



Peer review status:

This is a non-peer-reviewed preprint submitted to EarthArXiv.

1 **Large earthquakes are more predictable than smaller ones**

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16 **Author contributions statement**

17 P.V.-A.: Conceptualization, Formal analysis, Investigation, Methodology, Project
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21 Software, Validation, Visualization, Writing – original draft and review & editing.

22

23 **Abstract**

24 Large earthquakes have been viewed as highly chaotic events regardless of their
25 magnitude, making their prediction intrinsically challenging. Here, we develop a
26 mathematical tool to incorporate multiscale physics, capable of describing both
27 deterministic and chaotic systems, to model earthquake rupture. Our findings
28 suggest that the chaotic behavior of seismic dynamics, that is, its sensitivity to
29 initial and boundary conditions, is inversely related to its magnitude. To validate
30 this hypothesis, we performed numerical simulations with heterogeneous fault
31 conditions. Our results indicate that large earthquakes, usually occurring in
32 regions with higher residual energy and lower b-value (i.e., the exponent of the
33 Gutenberg-Richter law), are less susceptible to perturbations. This suggests that
34 a higher variability in earthquake magnitudes (larger b-values) may be indicative
35 of structural complexity of the fault network and heterogeneous stress conditions.
36 To further validate our findings, we compare our theoretical predictions with real
37 seismicity in Southern California; specifically, the relationship between the b-
38 value and the fractal dimension of hypocenters with our model predictions finding
39 good agreement. The statistical similarities observed between the simulated and
40 real earthquakes support the hypothesis that large earthquakes may be less
41 chaotic than smaller ones; hence, more predictable.

42

43 **Keywords:** Earthquake predictability; Seismic rupture; Chaos theory; Residual
44 energy; b-value; HE-B method.

45

46

47 **1. Introduction**

48 Earthquakes are a persistent threat to human society, capable of causing
49 widespread devastation (e.g., [Kahandawa et al., 2018](#)). The rapid release of
50 accumulated tectonic stress can result in catastrophic natural disasters with
51 severe human and economic consequences ([Knopoff, 1958](#); [Vassiliou and](#)
52 [Kanamori, 1982](#); [Gudmundsson, 2014](#); [Aksoy et al., 2024](#); [Silverio-Murillo et al.,](#)
53 [2024](#)). To efficiently face seismic risk, a deeper understanding of seismicity is
54 needed. Particularly, a fundamental aspect of earthquake studies is the
55 examination of rupture processes along geological faults (e.g., [Christensen and](#)
56 [Beck, 1994](#); [Kintner et al., 2018](#); [Otarola et al., 2021](#); [Martínez-Lopez, 2023](#);
57 [Wang et al., 2023](#)), as these can induce notable changes in the soil's physical
58 characteristics, such as variations in ground velocity, acceleration and frequency
59 ([Colavitti et al., 2022](#); [Li et al., 2022](#); [Venegas-Aravena, 2024](#)). Evidence from
60 various studies points to the possibility that seismic rupture processes may exhibit
61 the hallmarks of chaotic systems, suggesting a complex and unpredictable nature
62 of these events. Some perspectives on earthquake generation are rooted in
63 simplified spring-block models which exhibit these chaotic dynamics (e.g., [Huang](#)
64 [and Turcotte, 1990](#); [Gualandi et al., 2023](#)). This is reflected in computational
65 simulations where small variations in the initial conditions generate completely
66 different rupturing scenarios (e.g., [Erickson et al., 2011](#)). That is, causing no
67 correlation between a priori and a posteriori parameter ([Venegas-Aravena et al.,](#)
68 [2024](#)). Despite this complexity, a consistent finding from these simulations is the
69 emergence of a single dominant parameter: residual energy. The importance of
70 the accumulated/residual strain is also coherent with recent results in the
71 modeling of paleoseismic recordings ([Salditch et al., 2020](#)). This parameter,

72 which defines zones where ruptures are prone to occur (Noda et al., 2022),
73 appears to exert a controlling influence on both the spatial extent and temporal
74 evolution of ruptures (Venegas-Aravena, 2023; Venegas-Aravena et al., 2024).
75 Its value is dependent on both the available energy, which is determined by the
76 initial stresses, and the fracture energy, which is associated with the energy
77 required to continue propagating the rupture (Noda et al., 2022). Therefore, when
78 the rupture front approaches a zone with negative (positive) stress release rate,
79 more (less) energy is consumed in generating the rupture, causing the rupture to
80 arrest (continue propagating). Other formalisms associated with friction have also
81 found that rupture arrest can be related to stresses and fracture energy (Barras
82 et al., 2023).

83 Residual energy has been linked to a parameter called thermodynamic fractal
84 dimension D (Venegas-Aravena and Cordaro, 2023a). This quantity is useful for
85 characterizing fault distribution (Zou and Fialko, 2024) and the spatial distribution
86 of global seismicity (Perinelli et al., 2024). For instance, it has been observed that
87 the fractal dimension decreases prior to a major earthquake, suggesting a
88 transition from a more diffuse, three-dimensional seismicity distribution to a more
89 localized, planar distribution along the fault (Murase, 2004; Wyss et al., 2004;
90 Iaccarino and Picozzi, 2023). Other studies have interpreted this decrease in D
91 as an indicator of an impending larger rupture due to the increase of shear
92 stresses (e.g., Ito and Kaneko, 2023; Venegas-Aravena and Cordaro, 2023a).
93 Furthermore, it has been linked to the b-value, a parameter describing
94 earthquake frequency (Venegas-Aravena and Cordaro, 2023b). Given the
95 proportional relationship between D and the b-value, a decrease in D is also
96 associated with a decrease in the b-value prior to large earthquakes.

97 Given the link between the parameter D and properties associated with chaotic
98 systems, as suggested by lower D values in less chaotic systems (Venegas-
99 Aravena and Cordaro 2024), the b-value is anticipated to provide insights into the
100 chaotic states of faults. To explore this connection, Section 2 delves into the
101 fundamental principles of multiscale thermodynamics applied to faults. Section 3
102 presents various simulations of heterogeneous ruptures, facilitating the
103 interpretation of parameters such as D and the b-value within the framework of
104 multiscale thermodynamics and chaotic systems. In Section 4, we apply these
105 concepts to a real fault system, specifically in Southern California, to support our
106 theoretical and numerical results. A discussion and conclusion are presented in
107 Sections 5 and 6 respectively.

108

109 **2. Theoretical background: multiscale thermodynamics**

110 As earthquakes are essentially multi-scale events that may exhibit chaotic
111 behavior, a physical framework is required to fully understand their dynamics. In
112 this regard, Venegas-Aravena and Cordaro (2024) have developed a quantitative
113 relationship linking the sum of the Lyapunov exponents Λ , to the thermodynamic
114 fractal dimension D , expressed as:

$$115 \quad \Lambda \sim e^{\frac{(D_E - D - 1)}{k_V}} \quad (1)$$

116 The Euclidean dimension is denoted by D_E , while k_V is a constant associated with
117 the system's scale. D is a parameter that characterizes the distribution of systems
118 exhibiting power-law behavior. For simplicity, k_V is assigned a value of 1. This
119 relationship can be observed in Figure 1a, where low (large) values of Λ are

120 associated with low (large) values of D . A higher value of Λ indicates that the
121 system is more susceptible to the influence of small changes in initial conditions
122 ([Ruelle, 1983](#); [Tabor, 1989](#)), whereas a more negative value of Λ suggests that
123 the system is less sensitive to these initial conditions, which could be considered
124 as being more regular. The sign of Λ provides an indication of whether a system
125 is non-reversible/dissipative (negative sum) or conservative (e.g., [Hoover and
126 Posch, 1994](#)). Given that the brittle crust is a system characterized by the
127 dissipation of stored energy, the sum of Λ is a more relevant metric than the
128 largest Lyapunov exponent, often used to determine the chaotic nature of a
129 system. Consequently, lower values of D correspond to less chaotic systems, i.e.,
130 less sensitive to initial conditions. In the case of earthquakes, D can be related to
131 the magnitude of seismic events through the equation ([Venegas-Aravena and
132 Cordaro 2023a](#)):

$$133 \quad M_W \sim \log_{10} e^{-\alpha(D)} \quad (2)$$

134 Where $\alpha(D) = pD/k_V$ and $p = 3/(5 - D)$. From Equation 1, D can be written in
135 terms of the sum of the Lyapunov exponent as $D = (D_E - 1) - k_V \ln \Lambda$. By
136 substituting this Equation into Equation 2, we can establish a direct relationship
137 between the magnitude and chaotic systems as:

$$138 \quad M_W = M_W(\Lambda) \quad (3)$$

139 This relationship is graphically depicted in the color-coded maps presented in
140 Figures 1a, 1b, and 1c. In these maps, red tones correspond to earthquakes of
141 greater magnitude, while blue tones represent smaller earthquakes. Notably, in
142 Figure 1a, large earthquakes are correlated with lower values of both D and λ .
143 This finding implies that larger seismic events exhibit a reduced sensitivity to

144 initial conditions. In contrast, smaller earthquakes (represented by blue hues in
145 Figure 1a) are associated with a higher degree of chaos, suggesting that these
146 events originate in a more chaotic environment where even small perturbations
147 can lead to seismic activity of varying scales, from small to intermediate.

148 This phenomenon is coherent with several observations suggesting the strong
149 sensitivity of small seismicity to stress perturbations, e.g., tides and hydrological
150 modulations (Rubinstein et al., 2008; Petrelis et al., 2021), which are not reported
151 for major events (Vidale et al., 1998; Kossobokov and Panza, 2020). To delve
152 deeper into this phenomenon, it is imperative to examine the energy conditions
153 within the fault, specifically the concept of residual energy, E^{res} (Noda et al.,
154 2021). This energy parameter serves as a criterion for the initiation of ruptures,
155 indicating that a positive E^{res} value signifies a greater propensity for a fault to
156 generate ruptures, while negative values diminish this likelihood. Mathematically,
157 this energy can be expressed as:

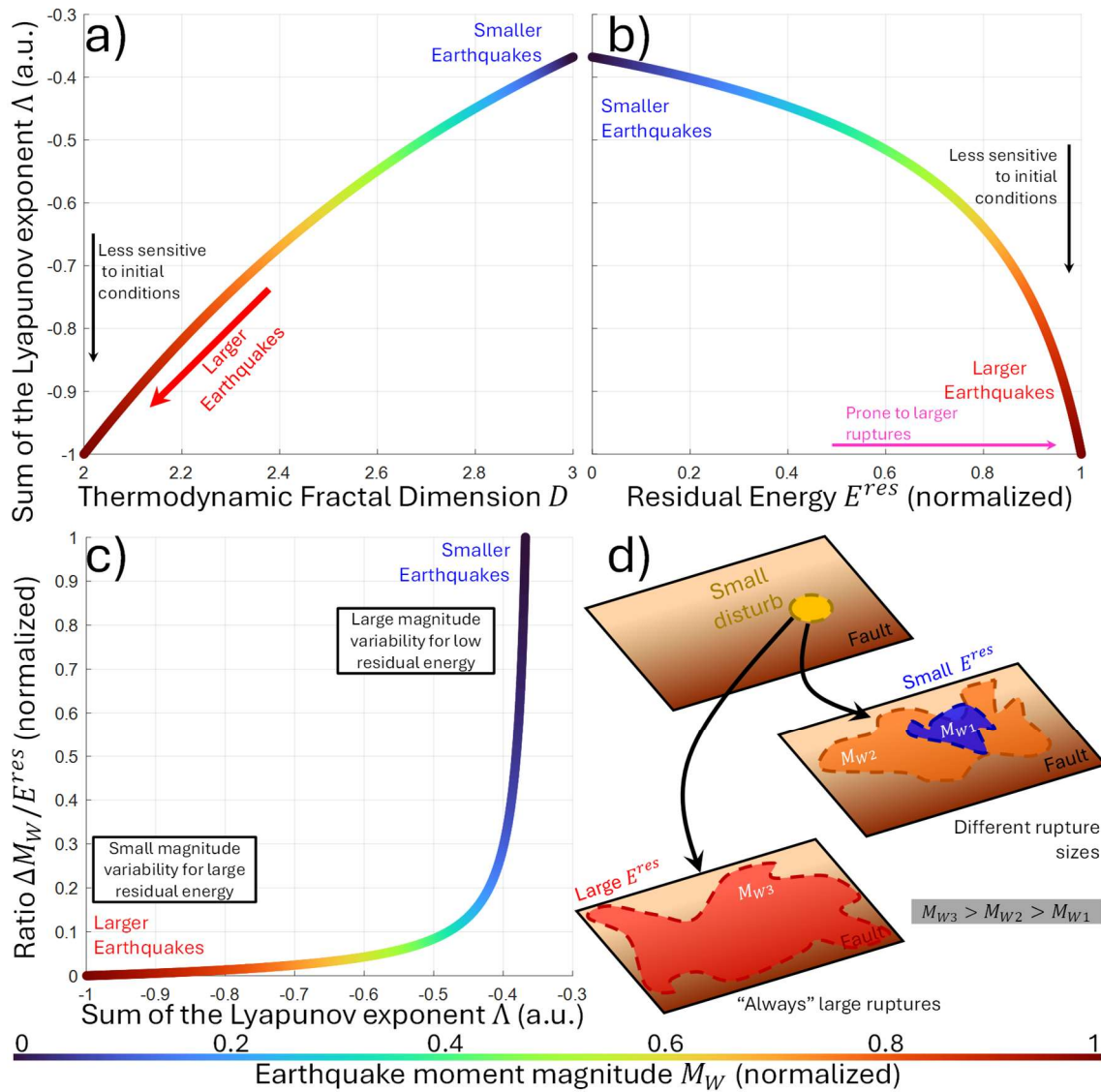
$$158 \quad E^{res} = \Delta W_0 - G_C \quad (4)$$

159 Where ΔW_0 represents the available energy, which can be correlated with the
160 elastic energy stored within the system, and G_C denotes the fracture energy,
161 characterizing the resistance to rupture propagation. This equation can be
162 expressed in terms of D as follows:

$$163 \quad E^{res} \sim e^{-\frac{D}{2k_V}} - d_0 10^D \quad (5)$$

164 Where d_0 is constant. Figure 1b illustrates the relationship between E^{res} , Λ , and
165 earthquake magnitude (color-coded map). Higher E^{res} values correlate with lower
166 Λ values, suggesting that regions more prone to rupture are also less sensitive to

167 initial conditions. Given that these regions are associated with large earthquakes
 168 (red colors), it is proposed that areas with high residual energy have higher
 169 chances to host a major event as a response to stress perturbations.
 170



171

172 **Figure 1: a) Equation 1 reveals a relationship between the sum of Lyapunov**
 173 **exponents Λ and the thermodynamic fractal dimension D . Systems with low**
 174 **sensitivity to initial conditions (highly negative Λ values) correspond to low**
 175 **D values. Colors indicate event magnitudes as calculated by Equation 2.**
 176 **Large events (red hues) are associated with low D and low Λ . b) Equation 5**

177 **relates Λ to residual energy (E^{res}). Higher E^{res} values correlate with a higher**
178 **probability of large earthquakes, which in turn are linked to lower chaos**
179 **and larger events. c) The plot of magnitude changes for a given E^{res} versus**
180 **Λ shows that small (large) earthquakes exhibit greater (lesser) magnitude**
181 **variability for low (high) E^{res} , as indicated by blue (red) hues. d) A schematic**
182 **illustrates how perturbations can trigger small-to-medium or large**
183 **earthquakes depending on E^{res} .**

184

185 To visualize this, Figure 1c shows the variation in magnitude ΔM_W relative to
186 residual energy. This indicates that the variation in M_W is small (large) when E^{res}
187 is high (low). This supports the notion that there is a more restricted range of
188 possible earthquakes when the residual energy in the fault is higher. Figure 1d
189 schematically depicts this concept: a fault with a small E^{res} can generate
190 earthquakes of magnitudes M_{W1} (blue area) and M_{W2} (orange area), whereas in
191 the case of a large E^{res} , only earthquakes of magnitude M_{W3} (red area) can be
192 generated, which is larger than both M_{W1} and M_{W2} .

193

194 **3. Simulations**

195 **3.1 Heterogeneous Energy-Based method**

196 The heterogeneous energy-based method (HE-Bm) posits that seismic rupture
197 propagation is governed by the heterogeneous distribution of residual energy
198 (Venegas-Aravena, 2023). This model suggests that rupture velocity and slip
199 magnitude at each point on the fault are directly correlated with the residual

200 energy. Consequently, regions with high residual energy are more prone to
201 experiencing large slip u_f and high rupture velocities v_r , which can potentially
202 lead to larger magnitude earthquakes. Thus, the relation between slip and
203 residual energy is $u_f \propto E^{res}$.

204 According to the framework of HE-Bm, E^{res} can be linked to the distribution of
205 interseismic coupling on a fault through the concept of available energy, while
206 fault geometry is related to residual energy via fracture energy (Venegas-Aravena
207 and Cordaro, 2023a). The two-dimensional fractal dimension D of natural
208 fractures have been determined to be ~ 2.3 (Huang et al., 1992). Consequently,
209 due to the proportional relationship between the fracture energy G_C and the
210 geometric variations of the fracture (e.g. Xie, 1993), it is expected that G_C will also
211 possess a fractal dimension of 2.3.

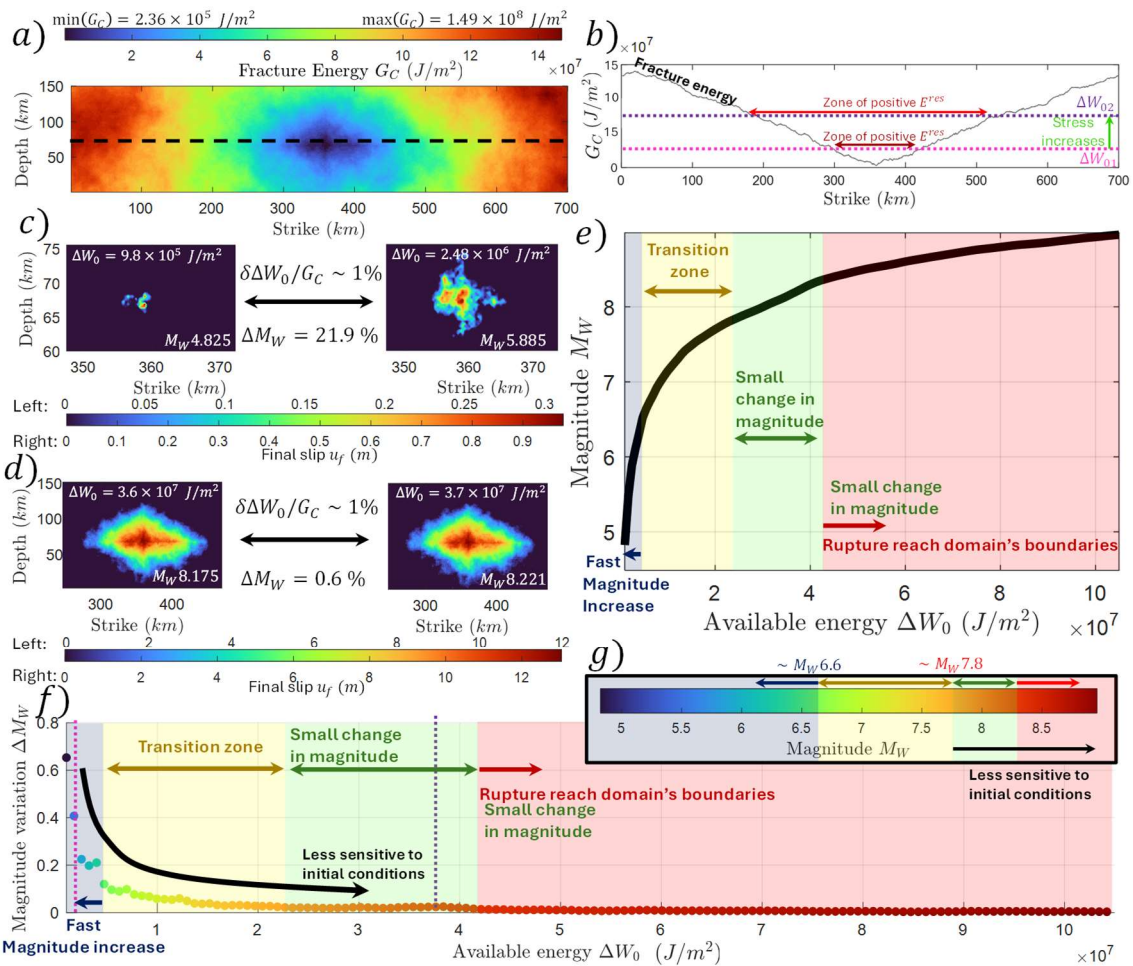
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213 **3.2 Ruptures in a single distribution of G_C**

214 Figure 1a presents an exemplar G_C distribution exhibiting a fractal dimension of
215 2.3. The strike and depth are 700 km and 150 km respectively. The spacing is
216 349.5 m for the strike and 371.6 m for the depth. This distribution was constructed
217 through the interpolation of random values, employing the methodology outlined
218 by Chen and Yang (2016). Given the established inverse correlation between
219 elevated G_C values and rupture size (Renou et al., 2022), attributed to the self-
220 arresting nature of ruptures induced by energy depletion, the central region
221 (depicted in blue in Figure 1a) was intentionally constrained to be three orders of
222 magnitude less than the peripheral regions (rendered in red). Note that the arrest
223 of ruptures due to geometric changes (fracture energy) has also been observed

224 in real faults (e.g., Padilla et al., 2024). Consequently, ruptures are invariably
 225 localized within the lower G_C value domains (represented by the blue hues in
 226 Figure 1a). As coupling seems to be related to stress (Wallace et al., 2012), this
 227 implies that for a given level of stress on a fault, it is the fault roughness that
 228 primarily determines residual energy. Smoother faults exhibit lower fracture
 229 energy, resulting in reduced resistance to rupture initiation and, consequently,
 230 larger rupture events. Conversely, rougher faults present higher fracture energy,
 231 limiting residual energy and thus constraining rupture size.

232



233

234 **Figure 2: a) Example of fracture energy distribution with $D=2.3$, where the**
 235 **central region has low values, and the edges have high values. b) Fracture**

236 energy profile corresponding to the segmented black line in a). Magenta
237 and purple segmented lines indicate two levels of available energy ΔW_0 .
238 Dark red and red double arrows indicate the size of positive E^{res} , potentially
239 corresponding to rupture size. c) and d) shows the final slip distributions
240 for conditions of low and high available energy, respectively. c) reveals
241 larger changes in magnitude than those shown in d). e) The relationship
242 between moment magnitude and ΔW_0 is shown. The gray region highlights
243 a rapid increase in magnitude with increasing ΔW_0 , while the yellow and
244 green zones show a decreasing rate of increase. The red zone indicates
245 ruptures that reach the fault edges. f) The variation of magnitude with
246 available energy for different values of the parameter ΔW_0 is shown. Lower
247 values of ΔW_0 result in larger changes in magnitude. g) The color map used
248 in f), which indicates the sensitivity of earthquakes to initial conditions:
249 blue for smaller earthquakes ($M_w \sim 6.6$), yellow for a transition region (M_w
250 between 6.6 and 7.8), and red for larger earthquakes ($M_w > 7.8$) that are less
251 sensitive.

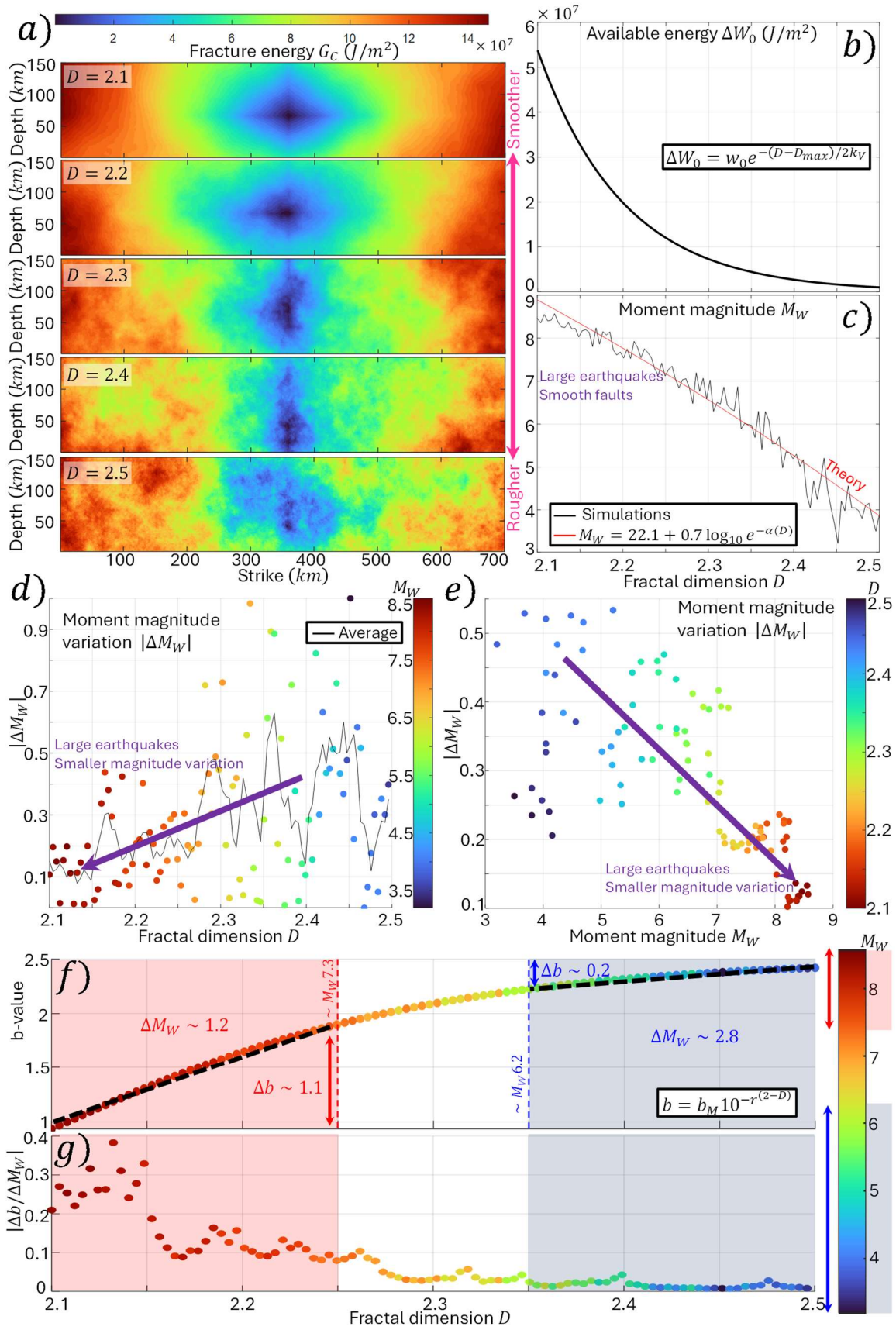
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253 Equivalently, for a given fracture energy distribution, the residual one will be
254 determined by the amount of available energy. This example is shown in Figure
255 1b. The black curve corresponds to a trace indicated by the black segmented line
256 in Figure 1a. The minimum value is $G_C = 2.36 \times 10^5 \text{ J/m}^2$, which is found
257 approximately at the midpoint of the fault (strike of 350 km). These values of G_C
258 tend to increase towards the strike equal to zero km and equal to 700 km. The
259 segmented magenta and purple lines represent two uniform distributions of
260 available energy, ΔW_{01} and ΔW_{02} , respectively. In this case, ΔW_{02} is greater than

261 ΔW_{01} , indicating that the first case has a smaller amount of accumulated stress
262 on the fault than the second case. The dark red double arrow would indicate the
263 zone with positive residual energy given the level of ΔW_{01} , which is equivalent to
264 a potential rupture zone. The red double arrow indicates the zone of positive
265 residual energy given a higher accumulated stress (given by ΔW_{02}). This zone is
266 wider than the region marked by the dark red arrow, highlighting the presence of
267 larger ruptures promoted by high stress values throughout the crustal volume.
268 The increase of available energy also translates into changes in earthquake
269 magnitudes. For instance, Figure 2c illustrates two ruptures initiated with similar
270 available energy ($\sim 10^6 J/m^2$), representing a one percent variation relative to the
271 maximum fracture energy. This excess in available energy defines the positive
272 residual energy area (rupture area A), which can be related to the seismic
273 moment M_0 through the empirical relationship $M_0 = \mu C_2 A^{3/2}$, where μ is the
274 shear modulus with a value of $4 \times 10^{10} Pa$ and C_2 is a dimensionless constant
275 equal to 3.8×10^{-5} (Leonard, 2010). The small variation available energy results
276 in a $\sim 22\%$ increase in earthquake magnitude (from $M_W \sim 4.8$ to $M_W \sim 5.9$).
277 Conversely, when the available energy is higher ($\sim 3.6 \times 10^7 J/m^2$), a similar 1%
278 increase produces earthquakes with nearly identical magnitudes ($M_W \sim 8.2$
279 representing a variation smaller than 1%, as shown in Figure 2d), suggesting that
280 faults with higher residual energy yield similar magnitude earthquakes.

281 The dependence of magnitude on available energy is depicted in Figure 2e. This
282 figure shows 140 simulations with different values of ΔW_0 . A significant increase
283 in magnitude is observed for low available energy values, indicated by the blue
284 region. In this region, a small increase in ΔW_0 (less than $10 MJ/m^2$) can elevate
285 an earthquake from magnitudes less than 5 to approximately $M_W 6.6$. The yellow

286 and green regions show a less pronounced increase in magnitude compared to
287 the blue region. The red range represents events where the ruptures approach
288 the fault boundaries. The magnitude change in this region appears unaffected by
289 fault boundary influences. Figure 2f quantifies these variations, revealing that the
290 blue region experiences ΔM_W values close to 0.7, while the green and red regions
291 show negligible changes. The color map in Figure 2g corroborates these findings,
292 with earthquakes smaller than $M_W 6.6$ predominantly falling within the blue region
293 and larger earthquakes (M_W larger than $M_W 7.8$) exhibiting minimal sensitivity to
294 variations in ΔM_W .



296 **Figure 3: a) Fracture energy of a fault with different fractal dimension (D)**
297 **values. The distribution becomes rougher as D increases. b) There is an**
298 **exponential relationship between available energy and D . As D decreases,**
299 **the available energy increases exponentially. c) The relationship between**
300 **moment magnitude and D is shown. The simulated data is represented by**
301 **the black curve, and the theoretical prediction by [Venegas-Aravena and](#)**
302 **[Cordaro \(2023a\)](#) is shown in red. d) and e) show the relationship between**
303 **fractal dimension, moment magnitude, and magnitude variation. The purple**
304 **arrow indicates that in both figures, low values of D are associated with**
305 **high-magnitude earthquakes and small magnitude variations. That is, a**
306 **small change in D , when ΔW_0 values are high, almost always generates**
307 **similar large earthquakes. When ΔW_0 values are low, there is a greater**
308 **variation in magnitude. f) Relationship between b-value and D . There is a**
309 **greater decrease in the b-value when there are larger earthquakes. g)**
310 **Variation of the b-value with changes in M_W . This variation is greater when**
311 **D is lower.**

312

313 **3.3 Ruptures with different G_C**

314 While natural faults can be characterized by a fractal dimension of approximately
315 2.2 (e.g. [Kagan, 1991](#)), variations in this value are possible. To investigate the
316 impact of fractal dimensions on fracture energy, 100 simulations were conducted
317 with fractal dimensions ranging from 2.1 to 2.5. Figure 3a illustrates examples of
318 fault geometries with fractal dimensions of 2.1, 2.2, 2.3, 2.4, and 2.5, respectively,
319 where maximum and minimum G_C values are consistent with Figure 2a. In these

320 simulations, lower G_C values are maintained at the fault center and higher values
321 at the edges. As shown in Figure 3a, the distribution of G_C becomes smoother as
322 the fractal dimension approaches 2.1.

323 As suggested by [Venegas-Aravena and Cordaro \(2023a\)](#), ΔW_0 is inversely
324 related to D , with the specific relationship being $\Delta W_0 \sim e^{-D/2k_V}$. Therefore, any
325 change in G_C must be accompanied by a corresponding change in ΔW_0 . Figure
326 3b visualizes this relationship using parameter values of $w_0 = 9.8436 \times 10^5 \text{ J/m}^2$,
327 $D_{max} = 2.5$ and $k_V = 0.05$. These values yield a range of ΔW_0 consistent with the
328 previous section, ensuring that ΔW_0 is sufficiently large to allow for rupture
329 initiation but not so large as to be influenced by domain boundaries. The figure
330 clearly demonstrates that lower values of D are associated with higher values of
331 ΔW_0 , indicating smoother spatial distributions of G_C . The magnitude of these
332 ruptures also varies as a function of D . In Figure 3c, the red curve represents the
333 theoretical relationship proposed by [Venegas-Aravena and Cordaro \(2023a\)](#),
334 given by $M_W \sim \log_{10} e^{-\alpha}$, where $\alpha = pD/k_V$ and $p = 3/(5 - D)$. The observed
335 magnitudes, depicted by the black curve, align well with the theoretical values.
336 However, a higher variability in magnitude $|\Delta M_W|$ is observed for larger values of
337 D (greater than 2.4), while lower values of D (less than 2.3) exhibit lower
338 variability. This variation is visualized in Figure 3d, where the color map indicates
339 magnitude. The black curve in Figure 3d represents a 5-point moving average of
340 $|\Delta M_W|$, with the purple arrow highlighting the trend towards lower magnitude
341 variability for smaller values of D . Figure 3e explicitly shows the relationship
342 between M_W and its average variability (in this case a 10-point moving average),
343 with the color map indicating D values. As the purple arrow suggests, there is an
344 inverse correlation between M_W and its average variability, where earthquakes

345 with magnitudes less than $M_W \sim 5$ can exhibit magnitude differences greater than
346 $0.5M_W$. In contrast, for earthquakes with magnitudes greater than $M_W \sim 8$, this
347 variability decreases to approximately $0.1M_W$.

348

349 **3.4 Chaos and b-value**

350 Both laboratory and field studies have shown a negative correlation between the
351 b-value, which quantifies the frequency of earthquakes of different magnitudes in
352 each region and increasing stress levels. This leads to a decrease in the b-value
353 and may be associated with large magnitude earthquakes (Scholz, 2015; Dong
354 et al., 2022). Studies have established a theoretical link between the b-value and
355 fractal dimension, suggesting that lower b-values correspond to lower D values
356 and vice versa (Aki, 1981; Venegas-Aravena and Cordaro, 2023b). Specifically,
357 this relationship is expressed as $b - value = b_M 10^{-r(2-D)}$, where b_M is 2.5 and r
358 is a constant between 10^3 and 10^4 (in this study, r is set to medium value 5000
359 for simulations). This law is illustrated in Figure 3f. The color map indicates
360 earthquake magnitudes, with blue transparency representing events from
361 $M_W \sim 3.4$ to $M_W \sim 6.2$, corresponding to a D variation of 1.5. The magnitude
362 variation within this zone is $\Delta M_W \sim 2.8M_W$, while the b-value decrease is $\Delta b \sim 0.2$.
363 Red transparency indicates the same variation of D , but with earthquakes ranging
364 from $M_W \sim 7.3$ to $M_W \sim 8.5$, corresponding to a $\Delta M_W \sim 1.2$ and a $\Delta b \sim 1.1$
365 decrease. The reduction in the rate of change of the b-value with respect to
366 magnitude is clearly displayed in Figure 3g. The blue transparent area
367 emphasizes a region where the absolute value of the b-value remains relatively
368 constant, even as the magnitude of earthquakes fluctuates. It is important to note

369 that a 5-point moving average was applied to the data. This suggests that the b-
370 value is less sensitive to changes when the fault system predominantly generates
371 smaller earthquakes. In contrast, the red transparent area reveals a more
372 pronounced relationship between b-value and magnitude, with the b-value
373 fluctuating more rapidly as the magnitude increases. These results imply that a
374 more abrupt decrease in the b-value is associated with smaller changes in
375 magnitude but, likely, also with changes in fault conditions leading to less chaotic
376 behavior.

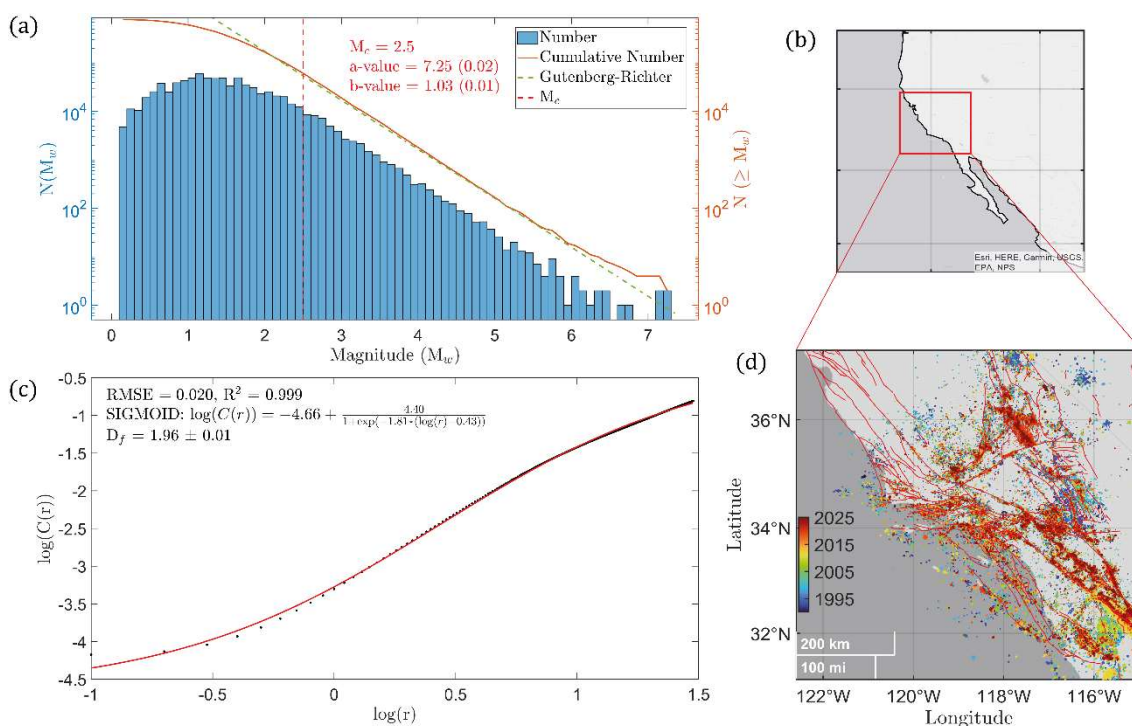
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378 **4. A reality check: comparison with seismicity in Southern California**

379 We have already compared our theoretical predictions with the output of dynamic
380 simulations of earthquakes; here we make a reality check with the statistical
381 properties of seismic catalogs. Specifically, we validate the compatibility of the
382 relationship $b - value = b_M 10^{-r(2-D)}$, between the b-value of the Gutenberg-
383 Richter law and the fractal dimension of faulting. Since it is not possible to directly
384 investigate the fractal properties of faults, we calculate the fractal dimension of
385 hypocenters (hereafter referred as D), which are expected to be distributed within
386 a subset of the fracture network; hence, D is equal or lower than the value for the
387 fault system. Nevertheless, even with different coefficients (b_M and r), the
388 empirical law of the b-value follows the same trend because seismic events are
389 supposed to occur throughout the whole investigated crustal volumes. Thus, we
390 specify that we are not interested either in assessing the true fractal dimension
391 of the networks of faults hosting seismicity (an accurate estimation is not feasible)
392 nor the true fractal dimension of seismic events in their long-term behavior (which

393 would require much longer catalogs than available nowadays and accurate
 394 declustering). Here, our goal is just the observational validation of the
 395 mathematical relationship $b - value = b_M 10^{-r(2-D)}$. It requires a high-quality
 396 relocated seismic catalog produced by a roughly uniformly distributed network of
 397 seismic stations (i.e., uniform completeness magnitude). Both the background
 398 and triggered components are considered, otherwise the spatial variations of the
 399 b-value and fractal dimension vanish preventing any investigation of their
 400 relationship with the available catalogs.

401



402

403 **Figure 4: Seismicity in Southern California (SCEC Catalog, 1990-2025,**
 404 **latitude 31°-37° N, longitude 115°-122° W, depth lower than 30 km). (a)**
 405 **Frequency-magnitude distribution of seismicity, the completeness**
 406 **magnitude is highlighted by the red vertical dashed line. In (b) and (d) is the**
 407 **map with the spatial distribution of seismicity. (c) log-log representation of**

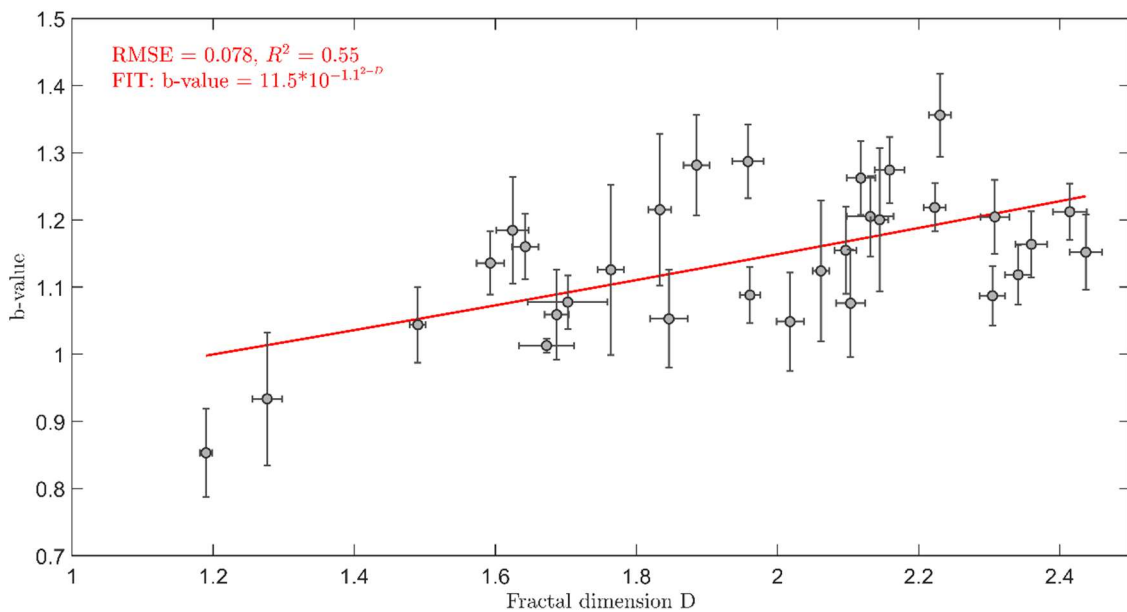
408 **the correlation function vs the threshold radial distance (km). The plot**
409 **shows a range of scales where the curve is well approximated by a line, i.e.,**
410 **hypocenters follow a fractal distribution in space.**

411

412 We analyze the shallow crustal seismicity (depth lower than 30 km) in Southern
413 California between 1/1/1990 and 20/1/2025 listed in the Waveform Relocated
414 Earthquake Catalog for Southern California ([Hauksson et al., 2012](#)). A visual
415 representation of seismicity considered in this study is given in Figure 4a, b and
416 d. The catalog is divided into several squared regions. The number and selection
417 procedure used to define the structure of the subsets do not significantly affect
418 the final output provided that the fractal probability and the b-value be
419 investigated only for regions with at least 500 events to get stable and reliable
420 results. Only events above the completeness magnitude are considered, with M_c
421 = 2.5. It is estimated according to the EMR method ([Woessner and Wiemer,](#)
422 [2005](#)). Since short-term aftershocks incompleteness (STAI) after the occurrence
423 of major events is still present even if a great part of the catalog contains reliable
424 information, the b-value is calculated using the b-positive algorithm ([van der Elst,](#)
425 [2021](#)) with the b-more positive correction ([Lippiello and Petrillo, 2024](#)) to avoid
426 bias. The uncertainty of the b-value is found using bootstrapping over 100
427 simulations with acceptance probability equal to 0.5. The fractal dimension of the
428 hypocenters is measured using the Grassberger and Procaccia algorithm
429 ([Grassberger and Procaccia, 1983](#)). Here, we introduce a new method to remove
430 possible sources of bias in its estimation due to the arbitrary selection of the lower
431 and upper cut-offs for the linear region in the log-log plot. The curve of the
432 correlation function $C(r)$ as a function of the threshold radius r is fitted using the

433 sigmoid function $y = y_0 + k/(1 + e^{-\alpha x})$, where $y = \log(C(r))$ and $x = \log(r)$,
 434 while k , α and y_0 are left as free parameters, so that the fractal dimension
 435 (i.e., the derivative of the sigmoid in its symmetry saddle point) is given by $D =$
 436 $k\alpha$. The uncertainty is calculated by propagating the fit errors of k and α . The
 437 estimation of the fractal dimension of hypocenters for the whole catalog is in
 438 Figure 4c. The analysis performed over a wide range of possible grids (both
 439 uniformly spaced and nested according to the number of seismic events within
 440 them) shows that the b-value and the fractal dimension of hypocenters are
 441 positively correlated; moreover, the curve $b - value = b_M 10^{-r^{(2-D)}}$ provides a
 442 good fit of the relationship between b and D, in agreement with our theoretical
 443 prediction. The output of our investigation is summarized in Figure 5.

444



445

446 **Figure 5: b-value vs fractal dimension of hypocenters (D) in Southern**
 447 **California. The b-value of the Gutenberg-Richter law is found to be**
 448 **positively correlated with the fractal dimension of hypocenters in Southern**

449 California (SCEC Catalog, [Hauksson et al., 2012](#)) coherently with previous
450 literature on the topic. The plot represents shallow crustal seismicity from
451 1/1/1990 to 20/1/2025 (depth lower than 30 km) in between latitude 31°-37°
452 N and longitude 115°-122° W and above the completeness magnitude $M_c =$
453 2.5. Error bars represent 1σ uncertainty. The b-value is estimated using the
454 b-more-positive approach, while the fractal dimension of hypocenters is
455 found by applying the Grassberger & Procaccia algorithm ([Grassberger](#)
456 [and Procaccia, 1983](#)). The red line is the output of the non-linear fit b-value
457 $= b_M 10^{-r^{2-D}}$ whose trend is predicted in our model and derived in [Venegas-](#)
458 [Aravena and Cordaro, 2023b](#).

459

460 5. Discussions

461 The chaotic nature of earthquakes has been a subject of intense debate within
462 the scientific community (e.g., [Scholz, 1990](#); [Huang and Turcotte, 1992](#); [Goltz,](#)
463 [1997](#); [Vieira, 1999](#); [Yilmaz et al., 2023](#)). Traditionally, earthquakes have been
464 considered highly unpredictable due to the complexity of the processes involved
465 in fault rupture (e.g., [Geller et al., 1997](#); [Kagan 1997](#)). However, a growing body
466 of research, spanning conceptual frameworks, crustal stress, thermodynamics,
467 artificial intelligence, and GNSS measurements, suggests that fault stability may
468 be investigated, with a potential influence on precursory activity, may be achieved
469 (e.g., [Wyss, 1997](#); [Crampin and Gao, 2010](#); [Posadas et al., 2021](#); [Bhatia et al.,](#)
470 [2023](#); [Bletery and Nocquet 2023](#); [Devi et al., 2024](#)), especially in the case of
471 larger magnitude seismic event ([Kaveh et al., 2024](#)). In this study, we propose a
472 novel perspective, grounded in multi-scale thermodynamics, suggesting that the

473 chaotic nature of earthquakes may be modulated by the residual energy stored
474 within faults volumes ([Venegas-Aravena and Cordaro 2023a](#)). These primarily
475 theoretical developments suggest that as the thermodynamic fractal dimension
476 (D) decreases, the residual energy in the system increases. This correlation is
477 further supported by the framework presented in Section 2, where we
478 demonstrate that a higher residual energy is linked to a lower sum of Lyapunov
479 exponents Λ (red regions in Figure 1b), a hallmark of reduced chaoticity.
480 Consequently, earthquakes with higher residual energy exhibit more deterministic
481 behavior, as larger magnitude earthquakes (lower Λ) show a smaller change in
482 magnitude per unit of residual energy (Figure 1c). This suggests that these
483 earthquakes are less susceptible to the exponential growth of small
484 perturbations, a hallmark of chaotic systems.

485 Section 3 evaluates the proposed hypothesis through numerical simulations of
486 kinematic rupture using the HE-B method. This method has been effective in
487 modeling self-arrested earthquakes that comply with observational constraints,
488 such as the asperity criterion of [Somerville et al. \(1999\)](#), which links the zone of
489 high slip (known as asperity) to the release of a large amount of seismic energy.
490 The results obtained show a clear relationship between residual energy and the
491 variability in earthquake magnitude. For instance, the blue zone in Figure 2e
492 demonstrates that smaller earthquakes ($M_W < 6.6$) exhibit a high sensitivity to
493 variations in available energy. This means that small increases in available
494 energy can lead to significant increases in the magnitude of the resulting
495 earthquake, as evidenced by the high values of ΔM_W and the regions of high
496 magnitude variability (transparent blue zone in Figure 2f). Specifically, this
497 indicates that a 1% increase in available energy can result in an increase in the

498 magnitude of the resulting earthquake greater than $0.5M_W$. In other words, when
499 a fault has a specific low value of available energy (e.g., $0.1 \times 10^7 J/m^2$) and has
500 the potential to generate a magnitude $M_W = 5$ earthquake, it can produce a larger
501 earthquake ($M_W = 5.5$) or a smaller one ($M_W = 4.5$) if the available energy is
502 slightly increased or decreased. In contrast, larger earthquakes ($M_W > 7.8$) show
503 a notable lack of sensitivity to changes in available energy (transparent green
504 and red zones in Figure 2f). These zones indicate that a slight increase or
505 decrease in available energy does not produce significant changes in the
506 magnitude of the resulting earthquake. This low sensitivity demonstrates less
507 chaotic behavior in the simulations, where larger earthquakes are more
508 predictable and less influenced by small perturbations. This also suggests that
509 these larger events may also be more predictable. In line with this, analyses
510 conducted in seismic rupture simulations with a rate-and-state friction law on
511 simple faults (Kaveh et al., 2024) are consistent with the results shown in this
512 work, which is based on the distribution of residual energy. An interesting aspect
513 of the work by Kaveh et al. (2024) comes from the threshold above which
514 predictions can be made. They observed that it was possible to make forecasts
515 of earthquakes with magnitudes greater than approximately $M_W 6.9$. In contrast,
516 Figures 2e and 2f indicate that the variation in magnitude due to a change in
517 available energy begins to be less than $0.1M_W$ when earthquakes begin to have
518 magnitudes greater than $M_W 6.6$ (transparent yellow zone). This suggests a
519 similar threshold for predictability in both studies.

520 To further assess the link between residual energy and chaos in seismic activity,
521 numerical simulations were performed varying both the fracture energy
522 distribution and the system's available energy, in accordance with Equation 5.

523 The findings strongly corroborate the proposed hypothesis. Figure 3c, for
524 example, demonstrates an excellent agreement between the simulated
525 magnitude-parameter D relationship (black curve) and the theoretically predicted
526 one (Equation 2, red curve), affirming the established theoretical connection
527 between parameter D , residual energy, and simulated earthquake magnitudes.
528 Additionally, Figures 3d and 3e corroborate the trend observed in the earthquakes
529 of Figure 2: larger magnitude earthquakes exhibit a smaller variation in their
530 magnitude, suggesting a lower degree of chaos in these events. To gain further
531 credibility, a parameter more commonly used in seismology was needed. The b -
532 value of the Gutenberg-Richter law has traditionally been used as an indicator of
533 the relative occurrence rate of earthquakes of different magnitudes (Ito and
534 Kaneko, 2023; Lacidogna et al., 2023). In this study, the relationship between the
535 b -value and the degree of chaos in seismicity has been explored. In particular,
536 Figure 3f shows that the b -value decreases more abruptly for earthquakes with
537 magnitudes greater than $M_w 7.3$ (red zone), indicating a decrease in the
538 occurrence rate of smaller earthquakes relative to larger ones. This behavior is
539 associated with b -value changes on the order of 0.4 and suggests a less chaotic
540 regime. On the other hand, for earthquakes with magnitudes less than $M_w 6.3$
541 (blue zone), the b -value exhibits less pronounced changes, indicating a greater
542 variability in the occurrence rate of earthquakes of different magnitudes, and
543 therefore, a more chaotic regime. This relationship between the b -value and
544 seismic chaos is consistent with the interpretation of parameter D . Low D values
545 (associated with higher-magnitude earthquakes) imply a more homogeneous
546 distribution of residual energy over a larger area within faults, thereby reducing
547 the probability and number of smaller events. Consequently, the ratio of large to

548 small events, known as the b-value, is directly influenced by D . Therefore, it can
549 be argued that a typical decrease in the b-value (e.g., [Rivière et al., 2018](#); [Sharon
550 et al., 2022](#); [Chan et al., 2024](#)) is a measure of the chaos of a system, supporting
551 the notion that low b-values are associated with imminent larger magnitude
552 earthquakes, coherently with previous research (e.g., [Gulia and Wiemer, 2019](#)).

553 We also prove the theoretical relationship between the b-value and fractal
554 dimension $b - value = b_M 10^{-r(2-D)}$, also supporting the idea that large
555 earthquakes tend to occur within networks with low fractal dimensions (i.e., along
556 major faults). See Figure 5 for the output in the case of the SCEC Catalog in
557 Southern California (1990-2025). Observations show good agreement with
558 theory and, even though with relatively large uncertainties, are statistically robust.
559 It is important to note that, according to [Venegas-Aravena and Cordaro \(2023b\)](#),
560 the case where b-value and D are approximately proportional can be obtained,
561 as shown in Figure 5. The direct effect of the fractal dimension of faulting on the
562 maximum magnitude is more difficult to observe since large earthquakes are rare
563 events and the available seismic catalogs only contain a few cases, if any, of
564 events with the largest expected magnitude for each fault system, preventing a
565 reliable analysis. Moreover, the results would be rescaled for the size of the
566 largest seismogenic source in each fault network, which is tricky to estimate.
567 Conversely, the b-value can nowadays be evaluated by robust and unbiased
568 estimators. This is the reason why we choose to validate directly the relationship
569 between b and D .

570 The analyses conducted in this study, which involves theoretical, numerical and
571 observational data, contrast with the traditional view of earthquakes as highly
572 chaotic systems. However, it is important to note that our proposal does not

573 dismiss the role of chaos in seismic rupture dynamics. Rather, it suggests that
574 the chaotic behavior may be modulated by the amount of residual energy stored
575 in the fault. In other words, when residual energy is high, the probability of
576 releasing all that energy in a sudden event increases due to the coalescence of
577 different fault segments ready to nucleate or that can be dynamically activated
578 during the coseismic phase. Conversely, when residual energy is low, the fault
579 has more options for releasing that energy, which can lead to seismic ruptures of
580 varying sizes. Our findings have significant implications for understanding the
581 precursor seismicity, known as foreshocks (e.g., [Lippiello et al., 2019](#); [Bolton et
582 al., 2023](#)). If we can accurately quantify the residual energy in a fault, we may be
583 able to estimate the probability of large magnitude earthquakes and assess their
584 destructive potential.

585 This is particularly significant as recent research indicates that approximately half
586 of large-magnitude seismic events may be preceded by precursor seismic
587 activity, although the magnitude difference between foreshocks and mainshocks
588 does not appear to be substantial ([Wetzler et al., 2023](#)). This suggests that states
589 of higher energy preferentially evolve into large earthquakes. These states are
590 associated with a greater amount of available energy, which is directly related to
591 stress. Here, it is important to note that low values of D also indicate an
592 accumulation of stress in localized areas ([Venegas-Aravena et al., 2022](#)).
593 Geodetic measurements have confirmed this accumulation of localized stresses
594 between earthquakes of magnitudes greater than $M_w 7$ ([Kato and Ben-Zion,
595 2020](#)). Additionally, foreshocks also appear to be related to the geometric
596 conditions of faults (e.g., [McLaskey and Kilgore, 2013](#); [Cattania and Segall,
597 2021](#)), which can be incorporated into the residual energy through fracture

598 energy. Subsequently, residual energy can be used to estimate the physics of
599 seismic precursors. For instance, it has been estimated that foreshocks may not
600 be reliable when estimating the probability of subsequent mainshocks
601 ([Zaccagnino et al., 2024](#)). Here, residual energy in the context of multi-scale
602 thermodynamics can offer two explanations for the lack of clarity regarding
603 foreshocks. Firstly, foreshock-type activity should arise as a stress perturbation,
604 which, when considering a state of residual energy, can trigger events of different
605 magnitudes but within a range of magnitudes close to that of the potential future
606 mainshock. However, the magnitude of these foreshocks can be chaotic, limiting
607 the ability to conduct statistical analyses and thus declaring them as foreshocks
608 in real-time measurements. Secondly, the increase in residual energy implies a
609 lower variability in the magnitude of earthquakes, in agreement with [Lippiello et](#)
610 [al., 2024](#). This suggests that when residual energy may be very high in a fault,
611 stress perturbation has higher chances to trigger large earthquakes, limiting the
612 existence of foreshocks. That is, the probability that larger earthquakes are
613 affected or associated with foreshocks could decrease with an increase in the
614 magnitude of the mainshock.

615 Here, one way to associate foreshocks with large magnitude mainshocks is if the
616 rupture area of a foreshock reaches a zone of the fault with very high residual
617 energy, which could be seen as a perturbation leading to a single large
618 subsequent earthquake. This scenario could occur in the so-called "Mogi
619 Doughnut" of subduction zones ([Mogi, 1969](#)), where the shallowest zones of the
620 plate interface accumulate large amounts of energy while most seismicity occurs
621 at deeper locations with lower accumulated energy ([Schurr et al., 2020](#)).

622 The findings and interpretation carried out in this work agree with recent modeling
623 and observational findings in the literature (e.g., [Longobardi et al., 2024](#); [Nielsen,](#)
624 [2024](#)) and could also influence future research, which may focus on developing
625 methods to directly measure residual energy in natural faults, creating more
626 sophisticated models that incorporate parameter D and allow for the simulation
627 of the evolution of residual energy over time. Finally, exploring the implications of
628 our results for seismic risk assessment and the design of earthquake-resistant
629 structures.

630

631 **6. Conclusions**

632 The results obtained in this study suggest that the chaotic behavior of
633 earthquakes can be modulated by the amount of residual energy stored in the
634 fault. Our findings indicate that larger earthquakes, associated with higher
635 residual energy, exhibit less chaotic behavior. This new perspective challenges
636 traditional conceptions about the nature of earthquakes and opens new avenues
637 of research in seismology. While these results are promising, further research is
638 required to confirm and deepen our findings. Specifically, methods need to be
639 developed to directly measure residual energy in natural faults and to construct
640 more sophisticated models that incorporate the parameter D . This research will
641 allow us to advance our understanding of the mechanisms governing the
642 generation and propagation of earthquakes and pave the way for a better
643 assessment and mitigation of seismic risk.

644

645

646 **Data and Code Availability**

647 Data and codes are available upon reasonable requests to both the authors.

648

649 **Acknowledgment**

650 The authors wish to thank the editor for handling the manuscript and reviewers
651 whose comments and suggestions greatly improved the first version of the paper.

652 Discussions with Valeria Becerra-Carreño, Enrique G. Cordaro, Matteo Picozzi
653 and Didier Sornette have been inspiring for the work.

654

655 **Competing interests**

656 The authors declare no competing interests.

657

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