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Resolved tropical cyclones trigger CO₂ uptake and phytoplankton bloom in an Earth system model simulation

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¹ Abstract

The ocean carbon cycle is directly impacted by storms in the atmosphere. Tropical cyclones 2 (TCs), particularly, are known to drive intense air-sea CO_2 fluxes and to trigger phytoplankton blooms. 3 However, the latest generation of Earth system models (ESM) cannot realistically represent TCs due to 4 their coarse spatial resolution (typically 100-200 km grid spacing). Here, we present the first km-scale 5 coupled, global, storm- and eddy-resolving (5 km ocean, 5 km atmosphere) ESM simulation including 6 ocean biogeochemistry that is able to resolve TCs, and the cascade of physical-biogeochemical mecha-7 nisms that unfold in their response. Our simulated TCs enhance CO₂ fluxes by 20-40 times and cool the 8 surface ocean by 2-3°C, thus contributing to inverting the CO₂ flux direction from ocean outgassing to 9 uptake. Our TCs furthermore trigger a phytoplankton bloom in autumn in the western North Atlantic, 10 which is missed by coarser ESMs. While our TCs boost primary production, they also warm the ocean 11 subsurface, which enhances organic matter remineralization and attenuates their impact on carbon ex-12 port to depth. Our novel model configuration reproduces mechanisms underlying the ocean carbon cy-13 cle variability that remained so far unresolved in ESMs. By representing fine-scale atmosphere-ocean 14 biogeochemistry interactions in our ESM, we pave the way for future work to constrain uncertainties in 15 the role of km-scale events in the ocean carbon cycle at global and climatic scales. 16

17 Significance

Earth system models work with relatively low spatial resolution, typically 100-200 km in grid 18 spacing. This is too coarse for the representation of small-scale mechanisms and extreme events, such 19 as intense tropical cyclones (e.g. hurricanes). Nevertheless, such events trigger a number of physical 20 and biological processes that are relevant for climate, but missed by coarser models. Here, we present 21 the first km-scale global Earth system model simulation that resolves interactions between extremely 22 intense cyclones and the ocean carbon cycle with unprecedentedly high resolution (5 km ocean and 23 atmosphere). This will allow future work to reduce uncertainties in the role of km-scale processes in 24 the ocean carbon sink and in its response to a changing climate. 25

26 **1 Main**

The ocean absorbs about 25% of anthropogenic CO2 emissions every year [1, 2]. However, much 27 of this air-sea exchange of CO₂ is mediated by processes in the surface ocean and atmosphere that are 28 too small to be fully resolved in Earth system models (ESMs). State-of-the-art ESMs, composing 29 the Coupled Model Intercomparison Project (CMIP6) [3] and used to inform the Intergovernmental 30 Panel on Climate Change (IPCC), work with spatial resolution of tens to hundreds of kilometers, which 31 is too coarse to realistically represent tropical cyclones (TCs) [4, 5]. Intense TCs of category 4-5 32 in the Saffir-Simpson scale (winds > 58 m/s) are particularly misrepresented, or totally absent, in 33 coarse-resolution CMIP6-type models (100-200 km grid spacing) [4, 6]. The higher-resolution subset 34 of CMIP6 models (HighResMIP) [7], reaching up to 25 km in grid spacing, only "begin to capture 35 some structures of TCs more realistically" [4], but biases in TC characteristics persist [8, 9]. Here, 36 we present the first coupled storm- and eddy-resolving simulation (5 km ocean, 5 km atmosphere) 37 including the ocean biogeochemistry component HAMOCC (HAMburg Ocean Carbon Cycle) [10] in 38 the ICON (ICOsahedral Non-hydrostatic) model [11]. Our model configuration has been shown to 39 reproduce important TC characteristics, such as intensification rates, core structure [9], and interactions 40 with ocean eddies [12]. Here, we investigate the two intense TCs (i.e. category 4 or higher) in the 41 western North Atlantic occurring in our simulation, and their impact on the ocean carbon cycle. 42

TCs are known to drive high CO₂ fluxes between the ocean and atmosphere [13, 14]. TCs occur 43 about as frequently in ocean outgassing regions as in uptake regions, which mitigates their global annual 44 net effect on the ocean CO₂ uptake [15]. Nevertheless, TCs are responsible for moderating about 20-45 60% of all air-sea CO₂ exchange in their most prevalent regions and seasons, according to mooring 46 observations and regional modelling studies [16-19], thereby playing an important role in modulating 47 the ocean carbon sink variability. TCs impact the ocean CO_2 sink for weeks after their passage by 48 causing intense ocean heat loss, and hence cooling the surface ocean by 1-6°C [20, 21]. Cooling 49 decreases the ocean partial pressure of CO₂ (pCO₂) [22], thereby favoring an increasing ocean CO₂ 50 uptake (or decreasing CO₂ outgassing). Therefore, a realistic representation of the impact of TCs on 51 the ocean CO₂ sink must include both their instantaneous and longer-lasting components [15]. 52

A cascade of coupled physical-biogeochemical processes and feedbacks occur in response to 53 TCs. First, not only do air-sea CO₂ fluxes increase exponentially with surface wind speeds [23]. High 54 wind speeds during TC events also mix the upper ocean, contributing to cooling the surface ocean, and 55 warming the subsurface [20, 24]. Ocean mixing also replenishes the sunlit upper oceans with nutrients 56 from below, which enhances marine net primary production by phytoplankton [25-29]. Consequently, 57 intense TCs also mediate the ocean CO2 sink by enhancing the export of organic carbon to the ocean 58 interior [30]. The phytoplankton bloom following TCs also affects the penetration of shortwave radia-59 tion in the ocean, which further amplifies the TC-driven surface cooling effect [31, 32]. Lastly, the cold 60 wake of TCs is observed to negatively feedback on their own intensity and further development [33]. 61

Investigating the impact of km-scale extreme storms on the Earth's climate and carbon cycle remains challenging. Future projections of intense TC impacts, for example, rely on statistical or downscaling techniques [e.g. 34–36], which cannot consistently simulate the feedbacks between scales (i.e. the small scales of TCs and the global and climatic scales) and neither between Earth system components (i.e. ocean-atmosphere feedbacks). Our novel km-scale global ICON configuration is thus an invaluable tool that allows us to resolve ocean-atmosphere interactions triggered by extreme storms within the Earth system in a consistent manner for the first time. We explore the cascade of coupled ⁶⁹ physical-biogeochemical mechanisms that unfold after the passage of the two intense TCs occurring in ⁷⁰ the western North Atlantic in our simulations – the first intense TCs ever resolved in a fully coupled ⁷¹ global ESM including ocean biogeochemistry. We show how such extreme storms can increase the ⁷² ocean CO₂ uptake and trigger an autumn phytoplankton bloom in the western North Atlantic.

73 2 Results and Discussion

Our simulation produces two intense TCs in the western North Atlantic. The first one (TC1) occurs between 1-10 September, and the second one (TC2) occurs between 10-17 September. Both TCs intensify to become hurrisones (i.e. surface wind speeds >32 m/s) and reach hurrisone setescery.

TCs intensify to become hurricanes (i.e. surface wind speeds \geq 33 m/s) and reach hurricane category



Figure 1: **Simulated tropical cyclones characteristics.** Along-track maximum surface wind speed (**A**), translation speed (**B**), and changes in sea surface temperature (SST) due to TC passage (**C**) expressed as differences in SST from 24 hours after minus 24 hours before the TC averaged over a 200 km-radius circle around track points. Track points are spaced by 3 hours in A, B and C for visualization. Maps of surface latent heat flux (LHF, **D**, instantaneous) on 10 September at 16h UTC (12-noon at the longitude of TC1 location), and air-sea CO₂ flux (**E**, hourly mean) at 11h UTC (near TC1 at its maximum intensity). Both LHF and CO₂ fluxes are positive upwards (from ocean to the atmosphere), thus positive values show ocean heat loss and ocean CO₂ outgassing. Zoomed-in views of air-sea CO₂ fluxes (**F**) and SST (**G**), highlighting the cold wake of TC1 after its maximum intensity. Black cyclone symbols in E, F and G mark the position of TC1 at the time corresponding to the data shown, while grey cyclone symbols mark its previous positions at 12 and 24 hours before.

4 (>58 m/s) in latitudes north of 30°N (Fig. 1A). Our two TCs follow different paths in the North 77 Atlantic: Both TCs initially travel with a typical north-westward direction, but TC1 takes a sharp turn 78 eastwards after crossing 30°N while intensifying, somewhat similar to 2020 hurricane Paulette [37], 79 but without making landfall in Bermuda. During its sharp turn, TC1 travels with low translation speeds 80 below 6 m/s (Fig. 1B). On the other hand, TC2 follows a predominantly northward track after reaching 81 hurricane intensity, and faster than TC1, with translation speeds always above 6 m/s. Slow translation 82 speeds enable TCs more time to exert sustained wind stress on the ocean [38, 39], leading to enhanced 83 mixing and sea surface cooling. 84

TC1 cools the surface ocean by up to 2° C, while TC2 cools the surface ocean by about 1° C 85 where wind speed is highest (Fig. 1C). The largest cumulative sea-surface cooling (2-3°C) takes place 86 in the extra-tropics, where the two TC tracks approach each other and their areas of impact overlap 87 (Fig. 2D,F). This is also the warmest region along their tracks, which favors their own intensification. 88 Sea surface cooling is driven by the combined effect of surface heat fluxes and wind-driven vertical 89 diffusion, and only slightly counterbalanced by warm advection (Fig S1, see 'Ocean tracer budget 90 terms' in Methods). Horizontal advection plays a larger role in warming the wake of TC1, into which 91 the Gulf Stream brings warm tropical waters. 92

Our simulated TCs yield extreme latent heat fluxes (LHF) of $>1200 \text{ W/m}^2$ at their maximum wind speeds (Fig. 1D), in line with observations [e.g. 40], and larger than anywhere else on the Earth's surface, which contributes to cooling the surface ocean. For context, maximum surface LHF values in the Amazon forest range in 200-500 W/m² during peak incoming solar radiation at noon, when land evaporation and transpiration by vegetation are at largest [41, 42]. Not only do hurricanes drive intense surface heat exchange, but also intense CO₂ exchange rates between the ocean and atmosphere.

99 2.1 Hurricanes impact the ocean carbon sink

Our simulated TCs impact the ocean CO_2 uptake in two ways. First, TCs drive extremely large air–sea CO_2 fluxes *instantaneously* – during their passage – due to high wind speeds, thereby enhancing ocean outgassing (Fig. 2A,B, Fig. S2A,B). These fluxes are up to 20–40 times larger than CO_2 fluxes during non-TC conditions (1–31 October), in line with observations and previous regional modeling work [13, 14, 16]. The largest instantaneous CO_2 fluxes occur within a radius of 150-200 kilometers from the TC center at their full development, where wind speeds are the strongest (Fig. 1E-F, Fig. S3).

Second, our TCs cause a *longer-lasting* effect driven by sea surface cooling. The TC-driven rapid decrease in sea surface temperature (SST) causes an also rapid decrease pCO₂ that persists after their passage (Fig. 2C-F), in line with observations and regional modeling work [e.g. 13, 43]. The link between decreasing SST and CO₂ flux anomalies is visible in the wake of TC1 (Fig . 1F,G). The cool and low-pCO₂ wake is shifted to the right of the TC track, which is a characteristic feature of fast-translating TCs [21, 24]. The rightward shift is explained by the fact that TC wind vectors rotate in phase with the inertial current vectors to the right of the TC track [20, 44].

Sea surface cooling reduces surface ocean pCO₂ by \sim 10-40 ppm after each TC in the extratropics (Fig. 2F, Fig. S4, see '*Decomposing changes in pCO*₂' in Methods). The cooling impact on decreasing surface pCO₂ is only slightly counterbalanced by entrainment of dissolved inorganic carbon (DIC), and by increasing surface salinity, both increasing surface pCO₂, especially at latitudes north of 35°N (Fig. S4B,C). The latter is explained by vertical diffusion of salinity, which dominates over freshening due to surface fluxes (precipitation > evaporation during both TCs). The relative role of temperature, DIC and salinity in changing surface pCO₂ in response to TCs varies substantially



Figure 2: Along-track time-series diagrams of TC2 impacts on CO₂. Time series averaged over a moving 200 km-radius circle around the center of TC2 along its track. Shown quantities are air-sea CO₂ fluxes (**A**, negative indicates ocean uptake), and the fold-increase in air-sea CO₂ flux during the passage of the TCs along their tracks in comparison to 1-31 October, when no other TCs have occurred in this region (**B**). TC2 along-track absolute values of surface ocean pCO₂ (**C**), sea-surface temperature (SST, **D**), and their changes with respect to 12 hours before the TC2 passage (**E** and **F**, respectively). The vertical axes show the latitude coordinates of the center of TC2 along its track. The two-step reduction in pCO₂ and SST one followed 5-6 days after the other, shows the overlapping impacts of TC1 on the track of TC2 especially above 30°N. One can also see the daily cycle in the background as vertical stripes in all time-series diagrams. TC1 along-track time-series diagram for CO₂ is shown in Fig S2.

in different regions and between individual TC episodes [e.g. 16–19, 43, 45]. In the North Carolina estuary, for example, TCs have been observed to cause net CO₂ outgassing, where increasing surface DIC outweighs the cooling effect [46–48]. Pre-existing gradients in temperature, DIC, alkalinity and salinity, as well as cyclone intensity and translation speed, and thus their ability to induce vertical mixing, will altogether determine their net effect on pCO₂ anomalies [49, 50].

The CO₂ flux response to the longer-lasting cooling effect depends on the resulting local imbalance in pCO₂ between the surface ocean and atmosphere. The western North Atlantic is supersaturated in pCO₂ before the arrival of our TCs (i.e. $pCO_2^{ocean} > pCO_2^{atmosphere}$), which causes them to induce ocean CO₂ outgassing during their passage. Reducing surface pCO₂ either decreases ocean CO₂ outgassing, or induces uptake if the pCO₂ imbalance is inverted to undersaturated conditions (i.e. $pCO_2^{ocean} < pCO_2^{atmosphere}$).

Accumulating CO₂ fluxes from 24 hours preceding the TCs reveals that the longer-lasting cooling effect outweighs the instantaneous outgassing by mid September, especially after TC2 and at latitudes higher than 30°N, resulting in a net increase in ocean CO₂ uptake (Fig. S1C). The opposing instantaneous and longer-lasting TC impacts on air-sea CO₂ fluxes have been suggested to nearly cancel each other out at the global and annual scales, as well as regionally in the North Atlantic [15]. In other tropical regions, this effect may be less well-balanced. Previous work performed with a global ocean
model forced by atmospheric reanalyses showed that TCs would cause net ocean CO₂ outgassing in the
eastern North Pacific, but net uptake in the western North Pacific, and account for 2% and 33% of the
net CO₂ flux regionally, respectively [15].

Considering only their instantaneous effect, our two TCs together induce about 0.9 TgC in CO2 140 outgassing (1 TgC = 10^{12} g or 1 million metric tons of carbon) during their passage (integrated over a 141 200 km-radius circle along their tracks). This is equivalent to $\sim 18\%$ of the outgassing that takes place 142 in their track region during the warm season (July-September, +5.1 TgC), while also equivalent to less 143 than 2% of the CO₂ uptake that takes place in the same region over the rest of the year (-63.8 TgC). 144 While it is straightforward to quantify the instantaneous outgassing impact of our TCs by integrating 145 CO₂ fluxes in their track region during their passage, quantifying their longer-lasting cooling impact in 146 isolation is sensitive to the integration time window once chooses. Instead, here we underline that our 147 TCs contribute to inverting the air-sea CO₂ flux direction – from ocean CO₂ outgassing to uptake – by 148 modulating surface pCO₂ variability that has remained so far unresolved in global coupled ESMs. 149

150 2.2 Hurricanes trigger second phytoplankton bloom in autumn

We simulate a local increase in net primary production (NPP) by 1-2 orders of magnitude in 151 response to our TCs (Fig. 3A), in agreement with satellite observations [25–27]. Similarly to the 152 cooling effect, the NPP increase is also shifted to the right of the TC tracks, according to theory and 153 observations [20, 52], which is especially true for fast-translating TCs [25, 39]. However, the TC-driven 154 increase in phytoplankton concentration is not spatially restricted to the vicinity of their tracks. The Gulf 155 Stream and its mesoscale meanders play an important role in transporting phytoplankton away from the 156 track of TC1 after its passage (Fig. 3B-C). In fact, TCs themselves also intensify surface currents [53, 157 54], contributing to further spread their biological impact. The spatial spread of phytoplankton is mainly 158 done by geostrophic currents, which are characterized by flowing along contours of sea-surface height 159 anomalies (Fig. 3B-C). This highlights the role of the pronounced mesoscale, high eddy kinetic energy 160 currents in advecting phytoplankton horizontally away from the TC track, which continues for several 161 days after the TC passage. Therefore, the increase in phytoplankton concentration in the wake of TCs 162 is thus partly counterbalanced by horizontal advection (Fig. S5A-C). Consequently, our simulated TCs 163 trigger a phytoplankton bloom that spreads across the western North Atlantic far beyond the region of 164 hurricane-force winds. 165

Our simulated TCs contribute to the development of a phytoplankton bloom in autumn, distinct 166 from the spring bloom (Fig. 3D). The observed seasonal cycle of chlorophyll-a in the western North 167 Atlantic shows a peak in spring, a minimum by the end of summer, and a growing phase during au-168 tumn (Fig. 3E). Although our simulation does not yet perfectly reproduce the observed phytoplankton 169 phenology, it represents an improvement compared to state-of-the-art ESMs, in that our TCs trigger a 170 phytoplankton bloom in autumn. ESMs are able to represent the spring bloom from observations, but 171 they generally fail to capture the autumn recovery (Fig. 3F). A realistic representation of phytoplankton 172 phenology by ESMs has been a long-lasting challenge, especially in capturing such secondary blooms 173 in mid latitudes [55]. The occurrence of distinct autumn blooms is observed along the North Atlantic 174 current further north at 40-50°N [56], generally triggered by mixed-layer deepening by storms, not nec-175 essarily only TCs [57–59]. Nevertheless, the role of TCs in triggering autumn blooms has been reported 176 in observations [25, 39, 60], and TCs are suggested to drive interannual variability of chlorophyll-a in 177 the western North Atlantic [25, 57]. In 2020, hurricanes Paulette and Teddy followed paths similar 178



Figure 3: **Growth and spread of phytoplankton in response to TCs. A)** Vertically integrated net primary production (NPP) on 28 August 2020 (upper panel, before the hurricanes), on 10 September 2020 (middle panel, after TC1), and on 17 September (lower panel, after TC2). **B**) Zoomed-in map of surface velocities and sea-surface height anomalies (contours) in the Gulf Stream box shown in A. **C**) Temporal evolution of vertically integrated phytoplankton in the Gulf Stream box on 6 September (upper panel, before TC1), on 13 (middle panel) and 16 September (lower panel) after TC1, whose track is shown in red. **D**) Time series of Chlorophyll-a concentrations averaged over the western North Atlantic (triangle in A) in our 1-year ICON simulation, **E**) from satellite observations in 2018-2022, and **F**) in a CMIP6 multi-model mean in 2010-2014. CMIP6 models with daily chlorophyll data available are MPI-ESM-LR (HAMOCC), IPSL-CM6A-LR (NEMO-PISCES), and CESM2 (MARBL). Chlorophyll-a data in E are obtained from NOAA CoastWatch at coastwatch.noaa.gov/cwn/index.html, which are derived from multiple satellites (VIIRS-SNPP, VIIRS-NOAA-20, and OLCI-Sentinel-3A) and processed with the DINEOF gap-filling method [51]. Thin lines in D-F show daily data, while thick line are a 31-day running mean.

- ¹⁷⁹ to those of our simulated TCs, and triggered a September peak in phytoplankton concentrations in the
- western North Atlantic (Fig. 3E, Fig S6). By resolving the km-scale interactions between atmosphere
- ¹⁸¹ and ocean biogeochemistry, we also increase the realism of the biological response to intense TCs.

¹⁸² 2.3 Cascading physical-biogeochemical response to hurricanes

Our TCs induce a sudden decrease in upper-ocean stratification, thereby increasing the mixed-183 layer depth by 20-40 m at their maximum intensity (Fig. 4A,B, Fig. S7A). Strong surface wind speeds 184 also induce an instantaneous pulse of nutrient entrainment into the euphotic zone by turbulent mixing 185 (Fig. S8A,B). While NPP responds immediately to the presence of nutrients (Fig. 3A), maxima in phy-186 toplankton concentration occur days to weeks later (Fig. 4C, Fig. S5A). Phytoplankton is continuously 187 diffused from the upper layers downwards by wind-driven mixing (Fig S8D, Fig. S9A,B). The vertical 188 diffusion and advection of phytoplankton is mostly contained within the upper 100 meters in the wake 189 of TCs (Fig. S5D,E). 190



Figure 4: **Cascading physical and biogeochemical response to hurricane passage.** Time series from a fixed point in space (about 37°N and 67°W, see inset map) on the track of TC1, depicting conditions from before and after its passage as a category 4 hurricane. Variables shown are surface wind speed (A), and vertical profiles of density stratification expressed as the Brunt-Väisälä frequency (**B**), phytoplankton concentration (**C**), temperature (**D**), and aerobic organic matter remineralization (**E**). Small white axes indicate the local inertial period at this latitude.

Near-inertial oscillations are triggered by the TCs, and become the prevailing mode of variability 191 in the surface and subsurface ocean in their wake region, propagating horizontally and vertically down-192 wards [20]. The horizontal advection of phytoplankton takes place mostly at near-inertial frequencies 193 (Fig. S8C, Fig. S9C), which underscores their role in spatially spreading the TC-driven bloom. The 194 downward propagation of near-inertial waves also transport kinetic energy [52] and further mix the 195 thermocline, thus contributing to sustaining subsurface warming for weeks after the TC passage (Fig. 196 4D) [61, 62]. The surface cooling and subsurface warming - a typical dichotomy in the wake of TCs 197 [20, 63] – is reproduced in our simulations, where 50m-temperatures increase by 1-1.5°C in the wake 198 of TCs (Fig. S7B). 199

Our simulated TC-driven subsurface warming contributes to enhance organic matter remineral-200 ization (Fig. 4D-E), in addition to the increase in biological production itself. Increasing remineral-201 ization, in turn, decreases the amount of carbon exported to depth. The impact of TCs on increasing 202 organic matter export is, therefore, not as pronounced as their impact organic matter production by 203 phytoplankton (Fig. S10A). Only about 8-10% of the organic matter produced in the euphotic zone is 204 exported to depths below 100 m in the western North Atlantic in our simulation (Fig. S10B). We do 205 not detect a peak in export ratios in response to our TCs passage, as opposed to recent observations 206 [30]. Nevertheless, our export ratios compare well with osbervations-based estimates near Bermuda 207 (6.3-8.4% at 150m [64]). In the western North Atlantic, export ratio estimates range at least between 208 5-30%, with values generally increasing form the open ocean towards coastal zones (see [65, 66] and 209 references therein). 210

211 **3** Conclusion

We have conducted the first storm- and eddy-resolving (5 km ocean, 5 km atmosphere) global coupled ESM simulation including ocean biogeochemistry, which allowed us to resolve atmosphereocean biogeochemistry interactions in the km-scale within the Earth system for the first time. We have investigated two particularly intense TCs reaching category-4 hurricane force (surface wind speed >58 m/s) in the North Atlantic, which state-of-the-art ESMs have been, so far, unable to capture.

We have resolved the cascade of physical-biogeochemical effects that unfold in response to TCs 217 in a global coupled ESM the first time. Extreme wind speeds enhance CO₂ outgassing by 20-40 times 218 from the western North Atlantic, initially supersaturated in pCO2. Our TCs also cool the surface ocean 219 by 1-2°C each, decreasing the surface ocean pCO₂ by 10-40 ppm to undersaturated levels. Upwelling 220 and mixing of DIC only slightly counteract the cooling effect on surface pCO_2 . Consequently, our TCs 221 contribute to inverting the CO₂ flux direction from ocean outgassing to ocean uptake. Moreover, our 222 intense TCs increase NPP locally by 10-100 times, triggering a phytoplankton bloom in autumn, which 223 typical ESMs so far fail to represent. This distinct autumn bloom persists in the western North Atlantic 224 for weeks after the TC passage. Phytoplankton is spread horizontally by ocean mesoscale currents, 225 especially in the vicinity of the Gulf Stream. TC-driven subsurface warming increases organic matter 226 remineralization, thus impeding a larger impact of the TC-driven bloom on organic carbon export to 227 depth. 228

The frequency and intensity of extreme weather events increase in response to climate change. 229 The proportion of intense TCs is projected to increase in the future with "high confidence" according to 230 the latest IPCC report [4], and so is their impact on the ocean carbon cycle. TCs intensify particularly 231 fast in response to marine heatwaves [67, 68], which will also tend to become more intense in the future 232 [69, 70]. The large variability between TCs, and extreme events overall, prevents generalizations of our 233 findings in time or space, for which longer simulations are needed. Longer km-scale ESM simulations 234 (with grid-spacing of 5 km or finer) are necessary to realistically represent such future changes, their 235 impacts and feedbacks to climate by modulating the ocean CO_2 sink. Our work shows that this is now 236 possible. By resolving small-scale interactions between atmosphere and ocean biogeochemistry in a 237 global coupled ESM, we bridge the gap between km-scale mechanisms and the ocean carbon cycle at 238

239 global and climatic scales.

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Author Contributions

TI and CH conceived the project "Coupled Climate Carbon Cycle Sapphire (C4S)". DMN, FC, NS and TI conceived this study's idea. FC developed the coupled ICON model configuration with HAMOCC. DMN conducted the ocean-only spinup, and FC conducted the coupled ICON experiments. FC, NS and NB implemented the online calculation of budget terms. AK implemented the tracking algorithm and characterized the tropical cyclones. DMN analyzed the model output with support from all authors. DMN wrote the manuscript, revised by all authors.

259 Competing Interests

The authors declare no competing interests.

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467 Methods

ICON model and simulations

We use the ICOsahedral Non-hydrostatic (ICON) coupled model in the "Sapphire" configuration 469 [11], which includes atmospheric, ocean, land and ocean biogeochemistry components. ICON uses a 470 unstructured triangular grid, here configured with nominal mean horizontal resolution of 5 km in all 471 components. The atmospheric dynamics are based on the numerical weather prediction (NWP) model 472 by the German Weather Service (DWD) [71]. In the atmosphere, convection is explicitly resolved, while 473 parameterizations are only used in processes that cannot be resolved in the km-scale, namely radiation, 474 microphysics and turbulence [11]. The atmosphere is coupled to the ocean and sea-ice model ICON-475 O [72]. Here, ICON-O is configured with 72 layers in z^* coordinates, which distributes sea-surface 476 height anomalies at every time step in all model layers. In our configuration, ICON-O resolves ocean 477 eddies explicitly, and therefore parameterizations for eddy-induced horizontal diffusion and advection 478 are switched off (i.e. GM-Redi [73, 74]), and only parameterizations for vertical turbulent mixing and 479 dissipation are maintained [72]. Land biogeophysics and biogeochemistry are represented in JSBACH 480 [75] version 4. Ocean biogeochmistry is included with the HAMburg Ocean Carbon Cycle (HAMOCC) 481 model [10]. 482

HAMOCC simulates biogeochemical processes in the water column, including interactions with 483 the atmosphere and sediment. HAMOCC was initially developed based on a nutrient-phytoplankton-484 zooplankton-detritus (NPDZ) framework [76], but is now expanded to include the representation of 485 20 biogeochemical tracers [10]. Marine biology in HAMOCC is represented by two tracer variables: 486 bulk phytoplankton, which includes calcifying and silicifying species, and cyanobacteria (diazotrophs) 487 [77]. Phytoplankton growth is co-limited by nitrate, phosphate, and iron, while cyanobacteria can 488 utilize dinitrogen. Our HAMOCC configuration includes the representation of dynamic particle sinking 489 speeds in the water column through the M4AGO scheme [65], and temperature-dependent organic 490 matter remineralization. A constant carbon-to-nutrient ratio is assumed for all marine organic matter 491 following a global Redfield ratio concept [78]. HAMOCC is embedded in ICON-O, and therefore they 492 share the same grid configuration. 493

Our coupled ICON configuration (ocean and atmosphere physics) is the same as that from Baker et al. [9], where they show that key TC characteristics are simulated realistically compared to observations, namely TC intensification rates, TC inner-core structure, and the relationship between TC size and intensification. Our ICON configuration only differs from that of Baker et al. in that we include the representation of ocean biogeochemistry with HAMOCC, and therefore also the feedback to the ocean physics through the dynamic representation of sunlight absorption by phytoplankton [79].

We use a step-wise strategy of model complexity and resolution to arrive at the initial conditions 500 for our coupled model simulation. This strategy consists in initializing our model from previously tuned 501 and spun-up simulations at coarser resolutions (Fig. 5), providing enough time for the next simulation 502 to adapt to the given initial conditions. Our coupled simulation starts from a transient ocean 5 km 503 ICON-O/HAMOCC simulation, and from the Integrated Forcast System (IFS) atmospheric analysis 504 at 01-01-2020 [80]. The ocean biogeochemistry is taken from a well-tuned and well spun-up coarser 505 ICON-O/HAMOCC simulation (40 km spatial resolution, 64 vertical layers) which, in turn, follows 506 from a >3000 years-long ICON-O physics-only spin-up. The high-resolution ocean physics is taken 507 from an ICON-O 5 km transient simulation started in 2010 forced with the ERA5 reanalysis [81] and 508 increasing atmospheric CO₂ concentrations. We remap the 40 km HAMOCC variables onto the 5 km 509

ICON-O grid and run ICON-O with HAMOCC at 5 km from 01-01-2017 for three years to allow for 510 enough time for the biogeochemical tracers in the upper ocean layers to adapt to the ocean mesoscale 511 field. Spinning-up HAMOCC to reach quasi-stability in all layers and biogeochemical tracers, including 512 in the ocean sediment, requires hundreds or thousands of years of simulations. Such long spin-ups are 513 still impractical at 5 km resolution. The here-applied cascading strategy allows us to generate initial 514 conditions with relatively stable conditions (sufficiently small drifts) in the euphotic zone - where we 515 focus on in this study - and that are consistent with the ocean physics from the 5 km ocean model. 516 Moreover, any long-term trends are small relative to the variability in our short 1-year simulation. After 517 running the coupled model for one year with daily output, we repeat the simulation from 20-August to 518 31-October with hourly output. 519

520 TC tracking

We deploy a tracking algorithm that only requires 2D surface fields and closely follows the 521 approach taken by Shimura et al. (2017) [82]. The algorithm takes the hourly mean sea-level pressure 522 (MSLP) and near-surface wind speed over both land and ocean as inputs and consists of two main steps: 523 (1) identifying candidate tropical cyclone points and (2) stitching together the points to create tracks. 524 For each time step, we begin by finding the grid cell with the lowest MSLP. If the MSLP is less than 525 1010 hPa, the grid cell is stored as a candidate TC point and all grid cells within a 800 km radius are 526 masked. This is repeated until the MSLP of all remaining grid cells exceeds 1010 hPa. In the second 527 step, the tracks are formed by connecting candidate points with a surface wind speed of at least 17 m/s 528 that are sufficiently close in time and space. To this end, we start by selecting the candidate point at 529 the earliest possible time step within the region 0° N to 35° to mark the start of a track. We then look 530 for the next candidate point to add to the track, which must lie within 100 km and the next 24 hours. 531 Priority is given to points closer in time. In the case where multiple points are found for the same 532 time step, the one with the lowest MSLP is selected and the remaining points are removed. The track 533 is successively extended until no more candidate points satisfy the above conditions. This process is 534 repeated until all candidate points are added to tracks or removed. In a further post-processing step, 535 we retain only those tracks with a lifetime greater than 24 hours and lifetime maximum intensity that 536 exceeds hurricane strength (>33 m/s). 537



Figure 5: **Step-wise strategy for initialization.** Schematics of previous simulations performed for tuning and spin-up, leading to the initial conditions used for our coupled simulation.

⁵³⁸ Our algorithm identifies 42 TCs that reach at least category-1 hurricane intensity in our 1-year ⁵³⁹ global simulation. We do not consider less intense tropical storms and depressions. Three of the 42 ⁵⁴⁰ hurricane-force TCs take place in the North Atlantic. Here we focus on the two *intense* TCs (i.e. ⁵⁴¹ category 4 or higher), since they are those typically missed by ESMs. A third and weaker TC occurs in ⁵⁴² the Gulf of Mexico, reaching hurricane category 1 shortly before making landfall.

543 Ocean tracer budget terms

We explicitly calculate the contribution of each term composing the tendencies of most biogeochemical tracers, temperature, and salinity, during the ICON model integration. This tendency decomposition is calculated online at each time step, and stored as daily or hourly means, depending on the output frequency. The total tendency of tracer C is calculated as:

$$\frac{\partial C}{\partial t} = -\mathbf{U} \cdot \nabla C + \frac{\partial}{\partial z} \left(k_v \frac{\partial C}{\partial z} \right) + \text{Surface fluxes} + \text{Biogeochemical sources and sinks}, \quad (1)$$

where the terms on the right-hand-side are the contributions from advection (we calculate horizontal and vertical advection separately), vertical diffusion by turbulent mixing, and surface fluxes, respectively. U = (u, v, w) is the ocean velocity vector, and k_v is the vertical diffusion coefficient obtained using a turbulent kinetic energy (TKE) closure parameterization (see [72]).

The surface flux of CO₂ is calculated in HAMOCC as $F = k_w S_{CO_2} (pCO_2^{ocean} - pCO_2^{atmo})$, 552 where k_w is the gas transfer velocity, a function of temperature and the squared 10m wind speed, calcu-553 lated according to [23], S_{CO_2} is the solubility of CO₂ in seawater, a function of temperature, calculated 554 according to [83], pCO_2^{ocean} is the surface ocean pCO₂ calculated at each time step from DIC and 555 alkalinity concentrations, and pCO_2^{atmo} is the surface atmospheric pCO₂ corrected for the total atmo-556 spheric pressure following the protocols for the Ocean Model Intercomparison Project (OMIP) [84]. In 557 our simulation, atmospheric CO₂ concentrations are uniformly prescribed over the entire ocean surface 558 following year 2020 values, and do not change in response to air-sea CO₂ fluxes. 559

560 Decomposing changes in pCO₂

⁵⁶¹ Following Takahashi et al.[22], we decompose the total change in surface pCO₂ triggered by ⁵⁶² the passage of the TCs between contributions from changes in sea surface salinity (SSS), sea surface ⁵⁶³ temperature (SST) and surface DIC and alkalinity concentrations:

$$\delta p CO_2^{Total} = \delta p CO_2^{SST} + \delta p CO_2^{SSS} + \delta p CO_2^{DIC} + \delta p CO_2^{Alk}.$$
(2)

Here " δ " denotes changes in time with respect to the TC passage, not to be confused with ΔpCO_2 , which commonly represents the pCO₂ imbalance between ocean and atmosphere. Each term is calculated by scaling the simulated surface pCO₂ with changes in SST, SSS, DIC, and alkalinity:

$$\delta p CO_2^{SST} = \gamma_{SST} \ p CO_{2ref} \ \delta SST$$

$$\delta p CO_2^{SSS} = \gamma_{SSS} \ p CO_{2ref} \ \delta SSS/SSS_{ref}$$

$$\delta p CO_2^{DIC} = \gamma_{DIC} \ p CO_{2ref} \ \delta DIC/DIC_{ref}$$

$$\delta p CO_2^{Alk} = \gamma_{Alk} \ p CO_{2ref} \ \delta Alk/Alk_{ref}$$
(3)

where "ref" denotes a local reference value taken 24 hours prior to the TC passage, $\gamma_{SST} = 0.0423^{\circ}C^{-1}$, $\gamma_{SSS} = 1$ and, from our model data, we use the approximation presented by [85] to calculate:

$$\gamma_{DIC} = \frac{3 \ Alk \ DIC - 2 \ DIC^2}{(2 \ DIC - Alk)(Alk - DIC)}$$

$$\gamma_{Alk} = \frac{Alk^2}{(2 \ DIC - Alk)(Alk - DIC)}$$
(4)

The coefficient γ_{DIC} – also known as the Revelle factor[86] – allows for a linear transformation between percentage changes in DIC and pCO₂, that is, a scaling factor between $\delta DIC/DIC_{ref}$ and $\delta pCO_2/pCO_{2ref}$. Similarly for γ_{Alk} and γ_{SSS} . On the other hand, γ_{SST} yields a 4.23% change in pCO₂ in response to a 1°C change in SST in absolute terms directly. Differences between the reconstructed δpCO_2^{Total} from Eq. 2 and its simulated value by ICON/HAMOCC are relatively small (Fig S4-E), which confirms the robustness of this widely-used method.

575 Data Availability

ICON output necessary to reproduce figures will be publicly available at the World Data Climate
 Center (WDCC) before this paper is accepted for publication.

578 Code Availability

579 Scripts used to reproduce figures will be made available at the Max Planck Society Publication 580 Repository, before this paper is accepted for publication, at pure.mpg.de/pubman/faces/HomePage.jsp.

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Supplementary file for "Resolved tropical cyclones trigger CO₂ uptake and phytoplankton bloom in Earth system model simulation"

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Figure S1: Changes in sea-surface temperature (SST) relative to the passage of TCs. Total change in SST (A), expressed as accumulating temperature tendencies from 12 hours before the TC passage, cumulative tendency terms due to the isolated effects of surface heat fluxes and vertical diffusion by turbulent mixing (B), and due to the effects of vertical and horizontal advection (C).



Figure S2: Impact of TCs on air-sea CO_2 fluxes. Along-track time series of air-sea CO_2 fluxes averaged over a 200-km radius from the TC center (here as in Fig. 2 in the main text, but for TC1) (A), ratio between magnitude of air-sea CO_2 flux during TC1 and baseline, which here is the 1-month period from Oct-1 and Oct-31 excluding the TC days (B), and the cumulative air-sea CO_2 flux starting at 12 hours prior to the TC passage at each location along their tracks (C). A switch from positive to negative cumulative flux in (C) indicates the moment when the CO_2 uptake following the TC passage outweighs the outgassing triggered during their passage.



Figure S3: Temporal evolution of surface wind speeds along the track of TC1 (left) and TC2 (right) as a function of radial distance from cyclone center.



Figure S4: Along-track time series of changes in pCO₂ due to changes in sea-surface temperature (SST, **A**), DIC and alkalinity (**B**), sea-surface salinity (SSS, **C**), their reconstructed combined effect (**D**), and the difference between this reconstruction and the pCO₂ changes simulated directly with our model (reconstruction bias, **E**). On the left-hand-side are shown time series along the track of TC1, while on the right-hand-side are shown time series along the track of TC1.



Figure S5: Along-track time series of cumulative tendencies of integrated (0-100 m) phytoplankton concentrations (A) and its components, that is, changes in phytoplankton concentrations due to biological sources and sinks (B), horizontal advection (C), vertical advection (D), and wind-driven vertical diffusion by mixing (E). Accumulating periods start 12 hours before the TCs passage.



Figure S6: Satellite-derived Chlorophyll-a observations on 01-09-2020, 19-09-2020, and 26-09-2020, overlaid with the tracks of TCs Paulette and Teddy. TC Paulette reached peak intensity to hurricane category 2 on 10-09-2020, and made landfall in Bermuda on 14-09-2020. TC Teddy reached peak intensity on 17-09-2020 to hurricane category 4. Black lines show the TC tracks. Chlorophyll-a data are obtained from [1], and TC tracks are obtained from NOAA's Atlantic hurricane database (HURDAT2) [2].



Figure S7: Along-track time series of changes in mixed-layer depth (MLD, A) and in temperature at 50 m (B) with respect to the passage of TCs.



Figure S8: Time series at 37°N 67°W, continuation of Fig. 4 in the main text. Here shown are changes in phosphate concentrations due to vertical diffusion by wind-driven mixing (A), phosphate concentrations (B), changes in phytoplankton concentration due to horizontal and vertical advection (C) and due to vertical diffusion by mixing (D)



Figure S9: Along-track time series of changes in phytoplankton concentration due to horizontal advection and vertical diffusion. Quantities are averaged from a 150 km-radius circle following the center of the TCs along their tracks. The left vertical axes show the along-track distance from the first point where the TC is detected. The right vertical axes show the latitude coordinate of the center of TCs along their tracks. The vertical advection components of phytoplankton tendency show a similar pattern to those of horizontal advection, but with opposite sign, due to continuity. Note the dominant role of inertial oscillations, and their increasing frequency with latitude, in the variability of the advection terms. Data are *not* filtered.



Figure S10: A) seasonal cycle of organic matter export at 90m averaged over the western North Atlantic (triangular region in Fig. 3A), and B) the ratio between organic matter export and NPP in this same region. Thin lines show daily means, and thick lines show 31-day running means of daily means. The passage of our simulated TCs are marked by the red bars.

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