults show different stressing rates. We estimate the
14 present-day, evolved fault tractions along faults of the San Gorgonio Pass region using the time
15 since last earthquake, fault stressing rates (which account for fault interaction), and co-seismic
16 models of the impact of recent nearby earthquakes. The evolved tractions differ significantly
17 from the resolved regional tractions, with the largest dextral traction located within the
18 restraining bend comprising the pass, which has not had recent earthquakes, rather than outside
19 of the bend, which is more preferentially oriented under tectonic loading. Evolved fault tractions
20 can provide more accurate initial conditions for dynamic rupture models within regions of
21 complex fault geometry, such as the San Gorgonio Pass region. An analysis of the time needed to
22 accumulate shear tractions that exceed typical earthquake stress drops shows that present-day
23 tractions already exceed 3 MPa along portions of the Banning, Garnet Hill, and Mission Creek
24 strands of the San Andreas fault. This result highlights areas that may be near failure if
25 accumulated tractions equivalent to typical earthquake stress drops precipitate failure.

26

27 **Keywords:** San Andreas fault, San Gorgonio Pass, fault tractions, fault stressing rates, fault
28 interaction, seismic hazard

29

30 **1. INTRODUCTION**

31 The southern San Andreas fault system consists of multiple active faults that
32 accommodate the deformation between the North American and Pacific plates. Accurate
33 estimates of the earthquake hazard in California require an accurate assessment of the potential
34 for large through-going earthquakes and the ability for ruptures to propagate through fault
35 intersections and complexities (e.g., Field et al., 2013). One region of such complexity is the San
36 Gorgonio Pass region (SGPr), a restraining stepover along the southern San Andreas fault
37 (Figure 1). Accurate dynamic rupture models of the SGPr that simulate potential rupture paths
38 will help us assess the potential for large and damaging earthquakes through this region (e.g.,
39 Tarnowski, 2017; Douilly et al., 2017).

40 Dynamic rupture models show that, in general, the size and extent of earthquake ruptures
41 can depend highly on the initial conditions of the model (e.g., Oglesby et al., 2005). These
42 conditions include physical aspects, such as fault geometry and location of rupture nucleation
43 (e.g., Lozos et al., 2012; Lozos, 2016; Tarnowski, 2017), and time-dependent aspects, such as
44 state of stress and frictional parameters (e.g., Kame et al., 2003; Aochi and Olsen, 2004; Kase
45 and Day, 2006; Duan and Oglesby, 2007). Dynamic rupture models typically prescribe initial
46 shear and normal tractions by resolving the remote stress tensor, constrained from focal
47 mechanism inversions, onto individual fault elements (e.g., Kame et al., 2003; Oglesby et al.,
48 2003). This approach provides spatially variable ‘resolved’ tractions that capture the first-order
49 loading of the faults but does not take into account the loading history, nor the prior stress

50 interactions between faults. Not only can individual earthquake events change tractions along
51 nearby faults, advancing or retarding each faults' earthquake clock (e.g., King et al., 1994; Stein,
52 1999; Duan and Oglesby, 2005), but interaction among neighboring active faults influences their
53 long-term slip rates and stressing rates (Willemse and Pollard, 1998; Maerten et al., 1999;
54 Loveless and Meade, 2011). Stressing rates on any given fault can be estimated using geodesy
55 (e.g., Smith and Sandwell, 2006). However, the total accumulated traction along any given fault
56 segment depends on the accumulated tractions during the interseismic period as well as nearby
57 rupture history (e.g., Smith-Konter and Sandwell, 2009; Richards-Dinger and Dieterich, 2012;
58 Tong et al., 2014).

59 To account for loading history and fault interaction and produce more accurate estimates
60 of fault stress, we simulate deformation within the San Gorgonio Pass region using three-
61 dimensional forward models that provide both slip rates over multiple earthquake cycles and
62 stressing rates between earthquake events. Because we use slip rates over multiple earthquake
63 cycles to drive models that simulate interseismic deformation, the resulting shear stressing rates
64 incorporate the interactions between faults of the southern San Andreas fault system. The
65 interseismic shear stressing rates along with information about time since last earthquake event
66 can be used to estimate the shear traction on faults through the SGPr following the approach
67 employed by Tong et al. (2014). The resulting estimates of shear traction may differ from
68 resolving the remote stress tensor onto faults in that our models explicitly include fault
69 interaction and fault loading from depth during the interseismic period. Furthermore, we
70 incorporate the effects of recent earthquakes on faults near the SGPr to produce a more accurate
71 estimate of the current stress state of this system. Using tractions that incorporate fault
72 interaction and loading history may enhance the accuracy of dynamic rupture models, refining

73 our insight into the nature of potential earthquake rupture propagation within the San Gorgonio
74 Pass region.

75

76 **2. REGIONAL GEOLOGY – THE SAN GORGONIO PASS REGION**

77 Through the San Gorgonio Pass region (SGPr), deformation is partitioned onto multiple,
78 active and nonvertical fault strands (e.g., Matti et al., 1992; Figure 1). The San Bernardino strand
79 of the San Andreas fault lies at the northwest end of the San Gorgonio Pass. Two potential
80 rupture pathways go through the restraining bend connecting the San Bernardino strand to the
81 Coachella segment of the San Andreas fault. The southern pathway consists of the San Gorgonio
82 Pass thrust, Garnet Hill strand, and Banning strand of the San Andreas fault (Figure 1). The San
83 Gorgonio Pass thrust dips to the north and has a corrugated geometry near the Earth's surface
84 (e.g., Matti et al., 1992). The eastern end of the San Gorgonio Pass thrust connects to the Garnet
85 Hill and Banning strands. The north-dipping Garnet Hill and subparallel Banning strands have
86 approximately the same strike. The northern pathway through the SGPr consists of the Mill
87 Creek, Mission Creek, and Galena Peak strands of the San Andreas fault (Figure 1). Ongoing
88 debate centers on the geometry and activity of these fault strands through the northern part of the
89 SGPr (e.g., Kendrick et al., 2015; Beyer et al., 2018; Fosdick and Blisniuk, 2018). The San
90 Jacinto fault is sub-parallel to the San Andreas fault and extends to within 2 km of the San
91 Andreas fault at Cajon Pass. While the San Jacinto fault lies outside of the SGPr, it interacts with
92 the San Andreas fault and consequently impacts both long-term slip rates (e.g., Herbert et al.,
93 2014) and earthquake rupture paths (Lozos, 2016) on the San Andreas fault.

94 Recent paleoseismic data within the SGPr suggest that previous large through-going
95 earthquakes have a recurrence interval of ~ 1000 years, with the most recent earthquake rupture

96 through the San Gorgonio Pass along the southern pathway in 1400 AD (Heermance and Yule,
97 2017). Earthquakes along the San Bernardino strand (north of the restraining bend) and
98 Coachella segment of the San Andreas fault (south of the bend) occur more frequently, with
99 recurrence intervals of 200-300 years (e.g., Philibosian et al., 2011; Field et al., 2013; Onderdonk
100 et al., 2018). The difference in recurrence intervals outside of and inside of the restraining bed
101 suggests that previous earthquakes that have ruptured along the San Bernardino and Coachella
102 segments terminated at the restraining bend, which may be acting as an ‘earthquake gate’.
103 During the interseismic period since the last rupture event through the bend, shear tractions have
104 been accumulating along faults within the SGPr. Furthermore, recent earthquakes along faults
105 surrounding the SGPr could impact the state of stress within the San Gorgonio Pass, and thus the
106 shear and normal tractions along the faults.

107

108 **2.1 Recent earthquakes near the San Gorgonio Pass region**

109 To calculate the stress interaction effects from past earthquakes, we consider records of three
110 ground-rupturing earthquakes that occurred within the past 300 years near the SGPr. While many
111 smaller earthquakes have occurred within this region, these larger ground-rupturing events have
112 the greatest potential to impact tractions along nearby faults.

113 **2.1.1 1992 Landers Earthquake**

114 The Landers earthquake occurred on June 28, 1992, rupturing five fault segments,
115 striking northwest-southeast, in the Eastern California Shear Zone (Hart et al., 1993). The
116 interaction of these faults created a linked fault network that generated an M7.3 earthquake,
117 which is larger than expected for any single fault involved in the rupture (Aydin and Du, 1995).
118 The total rupture length is estimated at 85 km on the primary rupture trace (Sieh et al., 1993).

119 The epicenter was located on the south portion of the Johnson Valley Fault, and the rupture
120 traveled northward along the Landers-Kickapoo, Homestead Valley, Emerson, and Camp Rock
121 faults, crossing two extensional stepovers and one compressional stepover (e.g., Aydin and Du,
122 1995; Madden and Pollard, 2012), while only rupturing parts of the Johnson Valley, Emerson,
123 and Camp Rock faults (Sieh et al., 1993). All of the involved faults were previously mapped,
124 with the exception of the Landers-Kickapoo fault (Hart et al., 1993). The Johnson Valley and
125 Landers-Kickapoo faults each slipped locally more than 2 m and the central portion of the
126 Homestead Valley fault slipped more than 3 m (Sieh et al., 1993; Aydin and Du, 1995). Slip
127 exceeded 4 m on the Emerson fault, and a maximum dextral slip of approximately 6 meters
128 occurred on the North Emerson Fault (Bryant, 1992; Sieh et al., 1993; Bryant, 1994).

129

130 **2.1.2 1812 Wrightwood Earthquake**

131 The ~M7.5 earthquake that occurred on December 8, 1812 (here referred to as the
132 Wrightwood earthquake), is one of the earliest earthquakes documented in the historical records
133 of California; the rupture origin and extent are still uncertain. Evidence of this event has been
134 observed within several paleoseismic trench sites along the San Andreas fault north of Cajon
135 Pass (Weldon and Sieh, 1985; Seitz et al., 1997; Biasi et al., 2002; Fumal et al., 2002; Weldon et
136 al., 2002), with a maximum dextral slip of 4-6 m and possible northern rupture extent ~100 km
137 north of the Cajon Pass (Bemis et al., 2016). The southern extent of rupture is not well
138 constrained. The most recent event recorded at Plunge Creek on the San Bernardino strand (site 3
139 Fig.1) is dated to within the 1600s (McGill et al., 2002), but minor slip on secondary structures
140 farther south along the San Bernardino strand, near Burro Flats (site 4 Fig. 1), dates to the early
141 1800s (Yule and Howland, 2001). Several paleoseismic sites along the northernmost strand of

142 the San Jacinto fault record 1.8-3 m of slip during an early 1800s earthquake event (Kendrick
143 and Fumal, 2005; Onderdonk et al., 2013; Onderdonk et al., 2015). Several have suggested the
144 plausibility of the 1812 earthquake jumping the < 2 km extensional stepover between the San
145 Andreas and San Jacinto faults (Figure 1) and involving both faults (Onderdonk et al., 2013;
146 Onderdonk et al., 2015; Rockwell et al., 2015). Lozos (2016) used dynamic rupture models to
147 investigate rupture scenarios that best fit the paleoseismic evidence and historical accounts of the
148 Wrightwood earthquake. The models of Lozos (2016) suggest that the Wrightwood earthquake
149 nucleated near Mystic Lake on the San Jacinto fault (site 8 Fig. 1), produced a maximum of 6 m
150 of slip near Colton, and propagated north onto the San Andreas fault (maximum of 4-5 m of slip
151 between Cajon Pass and Wrightwood).

152

153 **2.1.3 1726 Coachella Valley Earthquake**

154 The Coachella segment of the southern San Andreas fault has not experienced a large
155 earthquake in historical time. Paleoseismic studies reveal that the most recent earthquake is dated
156 to 1726 ± 7 (Rockwell et al., 2018), with a possible rupture trace extending from Salt Creek site
157 along the Salton Sea (site 7 Fig. 1; Sieh and Williams, 1990) to the Thousand Palms oasis site on
158 the Mission Creek strand (site 5 Fig. 1; Fumal et al., 2002), with at least 2 m of dextral offset at
159 the Indio site on the Coachella segment (site 6 Fig. 1; Sieh, 1986).

160

161 **3. METHODS**

162 We use Poly3D, a quasi-static, three-dimensional boundary element method code, to
163 simulate loading and interseismic deformation along the southern San Andreas fault system.
164 Poly3D solves the relevant equations of continuum mechanics to calculate stresses and

165 displacements throughout the model (e.g., Thomas, 1993; Crider and Pollard, 1998). Faults are
166 discretized into triangular elements of constant slip (no opening/closing is permitted) within a
167 linear-elastic half-space. The element size along the faults of the San Gorgonio Pass region
168 (SGPr) average ~4 km and allow for our models to capture fault irregularities as small as ~10
169 km. We simulate the active fault geometry of the southern San Andreas fault, the San Jacinto
170 fault, and the Eastern California Shear Zone (Figure 2) based on the Southern California
171 Earthquake Center's Community Fault Model (CFM) version 4.0 (Plesch et al., 2007; Nicholson
172 et al., 2017). The CFM is compiled from geologic mapping, seismicity, and geophysical data.
173 While the CFM has been updated to version 5.2, the interpreted active fault geometry of the San
174 Gorgonio Pass region is still under debate (e.g., Kendrick et al., 2015; Beyer et al., 2018; Fosdick
175 and Blisniuk, 2018). We use version 4.0 of the CFM but include fault geometry modifications
176 that serve both to improve the representation of the mapped active fault geometry and to improve
177 the match of model and geologic uplift patterns and slip rates in the San Gorgonio Pass region
178 (e.g., Cooke and Dair, 2011; Herbert and Cooke, 2012; Fattaruso et al., 2014; Beyer et al., 2018).

179 Faults in the CFM are defined to the base of the seismogenic crust. To simulate long-term
180 and interseismic deformation, we extend the faults down to a freely slipping, horizontal basal
181 crack at 35 km depth that simulates distributed deformation below the seismogenic zone
182 (Marshall et al., 2009). This modification eliminates artifacts that develop when the long-term
183 slip rates go to zero at the base of the CFM-defined faults (Figure 2). Across many earthquake
184 cycles, deformation patterns are primarily controlled by fault geometry (e.g., Dawers and
185 Anders, 1995; Fay and Humphreys, 2005; Herbert and Cooke, 2012). Therefore, to capture the
186 first-order loading of active faults, we do not consider potential secondary impacts of
187 heterogeneous and/or anisotropic rock properties.

188 We prescribe the tectonic loading on the boundaries of the model base, far from
189 investigated faults. Following Beyer et al. (2018), we implement an iterative technique that
190 ensures a uniform tectonic velocity, determined from geodetic estimates (DeMets and Dixon,
191 1999) at the model edges that are sub-parallel to the plate boundary (sides labeled I on Figure 2)
192 and a linear velocity gradient at the models edges that cross the plate boundary (sides labeled II
193 on Figure 2). The iterative approach of Beyer et al. (2018) ensures that applied velocities are
194 within ~1% of the desired tectonic loading. For faults that extend beyond our model area (San
195 Andreas, San Jacinto, and Cucamonga-Sierra Madre fault systems), we apply slip rates to distal
196 edge patches of these faults to prevent non-zero slip rates on these faults at the edge of our
197 model. We apply 35 mm/yr dextral slip to the San Andreas fault at the northwestern edge of the
198 model (Weldon and Sieh, 1985). At the southeastern edge, we prescribe 25 mm/yr dextral slip to
199 the San Andreas fault and 10 mm/yr dextral slip to the San Jacinto fault (e.g., Sharp, 1981;
200 Becker et al., 2005; Fay and Humphreys, 2005; Meade and Hager, 2005). Because of complex
201 fault geometry and interaction among faults, deformation within the SGPr is not impacted
202 significantly by variations in the partitioning of slip rates between the San Andreas and San
203 Jacinto faults at this model edge (Fattaruso et al., 2014). We apply 1.6 mm/yr reverse slip
204 (McPhillips and Scharer, 2018) to the western edge of the modeled Cucamonga fault to account
205 for deformation along the Sierra Madre fault not included in our model.

206 We use a two-step modeling approach to estimate the interseismic stressing rates along
207 the southern San Andreas fault system. The first model simulates deformation over many
208 earthquake cycles (steady state model) providing slip-rate information to a second model
209 (interseismic model) that simulates the build-up of stress between earthquakes due to constant
210 slip below the locking depth. In the steady state model, tectonic loading is prescribed along the

211 model edges at the base of the model, far from the investigated faults. The faults throughout the
212 model have zero shear traction and slip freely in response to tectonic loading and fault
213 interaction. This zero-shear traction simulates the low dynamic strength of faults during rupture
214 (e.g., Di Toro et al., 2006; Goldsby and Tullis, 2011). We simulate interseismic deformation by
215 applying the distribution of slip rates determined with the steady state model to fault surfaces
216 below the prescribed locking depth and lock fault elements above the locking depth. The abrupt
217 transition from locked to slipping at the specified locking depth used here produces stresses that
218 are unreliable within one element of the transition, or ~ 5 km. We use a locking depth of 25 km
219 to ensure that our model results provide reliable fault tractions to about 20 km depth that can be
220 used within dynamic rupture simulations within the full depth of the seismogenic crust.

221

222 **3.1 Estimating the impact from nearby recent earthquakes**

223 We simulate the 1992 Landers earthquake, 1812 Wrightwood earthquake, and 1726
224 Coachella Valley earthquakes by prescribing the interpreted co-seismic slip distribution
225 associated with each earthquake (e.g., Sieh, 1986; Hart et al., 1993; Onderdonk et al., 2015) to
226 the modeled fault surfaces. We segment the rupture surface into multiple vertical segments and
227 prescribe each segment a uniform slip according to the observations at the rupture trace (Figure
228 4). All other faults in the model are locked, and we do not consider the effect of tectonic loading
229 while simulating each earthquake due to the short rupture time. The resulting static stress
230 changes due to each earthquake alters tractions along the faults within the SGPr.

231

232 **3.2 Estimating evolved tractions**

233 The interseismic model determines stressing rates due to deep movement below the

234 seismogenic crust and uses these stressing rates to calculate current shear tractions along the fault
235 segments of the southern San Andreas fault within the SGPr using the time since the last rupture.
236 Estimating both shear and normal tractions from stressing rates requires an assumption on how
237 such accumulated tractions may dissipate with time. This approach relies on the premise that
238 shear tractions that accumulate during the interseismic period are released during earthquake
239 events. Because normal tractions that accumulate in the interseismic period, such as within
240 restraining bends, are not necessarily relieved upon fault slip, models of earthquake cycles
241 require dissipation mechanisms in order to avoid singular-valued normal tractions. Duan and
242 Oglesby (2006) simulate multiple earthquake cycles by coupling a viscoelastic interseismic
243 model with an elastic dynamic rupture model, such that normal stresses are relaxed during the
244 interseismic period in the viscoelastic model and used as input to the dynamic rupture model.
245 Alternatively, the Rate and State Earthquake Simulator (RSQSIM) employs a constant, but
246 spatially variable normal stress distribution and disregards accumulated normal tractions (e.g.,
247 Richards-Dinger and Dieterich, 2012). Here, we follow the approach of RSQSIM and do not
248 carry the normal stressing rates through the rest of the analysis.

249 To estimate the shear traction that evolves over the earthquake cycle, henceforth called
250 the evolved shear traction, we follow Tong et al. (2014) and use the stressing rate information
251 from the interseismic model and the time since last event for each fault. In this approach, we
252 only consider large ground-rupturing events that are preserved in the paleoseismic record. The
253 approach analyzes a coseismic stress drop corresponding to the change from static friction to
254 dynamic friction during large ground-rupturing earthquakes (shaded region in Figure 3). If the
255 dynamic strength of the fault is near zero, a complete stress drop is associated with these events.
256 Such complete stress drop is consistent with recent field measurements of low temperatures

257 along recently ruptured fault surfaces, a result of a very low dynamic friction (e.g., Carpenter et
258 al., 2012, Fulton et al., 2013, Li et al., 2015), as well as high-speed laboratory frictional
259 experiments cited above. Consequently, the associated shear traction at any time in the
260 earthquake cycle is

$$261 \quad \tau = \dot{\tau} \cdot t \quad (1)$$

262 where τ is the evolved shear traction, $\dot{\tau}$ is the shear stressing rate and t is the time since last
263 event. We sum the evolved shear tractions calculated by Equation 1 and the static stress changes
264 due to nearby earthquakes to produce the present-day evolved shear tractions along the faults
265 within the SGPr. This simplified approach to estimate the distribution of present-day shear
266 tractions may provide more accurate initial conditions than the approach employed by dynamic
267 rupture models of estimating tractions by resolving the remote loading onto the faults because
268 the evolved tractions also incorporates both loading history and fault interaction, by including
269 long term slip rates at depth and recent nearby earthquake events (Figure 3).

270

271 **3.3 Consideration of geometric and tectonic uncertainty**

272 Using models of deformation over multiple earthquake cycles, Beyer et al. (2018)
273 compared the slip rate distribution from six plausible active fault configuration models to
274 available geologic slip rate data. The analysis revealed that two active fault configurations
275 provide the best fit to the geologic observations. For this study, we use both of the two best-fit
276 models: the *Inactive Mill Creek* and *West Mill Creek* models from Beyer et al. (2018). The most
277 pronounced fault geometry difference between the two configurations is the addition of a
278 through-going Mission/Mill Creek strand through the northern part of the SGPr in the *West Mill*
279 *Creek* model. Here, we present the results of the *Inactive Mill Creek* model geometry, and the

280 Supplemental Material contains the results of the *West Mill Creek* model geometry. Following
281 Herbert and Cooke (2012), Beyer et al. (2018) also tested each of the plausible fault
282 configurations under a range of reasonable tectonic loading (45-50 mm/yr at 340°-345°; DeMets
283 and Dixon, 1999); for this study, we use the mean slip rate from the end members of permissible
284 tectonic loadings.

285

286 **4. RESULTS**

287 We present the interseismic stressing rates for faults of the San Gorgonio Pass region
288 (SGPr) and show the impact of fault interaction on these rates. We analyze the results of the
289 models that simulate three recent ground-rupturing earthquakes and the impact of these
290 earthquakes on fault tractions within the SGPr. We then calculate the total evolved shear
291 tractions that incorporate the impacts of both fault interaction and loading history.

292

293 **4.1 Stressing rates**

294 Maps of interseismic shear stressing rate along the southern San Andreas fault reveal how
295 the fault geometry controls the stressing rate distribution (Figure 5). Figure 5 shows stressing
296 rates for the *Inactive Mill Creek* model configuration (Figure 5A and 5C) and the difference in
297 stressing rates between the two plausible active fault geometries (Figure 5B and 5D). Dextral
298 shear stressing rates are larger (maximum 12 kPa/yr) than the reverse-shear stressing rates
299 (maximum ~3 kPa/yr) along the San Andreas fault. Furthermore, portions of the faults parallel to
300 the overall plate motion, outside of the restraining bend, have greater dextral stressing rate than
301 faults within the bend. Dextral shear stressing rates are largest along the San Bernardino and
302 Mission Creek strands of the SAF and decrease within the restraining bend of the SGPr (Figure

303 5A). The San Gorgonio Pass thrust has an undulating strike and small sinistral shear stressing
304 rates occur locally along patches of the western San Gorgonio Pass thrust where the strike is less
305 than $\sim 265^\circ$. The reverse-shear stressing rates are near zero outside the restraining bend and
306 increase within the bend along north-dipping fault strands that strike obliquely to the plate
307 motion and accommodate uplift (Figure 5C). Stressing rates increase with depth, consistent with
308 the deep slip that is applied to faults in the interseismic model. The difference in stressing rates
309 between the two best-fitting model geometries (Figure 5B and 5D) are lower than 1 kPa/yr and
310 indicate that the *West Mill Creek* fault geometry produces higher dextral stressing rates (blue)
311 throughout most of the region. The greatest difference in reverse-shear stressing rates is limited
312 to within and just outside the bend. We only consider the *Inactive Mill Creek* fault geometry for
313 the rest of our analysis, and consequently, reported shear tractions may underestimate by $\sim 2\%$
314 shear tractions if the true active fault geometry is closer to the configuration of our *West Mill*
315 *Creek* model.

316 To the first order, the strike-parallel shear stressing rate along the southern San Andreas
317 fault correlates with the orientation of the fault segments relative to the applied model loading
318 that simulates plate motions. Previous models of the region have shown significant interaction
319 between the San Andreas and San Jacinto faults (Herbert et al., 2014; Fattaruso et al, 2014), so
320 we expand the analysis of stressing rates to include both of these faults in order to investigate the
321 influence of fault interaction on stressing rates. The model produces different interseismic
322 dextral shear stressing rates along similarly oriented portions of the San Andreas and San Jacinto
323 faults. Figure 6 shows a gridded surface fit through model data points (white circles) of dextral
324 stressing rates for different strikes and depths of the San Andreas (Figure 6A) and San Jacinto
325 faults (Figure 6B). Stressing rates for both faults increase with depth, and in general, dextral

326 stressing rates are higher on the San Andreas fault than on the San Jacinto fault for locations with
327 the same strike and depth. For the San Andreas fault, maximum dextral stressing rate occurs at
328 strikes between 300° - 305° . Relative to the San Andreas, the variation of dextral stressing rate
329 with strike along the San Jacinto fault is more subdued, but the distribution shows a maximum
330 strike stressing rate along segments that strike $\sim 310^{\circ}$. Both of these maximum shear orientations
331 differ from the orientation expected from resolving the regional stress tensor (black line in Figure
332 6). The difference between dextral stressing rates along the San Andreas and San Jacinto faults
333 and the expected distribution from resolved tractions demonstrates the strong impact of fault
334 interaction on the distribution of fault stressing rates.

335 The impact of fault interaction is also demonstrated in the relative stressing rates on the
336 Banning and Mission Creek strands, which differs between the two plausible fault configurations
337 (Figure 5). The presence of a through-going Mission Creek fault in the *West Mill Creek* model
338 geometry (Figure S1) shifts strike-parallel shear stressing rates from the Banning strand to the
339 Mission Creek strand by ~ 0.1 kPa/yr. These differences in stress accumulation rates over the
340 interseismic period demonstrates that fault interaction impacts the distribution of accumulated
341 tractions along faults within a complex system.

342

343 **4.2 Impact of stresses from regional earthquakes**

344 To assess the impact of the recent nearby earthquakes along faults within the SGPr, we
345 numerically simulate three ground-rupturing earthquakes. We simulate the Landers Earthquake,
346 Wrightwood Earthquake, and Coachella Valley Earthquake and examine the static stress change
347 due to each event (Figure 7). Because we do not consider the potential relaxation of these crustal
348 stresses over time (e.g., Pollitz and Sacks, 2002), the fault tractions from earthquakes modeled

349 provide an upper bound to expected tractions.

350 The modeled static stress change from the 1992 Landers earthquake impacts tractions
351 along faults within the San Gorgonio Pass restraining bend. The change in dextral tractions
352 (positive) reach a maximum of ~ 0.1 MPa, along the San Gorgonio Pass thrust and a change in
353 sinistral tractions (negative) of up to 0.15 MPa on the southern Garnet Hill, Banning, and
354 Mission Creek strands of the SAF (Figure 7A). The San Gorgonio Pass region lies in the
355 extensional quadrant of the Landers rupture, and as a result, the fault strands within the bend are
356 loaded with normal dip-slip tractions of ~ 0.13 MPa. This dip-slip traction change effectively
357 reduces the accumulated long-term reverse dip-slip traction on these faults.

358 Change in co-seismic dextral tractions due to the 1812 Wrightwood earthquake increase
359 the most (1.3 MPa) just south of the rupture limit on the San Bernardino strand, while the
360 southernmost portion of the San Bernardino strand experiences sinistral traction changes of \sim
361 0.25 MPa (Figure 7B). The western San Gorgonio Pass thrust is loaded with dextral tractions (\sim
362 0.4 MPa), while the Garnet Hill, Banning and Mission Creek strands experience slight (< 0.1
363 MPa) increases in sinistral shear tractions. Furthermore, the western-most extent of the San
364 Gorgonio Pass thrust has normal dip-slip shear, while the rest of the thrust and Garnet Hill strand
365 has reverse dip-slip shear. These complex fault stressing patterns result from the location and
366 orientation of the faults in relation to the Wrightwood rupture path. The close proximity of the
367 dextral slip on the SJF to the subparallel fault strands of the SAF results in sinistral co-seismic
368 traction changes on the SAF, which sits in the stress shadow of the Wrightwood earthquake.

369 Traction changes imposed on the San Gorgonio Pass fault strands due to the 1726
370 Coachella Valley earthquake reach ~ 1 MPa on the Mission Creek ahead of the rupture
371 termination and ~ 0.8 on the Banning strands near the junction with the Coachella segment

372 (Figure 7C). Dextral tractions of up to ~ 0.1 MPa extend into the restraining bend. While these
 373 nearby earthquakes are shown here to impact the SGPr, all the resulting static stress changes due
 374 to these earthquakes are small compared to the total tractions accumulated along these faults
 375 during the interseismic period.

376

377 **4.3 Estimate evolved stresses**

378 Paleoseismic data provide estimates for the time since last event, t , along active faults
 379 (e.g., Biasi et al., 2009; Table 1). We estimate the total present-day traction along each fault
 380 segment by summing the tractions from Equation 1 with the static traction change of each nearby
 381 earthquake. For the San Bernardino segment, we use time since last event from the compiled
 382 earthquake data of Biasi et al. (2009). Paleoseismic sites at Pitman Canyon (site 2 Fig. 1), Plunge
 383 Creek (site 3 Fig. 1), and Wrightwood (site 1 Fig. 1) provide a mean t of 207 years for the San
 384 Bernardino segment. Paleoseismic constraints from the Thousand Palms Oasis site (site 5 Fig. 1;
 385 Fumal et al., 2002) is used for the Mission Creek strand and the Coachella site (site 6 Fig. 1;
 386 Philibosian et al., 2011) for the Coachella segment. These studies are in agreement that the last
 387 rupture event occurred circa 1680. This event has been re-dated by Rockwell et al. (2018) to be
 388 around the year 1726 ± 7 . For the San Gorgonio Pass thrust, Banning, and Garnet Hill strands of
 389 the SAF, we use an earthquake rupture year of 1400 (Heermance et al., 2017; Yule et al., 2014).

390

Fault strand(s)	Paleoseismic Site(s)	Most recent EQ Year (AD)	Time since last event (yr)
San Bernardino	Pitman Canyon/Plunge Creek/Wrightwood (<i>Biasi et al., 2009</i>)	1812	207
Banning/SGPT/Garnet Hill	Millard Canyon (<i>Heermance et al., 2017, Yule et al., 2014</i>)	1400	619
Mission Creek/Coachella	1000 Palms/Coachella (<i>Fumal et al., 2002; Philibosian et al., 2011;</i>	1726	293

	Rockwell et al., 2018)		
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391 Table 1: Time since last event data used to calculate absolute shear stress from the
392 interseismic stressing rate.
393

394 Due to the variable time since last earthquake event across faults of the SGPr, the evolved
395 shear traction distribution along the fault surfaces (Figure 8A) differs significantly from the shear
396 stressing rate distributions (Figure 5). Whereas dextral-shear stressing rates are lower along the
397 north-dipping fault surfaces within the SGPr restraining bend than on fault surfaces outside of
398 the bend (Figure 5A), the longer t for the faults within the restraining bend increases the total
399 accumulated dextral shear traction within the bend relative to other faults (Figure 8A). Similarly,
400 although the Coachella segment and the San Bernardino strand of the SAF have greater dextral
401 stressing rates than the restraining segment, the more recent rupture of these segments in the
402 1726 and 1812 events, respectively, reduces the accumulated tractions outside the bend. The
403 largest evolved dextral shear tractions arise along the Banning and Garnet Hill strands of the
404 SAF near the juncture with the Coachella segment of the SAF (Figure 8A). Regions of high
405 dextral shear traction also arise along portions of the San Gorgonio Pass thrust. The evolved
406 reverse shear tractions are greatest along the San Gorgonio Pass thrust within the restraining
407 bend. We note that if the true active fault geometry is closer approximated by the alternative
408 fault configuration (Figure S1), the total evolved shear tractions may be underestimated by up to
409 2%.

410 These evolved shear tractions take into account fault interaction (Figure 6), the t for each
411 fault strand (Table 1), and the impact of recent nearby earthquakes (Figure 7). To assess the
412 impact of fault history and interaction, we compare the evolved dextral shear traction (Figure
413 8A) to the fault tractions that result from resolving the regional stress tensor constrained from
414 focal mechanism inversions onto the faults (Figure 8B). Following Tarnowski (2017), we use the

415 orientation of the stress field and relative magnitude of the principal stress axes from Hardebeck
416 and Hauksson (2001) and the stress ratio, $A\phi$ (Simpson, 1997), of 1.5, which indicates a mixed
417 strike-slip and thrust stress regime. The resulting stress tensor is scaled such that the change from
418 static to dynamic friction results in a 3 MPa stress drop (e.g., Tarnowski, 2017) in order to
419 represent the fault loading conditions preceding a large earthquake rupture. The larger magnitude
420 of the resolved dextral shear tractions compared to the evolved tractions is due to the scaling of
421 the regional stress to produce failure and a 3 MPa stress drop with a dynamic friction of 0.1. The
422 evolved tractions to the current year (2019) are not explicitly at failure and these tractions
423 exclude those required for dynamic sliding (non-shaded region of Figure 3). Consequently, we
424 focus our comparison on the patterns of the evolved and resolved stresses rather than the
425 absolute level of stress. The resolved tractions have greater lateral heterogeneity as the tractions
426 range from 10 MPa dextral to -5 MPa sinistral shear traction, where portions of the San
427 Gorgonio Pass thrust receives sinistral shear from resolved loading. In contrast, the evolved
428 tractions show dextral shear tractions everywhere on faults of the SGPr. Whereas the evolved
429 shear tractions increase with depth, the remote stress tensor is not resolved for different depths.

430

431 **5. DISCUSSION**

432 Here we discuss the impact of including fault interaction and the effects of recent nearby
433 earthquakes on fault tractions, and the implications of this study's findings for seismic hazard
434 assessment.

435

436 **5.1 Resolved tractions likely oversimplify initial conditions for rupture**

437 Rupture propagation within dynamic models highly depends on the initial conditions

438 used for the model (e.g., Kame et al., 2003; Duan and Oglesby, 2007; Lozos et al., 2012). These
439 models have the power to simulate potential rupture size and extent, as well as potential rupture
440 paths (e.g., Oglesby et al., 2003). Most rupture dynamic studies either use a homogeneous
441 regional stress field (e.g., Lozos et al., 2012) or a regional stress field with spatially rotating
442 principal stress (e.g., Aochi and Fukuyama, 2002) to estimate the initial shear tractions on fault
443 segments. Resolving the regional stress tensor onto faults does not account for fault interaction
444 or the rupture history of each fault. Whereas these effects may be minimal in regions with planar
445 faults, we show here that within regions of fault complexity, fault interaction and loading history
446 can advance, or retard, the fault towards failure. Prescribing tractions that incorporate the effects
447 of fault interaction and the loading history of each fault may improve the accuracy of dynamic
448 rupture models.

449 Due to fault interactions over multiple earthquake cycles and the variable time since last
450 earthquake event across faults of the SGPr, the evolved shear traction distribution along the fault
451 surfaces (Figure 8A) differs from the tractions resolved from the regional stress state (Figure
452 8B). Whereas resolved dextral-shear tractions are lower along the north-dipping fault surfaces
453 within the SGPr restraining bend than on faults outside of the bend (also seen in the stressing
454 rates in Figure 5), the longer t for the faults within the restraining bend increases the total dextral
455 shear traction compared to other faults (Figure 8A). Similarly, the more recent ruptures of the
456 1726 and 1812 earthquake events reduce the accumulated tractions outside the bend, but not
457 within. The largest dextral shear tractions arise along the Banning and Garnet Hill strands of the
458 SAF, especially near the juncture of the Banning strand with the Coachella segment of the SAF
459 (Figure 8A). Furthermore, the evolved stresses do not produce enigmatic left-lateral loading on
460 portions of the San Gorgonio Pass thrust of the resolved stresses. These structures do not show

461 surface evidence for left-lateral slip (Yule and Sieh, 2003). Consequently, the pattern of
462 consistent dextral shear traction produced by the evolved stresses that include fault interaction
463 and interseismic loading agrees with geologic evidence. Including evolved stress state for initial
464 conditions within dynamic rupture models can produce more accurate assessment of rupture
465 behavior along complex fault systems.

466

467 **5.2 Implications for seismic hazard**

468 To assess how close the faults of the San Gorgonio Pass region (SGPr) are to failure
469 within our model of evolved fault tractions, we analyze the time until failure for each fault
470 element. To consider this, we calculate how many years are required from the present day to
471 accumulate evolved net shear tractions equivalent to typical earthquake stress drops of 3 MPa
472 (Figure 9A) and 10 MPa (Figure 9B) (e.g., Allman and Shearer, 2009; Goebel et al., 2015).
473 Using this criterion, fault elements on the easternmost Banning, Garnet Hill, and Mission Creek
474 strands currently exceed 3 MPa at the base of the model (ellipse in Figure 9A). When using 10
475 MPa for the accumulated traction required to trigger the next earthquake, the first fault element
476 to fail is on the San Bernardino strand at 584 years from now (ellipse in Figure 9B). The
477 difference in vulnerable position within the fault system owes to the greater stressing rate along
478 the San Bernardino strand compared to the other faults.

479 The time since last ground rupturing earthquake event for fault strands within and just
480 outside the restraining bend are near or greater than the estimated recurrence interval for these
481 fault strands, and paleoseismic studies in this region suggest these faults are probably overdue, or
482 close to failure (Philibosian et al., 2011; Rockwell et al., 2018; Onderdonk et al., 2018).
483 Consequently, while 3 MPa represents a low stress drop in the SGPr (Goebel et al., 2015), this

484 lower stress drop implies that these faults are close to failure, which is consistent with earthquake
485 clock and recurrence intervals of these faults.

486

487 **6. CONCLUSIONS**

488 We use three-dimensional crustal deformation models to estimate present-day fault
489 tractions in the San Gorgonio Pass region. The models that estimate interseismic stressing rates
490 are loaded with deep slip rates determined from a multiple-earthquake-cycle model that
491 explicitly includes fault interaction. Consequently, the interseismic stresses incorporate both
492 regional tectonic loading and fault interaction. A gradient of increasing shear stressing rates with
493 depth emerges from our models that is consistent with deep interseismic deformation. To
494 investigate the role of fault interaction within our models, we compare our modeled dextral shear
495 stressing rates for the San Andreas and San Jacinto faults, which have similar orientation.
496 Subsequently, the interseismic stressing rates would be similar on the two faults if based solely
497 on orientation with respect to remote loading. Significant differences in the patterns of stressing
498 rates along patches of the San Andreas and San Jacinto faults with similar orientation and depth
499 arise due to the interaction between these two faults.

500 The total evolved present-day shear tractions along the fault include both the accumulated
501 stressing rates since the last earthquake event and the impact of nearby earthquakes. We simulate
502 recent nearby ground-rupturing earthquakes with co-seismic models to investigate the impact of
503 these rupture events on the stress state along the San Andreas fault within the San Gorgonio Pass
504 region. The pattern of total evolved fault tractions differs from that of the interseismic stressing
505 rates. Tractions are higher within the restraining bend than outside the bend because of the longer
506 time since last event on these faults. Fault strands within the restraining bend have been loading

507 for twice as long as the Coachella segment to the south and three times as long as the San
508 Bernardino strand to the north. Comparison of our evolved tractions to the tractions resolved
509 from the local stress field shows distinct differences. While the linear gradient with depth
510 emerges from our models, the resolved tractions are not depth dependent, and the gradient must
511 be added. Because the evolved tractions account for loading history, the largest tractions occur
512 within the restraining bend, which has the longer time since last event.

513 We investigate the time needed for the accumulated net shear traction on each fault
514 element to exceed 3 MPa and 10 MPa, typical coseismic stress drop values. Because the
515 interseismic stress rates differ for faults throughout the San Gorgonio Pass region, the location
516 and timing of potential failure depends on the stress drop value used as the shear traction
517 threshold. Assuming a lower stress drop value shows that faults in the San Gorgonio Pass are
518 currently at failure, whereas higher stress drop values do not.

519 This approach provides a more heterogeneous, more accurate representation of the
520 current stress state along the southern San Andreas fault than a simple regional stress tensor. In
521 regions of complex fault geometry such as the San Gorgonio Pass region, an ‘earthquake gate’,
522 the potential for a through-going rupture is unclear and stress state may have a large control on
523 rupture behavior. Our evolved fault tractions can provide more realistic initial conditions for
524 dynamic rupture models of these regions, and therefore improve seismic hazard assessments.

525

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531

532

533 **FIGURE CAPTIONS**

534 **Figure 1.** Map of the San Gorgonio Pass region (SGPr). While there has not been a rupture event
535 in the SGPr since ~1400 AD, recent nearby earthquakes may impact the stress within the bend.
536 Fault traces of the San Andreas fault are labeled (GP – Galena Peak, SGPT – San Gorgonio Pass
537 Thrust). Rupture traces considered in this study are highlighted: Landers earthquake (dark red),
538 Wrightwood earthquake (red), Coachella Valley earthquake (orange), and 1400 event (yellow).
539 Paleoseismic sites are numbered as such: 1 – Wrightwood, 2 – Pitman Canyon, 3 – Plunge
540 Creek, 4 – Burro Flat, 5 – Thousand Palms, 6 – Indio, 7 – Salt Creek, 8 – Colton, and 9 – Mystic
541 Lake.

542
543 **Figure 2.** Northward oblique view of the model setup. Tectonic loading is prescribed at the
544 boundaries of the model base, such that sides labeled (I) have a uniform tectonic velocity parallel
545 to the plate boundary and sides labeled (II) are prescribed a linear gradient in the tectonic loading
546 across the plate boundary. The shear traction-free faults slip freely in response to the loading and
547 fault interaction. Black box outlines the San Gorgonio Pass region. SAF - San Andreas fault; SJF
548 - San Jacinto fault.

549
550 **Figure 3.** Schematic sketch of fault loading through time. During the interseismic period, the
551 earthquake clock of a fault can be advanced or retarded by earthquakes on other nearby faults. If
552 the dynamic strength of the fault is near zero, then the stress drop associated with the change
553 from static to dynamic friction is a complete stress drop.

554
555 **Figure 4.** Northward oblique view of the San Gorgonio Pass region showing the distribution of
556 the applied slip (in meters) associated with the nearby 1992 Landers (dark red), 1812
557 Wrightwood (red), and 1726 Coachella Valley (orange) earthquakes.

558
559 **Figure 5.** Modeled interseismic stressing rates along faults within the San Gorgonio Pass. This
560 region is primarily loaded in dextral shear (red in A). North-dipping faults are also loaded in
561 reverse dip slip (red in C). B and D show the difference in stressing rates between the two
562 plausible fault configurations of Beyer et al. (2018); difference is slip rate from *Inactive Mill*
563 *Creek* minus slip rate from *West Mill Creek* model configuration.

564
565 **Figure 6.** Modeled dextral stressing rates plotted with depth and fault strike for the A) San
566 Andreas fault and B) San Jacinto fault. Model results are plotted as white circles, and a best-fit
567 surface is fitted through the data (background grid). Patches along the two faults with similar
568 orientation and depth have different values of shear stressing rates due to fault interaction. The
569 orientations of magnitude of dextral shear stressing rate for both faults differ from the maximum
570 shear direction predicted from the regional stress tensor (vertical black lines).

571
572 **Figure 7.** Static stress changes from modeled recent, nearby earthquakes resolved as right-lateral
573 tractions along faults of the SGPr. The Landers earthquake (A) increased dextral shear tractions
574 along the San Bernardino strand and San Gorgonio Pass thrust, and decreased dextral shear
575 tractions along the Garnet Hill, Banning, and Mission Creek strands. The Wrightwood
576 earthquake (B) produces a complex change in tractions. The earthquake increased dextral shear
577 tractions just east of the rupture termination on the San Bernardino strand but decreased dextral
578 shear tractions further east due to the interaction with the neighboring San Jacinto fault, which
579 had dextral slip. The Coachella Valley earthquake (C) increased dextral shear tractions on the

580 easternmost Mission Creek and Banning strands.

581
582 **Figure 8.** Evolved (A) and Resolved (B) right-lateral tractions along faults of the SGPr. The
583 increasing shear traction with depth emerges from the evolved stresses due to deep slip in
584 the interseismic models. The resolved tractions show greater lateral variation than the evolved
585 tractions. Arrows indicate the direction of principle compression.

586
587 **Figure 9.** Faults of the SGPr colored by years until failure. (A) shows time until net shear
588 tractions since the last earthquake exceed 3 MPa. With this criterion, fault elements on the
589 Banning, Garnet Hill, and Mission Creek strands are currently at failure. (B) shows time until net
590 shear tractions since the last earthquake exceed 10 MPa. Under this assumption, the first fault
591 element to fail is on the San Bernardino strand at 584 years from present. Ellipses highlight the
592 first elements to fail.

593
594

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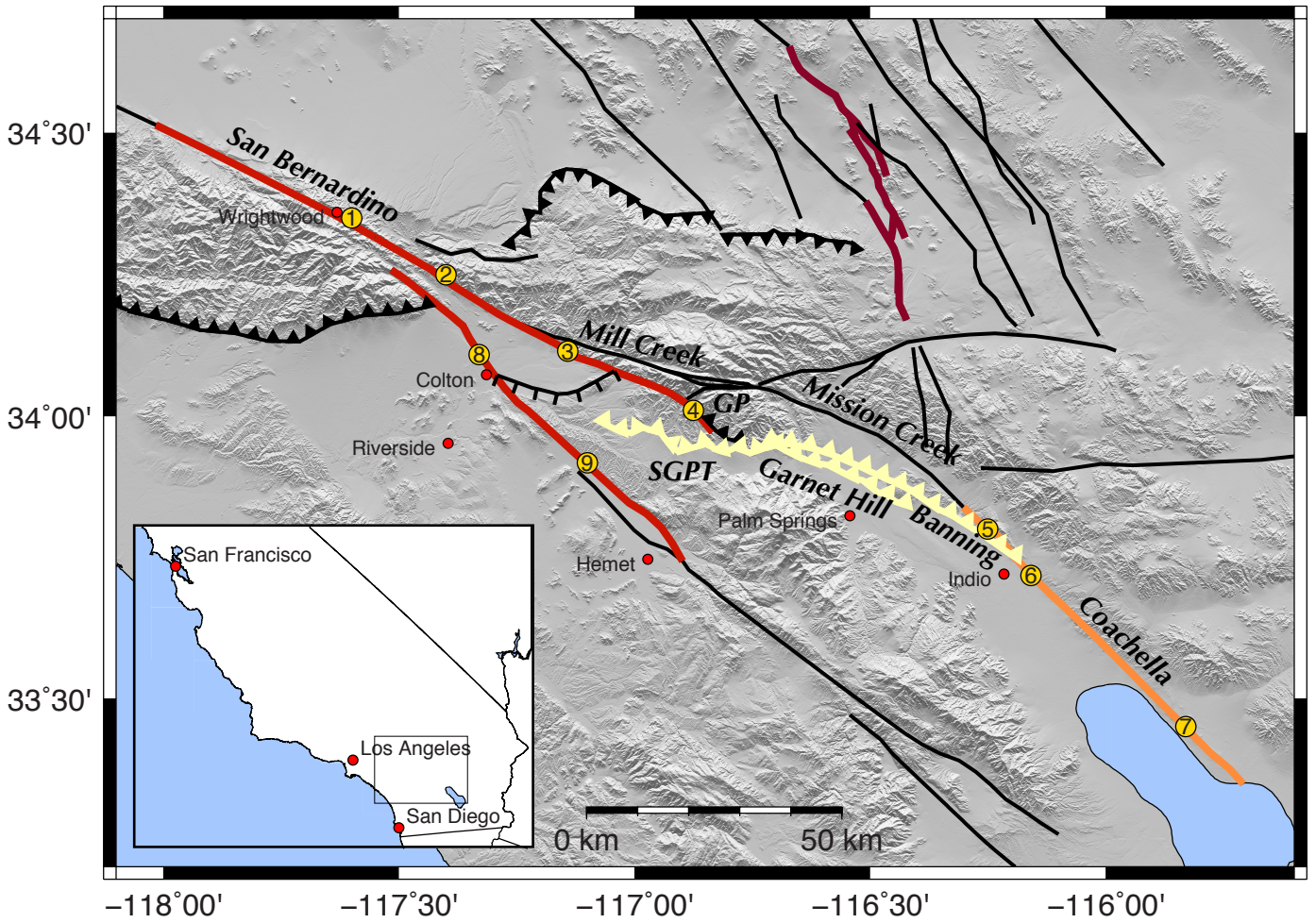
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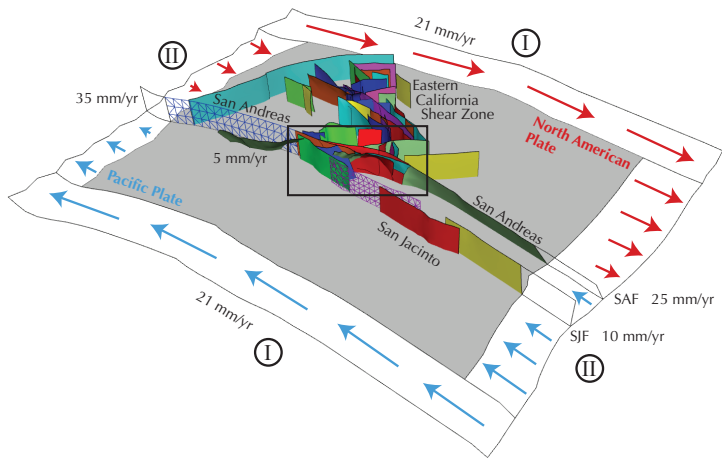
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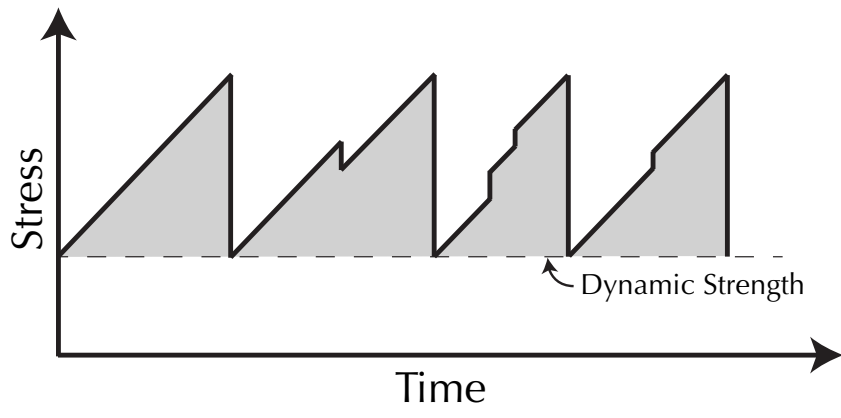
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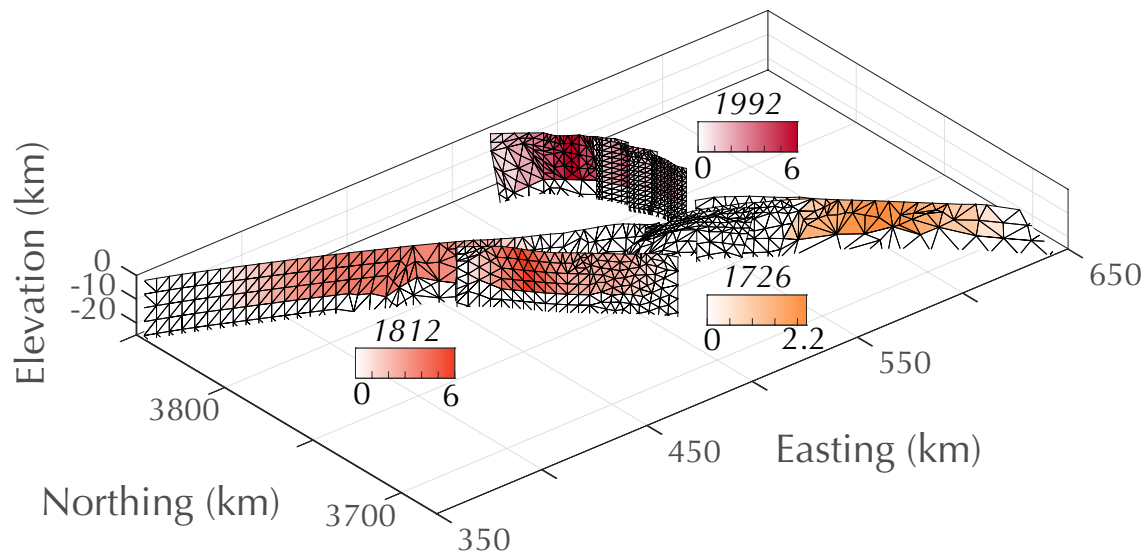
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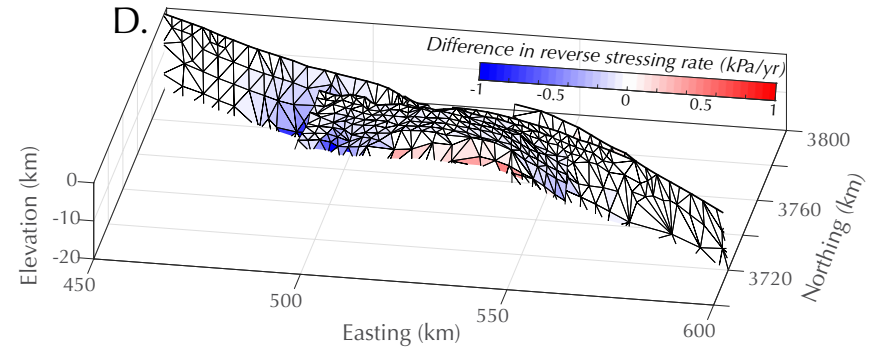
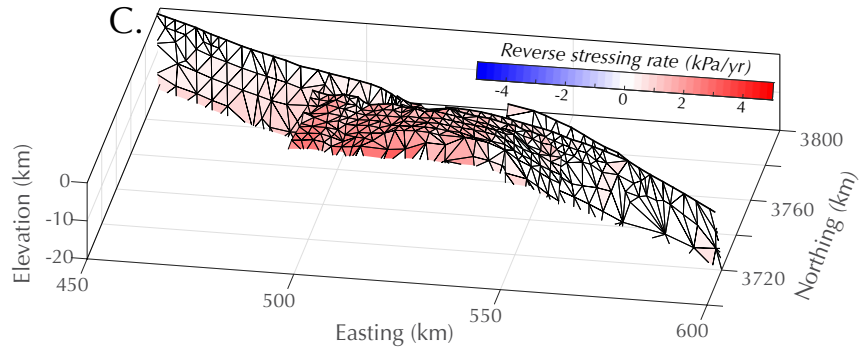
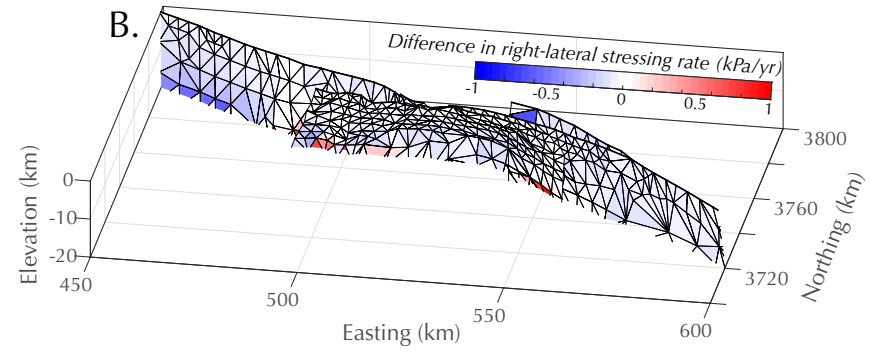
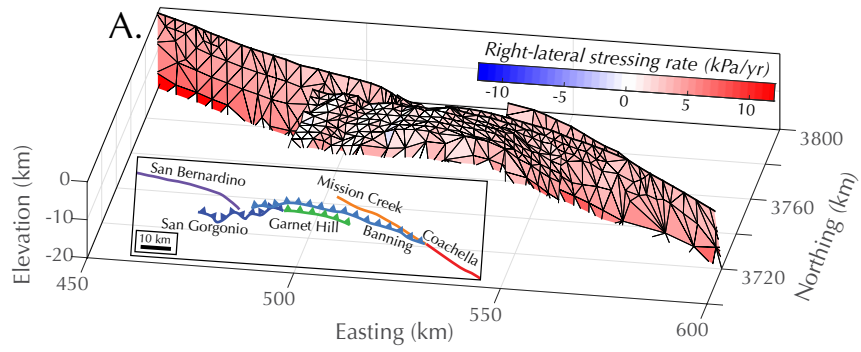
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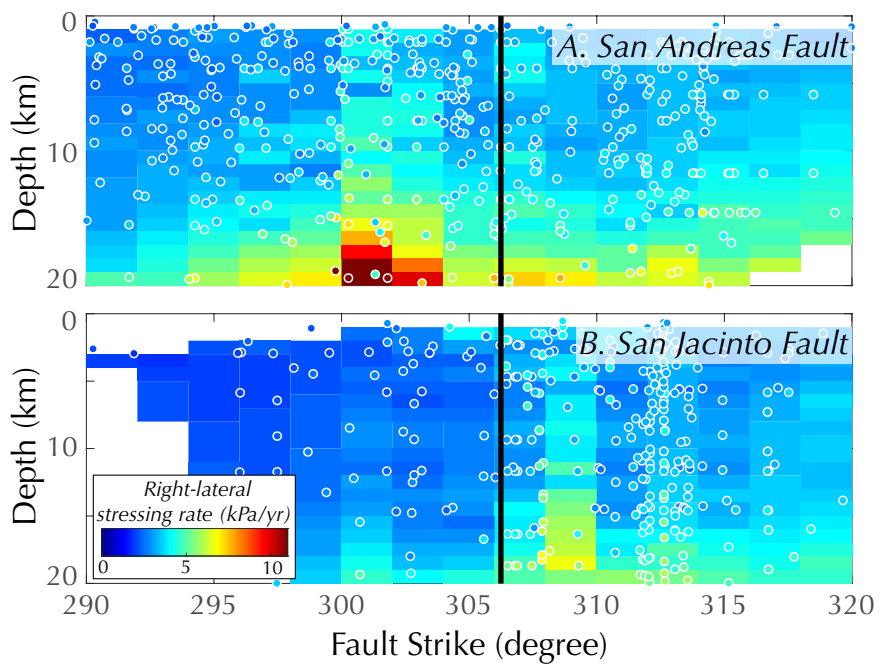


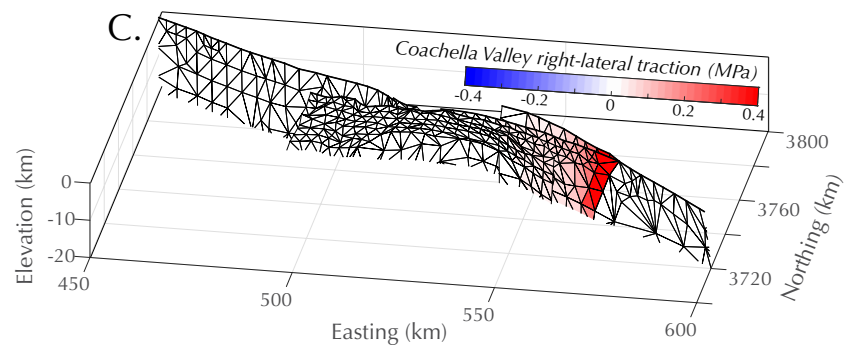
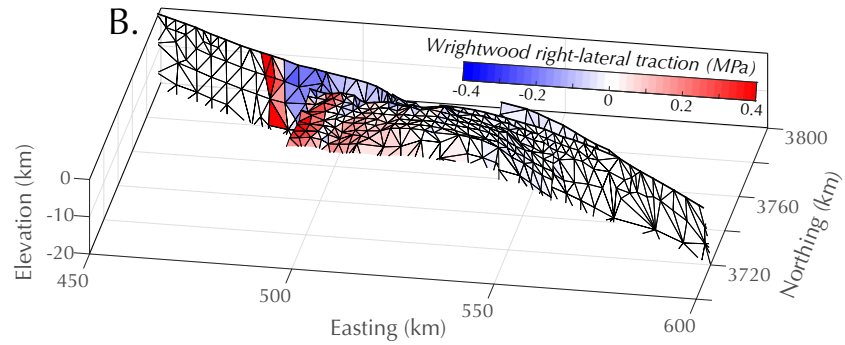
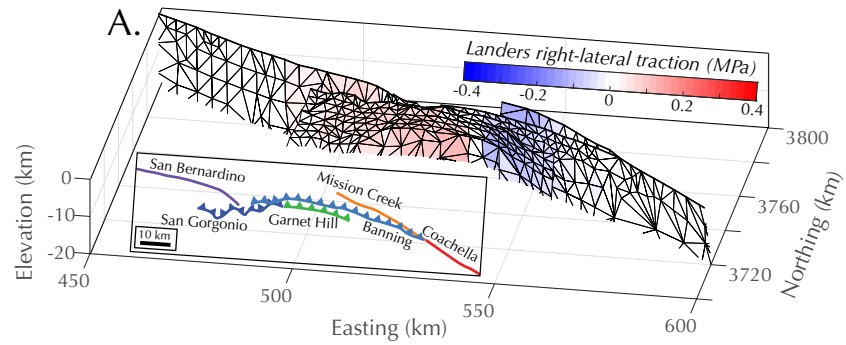


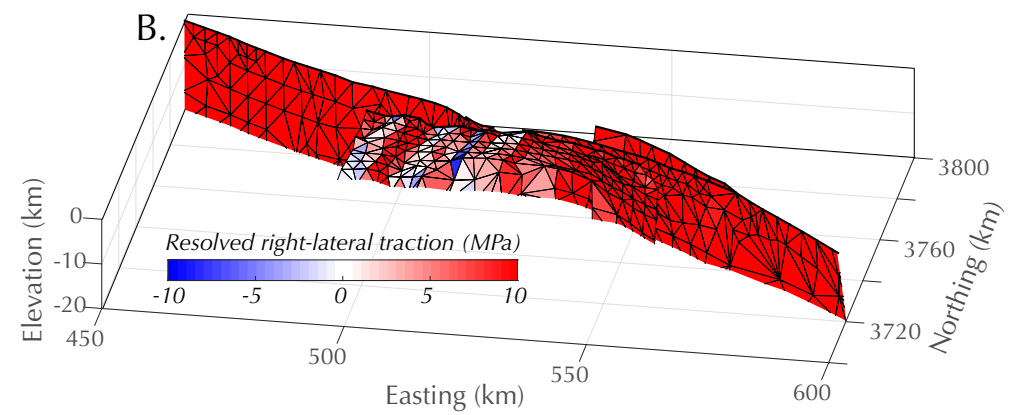
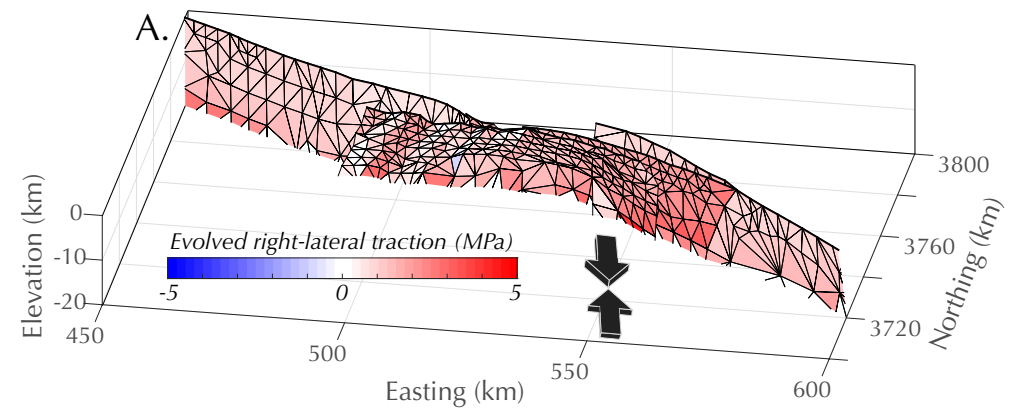


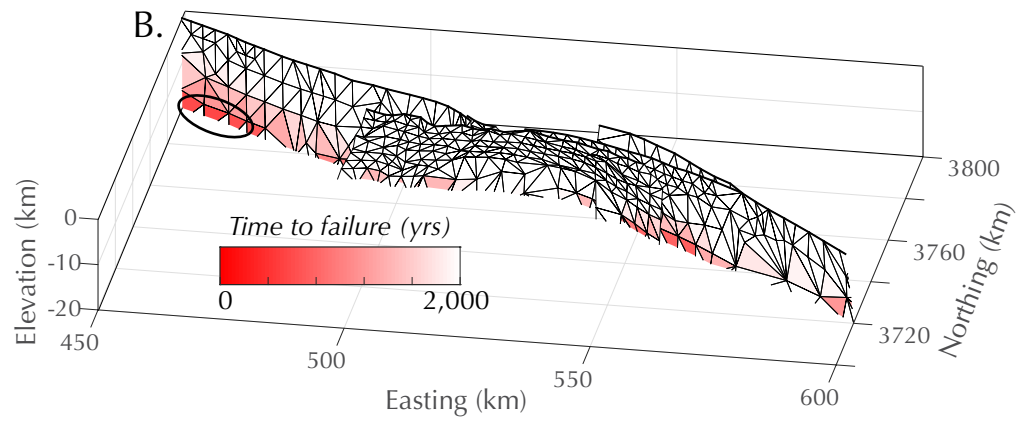
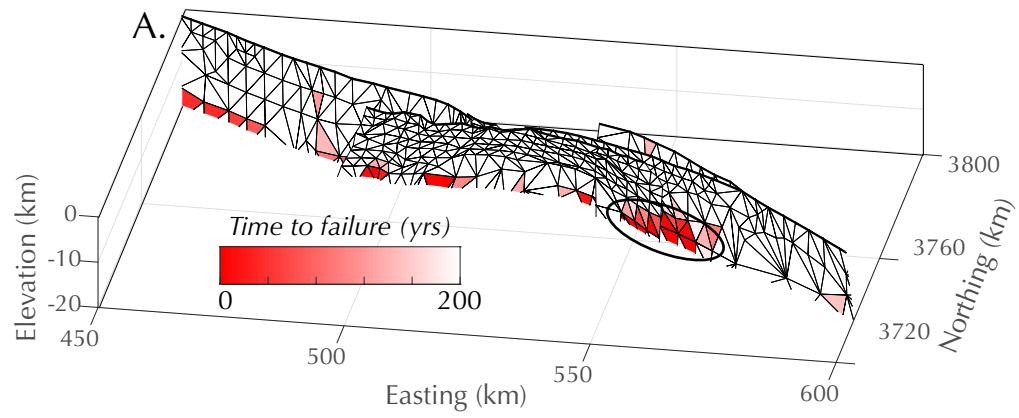












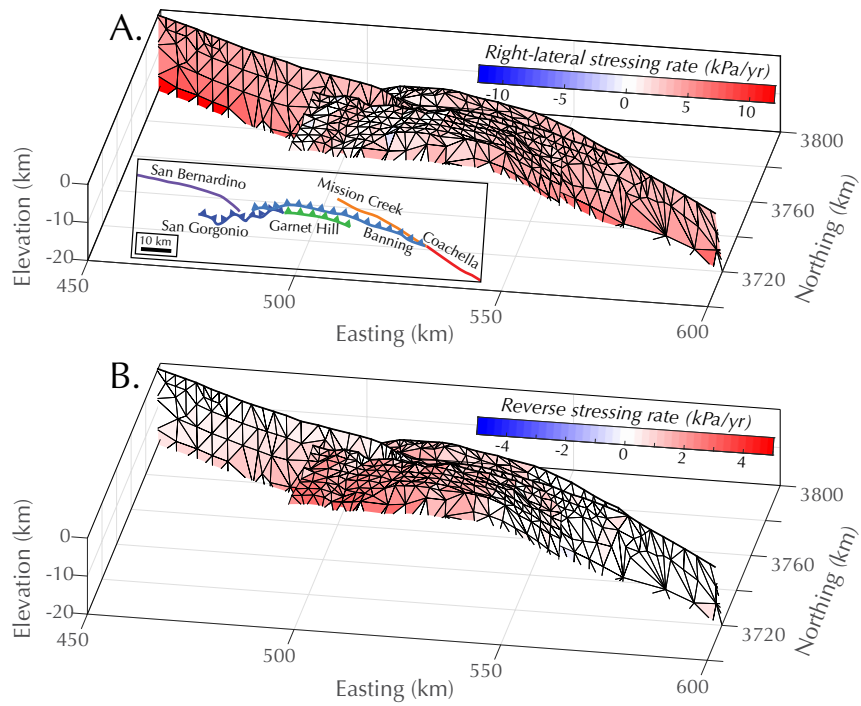


Figure S1. Right-lateral (A) and reverse dip slip (B) stressing rates for the alternative fault configuration from Beyer et al. (2018).