



Peer review status:

This is a non-peer-reviewed preprint submitted to EarthArXiv.

Learning wave scattering properties from seismograms

Reza Esfahani¹, Mirko Bracale¹, Léonard Seydoux², Michel Campillo¹

¹Institut des sciences de la Terre, Université Grenoble Alpes, Grenoble, France

²Institut de physique du globe de Paris, UMR CNRS 7154, Université Paris-Cité, Paris, France

Key Points:

- We introduce an unsupervised learning approach to directly evaluate near-source heterogeneity from seismograms
- We propose a proxy termed *local complexity* to quantify near-source structural heterogeneity.
- We demonstrate that our proposed model accurately captures statistical properties of the medium near seismic sources

Corresponding author: R. Esfahani:, dokhtdor@univ-grenoble-alpes.fr

Abstract

Heterogeneities in the Earth’s crust scatter seismic waves at many scales, trapping seismic energy and producing coda waves that encode valuable information on geological structures. In regions such as volcanoes and fault systems, analyzing coda waves is essential for characterizing non-uniform subsurface heterogeneity, improving interpretation and seismic imaging. Here, we apply unsupervised learning to infer properties directly from seismograms. We simulate 7,800 source-receiver seismograms within a realistic physics-based volcanic model of a magmatic plumbing system with complex interactions between dykes and sills. Recent studies suggest that the spectral characteristics of these synthetic seismograms are controlled by the partial resonance of multiply scattered waves. We leverage a deep scattering transform to extract robust, time-invariant representations of seismograms recorded with multiple stations, and use a manifold learning algorithm to visualize and analyze patterns in the scattering coefficients. By examining the connections in the embedded manifold, we reveal how local medium complexity influences recorded wavefields. Our results demonstrate that the proposed method effectively captures local resonant frequency and modulation induced by heterogeneous structures near the sources. We show that the statistical properties of the medium align with the estimated local complexities derived from seismic signals. By analyzing complete seismograms in a data-driven way, our method enhances subsurface heterogeneity characterization and offers a promising approach for improving the space-time monitoring in highly heterogeneous regions.

Plain Language Summary

The Earth’s crust in volcanic and fault zones is made up of a mix of different materials, creating a highly irregular underground structure. Near earthquake and volcanic sources, rock damage and underground fluids amplify certain frequencies in seismic waves, providing insights about subsurface properties. In this study, we introduce a machine learning approach to analyze these signals and reveal underground complexity without needing a large training dataset. Using numerical simulations, we test how well this method can identify different materials and their effects on seismic waves. By recognizing patterns in the data, our approach helps map variations in the subsurface and understand how they influence seismic signatures. Our results show that this technique effectively captures key properties of underground structures and could improve how we monitor earthquakes and volcanic activity over time. While challenges remain in adapting it to different geological settings, this

44 method offers a promising, data-driven way to study the Earth’s interior and track changes
45 in its structure.

46 **1 Introduction**

47 Scattered seismic waves are the results of small-scale heterogeneity in the Earth’s crust,
48 generating single and multiple scattering waves called coda waves (Aki & Chouet, 1975).
49 In complex geological environments such as volcanoes and fault systems, seismic wavefields
50 recorded at a surface are mixed with source effects, path effects, and effects of subsurface
51 heterogeneities. Recent studies have shown that in these environments, resonant frequencies
52 are controlled not only by the source but also by the heterogeneous media surrounding
53 it (Ben-Zion & Sammis, 2003; Barajas et al., 2023) (Bracale et al. [xxx]). Identifying the
54 local properties of seismic wavefield in a medium and linking resonant frequencies to local
55 properties can enhance our understanding of geological structures in active regions with
56 high heterogeneity.

57 Seismic wavefields can be considered a combination of ballistic waves, which propa-
58 gate straightforwardly, and scattered waves, which follow the ballistic waves and produce
59 the long-lasting coda. The scattered waves consist of combinations of single and multiple
60 scattering arrivals, including conversions between P and S waves. Several approaches have
61 been developed to assess the properties of heterogeneous media in the Earth’s interior by
62 analyzing coda regimes in the time-space domain (Aki & Chouet, 1975). First, under single-
63 scattering regime assumption, the coda can be modeled as a superposition of waves scattered
64 once within the medium. However, this oversimplification often fails in cases with strong
65 multiple scattering (e.g. Margerin, 2005; Hennino et al., 2001). As another end-member
66 model, the diffusion model approximates the end part of coda wave energy decay by treating
67 seismic energy as a random walk through the medium (Weaver, 1990; Margerin et al., 2009).
68 To bridge the gap between these two models, Wu and Aki (1988) introduced the radiative
69 transfer equation in seismology to model the decay of energy over time and space. This
70 approach is capable of modeling the seismic energy density in all seismic scattering regimes,
71 including the single scattering, multiple scattering, and diffusive regimes. Nevertheless, it
72 is important to note that these models make use of physically-meaningful quantities, such
73 as the scattering mean free paths (e.g. Margerin, 2005), that can be rigorously defined
74 statistically under the assumption of a uniform distribution of scatterers in the media.

75 In media with a non-uniform distribution of scatterers, such as volcanic environments,
76 studies have demonstrated a significant contribution of multiply scattered waves with a
77 small mean free path (Del Pezzo et al., 1996; Yamamoto & Sato, 2010a). The high level of
78 heterogeneity can even lead to rapid transition toward diffusion regime, producing spindle-
79 like envelopes with weak P -wave onsets and prolonged coda waves (Yamamoto & Sato,
80 2010b; Zieger et al., 2016). In addition, it has been suggested that some seismic signatures,
81 like volcanic tremors and long-period earthquakes (B. Chouet, 1992), can have amplified
82 partial resonance frequencies due to damaged and heterogeneous media near sources within
83 the magmatic plumbing system (Barajas et al., 2023), Bracale et al. [xxx].

84 On the other hand, the ongoing growth in seismic data has empowered applications
85 of machine learning in various tasks in seismology (Mousavi & Beroza, 2022). Machine
86 learning approaches have been used to detect complex patterns in seismic data for tasks
87 such as automated data processing (Mousavi et al., 2020), event detection (e.g. Zhu &
88 Beroza, 2019, with PhaseNet), seismic event localization (Majstorović et al., 2021), seismic
89 denoising (Viens & Van Houtte, 2019), ground-shaking simulation (R. D. Esfahani et al.,
90 2023), exploratory analysis (Mousavi et al., 2019; Seydoux et al., 2020; R. Esfahani et al.,
91 2021), and seismic imaging (Siahkoohi et al., 2022). While most applications of machine
92 learning rely on supervised learning using labeled data, here we employ an unsupervised
93 machine-learning approach to analyze seismograms, extracting robust features to capture
94 resonant frequencies. We utilize a deep scattering transform—a wavelet-based convolutional
95 neural network—to extract robust, time-invariant representations from non-stationary time
96 series (Andén & Mallat, 2014). By using predefined wavelet filters, the deep scattering trans-
97 form requires no filter learning, making it suitable for small datasets. This approach has been
98 successfully applied to various problems such as music classification (Andén & Mallat, 2014),
99 seismic precursor detection (Seydoux et al., 2020), earthquake clustering (Steinmann, Sey-
100 doux, Beaucé, & Campillo, 2022), subsurface monitoring (Steinmann, Seydoux, & Campillo,
101 2022), and volcano monitoring using continuous seismic records (Steinmann et al., 2024).

102 The primary objective of this study is to characterize the scattering and elastic prop-
103 erties of a non-uniform heterogeneous medium using unsupervised machine learning. To
104 achieve this, we analyze synthetic seismograms obtained from a companion study by Bra-
105 cale et al. [xxx] simulated in a 2D realistic magmatic structure based on (Melnik et al.,
106 2021). This model was developed through the injection of dykes and sills into the crust and
107 describes the melt fraction, i.e., the percentage of partially molten material. The effective

108 elastic moduli and seismic velocities are computed using the method proposed by (Schmeling
 109 & Wallner, 2012). The wavefield simulation produced 7,800 seismograms recorded at the
 110 surface, generated by sources distributed in depth on a regular grid in the central part of
 111 the heterogeneous model.

112 We employ a deep scattering network (DSN) to extract scattering coefficients from
 113 synthetic seismograms recorded at multiple stations with two components. We emphasize
 114 that the DSN (and the scattering coefficients) in this study differ from the physical scattering
 115 of seismic waves, and use the term “coda waves” instead of “scattered waves” to avoid
 116 confusion. To analyze and visualize the scattering coefficients, we use a manifold learning
 117 algorithm called Uniform Manifold Approximation and Projection (UMAP) (McInnes et
 118 al., 2018), a method that embeds the input into a two-dimensional space. Within this
 119 embedded space, we perform a local Principal Component Analysis (PCA) on the nearest
 120 neighbors of each data point (where each point represents a seismic source) to estimate
 121 local complexity and assess local resonant frequencies. The results are consistent with
 122 the statistical properties of the medium. We further analyze the effects of polarization of
 123 generated seismic S -waves using vertical and horizontal sources on the local complexity
 124 analysis. We demonstrate that the proposed technique provides valuable insights into the
 125 variations in heterogeneous media.

126 **2 Physics-based simulated training set**

127 The dataset used in this study comes from a companion study that simulates and
 128 validates seismic waveforms in a highly heterogeneous medium (Bracale et al. [xxx]). A
 129 subregion of the S velocity model is shown in Fig. 1 (the complete velocity model is presented
 130 in Fig. S3 in the Supplementary Materials). The model represents a magmatic plumbing
 131 system with 10 km sides embedded in a homogeneous medium with flat topography. The
 132 average P -wave velocity is 5.5 km s^{-1} , and S -wave 3.2 km s^{-1} , with a density of 2300 kg/m^3 .
 133 We adjust the original model by introducing an offset in the minimum S -wave velocity, as
 134 the solver employed for the SEM2D simulations solves the elastic wave equation (Komatitsch
 135 & Vilotte, 1998; Trinh et al., 2019; Cao et al., 2022). To stabilize the solver, we introduce
 136 a minimum value of 200 m s^{-1} in the velocity model.

137 The velocity model used in this study (subregion of velocity model in Fig. 1 and com-
 138 plete velocity model in Fig. S3) consists of $5,200 \times 3,200$ elements with a side length of 5 m,

139 corresponding to a 26×16 km area, to which 4.5 km of absorbing boundaries were added to
140 avoid reflections. The uppermost section of the model contains a homogeneous layer with
141 an approximate thickness of 4 km. Below this depth, we define a heterogeneous shallow
142 zone up to 6 km, containing horizontal structures such as sills. In the deeper part, from 6 to
143 13 km, the model becomes more heterogeneous due to the presence of vertical dykes. This
144 medium is highly complex, as scattering parameters such as ϵ , which defines the intensity of
145 velocity fluctuations, and the correlation length of the medium (see Supplementary Materi-
146 als for details, Text S2), exhibit significant variations in space. Therefore, interpreting such
147 data based on scattering theory considerations is extremely challenging, if not impossible,
148 as extensively discussed in Bracale et al. [xxx].

149 Due to the high computational resources required to run a complete simulation, we take
150 advantage of the reciprocity of the wave equation (Aki & Richards, 2002). This property
151 allows the interchange of source and receiver locations in the source scenario. We then
152 obtain the two components of ground motion generated by 7,800 sources and five receivers
153 by performing only ten numerical simulations. We considered two source mechanisms: a
154 horizontal force and a vertical force, both generated by a Ricker wavelet centered at 3 |hertz.
155 The receivers were placed on the surface, and sources were deployed at depths ranging from
156 2.5 to 14 km, arranged on a regular grid with a 100 m spacing. The simulated signal duration
157 is 37.4 s.

158 **3 Methods**

159 **3.1 Deep scattering transform**

160 Extracting robust representations from seismic waveforms should be tailored to a spe-
161 cific task. The time-frequency domain is particularly suitable for analyzing non-stationary
162 seismograms. The wavelet transform simultaneously decomposes the non-stationary signal
163 into a time-frequency domain, but there is a lack of time-translation invariance. In a deep
164 scattering transform, this invariance is achieved by applying a pooling operation to the scalo-
165 gram over time. Part of the loss of information due to the pooling operator can be recovered
166 through a second-order wavelet transformation followed by another pooling operator. The
167 pooling operator is user-defined, and acts as a pooling layer in a conventional convolutional
168 neural network (CNN). By cascading wavelet transforms and pooling operations, we use an
169 architecture similar to a convolutional neural network, known as a deep scattering network

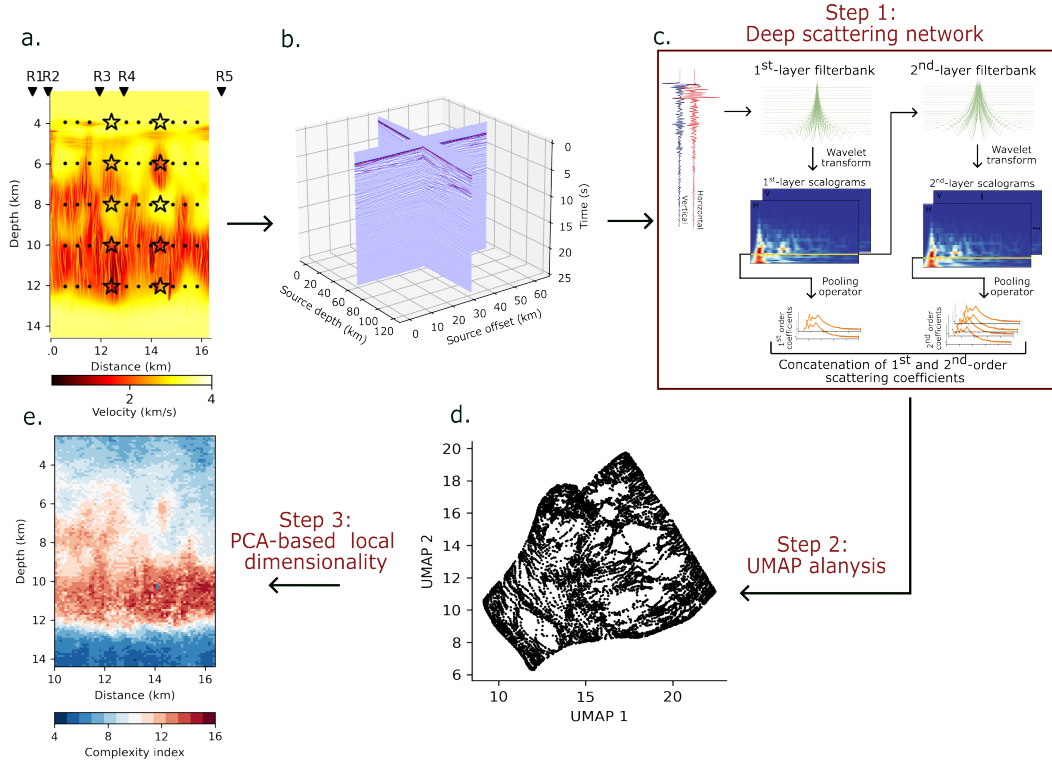


Figure 1. Workflow of the proposed approach. **a.** Shear wave velocity model used for seismograms simulation in a study by Bracale et al. [xxx]. **b.** Example of seismograms recorded at station R3 for sources at different offsets and depths (8,700 sources). **c.** Step 1: computation of first- and second-order scattering coefficients for a given station and source. **d.** Step 2: representation in the UMAP two-dimensional space. **e.** Step 3: analysis of local complexity in the embedded space using PCA-based local dimensionality.

170 (DSN, Andén & Mallat, 2014, see Supplementary Materials S1 for the formal details). The
 171 architecture is shown in Fig. 1c.

172 The DSN yields robust, time-translation-invariant representations that retain detailed
 173 information from signals in different scales over time and frequency. This architecture
 174 has been applied successfully and has outperformed other methods in tasks such as audio
 175 and music classification (Andén & Mallat, 2014), seismic signal exploration, and clustering
 176 (Steinmann, Seydoux, & Campillo, 2022; Steinmann, Seydoux, Beaucé, & Campillo, 2022;
 177 Seydoux et al., 2020).

178 In DSN, the first layer is defined by applying a wavelet transform to a signal to esti-
 179 mate the first-order scalogram and then the first-order scattering coefficients are calculated
 180 by applying a pooling operator. The second-order scalograms are defined by applying the
 181 second-order wavelet transforms to the modulus of the first-order scalogram. By applying
 182 the pooling operator to the modulus of the second-order scalograms, we calculate the second-
 183 order scattering coefficients. The first and second layers of the DSN can have varying num-
 184 bers of octaves and resolutions. The first-order scattering coefficients capture spectrum-like
 185 information (e.g., spectro-temporal energy fluctuations), while the second-order coefficients
 186 provide insights into the signal’s envelope modulations. Although the architecture of DSN
 187 resembles the one of a CNN, it differs in that each layer generates an output, not just the
 188 final layer (Andén & Mallat, 2014). Moreover, all filters in the DSN are predefined and not
 189 learned from training data. The following explains the details of the parameters used in our
 190 implemented DSN and explains how we applied it to seismic data streams.

191 The pooling operator is key, and must be tailored for specific tasks. Our goal is to
 192 analyze the coda waves that follow P - and S -wave onsets, so the pooling operator should
 193 mitigate the signature of P - and S -wave arrivals, and emphasize coda waves. Depending
 194 on the task, various pooling operators have been employed in seismic data analysis. For
 195 earthquake detection, maximum and average pooling operators are commonly used (Seydoux
 196 et al., 2020; Steinmann, Seydoux, Beaucé, & Campillo, 2022), while average pooling helped
 197 to monitor freezing of the near-surface (Steinmann, Seydoux, & Campillo, 2022). In contrast,
 198 the median pooling operator is useful to study the evolution of active volcanic systems over
 199 time (Steinmann et al., 2024). The median operator particularly emphasizes background
 200 information like coda waves and noise, while minimizing the focus on P - and S -wave arrivals
 201 and short-duration events (Steinmann et al., 2024).

202 Fig. 2a, b presents two synthetic seismograms, one simulated in a high-velocity zone with
 203 a local $v_s = 3000 \text{ m s}^{-1}$ and subsequently generates less coda waves. The other seismogram is
 204 simulated in a dyke with local $v_s = 1000 \text{ m s}^{-1}$, which leads to strong coda waves. Fig. 2c, d
 205 illustrates the effect of different pooling operators on the first-order scattering coefficients.
 206 Compared to the average pooling operator, the median operator more effectively captures
 207 information related to coda waves and reduces the influence of P - and S -wave arrivals.
 208 The second-order scattering coefficients derived from the median operator (Fig. 2e, f) and
 209 the average operator (Fig. 2g, h) demonstrate that the median operator provides better
 210 discrimination between seismograms with and without coda waves.

211 3.2 Uniform Manifold Approximation and Projection (UMAP) analysis

212 UMAP is a manifold learning technique designed for dimensionality reduction and vi-
 213 sualization of high-dimensional data by projecting data onto a low-dimensional embedded
 214 space. The UMAP algorithm constructs a manifold that preserves both local and global
 215 structures of the data using a graph-based algorithm. This approach focuses on the extrac-
 216 tion of the most relevant features from data and is based on algebraic topology and graph
 217 theory (more details in McInnes et al., 2018). UMAP is scalable and efficient for analyzing
 218 large datasets, and is similar to the t-distributed Stochastic Neighbor Embedding (t-SNE,
 219 Hinton & Roweis, 2002) and sequencing approaches (Baron & Ménard, 2019; Kim et al.,
 220 2020), which are also used for visualizing features and latent spaces in low-dimensional
 221 spaces.

222 UMAP involves three main hyperparameters: the number of nearest neighbors (n_n),
 223 the target embedding dimensions (d), and the minimum distance between points in the
 224 low-dimensional space (δ). In our case, we set $d = 2$, a common default value. The number
 225 of neighbors represents a trade-off between capturing small-scale and large-scale manifold
 226 structures. A smaller number of neighbors captures more detailed and localized struc-
 227 tures, while a larger number emphasizes global structures but may lose finer details. The
 228 δ hyperparameter controls how tightly points are packed in the low-dimensional represen-
 229 tation. A smaller δ results in a more densely packed representation that better preserves
 230 the local structure, whereas a larger δ spreads out the points, which can aid in visualiza-
 231 tion. In this study, we use UMAP to visualize high-dimensional scattering coefficients in
 232 a two-dimensional space (see Fig. 1c), with a choice of hyperparameters discussed in the
 233 Supplementary Materials (Text S3).

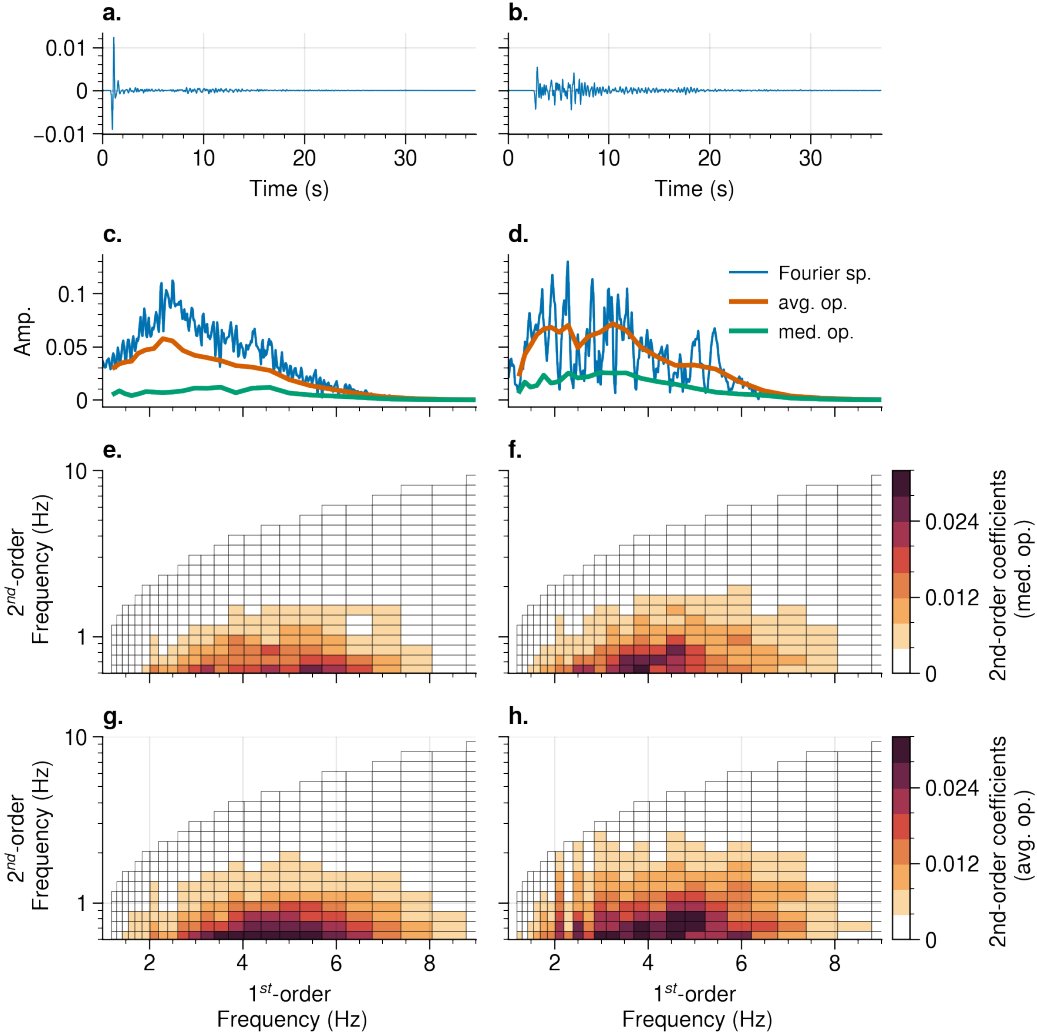


Figure 2. Effects of pooling operators of DSN on scattering coefficients. Comparison between median and average pooling operators on two seismograms with sources in (a.) high-velocity zone (low codas) and (b.) low-velocity zone (strong codas). c. and d. First-order scattering coefficients based on median (green) and average (red) pooling operators compared with the Fourier spectra (blue). The second-order scattering coefficients based on (e. and f.) median pooling and (g. and h.) average pooling operators.

234 **3.3 PCA-based local dimensionality**

235 The balance between local and global structures in UMAP space is achieved through
 236 the number of nearest neighbors (n_n). By leveraging the neighboring points of a given
 237 data point, we apply a local Principal Component Analysis (PCA) to each point and its n_n
 238 surrounding points (see Fig. 1d). We estimate the number of principal components required
 239 by performing a local PCA that satisfies a predefined variance threshold, referring to this
 240 as local complexity (or local dimensionality). The local complexity shows the embedded
 241 dimension of each point and its neighboring points to explain a specific variance (Brodu
 242 & Lague, 2012) and was applied to the Fourier representation in (Seydoux et al., 2016) to
 243 detect seismic activity from seismic networks. Ideally, the local complexity should align with
 244 the embedded dimension of the UMAP space; however, when the local complexity is high,
 245 embedding the data in two dimensions becomes challenging for UMAP to learn a manifold
 246 in a 2D space. The local complexity is also connected to the local entropy of the seismic
 247 wavefield, meaning that higher complexity indicates an increase in coda waves within the
 248 wavefield and, as a result increase in entropy. So the local PCA requires more components
 249 to explain the wavefield. We emphasize that the PCA-based local dimensionality is directly
 250 applied to scattering coefficients, while the UMAP algorithm provides the local connectivity
 251 structure via a k-nearest neighbor graph.

252 In local complexity analysis, the maximum number of principal components is deter-
 253 mined by $\min(n_n, n_{\text{features}}) - 1$, where n_n is the number of nearest neighbors and n_{features}
 254 is the number of scattering coefficients. The minimum number of principal components
 255 depends on the parametrization of the UMAP space.

256 **4 Results**

257 **4.1 Deep Scattering Representation of Simulated Seismograms**

258 We apply the DSN to the seismograms and extract a robust, time-invariant representa-
 259 tions. The implemented DSN consists of two layers. The first layer wavelet bank includes 32
 260 Gabor wavelets with center frequencies ranging from 1.2 to 9 Hz. The second-layer wavelet
 261 bank comprises 20 Gabor wavelets with center frequencies between 0.6 and 10 Hz. We use 8
 262 wavelets per octave at the first layer (dense representation) and 5 for the second layer (sparse
 263 representation). We use a quality factor of 2 for both layers. The wavelet bank is shown
 264 in Fig. S4. For a single station with two channels (horizontal and vertical components), we

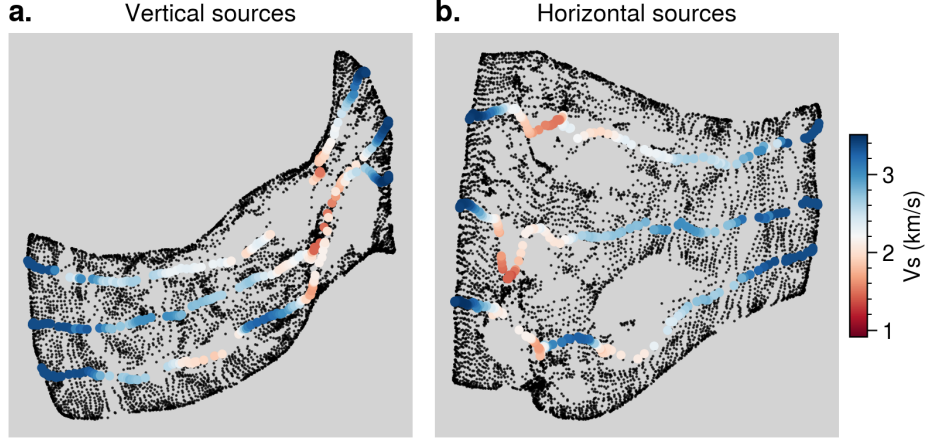


Figure 3. Uniform Manifold Approximation and Projection analysis for **a.** vertical and **b.** horizontal sources. The color represents the average shear-wave velocity surrounding the source locations in vertical segments at distances of 11, 13, and 14.5 km.

265 obtain 64 scattering coefficients at the first layer and 1,280 at the second. After computing
 266 the scattering coefficients, we mask the aliased second-order coefficients with second-order
 267 center frequencies above first-order center frequency (following Andén & Mallat, 2014). We
 268 analyze vertical and horizontal sources by combining data from all stations, assuming that
 269 different stations represent different realizations of the same event. In both cases, the seis-
 270 mograms are normalized before calculating the scattering coefficients. When concatenating
 271 data from all stations and all components, the total number of scattering coefficients be-
 272 comes 320 for the first layer and 6,400 for the second layer. We emphasize that the numbers
 273 of scattering coefficients are for records of five stations with two components.

274 4.2 Scattering coefficients in the UMAP space

275 Fig. 3 shows the low-dimensional representation from UMAP for vertical and horizontal
 276 sources. Each point in the UMAP space corresponds to a source in the subsurface. To better
 277 visualize the local and global structures within the UMAP space, we display the shear-wave
 278 velocity around each source for three segments at distances of 11, 13, and 14.5 km.

279 Fig. S5 shows the behavior of the UMAP manifold for vertical sources with different
 280 hyperparameters. While the overall structures of the UMAP spaces remain consistent across
 281 different parameters, it demonstrates how UMAP alters local and global structures as the

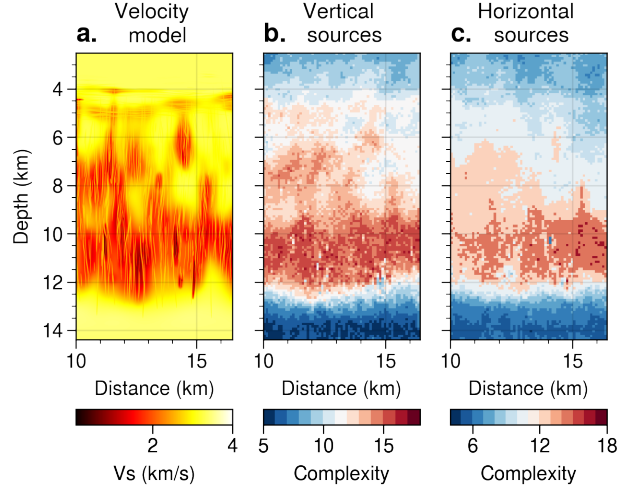


Figure 4. Local complexity analysis. **a.** Shear wave velocity model. Local complexity for the **b.** vertical sources and **c.** horizontal sources. The PCA explained variance is 80% for both sources.

282 number of nearest neighbors changes, especially in low- and high-velocity zones. We test n_n
 283 values of 10, 30, and 50 points. As the number of neighbors increases, UMAP tends to group
 284 all low-velocity and high-velocity zones at the top and bottom of the model, respectively,
 285 aligning with the global structure. We also test δ values of 0.1 and 0.9. Smaller δ values
 286 compress the data structure, resulting in more reliable representations. Based on visual
 287 inspection, we trade off hyperparameters that preserve both the local and global structures
 288 of the data. All analyses in Fig. 3 were performed with consistent hyperparameters: $\delta = 0.1$
 289 and $n_n = 40$ points.

290 4.3 Local complexity analysis

291 Our primary goal is to analyze the local structure of the data to characterize subsurface
 292 heterogeneities. To achieve this, we evaluate the local complexity using UMAP space with
 293 n_n of 20. The results of local PCA are shown in Fig. 4b,c for vertical and horizontal
 294 sources, respectively. The results indicate the number of principal components required by
 295 local PCA to explain 80% of the variance of the selected data, based on the local behavior
 296 of the UMAP space. The comparison between the local complexity and the velocity model
 297 shows that the local PCA captures the details of the velocity model for both sources.

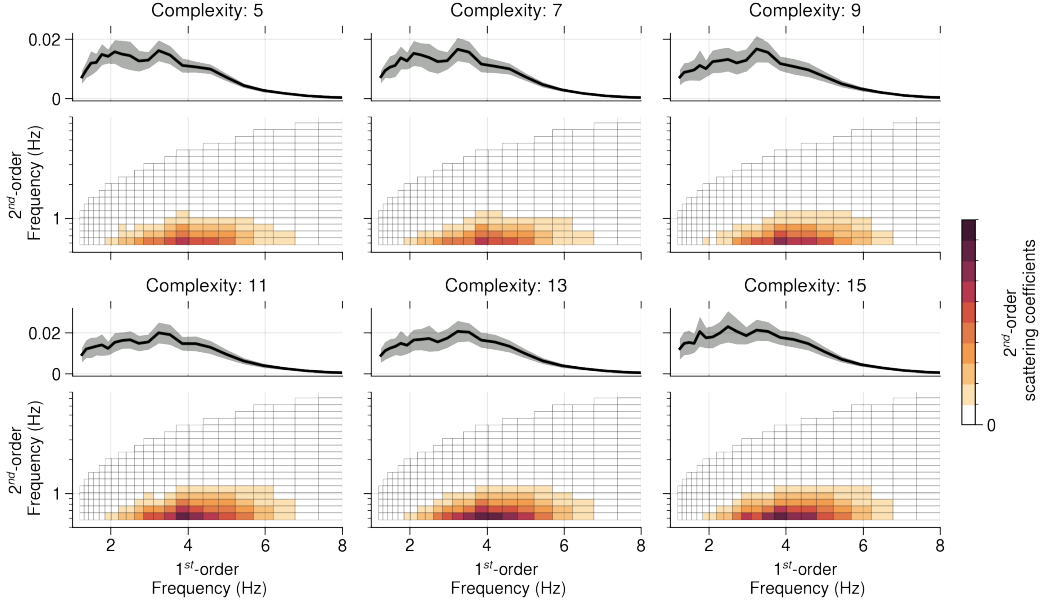


Figure 5. Average scattering coefficients for different complexities. The scattering coefficients are calculated for vertical sources at station R3 (see Fig. 3) and averaged for different values of complexity. The first-order scattering coefficients are shown on the top panels. The black curves show the averaged first-order coefficients and the gray areas show their standard deviations. The second-order scattering coefficients are presented in the bottom panels.

298 Local complexity analysis in the UMAP manifold is influenced by the number of neigh-
 299 bors and the predefined explained variance. Fig. S5 illustrates the sensitivity of local com-
 300 plexity to these parameters. As the number of neighbors increases, a higher local complexity
 301 is required to explain the selected variance. The number of neighbors primarily determines
 302 the minimum number of components needed for local PCA analysis. Increasing the ex-
 303 plained variance involves more PCA components in the final result, leading to the capture
 304 of more details from the scattering coefficients.

305 Fig. 5 shows the first- and second-order scattering coefficients for different complexity
 306 indexes for vertical sources and station R3. It shows that local complexity encodes infor-
 307 mation about the spectral peaks and frequency modulation. As local complexity increases,
 308 it leads to more complex spectral content and a shift toward higher frequencies (observed
 309 in the first-order) with larger frequency modulation (second-order).

5 Discussion

Here, we discuss the results of the UMAP analysis and local complexity analysis with a focus on vertical sources. In all of our analyses, we use a seismic network with five stations, as shown in Fig. 1a. The concatenation of records from multiple stations can be different realizations of a source, each affected by a unique path. These effects are mitigated through dimensionality reduction analysis (UMAP) since there is less similarity between them, and leads to accentuating the effects of heterogeneities in the vicinity of sources

5.1 Decoding Subsurface Structures: Interpretation of UMAP Space

UMAP space provides insights into the global and local structures of the scattering coefficients. Fig. 3 shows that the embedded manifold exhibits distinct patterns corresponding to high- and low-velocity zones. Sources in high-velocity zones are more widely distributed, while those in low-velocity zones appear densely embedded. This observation remains valid even by using different UMAP hyperparameters as shown in the Supplementary Materials Fig. S6.

5.2 Relation between Local Complexity and Statistical Properties of the Medium

Fig. 6 shows the relationship between local complexity and the average shear-wave velocity around each source for horizontal and vertical sources. The results indicate that high complexity levels capture low-velocity zones in the medium for both types of sources. This can be interpreted as an increase in entropy in the wavefield related to the seismic wave scattering in the low-velocity zones. As a result, the local PCA requires more components to capture the designated variance. The absolute residual (spectra peak anomaly) of the first layer is estimated between the first-order scattering coefficients for each complexity level and the average across all coefficients in Fig. 6. The results show that in the high-velocity zone between 12 and 14 km, the complexity is lowest, with residuals of scattering coefficients between 1 and 3 Hz. In the high-velocity zone at the top of the model, the residuals of scattering coefficients are about 2 Hz. Finally, in the low-velocity zone, the residuals shift to higher frequencies, around 3 to 5 Hz. Fig. 6c and f shows examples of waveforms for station R3 for each complexity level, as it increases, the coda waves in the waveform become more dominant.

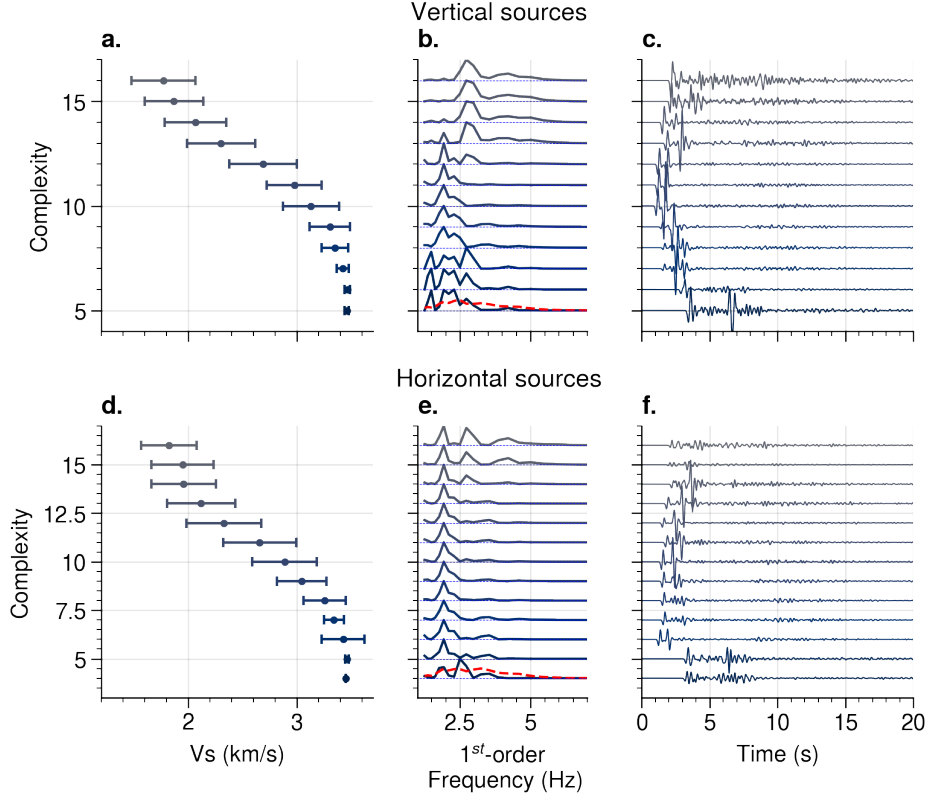


Figure 6. Relationship between complexity analysis and shear-wave velocity. The top panels show results from vertical sources, the bottom ones from horizontal sources. **a.**, **d.** Shear velocity distribution as a function of complexity. **b.**, **e.** First-order coefficients. **c.**, **f.** Examples of vertical components seismograms for each complexities.

340 Vertical and horizontal seismic sources generate waves with distinct polarizations. A
 341 comparison between the local complexity analysis of vertical and horizontal sources in
 342 Fig. 4b and c shows that source polarization has marginal effects on the final results. Nev-
 343 ertheless, the result of vertical sources yields better resolution in the top part of the model
 344 in depth between 4.5 and 5.5 km compared to the horizontal sources. It is important to note
 345 that in real-world scenarios, the polarization of seismic sources is more complex and vari-
 346 able than in this synthetic study, and their effects might be mitigated in the dimensionality
 347 reduction step.

348 The local complexity index captures the statistical properties of the medium as shown in
 349 Fig. 7. The correlation length and the intensity of velocity fluctuations ϵ (see Supplementary
 350 Materials Text S2) can be considered as a proxy for the mean free path. The correlation

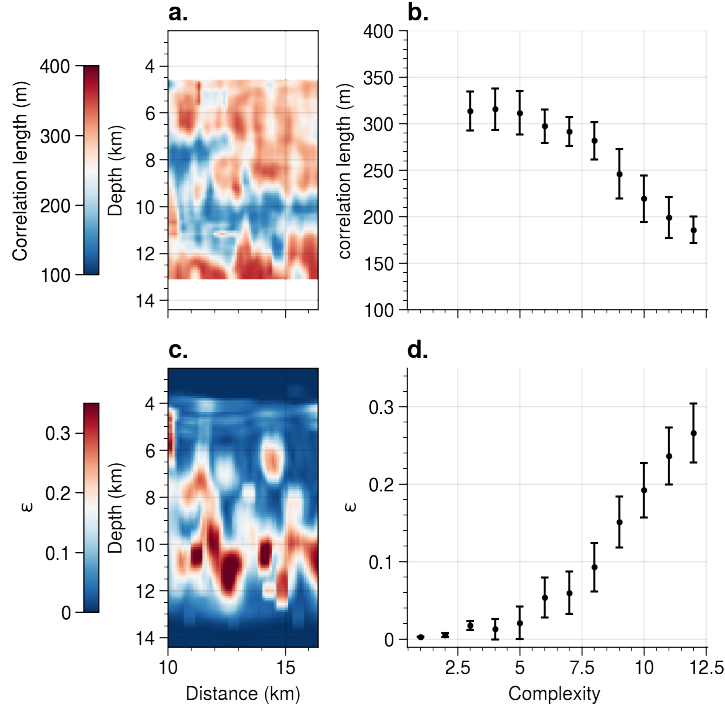


Figure 7. Comparison of medium properties with local complexity. **a.** Estimated correlation length in the medium. **b.** Relationship between correlation length and local complexity. **c.** Intensity of velocity fluctuations ϵ in space. **d.** Comparison between the intensity of velocity fluctuations and local complexity.

length decreases in the low-velocity regions, indicating that heterogeneities are stronger in these areas. The constant velocity regions at the top and bottom of the model are eliminated here since the velocity is constant. Fig. 7b further demonstrates that as local complexity increases, the properties of the medium change from an elastic to a scattering regime, leading to a decrease in correlation length. Fig. 7c shows the intensity of velocity fluctuations within a 500 m square window, normalized by the average velocity in the window. Higher values of ϵ indicate stronger local heterogeneity. Fig. 7d shows the relationship between the intensity of velocity fluctuations and local complexity and as ϵ increases, the local complexity increases in the model.

It is important to note that in heterogeneous media with a non-uniform distribution of heterogeneity, like the present study, seismic wave scattering regimes are spatially and temporally variable. The mean free path is intertwined with various parameters such as frequency, correlation length, wavenumber, and the intensity of velocity fluctuations. Lo-

364 cal complexity can be considered a data-driven proxy for evaluating heterogeneity in the
365 subsurface for such environments.

366 Given that source locations are required to apply the proposed method, it is also possible
367 to estimate local complexity directly within the spatial domain of the model rather than
368 relying on the combination of the UMAP space and the source locations. We did further
369 analysis applying local complexity using physical locations of sources (grid space with a
370 size of 7×7 around each source, a total of 49 sources) in Fig. S6. The observed local
371 complexities are consistent with results estimated in the UMAP space, suggesting that local
372 complexity behaves similarly across both real location and embedded manifold. It suggests
373 the structure-preserving mapping (isomorphism property) between the embedded domain
374 and the real locations of earthquakes. This is mainly due to the uniform distribution of
375 sources or simplicity in the source mechanism in this study. In real data, earthquakes
376 usually distribute non-uniformly, which can break this property.

377 **5.3 Limitations and outlook for future research**

378 In this work, we have focused solely on synthetic seismic data. However, the proposed
379 method has the potential to be applied to real data in various geological environments. The
380 proposed approach can also help to determine whether spectral peaks in volcanic tremors
381 and low-frequency earthquakes are the result of heterogeneity around the source (Barajas
382 et al., 2023) or due to source processes (B. A. Chouet, 1996).

383 Various studies have shown the complexity of fault zones using energy entrapment
384 due to significant heterogeneities based on teleseismic records (Share et al., 2017), ambient
385 noise analysis (van Dinther et al., 2020), and the v_P/v_S ratio (Huang et al., 2025). The
386 proposed method can provide insights into the complexity and heterogeneity level of active
387 fault systems by analyzing small earthquakes with accurately determined source locations.
388 We should note that applying the proposed approach to real data faces challenges like the
389 effects of low-velocity and weathering layers near the surface, the influences of topography
390 and the source mechanism.

391 **6 Conclusion**

392 In this study, we propose an unsupervised learning approach to assess heterogeneity
393 in complex geological media directly from seismograms. The proposed method is based on

394 deep scattering transformation and dimensionality reduction techniques. We introduce the
395 complexity analysis as a proxy for the level of heterogeneity and scattering characteristics
396 in media. We show that the complexity of the signal is linked to statistical properties in the
397 vicinity of the source, such as local correlation length and intensity of velocity fluctuations.
398 This approach could have valuable implications for improving the resolution of velocity
399 models and enhancing our understanding of geological structures in active regions, such as
400 fault systems and volcanoes.

401 **Open Research Section**

402 In this paper, we used the Scatseisnet package for deep scattering transformation (Seydoux
403 et al., 2025) and the Scikit-learn package for PCA analysis (Pedregosa et al., 2011). The
404 UMAP algorithm (McInnes et al., 2018) is available in [https://github.com/lmcinnes/
405 umap](https://github.com/lmcinnes/umap). We visualized our results using the Matplotlib (Hunter, 2007) and Proplot (Davis,
406 2021) packages. The Scipy package is used for other computational analyses (Virtanen et
407 al., 2020). The Python package to reproduce the results in this study is available in the
408 GitHub repository <https://www.github.com/resfahani/scatnet4scatwaves.git/>. The
409 simulated data used in this study is available at <https://zenodo.org/records/15118518>.

410 **Acknowledgments**

411 The authors declare no conflict of interest. This study was supported by the Multidis-
412 ciplinary Institute in Artificial Intelligence, MIAI, (for R.E., M.C., and L.S.), SPIN (for
413 M.B.), and F-IMAGE (for M.C. and L.S.).

414 **References**

- 415 Aki, K., & Chouet, B. (1975). Origin of coda waves: source, attenuation, and scattering
416 effects. *Journal of geophysical research*, *80*(23), 3322–3342.
- 417 Aki, K., & Richards, P. G. (2002). *Quantitative seismology*.
- 418 Andén, J., & Mallat, S. (2014). Deep scattering spectrum. *IEEE Transactions on Signal
419 Processing*, *62*(16), 4114–4128.
- 420 Barajas, A., Journeau, C., Obara, K., & Shapiro, N. M. (2023). Comparison of continu-
421 ously recorded seismic wavefields in tectonic and volcanic environments based on the
422 network covariance matrix. *Journal of Geophysical Research: Solid Earth*, *128*(12),
423 e2023JB026784. Retrieved from <https://agupubs.onlinelibrary.wiley.com/doi/>

- 424 [abs/10.1029/2023JB026784](https://doi.org/10.1029/2023JB026784) (e2023JB026784 2023JB026784) doi: [https://doi.org/](https://doi.org/10.1029/2023JB026784)
425 [10.1029/2023JB026784](https://doi.org/10.1029/2023JB026784)
- 426 Baron, D., & Ménard, B. (2019). Black hole mass estimation for active galactic nuclei from
427 a new angle. *Monthly Notices of the Royal Astronomical Society*, *487*(3), 3404–3418.
- 428 Ben-Zion, Y., & Sammis, C. G. (2003). Characterization of fault zones. *Pure and applied*
429 *geophysics*, *160*, 677–715.
- 430 Brodu, N., & Lague, D. (2012). 3d terrestrial lidar data classification of complex natural
431 scenes using a multi-scale dimensionality criterion: Applications in geomorphology.
432 *ISPRS journal of photogrammetry and remote sensing*, *68*, 121–134.
- 433 Cao, J., Brossier, R., Górszczyk, A., Métivier, L., & Virieux, J. (2022). 3-d multiparameter
434 full-waveform inversion for ocean-bottom seismic data using an efficient fluid–solid
435 coupled spectral-element solver. *Geophysical Journal International*, *229*(1), 671–703.
436 doi: [10.1093/gji/ggab484](https://doi.org/10.1093/gji/ggab484)
- 437 Chouet, B. (1992). A seismic model for the source of long-period events and harmonic
438 tremor. In *Volcanic seismology* (pp. 133–156).
- 439 Chouet, B. A. (1996). Long-period volcano seismicity: its source and use in eruption
440 forecasting. *Nature*, *380*, 309–316. Retrieved from [https://api.semanticscholar](https://api.semanticscholar.org/CorpusID:4341826)
441 [.org/CorpusID:4341826](https://api.semanticscholar.org/CorpusID:4341826)
- 442 Davis, L. L. B. (2021, October). *Proplot*. Zenodo. Retrieved from [https://doi.org/](https://doi.org/10.5281/zenodo.5602155)
443 [10.5281/zenodo.5602155](https://doi.org/10.5281/zenodo.5602155) doi: [10.5281/zenodo.5602155](https://doi.org/10.5281/zenodo.5602155)
- 444 Del Pezzo, E., Simini, M., & Ibanez, J. (1996). Separation of intrinsic and scattering q for
445 volcanic areas: a comparison between etna and campi flegrei. *Journal of volcanology*
446 *and geothermal research*, *70*(3–4), 213–219.
- 447 Esfahani, R., Vogel, K., Cotton, F., Ohrnberger, M., Scherbaum, F., & Kriegerowski, M.
448 (2021). Exploring the dimensionality of ground-motion data by applying autoencoder
449 techniques. *Bulletin of the Seismological Society of America*, *111*(3), 1563–1576.
- 450 Esfahani, R. D., Cotton, F., Ohrnberger, M., & Scherbaum, F. (2023). Tfcgan: Non-
451 stationary ground-motion simulation in the time–frequency domain using conditional
452 generative adversarial network (cgan) and phase retrieval methods. *Bulletin of the*
453 *Seismological Society of America*, *113*(1), 453–467.
- 454 Hennino, R., Trégourès, N., Shapiro, N., Margerin, L., Campillo, M., Van Tiggelen, B., &
455 Weaver, R. (2001). Observation of equipartition of seismic waves. *Physical review*
456 *letters*, *86*(15), 3447.

- 457 Hinton, G. E., & Roweis, S. (2002). Stochastic neighbor embedding. *Advances in neural*
458 *information processing systems*, 15.
- 459 Huang, Y., Ide, S., Kato, A., Yoshida, K., Jiang, C., & Zhai, P. (2025). Fault material het-
460 erogeneity controls deep interplate earthquakes. *Science Advances*, 11(9), eadr9353.
- 461 Hunter, J. D. (2007). Matplotlib: A 2d graphics environment. *Computing in Science &*
462 *Engineering*, 9(3), 90–95. doi: 10.1109/MCSE.2007.55
- 463 Kim, D., Lekić, V., Ménard, B., Baron, D., & Taghizadeh-Popp, M. (2020). Sequencing seis-
464 mograms: A panoptic view of scattering in the core-mantle boundary region. *Science*,
465 368(6496), 1223–1228.
- 466 Komatitsch, D., & Vilotte, J.-P. (1998, 04). The spectral element method: An efficient tool
467 to simulate the seismic response of 2d and 3d geological structures. *Bulletin of the*
468 *Seismological Society of America*, 88(2), 368-392. Retrieved from [https://doi.org/](https://doi.org/10.1785/BSSA0880020368)
469 [10.1785/BSSA0880020368](https://doi.org/10.1785/BSSA0880020368) doi: 10.1785/BSSA0880020368
- 470 Majstorović, J., Giffard-Roisin, S., & Poli, P. (2021). Designing convolutional neural network
471 pipeline for near-fault earthquake catalog extension using single-station waveforms.
472 *Journal of Geophysical Research: Solid Earth*, 126(7), e2020JB021566.
- 473 Margerin, L. (2005). Introduction to radiative transfer of seismic waves. *Geophysical*
474 *Monograph-American Geophysical Union*, 157, 229.
- 475 Margerin, L., Campillo, M., Van Tiggelen, B., & Hennino, R. (2009). Energy partition of
476 seismic coda waves in layered media: theory and application to pinyon flats observa-
477 tory. *Geophysical Journal International*, 177(2), 571–585.
- 478 McInnes, L., Healy, J., & Melville, J. (2018). Umap: Uniform manifold approximation and
479 projection for dimension reduction. *arXiv preprint arXiv:1802.03426*.
- 480 Melnik, O. E., Utkin, I. S., & Bindeman, I. N. (2021). Magma chamber formation by
481 dike accretion and crustal melting: 2d thermo-compositional model with emphasis on
482 eruptions and implication for zircon records. *Journal of Geophysical Research: Solid*
483 *Earth*, 126(12), e2021JB023008. Retrieved from [https://agupubs.onlinelibrary](https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2021JB023008)
484 [.wiley.com/doi/abs/10.1029/2021JB023008](https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2021JB023008) (e2021JB023008 2021JB023008) doi:
485 <https://doi.org/10.1029/2021JB023008>
- 486 Mousavi, S. M., & Beroza, G. C. (2022). Deep-learning seismology. *Science*, 377(6607),
487 eabm4470.
- 488 Mousavi, S. M., Ellsworth, W. L., Zhu, W., Chuang, L. Y., & Beroza, G. C. (2020).
489 Earthquake transformer—an attentive deep-learning model for simultaneous earth-

- 490 quake detection and phase picking. *Nature communications*, *11*(1), 3952.
- 491 Mousavi, S. M., Zhu, W., Ellsworth, W., & Beroza, G. (2019). Unsupervised clustering of
492 seismic signals using deep convolutional autoencoders. *IEEE Geoscience and Remote
493 Sensing Letters*, *16*(11), 1693–1697.
- 494 Pedregosa, F., Varoquaux, G., Gramfort, A., Michel, V., Thirion, B., Grisel, O., ... Duch-
495 esnay, E. (2011). Scikit-learn: Machine learning in Python. *Journal of Machine
496 Learning Research*, *12*, 2825–2830.
- 497 Schmeling, H., & Wallner, H. (2012). Magmatic lithospheric heating and weakening during
498 continental rifting: A simple scaling law, a 2-d thermomechanical rifting model and
499 the east african rift system. *Geochemistry, Geophysics, Geosystems*, *13*(8).
- 500 Seydoux, L., Balestrierio, R., Poli, P., Hoop, M. d., Campillo, M., & Baraniuk, R. (2020).
501 Clustering earthquake signals and background noises in continuous seismic data with
502 unsupervised deep learning. *Nature communications*, *11*(1), 3972.
- 503 Seydoux, L., Shapiro, N. M., de Rosny, J., Brenguier, F., & Landès, M. (2016). Detecting
504 seismic activity with a covariance matrix analysis of data recorded on seismic arrays.
505 *Geophysical Journal International*, *204*(3), 1430–1442.
- 506 Seydoux, L., Steinmann, R., Gärtner, M., Tong, F., Esfahani, R., Mouaoued, S., ...
507 Campillo, M. (2025, March). *Scatseisnet, a scattering network for seismic data anal-
508 ysis*. Zenodo. Retrieved from <https://doi.org/10.5281/zenodo.15110686> doi:
509 10.5281/zenodo.15110686
- 510 Share, P.-E., Ben-Zion, Y., Ross, Z. E., Qiu, H., & Vernon, F. L. (2017, 05). Internal
511 structure of the San Jacinto fault zone at Blackburn Saddle from seismic data of a
512 linear array. *Geophysical Journal International*, *210*(2), 819–832. Retrieved from
513 <https://doi.org/10.1093/gji/ggx191> doi: 10.1093/gji/ggx191
- 514 Siahkoobi, A., Rizzuti, G., & Herrmann, F. J. (2022). Deep bayesian inference for seismic
515 imaging with tasks. *Geophysics*, *87*(5), S281–S302.
- 516 Steinmann, R., Seydoux, L., Beaucé, E., & Campillo, M. (2022). Hierarchical exploration of
517 continuous seismograms with unsupervised learning. *Journal of Geophysical Research:
518 Solid Earth*, *127*(1), e2021JB022455.
- 519 Steinmann, R., Seydoux, L., & Campillo, M. (2022). Ai-based unmixing of medium and
520 source signatures from seismograms: Ground freezing patterns. *Geophysical Research
521 Letters*, *49*(15), e2022GL098854.
- 522 Steinmann, R., Seydoux, L., Journeau, C., Shapiro, N. M., & Campillo, M. (2024). Machine

- 523 learning analysis of seismograms reveals a continuous plumbing system evolution be-
 524 neath the klyuchevskoy volcano in kamchatka, russia. *Journal of Geophysical Research:*
 525 *Solid Earth*, 129(3), e2023JB027167.
- 526 Trinh, P.-T., Brossier, R., Métivier, L., Tavaré, L., & Virieux, J. (2019). Efficient time-
 527 domain 3d elastic and viscoelastic full-waveform inversion using a spectral-element
 528 method on flexible cartesian-based mesh. *GEOPHYSICS*, 84(1), R61-R83. Retrieved
 529 from <https://doi.org/10.1190/geo2018-0059.1> doi: 10.1190/geo2018-0059.1
- 530 van Dinther, C., Margerin, L., & Campillo, M. (2020, 12). Laterally varying scattering
 531 properties in the North Anatolian Fault Zone from ambient noise cross-correlations.
 532 *Geophysical Journal International*, 225(1), 589-607. Retrieved from [https://doi](https://doi.org/10.1093/gji/ggaa606)
 533 [.org/10.1093/gji/ggaa606](https://doi.org/10.1093/gji/ggaa606) doi: 10.1093/gji/ggaa606
- 534 Viens, L., & Van Houtte, C. (2019, 12). Denoising ambient seismic field correlation functions
 535 with convolutional autoencoders. *Geophysical Journal International*, 220(3), 1521-
 536 1535. Retrieved from <https://doi.org/10.1093/gji/ggz509> doi: 10.1093/gji/
 537 ggz509
- 538 Virtanen, P., Gommers, R., Oliphant, T. E., Haberland, M., Reddy, T., Cournapeau, D.,
 539 ... SciPy 1.0 Contributors (2020). SciPy 1.0: Fundamental Algorithms for Scientific
 540 Computing in Python. *Nature Methods*, 17, 261–272. doi: 10.1038/s41592-019-0686-2
- 541 Weaver, R. L. (1990). Diffusivity of ultrasound in polycrystals. *Journal of the Mechanics*
 542 *and Physics of Solids*, 38(1), 55–86.
- 543 Wu, R.-S., & Aki, K. (1988). Multiple scattering and energy transfer of seismic
 544 waves—separation of scattering effect from intrinsic attenuation ii. application of the
 545 theory to hindu kush region. *Scattering and Attenuations of Seismic Waves, Part I*,
 546 49–80.
- 547 Yamamoto, M., & Sato, H. (2010a). Multiple scattering and mode conversion re-
 548 vealed by an active seismic experiment at asama volcano, japan. *Journal of*
 549 *Geophysical Research: Solid Earth*, 115(B7). Retrieved from [https://agupubs](https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2009JB007109)
 550 [.onlinelibrary.wiley.com/doi/abs/10.1029/2009JB007109](https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2009JB007109) doi: [https://doi](https://doi.org/10.1029/2009JB007109)
 551 [.org/10.1029/2009JB007109](https://doi.org/10.1029/2009JB007109)
- 552 Yamamoto, M., & Sato, H. (2010b). Multiple scattering and mode conversion re-
 553 vealed by an active seismic experiment at asama volcano, japan. *Journal of*
 554 *Geophysical Research: Solid Earth*, 115(B7). Retrieved from [https://agupubs](https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2009JB007109)
 555 [.onlinelibrary.wiley.com/doi/abs/10.1029/2009JB007109](https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2009JB007109) doi: [https://doi](https://doi.org/10.1029/2009JB007109)

556 .org/10.1029/2009JB007109

557 Zhu, W., & Beroza, G. C. (2019). Phasenet: a deep-neural-network-based seismic arrival-
558 time picking method. *Geophysical Journal International*, *216*(1), 261–273.

559 Zieger, T., Sens-Schönfelder, C., Ritter, J. R., Lühr, B., & Dahm, T. (2016). P-wave scatter-
560 ing and the distribution of heterogeneity around etna volcano. *Annals of Geophysics*,
561 *59*(4).