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Numerical modelling of lithosphere-asthenosphere interaction and intraplate deformation in the Gulf of Guinea

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Abstract

To present day, the phenomenon of intraplate deformation and its associated earthquakes remain elusive. In this work, we argue that intraplate deformation may result from the interaction between lithospheric and upper mantle dynamic processes. To this extent, we targeted the Gulf of Guinea and adjacent Western Africa, a region with both low plate velocities and clear asthenosphere dynamics, allowing us to isolate the individual underlying dynamic constraints which govern intraplate deformation. Thus, here we present 3D numerical geodynamic models of the asthenosphere-lithosphere interaction in the Gulf of Guinea, ran with the state-of-the-art modelling code LaMEM. We employ different initial/boundary conditions such as: (a) identical vs different spreading rates for the varying segments of the Atlantic mid-ocean ridge, (b) the presence/absence of weak zones (e.g., the Romanche/Central-African shear zones), and (c) the effect exerted by an active mantle plume, with a varying ascension velocity. Seismic data was used to evaluate the models and their validity. Our results suggest that intraplate deformation within the Gulf of Guinea is influenced by the spreading rate of the mid-ocean ridge, with stress being localized around the ocean-continent transition and existing shear zones. They also suggest that the existence of an underlying stress source (e.g., a mantle plume) beneath the Cameroon region is crucial to explain the epicenter distribution/deformation in the region.

35 **Keywords**

36 Intraplate deformation; Seismicity; Numerical modelling; Gulf of Guinea

37 **Highlights**

38 • The seismic event distribution in the Gulf of Guinea can be explained by the interaction between
39 two separate stress sources.

40 • The presence of lithospheric weak zones allows the stress from the mid-ocean ridge to be
41 transferred in-land, with the seismicity in the Southern part of Ghana partly reflecting this
42 process.

43 • A mantle plume underneath the Cameroon can explain the seismic events in the region.

44

45 1. Introduction

46 Intraplate deformation and associated intraplate earthquakes have been a widely studied problem in the
47 fields of geodynamics and seismology. While most large magnitude earthquakes occur mainly along plate
48 boundaries, some of these major events have occurred in regions usually associated with low stress and/or
49 lack of seismicity (such as along rifted passive margins, e.g. Olugboji et al., 2021). Their occurrence has
50 been explored over the past century, and many scenarios have been proposed to explain them, such as
51 anomalies in the thermal structure of the lithosphere (e.g., Eastern China; Weiran et al., 2009), distribution
52 of local geological structures such as (weak) fault zones (e.g., Bergman and Solomon, 1980; Talwani, 2017),
53 stress buildup (e.g., Calais et al., 2016), far-field stress transmission (e.g., Nkodia et al., 2020), gravitational
54 potential energy (Levandowski et al., 2017; Neres et al., 2018); and/or to deep lithosphere/upper
55 asthenosphere processes (e.g., He and Santosh, 2017; Pysklywec and Cruden, 2004; Wang and Currie,
56 2017). One region widely affected by intraplate seismicity, both on- and off-shore, is the Sub-Saharan
57 African margin. Along this transform continental margin (e.g., Sykes, 1978; Attoh et al., 2005; Jourdon et
58 al., 2021), there has been a significant amount of in-land seismic events since 1615 (Mohammadigheymasi
59 et al., 2023c), including 73 events of Magnitude ≥ 5 associated which caused local destruction of
60 infrastructures in the Accra (Ghana). Up to now, no relationship between the local stress fields and pre-
61 existing structures has been provided (Olugboji et al., 2021), despite extensive studies over the past century
62 (see Nkodia et al., 2022 and references therein). To further decipher the origin of this intraplate activity,
63 we explore a possible link between intraplate earthquakes/deformation along the Sub-Saharan African
64 margin and the dynamics of the underlying asthenosphere and upper mantle, specifically along the Gulf of
65 Guinea (GG, see Fig. 1). Within the GG, there are two possible stress sources, with one being the slow
66 spreading of the Mid Atlantic Ridge (MAR, c.a. 5-15 mm/yr, Müller et al., 2008) and its associated
67 compressive/shear stress (Turcotte and Schubert, 2002); while the second can be traced to the possible
68 source of the Cameroon Volcanic Line (CVL). While different hypotheses have been proposed for the
69 formation and maintenance of the CVL (these are extensively explored in Adams, 2022), the present work
70 adopted the working theory of the existence of mantle plume in the region (e.g., Adams, 2022; Celli et al.,
71 2020). Furthermore, we considered two major weak zones, which can act as stress localizers (Willis et al.,
72 2019; Bergman and Solomon, 1980; Talwani, 2017), the Romanche Fracture Zone (Attoh et al., 2004, 2005)
73 and the Central African Shear zone (Plomerová et al., 1993). Using the state-of-the-art geodynamic
74 modelling code LaMEM (Kaus et al., 2016), we developed a set of numerical geodynamical 3D models in
75 which we systematically explored the contribution of the stress sources and weak zones to intraplate
76 deformation in the GG area. Our modelling conditions include the systematic testing of different initial and
77 boundary conditions such as:

- 78 • The shear stress induced by the individual spreading ridge segments of the Atlantic mid-ocean
79 ridge. This parameter was controlled by testing identical vs different spreading rates in the

80 individual segments. With identical spreading rates, the oceanic plate is pushed equally toward
 81 the continent by each individual ridge segment, inducing identical compressive stresses into the
 82 continent. Given this, any heterogeneity will be the result of the geometry/rheology of the
 83 continent itself.

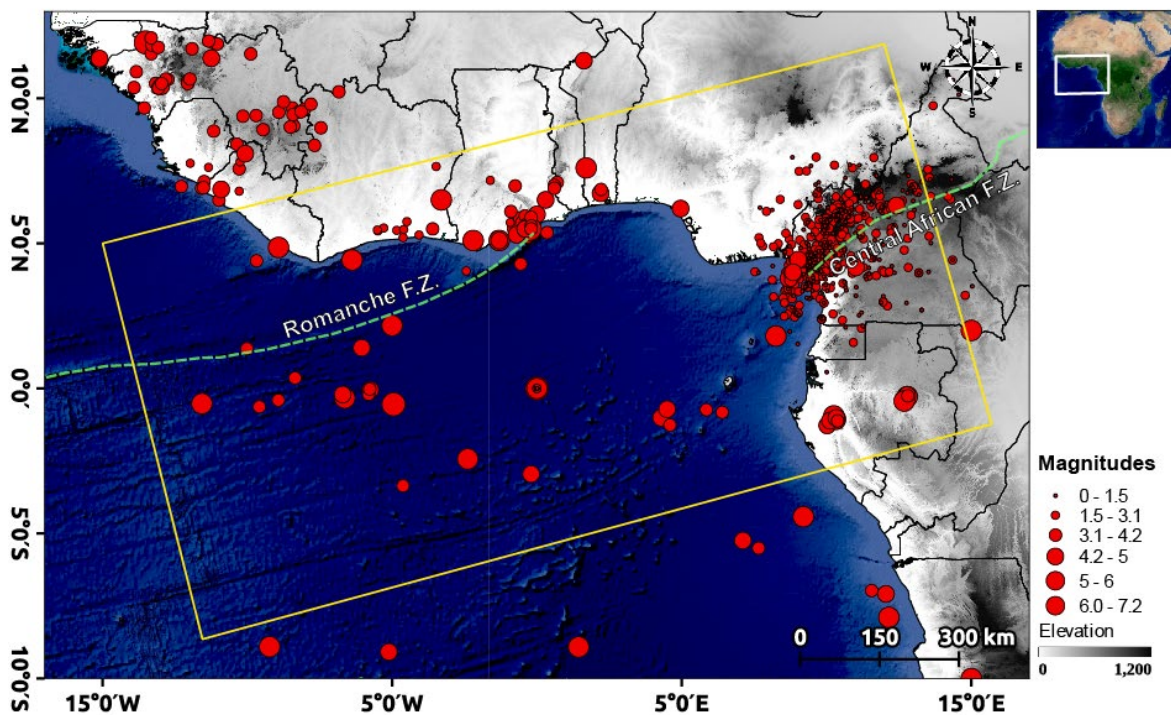
- 84 • The effects induced by the presence of weak zones (e.g., the Romanche/Central-African shear
 85 zones). To evaluate their effective contribution in lithospheric stress distribution, models were
 86 run with and without these weak structures while maintaining all other conditions.
- 87 • The regional stress changes induced by an encroaching active mantle plume. As with the
 88 previous condition, models were run with and without an active mantle plume.

89 **2. Geological and tectonic context**

90 The GG is a region famous for its coastlines, the circum-south Atlantic African coast, which perfectly match
 91 the coast of the south-America continent, forming a piece of the continental drift puzzle which contributed
 92 to the current theory of plate tectonics (Wegener, 1920; Le Pichon, 1968; Burke et al., 1971). Beyond the
 93 geometrical argument, geological and structural evidence based on the correlation between the Pan-
 94 African and Brazilian orogenic belts, shear zones and cratons were clearly established (Ledru et al., 1994;
 95 Caxito et al., 2020 and references therein) indicating that the two continental blocks were initially part of
 96 a same supercontinent, Gondwana. Moreover, these common structures and faciès present in both
 97 continents provide additional information on the phases of deformation that predate and participated to
 98 the formation of Gondwana dated to 630-500 Ma (Caxito et al., 2020 and references therein). These latter
 99 are confined in between the present-day cratons, namely the West Africa and Congo cratons for the African
 100 block and the Amazonian and São Francisco cratons for the Brazilian block, and testify of the past presence
 101 of oceanic basins, continental magmatism, multiple subduction arcs and orogens which took place between
 102 1000 and 630 Ma (Caxito et al., 2020 and references therein). Thus, along the GG coast, two main suture
 103 zones can be found, one along the eastern border of the West Africa craton. It is mainly known as the
 104 Dahomeyides belt, which crosses the countries of Ghana, Togo and Benin. A second one, the Central African
 105 Orogen, mainly located in Cameroon, borders the northern part of the Congo craton.

106 The breakup of Gondwana initiated after 250 Ma, was featured by multiple phases of rifting, themselves
 107 preceded by plume-related flood volcanism, leading to the progressive opening of the Central and South
 108 Atlantic (Fairhead and Green, 1989; Renne et al., 1992; Courtillot et al., 1999). This is during the period
 109 between the late stage of orogenes (<590 Ma) and progressive Gondwana breakup that an extensive
 110 network of shear zones developed, correlated in space and time with post-collisional granitic magmatism
 111 (Fairhead, 1988; Caxito et al., 2020). Among them, the Kandi-Transbrasiliano shear zone represents the
 112 main trans-continental structure which extends from the Hoggar region to Ghana in Africa and to the East

113 of the Amazonian craton for the Brasil. The corner of the GG is also featured by several shear zones and
 114 their splays, among which the Central African Shear Zone (CASZ) extends from the GG's corner in Cameroon
 115 to Darfur region in western Sudan. The GG also features more recent volcanic activity, that of the CVL which
 116 initiated about 30 Ma in between the oceanic domain with four islands - Pagalu, São Tomé, Príncipe and
 117 Bioko, and the continent domain, forming a 1000 km-long line (Déruelle et al., 1991). The CVL volcanism
 118 does not show any clear hotspot-like age progression, partly explaining why several hypotheses on its origin
 119 have been proposed (Déruelle et al., 1991; Burke, 2001). Presently, the main volcanic activity of CVL is
 120 focused in between Mount Cameroon volcano and Bioko Island (Tabod et al., 1992; DePlaen et al., 2014).
 121 As with most intraplate settings, GG shows a rather low seismicity which contrasts with the plate boundary
 122 formed by the Mid-Atlantic Ridge to the West of the GG (Meghraoui, 2016; Meghraoui et al., 2019). Despite
 123 this, a persistent seismic activity including the occurrence of a $M_w > 6$ earthquake has been recorded in
 124 South Ghana. There, the suture zone connects the continental tectonic structures to the oceanic domain
 125 with the Romanche fracture zone (Meghraoui et al., 2019; Mohammadigheymasi et al., 2023c and
 126 references therein).



127

128 **Figure 1: The Gulf of Guinea.** Modelled area is highlighted by a light yellow rectangle. Epicenters in the Mid
 129 Atlantic Ridge (MAR) are omitted for clearer visualization of MAR's structure. The grayscale background
 130 represents a Digital Elevation Model (DEM) provided by NASA Shuttle Radar Topography Mission (SRTM)
 131 (2013).

132

133

134 3. Methods

135 3.1. Data collection

136 As a tool to validate numerical models, we first compiled relevant existing data characterizing the crustal
137 structure, tectonic features, seismic anisotropy, and seismicity of the GG. This comprehensive compilation
138 draws from various studies and includes data from published papers and data centers. Noteworthy
139 parameters of our research encompass:

- 140 • **Crustal thickness, Moho depth, compressed and shear-wave velocities, and crustal stress field:**
141 Of particular significance is the extensive characterization the crustal structure of Nigeria, Ghana,
142 and Cameroon, encompassing information on V_p/V_s and Poisson's ratios, as well as Moho
143 depth. This data has been compiled from the following sources Elsheikh et al. (2014); Custódio et
144 al. (2022); Kamguia et al. (2005); Fairhead and Okereke (1988); Stuart et al. (1985); Tokam et al.
145 (2010); Akpan et al. (2016); Mohammadigheymasi et al. (2023c, 2024).
- 146 • **Seismic catalog:** The catalog we compiled integrates various sources, including the
147 comprehensive historical and instrumental catalog assembled by Musson (2014), event reports
148 from the Ghana Geological Survey (GHGS), and data from the International Seismological Center
149 (ISC). Furthermore, it reflects the results of recent efforts to reprocess existing digital seismic
150 data in the GG, utilizing deep learning methods (Mohammadigheymasi et al., 2022, 2023a,b;
151 Carvalho et al., 2023; Mohammadigheymasi et al., 2024). The catalog spans a period from 1615
152 to 2023, encompassing 1723 earthquakes with moment magnitudes M_w ranging from 1 to 7.2,
153 covering a depth range from 60 km below the Earth's surface up to the surface. A graphical
154 representation of this catalog is provided in Fig. 1.

155 A summary of the compiled crustal structure, thickness, and average crustal velocity in Cameroon, Ghana,
156 and Nigeria is presented through tables, figures, and a supplementary GIS database. The raw data and GIS
157 datasets are stored in the supplementary GitHub repository of this paper, a link for which is provided in
158 the section 7 – Data Availability.

159 3.2. Numerical approach

160 The conducted numerical models were run using the LaMEM code (Kaus et al., 2016), using internally
161 imposed kinematic conditions to simulate the MAR spreading and the Cameroon plume upwelling. No
162 compressibility was assumed for these modelling runs. LaMEM employs a finite difference staggered grid
163 discretization which is coupled with a particle-in-cell approach (Kaus et al., 2016) as to obtain numerical
164 solutions for the conservation equations of mass, momentum, and energy (eq. 1-3)

165

$$\frac{\partial \mathbf{v}_i}{\partial x_i} = 0 \quad (1)$$

$$-\frac{\partial P}{\partial x_i} + \frac{\partial \tau_{ij}}{\partial x_j} + \rho \mathbf{g}_i = 0 \quad (2)$$

$$\rho C_p \left(\frac{\partial T}{\partial t} + \mathbf{v}_i \frac{\partial T}{\partial x_i} \right) = \frac{\partial}{\partial x_i} \left(\kappa \frac{\partial T}{\partial x_i} \right) + H_R + H_S \quad (3)$$

166

167 Here, i and j are coordinate indexes, \mathbf{v}_i represents the velocity vector, x_{ij} the cartesian coordinates, P the
 168 pressure, τ_{ij} the shear stress, ρ the density, \mathbf{g} the gravitational acceleration vector, C_p the specific heat, T
 169 the temperature, t the time, κ the thermal conductivity, and H_R and H_S represent the radiogenic and shear
 170 heating components, respectively. The shear heating component is defined as:

$$H_S = \tau_{ij} (\dot{\epsilon}_{ij} - \dot{\epsilon}_{ij}^{\text{elastic}}) \quad (4)$$

171 with ϵ_{ij} as the total deviatoric strain rate tensor and $\dot{\epsilon}_{ij}^{\text{elastic}}$ the deviatoric elastic strain rate tensor.

172 All presented models were run using non-linear viscoelastoplastic rheology, with the following constitutive
 173 equations (Kaus et al., 2016; Piccolo et al., 2020):

$$\dot{\epsilon}_{ij} = \dot{\epsilon}_{ij}^{\text{viscous}} + \dot{\epsilon}_{ij}^{\text{elastic}} + \dot{\epsilon}_{ij}^{\text{plastic}} = \frac{\tau_{ij}}{2\eta_{eff}} + \frac{\dot{\tau}_{ij}}{2G} + \dot{\gamma} \frac{\partial Q}{\partial \tau_{ij}} \quad (5)$$

$$\dot{\tau}_{ij} = \frac{\partial \tau_{ij}}{\partial t} + \tau_{ik} \omega_{kj} - \omega_{ki} \tau_{kj} \quad (6)$$

$$\omega_{ij} = \frac{1}{2} \left(\frac{\partial \mathbf{v}_j}{\partial x_i} - \frac{\partial \mathbf{v}_i}{\partial x_j} \right) \quad (7)$$

174 with η_{eff} as the effective viscosity, $\dot{\tau}_{ij}$ the Jaumann objective stress rate, ω_{ij} the spin tensor, G the elastic
 175 modulus and Q the plastic flow potential.

176 The creep viscosity, η_{vs} , is calculated as:

$$\eta_{vs} = \frac{1}{2} A^{-\frac{1}{n}} \times \dot{\epsilon}_{II}^{\frac{1}{n}-1} \times \exp \left(\frac{E_a + V_a P}{nRT} \right) \quad (8)$$

177 with A as the diffusive or dislocation pre-exponential factor, n the stress exponent, $\dot{\epsilon}_{II}$ the square root of
 178 second invariant of the deviatoric strain rate tensor (eq. 5), E_a the activation energy, V_a the activation
 179 volume and R the gas constant.

180 Plastic flow is ensured by employing a Drucker-Prager yield criterion (Drucker and Prager, 1952):

$$\sigma_Y = C \cos(\phi) + P \sin(\phi) \quad (9)$$

181

182 with σ_y as the yield stress tensor, ϕ the internal friction angle and C the cohesion. The onset of plastic
 183 weakening takes place once mantle materials accumulate at least 10% of total plastic strain and this effect
 184 halts after at least 60% of total plastic strain has been accumulated. During softening, the materials'
 185 cohesion and internal friction angles are linearly reduced until they reach 1% of their initial values. The
 186 effective viscosity (η_{eff}) of the individual phases is obtained by calculating the minimum between the
 187 calculated viscoelastoplastic viscosity and the Newtonian viscosity.

188 The second invariant of the stress tensor (σ_{II}) is obtained as follows:

$$\sigma_{II} = \sqrt{\frac{1}{2} \sigma'_{ij}{}^2} \quad (10)$$

189 in which $\sigma'_{ij}{}^2$ is the sum of the square of all individual deviatoric stress tensor components.

190 The age dependence of the thermal profiles of the plates follows the half-space cooling model:

$$T = T_{surface} + (T_{mantle} - T_{surface}) \times \operatorname{erf}\left(\frac{y}{\sqrt{\kappa t}}\right) \quad (11)$$

191 Here, $T_{surface}$ represents the temperature at the surface of the model (273 K), T_{mantle} is the temperature
 192 at the lithosphere-asthenosphere boundary (1523 K), y the depth, κ the thermal diffusivity, and t the age
 193 of the plate. The effective (rheological) lithosphere thickness throughout the model is set by the 1523 K
 194 (1250 °C) isotherm. The upper mantle thermal profile follows the mantle adiabat, with a gradient of 0.5
 195 K/km. All material densities are temperature and pressure dependent:

$$\rho = \rho_0 + \alpha(T - T_0) + \beta(P - P_0) \quad (12)$$

196 Here, ρ_0 is the density of the material at the reference temperature T_0 , α is the thermal expansibility and β
 197 is the compressibility. All rheological parameters can be found in Table 1.

198 The calculation of the modelled stress regime follows the approach defined in Delvaux et al. (1997), by
 199 calculating a ration (R) between the primary components of the stress tensor ($\sigma_1, \sigma_2, \sigma_3$).

$$R = \frac{\sigma_2 - \sigma_3}{\sigma_1 - \sigma_3} \quad (13)$$

200 This ration is then used to calculate a stress index (R') which is defined based on the primary stress
 201 directions as follows:

$$\begin{cases} R' = R & , \sigma_1 \text{ vertical} \\ R' = 2 - R & , \sigma_2 \text{ vertical} \\ R' = 2 + R & , \sigma_3 \text{ vertical} \end{cases} \quad (14)$$

202 The specific stress regime is based on the value obtained for R' with $R' \in [0, 1]$ describing extensive, $R' \in$
 203 $]1, 2[$ strike-slip, and $R' \in [2, 3]$ compressive conditions, respectively.

204 **3.3. Initial setup and modelling approach**

 205 A set of 3D models were performed to simulate the kinematics of the sub-Saharan African margin. To assess
 206 the influence of the individual controlling factors, we systematically increased the complexity of our models
 207 by adding one factor per model, resulting in a total of four exploration models. We ran an additional model
 208 in which we test a mid-ocean ridge with no spreading rate variability along the individual segments, as well
 209 as a model with a wider plume head, for a total of six models (see Table 2).

 210 The prescribed model domain was 3200 km long, 1500 km wide and 710 km thick (see Fig. 2) and was
 211 discretized along a 384x192x128 resolution grid, resulting in a c.a. 8x8x6 km cell. This resolution allowed
 212 us to prescribe two of the major weak zones present in the Gulf of Guinea (namely, the Romanche and
 213 Central African Shear zones, see Fig. 2) and ensure they have a minimum width of 16 km. The model
 214 included a 50 km thick sticky-air layer which acts as a free surface, allowing for the formation of
 215 topography. Furthermore, the top boundary is open, ensuring free movement of this layer. All other model
 216 boundaries were defined as free slip, which allows for motion along the direction of the boundary but not
 217 across it.

218

 219 **Table 1** – Physical parameters applied in the models, for each of the different phases. UC -
 220 (Continental) upper crust, LC - (Continental) lower crust. Creep and thermal parameters for the
 221 oceanic and upper mantle phases adapted from Kohlstedt et al. (1995), Ranalli and Karato (1995), and
 222 Ranalli (1997). Creep and thermal parameters for the continental phases from Ranalli and Karato
 223 (1995). Elastic parameters from Kaus et al. (2015). Plastic flow parameters from Ranalli and Karato
 224 (1995) and Ranalli (1997) and Li et al. (2010).

225

Properties	Parameters	Continent		Oceanic Crust	Upper Mantle
		UC	LC		
Density	ρ_0 - kg/m ³	2800	3000	3000	3300
Dislocation creep	B_n - Pa ⁻ⁿ /s ⁻¹	6.40x10 ⁻⁶	3.30x10 ⁻⁴	2.50x10 ⁴	2.50x10 ⁴
	n	2.40	3.20	3.5	3.5
	E_n - J/MPa/mol	1.56x10 ⁵	2.38x10 ⁵	5.30x10 ⁴	5.30x10 ⁴
	V_n - m ³ /mol	0	0	1.35x10 ⁻⁵	1.35x10 ⁻⁵
Diffusion creep	B_I - Pa ⁻ⁿ /s ⁻¹	-	-	1.50x10 ⁹	1.50x10 ⁹
	n	-	-	1	1
	E_I - J/MPa/mol	-	-	3.75x10 ⁵	3.75x10 ⁵
	V_I - m ³ /mol	-	-	1.35x10 ⁻⁵	1.35x10 ⁻⁵
Plastic flow	Cohesion - MPa	20	20	30	30
	Softened cohesion - MPa	0.2	0.2	0.3	0.3
	ϕ - degrees	6.8	6.8	20	20
	Softened ϕ - degrees	0.068	0.068	0.2	0.2
Elasticity	G - Pa	5.00x10 ¹⁰	5.00x10 ¹⁰	5.00x10 ¹⁰	5.00x10 ¹⁰
Thermal	Expansivity (α) - 1/K	3x10 ⁻⁵	3x10 ⁻⁵	3x10 ⁻⁵	3x10 ⁻⁵

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Properties	Compressibility (β) – Pa ⁻¹	1x10 ⁻¹¹	1x10 ⁻¹¹	1x10 ⁻¹¹	1x10 ⁻¹¹
	Conductivity (κ) - W/m/k	3	3	3	3
	Heat capacity - J/kg/K	1050	1050	1050	1050
	Radiogenic heat production – W/kg	5.36x10 ⁻¹⁰	2.32x10 ⁻¹⁰	6.67 x10 ⁻¹²	6.67x10 ⁻¹²

226

227 **Table 2:** List of models run in the present study

Model number	Active ridge	Active plume	Weak zones	Additional constraints
1	Yes	No	No	-
2	Yes	No	Yes	-
3	Yes	No	No	No ridge segment variability
4	Yes	Yes	No	-
5	Yes	Yes	Yes	-
6	Yes	Yes	Yes	Wider plume head

228

229 The rheology of the oceanic crust follows a dry olivine creep law (Ranalli, 1997) and had a variable thickness
 230 which depended on its distance to the spreading ridge centre (half-space cooling model, eq. 10). The
 231 lithospheric mantle also follows a dry olivine creep law and differs in behaviour from the crustal material
 232 due to a higher temperature (eq. 8).

233 Our modelled African continent was internally divided into blocks, each with a specific crust thickness and
 234 lithospheric mantle thicknesses, which was obtained from geophysical studies (see references in Section
 235 2.1). For simplicity, we assumed that the crust is vertically divided into equally thick upper and lower crust.
 236 We assumed a quartzitic creep (Ranalli and Karato, 1995; Ranalli, 1997) and plastic flow (Li et al., 2010)
 237 laws for the materials of the upper crust continental. This setting reproduces the overall heterogeneous
 238 nature of continental crust, allowing for a weak/brittle upper crust. The lower crust was assumed to follow
 239 a stronger plagioclase creep/plastic flow (Ranalli, 1995), ensuring a more consistent stress/strain
 240 delocalization and replicating the well-established jelly-sandwich rheological model for continental
 241 lithosphere (Burov, 2011).

242 The internal dynamics of the model are controlled by pre-imposed kinematic conditions, such as the
 243 spreading rate of the individual ridge segments along the MAR as well as the injection rate of the high
 244 temperature materials of the Cameroon Plume. The geometry and kinematics of the MAR are simulated
 245 by implementing four independent kinematically divergent zones, with variable spread rates (see Fig. 2 and
 246 Table 3). Each segment is laterally offset from the previous as to favour the establishment of orthogonal
 247 transform faults between them and to depict at best the dynamics and related stress distribution of the
 248 Atlantic basin for the GG region. The Cameroon plume was modelled by implementing an injection point
 249 at the bottom boundary of the model, in which hotter material is added at a constant rate over time. The
 250 composition of the plume material was assumed to be identical to the surrounding mantle (see Table 1)
 251 but, due to a positive buoyancy (derived from a thermally lowered density), it is forced to ascend. While
 252 the specific geometry of this plume cannot be predicted from the start, we imposed a width of 50 km for
 253 the stem and an injection velocity of 20 cm/yr. This ensured that the plume head reaches the base of the

254 lithosphere at approximately the same place as in nature (see Fig. 2) when accounting for the spreading
 255 rate of the MOR.

256 The weak zones implemented (which represent two major fault zones present in the region, the Romanche
 257 and Central African Shear Zones, see Fig. 1) were defined by an isoviscous rheology (i.e., with a constant
 258 low viscosity) of 10^{19} Pa·s as to achieve high strain localization.

259

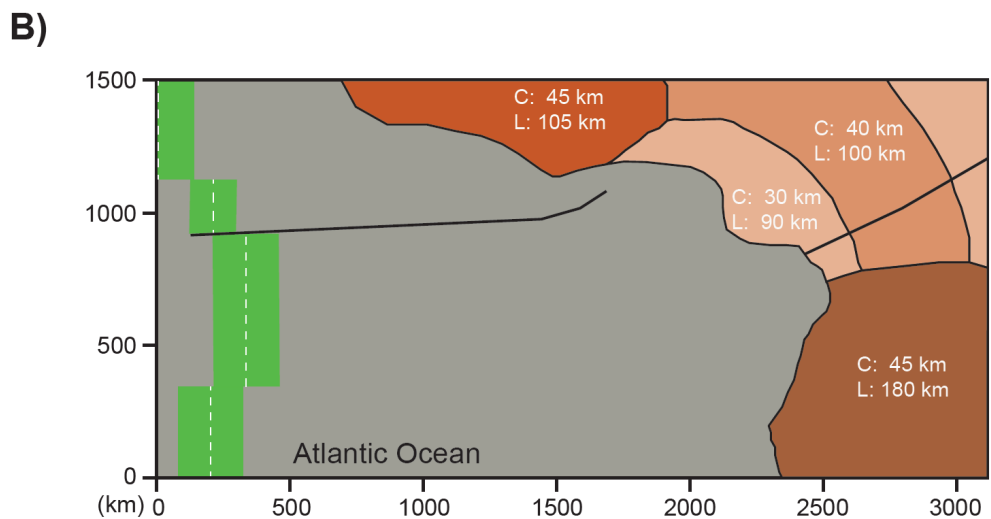
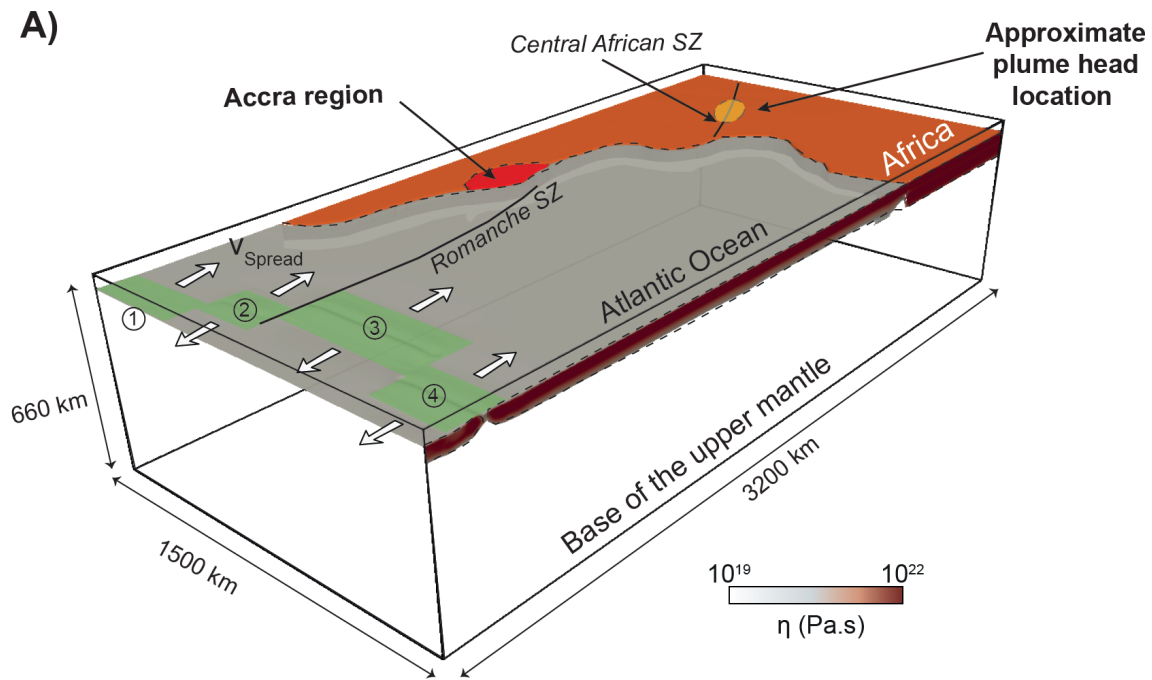
260 **Table 3 – Right and left spreading rates in the simulated ridge segments.** The segments are
 261 identified in Fig. 2-A. Here, positive spreading rate values indicate a spread towards the continent.
 262 Conversely, negative values indicate towards the wall.

263

Segment	Left spreading rate (cm/yr)	Right spreading rate (cm/yr)
1	-	0.50
2	-0.50	1.00
3	-1.25	1.25
4	-1.20	1.20

264

265



266

267 **Figure 2: A. Model setup for the experimental initial state:** schematic representation of the geometric
 268 configuration and model dimensions. The placement of the plume head is approximate and follows
 269 previous descriptions (Burke, 2001). The viscosity scale applies only to the front-facing wall. The numbered
 270 green polygons mark the position of the individual ridge segments, with their spreading direction indicated
 271 by the white arrows. Their half-spreading rates can be found in Supplementary Table A.2. The depth
 272 indicated along the z axis does not include the 50 km thick sticky air layer. **B. Continental crustal and**
 273 **lithospheric thicknesses:** the different shaded regions indicate the individual cratonic regions present in
 274 our model. Crustal thickness data was described in Section 2.1. Lithospheric thickness from Globig et al.
 275 (2016).

276

277 **4. Results**

278 In this work, we investigated the role of MAR, the two major weak zones (see Fig. 1) and the CVL plume as
279 well as that of their possible interplay in the stress distribution and transfer within the GG. We depart from
280 a simple reference model (Model 1) where only the MAR is active. Then, we incrementally added
281 complexity in subsequent iterations. This ensured that any detected variation would be explained by the
282 added component. In both the single and dual stress source models, two major phases can be defined in
283 which the same broadly defined events occur, namely: Phase 1) establishment of the mid-ocean ridge, and
284 Phase 2) stress migration away from the ridge and towards the continent (see Figs. 3 and 4). An additional
285 phase is seen in the single source models which entails the development of combined continent-ocean
286 deformation.

287 **4.1. Phase 1 – Ridge establishment**

288 In all models (regardless of the type and number of stress sources), the first 0.3 Myr consisted of the
289 establishment of the spreading centers along the ridge. During this time period, stress is almost exclusively
290 confined to these extensional structures (see Figs. 3 and 4, left-most column). Although all models show
291 similar results, the models with pre-existing weak zones (models 2, 5 and 6, in Figs. 3 and 4) show a lower
292 stress accumulation at the MAR at this stage. Furthermore, these same models also show a pronounced
293 stress localization zones both in the ocean-continent transition (OCT) zone and within the continent, with
294 the former showing higher stress values (compare in Figs. 3 and 4 the top with the middle rows).

295 **4.2. Phase 2 – Deformation migration**

296 Once the ridge system is established ($\Delta t > 0.3$ Myr, middle columns in Fig. 3), stress begins to be transmitted
297 to the more distant oceanic plate. Additionally, we also observe incipient stress localization within the
298 continent, specifically in and around the thinner continental blocks (see Fig. 2, with the thicker regions still
299 showing low to no stress accumulation).

300 The addition of pre-existing weak zones directly affects both the stress distribution and migration velocity.
301 Under their effects (Model 2, middle row in Fig. 3), not only does stress localize quicker within the
302 continent, but it is also maintained for a longer period, with distinct higher stress regions being observed
303 for the rest of the model.

304 Finally, when the individual segments of the MOR have identical spreading rates (Model 3, Fig. 3, bottom
305 row), the stress propagation in the oceanic plate appears to be less efficient, with the latter model showing
306 overall lower stress values, with more pronounced stress contrasts within the continent (compare in Fig. 3
307 the top and bottom rows). Nevertheless, both Model 1 and 3 show the same tendency towards the
308 homogenization of stress within the continent, as opposed to the formation of stress bands shown in Model
309 2. In the models with two stress sources, this phase coincides with the arrival of the plume head to the

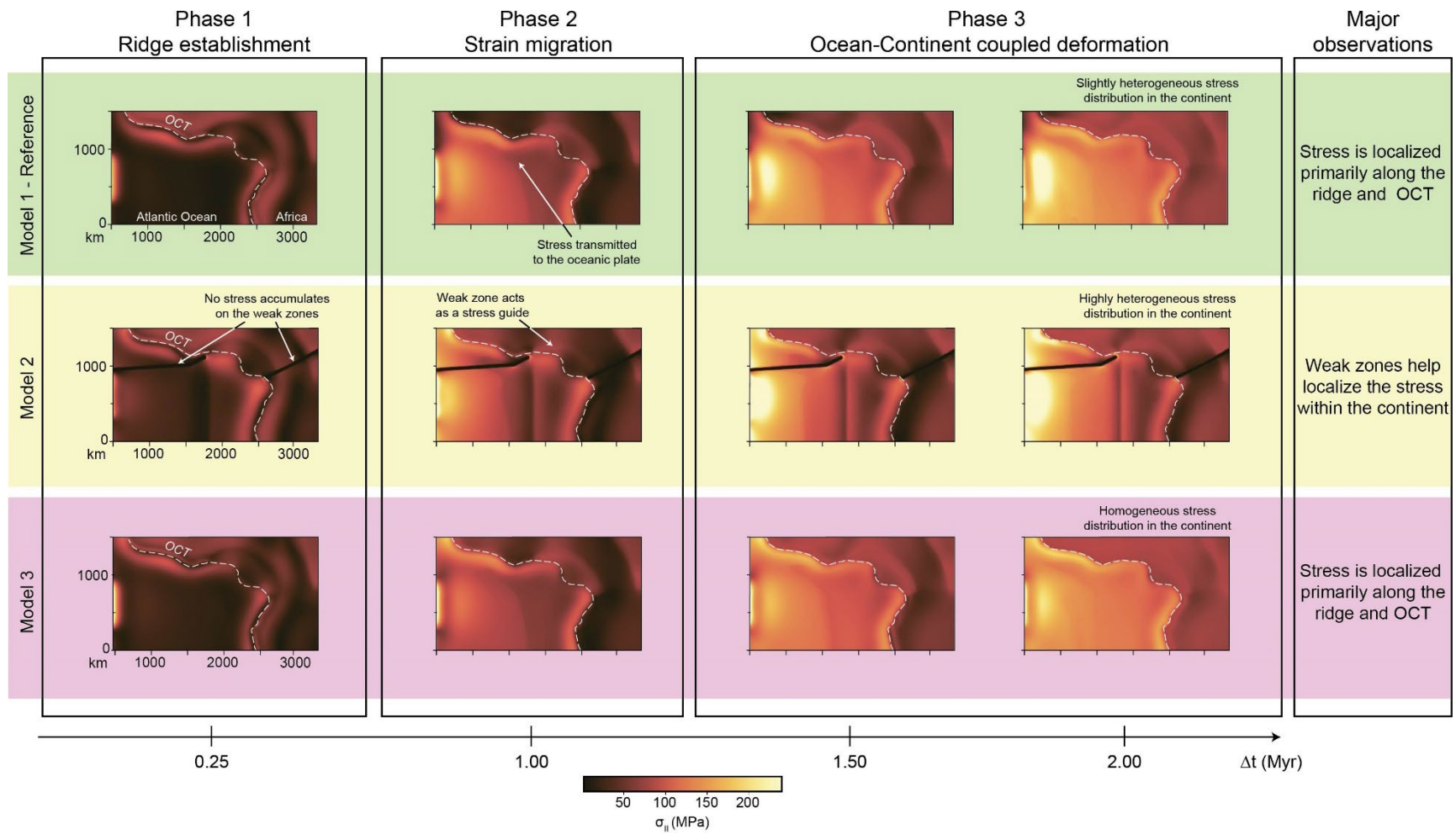
310 base of the lithosphere, as observed in Fig. 4. In all three of these models, the plume encroaches the base
311 of the lithosphere, straining it the base of the lithosphere (see Fig. 5), which effectively lowers the viscosity
312 at this location (note in Fig. 4 the lower viscosity values above the 1000°C line). The surface expression of
313 this plume arrival, however, is not identical for all scenarios. We can observe in Fig. 4 that while the two
314 models with pre-established weak zones show some surface expression of the plume arrival, the model
315 without weak zones does not. In addition, this surface expression is more pronounced in the model with a
316 wider mantle plume head.

317

318 **4.3. Phase 3 – Combined ocean-continent deformation**

319 This last phase is observed in the models with a single stress source (i.e., Fig. 3), taking place after 2 Myr of
320 model time. After this period, we observe the establishment of the final stress localization sites, which
321 appear to be controlled by the presence of the weak zones. When these are not present, stress distribution
322 shows a tendency towards homogenization within the continent, where only relatively weak stress
323 contrasts are seen along craton boundaries (see in Fig. 3 the top and bottom rows). By contrast, weak zones
324 allow for the maintenance of stress localization sites which are still observable within the continent. These
325 sites appear along the weak zone tips, as well as along the boundaries of the cratons (see Fig. 3, middle
326 row).

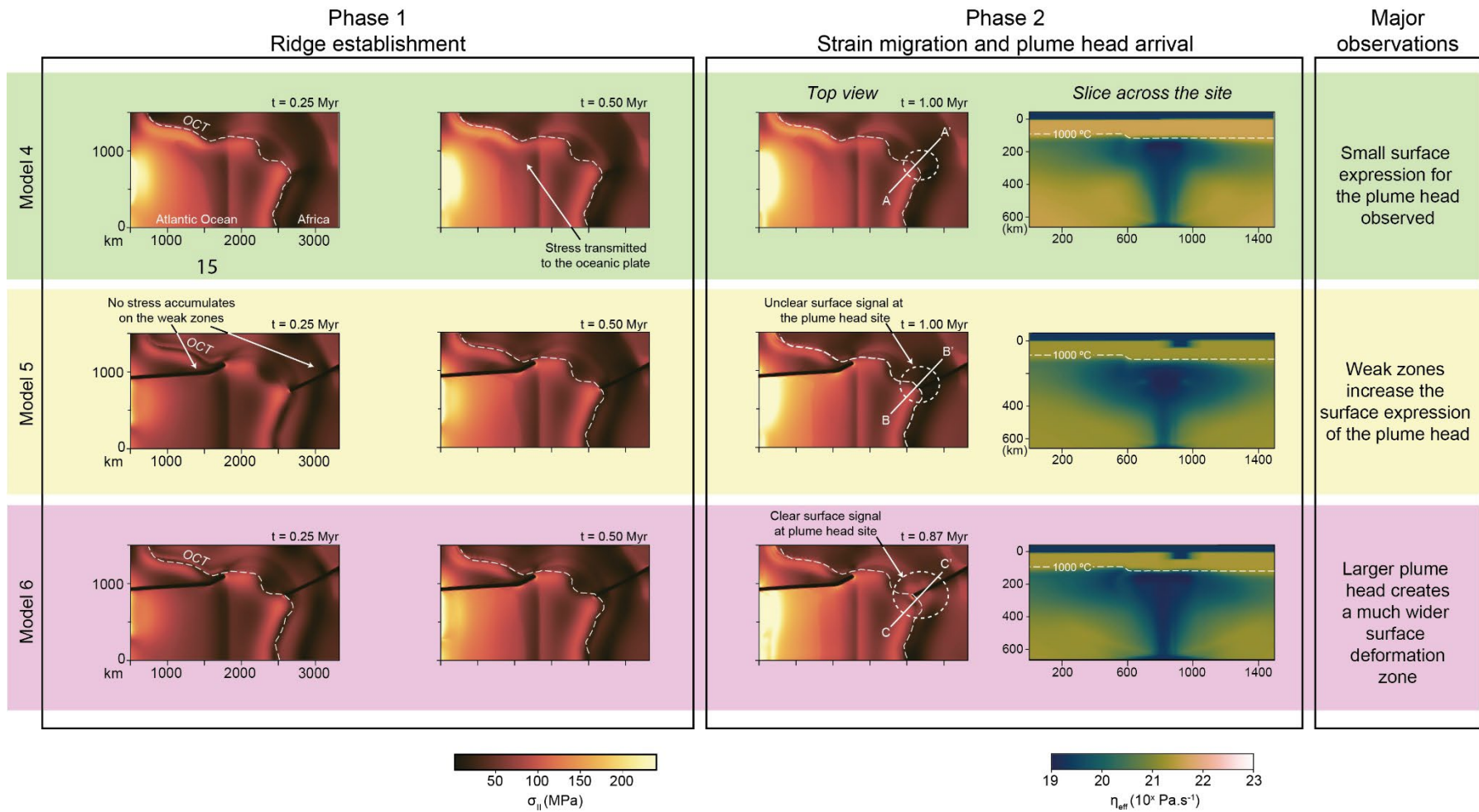
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328 (See caption in the next page)

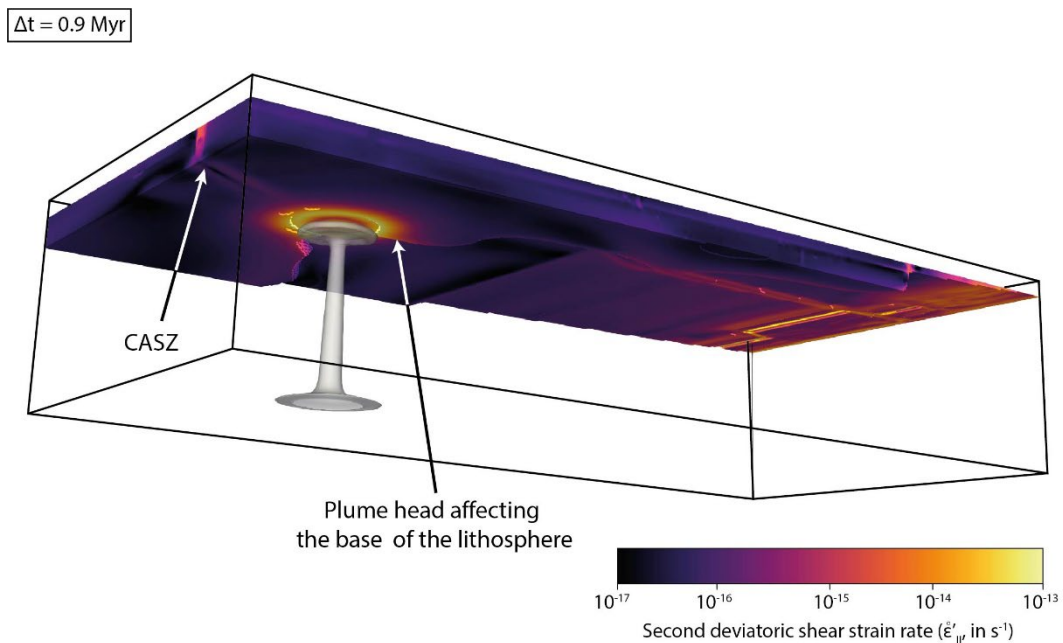
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330 **Figure 3: Major phases and events observed in the single stress source models in terms of deviatoric**
331 **stress (σ_{II}) distribution.** Each column of figures corresponds to identical model times, while each row
332 describes the evolution of a specific model. The right-most column summarizes the key observations for
333 each model. Phase 1 (left-most column) shows first-order identical stress distributions between the three
334 models, with Model 2 (middle row) showing a more effective dissipation of the ridge stress (lower overall
335 σ_{II}). The same is mostly observed in Phase 2 (second column to the left), in which all models show a simple
336 an increase of stresses within the oceanic plate and, in the case of Model 2, increased stresses within the
337 continent. Lastly, Phase 3 shows a more advanced stage of the system evolution, highlighting the major
338 differences between the three models: both models 1 and 3 show the eventual establishment of (mostly)
339 homogeneous stress distribution in the continent while, by contrast, Model 2 shows a heterogeneous
340 stress state. The white dashed line represents the transition between the ocean and the continent, or OCT.
341 Difference from reference model can be found in Supplementary Fig. 1.



343 **Figure 4: Major phases and events observed in the dual stress source models in terms of**
 344 **deviatoric stress (σ_{II}) distribution.** The three left columns of figures correspond to identical
 345 model evolution stages (specific model time is indicated for each), while each row describes
 346 the evolution of a specific model. In Phase 1 (two left-most columns), the model results are
 347 broadly similar to the ones shown in Fig. 3, with the models with weak zones showing more
 348 effective stress dissipation at the ridge (lower overall σ_{II} values). In Phase 2, the plume has
 349 arrived at the base of the lithosphere in all three models (see the right column in Phase 2,
 350 showing a cross section across the plume site). The surface expression of this plume is a
 351 function of both the size of the plume head and the presence of weak zones, as (a) only the
 352 models with weak zones show any surface expression (middle and bottom rows), and (b) this
 353 expression is much more pronounced in the model with a wider plume (bottom row). The white
 354 dashed line represents the transition between the ocean and the continent, or OCT. Difference
 355 from reference model can be found in Supplementary Fig. 1.

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Figure 5: Side view of model 6, after a model time of 0.9 Myr. Here, we observed that a wide high strain ring formed around the top of the plume head. This provided a strong indication that the base of the lithosphere was under a strong strain induced by the arrival of the plume head. Furthermore, by observing the left side of the figure, we also observed that the entirety

365 of the Central African shear zone (CASZ) is also affected, showing relatively higher strain values
366 along its path. The plume contour is done by mapping the 1700°C isotherm.

367

368

369 **5. Discussion**

370 To better decipher the processes leading to the intraplate seismicity and deformation in the
371 GG, our approach consisted of: 1) evaluating the role of the oceanic lithosphere in the stress
372 distribution in the GG region and especially at OCT; and 2) independently test the effect of two
373 sources of stress, namely the MAR and the Cameroon Plume.

374 **5.1. Single stress source**

375 In the models in which the MAR was the sole stress source (i.e., Fig. 3, models 1 and 3), we
376 observe a homogeneous low stress state within the modelled African continent. While these
377 stress states can generate complex strain distributions, when applied in heterogeneous media
378 (e.g., Gerya, 2009), our simulated African continent assumes a constant (i.e., homogeneous)
379 composition with only age differences. Thus, under present modelling conditions, most of the
380 stress will be localized along rheological boundaries, such as the OCT and cratonic boundaries,
381 following the observations of prior studies on deformation of continental lithosphere (e.g.,
382 Vauchez et al., 1998; Calignano et al., 2015).

383 The addition of weak zones in this system has a considerable effect in the stress distribution
384 (compare in Fig. 3, models 1 and 2). Along the ridge-distal tip of our modelled Romanche shear
385 zone, we can observe a significant increase of the stress within the oceanic and continental
386 domains, as well as clear stress contrast along the cratonic boundaries in the region
387 (boundaries shown in Fig. 2). This in-land stress propagation is consistent with previous
388 transform margin modelling studies in which the formation and development of moderately
389 oblique transform faults induces deformation in the continent (e.g., Jourdon et al., 2021).
390 While the low stress within the weak zone (see Fig. 3, Model 2) indicates that the majority of
391 the deformation is being accommodated here, the differential stress along the tip would likely
392 induce the activation of the numerous smaller scale existing fault zones in the region, an effect

393 which has been suggested to be in effect in a prior study of the deformation in Ghana (Nkodia
394 et al., 2020).

395 **5.2. Dual stress source**

396 The addition of an active mantle plume beneath the Cameroon region (see Fig. 2) represented
397 an additional stress source in the system. A comparison between models 1 and 4 revealed that
398 this new stress source produces little-to-no difference on the overall σ_{II} within the continent
399 after 1 Myr (compare the top rows in Figs. 3 vs 4), which suggests that, under present modelling
400 conditions, it might not represent a significant change in the system. However, mantle plumes
401 have been shown to induce both short- and large-scale stress changes both at the surface (e.g.,
402 Wang and Li, 2021; Burov and Guillou-Frottier, 2005) and at the base of the lithosphere (e.g.,
403 Gedamu et al., 2023). Thus, some additional component was required to obtain the expected
404 surface expression of the plume head. To that extent, we found that the addition of large-scale
405 continental weak zones significantly contributed towards the magnification of plume surface
406 expression, while oceanic based weak zones contributed towards the in-land propagation of
407 the stress along their tips (as previously detailed). Despite the plume head being constrained
408 to the Cameroon region (Fig. 4), it is affecting a much wider region, derived from the thermal
409 weakening of the base of the lithosphere (Fig. 4, right column), but also separating the
410 continent into a north and south domain by the CASZ. By having this precisely placed weak
411 boundary, the plume is able easily force the two regions to spread apart, increasing the surface
412 expression. The widespread weakening by the plume effectively increases the rheological
413 contrast to the east of the Accra region, which in turn leads to a more efficient localization of
414 stress along the cratonic boundary (e.g., Vauchez et al., 1998; Calignano et al., 2015). This effect
415 is also a function of plume head width, as after an identical model time, model 6 shows a much
416 wider surface expression when compared to Model 5 (see Fig. 4).

417 **5.3. Implications for the observed intraplate deformation and** 418 **seismicity**

419 Within the GG, the seismic activity is mostly concentrated in two major clusters (see Fig. 1),
420 namely the Accra and Cameroon clusters, as well as some scattered events within the Atlantic
421 plate. To evaluate our model results we use two approaches, namely by comparing: (1) the

422 modelled regions of high stress (as shown in Figs. 3 and 4) with the epicenter distribution within
 423 the GG (e.g., Mohammadigheymasi et al., 2023c); and (2) between the established deformation
 424 regimes for the Gulf of Guinea (e.g. Nkodia et al., 2022) and the modelled deformation regimes.

425 It has been shown in previous sections that in the scenarios in which the MAR is the only stress
 426 source, that most of the available stress localizes along the ocean-continent transition area,
 427 forming an envelope around the continental region (see Fig. 3). While this stress distribution
 428 could explain the seismic events within the oceanic plate, we argue that even though geological
 429 formations are various in this part of the African continent, they are not associated with strong
 430 enough rheological contrasts to explain the 2 seismic clusters.

431 Prior studies in region concerning the origin of the Accra cluster have proposed that these
 432 events could be related to either the western Saint Paul FZ and isostatic motions (e.g., Attoh
 433 et al., 2005), or to the presence of the Romanche FZ (e.g., Kutu, 2013). We tested the latter
 434 hypothesis through Model 2, where we assessed the efficiency of two large-scale weak zones
 435 at nucleating stress. As described in a previous section, the addition of the Romanche FZ was
 436 enough to induce a pronounced heterogeneity within this region of African continent which
 437 can be correlated with a pre-existing rheological boundary (namely a transition between two
 438 different continental blocks, corresponding to a suture zone, see Fig. 2-B). The location of this
 439 in-land stress propagation also broadly correlates to the location of active structures in the
 440 region (see Fig. 6-A and B). Given these results, we argue that it is likely that the Romanche FZ
 441 participates in inducing the reactivation of pre-existing structures around Accra, such as the
 442 Dahomeyides belt (see Fig. 1), which would explain the observed seismicity in this region. In
 443 addition, despite the presence or absence of the Romanche FZ, all models predict the
 444 occurrence of broadly compressive to transpressional stress conditions for ocean-continent
 445 transition within this region of the GOG (see Area 1 in Fig. 7-A and 7-B, as well as Supplementary
 446 Figure 2). Furthermore, as these localized transitions from compressive to transpressive are
 447 associated with concavities in the shoreline, it suggests that that the geometry of the
 448 continental margin can play a pivotal role in the distribution of the stress regimes in the area.
 449 Both observations are consistent with prior studies on the deformation within the GOG (e.g.,
 450 Nkodia et al., 2022). However, as in Model 2 no increased stress was observed within the
 451 Cameroon region, it suggests that the MAR and weak zones alone are insufficient to explain

452 both seismic clusters within the GG, reinforcing the idea of the role played by the OCT
453 geometry.

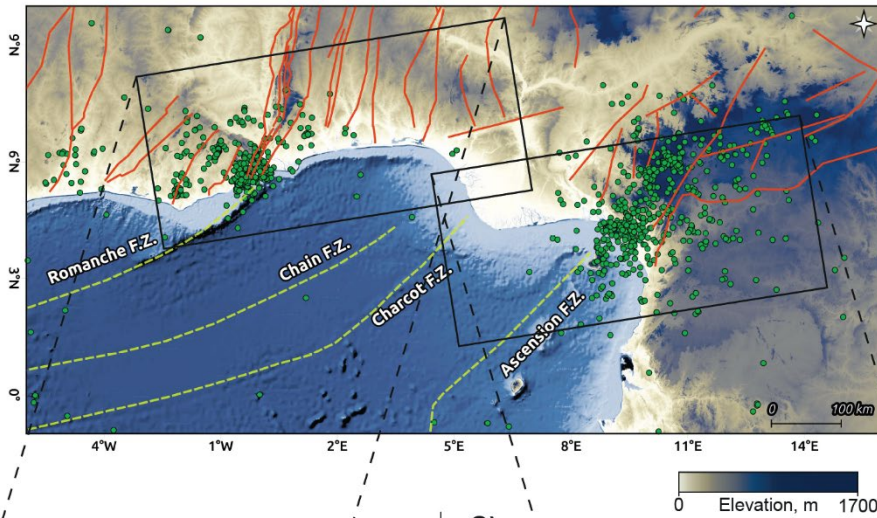
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455 One of the currently working hypotheses for the formation of the CVL is a complex interaction
456 between the CASZ and an upper mantle plume (e.g., Adams, 2022). Although yet unconfirmed,
457 the region beneath Cameroon has shown to have some of the signatures of a mantle plume,
458 such as a low shear wave velocity which may correspond with an anomalously high thermal
459 signature (e.g., Celli et al., 2020) and OIB-style isotopic signatures (e.g., Tanyileke et al., 1996).
460 Independently of the source of the CVL mantle plume, we aimed at investigating the influence
461 of a (local) thermal anomaly in the distribution of stress in the surrounding/region. Under these
462 constraints, both models 5 and 6 show the emergence of a large high stress zones centered
463 around the plume head, as well as a tendency towards extensive stress regimes within the
464 vicinity of the CASZ (see Fig. 4 and Fig. 7). This high stress region could participate in
465 redistributing stress in a wider region than that directly influenced by the current active centres
466 of the CVL within the Cameroon cluster (see Fig. 6-C), but the localized extension (see Fig. 7-B)
467 and thermal anomaly itself (see Fig. 5) could also help to explain the volcanic activity. One can
468 note that the present-day magmatic and volcanic activity is confined in between Mt Cameroon
469 area and Bioko island, with earthquakes as deep as 50-55 km (Ekodo et al 2023; DePlaen et al
470 2014). Despite located in this volcanic area, the alignment detected for some seismic events
471 with regional fault systems, attest of complex magmato-tectonic interaction, which could be
472 enhanced in this sector by processes of various origin, as suggested by our models.

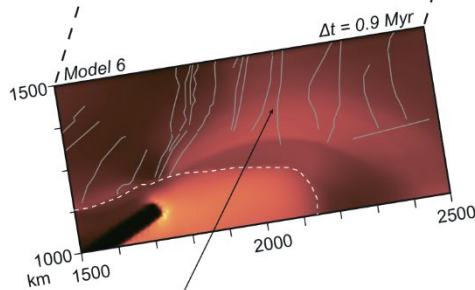
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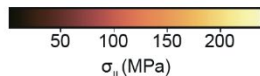
A)



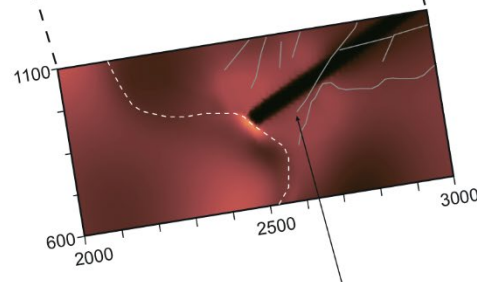
B)



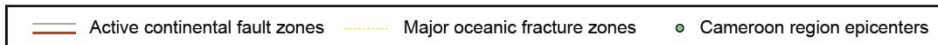
Active fault zones broadly constrained to higher stress regions



C)



Higher stress zone matches the cluster of epicenters observed



475

476 **Figure 6: A. Detail map of the target comparison region.** This map shows the target area to be
 477 used as a comparison with the model results, including both the Accra and Cameroon regions.
 478 **B. Model inset of the Accra Region.** Within this model region, it is shown that the active
 479 continental fault zones (marked in grey) fall at first order within the higher stress bands. **C.**
 480 **Model inset of the Cameroon region.** This model region, located just above the plume head
 481 (as detailed in Fig. 5), shows a large high stress region around the tip of the CASZ, which globally
 482 engulfs most of the epicenters found in and around Cameroon.

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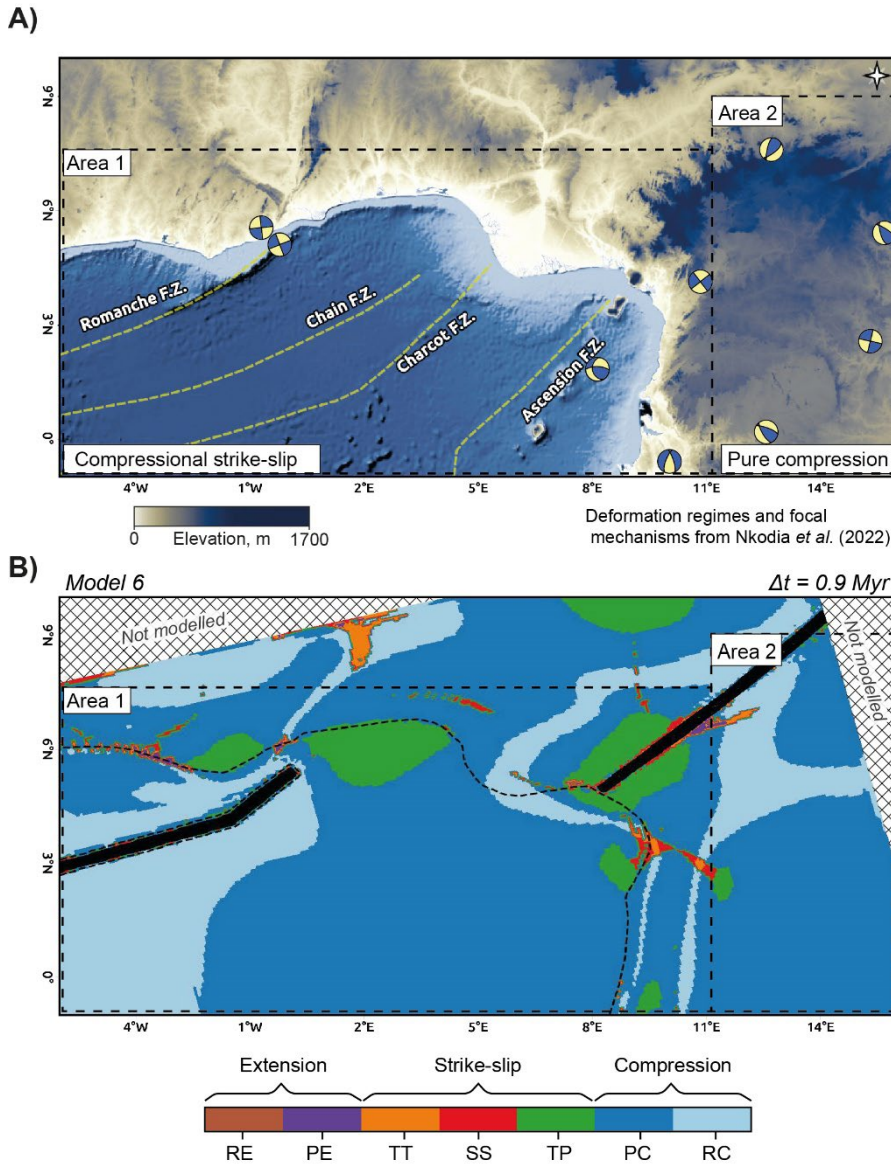


Figure 7: A. Detail map of the target comparison region. This map shows the target area that is compared with the model results and which includes stress regime domains and known focal mechanisms (both from Nkodia *et al.*, 2022). **B) Calculated stress regime conditions for Model 6.** The model shows a prevalence of mostly compressive stress conditions throughout, with highly localized strike-slip/extensional conditions (e.g., around the plume head and in the junction between the Romanche FZ and the continent). These conditions show a good first-order agreement with the known natural stress regimes for the GG, despite local mismatches due to minor (and not modelled) structures. Due to their inherent highly mobile nature, the internal regions of the weak zones show random distributions of stress regimes and, thus, have been blacked out. FZ: fracture zone; RE: radial extension; PE: pure extension; TT: transtension; SS: pure strike-slip; TP: transpression; PC: pure compression; and RC: radial compression.

486 **5.4. Model limitations**

487 The presented models provide insight on the drivers for intraplate deformation and seismicity
488 within the GG. Nevertheless, as with most numerical models, the observations made must take
489 into consideration the limitations of the models and the data quality and coverage.

490 A first limitation pertains to an uncertainty regarding the condition selection characterizing the
491 rheology of the upper mantle (e.g. Jain and Korenaga, 2020; King, 2016, and references
492 therein). Given this, while there is some liberty regarding the choice of rheological parameters
493 (such as the ones presented in Table A.1), it should be expected that slight variations of these
494 parameters could lead to different results. One possible example of these changes may be
495 effective viscosity of the ascending plume materials which could, for instance, rise faster or
496 slower through the mantle and/or form a narrower/wider plume head beneath the
497 lithosphere. Another parameter which may impact the results is the geothermal gradient.
498 Different thermal profiles would influence the deformation and stress pattern for the different
499 region and sub-regions modelled. Although our thermal profiles are approximated to a
500 continental geotherm, a colder continental lithosphere would result in a stronger prevalence
501 of compressive regimes (an additional model with the initial geometry of Model 2 but older
502 initial lithosphere was run, the results of which can be found in Supplementary Fig. 3).

503 Secondly, these models do not represent the full extent of the structural complexity of the
504 region as, for instance, we limit our weak zones to that of the CASZ and the Romanche SZ.
505 While there are other large transform faults in the region (e.g., the Ascension SZ), none has the
506 proximity and connection to the known in-land deformation as the Romanche. Thus, only the
507 latter was implemented. It is however plausible that the addition of the St. Paul SZ would
508 increase stresses west of Accra and add further stresses along the Dahomeyides belt.
509 Furthermore, the size of the continental blocks (as well as their thicknesses) used in this work
510 are of first-order and do not reflect the real complexity of the African continent. Nevertheless,
511 our objective was to provide a first order approach concerning the role of the major features
512 on stress localization

513 Additionally, the models presented in this work do not employ two-phase flow. Mantle plumes
514 are known to induce partial melting within the lithosphere (e.g., Manglik and Christensen,
515 2006), which not only introduces a liquid phase into the system but also weakens the local

516 rheology (e.g., Whattam and Stern, 2015). A weaker rheology implies that the stress needed to
 517 induce deformation is also lower which, therefore, would further imply that a lower stress
 518 accumulation would be seen above the plume site. Furthermore, the ascension of the plume
 519 in the model is controlled by the interplay between constant injection velocity at the base of
 520 the model and thermally lowered density, while in nature the positive buoyancy of the plume
 521 is the more important controlling factor. While our imposed injection velocity is in line with the
 522 estimated rates for mantle plumes, it is likely that the combined effect of the two factors may
 523 be resulting in an exaggerated velocity. In short, it is possible that our models present an
 524 overestimation of the accumulated stress above the plume site but are likely observing an
 525 underestimation of the deformation in this region.

526 Finally, one limitation of this work is the identification of possible seismogenic sources and
 527 processes in this region. This poor constraint stems from the natural complexity of intraplate
 528 settings (discussed in Talwani, 2014), as well as sparse instrumental coverage (Ahulu and
 529 Danuor, 2015). Thus, there may be a mismatch between the real distribution of seismicity (both
 530 in space and time) and the one shown in the present work, which could compromise our
 531 interpretation.

532 **6. Conclusions**

533 With this work, we investigated how the interplay between two major stress sources (i.e., the
 534 MAR and a possible Cameroon mantle plume) and large-scale weak zones can explain the
 535 intraplate deformation and seismic event distribution within the Gulf of Guinea region using a
 536 geodynamic modelling approach.

537 Our results (as shown in Figs. 3 and 4) suggest that, within the GG, stress distribution can be
 538 explained by the contribution of at least two stress sources: the MAR, which contributes to
 539 transfer stresses along the OCT and the mantle plume, which localizes stresses in the
 540 continental lithosphere in its direct vicinity.

541 They further suggest that large-scale weak zones are crucial for transferring stress in-land from
 542 the MAR, allowing for intraplate stress to accumulate in the continent, following prior field
 543 studies in the region such as Nkodia et al., (2022). This leads us to suggest that the maintained

544 seismicity in the sector of Accra results from a complex interaction between MAR, weak
545 fracture zones and OCT, while the seismicity observed in Cameroon, appears to be mainly
546 related to magmatism including active volcanism, and tectonics in the predominantly in the
547 continental domain.

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562 **8. Data Availability**

563 The raw data and GIS datasets for the compiled crustal structure, thickness, and average crustal
564 velocity in Cameroon, Ghana, and Nigeria can be accessed at
565 <https://github.com/SigProSeismology/Crustal-data-Gulf-of-Guinea.git>. The raw model output
566 can be found at <https://doi.org/10.17605/OSF.IO/6S4QD>.

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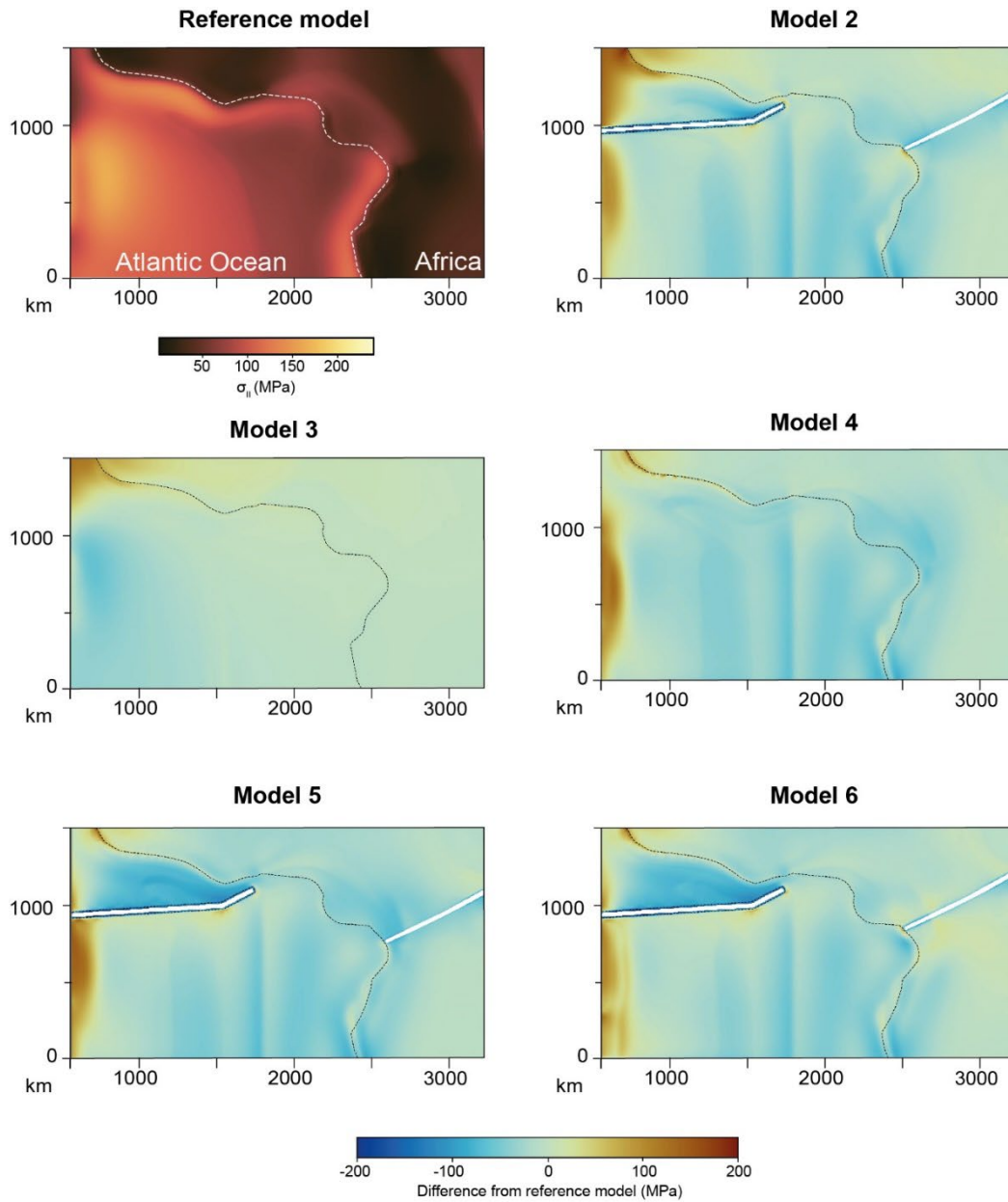
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805 **CRedit authorship contribution statement**

806 **Almeida, J.:** Model conceptualization and design, post-processing analysis, discussion
807 of the results, writing of the manuscript. **Mohammadigheymasi, H.:** Seismic data collection
808 and cataloging, discussion of the results, writing and review of the manuscript. **Neres, M.:**
809 Model conceptualization, discussion of the results, writing and review of the manuscript.
810 **Dumont, S.:** Model conceptualization, discussion of the results, writing and review of the
811 manuscript.

812 **Appendix I**

813



814 **Supplementary Figure 1 – Variation in deviatoric stress magnitude (σ'_{II}) between models.** This comparison
815 allows for the clear observation of zones of stress “shadow” to the north of the Romanche weak zone,
816 which result from its curved geometry.

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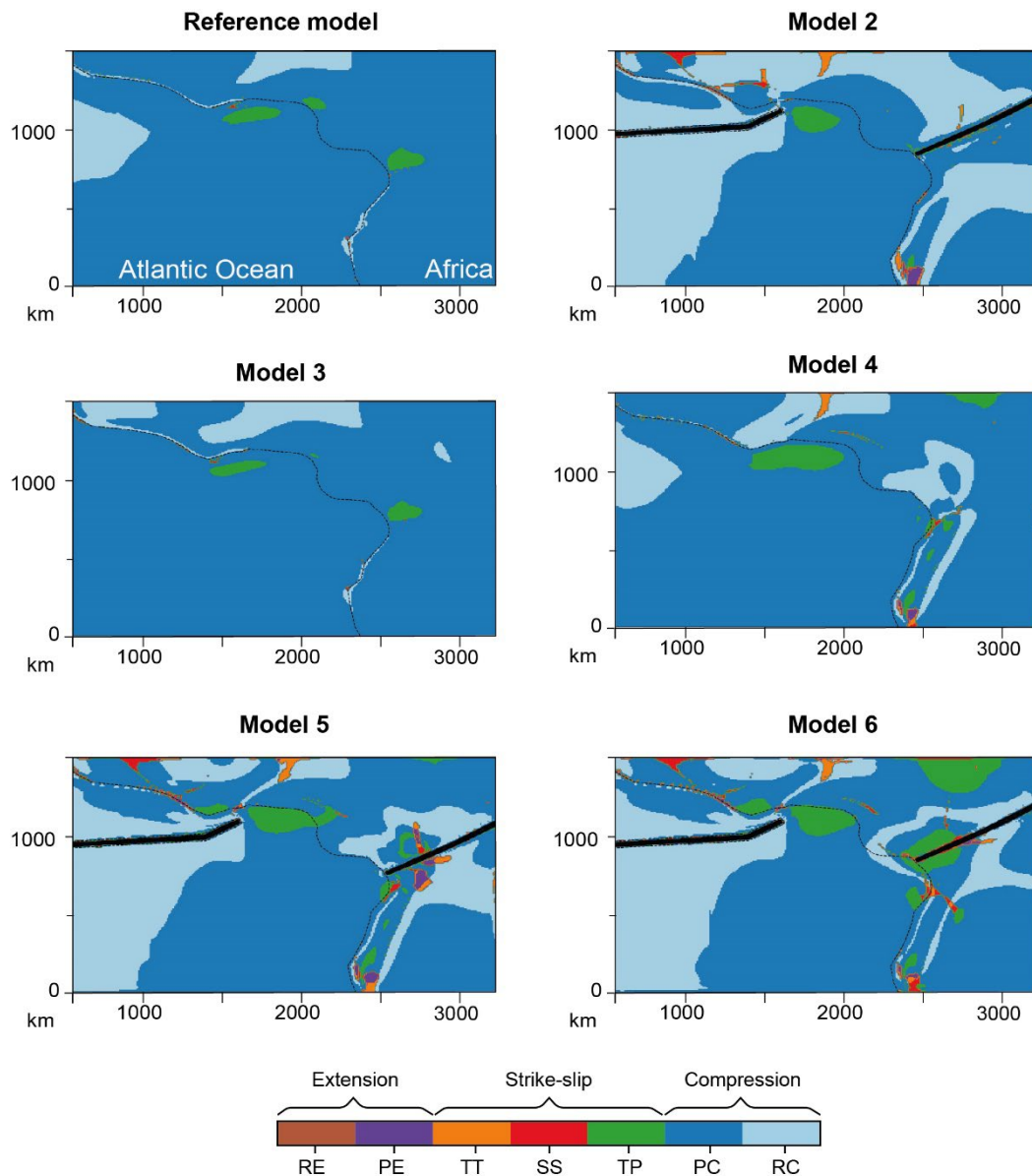
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Supplementary Figure 2 – Final stress regimes for the different models. All models broadly predict compressive conditions for both the oceanic and continental plates, which is in-line with the known stress regimes for the region (e.g. Nkodia et al., 2022). It is noticeable that inflexion of the African coastline produces localized regions of transpression, indicating a strong influence from its shape on the overall deformation. Lastly, models with an active plume (Models 4, 5 and 6) show clear changes in the surface deformation imposed by the plume head.

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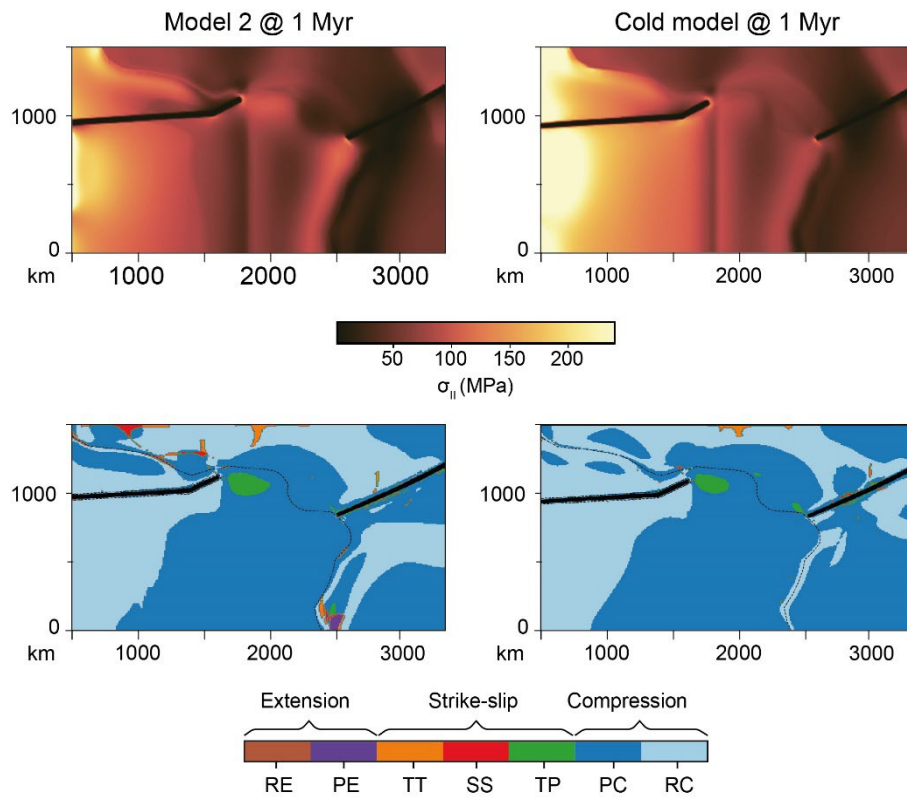
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Supplementary Figure 3 – Comparison between Model 2 and a colder lithosphere model. The model on the right was created by taking Model 2 and doubling the initial thermal age of every continental lithosphere block. This further approximates the thermal profile to a geotherm (e.g., Chapman et al., 1997). While the generic pattern distribution is identical (i.e., a prevalence of compressional regimes), the second model loses most non-compressional sites. The older (and colder) regions of the model, located south of the CASZ are under pure compression, with the small extensional site disappearing.

830