

1 **The complex rupture evolution of the long and slow, tsunamigenic 2021 South**
2 **Sandwich Islands earthquake**

3
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17 **Abstract**

18 On August 21, 2021 a large earthquake occurred in the South Sandwich subduction zone, and the associated tsunami
19 was widely observed. In order to robustly analyse the detailed seismic source process of this long-source duration
20 (over 200 s) event occurring in a convexly shaped subduction zone, we applied the Potency Density Tensor Inversion
21 with a non-rectangular and non-planar source surface to the broadband teleseismic *P*-waves. This method allows us
22 to suitably reduce the effect of the time-increasing uncertainty of the Green's function and the effect of modelling
23 errors related to the fault geometry. We found the slip vectors of the 400 km long rupture area rotate clockwise to the
24 south, corresponding to the clockwise rotation of the trench strike. Our results reveal that the rupture propagated up-
25 dip to the shallow region, then propagated to the south-southeast along the trench, and stagnated for about 30 s at
26 around 130 km south-southeast from the epicentre. After the stagnation of the rupture front, the moment-rate
27 gradually increased with time, although a clear rupture area could not be identified for about 45 s. Afterwards the
28 rupture re-propagated south-southwest along the trench from the stagnation area. The slow rupture growth following
29 the stagnation of the fast rupture triggered a new fast rupture, which led to the 2021 South Sandwich Islands
30 earthquake having the typical characteristics of a tsunami earthquake, with a long rupture duration and a slow average
31 rupture front velocity.

32 Introduction

33 Understanding the rupture process of tsunami-generating earthquakes is important for assessing the
34 potential for future tsunami events. A “tsunami earthquake” is an event which has rupture characteristics that lead to
35 the generation of tsunamis that are unusually large given the event’s magnitude^{1,2}. It has been pointed out that tsunami
36 earthquakes have a slow rupture front velocity and long rupture duration^{3,4}. In the case of the 1992 Nicaragua
37 earthquake, one of the typical tsunami earthquakes, the rupture front velocity and the rupture duration were reported
38 to be 1.5 km/s, or less, and ~110 s, respectively²; it has been suggested that the slow slip along the plate interface
39 occurred due to accumulated soft subducted sediments². The mechanism of tsunami earthquakes are disputed in the
40 scientific literature: variable frictional properties on the plate interface⁵, off-fault rupture⁶, accretionary prism rupture
41 caused by rapid stress loading⁷, thermal pressurization⁸ and so on have been suggested being responsible for the
42 heterogeneous rupture behaviour, typically characterized by a slow and long rupture process.

43 On August 12, 2021, the 2021 South Sandwich Islands (SSI) earthquake struck offshore the South
44 Sandwich Islands, in the shallow part of the South Sandwich subduction zone^{9,10} classified as tectonic-erosion type<sup>11-
45 13</sup>. In this subduction zone, the South American plate is subducting under the Sandwich plate at a rate of 70–78
46 mm/yr, in a west-southwest direction¹⁴ (Fig. 1). The South Sandwich trench curves convexly to the east¹⁵ (Fig. 1),
47 and the South American plate is subducting obliquely in the southern region of this subduction zone. The bathymetric
48 data on the South American plate¹⁶ shows several chains of seamounts extending in east-southeast to west-northwest
49 directions. The aftershock distribution determined by the U.S. Geological Survey (USGS) extends from ~100 km
50 north to ~300 km south of the mainshock epicentre along the South Sandwich trench¹⁷.

51 The USGS identified three independent events corresponding to the 2021 SSI earthquake (Fig. 1). At
52 18:32:52 (UTC), a moment magnitude (M_W) 7.5 event occurred at 57.567° S, 25.032° W was followed 145 s later by
53 an M_W 8.1 event occurred about 90 km south of the epicentre of the first event; finally, 102 s later, an mb 6.7 event
54 occurred about 280 km south-southeast (SSE) of the epicentre of the first event (Fig. 1). The Global Centroid Moment
55 Tensor (GCMT) project^{18,19} reported only a M_W 8.3 earthquake with a centroid at 59.48° S, 24.34° W (Fig. 1). Jia et
56 al. (2022)⁹ applied the multiple point source inversion to the 2021 SSI earthquake using the low-frequency waveforms
57 (0.002–0.05 Hz) and identified five sub-events: two thrust sub-events in the first 50 s, one long-period thrust sub-
58 event for about 180 s with a centroid time of 90 s after the beginning of the rupture, and two thrust sub-events which
59 occurred 3 mins after the beginning of the rupture. Metz et al. (2022)¹⁰ carried out moment tensor inversion and finite
60 fault inversion (FFI), and reported four sub-events. They also performed a teleseismic back-projection (BP) analysis
61 and reported high-frequency (0.5–2 Hz) energy emitters migrating southward. Previous studies used long-period (20
62 s and longer) waveforms to perform source inversion and have shown that the 2021 SSI earthquake is composed of
63 multiple sub-events^{9,10} and has the characteristics of a typical tsunami earthquake, with slow rupture front velocity
64 and long rupture duration⁹. On the other hand, it also involves spatiotemporally complex high-frequency rupture
65 inferred from the BP analysis¹⁰. To understand the rupture process of the 2011 SSI earthquake, which has the
66 characteristics of a tsunami earthquake, it is essential to construct a detailed seismic source model that is able to
67 explain the broadband teleseismic body waves.

68 In general, it is difficult to stably estimate the detailed seismic source process of an earthquake with

69 complex fault geometries and a long rupture duration, such as the 2021 SSI earthquake. The influence of the
 70 uncertainty of the Green's functions becomes dominant as the source duration gets longer²⁰. A planar rectangle fault
 71 plane that is often adopted in the FFI may not necessarily be suitable for the actual curved and convex source fault,
 72 which can also increase the modelling uncertainty^{21,22}. The recently proposed Potency Density Tensor Inversion
 73 (PDTI) has made it possible to estimated detailed source processes including information on focal mechanism by
 74 reducing the effects of the uncertainty of the Green's functions^{21,23}. In this study we apply the PDTI with a non-
 75 rectangular and non-planar source surface to the observed teleseismic waveforms of the 2021 SSI earthquake to
 76 estimate its source process, including spatiotemporal changes in fault geometry and slip vector. We discuss the
 77 complex rupture propagation during the 2021 SSI earthquake and propose the possibility of partial slip partitioning.

78

79 **Method**

80 In general, the earthquake source can be expressed by the volume density of the moment-rate tensor²⁴. The
 81 moment-rate volume-density tensor is calculated by multiplying the potency-rate volume-density tensors by the
 82 rigidity. In the case of the fault dislocation, the potency-rate tensor is represented by a linear combination of five
 83 basis double-couple components^{21,25}. Therefore, the teleseismic waveform of an earthquake observed at station j is
 84 given by:

85

$$86 \quad u_j(t) = \sum_{q=1}^5 \int_V \dot{P}_q(\xi, t) * G_{qj}(\xi, t) dV + e_{bj}(t) \quad (1)$$

87

88 where V is the 3-D source area, \dot{P}_q is a potency-rate volume-density function of q th basis double-couple component,
 89 G_{qj} is a true Green's function of the q th basis double-couple component, ξ is the location in the V , e_{bj} is the
 90 background and instrumental noise and $*$ denotes the convolution operator in the time domain.

91

92 In the FFI of seismic waveforms, the potency volume-density tensor is approximated by the fault slip vector
 93 (potency areal-density vector) on a pre-assumed fault plane^{26,27}. The FFI method can resolve the spatiotemporal
 94 distribution of the potency density-rate (slip-rate) on the fixed fault plane, however it cannot identify fault ruptures
 95 with a different fault geometry.

96

97 In the multiple point source inversion, the moment volume-density tensor is approximated by a sum of the
 98 moment tensors at multiple point sources^{25,28,29}. This approach makes no assumptions about the fault plane and thus
 99 allows discussion of the possibility of slip occurring on an unknown fault. However, the multiple point source
 100 inversion method is unable to estimate the detailed rupture process for each sub-event. In addition, only long-period
 101 waves can be analysed due to the effects of the modelling errors caused by the simplification of the source model.

102

103 In the PDTI method, the potency-rate volume-density tensor is approximated by the potency-rate areal-
 104 density tensor on a pre-assumed model surface²¹. Then, eq. (1) becomes

105

$$103 \quad u_j(t) \approx \sum_{q=1}^5 \int_S \dot{P}'_q(\xi, t) * G_{qj}(\xi, t) dS + e_{bj}(t) \quad (2)$$

104

105 where S is 2-D model surface and \dot{P}'_q is a potency-rate areal-density function. As a result, the ruptures occurring
 106 on various faults are projected onto the model surface as the potency-rate areal-density tensor. In the FFI method, the
 107 number of degrees of freedom of the potency tensor is reduced from five to two by specifying the fault plane^{26,27}.
 108 However, in the PDTI method, the number of degrees of freedom of the potency tensor remains at five because the
 109 fault plane is never specified²¹. In other words, by increasing the degrees of freedom of the model, PDTI is capable
 110 of estimating a detailed seismic source process model including information on the fault geometry.

111 In general, a high degree of model freedom can lead to problems such as overfitting and unstable solutions.
 112 To avoid these problems, the PDTI method incorporates the modelling error derived from the uncertainty of the
 113 Green's function into the data covariance matrix, as proposed by Yagi and Fukahata (2011)²⁰, and applies Akaike's
 114 Bayesian Information Criterion (ABIC)³⁰⁻³² to reasonably evaluate the strength of the smoothing constraint. In this
 115 study, we use the latest version of PDTI, which introduces a time-adaptive smoothing constraint²³ to avoid the
 116 problem of over-smoothing caused by fixing the smoothing strength.

117 For eq. (2) to work successfully, it is necessary to reduce the modelling error by making the model surface
 118 S closer to the earthquake source faults. In this study, we analysed the source process by projecting the potency
 119 density tensor distribution onto a non-planar model created by referring to the slab geometry data in Slab2.0³³.

120

121 **Data and model parameterization**

122 For the PDTI, we used the vertical components of the teleseismic body P -waves from 47 stations at
 123 epicentral distances of 30°–100° downloaded from the SAGE Data Management Center (Fig. 2a). We selected
 124 stations, ensuring a high signal-to-noise ratio and better azimuthal coverage (Fig. 2a). We converted the observed
 125 waveforms to velocity waveforms by removing the instrument response and then decimated the signal by using a 1.2
 126 s sampling. We calculated the theoretical Green's functions using the method of Kikuchi and Kanamori²⁵ with a
 127 sampling rate of 0.1 s. We used the 1-D averaged structure of the CRUST1.0³⁴ in the source region to calculate the
 128 theoretical Green's function (see Supplementary Table S1). Other structure models, CRUST2.0³⁵ (see Supplementary
 129 Table S2) and ak135-F spherical average model^{36,37} (ak135-F) (see Supplementary Table S3), are also used to
 130 evaluate the robustness of the modelling (Fig. 4).

131 According to the USGS aftershock distribution¹⁷, the 2021 SSI earthquake is assumed to have occurred
 132 near the subducted and curved slab surface. To cover the aftershock distribution within three days of the event (Fig.
 133 1), we followed Yamashita et al. (2021)³⁸ and set up a non-rectangular model plane with the strike and dip angles of
 134 190° and 10°, respectively, with a maximum width of 150 km and length of 405 km, which is expanded using a
 135 bilinear B-spline with a knot spacing of 15 km in both the strike and dip directions. Then the depth of each knot is
 136 adjusted based on Slab2.0³³ to minimize the location error in the Green's function. We adopted a hypocentre (initial
 137 rupture point) by using the USGS-determined epicentre¹⁷ of the first event (57.567° S, 25.032° W) and the
 138 corresponding depth of the Slab2.0³³ at 21.6 km. To enable fast rupture propagation immediately after an initial break
 139 and to reduce the number of model parameters, a hypothetical rupture initiation time was set for each spatial node.
 140 The hypocentre and the surrounding nodes were set to be able to rupture immediately after the initial break, and the

141 other nodes were set to be able to rupture immediately after the hypothetical 2.0 km/s front had passed. Based on the
142 reported long duration of the 2021 SSI earthquake^{9,10}, the rupture duration at each knot is 180 s with a sampling
143 interval of 1.2 s, and the total rupture duration is 280 s.

144 In general, the PDTI method adjusts the waveform length to mitigate the influence of the *PP*-waves on the
145 inversion results. However, when this scheme is applied to an earthquake with a long rupture duration, such as the
146 2021 SSI earthquake, only a few observation points contribute to the results after 200 s from the origin time. The
147 source time function of the 2021 SSI earthquake obtained using the empirical Green's function^{39,40} shows that the
148 peak value of the moment rate in the 80 s after the origin time is sufficiently smaller than the peak value in the 200 s
149 after the origin time (see Supplementary Fig. S1). Therefore, in this study, the waveform from the arrival of the *P*-
150 wave to 80 s after the arrival of the *PP*-wave was used for inversion to stabilize the solution after 200 s from the
151 origin time.

152

153 Results

154 The spatial distribution of the potency density tensor, calculated by taking the temporal integration of
155 potency-rate density functions at each knot, is dominated by thrust focal mechanisms that are similar to the total
156 moment tensor (Figs. 1, 2c). The source area spans for about 400 km in length, with a maximum potency of 6.4 m at
157 about 210 km south-southeast of the epicentre (Fig. 2c). The estimated total moment tensor solution, obtained by the
158 spatiotemporal integration of the potency-rate density tensors, shows a thrust focal mechanism including a 25% non-
159 double-couple component (Fig. 1). The total seismic moment is 6.47×10^{21} N m (M_w 8.5). The slip vector of the nodal
160 plane closest to the slab surface geometry shows a tendency to rotate clockwise to the south-southwest (SSW) (Fig.
161 5a). The moment rate function has ten or more spikes with the largest peak of 6.71×10^{19} N m/s at 201 s (Fig. 2b).
162 The synthetic waveforms calculated from the obtained source process model well reproduce the observed waveforms,
163 including the data points not used for the inversion (Supplementary Fig. S2). In this study, we define Episodes 1–4
164 based on the moment rate function (Fig. 2d) and snapshots (Supplementary Figs. S3, S4, S5).

165 In Episode 1, the thrust-type rupture propagated mainly on the up-dip side of the hypocentre and then
166 reached to the trench at 30 s after the origin time (OT) (Fig. 3, Supplementary Fig. S3). Figure 2d shows the time
167 evolution of the rupture area projected along the strike of the model plane (190°). As shown in Fig. 2d, the thrust-
168 type rupture also propagated asymmetrically in the SSW and north-northeast (NNE) directions until OT+30 s. The
169 rupture front velocity in the strike direction is about 2 km/s. The SSW rupture started at OT+8 s, and the rupture front
170 velocity in the strike direction is about 3 km/s.

171 In Episode 2, from OT+30 s, the thrust-type rupture propagated unilaterally in southward direction along
172 the trench, at the rupture front velocity in the strike direction of about 3 km/s from OT+30 s to OT+70 s and then the
173 southward rupture propagation stagnated at around 130 km SSW from the epicentre (Figs. 2d, 3, Supplementary Fig.
174 S3). Within this episode, a large fault slip event, with a peak along the trench at about 130 km SSE of the epicentre,
175 begins at OT+55 s and continues for about 35 s.

176 In Episode 3, the moment-rate function increases gradually from about OT+100 s (Fig. 2d). This gradual

177 increase continues for about 45 s. We note it is difficult to identify a clear rupture propagation during Episode 3 due
178 to the relatively small potency-rate density (Fig. 2d, Supplementary Fig. S4).

179 Episode 4 begins with the rapid rise in the moment-rate function around OT+145 s (Fig. 2d). This initiation
180 coincides with the origin time of the secondary M_w 8.1 event, as determined by the USGS (Fig. 2d). The thrust-type
181 rupture propagated in the SW direction from around 100 km SSW of the epicentre and reached to the southern end
182 of the model area at about OT+200 s (Fig. 2d, Supplementary Fig. S5). The maximum peak of the moment rate
183 function equals 6.71×10^{19} N m/s, at 202 s (Fig. 2d).

184

185 **Model sensitivity tests**

186 In this study, we calculated the Green's functions using the 1-D structure model that averaged the
187 CRUST1.0³⁰ structure in the source region. In general, the longer the rupture duration, the more likely it is to be
188 affected by the uncertainty of the Green's functions²⁰. The effect of the Green's function uncertainty is introduced to
189 the data covariance matrix according to Yagi and Fukahata (2011)²⁰, but this approach cannot evaluate the effect of
190 non-Gaussian errors originating from the model setting⁴¹. Although the effect of non-Gaussian errors originating
191 from the setting of the structure model can be evaluated using a Bayesian multi-model estimation^{41–43}, it is not
192 practical to apply this approach to PDTI in terms of computational cost⁴⁴. Therefore, we examined the behaviour of
193 the solutions for two additional structure models: CRUST2.0³⁵ and ak135-F^{36,37}. In addition to the structure sensitivity
194 test, we also examined the behaviour of the solution when the model surface was set to a plane with a strike of 190°
195 and dip of 10°, rather than the non-planar surface referenced in Slab2.0³³. In this test, a 1-D structure that averaged
196 the CRUST1.0 model in the source region is used.

197 Figure 4 shows the variation in the solution when the structure or depth of the model surface is perturbed.
198 The timing of the peaks of the moment-rate function, which can be seen in Episode 1, 2 and 4, is slightly perturbed
199 depending on the model, but similar results are obtained for all models (Fig. 4a). This result is consistent with the
200 results obtained by the model sensitivity test in previous research²¹. On the other hand, the amplitude of the moment-
201 rate function is almost the same for all models in Episode 1, but the model dependence of amplitude increases after
202 Episode 2 (Fig. 4a). This result suggests that the effect of model uncertainty increases over time. However, the key
203 features of the moment-rate function, such as the gradual increase in the moment-rate during Episode 3 and the rapid
204 increase in the moment-rate at the start of Episode 4, are reproduced in all models (Fig. 4a). The asymmetric bilateral
205 rupture propagating in the NNE–SSW direction in Episode 1, the unidirectional rupture propagating in the SSW
206 direction in Episode 2, the absence of a clear event in Episode 3, and the large slip event occurring 200 km SSW
207 from the hypocentre in Episode 4 are all reproduced in the three alternative models (Fig. 4b, c, d). The rupture area
208 at the start of Episode 4 differs depending on the model, but a relatively large potency-rate density region near the
209 epicentre of the USGS' secondary M_w 8.1 event is commonly obtained in all three alternative models (Fig. 4b, c, d).
210 The moment magnitudes are perturbed by the model setting, and the values range from M_w 8.3–8.5 when using the
211 different structural models and model geometries.

212

213 **Discussion**

214 We analysed the broadband teleseismic body *P*-waves from the 2021 SSI earthquake and found that the
215 rupture process can be divided into four episodes (Fig 2b, d). In Episode 1, the rupture propagated from the
216 hypocentre to the shallow region while expanding in the NNE–SSW direction, reaching the South Sandwich trench
217 (Supplementary Fig. S3). In Episode 2, the rupture propagated in the SSE direction along the trench, but after OT+70
218 s, it remained stagnant around 130 km SSE of the epicentre (Supplementary Fig. S3). In Episode 3, a clear rupture
219 area cannot be identified (Supplementary Fig. S4), but the moment-rate increases gradually with time (Fig. 2b). In
220 Episode 4, the rupture propagated towards SSW, until OT+200 s (Supplementary Fig. S5). Our seismic source model
221 explains the broadband teleseismic body *P*-waveforms, and its characteristics are reproduced even when other three
222 different model settings are used (Fig. 4). Evaluating the smoothing strength using ABIC^{30–32} prevents overfitting
223 and makes it possible to estimate robust results. It is worth noting that the PDTI results are smoothed according to
224 the amount of information in the data. In the following, we discuss a series of fast and slow rupture evolution in
225 particular comparing with the previous BP analysis¹⁰, exhibiting the typical tsunami-earthquake characteristics but
226 having a more heterogeneous rupture evolution. We also discuss a possibility of slip partitioning based on our finding
227 of rotation of slip vector azimuths.

228 The PDTI method employed in this study reduces the effects of modelling errors by increasing the number
229 of degrees of freedom of fault geometry and rupture front^{21,45,46}, while the BP method resolves the rupture propagation
230 process by tracking the wave radiation sources without requiring detailed model setting^{47–49}. The results of the BP
231 method using high-frequency waveforms¹⁰ show that the rupture propagated in the SSE direction from the epicentre
232 until OT+100 s. The BP signal remained at a low level from OT+100 to 160 s, while the strong BP signals are
233 distributed about 250 km south of the epicentre after OT+200 s¹⁰. Considering that the high-frequency waveforms
234 are generally sensitive to disturbances of slip-rate and the rupture front propagation^{50,51}, the results of this study can
235 be compared with the BP result. The SSE rupture propagation until OT+100 s obtained by the BP method corresponds
236 to Episodes 1 and 2 of this study, the low-level BP signal from OT+100 to OT+145 s corresponds to Episode 3, and
237 the strong BP signal from OT+200 s corresponds to the large rupture event of Episode 4. After the start of Episode
238 4, the BP signal increases, but it is weaker than the other major BP signals¹⁰. This may suggest that the initial rupture
239 of Episode 4 accelerated smoothly.

240 The averaged rupture front velocity of the 2021 SSI earthquake is about 1.5 km/s, estimated from the
241 distance and time difference from the initial break and the major rupture event of Episode 4 (Fig. 2d), in agreement
242 with previous research^{9,10}. This slow rupture front velocity appears to reflect the characteristics of a tsunami
243 earthquake^{2,52}. However, our results also show that the rupture front propagated relatively fast in Episodes 1, 2 and
244 4. In Episode 1, the rupture front velocity in the strike direction is about 2–3 km/s (Fig. 2d). Considering that the
245 rupture is not only propagating in the strike direction, but also towards the shallow region, the actual rupture front
246 velocity may reach about 2.8–4.2 km/s. In Episode 2, the rupture front velocity in the strike direction is about 3 km/s
247 (Fig. 2d). Adjusting for discrepancy between the strike direction and the actual rupture propagation direction, the
248 rupture front velocity in Episode 2 is about 3.2 km/s. The smoothing strength increases over time because of the
249 increase with time of the uncertainty of the Green's function, thus it is difficult to trace the rupture front in Episode

250 4. If we take the USGS hypocentre of the M_W 8.1 event as the starting point of Episode 4, the rupture velocity in
251 Episode 4 is about 3.2 km/s (Fig. 2d). The rupture front velocity that can be identified separately for Episodes 1,2
252 and 4 is about 60–90% of the S-wave velocity, which is consistent with the rupture front velocity of regular
253 earthquakes^{53,54}.

254 In Episode 3, the moment-rate increases gradually for about 45 s, but a clear rupture area cannot be
255 identified (Fig. 2b, d). On the other hand, the BP analysis results show that the source of the wave corresponding to
256 Episode 3 is stagnating at a point about 110 km south-southeast of the hypocentre¹⁰. The BP analysis results suggest
257 that the rupture area of Episode 3 is distributed around the rupture stagnation area of Episode 2 and the rupture
258 initiation area of Episode 4. The gradual increase in the moment rate reflects the gradual increase in the rupture area
259 and/or the gradual acceleration of fault slip-rate. Therefore, the shear stress loading due to the gradual expansion of
260 the rupture area and/or the gradual acceleration of slip-rate may trigger the rupture of Episode 4. Notably, the long-
261 period, slow rupture event reported by previous studies^{9,10} coincides with the slow rupture growth observed during
262 Episode 3 (Supplementary Fig. S6).

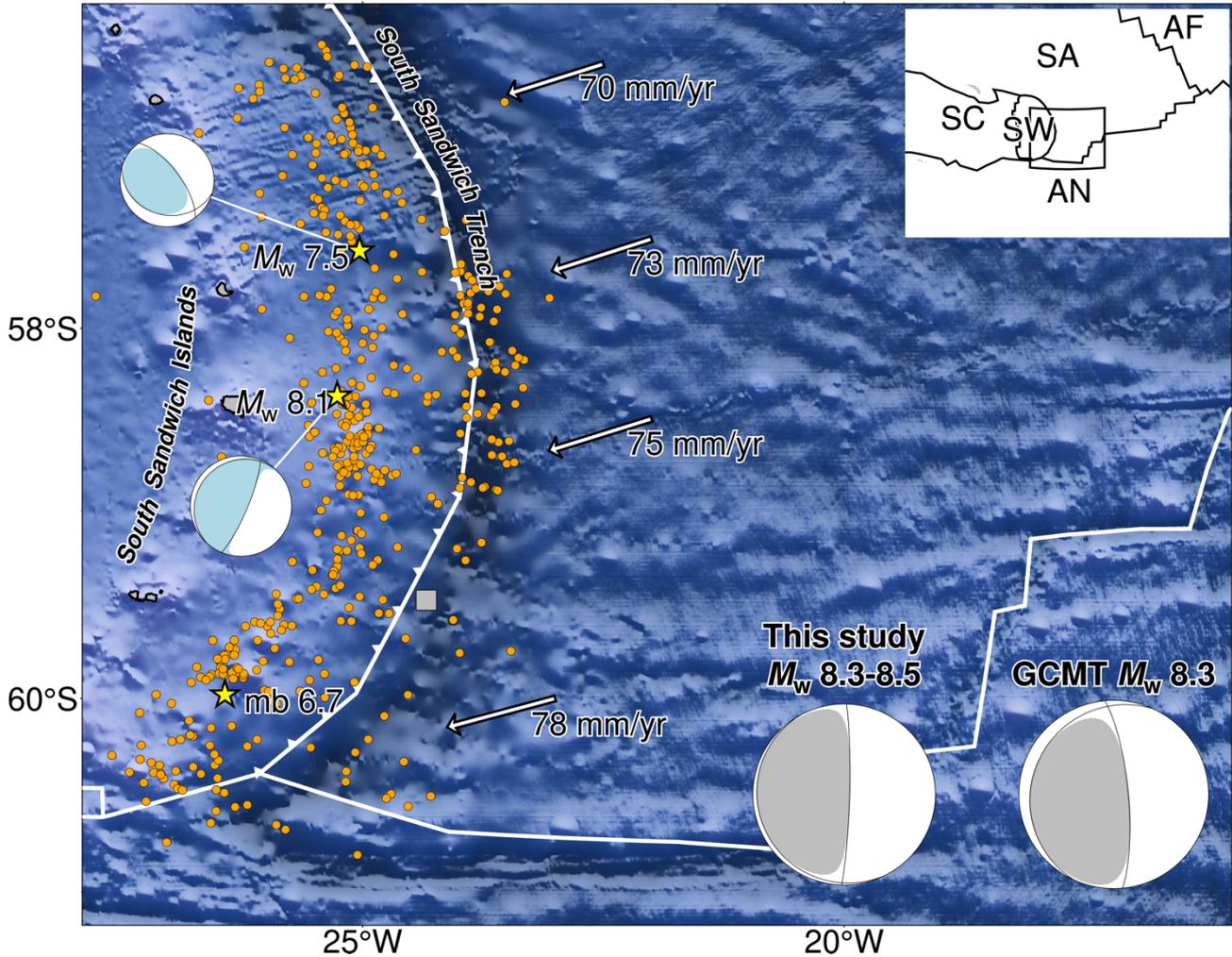
263 In oblique subduction zones, slip partitioning often occurs between dip-slip interplate faults and strike-slip
264 faults in the continental crust^{55,56}. When slip partitioning occurs, the slip vector of the dip-slip fault can point in a
265 direction between the relative plate subduction direction and the direction normal to the trench axis^{55,57,58}. Figure 5a
266 shows the spatial variation of the slip vector of the nodal plane closest to the slab surface geometry. The slip vectors
267 obtained in this study show a tendency to rotate clockwise as the slip propagates south-southwest in response to the
268 change in the trench strike. In the rupture region spanning the latitudes 57°S–58.4°S, the slip vector azimuth becomes
269 larger towards the south, while at latitudes 58°S–58.4°S it becomes consistent with the trench-normal direction (Fig.
270 5a). Between latitudes 58.4°S–59.3°S, the slip vector azimuth is about 15° greater than the direction of plate
271 subduction, and deviates from the trench-normal direction (Fig. 5a). From 59.3°S to 60.2°S, the slip vector azimuth
272 becomes larger towards the south, and from 59.8°S–60.2°S it approaches the trench-normal direction again (Fig. 5a).
273 A similar result can be observed using the background seismicity (earthquakes of $M_W > 5.5$) from the GCMT
274 solutions: the dip-slip vector azimuths seem to align with the trench-normal direction at latitudes 57°S–58°S, while
275 the azimuths tend to orient between the trench-normal direction and subduction direction at latitudes 58°S–60°S (Fig.
276 5b). The spatial variation of the dip-slip vectors, including our results, suggests that slip partitioning^{55,57,58} may play
277 a role in the convex South Sandwich subduction zone.

278

279 **Conclusion**

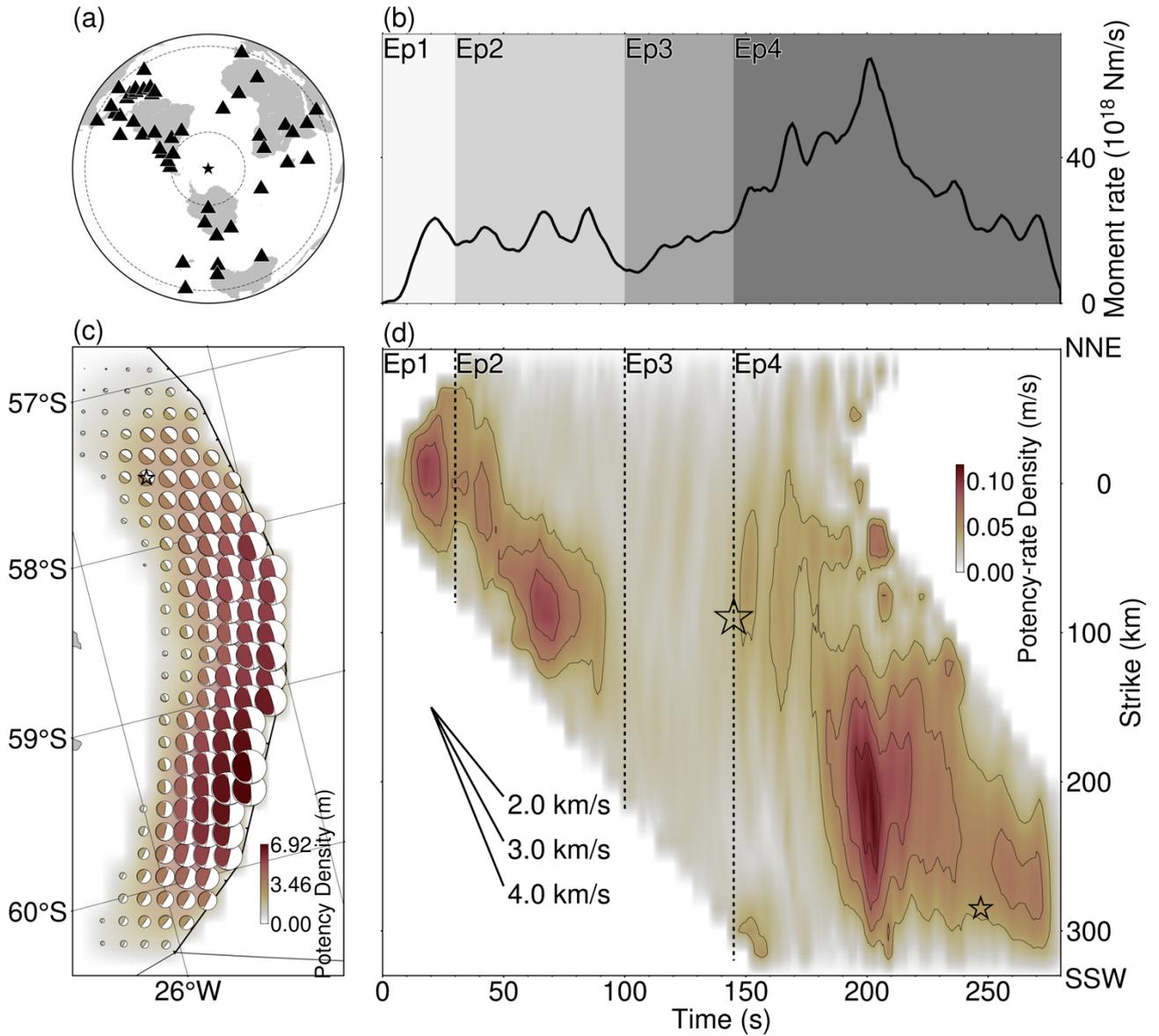
280 We estimated the spatiotemporal evolution of rupture propagation during the 2021 South Sandwich Islands
281 earthquake by using the PDTI method. The results of the model sensitivity test using the three alternative models
282 suggest that the rupture propagation process obtained in this study is stable. The rupture can be divided into four
283 episodes. In Episode 1 (0-30 s), the fast rupture propagated in a shallow direction, in Episode 2 (35-100 s) the fast
284 rupture propagated southeast along the trench and then stagnates, in Episode 3 (100-145 s) the rupture slowly
285 expanded around the stagnant area, and in Episode 4 (145-280 s), the fast rupture propagated in a south-southwest

286 direction along the trench. The results of this study demonstrate that the characteristics of tsunami earthquakes, such
287 as long rupture duration and a slow average rupture front velocity, are observed in the case of the 2021 South
288 Sandwich Islands earthquakes as a complex process: slow rupture growth after stagnation of the initial fast rupture
289 and the triggering of a final fast rupture. The spatial variation of the slip vectors shows a tendency to rotate clockwise
290 to the south, corresponding to the clockwise rotation of the strike of the convex South Sandwich trench to the south.
291 Our detailed source model successfully explains the broadband tele-seismic *P*-waves and can provide a unified
292 explanation of the previous back-projection result and the multiple point source inversion results of the long-period
293 waveforms, shedding light on the highly heterogeneous rupture process of the 2021 South Sandwich Islands
294 earthquake.



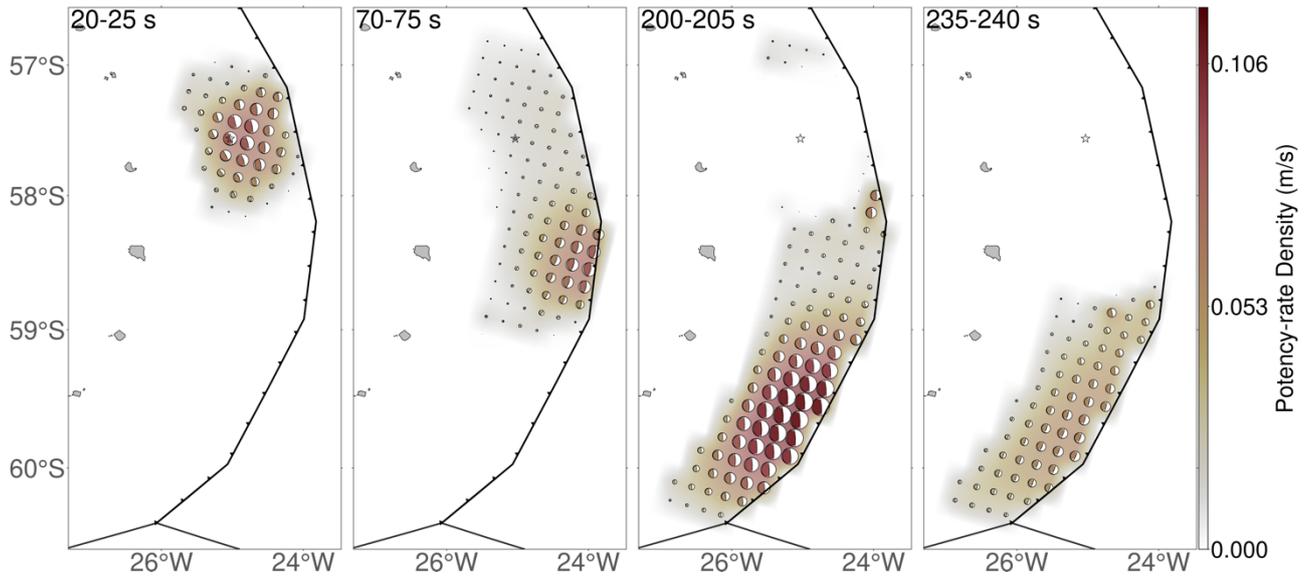
295

296 **Figure 1.** Summary of the study region. Three yellow stars are the epicentres of M_w 7.5, M_w 8.1 and mb 6.7 events
 297 reported by the USGS NEIC¹⁷, and two light blue beachballs are the moment tensors of M_w 7.5 and M_w 8.1 events.
 298 The grey square is the centroid location reported by the Global Centroid Moment Tensor (GCMT) project^{18,19}, and
 299 the grey beachballs on the bottom right corner are the GCMT solution and the total moment tensor solution obtained
 300 in this study. The orange dots are the aftershocks from 2021-08-12 18:32:52 to 2021-8-15 18:32:52 (within three
 301 days after the M_w 7.5 events) detected by the USGS NEIC. The vector shows the subduction velocity and azimuth
 302 of South American plate¹⁴. The white line indicates the plate boundary¹⁵, and the white front is the South Sandwich
 303 trench. The shaded background bathymetry is derived from SYN BATH¹⁶. The inset map shows the plates around the
 304 study region: Sandwich (SW), Scotia (SC), South American (SA), Antarctic (AN) and Africa (AF) plates¹⁵. The black
 305 lines are the plate boundaries¹⁵ and the box outlines the extent of this figure.

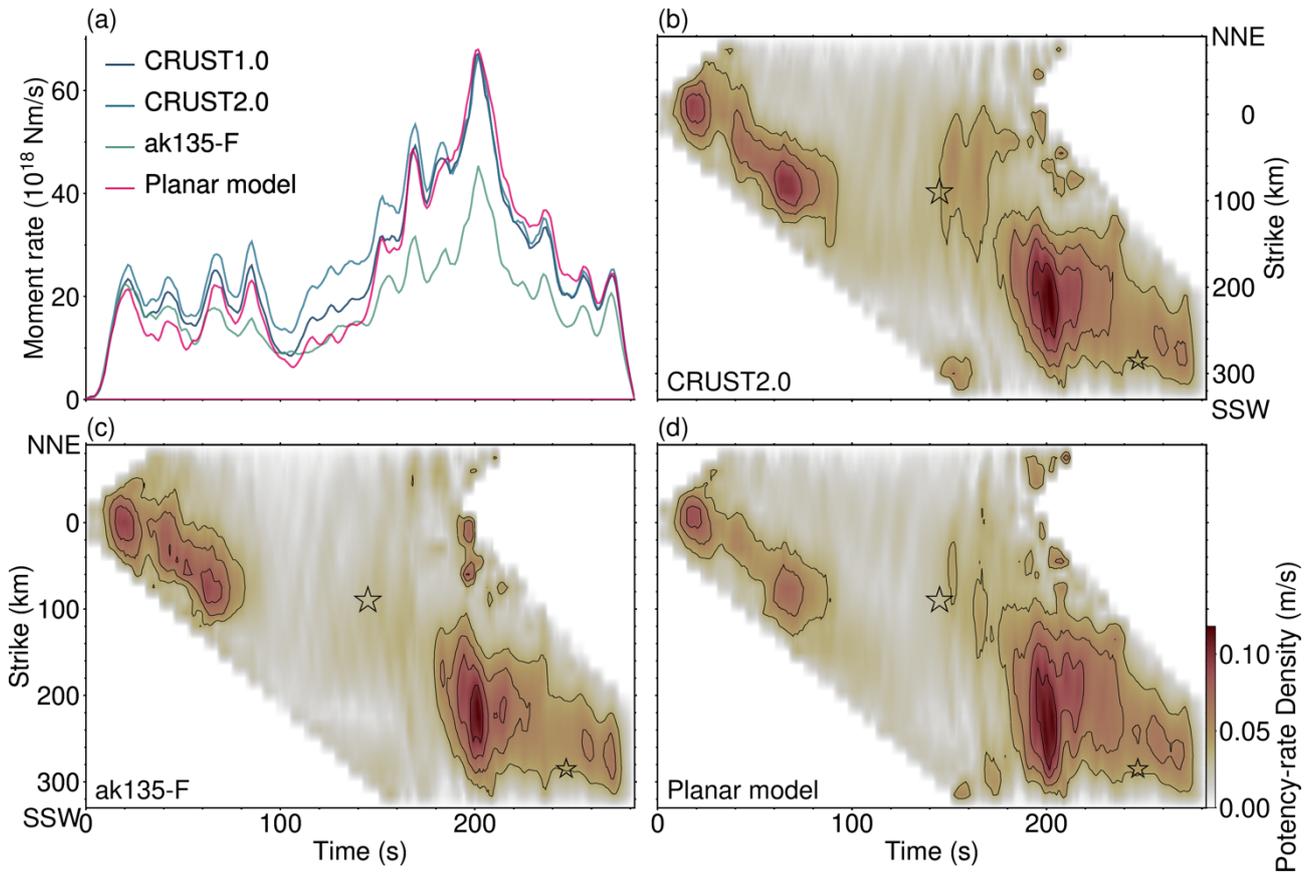


306

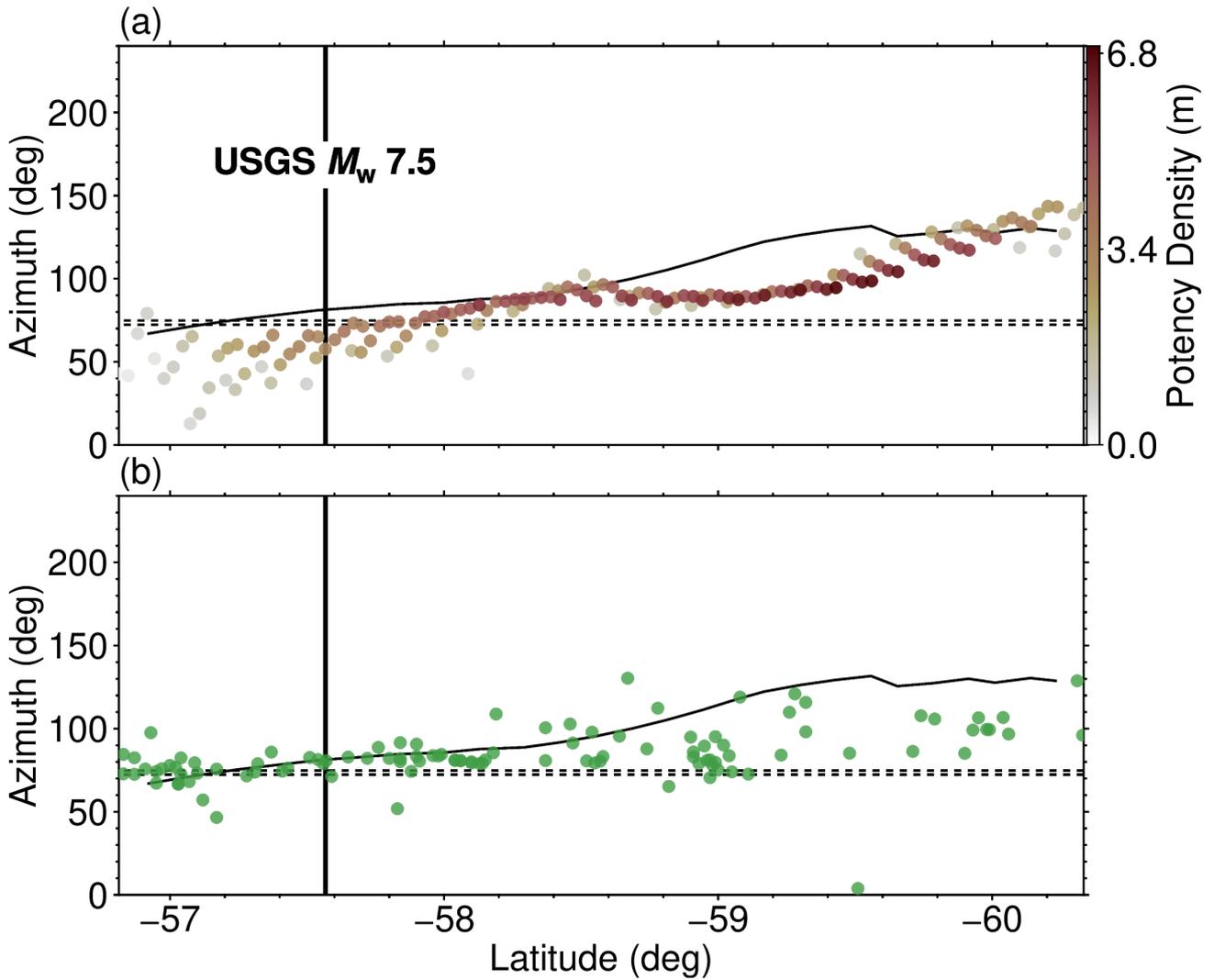
307 **Figure 1.** Summary of inversion results. (a) The station distribution used in the inversion, projected on the azimuthal
 308 equidistant map. The dotted circles show epicentral distances of 30° and 100° . (b) Moment rate function. The
 309 background grey colours show the time range of four episodes defined in this study. (c) Map projection of the potency
 310 density tensor distribution on the non-planar model. The star and black line indicate the epicenter¹⁷ and plate
 311 boundary¹⁵, respectively. (d) Potency-rate density evolution projected along strike. The contours of 0.04, 0.06, 0.08
 312 and 0.10 m/s are shown as black lines. The vertical axis is distance from the epicentre. The large and small stars are
 313 the epicentres of the M_w 8.1 and m_b 6.7 events detected by the USGS¹⁷.



314
315 **Figure 2.** Selected snapshots of the potency-rate density distribution (Fig. 2d). The beachballs show the potency-rate
316 density tensors. The star and black line indicate the epicentre¹⁷ and plate boundary¹⁵, respectively.



317
 318 **Figure 4.** Summary of results from model sensitivity test. (a) The moment-rate functions when the Green's function
 319 is calculated using two different structure models (CRUST2.0³⁵ and ak135-F^{36,37}; see Supplementary Tables S2 and
 320 S3) and when the potency-rate density tensor distribution is projected onto a planar model with a strike of 190° and
 321 dip of 10° . (b, c) Potency-rate density evolution obtained with CRUST2.0³⁵ (see Supplementary Table S2) and ak135-
 322 F^{36,37} (see Supplementary Table S3) structure models for calculating the Green's function, respectively. (d) Potency-
 323 rate density evolution obtained using the planar model. The large and small stars plotted in each figure are the
 324 epicentres of the M_w 8.1 and m_b 6.7 events, respectively detected by the USGS¹⁷.



325

326 **Figure 5.** Spatial variation of slip vectors. (a) The spatial variation of total slip vectors for each knot obtained by
 327 PDTI. The slip vectors are calculated for the preferred nodal plane. (b) The spatial variation of slip vectors for GCMT
 328 solutions^{18,19} of background events larger than M_w 5.5. In both (a) and (b), the black solid line is the direction
 329 perpendicular to the strike of the subducting slab near the trench. The two black dashed lines indicate the direction
 330 of plate subduction¹⁴, the upper line and the lower line are corresponding to the most northerly and to the most
 331 southerly subduction vector azimuths denoted as arrows in Fig. 1, respectively. The vertical line indicates the
 332 rupture point.

333 **Data availability**

334 All seismic data were downloaded from IRIS Wilber 3 system (<https://ds.iris.edu/wilber3/>) or IRIS Web Services
335 (<https://service.iris.edu/>), including the following station networks: (1) Caribbean Network (CU;
336 <https://doi.org/10.7914/SN/CU>); (2) GEOSCOPE (G; <https://doi.org/10.18715/GEOSCOPE.G>); (3) the Global
337 Telemetered Seismograph Network (USAF/USGS) (GTSN) (GT; <https://doi.org/10.7914/SN/GT>); (4) the IRIS/IDA
338 Seismic Network (II; <https://doi.org/10.7914/SN/II>) and (5) the Global Seismograph Network
339 (IU; <https://doi.org/10.7914/SN/IU>). The moment tensor solutions of the Global Centroid Moment Tensor (GCMT)
340 catalog are available through the website <https://www.globalcmt.org/CMTsearch.html>. The CRUST1.0³⁴,
341 CRUST2.0³⁵ and ak135-F^{36,37} are available through the websites <https://igppweb.ucsd.edu/~gabi/crust1.html>,
342 <https://igppweb.ucsd.edu/~gabi/crust2.html> and <http://rses.anu.edu.au/seismology/ak135/ak135f.html>, respectively.
343 The slab2 model is available at <https://doi.org/10.5066/F7PV6JNV>. Plate motion data are obtained from Plate Motion
344 Calculator (http://ofgs.aori.u-tokyo.ac.jp/~okino/platecalc_new.html). We used Generic Mapping Tools (v6.5.0)⁵⁹
345 (<https://docs.generic-mapping-tools.org/latest/index.html>).

346

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348 This work was supported by JSPS Grant-in-Aid for Scientific Research (C) 22K03751. The facilities of IRIS Data
349 Services and specifically the IRIS Data Management Center were used for access to waveforms, related metadata,
350 and/or derived products used in this study. IRIS Data Services are funded through the Seismological Facilities for
351 the Advancement of Geoscience (SAGE) Award of the National Science Foundation under Cooperative Support
352 Agreement EAR-1851048. We used FPSPACK⁶⁰ software to handle the focal mechanism obtained in the inversion.
353 The ray parameter and travel-time are calculated with the TauP Toolkit⁶¹ for the calculation of the Green's function.
354 All the figures were generated with Generic Mapping Tools (v6.5.0)⁵⁹.

355

356 **Contributions**

357 R.Y., Y.Y. and R.O. conceptualized this study. R.Y. and Y.Y. processed data and carried out the inversion analysis.
358 R.Y. generated the figures. R.Y., Y.Y., R.O. and B.E. interpreted and discussed the results and data. R.Y. wrote the
359 original manuscript. Y.Y., R.O. and B.E. substantially revised the manuscript. All authors approved the revised
360 manuscript.

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- 492

Supporting Information for
**The complex rupture evolution behind the long and slow, tsunamigenic 2021
South Sandwich Islands earthquake**

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Figure S6. Comparison of the potency-rate density evolution projected along strike and sub-events identified by previous studies.

Table S1. 1-D structure around the source used for calculating the Green's function, based on CRUST1.0 model¹

V_P (km/s)	V_S (km/s)	Density (10^3 kg/m ³)	Thickness (km)
1.50	0.00	1.02	3.80
5.75	3.20	2.70	3.20
7.10	4.05	3.05	4.80
8.04	4.47	3.32	- (below moho)

Table S2. 1-D structure around the source used for calculating the Green's function, based on CRUST2.0 model²

V_P (km/s)	V_S (km/s)	Density (10^3 kg/m ³)	Thickness (km)
1.50	0.00	1.02	3.61
5.00	2.50	2.60	1.70
6.60	3.65	2.90	2.30
7.10	3.90	3.05	2.50
8.15	4.65	3.35	- (below moho)

Table S3. 1-D structure around the source used for calculating the Green's function, based on ak135-F model^{3,4}

V_P (km/s)	V_S (km/s)	Density (10^3 kg/m ³)	Thickness (km)
1.45	0.00	1.02	3.00
5.80	3.20	2.60	6.70
6.80	3.90	2.92	8.00
8.04	4.48	3.64	- (below moho)

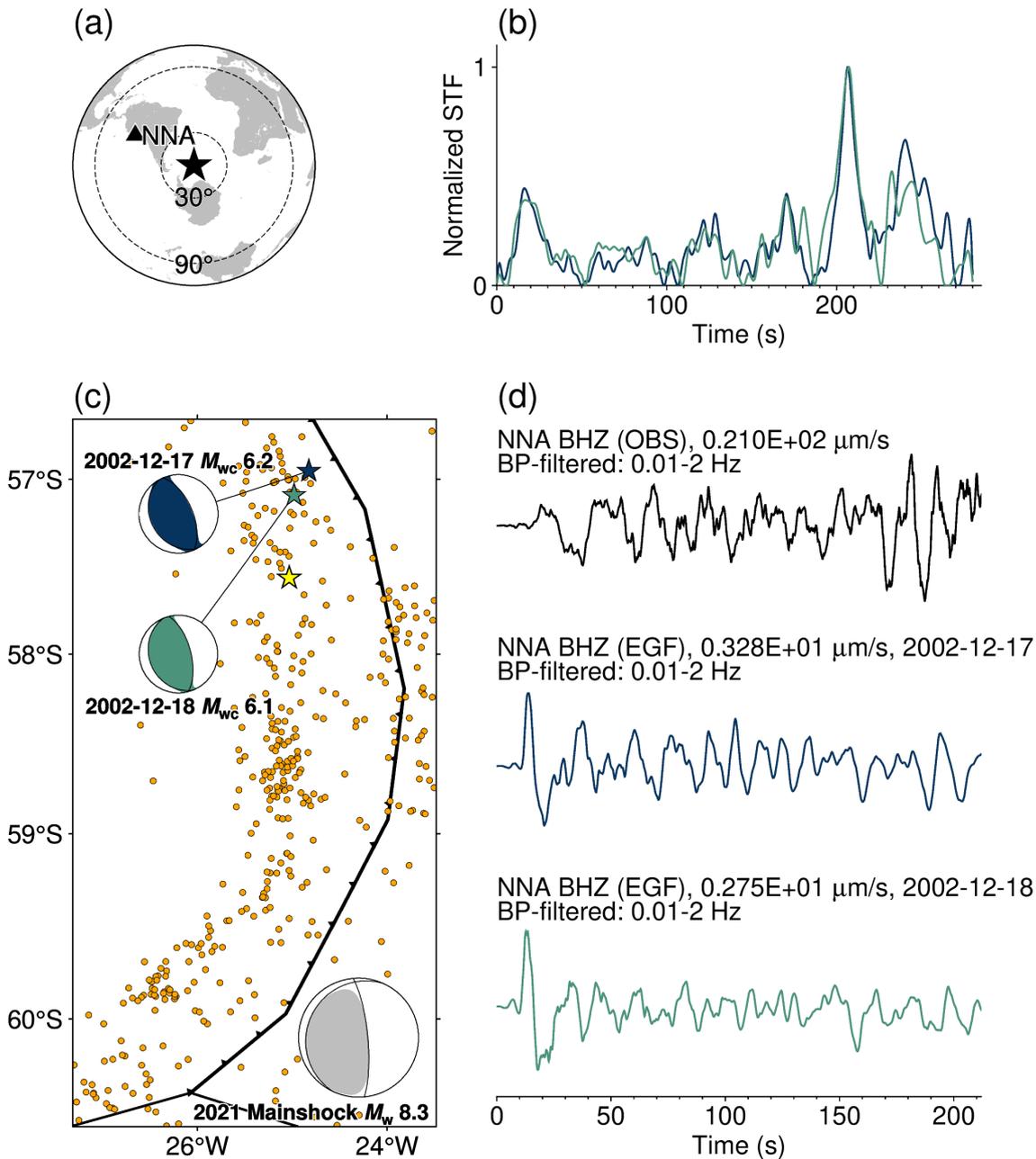


Figure S1. Summary of the empirical Green's function (EGF) analysis. (a) An azimuthal equidistant projection of the location of NNA station. (b) The normalized source time function obtained by using observed waveforms of the 2002-12-17 (blue) and the 2002-12-18 (green) event as EGFs. A band-pass filter of 0.01 to 2 Hz was applied to the observed waveform and EGFs. (c) Map projection of the GCMT^{5,6} solution of the mainshock (grey beachball) and the USGS⁷ two events used as EGFs. The blue and green stars are the epicentres of 2002-12-17 and 2002-12-18 event, respectively. The yellow star is the epicentre of the first event of the 2021 SSI earthquake. Orange dots are aftershocks that occurred within 3 days of the mainshock detected by the USGS. The black lines are the plate boundary⁸. (d) Observed waveform of the 2021 SSI earthquake (upper) and observed waveforms of the 2002-12-17 event (blue) and the 2002-12-18 event (green) used as EGFs. Both observed waveform and EGFs are band-pass (0.01–2 Hz) filtered. This figure was made with Generic Mapping Tools (v 6.5.0) software⁹.

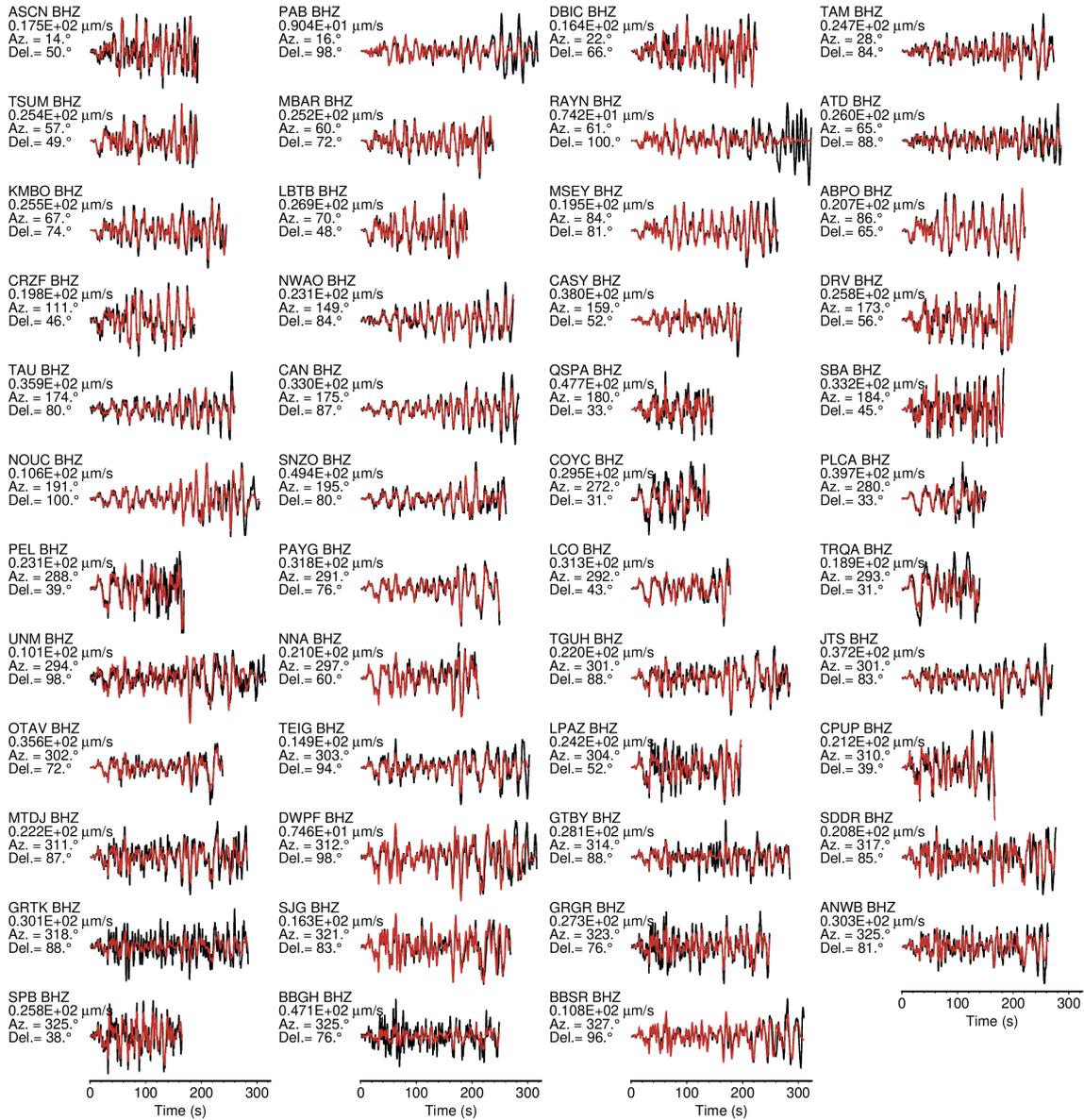


Figure S2. Comparison of observed waveforms (black) with Synthetic waveforms (red). Each panel is labelled with the station name, maximum amplitude, azimuth (Azi.), and epicentral distance (Del.) from the mainshock. These waveforms were resampled to 20 Hz for plotting. This figure was made with the Generic Mapping Tools (v 6.5.0) software⁹.

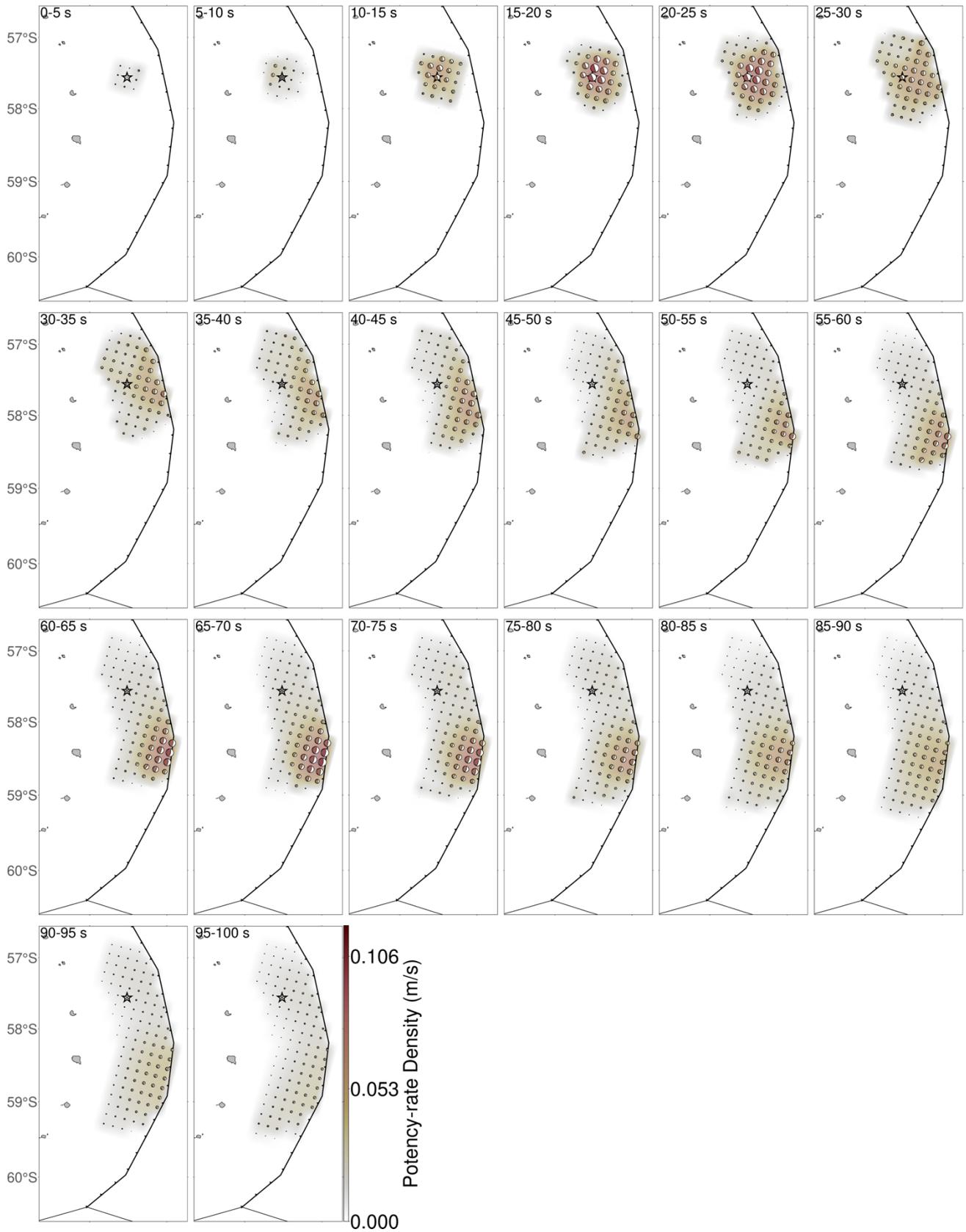


Figure S3. Full snapshots of potency-rate density evolution for Episode 1 and 2. The beachballs show the potency-rate density tensors. This figure was made with Generic Mapping Tools (v 6.5.0) software⁹.

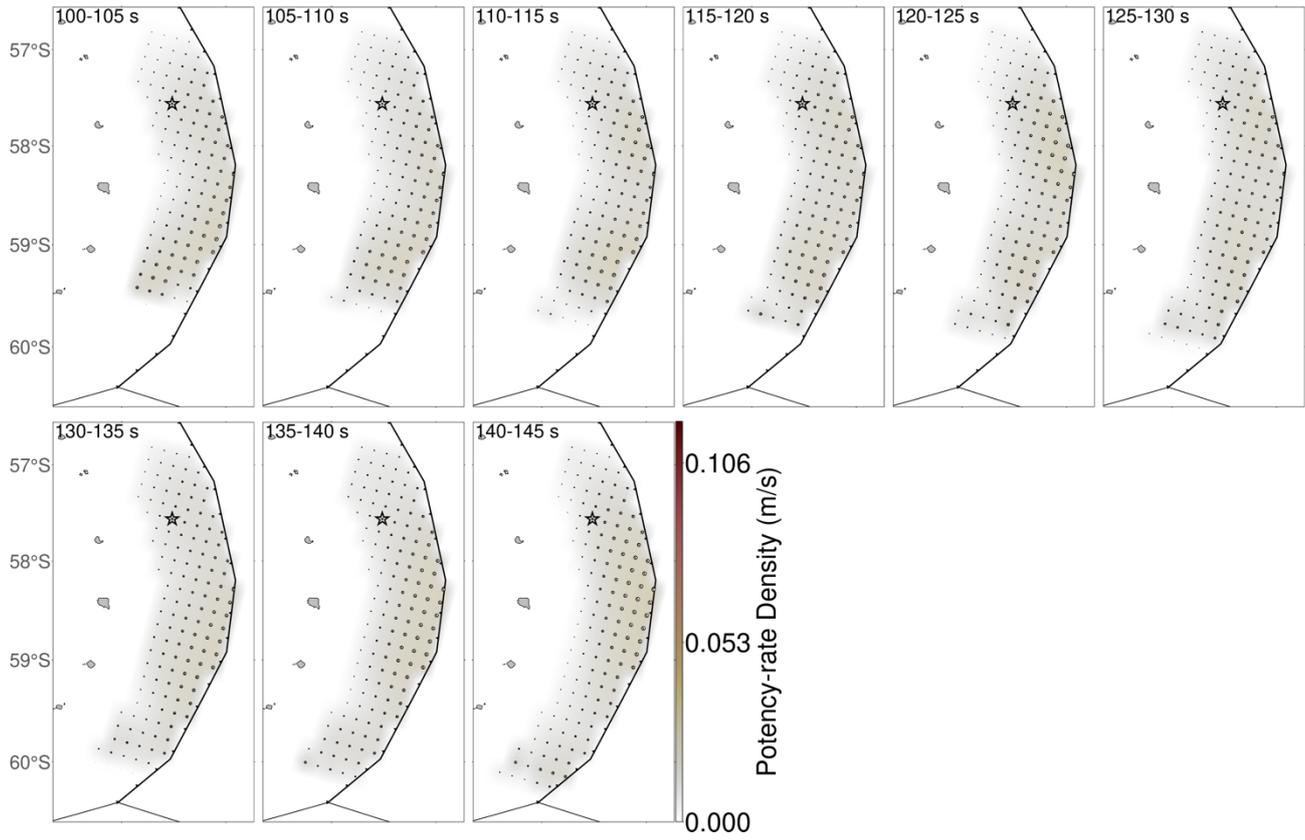


Figure S4. Full snapshots of potency-rate density evolution for Episode 3. The beachballs show the potency-rate density tensors. This figure was made with Generic Mapping Tools (v 6.5.0) software⁹.

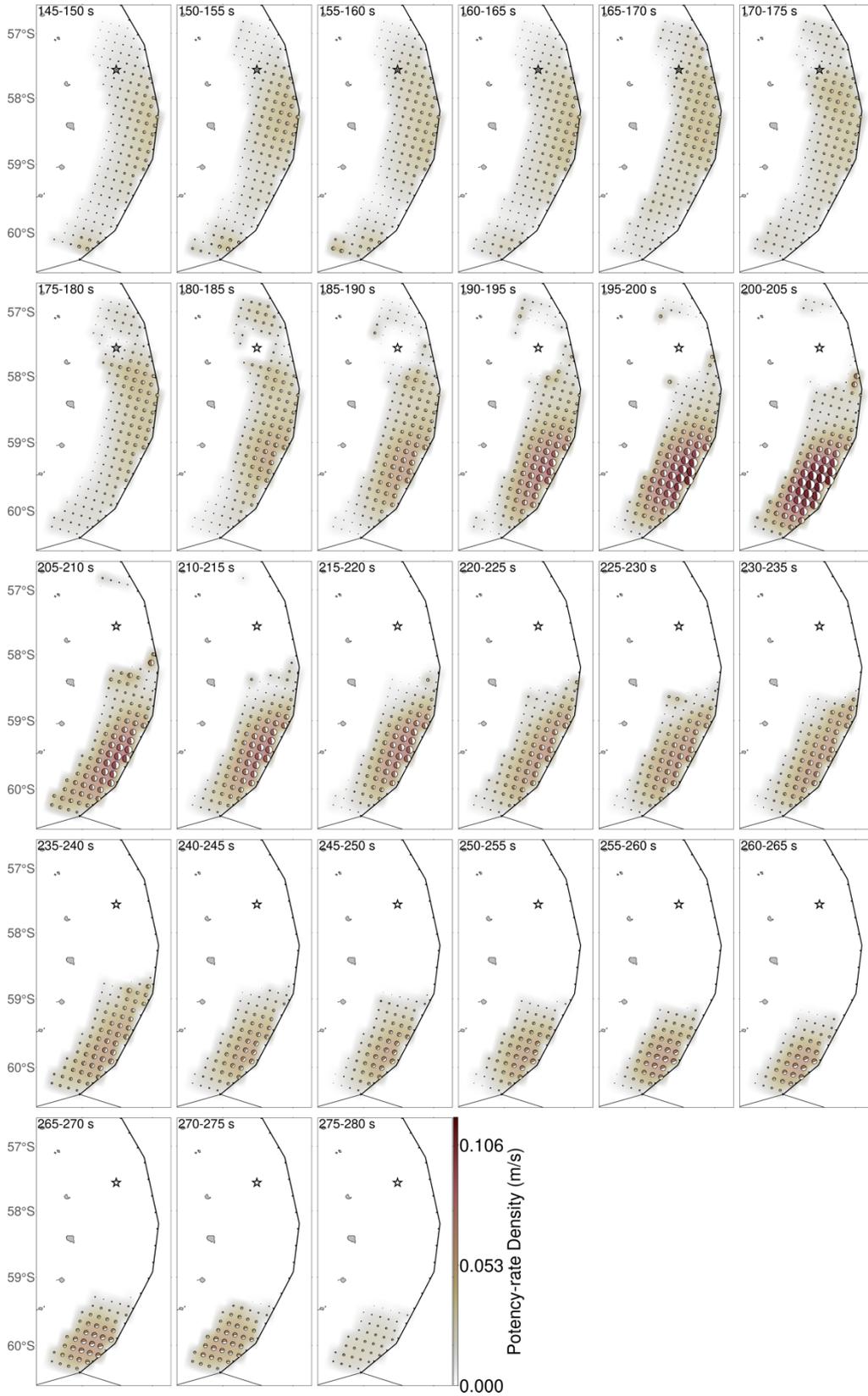


Figure S5. Full snapshots of potency-rate density evolution for Episode 4. The beachballs show the potency-rate density tensors. This figure was made with Generic Mapping Tools (v 6.5.0) software⁹.

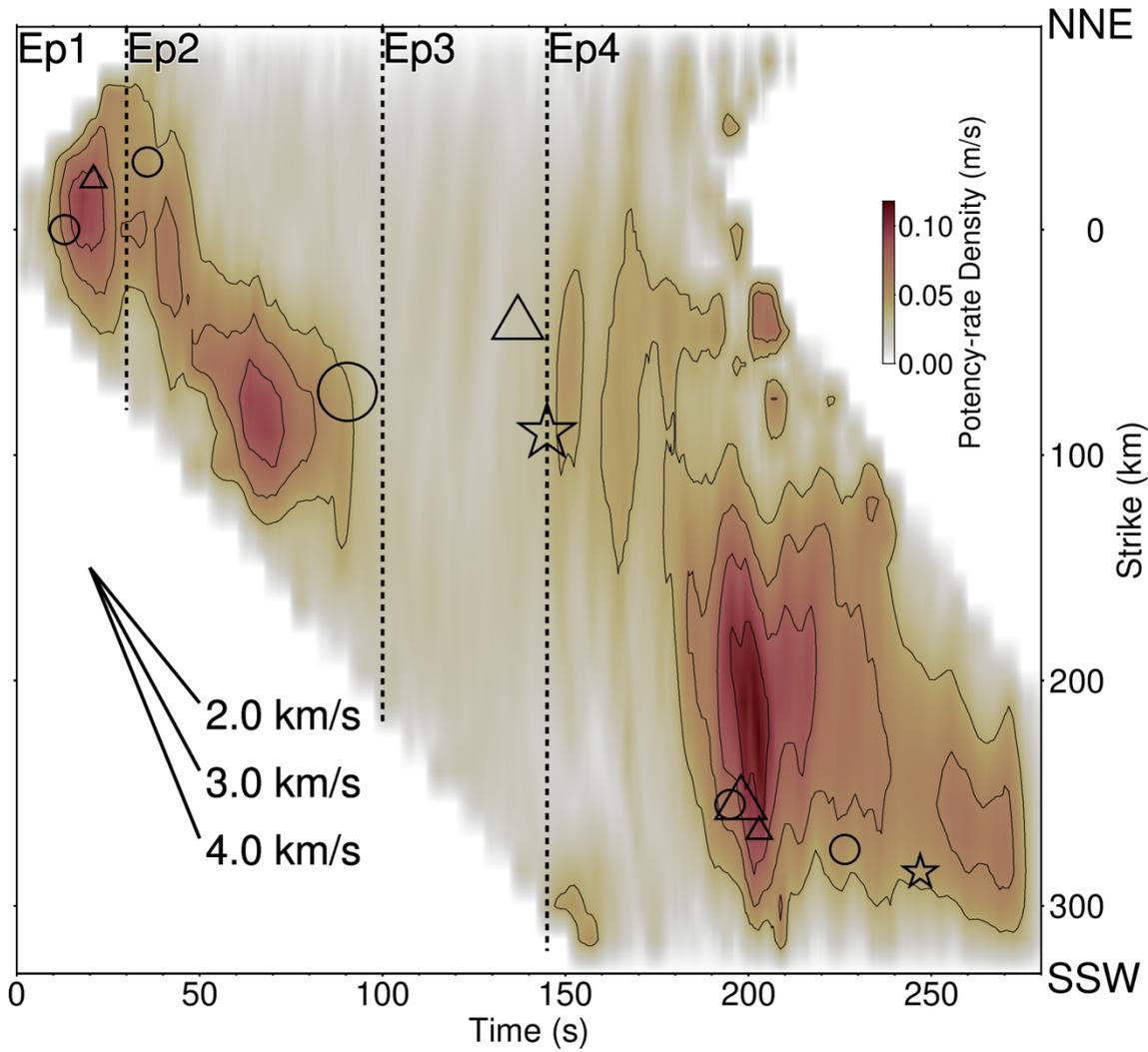


Figure S6

Comparison of the potency-rate density evolution projected along strike and sub-events identified by previous studies^{10,11}. The circles indicate the centroid of multiple sub-events reported by Jia et al.¹⁰, and the triangles indicate locations of the centroid moment tensor reported by Metz et al.¹¹ The bigger plots are longer events which have the duration longer than 100 s. The large and small star are the epicentres of the M_w 8.1 and mb 6.7 events detected by the USGS⁷, respectively. The contour is the same as Fig. 2d. This figure was made with Generic Mapping Tools (v 6.5.0) software⁹.

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