

# Determining cooling rates from mica $^{40}\text{Ar}/^{39}\text{Ar}$ thermochronology data: effect of cooling path shape

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## Abstract

Tectonic models are commonly underpinned by metamorphic cooling rates derived from diffusive-loss thermochronology data. Such cooling ages are usually linked to temperature via Dodson's 1973 closure temperature ( $T_C$ ) formulation, which specifies a 1/time-shaped cooling path. Geologists, however, commonly discuss cooling rates as a linear temperature/time shape. We present the results of a series of simple finite-difference diffusion models for Ar diffusion in muscovite and biotite that show that the difference in recorded age between 1/t and linear cooling paths increases significantly with hotter starting temperatures, slower cooling rates and smaller grain sizes. Our results show that it is essential to constrain the cooling path shape in order to make meaningful interpretations of the measured data.

## Introduction

34 The ratio of parent to daughter radiogenic isotopes has been used for over a hundred  
 35 years to constrain geological ages and timescales (e.g. as reviewed by Condon and Schmitz,  
 36 2013). Minerals that host the radioactive parent element are commonly referred to as either  
 37 “geochronometers”, which record the timing of their crystallisation or  
 38 “thermochronometers”, which record the timing of cooling through an estimated temperature  
 39 window at some point after their crystallisation (e.g. as reviewed by Reiners, 2005). The  
 40 record of different time-temperature pairs in any one rock or tectonic region helps to  
 41 constrain thermal history and thus provide clues about the mechanism(s) by which the rocks  
 42 were exhumed to the surface.

43 Many thermochronometers are based on the premise that some of the daughter isotope  
 44 concentration is lost via thermally-activated diffusion at high temperatures, and that the  
 45 resulting mineral age can be linked to temperature via the mathematics governing such  
 46 diffusion. The temperature of a thermochronometer-bearing rock at the time the  
 47 thermochronometer recorded its apparent (bulk, whole-grain average) cooling age is most  
 48 commonly estimated using Dodson’s closure temperature ( $T_C$ ) formulation (Dodson 1973),  
 49 which, for thermally activated diffusion described by

$$50 \quad D = D_0 e^{-E_a/RT} \quad [1]$$

51 is given by:

$$52 \quad T_C = R/[E_a \ln(A\tau D_0/a^2)] \quad [2]$$

53 Where  $D$  is the diffusion coefficient,  $D_0$  is the diffusion pre-exponential factor,  $R$  is  
 54 the gas constant,  $E_a$  is the activation energy,  $a$  is the diffusion (or grain) radius,  $A$  is a grain-  
 55 shape-related constant and  $\tau$  relates the  $T_C$  to cooling rate:

$$56 \quad \tau = \frac{R}{(E_a dT^{-1}/dt)} = - \frac{RT^2}{(E_a dT/dt)} \quad [3]$$

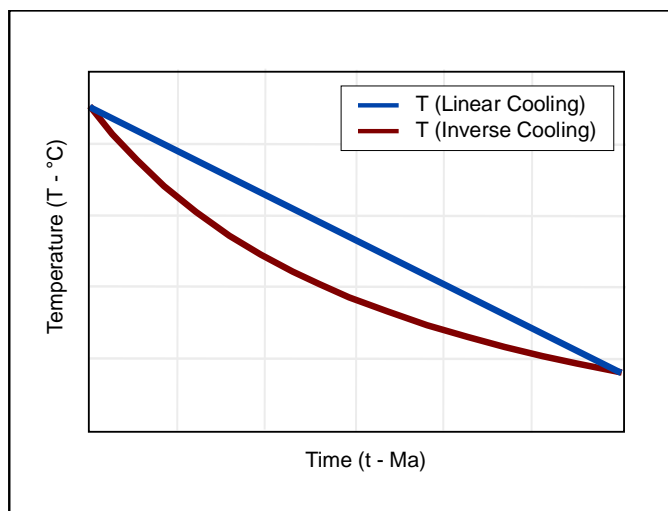
57 This result of an analytical solution to the diffusion equation has had an enduring  
 58 legacy due to its mathematical elegance and simplicity of application. However the Dodson  
 59  $T_C$  formulation is underpinned by several important assumptions and approximations:

- 60 (1) that thermally activated volume diffusion was the only mechanism by which the  
 61 daughter isotope was mobilised within the mineral;
- 62 (2) that the mineral crystallized with no inherited daughter isotope;
- 63 (3) that a daughter isotope concentration of zero was maintained at the mineral grain  
 64 boundary throughout cooling;
- 65 (4) that the starting temperature was high enough for diffusion of the daughter isotope  
 66 to be efficient, and removal from the grain to be geologically instantaneous, and

67 (5) that the cooling path from the time of crystallisation to the time of closure  
68 conformed to a  $1/t$  (time) -shape.

69 These approximations have a major impact on the applicability of the formulation to  
70 any particular geological scenario. The further any scenario deviates from these assumptions,  
71 the greater the (commonly un-quantified and un-reported) interpretational uncertainties on the  
72 link between age and temperature. A refinement of the  $T_C$  formulation to consider cases that  
73 did not conform to point (4) was proposed by Ganguly and Tirone, 1999, but has not been  
74 applied by the thermochronometer community to nearly the same extent that the original  
75 Dodson formulation has been.

76 The Dodson closure temperature formulation is most commonly used to constrain  
77 cooling rates by linking the  $T_C$  + time pair to a higher temperature + time pair linearly.  
78 However,  $T_C$  has been derived explicitly for temperature histories that involve cooling  
79 proportional to  $1/t$  (Figure 1) as this creates a linear time dependence in the exponent in  $\exp(-$   
80  $E_a/RT)$ ) and allows the analytic integration of the time dependence. To calculate a closure  
81 temperature using the Dodson  $T_C$  formulation and then to use that result to calculate a linear  
82 cooling rate is therefore both circular (as also noted by e.g. Ganguly and Tirone, 2009) and  
83 ultimately incorrect.



84

85 **Fig. 1.** A schematic representation showing the difference between linear (blue) and  
86  $1/t$  (red) cooling paths. Note that the  $1/t$ -shaped path initially cools faster, therefore reducing  
87 the opportunity for daughter product loss by diffusion.

88

89 Modern analytical equipment can now provide ever more precise isotope  
90 concentration (age) data, at ever increasing spatial resolution. Furthermore, the diffusion

91 equation can be solved numerically on any standard computer. Here we investigate the  
92 effects of linear and 1/t cooling path shapes on the bulk ages and core-rim age profiles of Ar  
93 in muscovite and biotite in grains of different radius that have cooled from different  
94 temperatures at different rates. The model results show that the ages recorded by muscovite  
95 and biotite that have cooled following these different simple end-member paths differ  
96 significantly, especially at higher peak temperatures increase and smaller grain sizes. The  
97 results also allow cooling rates to be determined directly if there is independent evidence for  
98 cooling path shape, so long as the time at which cooling started is known.

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100

### 101 **The DiffArgP\_inverse Code**

102 The finite-difference code DiffArgP\_inverse is a modified version of DiffArg  
103 (Wheeler 1996). It is written in Matlab 4.1 and solves the diffusion equation numerically.  
104 DiffArgP\_inverse differs from DiffArg in that it includes the effect of pressure on the  
105 diffusion of Ar in muscovite (Harrison et al., 2009) and the functionality to model 1/t-shaped  
106 thermal histories to match the analytical solution of Dodson, 1973, rather than only linear or  
107 piecewise-linear histories. DiffArg and its modified variants has previously been used to  
108 model Ar diffusion in different minerals that experienced complex metamorphic histories in a  
109 variety of tectonic environments (e.g. Mark et al., 2008; Warren et al., 2012a;b; Wartho et al.,  
110 2013; McDonald et al., 2016; 2018). The code allows the user to input any thermal and  
111 (de)compression history and produces outputs of integrated single grain (bulk) ages and core-  
112 rim age profiles. Any of the DiffArg versions are available from Hanke or Warren on request.  
113 Further details of the DiffArg\_Inverse code are presented in Supplementary Document S1.

114

### 115 **Methods**

116 The bulk (volume-integrated)  $^{40}\text{Ar}/^{39}\text{Ar}$  ages of muscovite and biotite of different  
117 grain size were modelled for a variety of different starting temperatures and linear vs. inverse  
118 (1/t) cooling histories. Muscovite and biotite were modelled with cylindrical geometry and  
119 grain radii of 1 mm, 0.5 mm, and 0.25 mm as these are the most typical grain sizes picked for  
120 metamorphic  $^{40}\text{Ar}/^{39}\text{Ar}$  analyses. The diffusion parameters applied to each mineral are  
121 outlined in Table 1.

122 All minerals were modelled as “crystallising” then instantaneously cooling from  
123 starting temperatures of 700°C, 600°C, 500°C, and 450°C at a starting pressure of 1 GPa to  
124 represent a variety of metamorphic terranes exhuming from mid-crustal conditions (Tables 1,

125 2, 4, Supplementary Tables S.2, S.4). A series of muscovite models was run at a starting  
 126 pressure of 2 GPa to more closely match conditions found in subduction zones (c.f. Warren et  
 127 al., 2012a; Table 3, Supplementary Table S.3), and a further series of muscovite models was  
 128 run with spherical geometry to allow comparison with the cylindrical geometry models  
 129 (Supplementary Table S.5 and Figure S.6). Linear cooling rates of 5, 10, 25, 50, and 70°CMa<sup>-1</sup>  
 130 were run in order to compare results for typical rates of cooling in different tectonic terranes.  
 131 1/t cooling rate models were run for equivalent “time to reach 0°C” as the linear models, in  
 132 order to compare results for different cooling path shapes. Model pressures were decreased to  
 133 0 GPa over the same time interval.

134 The grain boundary conditions in all models were modelled as zero daughter element  
 135 concentration, for the purposes of investigating behaviour in an open system. Model ages  
 136 were calculated for 2-dimensional (cylindrical) diffusion geometry (Hames and Bowring  
 137 1994) and the time integration was performed using the Crank–Nicholson solver, with a  
 138 recommended time step that is 10 times larger than the value suggested for a stable fully-  
 139 explicit method (Table 1; Wheeler 1996).

140 A series of models was run to test the effect of the published experimental  
 141 uncertainties on E<sub>a</sub> and D<sub>0</sub> (Harrison et al., 2009 for muscovite and Harrison et al., 1985 for  
 142 biotite) on the model results. The results are detailed in Supplementary Table S.7.

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144 **Table 1.** Diffusion and other model parameters used in this study.

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<b>Modelled diffusion parameters</b>						
<b>Mineral</b>	<b>System</b>	<b>E<sub>a</sub></b> Jmol <sup>-1</sup>	<b>D<sub>0</sub></b> mm <sup>2</sup> s <sup>-1</sup>	<b>V<sub>0</sub></b> cm <sup>3</sup> mol <sup>-1</sup>	<b>P<sub>0</sub></b> Gpa	<b>Reference</b>
Muscovite	<sup>40</sup> Ar/ <sup>39</sup> Ar	263592	2.30E+02	14	1	Harrison et al., (2009)
Biotite	<sup>40</sup> Ar/ <sup>39</sup> Ar	196648	7.70E+00	0	0	Harrison et al., (1985)
<b>Other model parameters</b>						
<b>Mineral</b>	<b>System</b>	<b>Grain shape</b>	<b>Radius range</b> mm	<b>Starting temp range</b> °C	<b>Linear cooling rate range</b> °C/Ma	<b>Starting pressure range</b> GPa
Muscovite	<sup>40</sup> Ar/ <sup>39</sup> Ar	Cylinder	1-0.25	700-450	5, 10, 25, 50, 70	2-1

Biotite	$^{40}\text{Ar}/^{39}\text{Ar}$	Cylinder	1-0.25	700-450	5, 10, 25, 50, 70	1
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**Global model parameters**

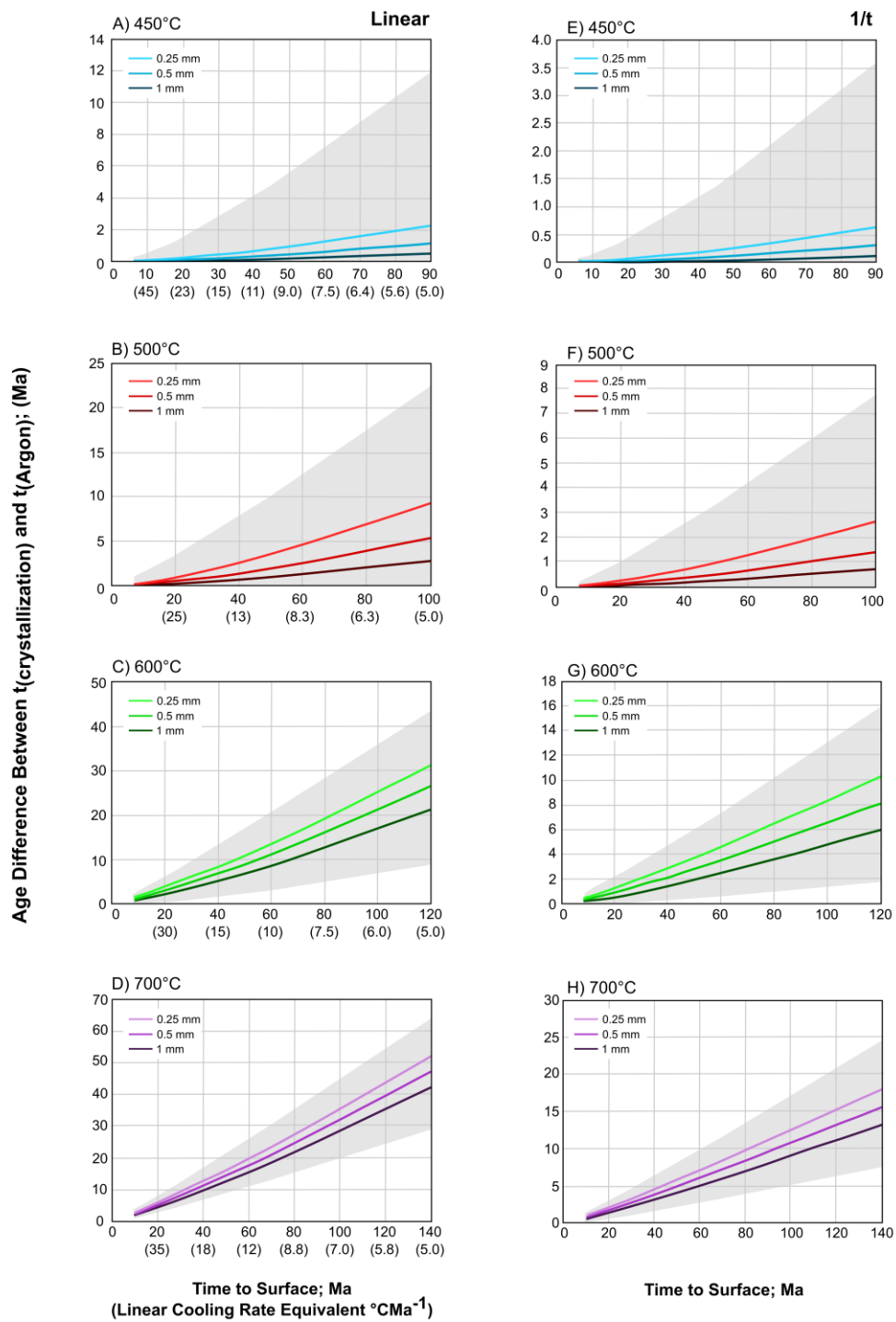
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Grain boundary:	Zero concentration
Solver:	Crank-Nicholson
Time step	10

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147 **Results**

148 The model results are plotted in Figures 2 and 3 (Muscovite modelled from pressures  
149 of 1 GPa and 2 GPa respectively) and 4 (Biotite from 1 GPa). Summary model results for the  
150 bulk (volume-averaged) ages are presented in Tables 2 (Muscovite), and 3 (Biotite). Full  
151 results including core-rim model age variations are presented in Supplementary Tables S.2  
152 (Muscovite 1 GPa) Table S.3 (Muscovite 2 GPa), Table S.4 (Biotite) and Table S.5  
153 (Muscovite 1 GPa with spherical geometry).



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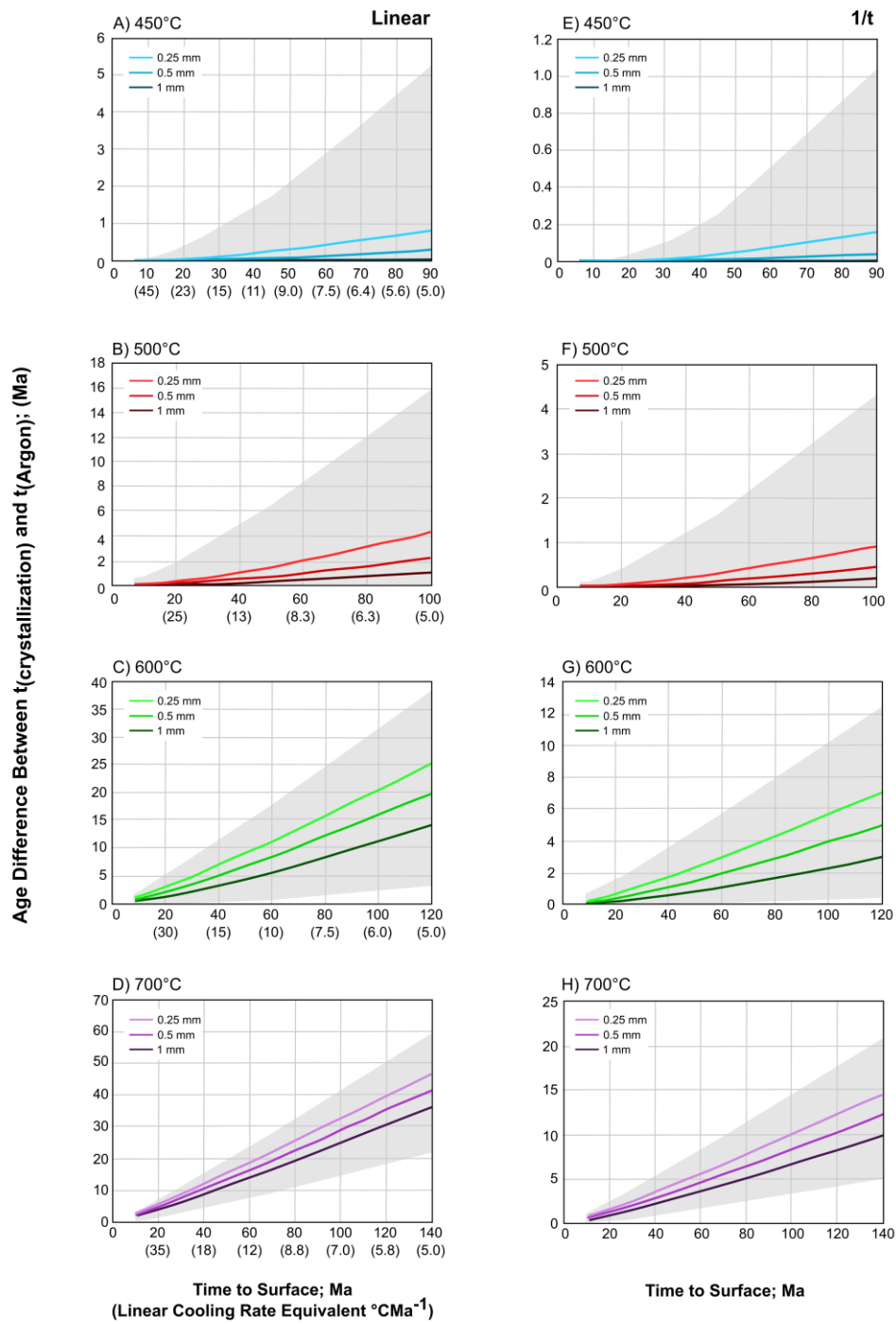
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**Fig. 2.** *Muscovite linear and 1/t results for models run at 1 GPa. Different coloured lines show different grain sizes. A-D show results for linear models at different starting temperatures; E-H show results for 1/t models that run over the same timescale. For ease of comparison, both sets of models run for the equivalent “time to surface” which is plotted on the x-axis. The equivalent linear rate is plotted underneath the “time to surface” value on*

160 *the linear model plots. The y-axis plots the difference between the time at which cooling*  
161 *starts and the recorded  $^{40}\text{Ar}/^{39}\text{Ar}$  age: if this is, the grain size and the starting temperature*  
162 *are known for the analysed samples, then the cooling rate can be read off the graph directly.*  
163 *Note the differences in the y-axis scale between the linear and 1/t results. The grey outline*  
164 *maps the maximum uncertainty associated with the experimental diffusion parameters of*  
165 *Harrison et al., 2009 for the 0.5 mm grain-size models (the results for the other grain sizes*  
166 *will scale accordingly).*

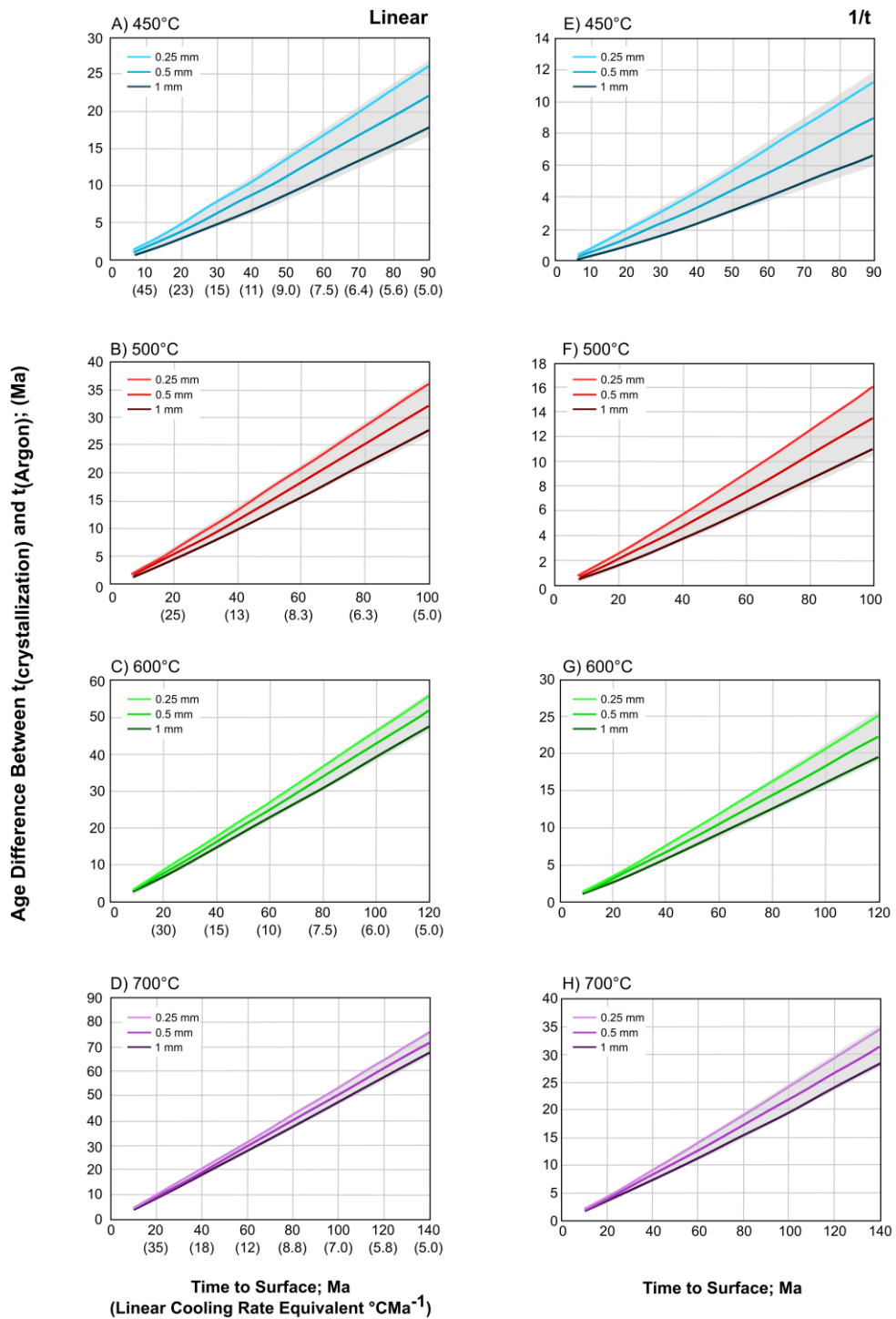




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**Fig 3.** *Muscovite linear and 1/t results for models run at 2 GPa. Different coloured lines show different grain sizes. A-D show results for linear models at different starting temperatures; E-H show results for 1/t models that run over the same timescale. For ease of comparison, both sets of models run for the equivalent “time to surface” which is plotted on the x-axis. The equivalent linear rate is plotted underneath the “time to surface” value on the linear model plots. The y-axis plots the difference between the time at which cooling*

174 starts and the recorded  $^{40}\text{Ar}/^{39}\text{Ar}$  age: if this is, the grain size and the starting temperature  
175 are known for the analysed samples, then the cooling rate can be read off the graph directly.  
176 Note the differences in the y-axis scale between the linear and  $1/t$  results. The grey outline  
177 maps the maximum uncertainty associated with the experimental diffusion parameters of  
178 Harrison et al., 2009 for the 0.5 mm grain-size models (the results for the other grain sizes  
179 will scale accordingly).  
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*Fig. 4. Biotite linear and 1/t model results. Different coloured lines show different grain sizes. A-D show results for linear models at different starting temperatures; E-H show results for 1/t models that run over the same timescale. For ease of comparison, both sets of models run for the equivalent “time to surface” which is plotted on the x-axis. The equivalent linear rate is plotted underneath the “time to surface” value on the linear model*

187 plots. The y-axis plots the difference between the time at which cooling starts and the  
 188 recorded  $^{40}\text{Ar}/^{39}\text{Ar}$  age: if this is, the grain size and the starting temperature are known for  
 189 the analysed samples, then the cooling rate can be read off the graph directly. Note the  
 190 differences in the y-axis scale between the linear and  $1/t$  results. The grey outline maps the  
 191 maximum uncertainty associated with the experimental diffusion parameters of Harrison et  
 192 al., 2009 for the 0.5 mm grain-size models (the results for the other grain sizes will scale  
 193 accordingly).

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196 **Table 2.** Model results for muscovite diffusion run with cylindrical geometry and at 1  
 197 GPa.

Linear Models		Cooling Rate ( $^{\circ}\text{C}\text{Ma}^{-1}$ )									
		5		10		25		50		70	
Grain Radius	T ( $^{\circ}\text{C}$ )	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma
0.25 mm	450	90	2.40	45	0.84	18	0.20	9	0.07	6.4	0.04
	500	100	10.10	50	3.91	20	1.05	10	0.37	7.1	0.23
	600	120	31.26	60	14.45	24	5.13	12	2.31	8.6	1.58
	700	140	52.08	70	24.87	28	9.31	14	4.40	10	3.06
0.5 mm	450	90	1.15	45	0.38	18	0.09	9	0.03	6.4	0.01
	500	100	5.81	50	2.10	20	0.54	10	0.19	7.1	0.12
	600	120	26.45	60	11.96	24	4.10	12	1.78	8.6	1.18
	700	140	47.35	70	22.44	28	8.30	14	3.88	10	2.67
1 mm	450	90	0.52	45	0.16	18	0.03	9	0.01	6.4	0.00
	500	100	2.99	50	1.04	20	0.25	10	0.09	7.1	0.06
	600	120	21.33	60	9.31	24	2.99	12	1.21	8.6	0.76
	700	140	42.29	70	19.82	28	7.20	14	3.32	10	2.26
1/t Models											
Grain Radius	T ( $^{\circ}\text{C}$ )	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma
0.25 mm	450	90	0.58	45	0.19	18	0.05	9	0.02	6.4	0.01
	500	100	2.61	50	0.95	20	0.25	10	0.09	7.1	0.05
	600	120	10.27	60	4.60	24	1.56	12	0.67	8.6	0.44
	700	140	18.06	70	8.43	28	3.06	14	1.41	10	0.96
0.5 mm	450	90	0.26	45	0.08	18	0.02	9	0.00	6.4	0.00
	500	100	1.36	50	0.48	20	0.12	10	0.04	7.1	0.03

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**Table 3.** Model results for muscovite diffusion run with cylindrical geometry and at 2

GPa.

Linear Models	Cooling Rate ( $^{\circ}\text{C}\text{Ma}^{-1}$ )										
	5		10		25		50		70		
Grain Radius	T ( $^{\circ}\text{C}$ )	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma
1 mm	600	120	8.12	60	3.52	24	1.13	12	0.46	8.6	0.29
	700	140	15.65	70	7.23	28	2.57	14	1.16	10	0.79
	450	90	0.10	45	0.03	18	0.00	9	0.00	6.4	0.00
	500	100	0.66	50	0.22	20	0.06	10	0.02	7.1	0.01
1 mm	600	120	5.97	60	2.45	24	0.70	12	0.26	8.6	0.16
	700	140	13.25	70	6.02	28	2.09	14	0.92	10	0.62

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Linear Models		Cooling Rate ( $^{\circ}\text{C}\text{Ma}^{-1}$ )										
		5		10		25		50		70		
Grain Radius	T ( $^{\circ}\text{C}$ )	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma	
0.25 mm	450	90	0.82	45	0.27	18	0.05	9	0.01	6.4	0.00	
	500	100	4.33	50	1.56	20	0.40	10	0.14	7.1	0.09	
	600	120	25.22	60	11.29	24	3.80	12	1.62	8.6	1.06	
	700	140	47.06	70	22.27	28	8.21	14	3.83	10	2.63	
0.5 mm	450	90	0.31	45	0.07	18	0.01	9	0.00	6.4	0.00	
	500	100	2.24	50	0.80	20	0.20	10	0.07	7.1	0.04	
	600	120	19.86	60	8.52	24	2.64	12	1.04	8.6	0.64	
	700	140	41.87	70	19.58	28	7.08	14	3.25	10	2.21	
1 mm	450	90	0.05	45	0.01	18	0.00	9	0.00	6.4	0.00	
	500	100	1.12	50	0.38	20	0.08	10	0.02	7.1	0.01	
	600	120	14.12	60	5.59	24	1.53	12	0.56	8.6	0.34	
	700	140	36.32	70	16.69	28	5.87	14	2.62	10	1.76	

1/t Models		Cooling Rate ( $^{\circ}\text{C}\text{Ma}^{-1}$ )										
		5		10		25		50		70		
Grain Radius	T ( $^{\circ}\text{C}$ )	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma	
0.25 mm	450	90	0.16	45	0.04	18	0.01	9	0.00	6.4	0.00	
	500	100	0.94	50	0.33	20	0.08	10	0.03	7.1	0.02	
	600	120	7.11	60	3.03	24	0.93	12	0.37	8.6	0.23	
	700	140	14.68	70	6.75	28	2.39	14	1.07	10	0.73	
0.5 mm	450	90	0.04	45	0.01	18	0.00	9	0.00	6.4	0.00	
	500	100	0.47	50	0.16	20	0.04	10	0.01	7.1	0.00	
	600	120	5.01	60	1.99	24	0.54	12	0.20	8.6	0.12	
	700	140	12.32	70	5.58	28	1.92	14	0.84	10	0.56	
1 mm	450	90	0.00	45	0.00	18	0.00	9	0.00	6.4	0.00	
	500	100	0.21	50	0.06	20	0.01	10	0.00	7.1	0.00	

600	120	3.01	60	1.10	24	0.28	12	0.10	8.6	0.06
700	140	9.98	70	4.41	28	1.46	14	0.61	10	0.40

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**Table 4.** Model results for biotite diffusion.Cooling Rate ( $^{\circ}\text{C}\text{Ma}^{-1}$ )

Linear Models		5		10		25		50		70	
Grain Radius	Starting T ( $^{\circ}\text{C}$ )	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma
0.25 mm	450	90	25.88	45	11.89	18	4.24	9	1.91	6.4	1.27
	500	100	35.88	50	16.92	20	6.24	10	2.91	7.1	2.00
	600	120	55.88	60	26.92	24	10.24	12	4.91	8.6	3.41
	700	140	75.88	70	36.92	28	14.24	14	6.91	10	4.87
0.5 mm	450	90	21.82	45	9.78	18	3.34	9	1.42	6.4	0.96
	500	100	31.82	50	14.82	20	5.34	10	2.49	7.1	1.63
	600	120	51.82	60	24.82	24	9.34	12	4.49	8.6	3.11
	700	140	71.82	70	34.82	28	13.34	14	6.49	10	4.56
1 mm	450	90	17.45	45	7.54	18	2.39	9	0.96	6.4	0.63
	500	100	27.45	50	12.51	20	4.39	10	1.97	7.1	1.28
	600	120	47.45	60	22.51	24	8.39	12	3.97	8.6	2.77
	700	140	67.45	70	32.51	28	12.39	14	5.97	10	4.21

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1/t Models		5		10		25		50		70	
Grain Radius	T ( $^{\circ}\text{C}$ )	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma	Time to $0^{\circ}\text{C}$ Ma	$\Delta t$ Ma
0.25 mm	450	90	11.10	45	4.97	18	1.67	9	0.75	6.4	0.49
	500	100	15.76	50	7.25	20	2.62	10	1.18	7.1	0.80
	600	120	25.13	60	11.85	24	4.42	12	2.08	8.6	1.42
	700	140	34.39	70	16.45	28	6.23	14	2.95	10	2.07
0.5 mm	450	90	8.77	45	3.83	18	1.19	9	0.52	6.4	0.33
	500	100	13.31	50	6.03	20	2.14	10	0.94	7.1	0.63
	600	120	22.31	60	10.45	24	3.83	12	1.76	8.6	1.23
	700	140	31.40	70	14.89	28	5.54	14	2.65	10	1.80
1 mm	450	90	6.45	45	2.64	18	0.79	9	0.31	6.4	0.19
	500	100	10.81	50	4.76	20	1.58	10	0.69	7.1	0.45
	600	120	19.54	60	9.06	24	3.24	12	1.47	8.6	1.04
	700	140	28.28	70	13.33	28	4.93	14	2.36	10	1.58

207

208 The graphs all show similar trends:

209 (1) Faster cooling results in a smaller difference in time between the timing of  
210 maximum temperature attainment (cooling initiation) and the recorded cooling age ( $\Delta t$ ).

211 (2) Colder initial “peak” starting temperatures result in smaller  $\Delta t$ .

212 (3) Smaller grain sizes result in larger  $\Delta t$ .

213 (4) Smaller  $\Delta t$  values are recorded for the 1/t models than for the linear models.

214 Results (1) and (3) are consistency checks to show that the models are behaving as  
215 expected. Result (2) similarly matches the predictions of the modified formulation of  
216 Ganguly and Tirone, 1999. Result (4) clearly shows the importance of the cooling path shape  
217 on the resulting thermochronometer age – this will be discussed further below.

218 Figure 2 shows that very little diffusive loss is expected in white mica grains that cool  
219 from relatively low peak temperatures of 450°C. The  $^{40}\text{Ar}/^{39}\text{Ar}$  age of a 0.25 mm radius white  
220 mica grain cooling linearly at a rate of 5°C $\text{Ma}^{-1}$  from 450°C and 1 GPa would be expected to  
221 be 2.4 Ma younger than the peak temperature age, whereas one cooling from 700°C would be  
222 expected to yield an age that is 52 Ma younger (Table 2). Similarly, the  $^{40}\text{Ar}/^{39}\text{Ar}$  age of a  
223 0.25 mm radius white mica grain cooling linearly at a rate of 5°C $\text{Ma}^{-1}$  from 600°C and 2 GPa  
224 would be expected to be ~25 Ma younger than the peak temperature age (Table 3). Similar-  
225 sized grains cooling to 0°C over the same time interval but following a 1/t path from 450°C  
226 or 700°C at 1 GPa would only yield ages that were 0.6 or 16 Ma younger than the peak  
227 temperature age. A 1 mm radius grain cooling from 450°C, however, would be expected to  
228 record an age within uncertainty of the timing of peak metamorphism.

229 Models run using spherical diffusion geometry yield slightly younger ages ( $\Delta t$  of 54  
230 Ma rather than 52 Ma for a 0.25 mm radius grain cooling from 700°C at 5°C $\text{Ma}^{-1}$  for  
231 example; Supplementary Table S.5; Supplementary Figure S.6).

232 Figure 3 shows that biotite should yield significantly younger ages than muscovite for  
233 grains of the same radius, cooling from the same starting temperature and following the same  
234 cooling path. For example a 0.25 mm radius grain cooling at 5°C $\text{Ma}^{-1}$  from 450°C would be  
235 expected to be 26 Ma younger than the age of peak temperature metamorphism, whereas one  
236 cooling from 700°C at the same rate would be expected to yield an age that was 76 Ma  
237 younger (Table 3).

238

## 239 Discussion

240 The results clearly show that the shape of the cooling path makes an increasingly  
241 important contribution to the recorded thermochronometer age as grain sizes and cooling

242 rates decrease and peak temperatures increase. The uncertainty inherent in using the Dodson  
243  $T_C$  formulation to estimate (linear) cooling rates therefore also magnifies accordingly.

244 The model results are more sensitive to systematic uncertainties in the  
245 experimentally-determined activation energy ( $E_a$ ) than in the exponential pre-factor ( $D_0$ ) for  
246 each mineral (Figures 2-4 and Supplementary Table S.7). These figures show that  
247 uncertainties in the diffusion parameters have a significant, but systematic, effect on the  
248 recorded thermochronological ages. These uncertainties apply equally to both cooling history  
249 shapes discussed here.

250 The most recent diffusion parameters for muscovite (Harrison et al., 2009) were  
251 calculated for isotropic 3-dimensional (spherical) diffusion geometry. It has been suggested  
252 that modelling muscovite as a cylinder but using diffusion parameters calculated for spherical  
253 geometry invalidates the results (Foster and Lister, 2017). However the overall difference in  
254 the diffusion coefficient is a factor 2 in  $D_0$ , which translates into an activation barrier of <0.6  
255 kcal/mol at 400 K. This is well below the uncertainty of 7 kcal/mol in the Harrison et al.,  
256 2009 diffusion parameters and thus adds no extra uncertainty to our overall results, as also  
257 suggested in other studies (e.g. Huber et al., 2011).

258

### 259 **Applying Model Results to Natural Systems**

260 The results presented here can be used to constrain the cooling rates of natural  
261 systems if the following pieces of information are known or can be estimated:

262 1) A petrographically-based interpretation of the temperature at which the dated  
263 grain(s) grew, and the portion of the metamorphic path along which the grain(s) grew (e.g.  
264 prograde peak or retrograde). This will inform and constrain the extent of diffusive  
265 opportunity that the grain could have experienced. For example a grain growing during the  
266 prograde history will have longer residence at high temperatures, therefore allowing it more  
267 opportunity to lose argon.

268 2) The peak temperature experienced by the grain(s), required for the ultimate  
269 determination of a cooling rate.

270 3) The time at which the grain reached its peak temperature (constrained or estimated  
271 by independent geochronometers), required for the ultimate determination of a cooling rate.  
272 This is further discussed below.

273 4) The thermochronometric ages of the grains of interest; different data collection  
274 methods are further discussed below.

275 5) The grain size(s) of the dated grains.



276           6) The assumption or knowledge that open grain-boundary, thermally-activated  
277 diffusion was the dominant process in determining the final Ar concentration. This  
278 approximation is difficult to assess (e.g. Warren et al., 2012a,b) but should be acknowledged  
279 in any thermochronological interpretation.

280           Note that only very simple cooling path shapes have been modelled here. Steady  
281 progress is being made in the development of modelling tools that can suggest a “best fit”  
282 cooling path to U-Th-He, fission track and U-Pb rutile data, but currently none of these tools  
283 explicitly incorporate  $^{40}\text{Ar}/^{39}\text{Ar}$  data: e.g. HeFTy (Ketcham, 2005), QTQt (Gallagher, 2012),  
284 UpBeat (Smye et al., 2018).

285

286           **Determining the timing of cooling initiation:** Direct determination of a cooling rate  
287 (and cooling rate shape) from thermo- and geochronological data requires that at least two,  
288 and possibly three, T-t pairs are known. Timing of peak T in metamorphic rocks is  
289 commonly constrained by U-Pb ages of zircon, monazite, garnet, allanite and/or rutile, with  
290 secondary (higher-temperature cooling) T-t pairs provided by U-Pb rutile and/or titanite data.  
291 There are, of course, multiple uncertainties inherent in linking these ages to peak temperature  
292 because all of these minerals may crystallise at different stages of the metamorphic PT path.  
293 Careful petrochronological investigation is required to confirm that the ages yielded by any  
294 of these minerals relate to the timing of attainment of peak temperatures or higher-than-  
295 argon-closure cooling (e.g. Kohn et al., 2017).

296

297            **$^{40}\text{Ar}/^{39}\text{Ar}$  data collection methods:**  $^{40}\text{Ar}/^{39}\text{Ar}$  mica data can currently be collected in  
298 many different ways: by multiple- or single-grain step heating experiments (e.g. Turner,  
299 1970), by single grain fusion methods (e.g. Fleck and Carr, 1990) or by laser ablation (e.g.  
300 Kelley et al., 1994). All methods have their advantages and disadvantages in terms of volume  
301 of material analysed, analytical precision and petrographic (location) control on age.

302           The model data presented here are compatible for assessment against the bulk  
303 (volume-averaged) ages – i.e. equivalent to single grain fusion  $^{40}\text{Ar}/^{39}\text{Ar}$  data. We caution  
304 against using multiple-or single-grain step heating  $^{40}\text{Ar}/^{39}\text{Ar}$  ages to compare against model  
305 results. Plateau ages imply no core-rim variation in Ar distribution, (and thus an  
306 interpretation of rapid cooling), but a plateau result does not in itself guarantee that the  
307 calculated age is geologically meaningful, especially in high pressure metamorphic rocks  
308 (e.g. Sherlock and Arnaud, 1999). Non-plateau spectra can be produced by a variety of  
309 factors that complicate linking spectrum shapes to within-grain Ar distribution. Single grain

310 fusion populations can help provide an assessment of how homogeneous Ar is distributed  
311 across mica grains within individual samples (e.g. Uunk et al., 2018).

312 In-situ, high-spatial precision  $^{40}\text{Ar}/^{39}\text{Ar}$  data such as collected by laser ablation  
313 methods, and collected in grains large enough and cooled slowly enough from a high enough  
314 temperature to be able to detect such changes, can also be assessed against the core-rim  
315 model age predictions for simple linear and  $1/t$  cooling histories presented in Supplementary  
316 Tables S.2 -S.4.

317

318 **Comparing analytical data to model results:** The time difference ( $\Delta t$ ) between the  
319 timing of the thermal peak (or to be absolutely correct, the timing of cooling initiation) and  
320 the age recorded by the thermochronometer (Figures 2-4 ) provides a basis for determining  
321 cooling rates under the fundamental approximations (1) that thermally activated volume  
322 diffusion was the only mechanism by which the daughter isotope was mobilised within the  
323 mineral; (2) that the mineral crystallized with no inherited daughter isotope; and (3) that the  
324 experimentally-derived diffusion parameters mimic what happens in nature. It is important to  
325 acknowledge that minerals may not degas in a high-vacuum environment in an experiment  
326 that lasts a few days in the same way that a mineral degasses in a rock over millions of years,  
327 however these experimental data are the best available at the present day.

328 For example, consider a scenario whereby a 0.5 mm radius muscovite in a rock that  
329 started cooling from 500°C at 100 Ma yields an age of 94 Ma.  $\Delta t$  is therefore 6 Ma. Table 2  
330 and Figure 2 suggest that those data are compatible with a linear cooling rate of  $5^\circ\text{C}\text{Ma}^{-1}$ .  
331 However this is not enough information to determine whether (a) the system was diffusively  
332 open (a fundamental requirement of any diffusive-based interpretative link between age,  
333 temperature and cooling rate is that effectively there is infinite sink for the daughter element  
334 diffusing out of the mineral grain) and/or (b) whether the cooling path was overall linear or  
335 some other shape. Both of these can be resolved following a match between data and model  
336 predictions.

337 For example, a rock cooling from 600°C might yield 1 mm radius biotite grains with a  
338  $\Delta t$  of 9 Ma, 0.5 mm radius grains with a  $\Delta t$  of 10.5 Ma and 0.25 mm radius grains with a  $\Delta t$   
339 of 12 Ma. These data would be compatible with a cooling path of  $1/t$  shape that cooled to  
340 0°C over 60 Ma. A minimum of two different ages – either different grain sizes of the same  
341 mineral or different minerals, should allow differentiation of the best-fit cooling path.

342 At rapid cooling rates, the difference between the cooling ages predicted by a linear  
343 temperature decrease and a  $1/t$ -shaped path would be indistinguishable within the typical

344 uncertainties in analytical results and in the experimental diffusion parameters. At cooling  
345 rates  $<10^{\circ}\text{C}\text{Ma}^{-1}$ , differences in the shapes of the cooling paths start to become important for  
346 distinguishing between exhumation mechanisms.

347 Small values of  $\Delta t$  e.g.  $< 1$  Ma are currently challenging to resolve analytically. The  
348 mica  $^{40}\text{Ar}/^{39}\text{Ar}$  models for low starting temperatures confirm previous suggestions that  
349 rapidly-cooled rocks that reached low peak temperatures (such as in subduction zones) will  
350 not yield ages that allow cooling rates to be determined.

351

352 **Other factors affecting daughter element distribution:** Inheritance or loss of  
353 daughter product during recrystallization and deformation during cooling can affect daughter  
354 element concentrations much more than diffusion (Villa 1998; Allaz et al., 2011; Villa et al.,  
355 2014). It is also obvious that re-crystallisation during exhumation means that the temperature  
356 that that particular grain cooled from may be lower than the peak temperature. In cases where  
357 thermochronometer minerals show signs of secondary recrystallization or other chemical  
358 modification, the model results are almost certainly not applicable, and a link between  
359 temperature and age may be more difficult to constrain. The diffusion models are *only*  
360 applicable to rocks in which an open system can be assumed, and where both the timing and  
361 pressure-temperature conditions of the last episode of mineral crystallisation are known or  
362 can be estimated.

363 If the results presented here are used to estimate cooling rates or constrain cooling  
364 path shapes, each practitioner will need to estimate the geological uncertainty for their  
365 particular study, noting that this is almost certainly the largest overall source of error in their  
366 interpretation. Our results are based on the assumption that cooling starts directly after the  
367 model grain has crystallised at peak temperatures. In reality, the minerals of interest may  
368 have grown along the prograde path and/or have resided at peak temperatures for a  
369 geologically-significant period of time before cooling started. If temperatures were low  
370 enough for diffusion to be inefficient, some of that pre-cooling history may be recorded in the  
371 thermochronometer minerals. Thermochronologists should model the effect of pre-peak  
372 thermal history for their particular geological location to convince themselves whether or not  
373 the thermochronometer minerals in their study area may record this.

374

## 375 **Conclusions**

376 The rates and timescales over which rocks are buried, transformed, deformed and  
377 exhumed help constrain the tectonic mechanisms that act on them.  $^{40}\text{Ar}/^{39}\text{Ar}$  data from micas

378 have long been used to link time to temperature and thus constrain cooling rates. The  
379 Dodson closure temperature formulation (Dodson, 1973) provides an elegant analytical  
380 solution to the diffusion equation but its application for determining cooling rates is  
381 commonly based on assumptions that are a poor match to geological reality. Our results of a  
382 series of diffusion models that quantify the differences in age expected from a simple linear  
383 and 1/t-shaped cooling histories show that the cooling path shape exerts considerable  
384 influence on the resulting age at hotter starting temperatures, slower cooling rates and smaller  
385 grain sizes. If the cooling path shape and timing of cooling initiation are known, then our  
386 results also provide a simple way of estimating cooling rates and cooling rate shapes from the  
387 difference between the timing of cooling initiation at maximum temperature and the yielded  
388 thermochronometer age. Future incorporation of  $^{40}\text{Ar}/^{39}\text{Ar}$  diffusion systematics into forward  
389 modelling packages that also incorporate other thermochronometers provides the best future  
390 solution for constraining cooling rates, with the caveat that more precise diffusion data are  
391 needed.

392

393

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401

402

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485

## 486 **Supplementary Data**

487

488 **S.1** Instructions for operating DiffargP\_inverse

489

490 **S.2** Full muscovite results for 1GPa models

491

492 **S.3** Full muscovite results for 2GPa models

493

494 **S.4** Full biotite model results

495

496 **S.5.** Full muscovite results for spherical geometry models

497

498 **S.6.** Muscovite linear and 1/t results for models run at 1 GPa with spherical geometry.

499 Different coloured lines show different grain sizes. A-D show results for linear models at  
500 different starting temperatures; E-H show results for 1/t models that run over the same

501 timescale. For ease of comparison, both sets of models run for the equivalent “time to  
502 surface” which is plotted on the *x*-axis. The equivalent linear rate is plotted underneath the

503 “time to surface” value on the linear model plots. The *y*-axis plots the difference between the  
504 time at which cooling starts and the recorded  $^{40}\text{Ar}/^{39}\text{Ar}$  age: if this is, the grain size and the

505 starting temperature are known for the analysed samples, then the cooling rate can be read off  
506 the graph directly. Note the differences in the *y*-axis scale between the linear and 1/t results.

507

508 **S.7.** Results of sensitivity tests

509

510

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512

