1 **Mapping textures of polar ice cores using 3D laboratory X-ray microscopy**

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11

12 **ABSTRACT**

13 Deep ice cores from polar ice sheets enable reconstructions of Earth's past climate. Ice-core

14 records are therefore crucial for projecting future climate change, however, our ability to

15 interpret them relies on our understanding of polycrystalline-ice microstructures and

16 mechanics. In turn, these microstructures enable modeling of ice flow and large-scale effects of

17 ice-sheet evolution. Since drilling began in the 1950s, the ice textures and climate proxies

18 developed to decipher ice-core records have been analyzed in one- or two-dimensional (2D)

19 spaces, necessitated by the analytical instruments of core-processing lines and laboratories.

20 Here we develop a three-dimensional (3D), non-destructive approach to textural analysis that

21 preserves the natural context of ice and complements standard methods. Our method

22 combines lab-based absorption and diffraction contrast tomography to simultaneously

23 visualize, measure, and spatially correlate ice grains and air bubbles from volumetric and 3D

24 crystallographic perspectives, both lost during traditional sample preparations. We evaluate the

25 representation of 3D versus 2D data and discuss how access to both *c*- and *a*-axis directions of

26 grains may help constrain micromechanical models. We also built a specially designed cooling

27 device for the laboratory X-ray system to extend observational volumes by several orders of

28 magnitude over previous synchrotron-based measurements.

29

1 INTRODUCTION

Continental ice sheets covering Greenland and Antarctica play pivotal roles in Earth's climate system. Deep drill cores of these ice sheets reveal they have captured and preserved their response to past climates continuously over the last one hundred thousand to one million years, providing a means to reconstruct the past to help predict future climate and project changes in ice-sheet mass, geometry, and sea-level (e.g., Faria and others, 2014; Lauritzen and others, 2025). Reconstructions of climate records and ice dynamics use the same archive of ice-core layers and are interdependent in many ways. That is, when these layers formed through the burial and compaction of recurrent snowfalls, they trapped ancient atmospheres in the form of air bubbles, airborne particles, and elemental and isotopic abundances, but interpreting their stratigraphic chronology relies on our understanding of ice-grain mechanics and interactions with these climate proxies (e.g., Faria and others, 2010, 2014; NEEM community members, 2013; Stoll and others, 2021; Westhoff and others, 2021). More specifically, how microstructural mechanisms facilitate the creeping flow of ice sheets at the crystal-scale not only informs the origin and integrity of ice-core stratigraphy but also allows large-scale ice-sheet evolution and its climatic effects to be modeled (Fan and Prior, 2023; Gerber and others, 2023; Zhang and others, 2024). Our ability to interpret such information, however, fundamentally relies on our approach to textural analysis (grain and bubble morphologies, sizes, orientations) and how we access ice layers to begin with.

Since polar ice core drilling began in the 1950s, scientists have accessed thousands of meters of ice-sheet stratigraphy by developing what are now standard core-processing techniques. These include an assembly line of instruments and saws to measure along-core properties and to send samples to various labs for gas, chemistry, and isotope analysis. Although traditional sample preparations and analytical techniques involve dissection of the ice cores and can remove important three-dimensional (3D) context (Westhoff and others, 2021; Disbrow-Monz and others, 2024), one- and two-dimensional (2D) analyses have led to detailed climate-record reconstructions and form the basis of most textural studies carried out to date (e.g., Gow and Williamson, 1976; Lipenkov and others, 1989; Thorsteinsson and others, 1997; Svensson, Baadsager, and others, 2003; Durand and others, 2006; Binder and others, 2013;

59 Montagnat, Azuma, and others, 2014; Fan and others, 2021; Stoll and others, 2025). A question

60 remains, however, as to what scientists can learn if they initially and routinely have access to

61 ice-core textures in 3D.

62   Progress towards 3D characterization of snow and ice textures has been made over the

63 last two decades using X-ray microtomography. Air bubbles in deep ice cores are often

64 relatively easy to wholly visualize by X-ray attenuation given their small scale and low density

65 compared to ice grains, making it possible to measure porosity and contextualize trapped

66 greenhouse gases (Burr and others, 2018; Baker, 2019). However, in both synchrotron- and lab-

67 based X-ray systems, challenges surrounding the necessary cold analytical environments and

68 sample-size limitations mean that sufficient grain numbers with crystal orientation

69 measurements have not been achieved (e.g., Flin and others, 2004; Schneebeli and Sokratov,

70 2004; Kaempfer and Schneebeli, 2007; Rolland du Roscoat and others, 2011; Heggli and others,

71 2011; Riche and others, 2013; Granger and others, 2021).

72   However, crystal orientations are thought to play a critical role in ice-sheet evolution,

73 which begins with the intrinsic visco-plastic anisotropy of each ice crystal's hexagonal structure

74 (Duval and others, 1983; Fan and others, 2021; Fan and Prior, 2023). The basal plane in ice has a

75 relatively low resistance to shear, thus the gravity-driven flow occurring in ice sheets causes

76 dislocation slip mainly along the basal plane, while the crystal *c*-axis tends toward the

77 compressive-stress direction (Gow and Williamson, 1976; Alley, 1992). This leads to a bulk

78 anisotropy manifesting in crystallographic fabrics that develop under continued snow

79 deposition, burial, and ice deformation (Gow and Williamson, 1976). In effect, feedback occurs

80 between fabrics and flow, such that crystal orientations control the bulk mechanical behavior of

81 the polycrystalline ice, but the fabrics also evolve as a function of the ice's thermomechanical

82 history. Therefore, crystal orientations exert an important rheological control on the large-scale

83 flow of ice sheets (Duval and others, 1983; Shoji and Langway, 1985; Castelnau and others,

84 1998; Gerber and others, 2023; Rathmann and others, 2025).

85   Depending on the extent and mechanisms of ice deformation, the fabric-induced

86 mechanical anisotropy of the ice layers can vary considerably over the depth of the ice column

87 (Thorsteinsson and others, 1997; Montagnat, Azuma, and others, 2014; Montagnat and others,

88   2020; Stoll and others, 2025). For this reason, crystal orientation measurements are routinely

89   performed with automatic fabric analyzers (AFAs) or other mapping methods that rely on

90   optical microscopy and cm-scale, 2D thin sections of ice cores (Wilson and others, 2003;

91   Kipfstuhl and others, 2006). Although these are high-resolution methods and can rapidly

92   measure *c*-axis directions for hundreds of grains, they cannot determine directions of the three

93   *a*-axes for each grain (Wilen and others, 2003; Wilson and others, 2014). This leaves

94   consequential gaps in textural data, such as whether grains are truly distinct from one another

95   outside the section plane (Disbrow-Monz and others, 2024), or whether *a*-axis directions can or

96   should be used as strain markers to understand ice flow (Qi and others, 2019; Hunter and

97   others, 2023; Lilien and others, 2023). Furthermore, an incomplete understanding of how to

98   efficiently parameterize the elastic and visco-plastic anisotropies of individual grains partly

99   prevents ice-flow models from accounting for their effects, at least to the desired accuracy

100   (Ranganathan and others, 2021; Richards and others, 2021; Lilien and others, 2023;

101   Ranganathan and Minchew, 2024). At present, coupling small-scale micromechanical models

102   with large-scale ice-flow models is rarely considered for two main reasons. First, it has proven

103   challenging to incorporate realistic, yet efficient, representations of crystal fabrics into large-

104   scale models (e.g., Montagnat, Castelnau, and others, 2014; Lilien and others, 2023). Second,

105   only a few ways have been proposed to parameterize the predictions made by

106   micromechanical models, each with their own caveats and limitations (Placidi and others, 2010;

107   Gillet-Chaulet and others, 2005).

108   In addition to crystal-orientation fabrics, more work is needed to understand the

109   influence of grain sizes, shapes, boundaries, and even bubbles on ice metamorphosis and how

110   they should also be parameterized in models (Kuiper and others, 2020). However, collecting

111   this microstructural information has proven challenging for polycrystalline ice. Optical and cryo-

112   electron backscatter diffraction (EBSD) techniques provide visibility of ice-grain boundaries, but

113   only in 2D (e.g., Iliescu and others, 2004; Piazolo and others, 2008; Prior and others, 2011,

114   2015). Whereas, X-ray tomography provides 3D context but traditionally necessitates X-ray

115   diffraction from mm-scale samples to distinguish individual grains (Baker, 2019). Ultimately,

116   gaining the ability to correlate 3D behaviors of microstructures and bubbles in ice cores over

117 large sample volumes could improve both paleoclimate reconstructions and the physical

118 realism of large-scale ice-flow models.

119   Here we present a new non-destructive 3D method for the textural analysis of ice cores

120 that complements standard methods used in ice-core studies. By implementing well-

121 established grain-mapping techniques that have been transferred from the synchrotron to the

122 lab (e.g., Juul Jensen and others, 2006; Ludwig and others, 2008; King and others, 2014;

123 McDonald and others, 2015), our method makes it possible to simultaneously map ice grains

124 and pore spaces in 3D while measuring their volumes and shapes with respect to complete

125 crystallographic orientations of the grains. Our method combines two lab-based X-ray

126 modalities—absorption (ACT) and diffraction contrast tomography (DCT)—which enable the

127 visualization, measurement, and spatial correlation of ice-grain and air-bubble textures, adding

128 new perspectives on ice-core stratigraphy and new ways to access modeling parameters. We

129 achieve this by pushing the spatial resolution and sample volume limits of DCT in the lab and by

130 mitigating challenges in X-ray diffraction posed by crystal deformation. Our measurement

131 approach is based upon a specially designed cooling device that accommodates ~3 x 3 x 3 cm^3

132 samples containing hundreds to thousands of grains, extending observational volumes of ice by

133 one to four orders of magnitude over those quoted 5 to 30 years ago (e.g., Liu and others, 1992,

134 1995; Rolland du Roscoat and others, 2011; Baker, 2019). To demonstrate the power of greater

135 spatial and statistical information, we evaluate the representation of grain size distributions

136 and orientations in 3D versus 2D, investigate grain and bubble relationships in 3D, and discuss

137 the relevance of complete 3D crystallographic information in constraining micromechanical

138 models that underpin current large-scale ice-flow models.

139

140 **2 METHODS**

141   **2.1 Building a freezer inside the laboratory X-ray microscope**

142   A cooling device was designed at Lund University under the requirement that it (1)

143 keeps ice at a temperature well below -10°C and (2) maximizes sample size while minimizing

144 grain diffraction overlap, which we simulated beforehand. A sketch of the experimental set-up

145 is shown in **Figure 1**. Each sample sat directly on a 34 mm-diameter aluminum plate that was

146   cooled by a Peltier element. The hot bottom-side of the Peltier was kept cold using an external

147   chiller, attached to the device through an integrated slip ring system that held fluid hoses

148   stationary during sample rotation. Sample housing above the plate included a 27 mm-tall

149   polymer (PMMA) tube with 3 mm-thick walls, as well as surrounding polyethylene foam with 13

150   mm-thick walls, bringing the total housing diameter to 66 mm. For 30 minutes, chiller fluid

151   cooling to 5°C was circulated through the device before powering the Peltier and setting it to -

152   35°C, an unobtainable target that prevented temperature regulation. Plate sensor readings

153   reached -29°C in 10 minutes and remained stable below each sample during data acquisition

154   (up to 24 hours). A sensor embedded 3 mm beneath the upper sample surface also read stable

155   at -21°C during a test. While not ideal, thermal gradients of ~0.3°C mm^{-1} across the ~3 cm

156   sample height had negligible effects on ice microstructures at these temperatures and on our

157   analytical timescales. Air temperatures between the upper sample surface and insulating foam

158   reached -14°C in 5-7 minutes and typically continued to decrease by a few degrees.

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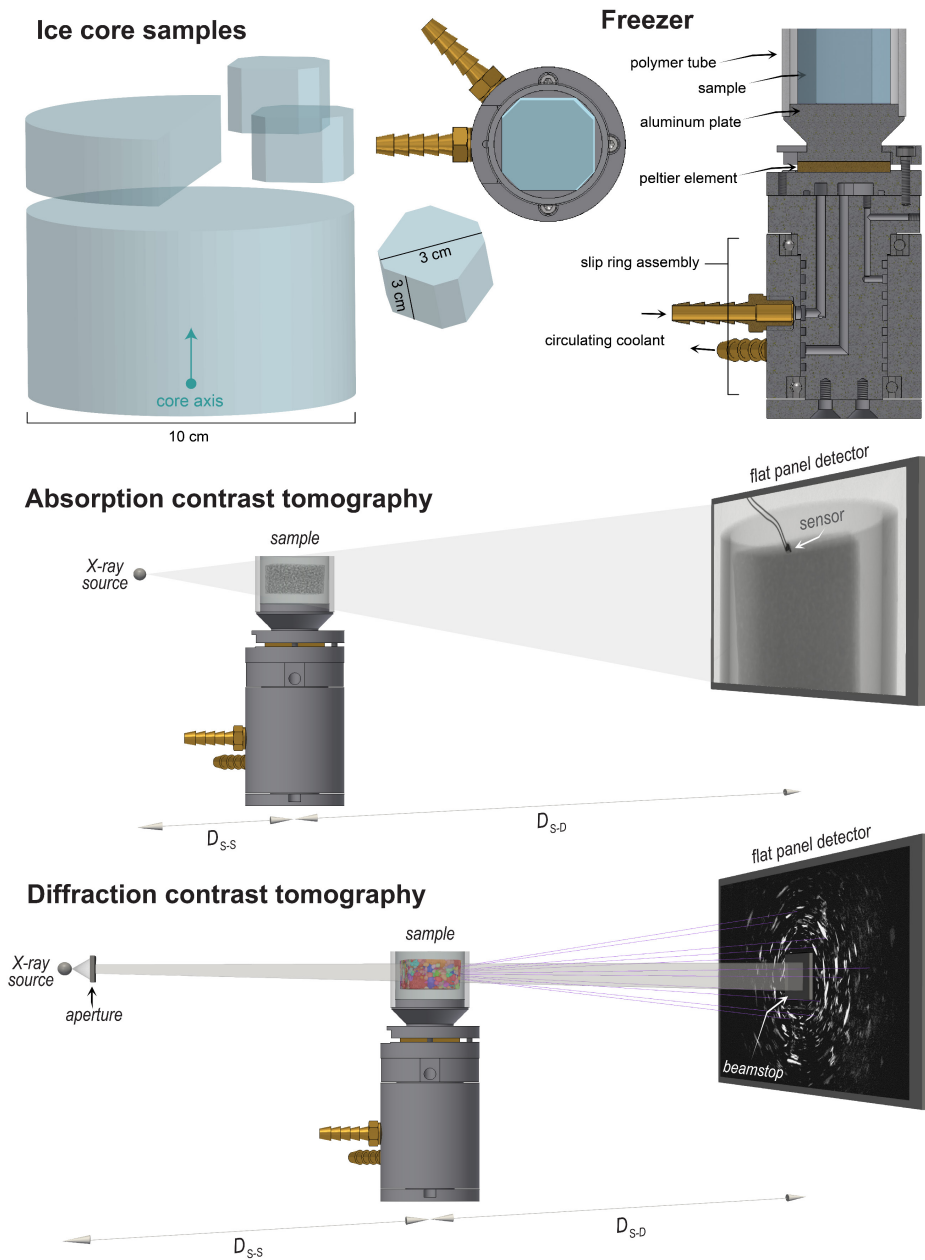
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162   **Figure 1.** The cooling device, sample and beam geometries used for lab-based multimodal X-ray imaging.

163   Examples of absorption and diffraction contrast projection data shown here correspond to a subvolume

164   of a firn sample described in Section 2.2. Source-to-sample and sample-to-detector distances are

165   referred to as D_{S-S} and D_{S-D}, respectively, and are provided in Section 2.3.

166

167

2.2 Selecting and working with ice core samples

169   We selected 10 samples from the Niels Bohr Institute's ice-core archive at the University

170   of Copenhagen. Samples include firn and Holocene-age deep ice from central, northwestern,

171   and southern drilling sites on Greenland (i.e., EGRIP, Eurocore, NEEM, and Dye-3; **Fig. 2**). For

172   experimental purposes, we targeted non-precious cores that had suffered drilling and recovery

173   complications in the field. Of the samples, we focused on three that cover a range in depth and

174   enabled us to develop methods and establish proof-of-concept (**Fig. 2**; **Table 1**). Each sample

175   captures less than one year, and in some cases less than one winter or summer season, based

176   on its thickness (0.03 m) and the annual ice-layer thickness at the depth it was extracted from.

177   The original cores have been stored at -30°C for durations that can be deduced from their

178   drilling dates in **Table 1**.

179   Sections of the 10 cm-diameter cores were sawn inside a walk-in freezer at the

180   University of Copenhagen and cut into smaller, ~3 x 3 x 3 cm^3 samples (**Fig. 1**). The samples

181   were then sealed in plastic bags and transported for 1 hour by cooler box to Xnovo Technology

182   in Køge, Denmark for X-ray imaging. The cooler box and its contents were stored for a few days

183   to months between -20 and -25°C inside a small, temperature-monitored freezer next to the X-

184   ray microscope. Ice exposure to room temperature lasted 2-4 seconds when samples were

185   transported by cold steel tweezers into the X-ray microscope cabinet. Mounted samples were

186   frozen to the aluminum plate using one drop of water, which temporarily increased the plate

187   temperature by 1°C. No ice metamorphosis occurred before or during data acquisition,

188   although an inconsequential <2 mm layer of frost commonly accumulated on sample surfaces
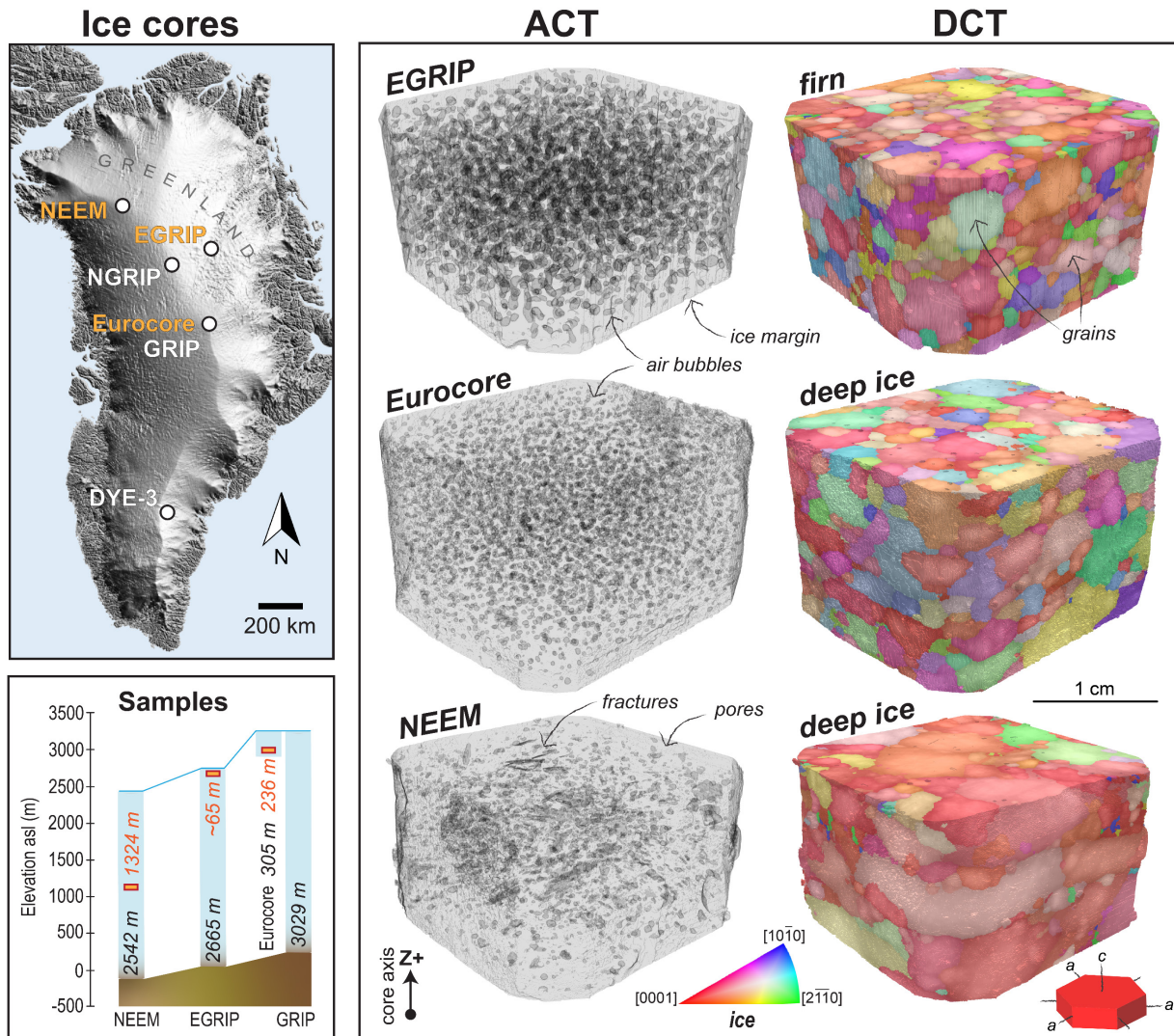
189   while inside the cooling device.

190

191

Figure 2. Overview of the three ice-core sample datasets used for method development. We screened samples from five ice-core drilling sites shown on the DEM of Greenland. In this paper, we present sample data from the sites labeled in orange (also see **Table 1**). Sample depths are provided with elevations and the final core depths at those drilling sites. EastGRIP was drilled between 2016 and 2023 within the active Northeast Greenland Ice Stream (NEGIS). The depth of the EGRIP firn sample is estimated based on sample porosity and firn density profiles of NEGIS (Vallelonga and others, 2014). Eurocore was drilled in 1989, just 30 m away from the GRIP site (GReenland Ice core Project, drilled 1990-92). Our deepest sample is from the NEEM core (North Greenland Eemian ice drilling project), drilled between 2009 and 2012. The observed differences in porosity and grain sizes between the three samples generally reflect their different depths within the ice sheet.

Table 1. Details of Holocene samples from Greenland ice cores

Sample	Core site	Year drilled	Depth (m)	Date[a-c]	Age (ka)	Annual layer thickness (m yr⁻¹)[a-c]
Firn	EGRIP	2016	~65	~1620 CE	~0.4	0.12
Late Holocene	Eurocore (GRIP)	1989	236	~1030 CE	~1.0	0.25
Early Holocene	NEEM	2010	1324	~7610 BCE	~9.6	0.06

[a]Mojtabavi and others (2020); [b]Seierstad and others (2014); [c]Rasmussen and others (2013)

204

2.3 Absorption and diffraction contrast tomography

206   Ice textures were characterized through correlative, multimodal X-ray imaging using a

207 Zeiss Xradia 520 Versa X-ray microscope equipped with a LabDCT Pro module. Absorption

208 contrast tomography was used to visualize sample morphology and determine the shapes of

209 pore spaces, including trapped (primary) air bubbles and secondary pores and fractures.

210 Whereas, diffraction contrast tomography was used to measure and map 3D crystallographic

211 orientations and shapes of grains. Grains can only be defined by DCT, because they lack

212 contrast in ACT due to their same composition and density. The crystal lattice distortion

213 inherent to glacial processes prompted us to explore different X-ray scanning conditions to

214 optimize data collection from grains that are plastically deformed, a general challenge in X-ray

215 diffraction methods (**Fig. S1**).

216   To optimize absorption contrast between ice and pore spaces, ACT scans were

217 performed at 80 kV and 7 W. We placed the X-ray source at 61 mm and detector at 245 mm

218 away from the sample to maximize spatial resolution while avoiding hardware collision with the

219 cooling device (**Fig. 1**). These distances created a geometrical magnification of 5x that

220 corresponded to a large, roughly 46 x 29 mm² field of view on the flat panel detector, which, as

221 a result, produced a tomography image containing 3064 x 1936 pixels at a nominal size of 75

222 μm and effective size of 14.95 μm. We collected 3201 projections (images), each with 0.5

223   seconds of exposure, as the sample completed a 360° rotation. This ACT acquisition step took

224   less than 2 hours.

225   Directly after the ACT scan, we performed the correlative DCT scan. Given that many

226   grains lie in the X-ray beam's path, we optimized diffraction patterns for crystallographic

227   reconstruction by performing DCT scans (110 kV, 10 W) with the flat panel detector in a

228   projection geometry close to Laue focusing (Bachmann and others, 2019), such that source-to-

229   sample (200 mm) and sample-to-detector (245 mm) distances were similar enough to produce

230   radially-focused, linear diffraction spots that reduced spot overlap (**Fig. 1**). Two different

231   source-beam apertures were used depending on the extent of grain deformation, which causes

232   a radial spread of cloudy and elongated diffraction spots, known as asterism (**Fig. S1**). A 375 x

233   375 μm^2 square aperture was used for firn and deep ice from <300 m depths. Whereas, a 200

234   μm circular pinhole aperture improved results for much deeper (1000s of meters) and more

235   deformed ice by (1) illuminating less sample volume per projection and (2) improving

236   diffraction spot definition by constraining spots to the aperture's pinhole shape. Depending on

237   the aperture used, our source distance gave roughly a 9 x 9 mm^2 to 5 x 5 mm^2 field of view. To

238   keep experiment times to a day, partial sample volumes of approximately 3 x 3 x 1.7 cm^3 were

239   covered in one advanced helical phyllotaxis raster scan (Oddershede and others, 2022)

240   consisting of 1119 to 5061 projections, each with an exposure time of either 5 or 10 seconds for

241   the larger aperture and 10 seconds for the pinhole. The DCT acquisition step took

242   approximately 5-6 hours when using the larger aperture and 20 hours using the pinhole. From

243   start to finish, each multimodal imaging measurement was completed in ~8-24 hours.

244

245   **2.4 Data reconstruction and processing**

246   *2.4.1 Ice and pore volumes*

247   We reconstructed sample volumes from the ACT data using cubic voxels with side-

248   lengths of 14.95 μm. This was performed using a filtered back projection routine and the

249   reconstructor tool in ZEISS Scout-and-Scan software. Because individual ice grains have the

250   same density, they cannot be distinguished by way of X-ray attenuation and thus form one

251   continuous volume with homogenous grayscale intensity. Much lower density air bubbles >40

252   μm in diameter and apparent microfractures were detectable at the reconstructed voxel size,

253   while dust impurities, which are typically on the order of a few micrometers, fell below the

254   detection limit.

255

256   *2.4.2 Ice grain crystallographic orientations and morphologies*

257   Using Xnovo Technology's GrainMapper3D software, crystallographic maps were

258   reconstructed from the DCT data with 60 μm-wide cubic voxels inside the volume mask

259   obtained from the ACT data (Bachmann and others, 2019; **Fig. 2**). Natural ice on Earth occurs as

260   ice Ih (ice one hexagonal), which belongs to the hexagonal crystal system and dihexagonal

261   dipyramidal 6/*mmm* Laue class (and point group). Lattice parameters for space group $P6_3/mmc$

262   of *a* = 4.489 Å, *c* = 7.327 Å and the four {*hkil*} families {$2\overline{1}\overline{1}0$}, {$10\overline{1}0$}, {$10\overline{1}3$}, and {$2\overline{1}\overline{1}2$} were

263   used to reconstruct ice grains from diffraction patterns. We use **Figure 3** to show how we

264   define grains and subgrains in our data. Adjacent regions in 3D space were assigned to the

265   same "grain" if the voxel-to-voxel, crystallographic angular misorientation was <2° (EGRIP firn,

266   Eurocore deep ice) or <3° (NEEM deep ice). The reason for using different thresholds for

267   different samples relates to their amount of deformation, its impact on diffraction pattern

268   quality, and therefore the confidence in grain detection and reconstruction. Lower

269   misorientation thresholds yielded large numbers of low confidence grains a few hundred

270   microns in size, which were interpreted to be unphysical. Therefore, grain maps were cleaned

271   by replacing any grain with a neighboring grain that it shared the largest boundary area with,

272   based on fulfillment of at least two of the following three criteria indicative of poor

273   reconstruction: (1) a mean completeness of <60% across the grain (completeness is defined as

274   the ratio between the number of observed and expected diffraction spots for a given

275   orientation); (2) a grain volume less than 8^3 = 512 voxels (equivalent sphere diameter <600 μm);

276   (3) a misorientation to a neighbor-of-a-neighbor grain of <2° (EGRIP firn, Eurocore deep ice) or

277   <4° (NEEM deep ice). These limits ensured the merging of high confidence grains with low

278 confidence grains located at their boundaries and with misorientations close to the set

279 threshold.

280   Using GrainMapper3D, we validate the reconstruction of both non-deformed and

281 deformed grains by inspecting forward projections of the {*hkil*} families listed above (see

282 Bachmann and others, 2019 for details). As seen in **Figure S1**, forward projection outlines cover

283 the spread in orientation represented by elongated diffraction spots that originate from

284 deformed grains. When we define a deformed grain (i.e., consider the deformed grain to be

285 one grain with one orientation), it becomes the average orientation (or the average forward

286 projection), which can also be seen in **Figure S1**. Although the current reconstruction model

287 does not account for lattice deformation, it can be tolerated so long as diffraction spots can be

288 clearly extracted and separated (Bachmann and others, 2019). Resolving subdomains or

289 subgrain boundaries is dependent on use of the Laue-focusing geometry (increases angular

290 sensitivity), working distances, and grain and sample sizes (Bachmann and others, 2019). We

291 calculated crystallographic misorientations using the difference between a reconstructed

292 grain's average orientation and the orientations of individual voxels composing that grain (i.e.,

293 the grain reference orientation deviation, or GROD). Because we considered adjacent voxels

294 with <2-3° misorientations to be the same "grain", the maximum misorientations of subgrains

295 with low-angle grain boundaries were defined to be <2° or <3°, depending on the specific

296 sample. We show in **Figure 3** that combined maps of completeness, inverse pole figures (IPFs),

297 and grain reference orientation deviations (GRODs) can be used to characterize subgrains. The

298 uncertainty in grain boundary position is approximately 5-10% of the grain size and is related to

299 the diffraction spot size and their binarization in our data (cf. Bachmann and others, 2019). For

300 spots originating from deformed grains, this uncertainty can be greater and needs to be

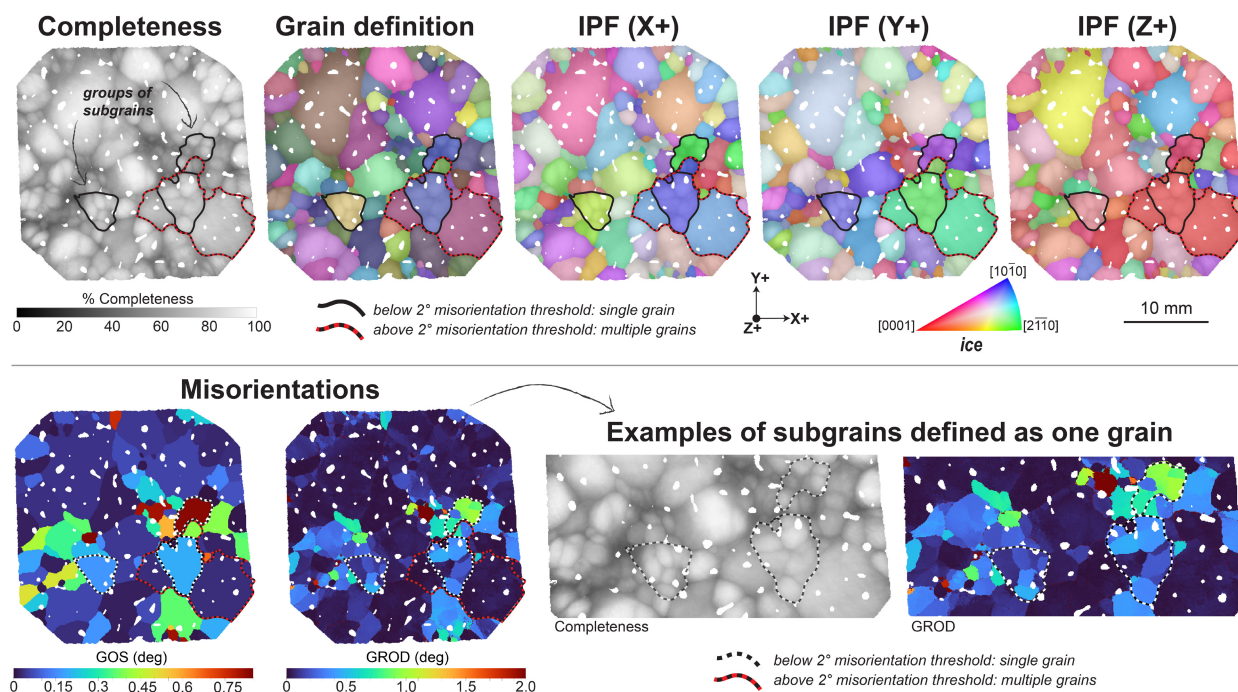301 quantified and statistically assessed in the future.

302

303

304

Figure 3. Grain and subgrain definition in this study. Grains were reconstructed based on our chosen scanning conditions and crystallographic misorientation thresholds described in the **Methods**. Virtual slices of 3D grain maps of the EGRIP firn sample are used here for demonstrative purposes, wherein we defined grains based on a 2° misorientation threshold. Grain completeness typically drops both at grain and subgrain boundaries, enabling their visualization. This is because of the 1:1 relationship between the shapes of diffraction spots on the detector and the shapes of the diffracting volumes (i.e., grains or subgrains), with the spot edges corresponding to grain and subgrain boundaries. As a note, the completeness map underlies all other maps shown in the top row. The grain definition map reflects our chosen misorientation threshold to define grains, resulting in groupings of adjacent regions defined by completeness drop-off and circled in black. When looking at IPF maps for XYZ directions, regions defined as one grain appear with the same IPF color, consistent with the grain definition map. IPF coloring may be similar for neighboring grains above the chosen misorientation threshold, such as those circled by red-dashed lines. The grain reference orientation deviation (GROD) and grain orientation spread (GOS, or average GROD) maps verify that apparent subgrains identified in previous maps can be interpreted as such, given that the locations of completeness drop-off correspond to low-angle misorientation boundaries.

321

322

2.5 Data analysis and visualization

323

324   Our analysis was focused on demonstrating the complementarity of 3D and 2D data, as

325 well as the new possibilities offered by 3D data available through our method. Therefore, the

326 analysis is limited in scope regarding ice mechanics and climate records. We instead focus on

327 comparing traditional and new perspectives using a few of the most common types of texture

328 analyses performed on ice cores. All analyses and 3D visualizations were generated using in-

329 house code at Xnovo Technology, GrainMapper3D (Bachmann and others, 2019), Paraview

330 (Ayachit, 2015), or the Python package Vedo (Musy and others, 2025).

331

2.5.1 Grain sizes

332

333   Grain size statistics were derived from both 3D reconstructed volumes and 2D cross-

334 sectional slices for comparison. Two grain metrics, the equivalent grain diameter and major-axis

335 length, were chosen to demonstrate the comparisons in **Figures 4 and 5** (also see **Fig. S2**). For

336 each grain and each slicing axis (X, Y, Z), the mean, minimum, and maximum values of the

337 chosen 2D metric across all slices intersecting that grain were calculated. In the generated

338 scatter plots, each point represents a single grain, with the 3D metric value compared to the

339 mean 2D metric value calculated from all slices intersecting the grain, either along a specified

340 slicing axis or XYZ axes combined. Vertical bars were added to each point (grain) to indicate the

341 minimum and maximum 2D metric value observed when slicing along a specified axis.

342 Additionally, grain size distributions based on the equivalent diameter were visualized through

343 histograms that emphasize the contribution of different sizes to the total volume (for 3D) or

344 total area (for 2D slices). For demonstrative purposes, grain-related analyses include grains that

345 intersect the sample margins due the presence of larger grain sizes (cf. Svensson, Baadsager,

346 and others, 2003).

347

2.5.2 Grain orientations

348

349   Pole figure analysis involved plotting separately: (1) the number-weighted grain

350 orientations; (2) the volume-weighted grain orientations; and (3) all grains compared to those

351   of separate grain size fractions for (1) and (2) (**Figs. 6 and 7**). To do this, we plotted and

352   symmetrized discrete *c*-axis and *a*-axes directions for each grain within the sample reference

353   frame (XYZ). The kernel density of the distribution of crystal axis directions was estimated using

354   a de La Vallee Poussin kernel with a halfwidth (kappa) of 10° and weighting them by number or

355   volume. A unit sphere grid with density weights was plotted in a filled tri-contour plot in an

356   equal-angle, lower-hemisphere stereographic projection.

357

358   *2.5.3 Pore positions and shapes*

359   Pore analysis in 3D leveraged complementary information from ACT and DCT datasets.

360   To gain insight into spatial relationships between pores and grains, we correlated ACT

361   segmentation and DCT grain labels. After segmenting pores from the ACT volume, a connected

362   components analysis (26-connectivity) and 5-voxel minimum was used to identify individual

363   pores and calculate their voxel count, volume, equivalent sphere diameter, and centroid. Pore-

364   grain adjacency was then determined relative to the DCT grain map, such that the

365   neighborhood of each identified ACT pore was examined within the corresponding region of the

366   DCT volume. The pore-grain adjacency was visualized by generating a surface mesh of grains in

367   the DCT volume (colored by IPF crystallographic orientations) and overlaying it with the ACT-

368   derived pore mesh (one color, gray), which made pore and grain boundary surfaces visible.

369   These surfaces provided a visualization of ACT pores colored by the IPF orientation of their host

370   grain(s). We then characterized pores based on their *position* and size with respect to grain

371   boundaries, classifying them as "intragranular" or "intergranular" based on the number of

372   unique adjacent DCT grain labels detected in a 3 x 3 x 3 (ACT) voxel-dilated boundary around

373   each pore. In the EGRIP firn sample, where the average pore size is 1 mm and most grains are

374   smaller than 10 mm, the pore classification is highly accurate considering the estimated

375   uncertainty on grain boundary position is 5-10% of the grain size. The pore classification is

376   considered less accurate in our Eurocore and NEEM data due to a combination of smaller pores

377   and larger, more deformed grains. Counts and sizes of total, intragranular, and intergranular

378   pores were analyzed statistically. Pores that intersect sample margins had a negligible effect on

379   the analytical results.

380   To characterize each pore based on its *shape* and size, we made principal

381   measurements using ACT data and performed a convex hull fitting with a "minimum volume

382   enclosing ellipsoid" containing a set of points. Fitted axis measurements were made for

383   ellipsoids following the axis convention $a \geq b \geq c$ [where ellipsoid factor, *EF = c/b − b/a,* and the

384   aspect ratio is *a/c*]. These measurements include axis length (L) and corresponding azimuth (φ)

385   and elevation (θ) after rotation. Analyses focused on measuring the extent and direction of

386   pore flattening or elongation with respect to the host grain's *c*-axis direction, as well as the

387   sample compression axis (Z), which is parallel to the ice-core axis. Only intragranular pores were

388   considered for the purpose of comparing results to those of Fegyveresi and others (2019), who

389   focused on intragrain deformation. Therefore, the intergranular pores touching grain

390   boundaries were excluded entirely from our shape-related analyses, while the intragranular

391   pores were also filtered to exclude those intersected by sample margins.

392

393   **3 RESULTS AND DISCUSSION**

394   **3.1 Visualizing ice grains and pores through correlative 3D images**

395   Herein we demonstrate some of the utility of lab-based, multimodal (ACT + DCT) X-ray

396   tomography for the 3D analysis of ice cores. For brevity, we focus on one firn sample and two

397   deep ice samples that cover a range in depth and metamorphosis, as exemplified by differences

398   in their porosities, grain sizes, shapes, and crystallographic orientations visible in **Figure 2**.

399   Again, reconstructing and correlating images of these parameters required the consecutive

400   collection of X-ray absorption and diffraction image data. When correlating the two image

401   types, pores and ice grains are numerable and their spatial characteristics and relationships

402   measurable in 3D. Notably, the correlated images enable fast qualitative assessments of ice

403   core textures that can guide and contextualize measurements.

404   Prior to any quantitative analysis, there is visible evidence of crystallographic preferred

405   orientations (CPOs) in the IPF maps in **Figure 2**. These maps are colored with respect to the

406   "up" or vertical (Z+) direction of ice cores, perpendicular to the glacier surface. The dominance

407   of red indicates a *c*-axis preferred orientation, such that many grains are oriented with the axis

408   near-vertical. These observations match those commonly made from automatic fabric analyzer

409   images of thin sections (e.g., Faria and others, 2014). The correlative maps also allow

410   qualitative observations of *a*-axis CPOs, as well as volumetric differences in porosity and grain

411   sizes between different samples.

412   For the samples we considered and the ~12-15 cm^3 volumes we correlated, we observed

413   representative air bubble sizes, shapes, and distributions, but capturing representative grain

414   sizes was challenging, as large grains can intersect sample margins or extend beyond imaged

415   volumes. Our observations are corroborated by previous documentations of bubbles, which

416   reach just tens of micrometers to millimeters beneath the lower part of the firn-column, where

417   bubbles become closed-off (e.g., Bendel and others, 2013; Westhoff and others, 2024). In

418   contrast, ice grains can reach centimeters at depth and even meters near warm bedrock (Baker,

419   2019). While we achieve a total sample size of ~3 cm tall x ~3 cm diameter, which is much

420   greater than sizes attempted previously in synchrotron- and lab-based multimodal experiments

421   (e.g., 2-mm thick up to 1 x 1 cm; Liu and others, 1992, 1995; Jia and others, 1996; Rolland du

422   Roscoat and others, 2011), we acknowledge that even larger samples would be ideal for

423   accurate characterizations of coarse-grained ice. However, larger samples equate to

424   measurement limitations, ranging from too many diffracting grains crowding the X-ray

425   detector, to resolution limits and long acquisition times.

426   Analyzing ice cores with grain sizes too large to capture whole is a problem for all

427   existing methods. When using optical microscopy and thin sections, the problem may be

428   compounded by irregular grain shapes, such that duplicated measurements of grains appearing

429   more than once in the section will influence CPO analysis (Svensson, Schmidt, and others, 2003;

430   Monz and others, 2021). Furthermore, although cryo-EBSD has been applied to serial sections

431   of coarse-grained ice to identify representative volumes, shapes, and sufficient grain numbers

432   needed for accurate CPO characterizations, this is labor-intensive and can create positional and

433   angular uncertainties when recombining sections and stitching data together (Monz and others,

434   2021). We note that a stitching approach can also be taken with our X-ray method using serially

435   cut ice blocks, which, when combined with recent advancements in 3D image correlation, could

436 eliminate many uncertainties and address large grain sizes. Effectively, the sample volumes we

437 imaged here balance the pros and cons of relatively large samples for lab-based DCT and

438 provide the option to collect lower- versus higher-resolution data to fit scientific questions.

439

440 **3.2 Contrasts between 3D and 2D grain size distributions**

441 Ice rheology has both crystal orientation and grain size dependencies (Cuffy and others,

442 2000; Durand and others, 2006). How these characteristics vary across ice layers largely reflects

443 the deformation mechanisms driving the viscous flow of ice sheets under their own weight

444 (Thorsteinsson and others, 1997; De La Chapelle and others, 1998; Ranganathan and others,

445 2021). Crystal orientation fabrics and grain sizes, however, are commonly decoupled (Svensson,

446 Baadsager, and others, 2003; Wang and others, 2003; Stoll and others, 2025). In thin sections,

447 the cross-sectional areas and shapes of grains change depending on burial depth, temperature,

448 amount of shear, or impurity concentrations. Observed trends between grain area (2D)

449 distributions and depth have led to informative grain growth models, although it has been

450 acknowledged that grain volume (3D) data would better suit these models (Svensson, Schmidt,

451 and others, 2003).

452 **Figures 4 and 5** demonstrate how volumetric data from 3D grain maps can be used to

453 assess the representativeness of 2D grain size distributions, which are the standard calculations

454 typically made using vertically, and sometimes horizontally, cut surfaces of ice cores (Faria and

455 others, 2014). While there are many ways glaciologists go about these calculations, we chose to

456 compare 3D and 2D grain size distributions by calculating equivalent sphere and circle

457 diameters of grains from their volumes and cross-sectional areas, respectively. We also extract

458 the major- and minor-axis lengths of grains based on ellipsoid (3D) and ellipse (2D) fitting. We

459 show in **Figures 4 and 5** how vertical and horizontal cuts can impose bias on grain size

460 measurements and related calculations (also see **Figs. S2 and S3**). This is a well-known problem

461 that continues to motivate refinement of statistical correction methods, as a single 2D

462 measurement cannot capture the full grain shape and must be based on additional assumptions

463 (e.g., Morgan and Jerram, 2006; Mangler and others, 2022; Bretagne and others, 2023). To

464 demonstrate this dimensional bias statistically, we digitally cut the sample volume into vertical

465   and horizontal serial sections at intervals equal to the 60 μm voxel size of the DCT grain map.

466   This enabled us to measure the mean circle (2D) diameter of each grain against its spherical

467   (3D) diameter, while noting the full range of circle diameters encountered (**Figs. 4 and 5**). We

468   find that, on average, grain sizes would be severely underestimated in vertical cuts of our

469   samples, whereas grain sizes in horizontal cuts would be comparatively more representative

470   but with many overestimations (i.e. with respect to equivalent sphere and circle diameters).

471   Using the major-axis, for example, consistently leads to underestimations of grain sizes in these

472   samples (**Fig. 5 and S2**). These results indicate the influence of grain shapes and, in particular,

473   the effect of oblate or flattened grains in the horizontal plane, which becomes prevalent in

474   deep ice due to high compressive stress. Horizontal cuts may be significantly biased when

475   bimodal grain sizes are distributed unevenly or in layers, as in the NEEM sample (**Fig. 5**). As one

476   would expect, the probability of making a cut that yields a 2D grain size distribution matching

477   the 3D one will be greater for samples with more uniform grain shapes and sizes, but this is

478   improbable for samples containing high aspect ratio grains or varied sizes. Here we collectively

479   show how 3D grain maps provide insight into the relative contributions of grain shape versus

480   grain size, which may yield sample-specific correction factors for routine 2D grain size

481   calculations and help refine models of grain development and their utility (cf. Thorsteinsson

482   and others, 1997; Svensson, Schmidt, and others, 2003; Rollet and others, 2017).

483
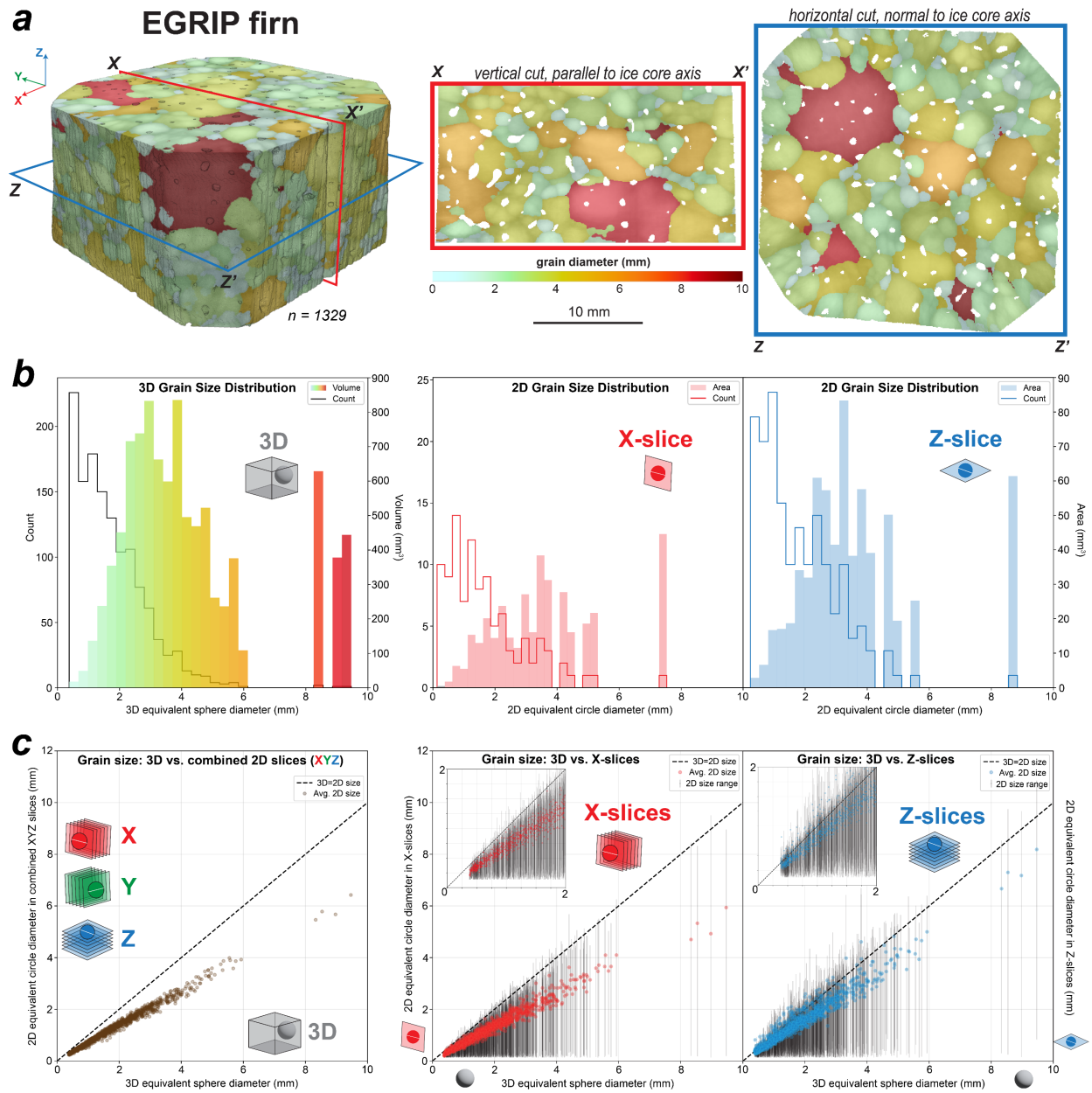
Figure 4. Comparing grain size information between 3D grain maps and virtual 2D slices of the EGRIP firn sample. **(a)** The 3D grain map, colored by grain diameter, was sliced orthogonally. **(b)** Grain size distributions based on 3D data are compared to those calculated in single, orthogonal X and Z section planes through the sample, which simulate vertical (X) and horizontal (Z) cuts typically made of ice cores. **(c)** Multiple XYZ slices in these plots represent 60 μm intervals (i.e., the reconstructed voxel size of DCT data). Grain sizes calculated from 3D volumetric data are plotted against grain sizes calculated from combined XYZ orthogonal slices, hence the 1:1 line, as well as the average 2D size of each grain found across XYZ slices combined. Each grain's average size based on either all X-slices or all Z-slices is also

493   compared to its volume-calculated size, with the full 2D size range for each grain plotted as vertical bars.

494   Note that all plot diameters represent equivalent sphere (3D) or circle (2D) diameters.

495

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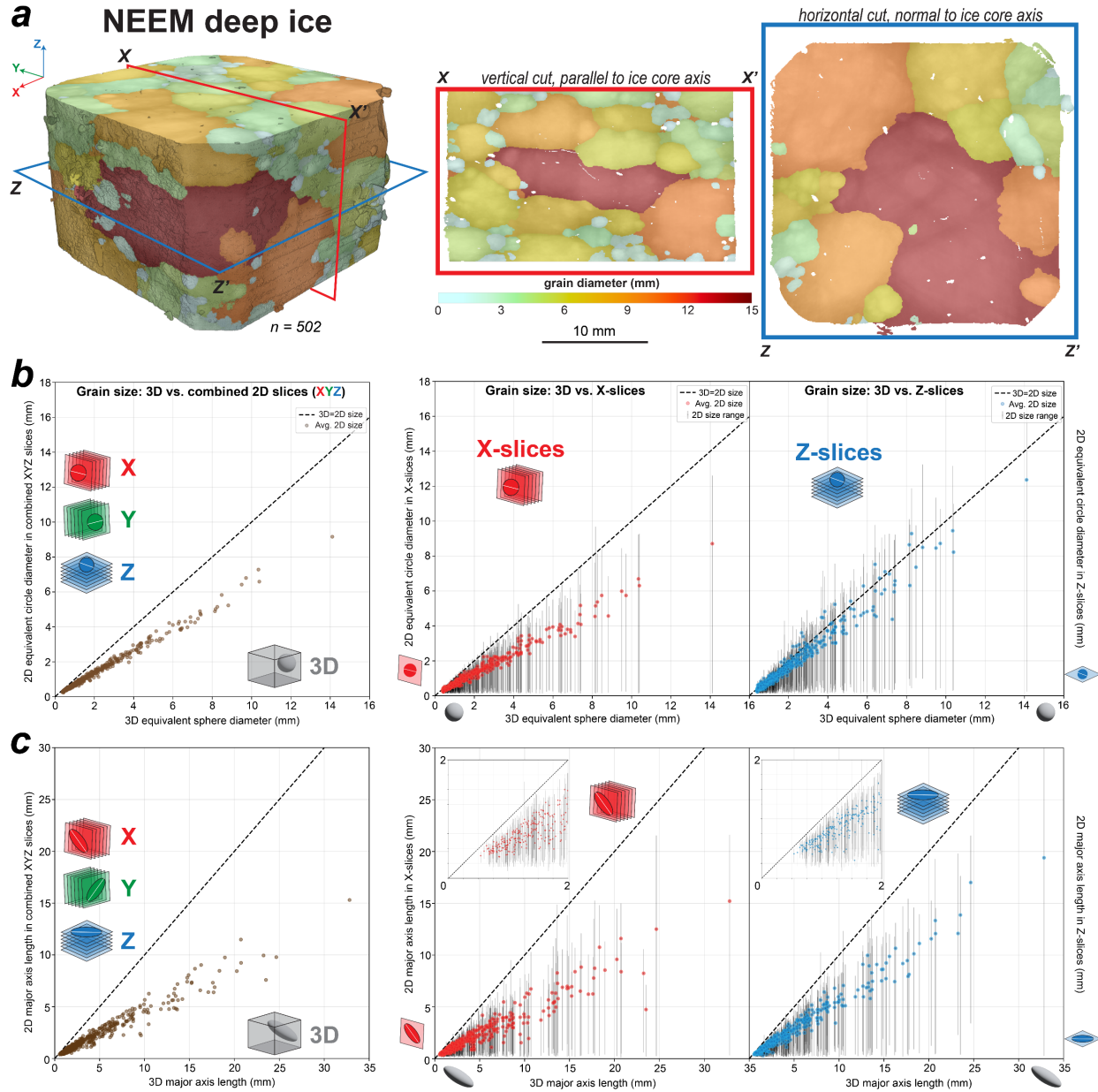
Figure 5. Comparing grain size information between 3D grain maps and virtual 2D slices of the NEEM deep ice sample. **(a)** The 3D grain map, colored by grain diameter, was sliced orthogonally. X and Z slices simulate the typical vertical (X) and horizontal (Z) cuts made of ice cores. **(b-c)** Multiple XYZ slices in these plots represent 60 μm intervals (i.e., the reconstructed voxel size of DCT data). Grain sizes calculated from 3D volumetric data are plotted against grain sizes calculated from combined XYZ

502   orthogonal slices, hence the 1:1 line, as well as the average 2D size of each grain found across XYZ slices

503   combined. Each grain's average size based on either all X-slices or all Z-slices is also compared to its

504   volume-calculated size, with the full 2D size range for each grain plotted as vertical bars. Note that **(b)**

505   plots diameters represent equivalent sphere (3D) or circle (2D) diameters, whereas **(c)** plots the major-

506   axis lengths of grains based on ellipsoid (3D) and ellipse (2D) fitting.

507

508   ## 3.3 Crystallographic fabrics characterized by grain size and *c*- and *a*-axes

509   Glacial ice behaves like a plastically deforming monomineralic rock. Its creeping flow

510   causes crystallographic fabrics and structures to develop to accommodate macroscopic strain,

511   similar to high-grade metamorphic rocks of Earth's crust and mantle. The fabrics are

512   manifestations of a mechanical anisotropy that develops slowly through intracrystalline glide on

513   the basal plane of ice, normal to the *c*-axis of its hexagonal crystal structure (Duval and others,

514   1983). While there are several dislocation slip systems in ice, deformation is dominated by basal

515   slip and explains why the crystal *c*-axis preferentially tends toward the direction of

516   compressional stress (Alley and others, 1988). As a result, polycrystalline ice becomes more

517   difficult to compress but easier to shear as the basal planes of more grains become

518   perpendicular to the compressive direction (Durand and others, 2006). It is therefore necessary

519   to track the strength and evolution of *c*-axis fabrics to determine the induced mechanical

520   anisotropy, which is relevant for modeling the large-scale flow of ice sheets (e.g., Azuma, 1994;

521   Montagnat, Castelnau, and others, 2014; Fan and others, 2021). However, because directions

522   of the three *a*-axes in the basal plane are also found to align depending on the strain history of

523   ice (Montagnat and others, 2015; Qi and others, 2019; Journaux and others, 2019; Monz and

524   others, 2021), *a*-axis measurements might also be needed to fully characterize deformation

525   (Hunter and others, 2023), despite that bulk viscosity is often thought to depend only on *c*-axis

526   directions. For example, *a*-axis measurements could help discern between models of fabric

527   evolution wherein plastic spin predicts different *a*-axis orientation dynamics. Yet, crystal *a*-axes

528   are largely uncharacterized for polar ice cores because standard optical techniques, being the

529   universal Rigsby stage and automatic fabric analyzer, simply cannot be used to measure them

530   (Rigsby, 1951; Langway, 1958; Russel-Head and Wilson, 2001; Wilen and others, 2003).

531   Methods used to simultaneously measure *c*- and *a*-axes in the past, like etching and spot-based

532   Laue X-ray diffraction, are time-intensive and impractical (e.g., Matsuda, 1979; Miyamoto and

533   others, 2011; Weikusat, Miyamoto, and others, 2011). Although modern cryo-EBSD can

534   measure *c*- and *a*-axes and offers advantages in speed, resolution, and angular precision, it is

535   destructive and does not currently offer volumetric information necessary to fully characterize

536   CPOs (Weikusat, de Winter, and others, 2011; Prior and others, 2015).

537   Our coupled measurements of grain volume and full crystallographic orientation enable

538   us to interrogate fabrics in ways not previously possible. For example, in **Figures 6 and 7**, we

539   use volume-weighted pole figures to characterize fabrics at different grain size fractions and

540   compare them to traditional point-scatter and number-weighted pole figures that include all

541   grains. The comparison shows that weighting by volume versus number of grains does not

542   result in a significant difference in the distribution of crystal-axis directions for our specific

543   samples, but the magnitudes of the distributions differ and may be relevant to different

544   questions concerning anisotropy. Such comparisons could be notably different for other ice

545   samples, like those with more diverse grain size distributions.

546   Notably, we observe *c*-axis distributions commonly documented in ice cores, including

547   multimaxima, a broad or tight maximum around vertical, and a vertical girdle distribution,

548   which can be superimposed when plotting all grain sizes together. In other words, we identify

549   that different grain-size populations can contribute differently to the overall fabric developed in

550   each of our samples (**Figs. 6 and 7**; also see **Fig. S4**). Therefore, 3D grain-population fabrics

551   could shed light on deformation or recrystallization mechanisms as they relate to grain size or

552   shape, for example, and may potentially be used to differentiate overprinted kinematics and

553   deformation drivers (Weikusat and others, 2017; Qi and others, 2019; Fan and others, 2020;

554   Stoll and others, 2025). Additionally, while several studies have used high grain number

555   statistics to conclude that fabric strength can develop independently of grain size with depth,

556   those relationships were assessed based on measured changes in the mean crystal area across

557   thin sections (e.g., Svensson, Schmidt, and others, 2003; Svensson, Baadsager, and others,

558   2003). Pulling apart fabrics using volume-based grain populations, in hindsight, may even

559   permit improved interpretations of optical thin-section images that cover centimeters to

560   meters worth of ice cores.

561   Segregating *c*-axis fabrics may ultimately help characterize *a*-axis fabrics, a challenge

562   inherent to the presence of three indistinguishable *a*-axes in the basal plane. Distributions of *a*-

563   axes appear considerably weaker by comparison, therefore we plot *a*-axes on a separate scale

564   unbound by the intensity of *c*-axis CPOs in **Figures 6 and 7**. For most grain size fractions in our

565   samples, the *a*-axes define a broad girdle that hosts multimaxima on a plane roughly normal to

566   the ice-core axis; whereas, *a*-axes of the largest grain sizes can deviate away from this plane

567   considerably. The presence of *a*-axis CPOs in our data corroborates those observed in

568   experiments and mountain glaciers, collectively indicating that slip is anisotropic in the basal

569   plane wherein *a*-axes lie (Montagnat and others, 2015; Qi and others, 2019; Journaux and

570   others, 2019; Monz and others, 2021; cf. Kamb, 1961), whether or not it has first-order effects

571   on the viscous anisotropy. We note, however, that because ice cores rotate when they are

572   brought to the surface, the unconstrained nature of pole figures with respect to the horizontal

573   direction may make it challenging to interpret *a*-axis CPOs geographically (Westhoff and others,

574   2021), likely requiring statistically robust datasets to fully evaluate their importance.

575
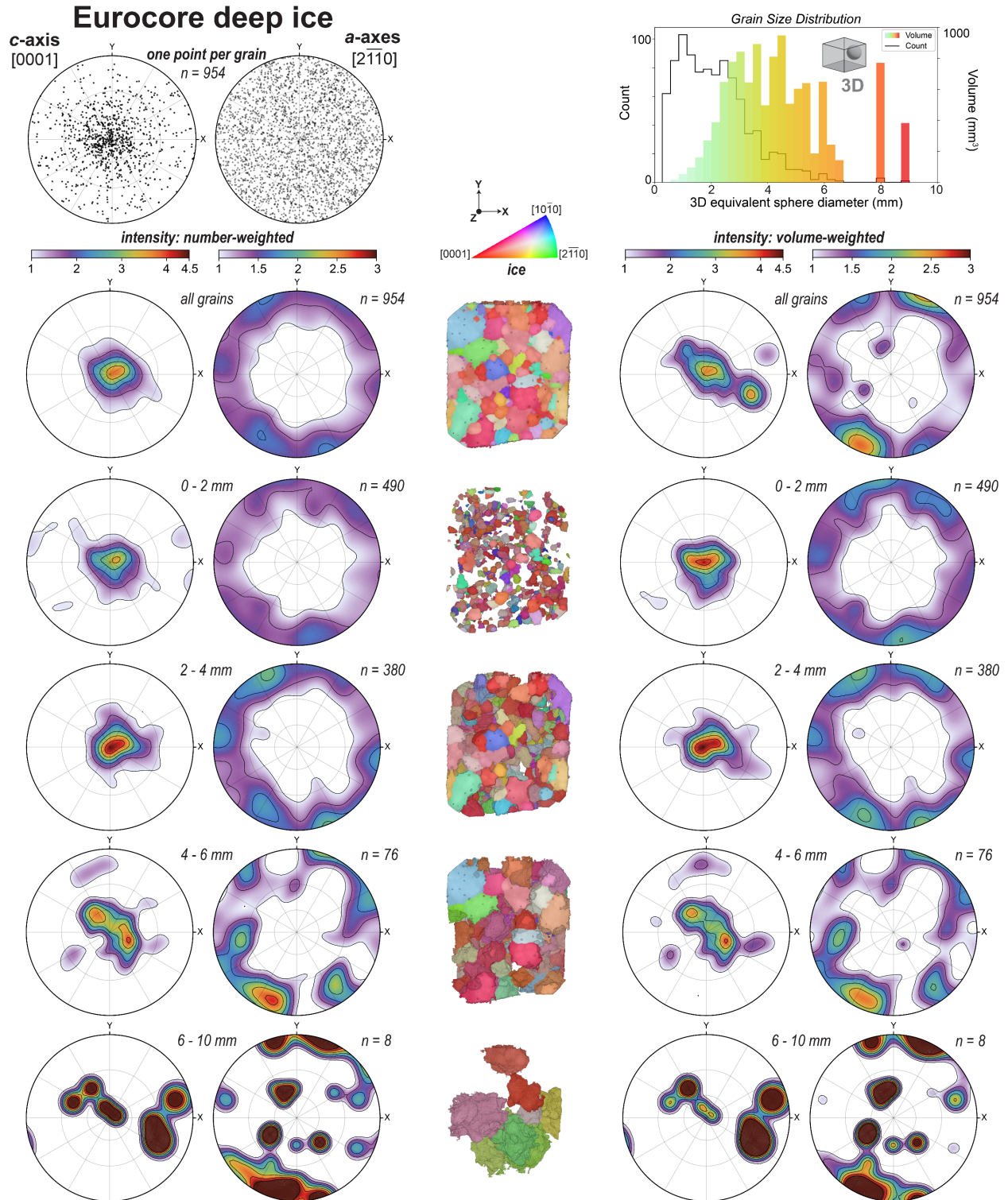
Figure 6. Comparing pole figure analyses for the Eurocore sample's grain orientations based upon count, volume, and grain size fraction. Pole figures provide views down the sample Z-axis, parallel to the ice-core axis. Analytical details are provided in the **Methods**. Note that glaciological studies employing EBSD

580 of ice may instead represent the *a*-axes as $[11\bar{2}0]$ (-a$_3$), which is symmetrically equivalent to $[2\bar{1}\bar{1}0]$ (+a$_1$)

581 (e.g., see Qi and others, 2019).

582
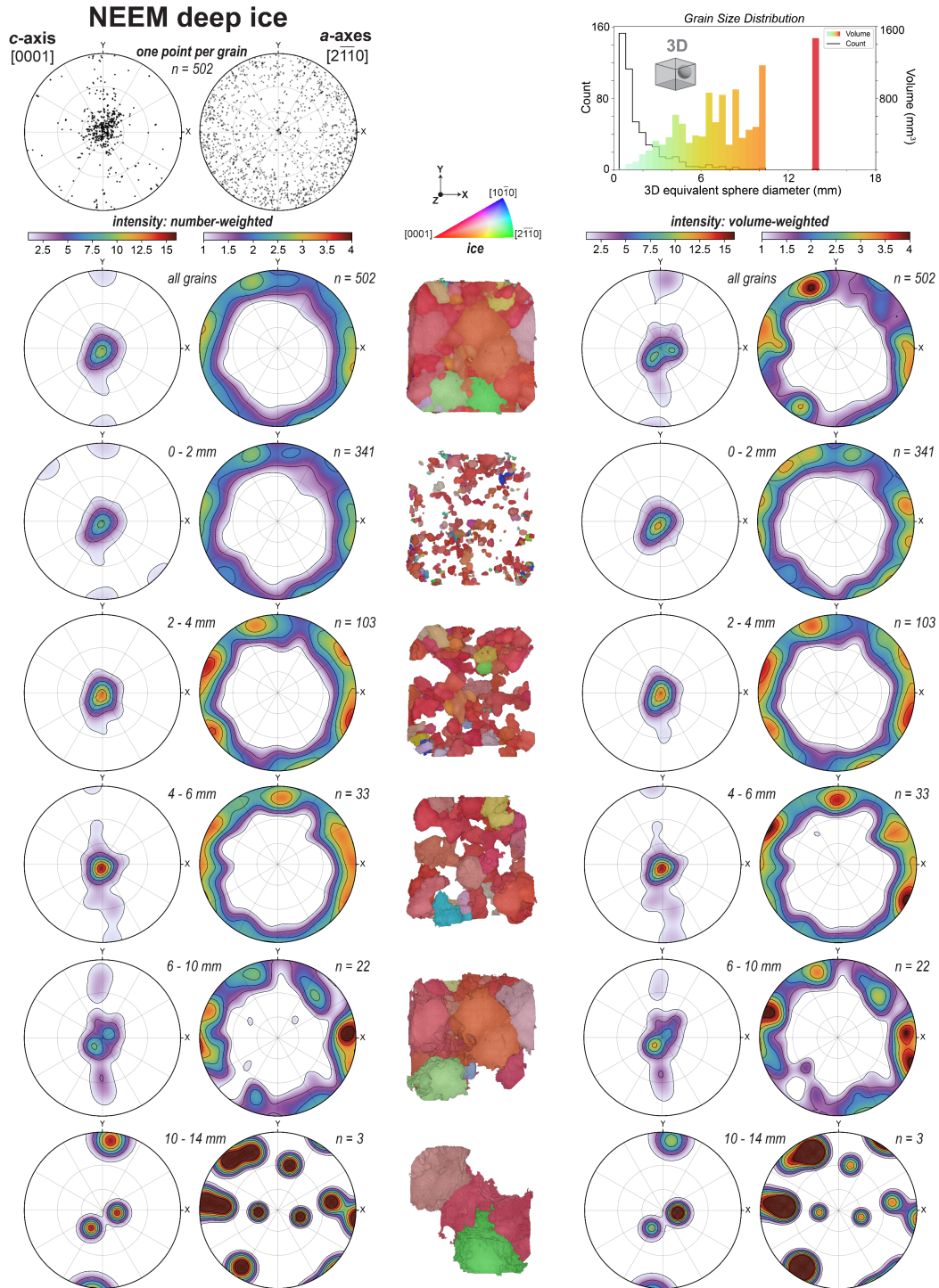
Figure 7. Comparing pole figure analyses for the NEEM sample's grain orientations based upon count,

volume, and grain size fraction. Pole figures provide views down the sample Z-axis, parallel to the ice-

core axis.

586

3.4 Misorientations across grains

588   Many microscopic mechanisms operate during ice deformation and to different extents

589 over the depth of ice sheets, ultimately affecting the ice's macroscopic flow behavior

590 (Thorsteinsson and others, 1997; De La Chapelle and others, 1998; Montagnat and others,

591 2015; Ranganathan and others, 2021). Recrystallization mechanisms activate to reduce the

592 strain energy that accumulates as neighboring grains interact and develop internal stress and

593 strain fields, which can be highly heterogeneous near grain boundaries (Duval and others, 1983;

594 Montagnat and others, 2015). These processes can lead to grain boundary migration, grain

595 growth, or nucleation and concentrate crystal lattice dislocations into substructures and

596 subgrain boundaries (Montagnat and others, 2015; Rollet and others, 2017).

597   Measuring crystallographic misorientation gradients across grains is commonly used to

598 understand dislocation density and recrystallization mechanisms, where the density may

599 comprise only some dislocation types based on their accessibility (e.g., Montagnat and others,

600 2015). In this context, cryo-EBSD has proven highly useful for mapping low-angle

601 misorientations at resolutions of 10 μm or higher, although some studies have employed lower

602 (e.g., 50 μm) resolutions to characterize subgrains (e.g., Montagnat and others, 2015, also see

603 Weikusat, Miyamoto, and others, 2011). Optical-based techniques, like AFAs, similarly provide a

604 high (6 μm) spatial resolution but a lower angular resolution (3°) compared to EBSD (0.1-1°)

605 (Prior and others, 1999; Winkelmann and others, 2020). However, cryo-EBSD and AFAs share

606 the same disadvantages in that they can only uncover information from cut surfaces, which

607 generates uncertainties in the representation of grain shapes and sizes. In contrast, lab-based

608 DCT can offer a spatial resolution down to about 30 μm for our specific sample sizes and a high

609 (0.1°) angular resolution (Sun and others, 2022). While the 2D-based techniques can cover large

610 sample areas (up to 7 x 3 cm for EBSD and 10 x 10 cm for AFA; Prior and others, 2015, Faria and

611 others, 2014), DCT provides the advantage of working with bulk samples that yield volumetric

612 information and maintains 3D contexts crucial to interpretations. Therefore, a correlative 3D-

613   2D microscopy method could generate complementary data sets that offer powerful multiscale,

614   multidimensional information.

615   Our large sample volume, experimental setup, and DCT measurements permitted

616   realistic grain map reconstructions down to 60 μm voxel sizes, comparable to EBSD resolutions

617   applied previously (e.g., Montagnat and others, 2015). At this resolution, we were able to

618   resolve what appear to be subgrains, whose boundaries strongly correspond to those seen in

619   grain completeness maps (**Fig. 3**), corroborating the evidence of intracrystalline deformation in

620   X-ray diffraction patterns (**Fig. S1**). Notably, DCT grain maps provide an indication of differences

621   in the relative amount of grain deformation (**Fig. 3**; also see section 3.5 and **Fig. S5**). We

622   observe, for example, that the amount of intragranular misorientation is weakly dependent on

623   the size of a grain, but this relationship varies with the degree of deformation in our samples

624   and is not straightforward. Specifically, smaller grains in the EGRIP firn sample tend to exhibit

625   higher misorientation angles and across distinct subgrain boundaries (see **Fig. 3**; also see **Fig.**

626   **S5**). Conversely, in the Eurocore and NEEM samples, larger grains tend to show a greater

627   amount of intragranular misorientation than the very smallest grains, but the spread in

628   misorientation angles varies between the larger grains (**Fig. S5**). The weak relationships may

629   indicate a dependency not only on grain size but other factors, such as the average grain

630   orientation with respect to far-field or even local stresses (see Section 3.5). At our resolution,

631   we also see that the largest grains of these deeper and more deformed samples lack the

632   distinct subgrain boundaries seen in the firn sample of **Figure 3** (compare to **Fig. 10** in Section

633   3.5). In the following section, we integrate these observations with those of bubble shapes and

634   discuss how 3D misorientation data could help constrain grain-scale deformation histories.

635

636   3.5 Spatial relationships between ice grains and pore spaces

637   The decrease in ice sheet porosity with depth reflects the densification of accumulated

638   snow and its progressive burial and deformation. An interplay exists between air-bubble and

639   grain-scale processes that must be understood when it comes to (1) measuring atmospheric

640   gases trapped in ice cores and (2) constraining controls on ice deformation. Knowing precisely

641   where and how air bubbles are occluded in the firn column is crucial for determining the age

642   difference between bubbles and the encapsulating ice (Burr and others, 2018; Westhoff and

643   others, 2024). Similarly, knowing where atmospheric gases become occluded in clathrates

644   under immense compression at depth, as well as how ice-core relaxation affects the

645   retransformation of clathrates to air bubbles, is crucial for understanding gas behavior (e.g.,

646   Pauer and others, 1996; Kipfstuhl and others, 2001). Further, knowing how grain-growth

647   mechanisms (that help to seal bubbles) are affected kinetically by the bubbles themselves is

648   important for understanding grain-scale deformation near the firn-ice transition (Roessiger and

649   others, 2014; Fegyveresi and others, 2019; Fan and others, 2023). Anisotropic bubble shapes in

650   deep ice may also record the stress distribution on individual grains and improve our

651   understanding of local and far-field stresses needed to model ice flow.

652   In **Figures 8-11**, we demonstrate how correlating bubbles (generally pores) with the 3D

653   grain structure allows an integration of qualitative and quantitative information that is difficult

654   to access from 2D sections of ice cores. By visualizing pores and grains as surface meshes

655   created from 3D maps, we can view bubbles from the perspective of their host-grain

656   orientations, which becomes a powerful guide for quantitative analysis.

657   Bubble position determined simply as intergranular versus intragranular is an important

658   parameter that provides insight on gas pathways, grain-growth behavior and ice-air boundary

659   mobility (e.g., grain-boundary migration and bubble pinning; Roessiger and others, 2014). In

660   **Figure 8,** we paired spatial classifications with bubble sizes and found expected differences

661   between the firn and deep ice samples. In the EGRIP firn sample, the size distribution of

662   bubbles that intersect grain boundaries is distinct from bubbles positioned wholly within grains.

663   Different types of bubbles are identifiable by their spherical, tube-like, or amoeboid shapes, the

664   latter two of which are visibly larger. Therefore, even without quantitative measures of the

665   complex bubble shapes, one can infer from the count and volume distribution data that more

666   non-spherical bubbles intersect grain boundaries (cf. Kipfstuhl and others, 2009). Differences

667   between inter- and intra-granular bubble size distributions in the Eurocore sample are more

668   subtle, consistent with grain growth increasing the numbers of intragranular bubbles in deeper

669   ice (De La Chapelle and others, 1998).

670
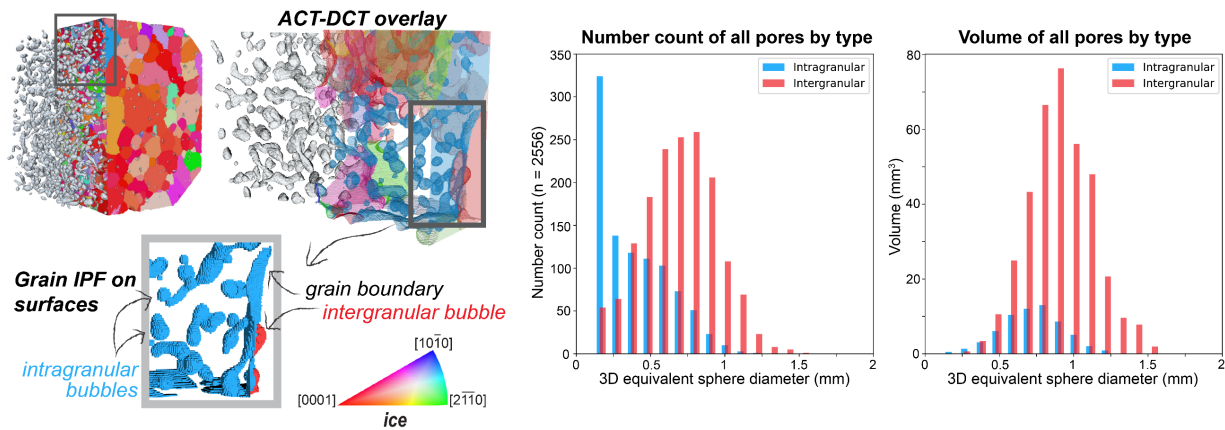
EGRIP firn bubbles
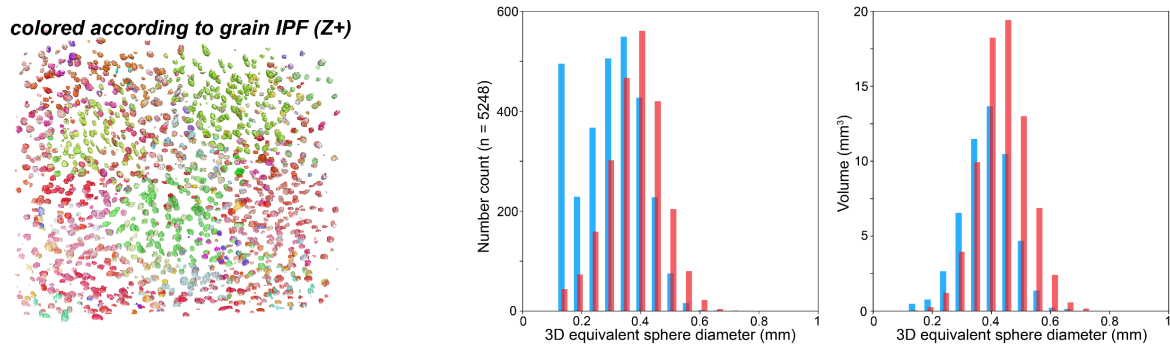
Figure 8. Size distributions of air bubbles categorized by their intergranular versus intragranular positions. Pore spaces in EGRIP firn and Eurocore deep ice samples represent primary air bubbles captured by ice during the densification process. The overlay of ACT and DCT surface meshes enable visualizations and correlations of bubbles and grain boundaries. Bubble surfaces are colored with respect to the IPF(Z+) color of their host grain(s) and therefore change color across the grain boundaries.

Bubble shapes in deep ice have been proposed as possible strain gauges, given bubbles deform ~1.7 times faster than the surrounding ice (Alley and Fitzpatrick, 1999; Fegyveresi and others, 2019). Previous data for elongated bubbles, as well as slightly nonspherical ones typical of deep ice cores, have led to the physical understanding that diffusion restores bubbles toward steady-state spherical shapes. Bubbles deform, however, if the diffusion rate is slower than that of deformation. Flattened or elongated bubbles are therefore thought to record strain rates and directions of deformation, which are expected to differ between grains and depend on both local and far-field stresses (Fegyveresi and others, 2019). A closer look at the Eurocore sample in **Figure 9** revealed bubble-shape anisotropies that are strongly grain-dependent. More

687   specifically, intragranular bubbles belonging to the same grain are visibly flattened or elongated

688   in roughly the same plane, but their sizes and aspect ratios vary (**Fig. 9**). Bubble-shape

689   orientations also tend to vary considerably near grain margins. Such observations are consistent

690   with those made previously in ice cores (e.g., Alley and Fitzpatrick, 1999; Fegyveresi and others,

691   2019), however, past studies employed 2D techniques that imposed measurement limitations,

692   which did not allow bubble shapes to be used to their full potential.

693   For the 954 grains in the Eurocore sample, we analyzed 5248 bubbles, of which 2647 are

694   intragranular and do not intersect the sample margins (see **Methods**). Of these intragranular

695   bubbles, approximately 65% are oblate (disk-shaped), 15% are prolate (rod-shaped), and 20%

696   are roughly spherical in shape (ellipsoid factor of -0.05 ≤ x ≤ 0.05) (**Figs. 9 and 10**). There are no

697   clear correlations between bubble shape, volume, and host-grain volume in this sample (**Fig.**

698   **10**; **Fig. S6**). There is a strong correlation, however, between the directions of bubble-

699   flattening/elongation and the *c*-axis and basal plane of the host grain (**Fig. 9**). An analysis of all

700   2647 intragranular bubbles indicates they have a strong preference to flatten and elongate in

701   the grain basal plane, the dominant slip plane in ice during deformation. Based on bubble

702   major-axis directions, we determined that 76% of them are elongated within 20° of the basal

703   plane (**Fig. 9**), while 57% are within 10°. The minor-axis directions indicate that 57% of bubbles

704   have a flattening, or shortening, direction within 20° of the *c*-axis. Furthermore, the angle

705   between the bubble direction and basal plane does not appear to be strongly influenced by the

706   orientation of the host grain with respect to the compressive direction (Z) (see **Fig. S7**), even for

707   grains whose *c*-axis is oriented at a high angle to Z, further supporting there is a strong

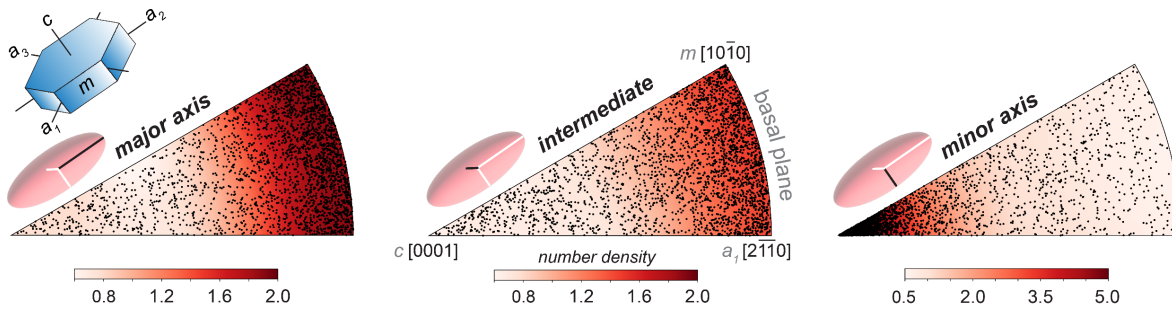708   crystallographic control on bubble deformation.
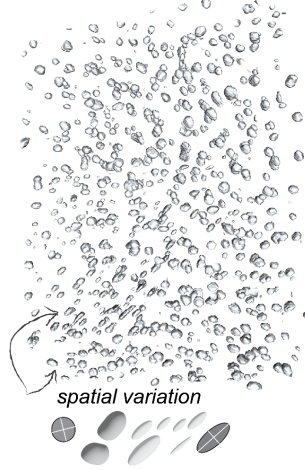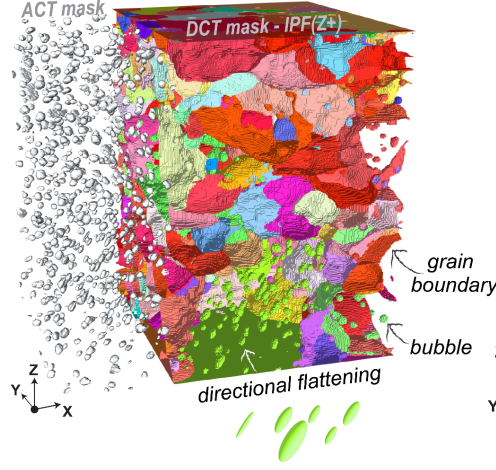
709
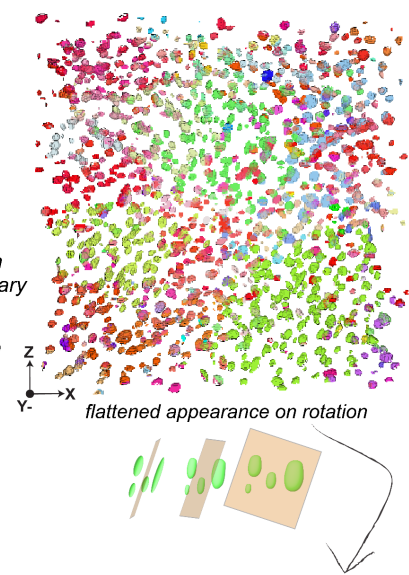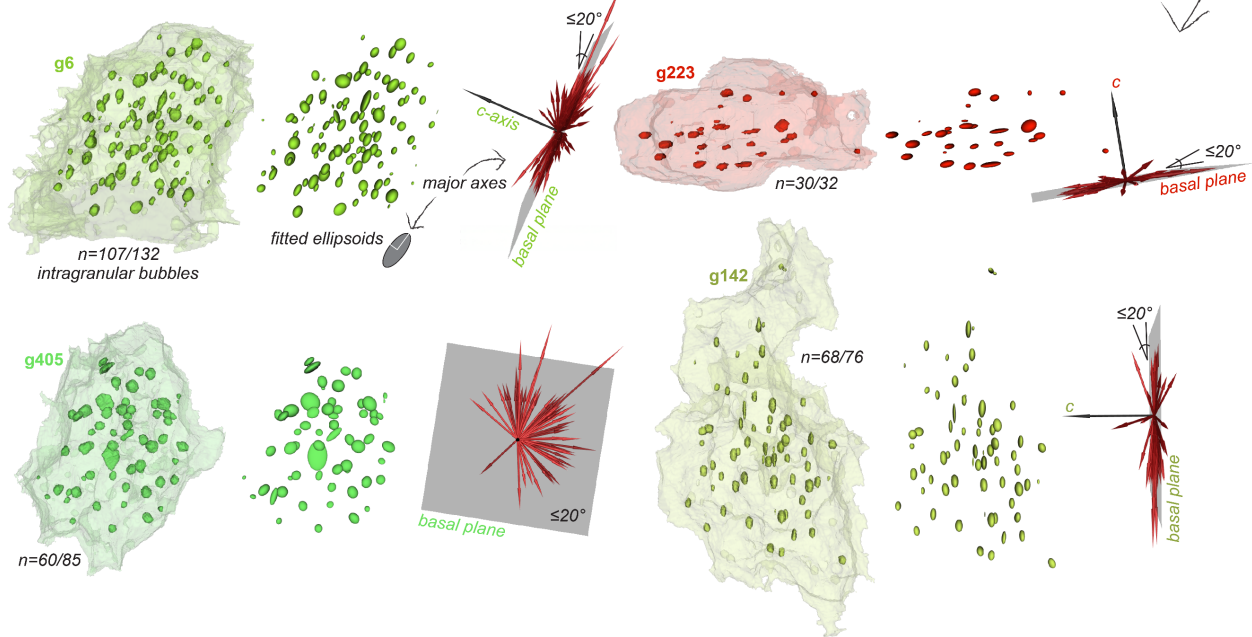
Figure 9. Bubble-grain relationships identified through correlative, multimodal imaging. Bubble-shape

anisotropies in the Eurocore sample were first observed using an ACT-derived surface mesh (gray). By

713 overlaying the ACT bubble-surface mesh with a grain-surface mesh generated from the DCT volume (IPF

714 color), the visualization of ACT bubbles colored by the IPF orientation of their host grain(s) led to the

715 observation that bubble anisotropy is grain-dependent. Rotation of the sample revealed most bubbles

716 inside the same grain exhibit flattening or elongation in roughly the same plane, although variability

717 occurs near grain boundaries. The three largest grains (g6, g405, g142) and another grain (g223) with its

718 *c*-axis near-vertical (Z) are shown with intragranular bubbles and their fitted ellipsoids (see **Methods**).

719 Major-axes of the ellipsoids are plotted as red vectors scaled by the ellipsoid aspect ratio, and only for

720 bubbles whose major-axis lies within 20° of the grain basal plane (the number fraction of these bubbles

721 is given as *n* for each grain). The dominant oblateness of bubbles with preferential flattening or

722 elongation in the grain basal plane is expressed collectively by IPFs of the fitted-ellipsoid axes. The IPFs

723 represent all intragranular bubbles plotted with respect to the crystallography of their host-grains,

724 which have been rotated into a single reference frame.

725

726 While some of our observations of bubble-grain relationships are comparable to those

727 of Fegyveresi and others (2019), their 2D optical technique offers measurements within one

728 section plane and typically of only *projections* of elongations. In **Figure 10**, we demonstrate the

729 limited information we would gather from a single section plane through different grains in our

730 Eurocore sample. Only by searching and selecting slices through 3D data could we find bubbles

731 with the highest *projected* "elongation" that could then be verified against a *projection* of the

732 host-grain *c*-axis (determined from pole figures). Accurate measures and the true spatial

733 variability of bubble shapes across the bulk sample would never be captured without 3D

734 information in this case. Nonetheless, we find it interesting that Fegyveresi and others (2019)

735 found correlations between bubble shape, size, and grain size in their Antarctica sample,

736 whereas we do not see clear trends. Although we can identify that grain size and the amount of

737 intragranular misorientation is weakly correlated (see section 3.4; **Fig. 10**; **Fig. S5**), we do not

738 see a clear relationship between these characteristics and the magnitude of bubble-

739 flattening/elongation, for example. One possibility is that grain-scale interactions generate

740 differences between local and far-field stresses, resulting in different deformation rates for

741 individual grains and variance in observable patterns (Fegyveresi and others, 2019). We suggest

742 another possibility is that the Eurocore sample's bubbles have undergone enough diffusional

743   restoration to erase any such correlations, if they ever existed. After all, the current bubble

744   configurations can be assumed to represent both the cumulative deformation and the balance

745   between diffusion and deformation over the entire history of the ice (Fegyveresi and others,
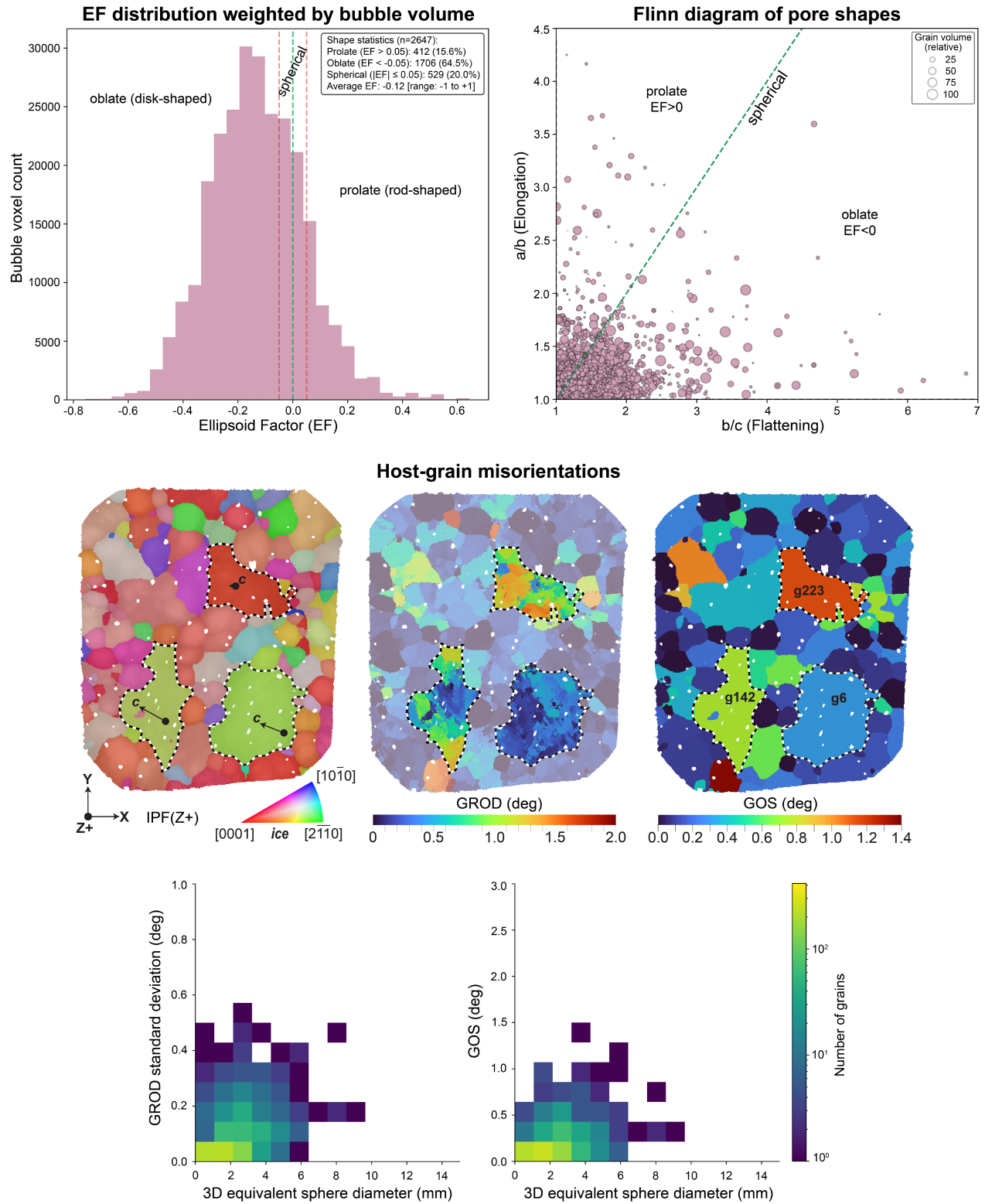
746   2019).

747

Figure 10. Examples of possible relationships and non-relationships among intragranular bubbles and their host-grains in the Eurocore sample. The ellipsoid factor (EF) was used to classify bubbles as prolate, oblate, or spherical based on their fitted ellipsoids. As depicted in the EF distribution plot, the

748

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751

752 boundaries for "spherical" bubbles were chosen at -0.05 ≤ EF ≤ 0.05. The Flinn diagram reflects these

753 shape classifications but adds information about host-grain volume, which has been normalized. A single

754 slice through the 3D grain map intersects three of the grains isolated in **Figure 9**, which are among the

755 eight largest grains. Notably, these grains show a relatively high amount of intragranular misorientation

756 but different GOS values. Plots of grain size versus GOS (average GROD) and the standard deviation of

757 GROD show weak correlations (also see **Fig. S5**).

758

759 On the other hand, at depths greater than the Eurocore sample, immense compression

760 causes gases to transform into clathrates and leads to the eventual loss of all original bubbles in

761 ice. However, following the recovery of bubble-free ice cores, the stress release can cause

762 fractures to form and bubbles to redevelop from clathrates (e.g., Pauer and others, 1996;

763 Kipfstuhl and others, 2001). These secondary processes have received considerable attention

764 because they may lead to gas escape and affect gas concentration measurements that are

765 notably destructive for ice cores.

766 The pore spaces in the NEEM sample are bubbles and cracks that have formed from

767 decomposed clathrates and relaxation since the core was drilled in 2010, although many pores

768 have a flattened, yet graupel-like appearance characteristic of primary clathrates (Pauer and

769 others, 1996; Kipfstuhl and others, 2001). These features are interesting because most pores

770 and apparent fracture planes have a strong preference to align with the basal plane of the host

771 grain (**Fig. 11**). We determined this by analyzing all 2424 intragranular pores (out of 3276 total

772 pores) for the 502 grains identified in the sample. The NEEM pore results (**Fig. 11**) are

773 remarkably similar to those of the Eurocore sample (**Fig. 9**), suggesting that the grain

774 crystallography may also exert a strong crystallographic control on relaxation processes and the

775 evolution or retransformation of clathrates. The visualization and statistical analysis of 3D sizes,

776 shapes, and spatial distributions of these features could be paired with Raman spectroscopy

777 (Pauer and others, 1996), for example, to contextualize and ensure accurate measurements of

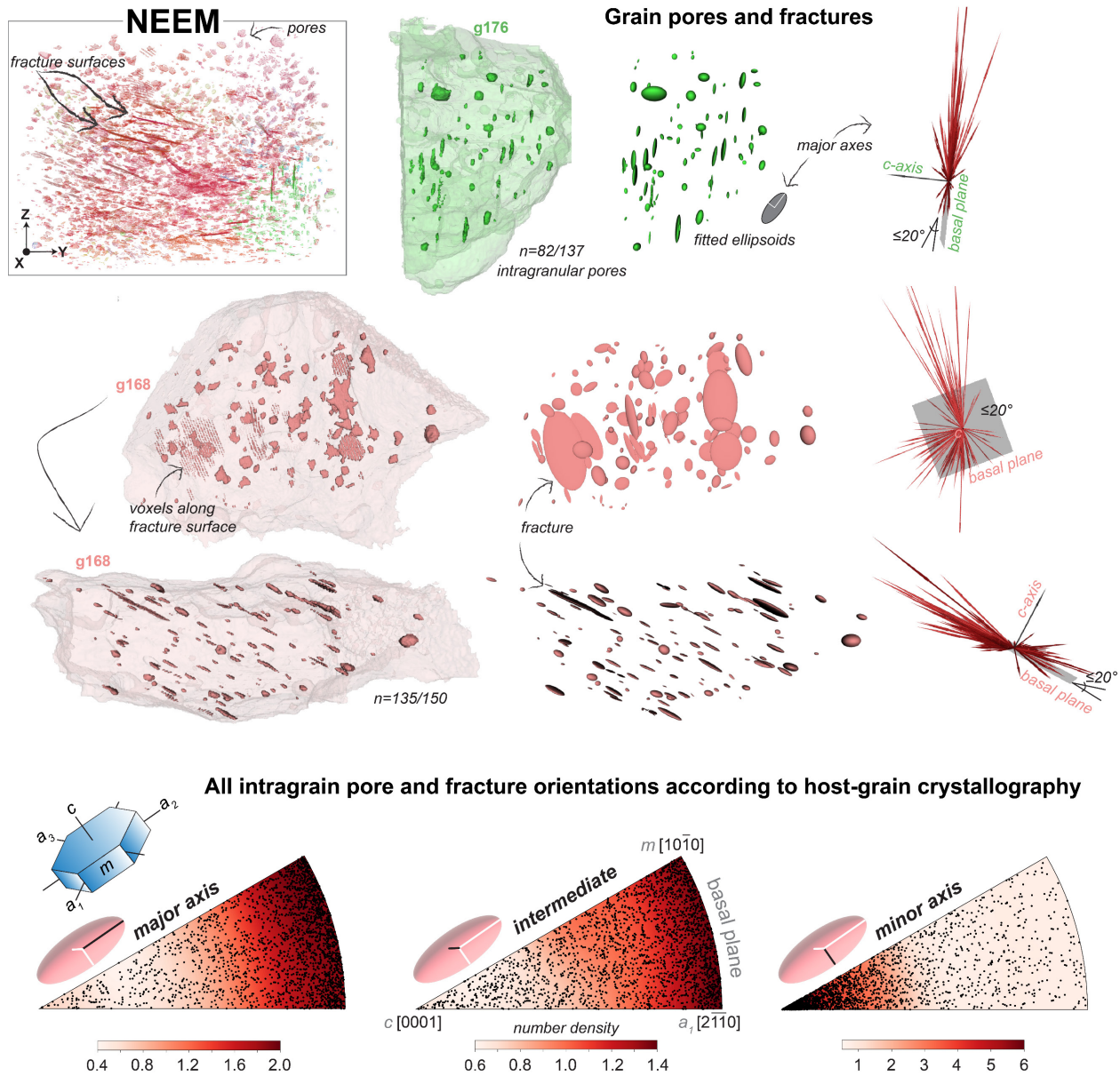778 the ancient atmospheric gases in ice cores.

779

Figure 11. Pore-grain relationships in the NEEM sample. The visualization of pore spaces and apparent microfractures throughout the sample (top left) were performed in the same way as described for the Eurocore sample in **Figure 9**. Two of the largest grains with the *c*-axis near-vertical (g168) or near-horizontal (g176) with respect to the ice-core axis are shown with intragranular pores and their fitted ellipsoids (see **Methods**). Grain 168 is shown in two different views; one view looks down the *c*-axis and the other looks normal to it. The features we refer to as microfractures are very thin and require careful segmentation of our data to preserve the fracture surfaces, which, at our data resolution, appear as adjacent lines of voxels. Our connected component analysis allowed the merging of these neighboring

789 voxels to fit an ellipsoid to each fracture. The vector plots and IPFs here are the same as described for

790 **Figure 9**.

791

4 CONCLUSIONS

793 Our pilot study demonstrates that ice core microstructures, fabrics, and air bubbles can

794 be imaged in 3D using multimodal X-ray imaging in a lab setting. We show that samples as large

795 as 3 x 3 x 3 cm^3 are achievable and balance the advantages and disadvantages of imaging what

796 are considered large volumes for the multimodal X-ray technique. These sample sizes allowed

797 us to achieve adequate resolutions and volume-based statistics under reasonable acquisition

798 times, and they provide the opportunity to vary the spatial resolution to investigate specific

799 scientific questions. Our novel method for measuring, mapping, and analyzing ice textures can

800 be combined with established ice-core analytical techniques to generate complementary 3D

801 and 2D datasets. Such datasets can maximize scientific insight through the benefits of 3D

802 sample curation, volumetric measurements, shape determinations with improved accuracy,

803 and the integration of complete ice-grain crystallographic orientations. Together, the analytical

804 method and our results demonstrate great potential for 3D textural data to advance ice-core

805 paleoclimate studies and the future development of more integrated models of grain-scale ice

806 deformation and large-scale ice flow.

807

5 OUTLOOK

809 Using 3D data to improve the accuracy of micromechanical models could benefit from

810 higher-resolution (<10 µm) multimodal scanning, which is needed to resolve intracrystalline

811 deformation features to a greater accuracy. Such high-resolution 3D datasets would help

812 constrain the complex behaviors of grains, bubbles, and potentially large impurities as they

813 interact during ice deformation, which is currently challenging to parameterize in models. One

814 caveat of going to higher resolutions is that the experiments will require smaller, less

815 representative sample volumes. This may require design modifications to the current cooling

816 device, and perhaps the development of a correlative, multiscale approach to data collection.

817   In general, our analytical method can advance the level of textural characterizations of

818 ice-core samples, which has the potential to improve paleoclimate reconstructions and

819 micromechanical models. Our results demonstrate the ability to detect 3D variations in ice-core

820 textures from different depths and climatic periods. Therefore, examples of future studies that

821 could apply our new method include stratigraphic sampling to investigate bubble closure,

822 clathrate formation, grain-scale deformation, or the 3D evolution of ice fabrics over the depth

823 of the ice sheet. Additionally, experimental studies could also employ our method to

824 understand, for example, (1) the effects of melting on ice textures using controlled temperature

825 changes and time-lapse experiments, or (2) the rates of microstructural evolution and

826 bubble/clathrate transformation under ambient pressure through interrupted time-series

827 experiments.

828

829 **AUTHOR CONTRIBUTIONS**

830 **Olivia A. Barbee**: Conceived the project, designed experiments, performed imaging, performed

831 data reconstruction, performed data visualization and analysis, created figure illustrations,

832 wrote the manuscript. **Jette Oddershede**: Performed simulations to guide experimental design,

833 supported imaging and data reconstruction, performed data post-processing, supervised the

834 project. **Ravi Raj Purohit Purushottam Raj Purohit**: Wrote code for data analysis and

835 visualization, performed data analysis; **Håkon W. Ånes**: Wrote code for data analysis and

836 visualization, performed data analysis. **Jonas Engqvist**: Designed experiments, built and tested

837 the cooling device. **Anders Svensson**: Selected and prepared samples, provided scientific insight

838 and discussion. **Nicholas M. Rathmann**: Provided scientific insight and discussion. **Thomas

839 Blunier**: Provided scientific insight and discussion. **Florian Bachmann**: Provided data

840 reconstruction and software support. **Stephen Hall**: Conceived and supervised the project,

841 designed experiments. All authors contributed to the final version of the manuscript.

842

843

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COMPETING INTERESTS

The authors declare no competing interests.

DATA AVAILABILITY

Data reported in this paper are available on request given individual file sizes reach 10s of GB.

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1159

1160 **Mapping textures of polar ice cores using 3D laboratory X-ray microscopy**

1161

1162 Olivia A. Barbee[1,2], Jette Oddershede[2], Ravi Raj Purohit Purushottam Raj Purohit[2], Håkon W.

1163 Ånes[2], Jonas Engqvist[1], Anders Svensson[3], Nicholas M. Rathmann[3], Thomas Blunier[3], Florian

1164 Bachmann[2], and Stephen Hall[1]

1165

1166 **SUPPLEMENTARY MATERIAL**

1167
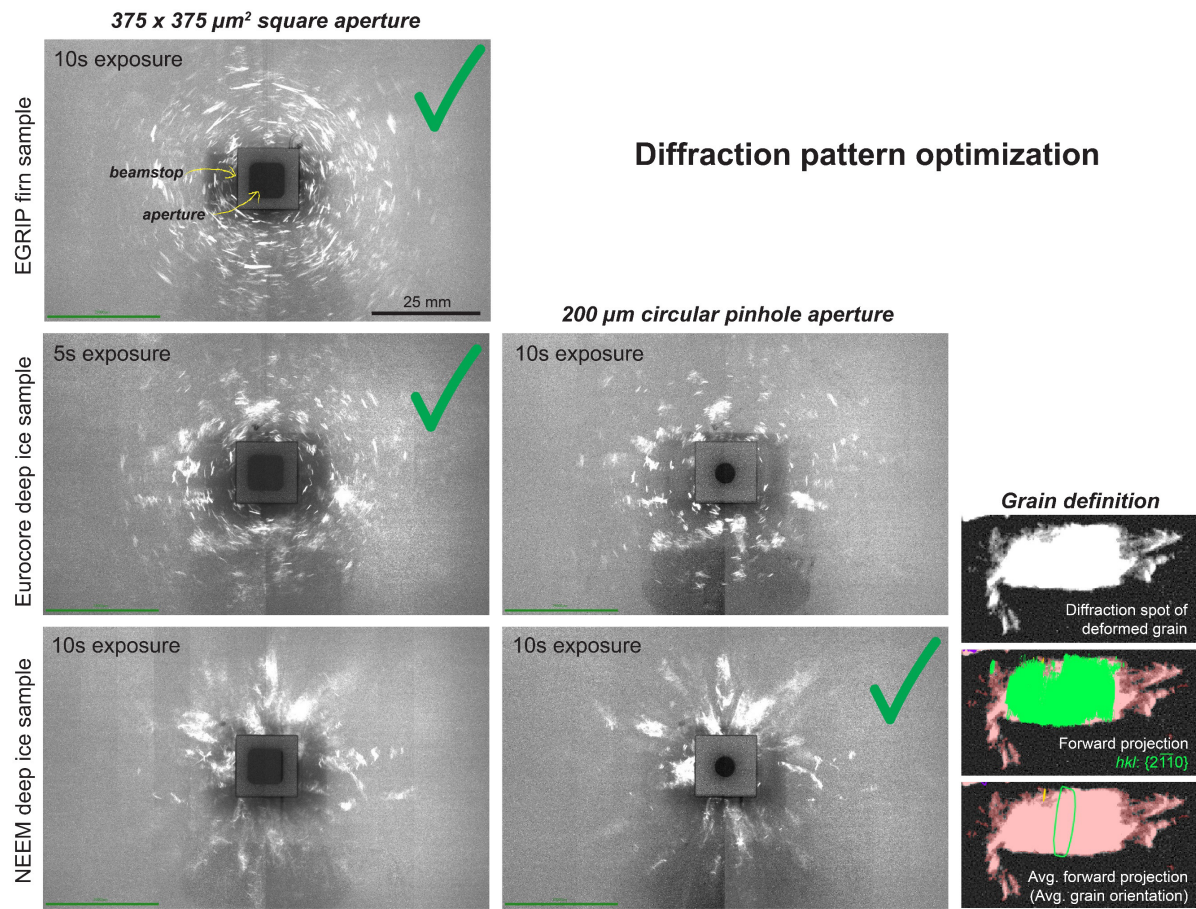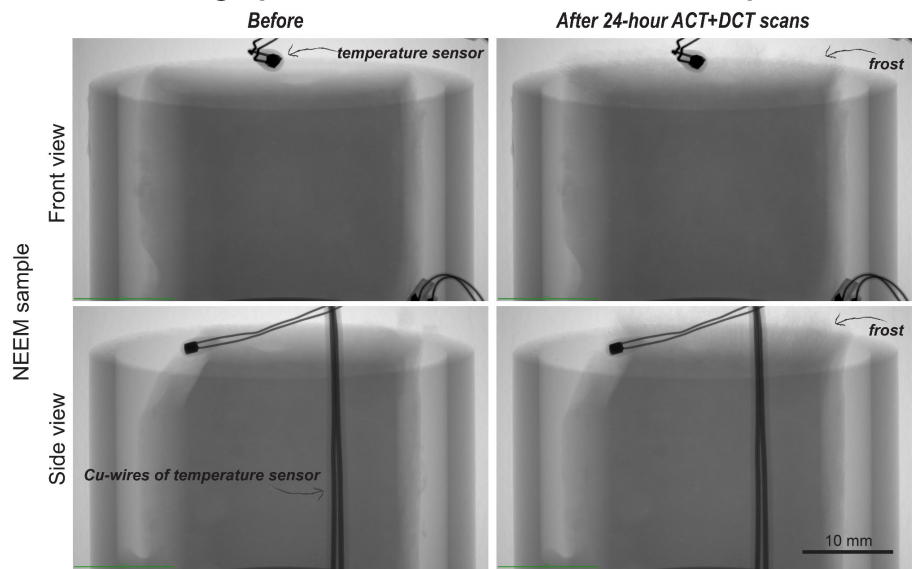
Figure S1. Examples of X-ray projections. At the top, raw X-ray diffraction patterns demonstrate how the source-beam aperture affects the shapes of diffraction spots when using the beam geometry and flat panel detector as described in the **Methods**. The green checkmarks indicate what aperture and

1172   exposure time proved optimal for DCT data acquisition on samples with different amounts of
1173   deformation. Binarized images to the right show how deformed grains are defined using a magnified
1174   area of a diffraction projection from the NEEM sample. The elongated diffraction spot associated with
1175   cloudy areas of intensity represents diffraction primarily from a large, distorted grain with orientation
1176   spread, and the linear upper and lower boundaries of the most intense area reflect confinement of the
1177   source-beam to the pinhole aperture. The pink transparent region represents the segmentation of
1178   diffraction spot(s), and the numerous green lines represent the forward-projected outlines of a single
1179   {*hkil*} (lattice plane) family (one of several used for indexing grains; Bachmann and others, 2019). When
1180   grains are defined during reconstruction, the average forward projection shown represents the average
1181   grain orientation. Below, the radiographs of the NEEM sample show its appearance before and after a
1182   24-hour-long period of ACT-DCT data acquisition. The only notable difference in the "after" image is the
1183   presence of a frost layer on the upper surface of the sample. The uppermost temperature sensor and its
1184   copper wires are in the field of view. The wires had negligible effects on diffraction projections.
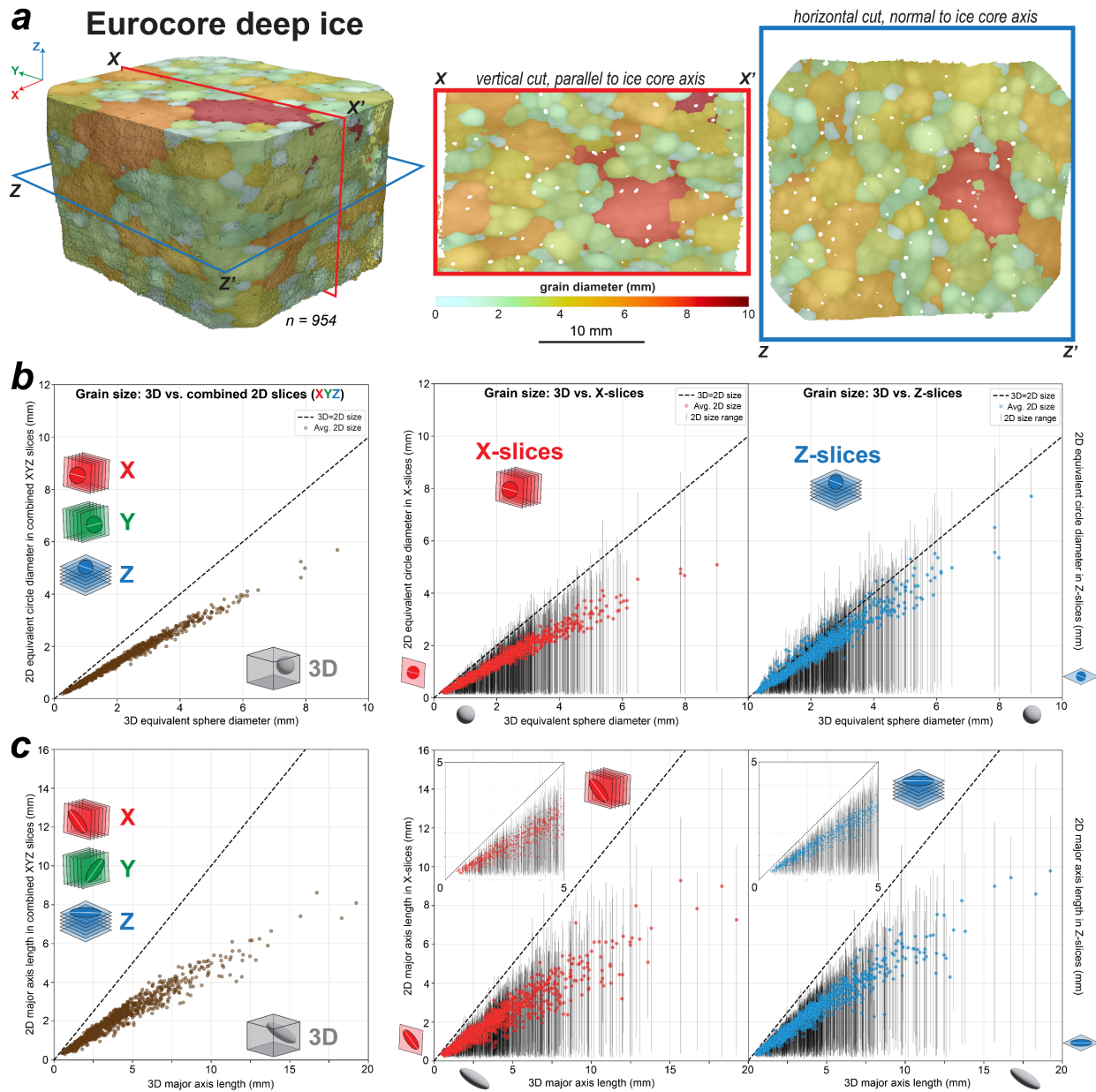1185

Figure S2. Comparing grain size information between 3D grain maps and simulated 2D slices of the Eurocore deep ice sample. **(a)** The 3D grain map, colored by grain diameter, was sliced orthogonally. X and Z slices simulate the typical vertical (X) and horizontal (Z) cuts made of ice cores. **(b-c)** Multiple XYZ slices in these plots represent 60 μm intervals (i.e., the reconstructed voxel size of DCT data). Grain sizes calculated from 3D volumetric data are plotted against grain sizes calculated from combined XYZ orthogonal slices, hence the 1:1 line, as well as the average 2D size of each grain found across XYZ slices combined. Each grain's average size based on either all X-slices or all Z-slices is also compared to its volume-calculated size, with the full 2D size range for each grain plotted as vertical bars. Note that **(b)** plots diameters represent equivalent sphere (3D) or circle (2D) diameters, whereas **(c)** plots the major-axis length of grains based on ellipsoid (3D) and ellipse (2D) fitting.
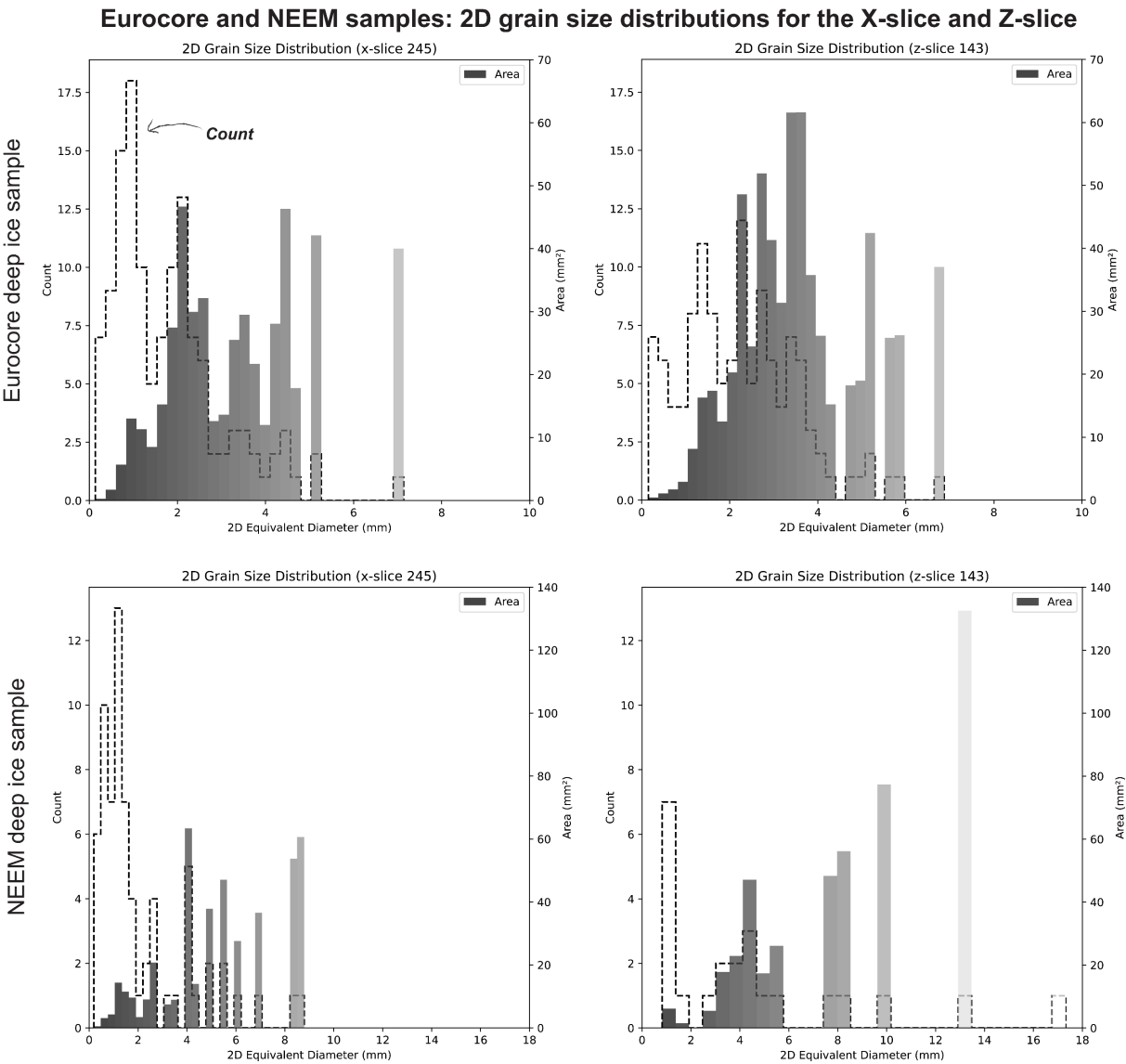
Eurocore and NEEM samples: 2D grain size distributions for the X-slice and Z-slice

1198
1199
1200 **Figure S3.** Grain size distributions for the single X-slice and Z-slice of the Eurocore and NEEM samples in
1201 **Figure 5** of the main text and **Figure S2**.
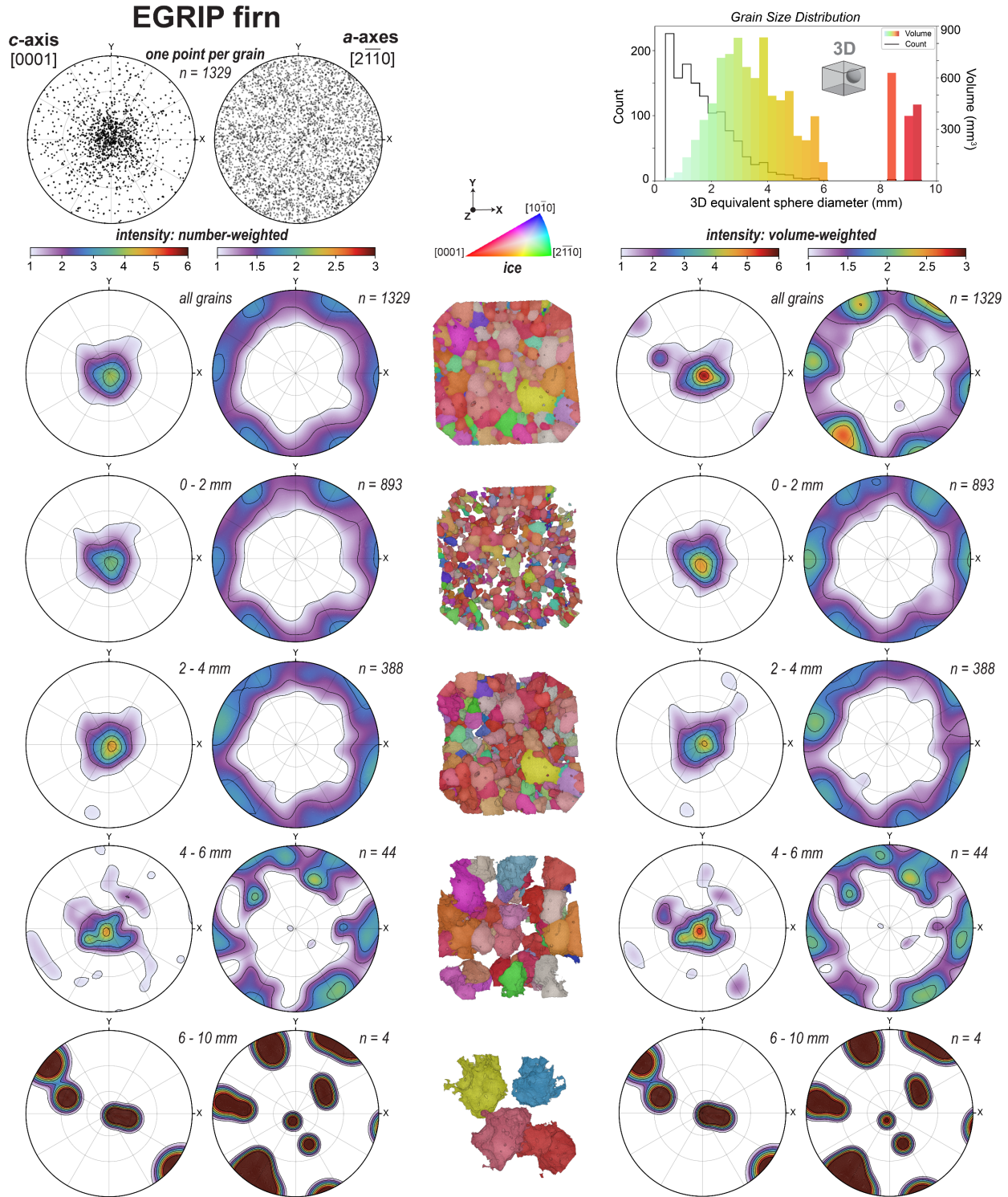
1202
1203
1204 **Figure S4.** Comparing pole figure analyses of firn grain orientations based upon count, volume, and grain
1205 size fraction. Pole figures provide views down the sample Z-axis, parallel to the ice-core axis. Analytical
1206 details are provided in the **Methods**. Note that glaciological studies employing EBSD of ice may instead

1207   represent the *a*-axes as $[11\bar{2}0]$ (-a_3), which is symmetrically equivalent to $[2\bar{1}\bar{1}0]$ (+a_1) (e.g., see Qi and
1208   others, 2019).
1209

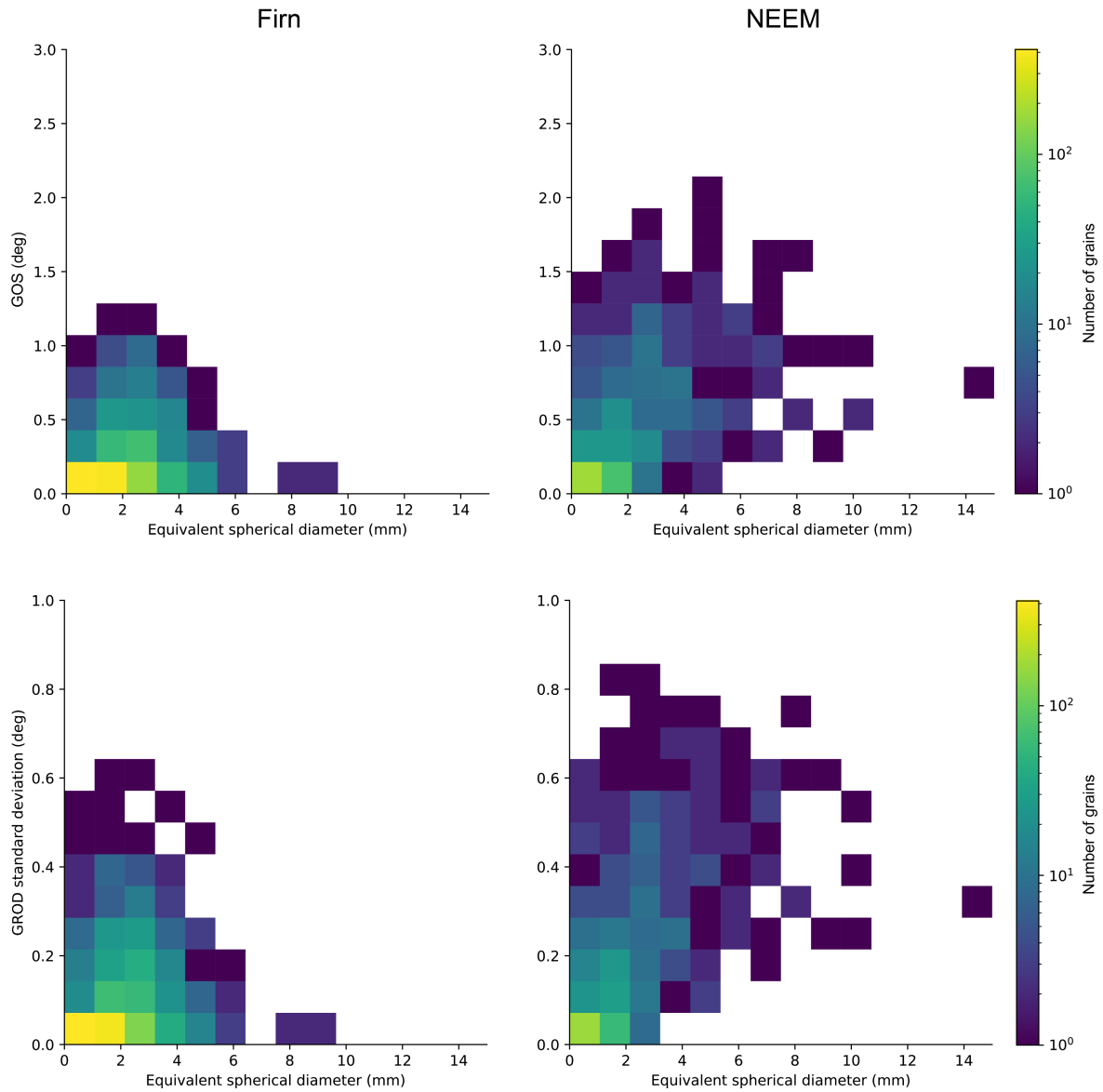# Firn and NEEM samples: Grain misorientation statistics

1210

1211   **Figure S5.** Grain-misorientation plots supporting observations reported for the EGRIP firn and NEEM
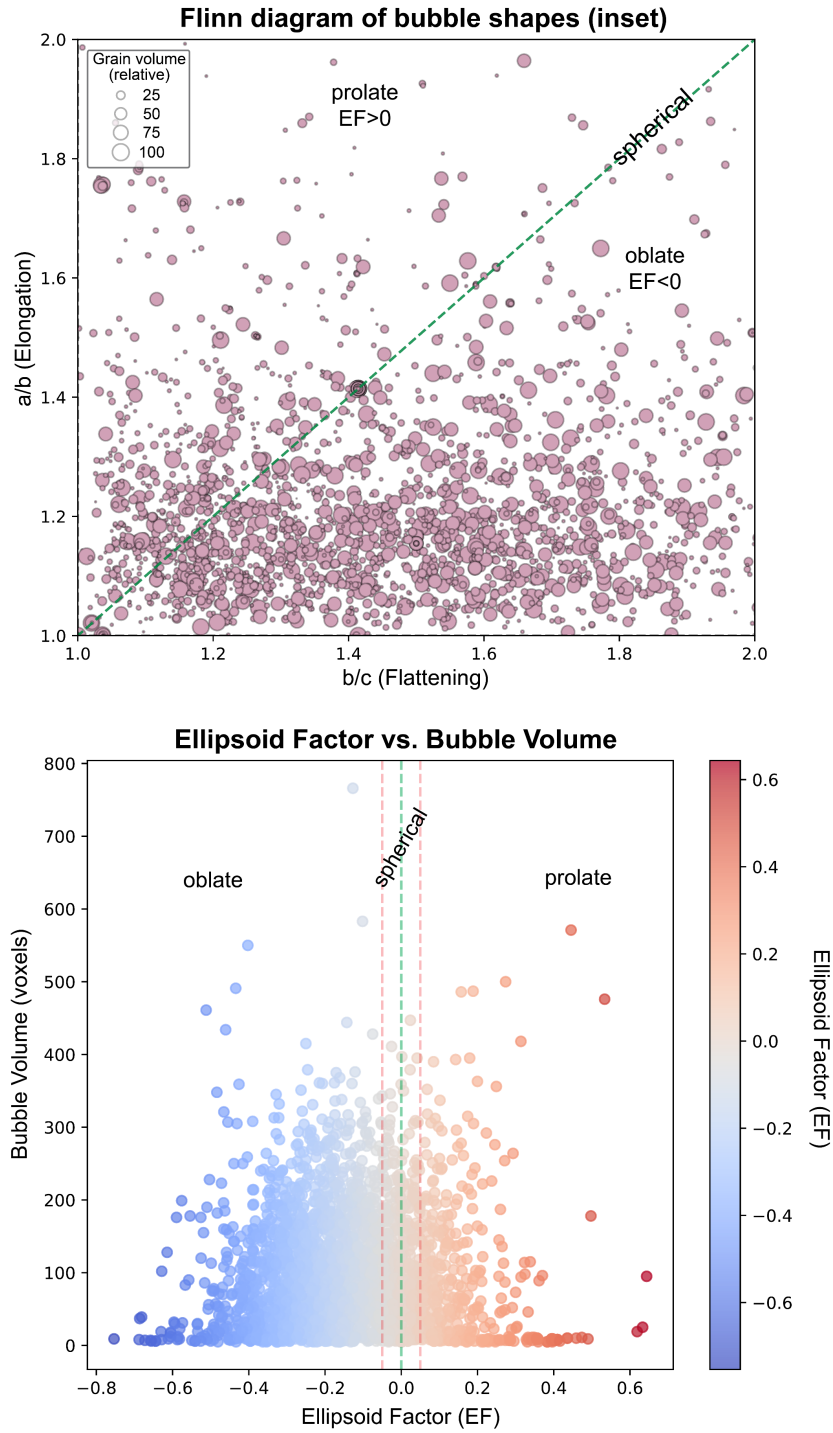1212   deep ice samples in Sections 3.4 and 3.5.

1213

Figure S6. Additional analytical plots supporting observations reported in Section 3.5. The Flinn diagram for bubbles in the Eurocore sample represents an inset of the diagram included in **Figure 10** of the main text. The lower plot shows the distribution of the bubbles based on bubble volume and ellipsoid factor.
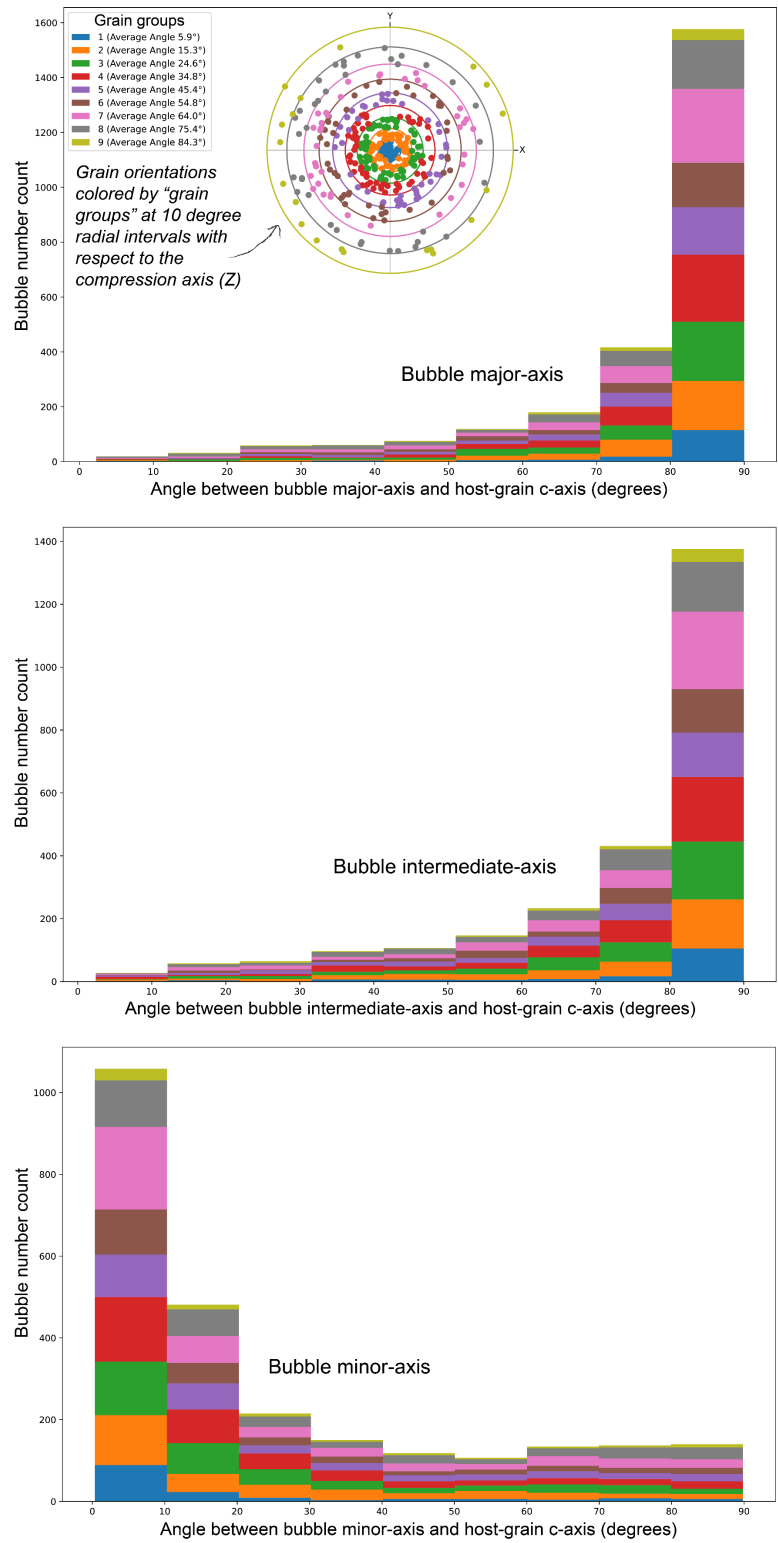
Figure S7. Stacked histograms showing the distribution of angle-differences between directions of bubble axes and the host-grain *c*-axis, with respect to how the host-grain *c*-axis is oriented relative to

1223   the compression axis (Z). The "grain groups" represent grains within 10-degree radial intervals away

1224   from the vertical (Z) direction, which represents the ice-core axis and direction of maximum

1225   compression. In the legend, the average *c*-axis angle misorientation from the Z-direction is reported for

1226   each "grain group".